

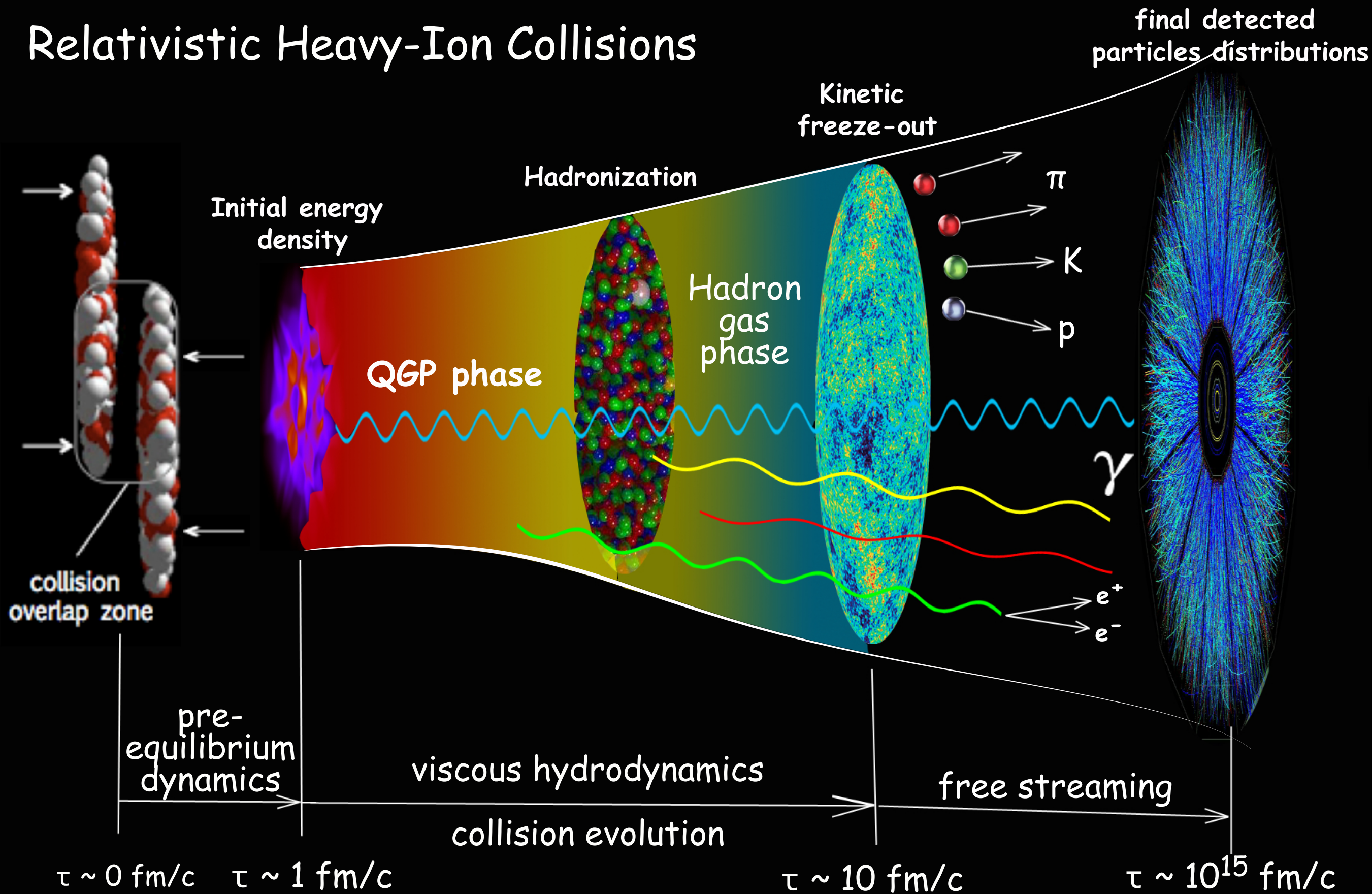
Thermal photons as a quark-gluon plasma thermometer revisited

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The Ohio State University

In collaboration with Ulrich Heinz, Charles Gale,
and Jean-Francois Paquet

Little Bang

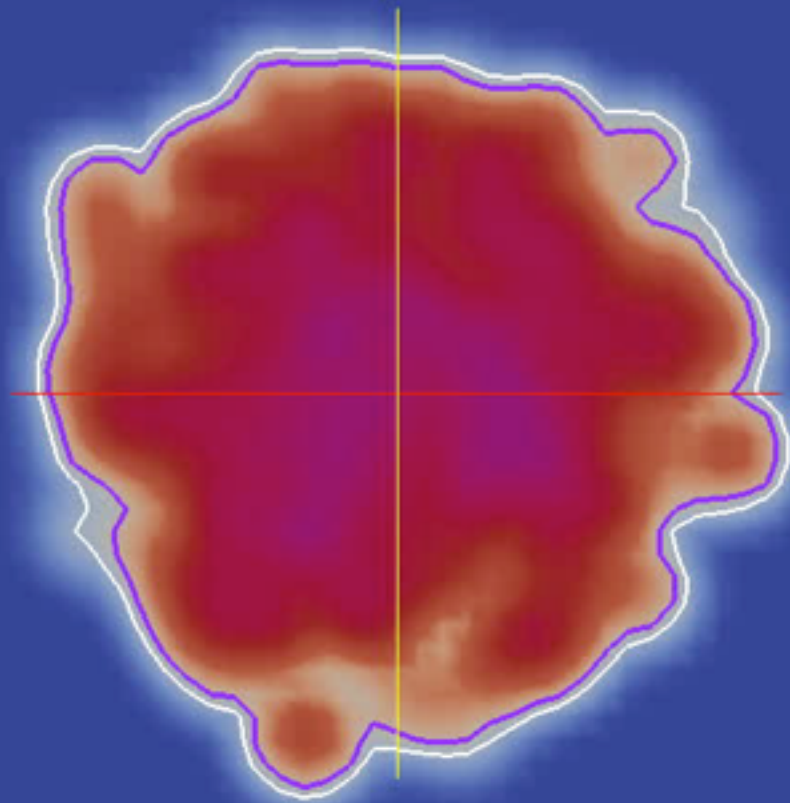
Relativistic Heavy-Ion Collisions



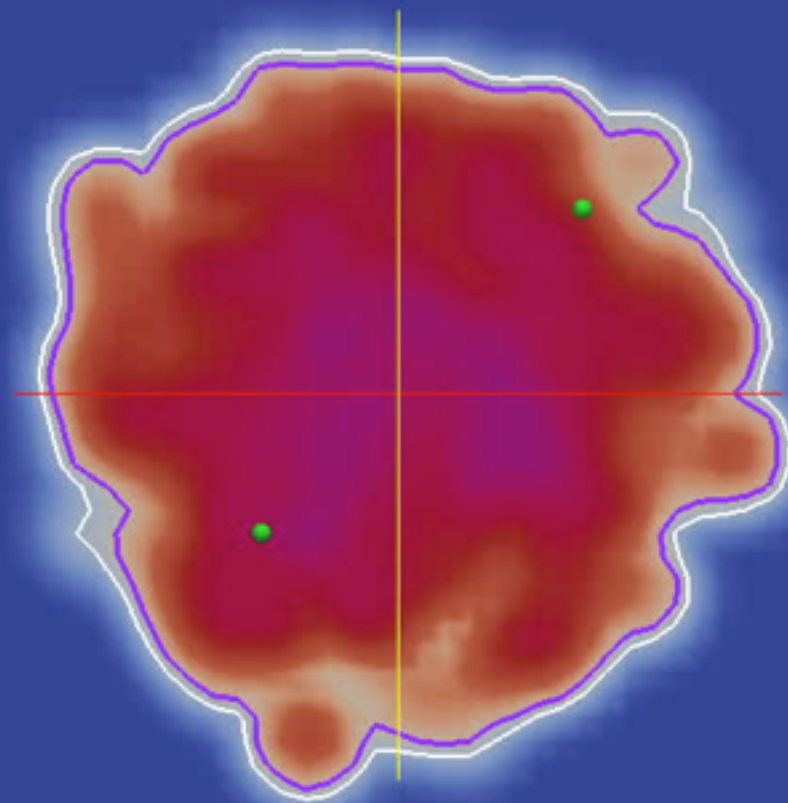
Photons from Heavy-ion Collisions

time = 0.6 fm/c

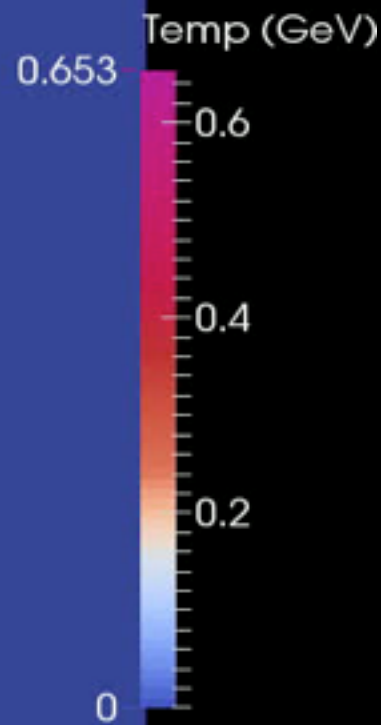
Pb+Pb @ 2.76 A TeV LHC



Charged Hadrons

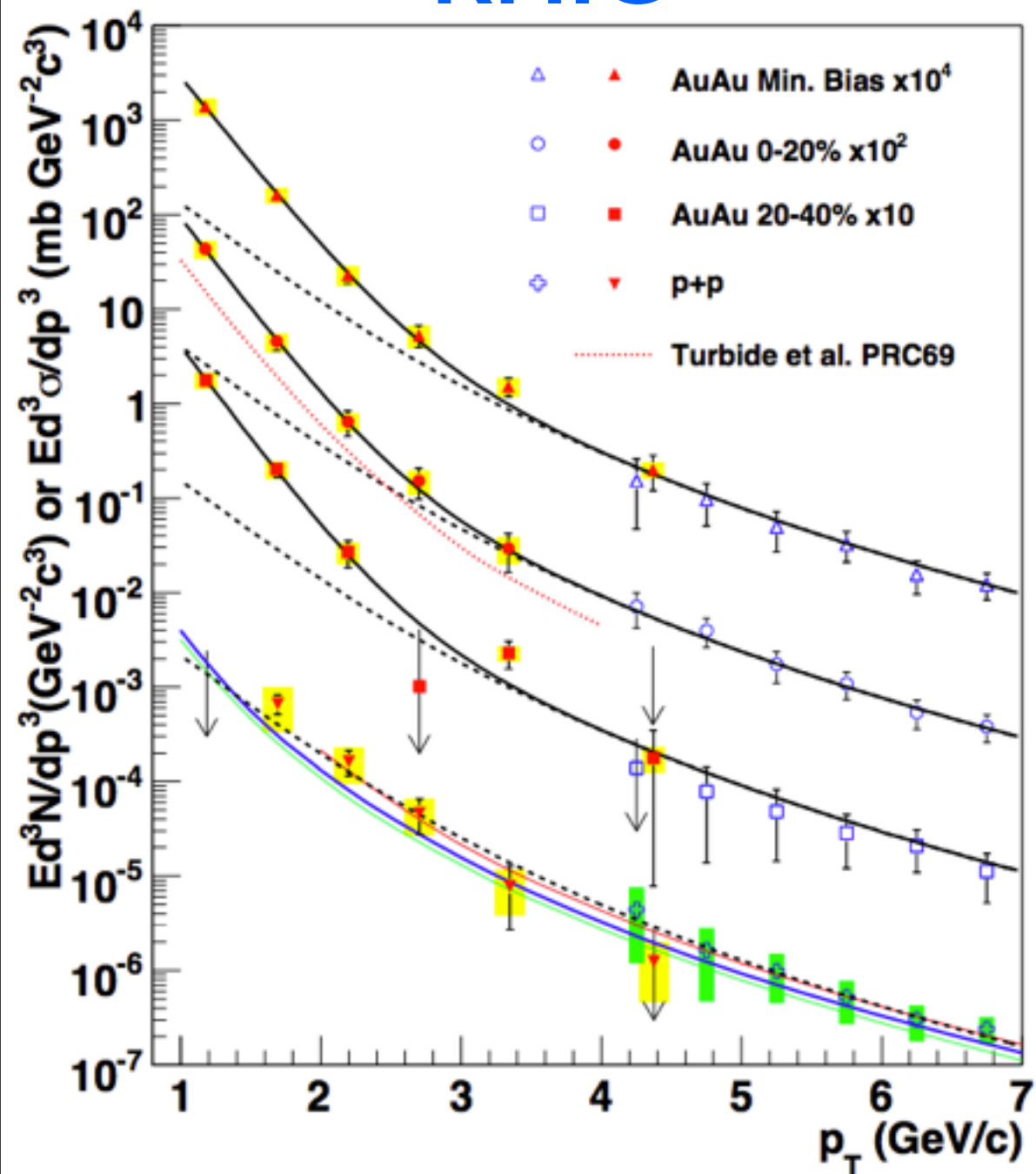


Thermal Photons oversample = 10



Fitted T_{eff} from Experiments

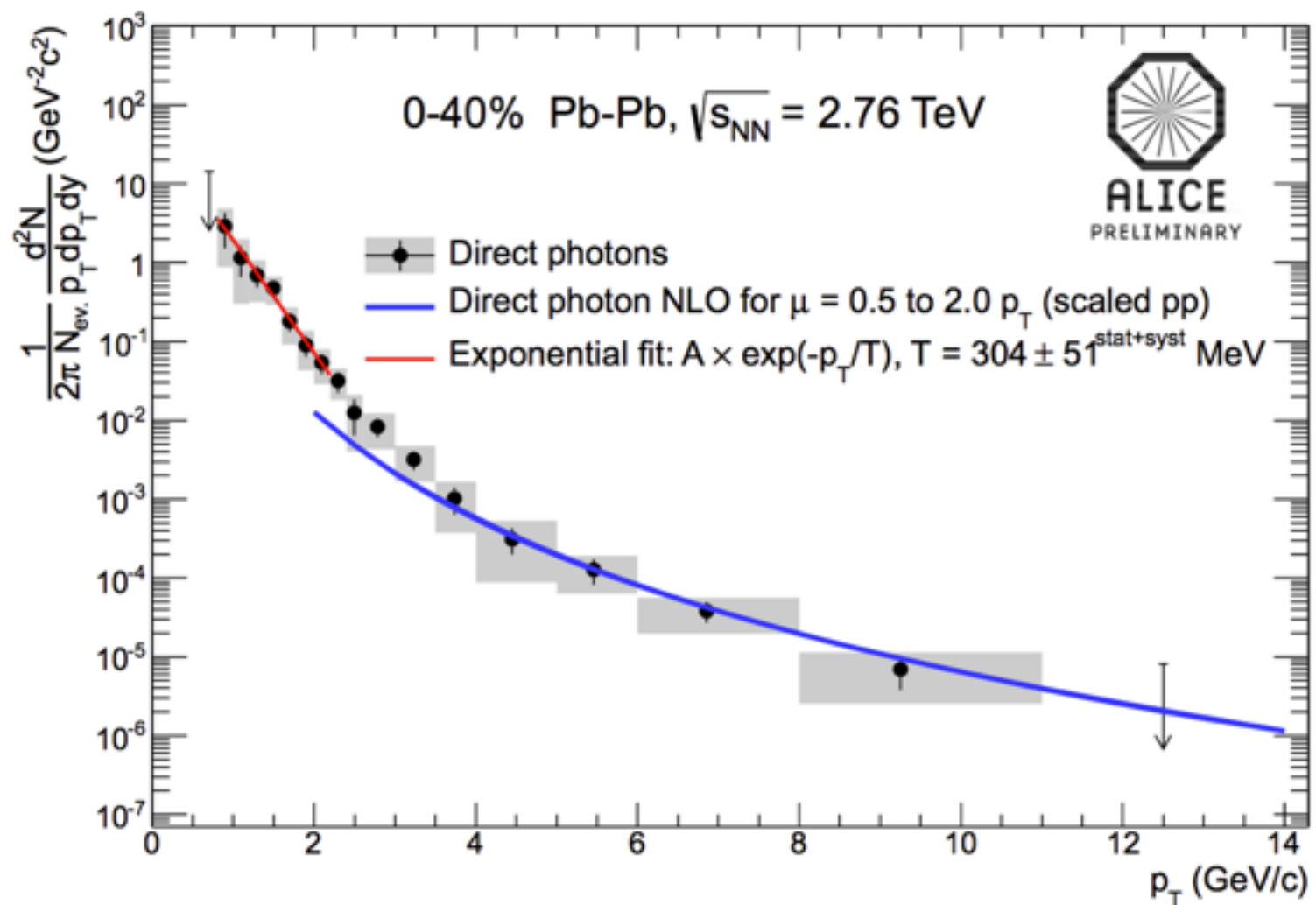
RHIC



0 – 20%

$$T = 221 \pm 19 \pm 19 \text{ MeV}$$

LHC



fit: $A \exp(-p_T/T)$

$$T = 304 \pm 51^{\text{stat+sys}} \text{ MeV}$$



What does this T mean



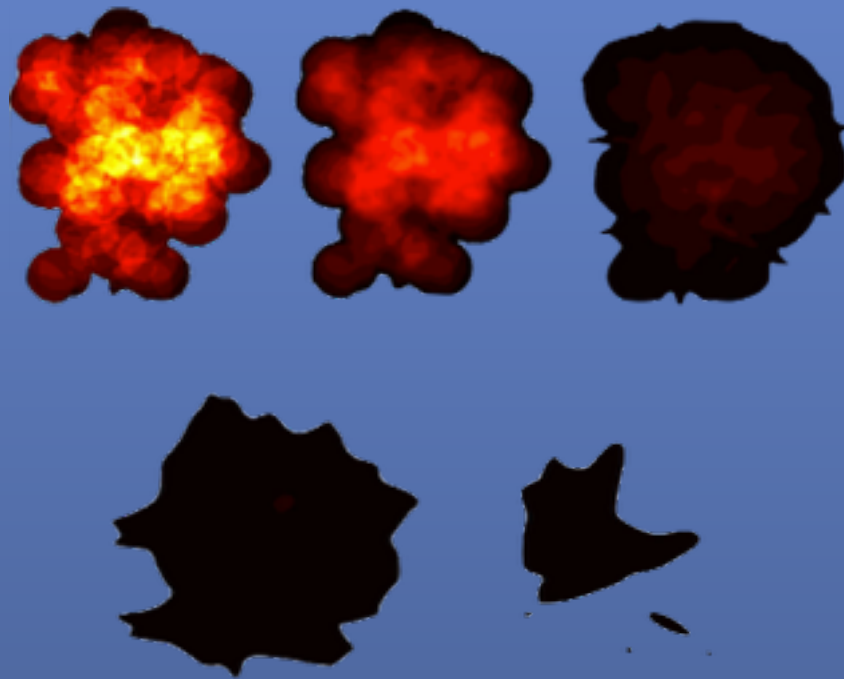
State-of-the-art hydrodynamic modeling

Initial Condition
Generators
(MC-KLN, MC-Glauber)

[https://github.com/
chunshen1987/iEBE.git](https://github.com/chunshen1987/iEBE.git)

Thermal Photon
Emission Rates

Hydrodynamic
Simulations
(VISH2+1)



HydroInfo
Package

$$e, s, p, T, \\ u^\mu, \pi^{\mu\nu}$$

Thermal Photon
Interface

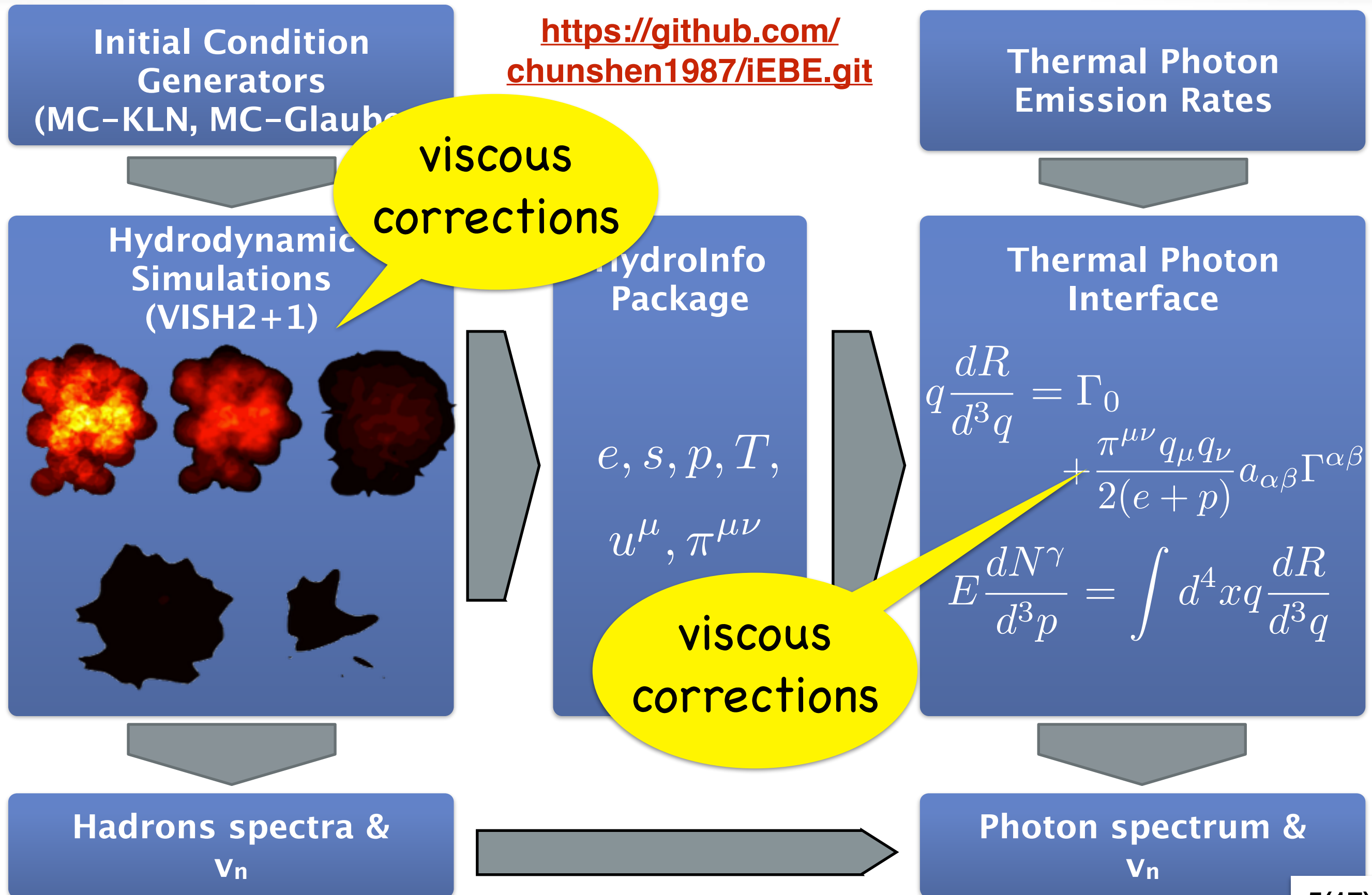
$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} q_\mu q_\nu}{2(e + p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$

$$E \frac{dN^\gamma}{d^3p} = \int d^4x q \frac{dR}{d^3q}$$

Hadrons spectra &
 V_n

Photon spectrum &
 V_n

State-of-the-art hydrodynamic modeling



Viscous Photon Emission Rates: General Formalism

Thermal photon emission rates can be calculated by

$$E_q \frac{dR}{d^3q} = \int \frac{d^3p_1}{2E_1(2\pi)^3} \frac{d^3p_2}{2E_2(2\pi)^3} \frac{d^3p_3}{2E_3(2\pi)^3} \frac{1}{2(2\pi)^3} |\mathcal{M}|^2 \\ \times f_1(p_1^\mu) f_2(p_2^\mu) (1 \pm f_3(p_3^\mu)) (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - q)$$

With

$$f(p^\mu) = f_0(E) + f_0(E)(1 \pm f_0(E)) \frac{\pi^{\mu\nu} \hat{p}_\mu \hat{p}_\nu}{2(e + p)} \chi\left(\frac{p}{T}\right)$$

We can expand photon emission rates around the thermal equilibrium:

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e + p)} a_{\alpha\beta} \Gamma^{\alpha\beta},$$

$$a_{\mu\nu} = \frac{3}{2(u \cdot \hat{q})^4} \hat{q}_\mu \hat{q}_\nu + \frac{1}{(u \cdot \hat{q})^2} u_\mu u_\nu + \frac{1}{2(u \cdot \hat{q})^2} g_{\mu\nu} - \frac{3}{2(u \cdot \hat{q})^3} (\hat{q}_\mu u_\nu + \hat{q}_\nu u_\mu).$$

Viscous Photon Emission Rates: General Formalism

Thermal photon emission rates can be calculated by

$$E_q \frac{dR}{d^3q} = \int \frac{d^3p_1}{2E_1(2\pi)^3} \frac{d^3p_2}{2E_2(2\pi)^3} \frac{d^3p_3}{2E_3(2\pi)^3} \frac{1}{2(2\pi)^3} |\mathcal{M}|^2 \\ \times f_1(p_1^\mu) f_2(p_2^\mu) (1 \pm f_3(p_3^\mu)) (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - q)$$

With

$$f(p^\mu) = \frac{\Gamma_0(q, T) + a_{\alpha\beta} \Gamma^{\alpha\beta}(q, T) \pi^{\mu\nu} \hat{p}_\mu \hat{p}_\nu}{2(e + p)} \chi\left(\frac{p}{T}\right)$$

We can expand Γ_0 and $\Gamma^{\alpha\beta}$ calculated in fluid local rest frame and the thermal equilibrium:

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e + p)} a_{\alpha\beta} \Gamma^{\alpha\beta},$$

$$a_{\mu\nu} = \frac{3}{2(u \cdot \hat{q})^4} \hat{q}_\mu \hat{q}_\nu + \frac{1}{(u \cdot \hat{q})^2} u_\mu u_\nu + \frac{1}{2(u \cdot \hat{q})^2} (\delta_{\mu\nu} - \hat{q}_\mu \hat{q}_\nu - u_\mu u_\nu).$$

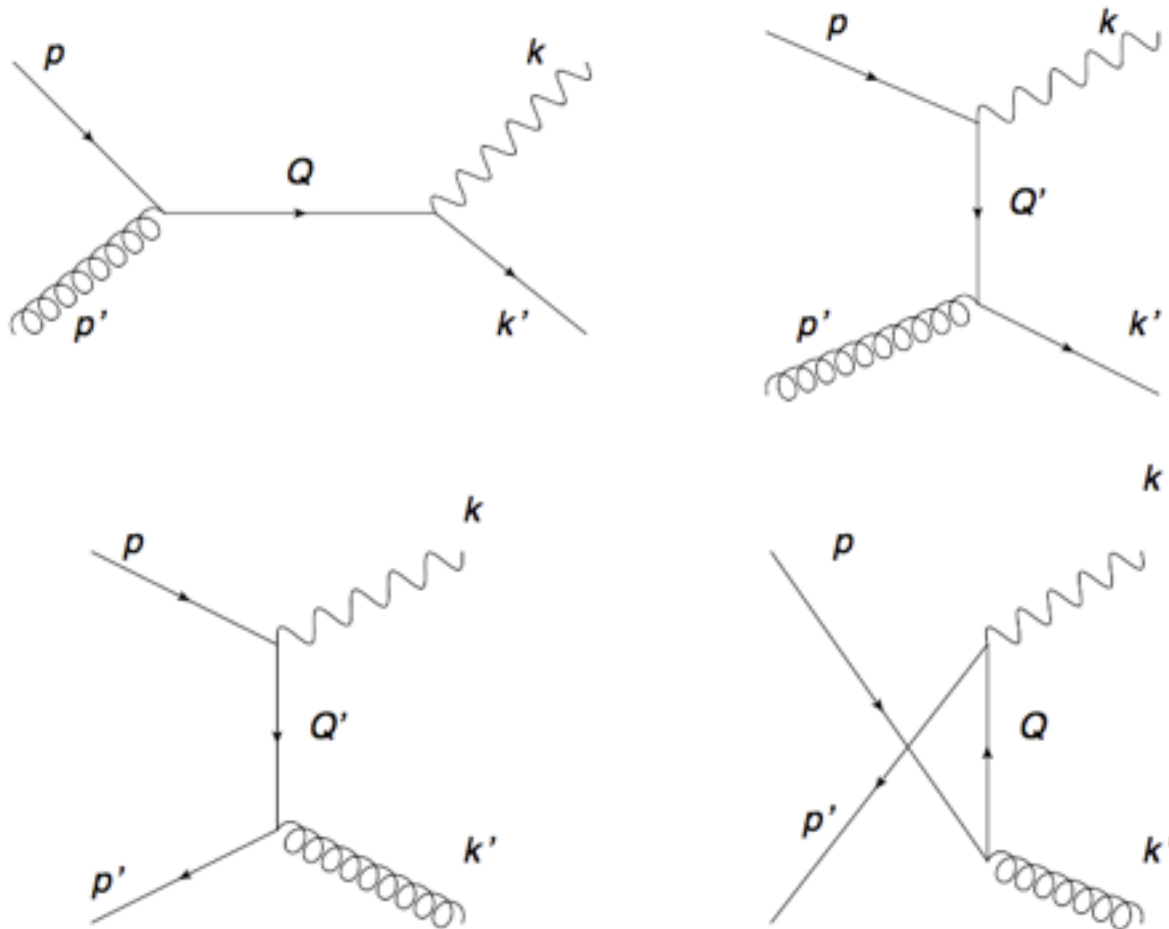
calculated in lab frame

Viscous Photon Emission Rates: General Formalism

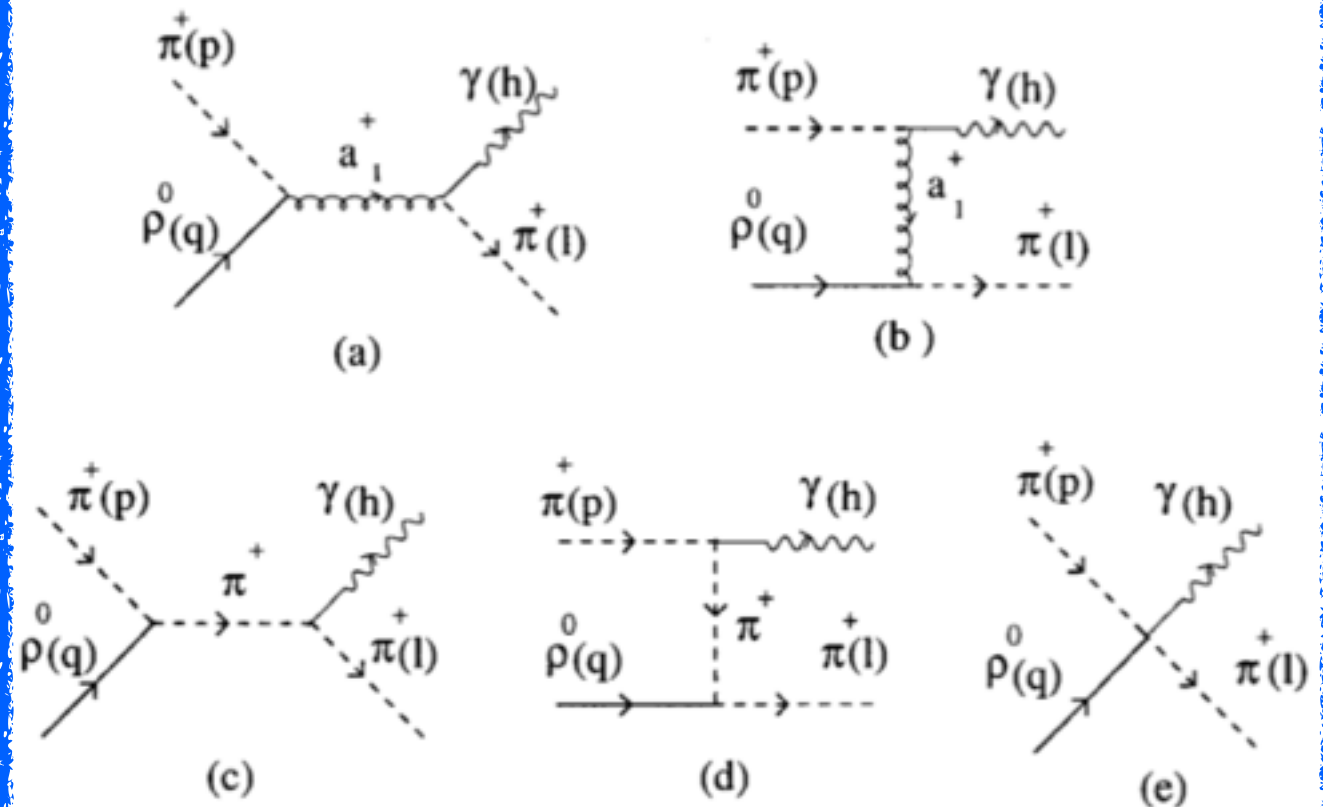
$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e + p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$

Equilibrium rates

QGP



Hadron Gas



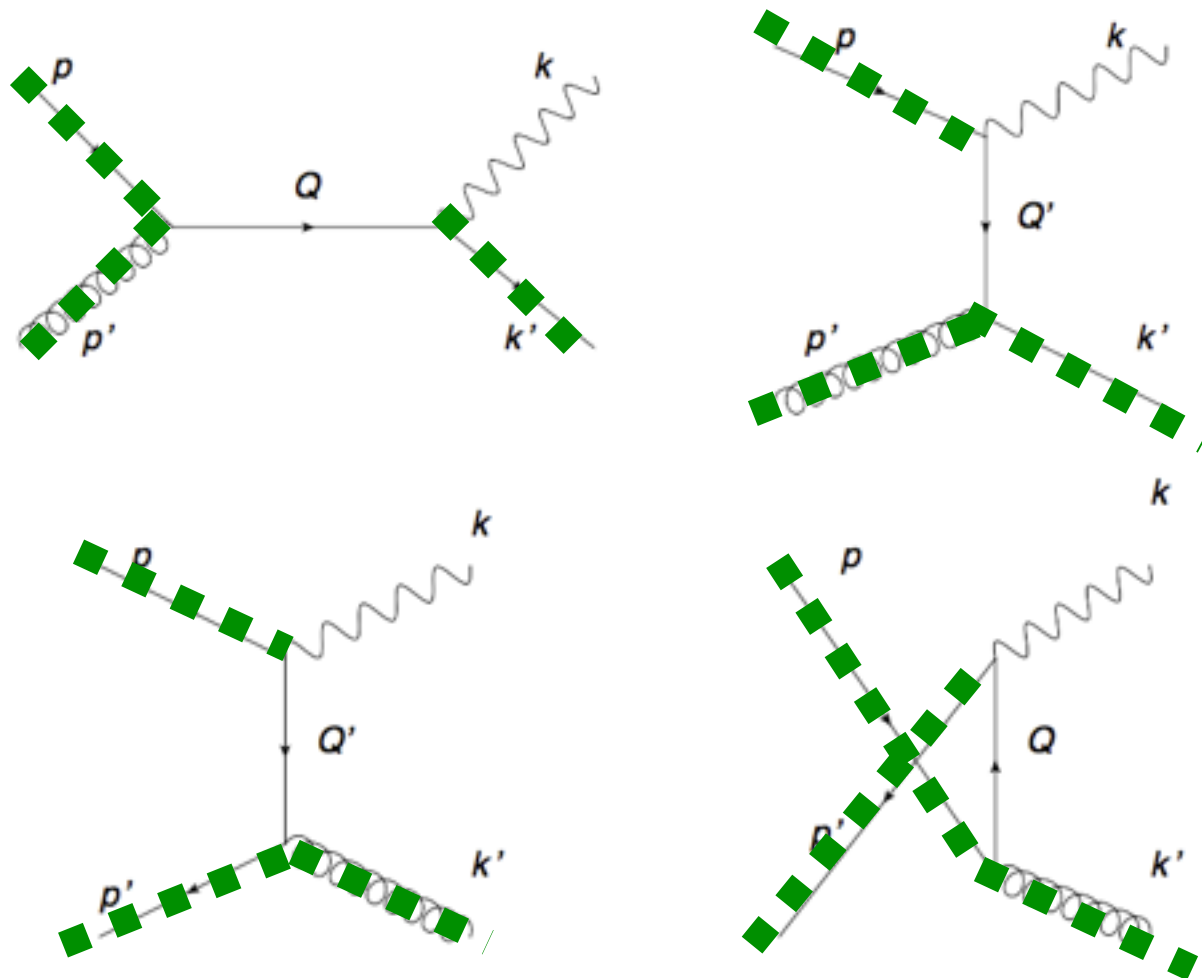
Viscous Photon Emission Rates: General Formalism

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e + p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$

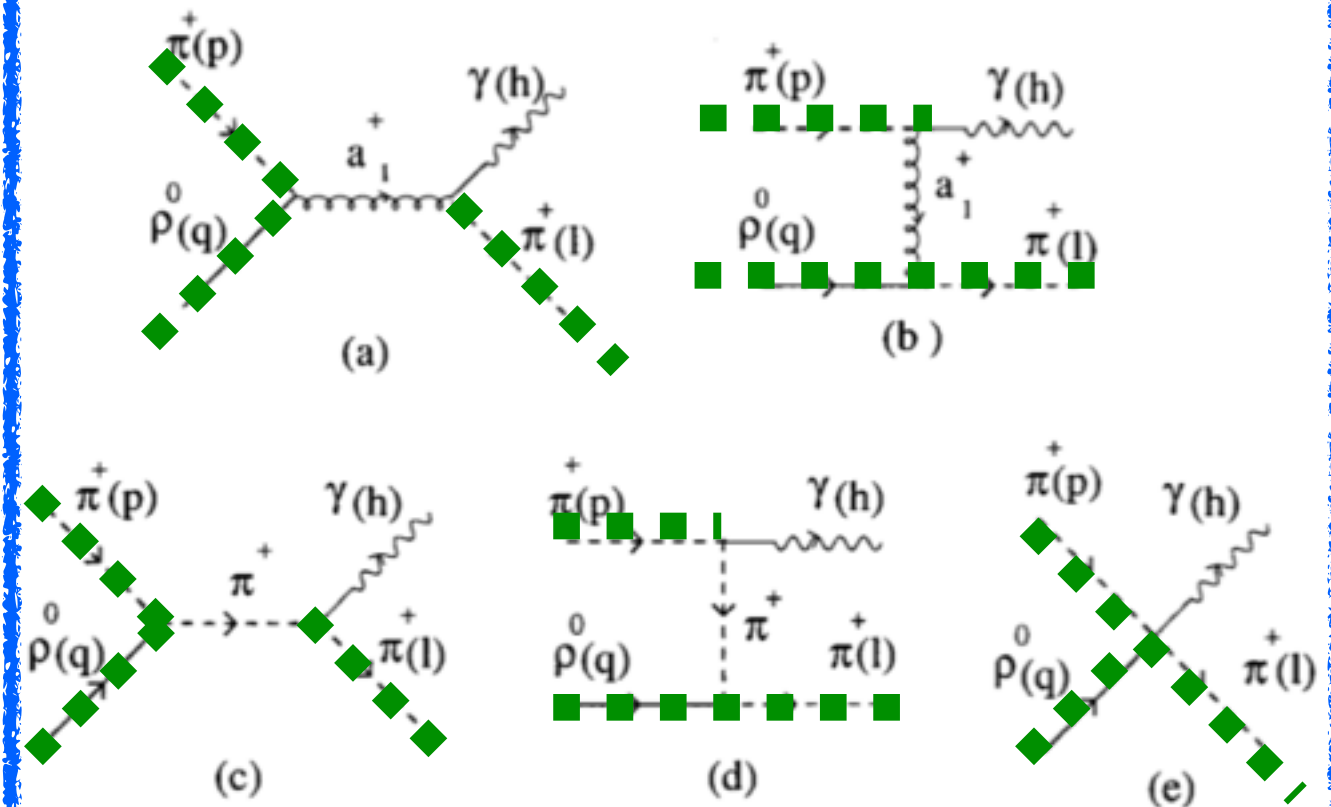
Equilibrium rates

off-equilibrium δf corrections

QGP



Hadron Gas



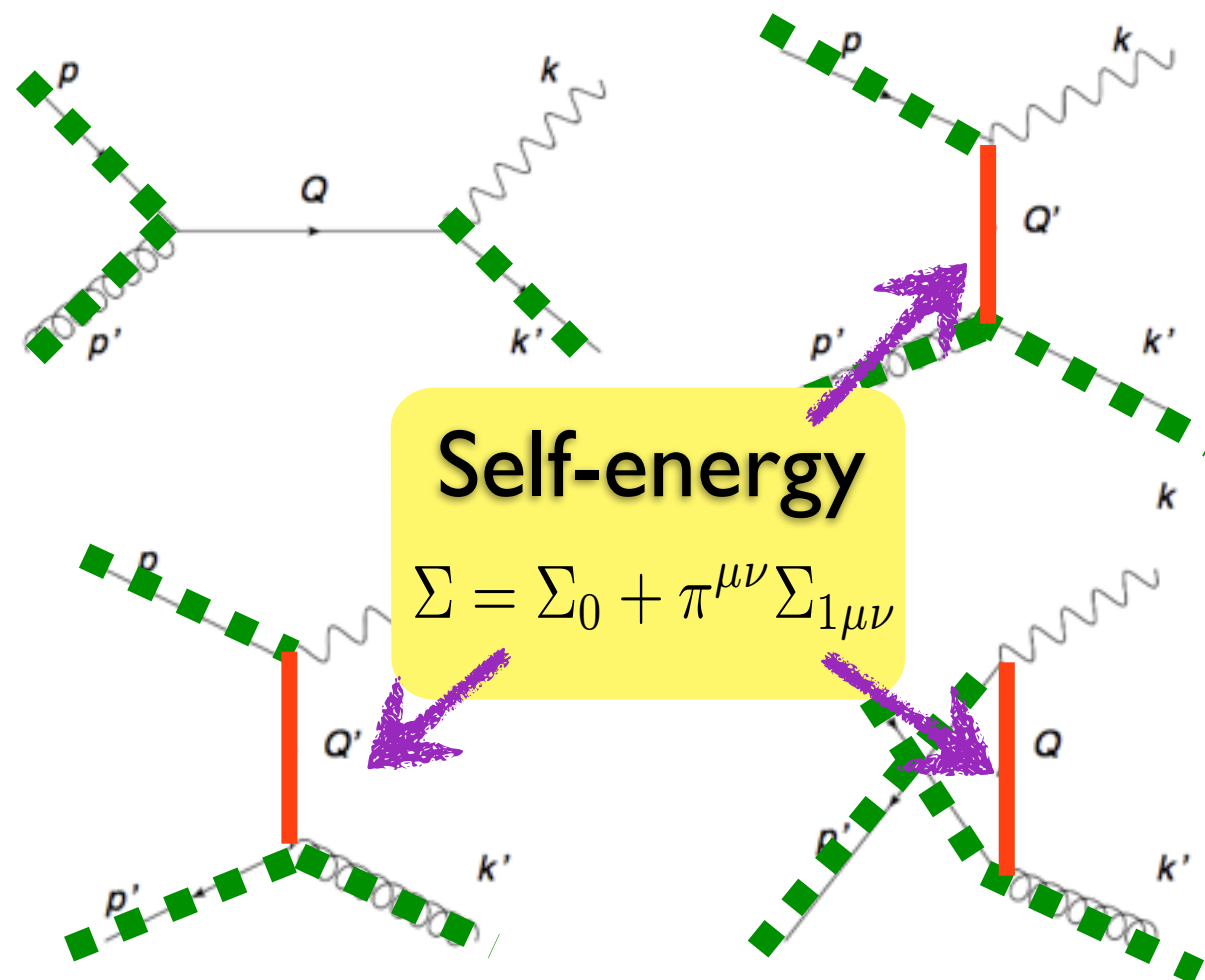
Viscous Photon Emission Rates: General Formalism

$$q \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu} \hat{q}_\mu \hat{q}_\nu}{2(e+p)} a_{\alpha\beta} \Gamma^{\alpha\beta}$$

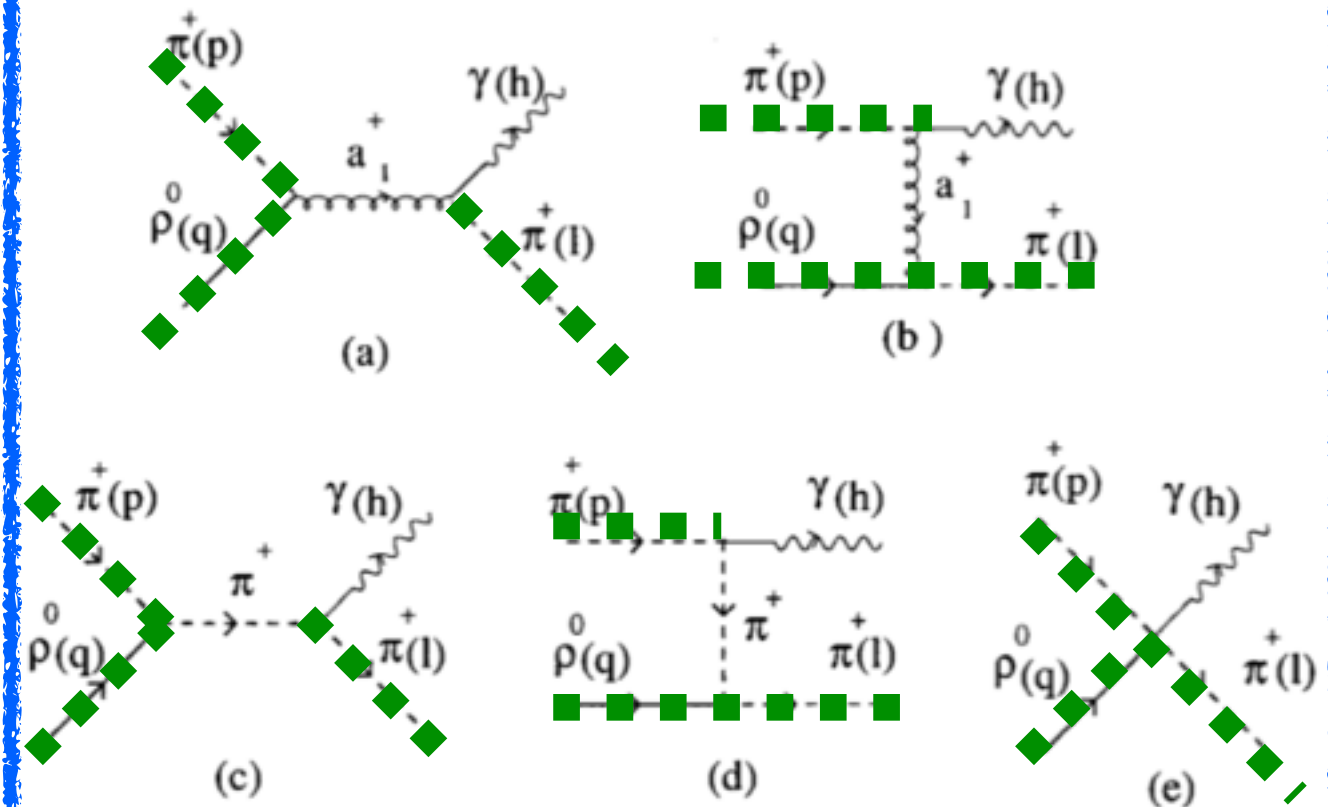
Equilibrium rates

off-equilibrium δf corrections

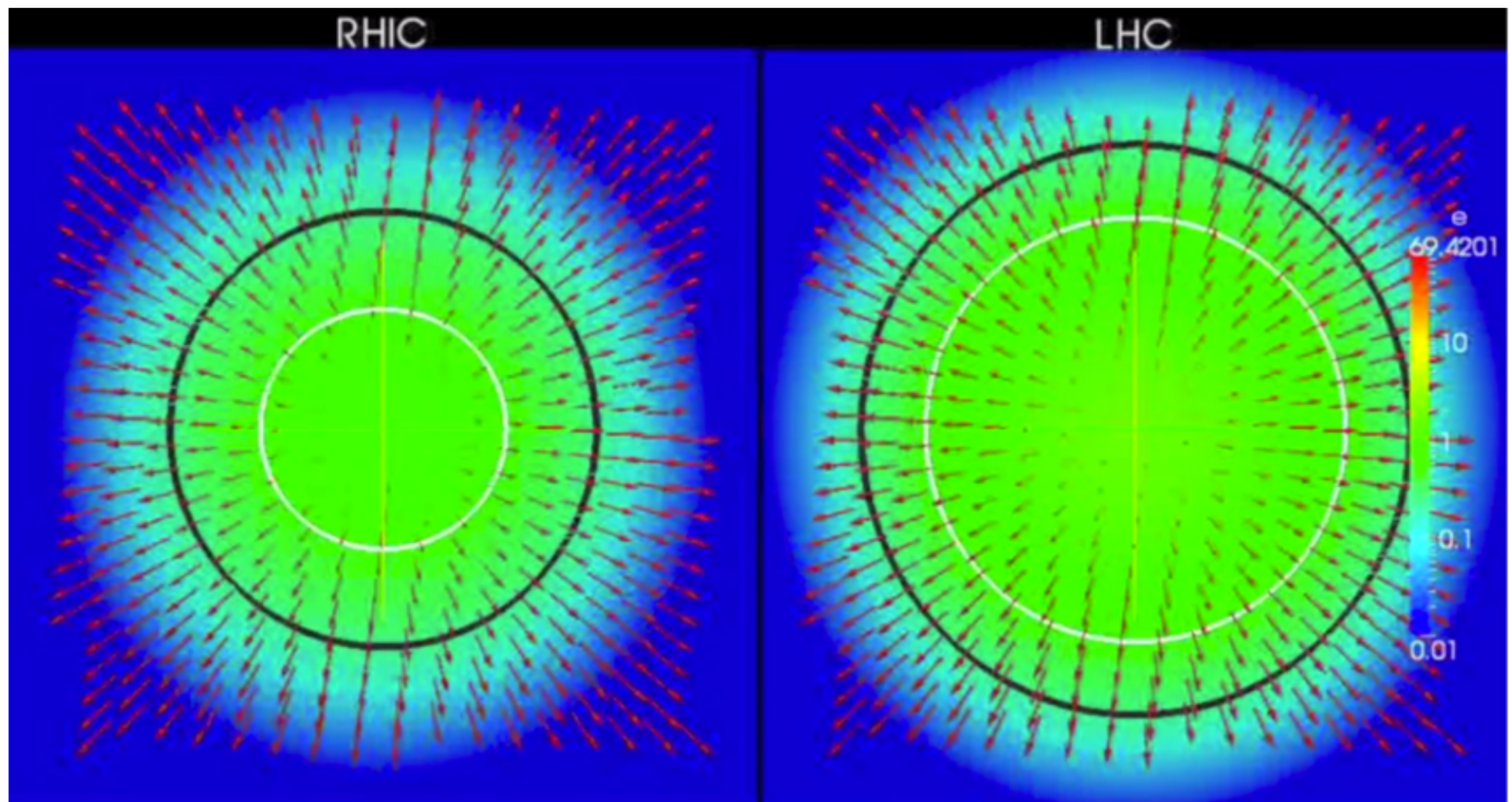
QGP



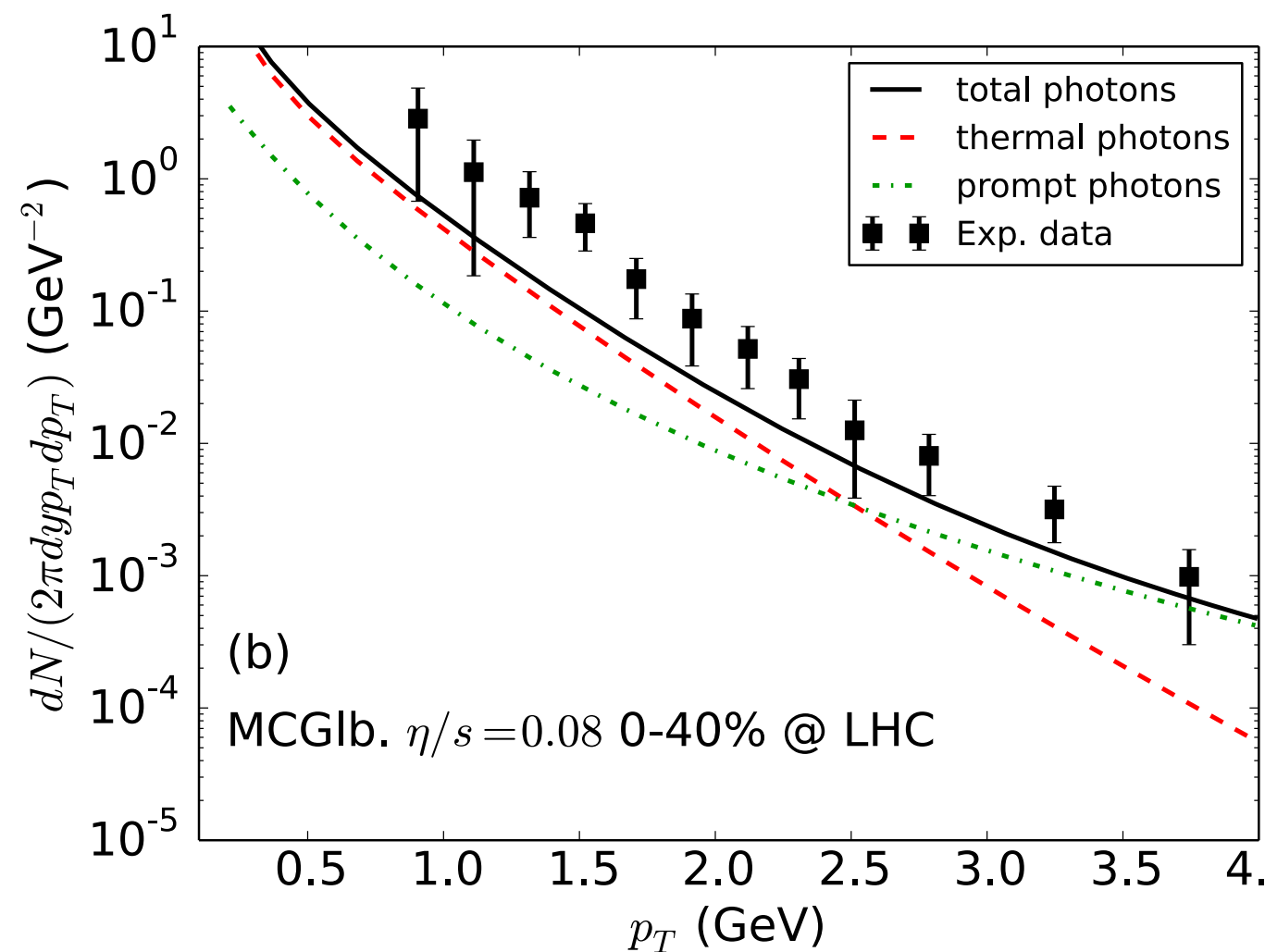
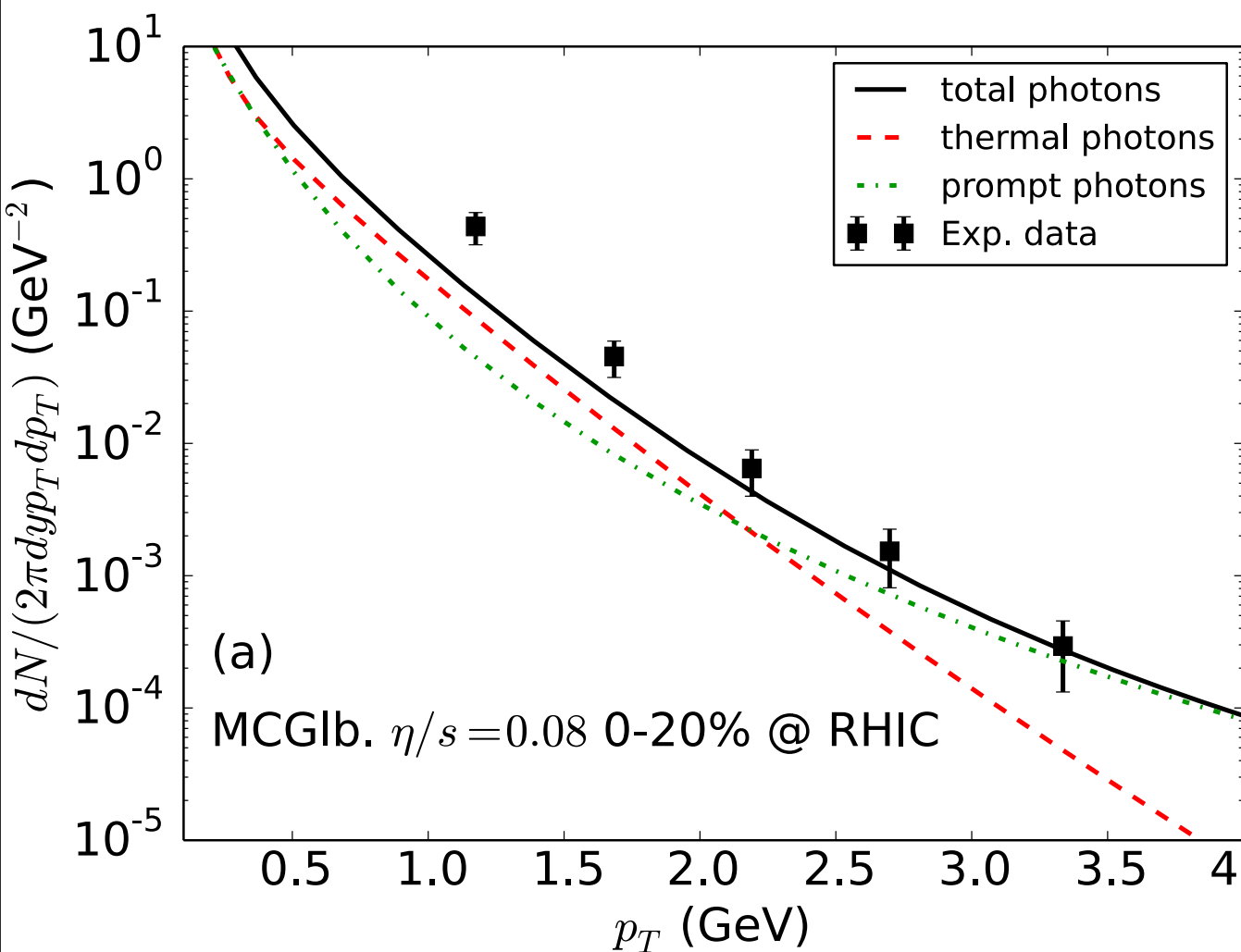
Hadron Gas



Photon spectra and radial flow

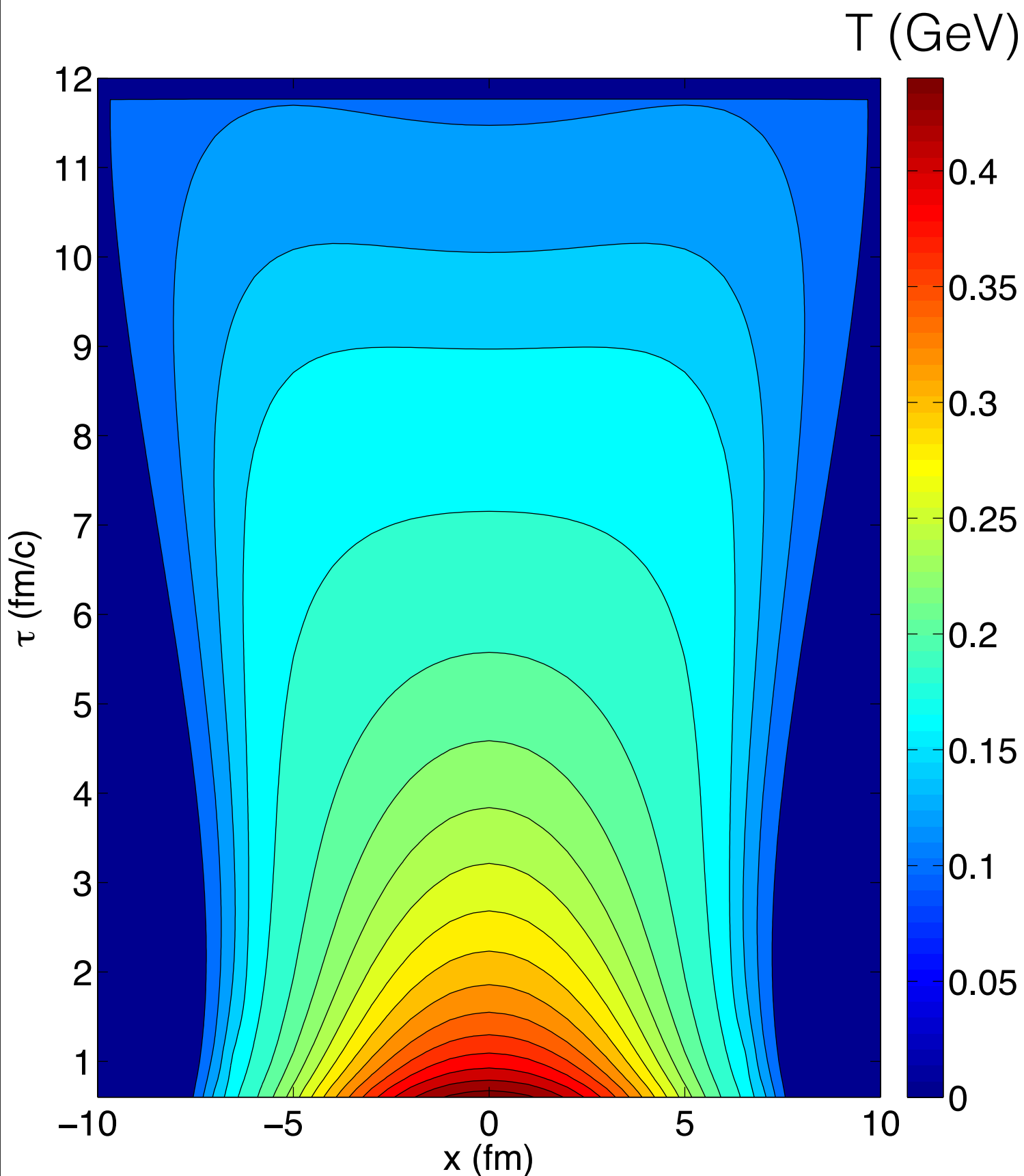


Thermal Photon Spectra

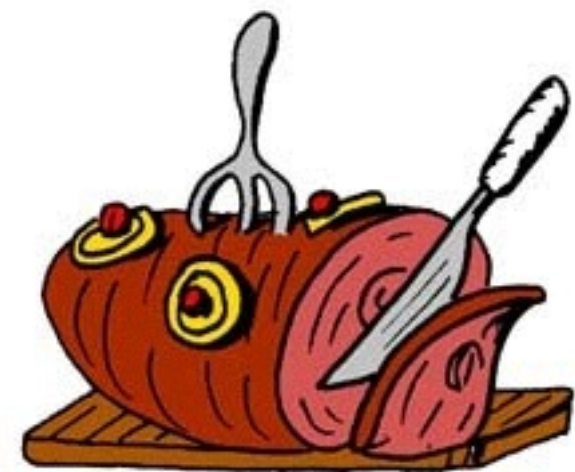


- With all available thermal emission sources, our current calculations still underestimate measured direct photon spectra at low p_T at both RHIC and LHC energies
- Additional emission sources need to be included to improve the agreement between theory and data

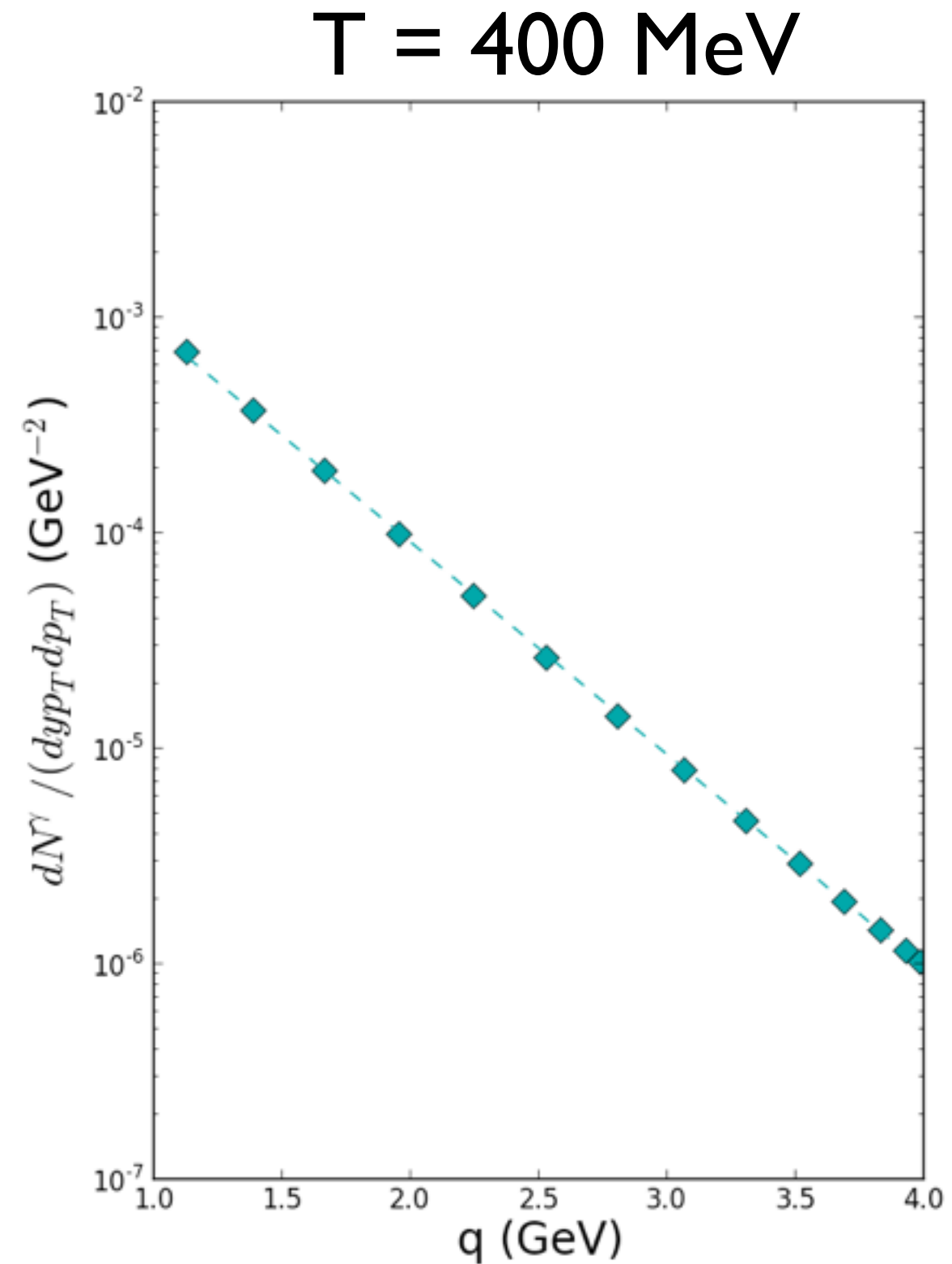
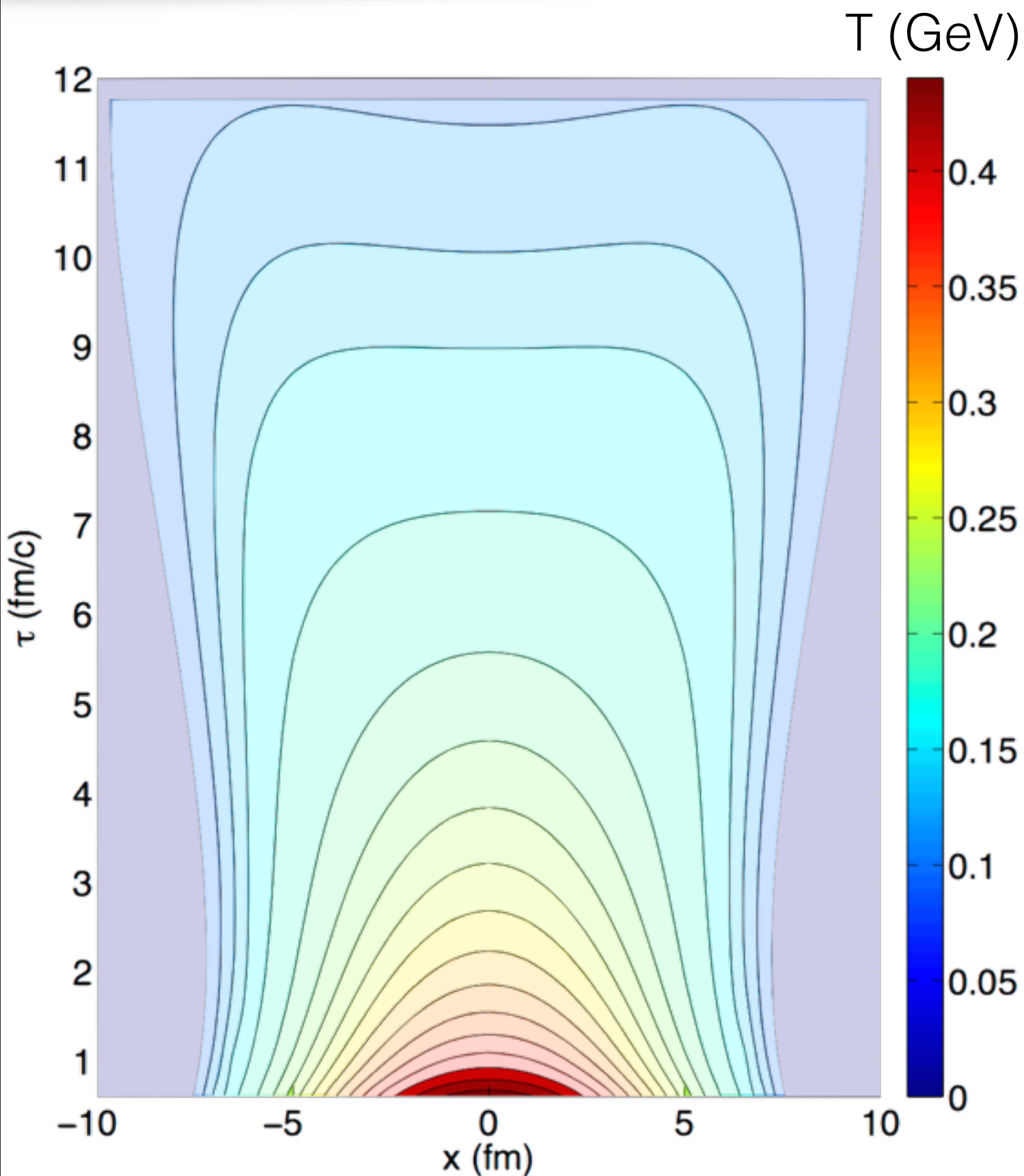
Slope of Photon Spectrum



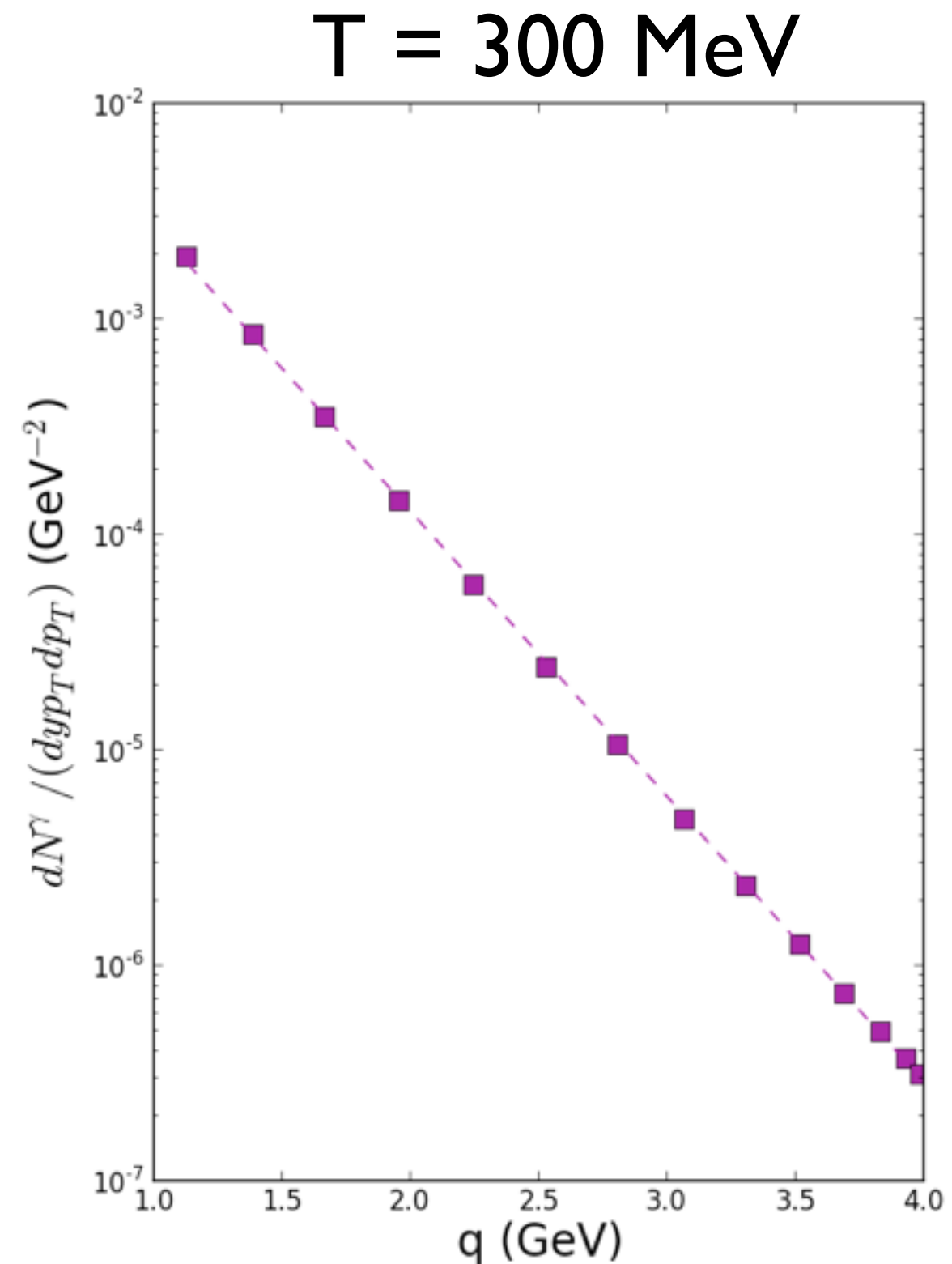
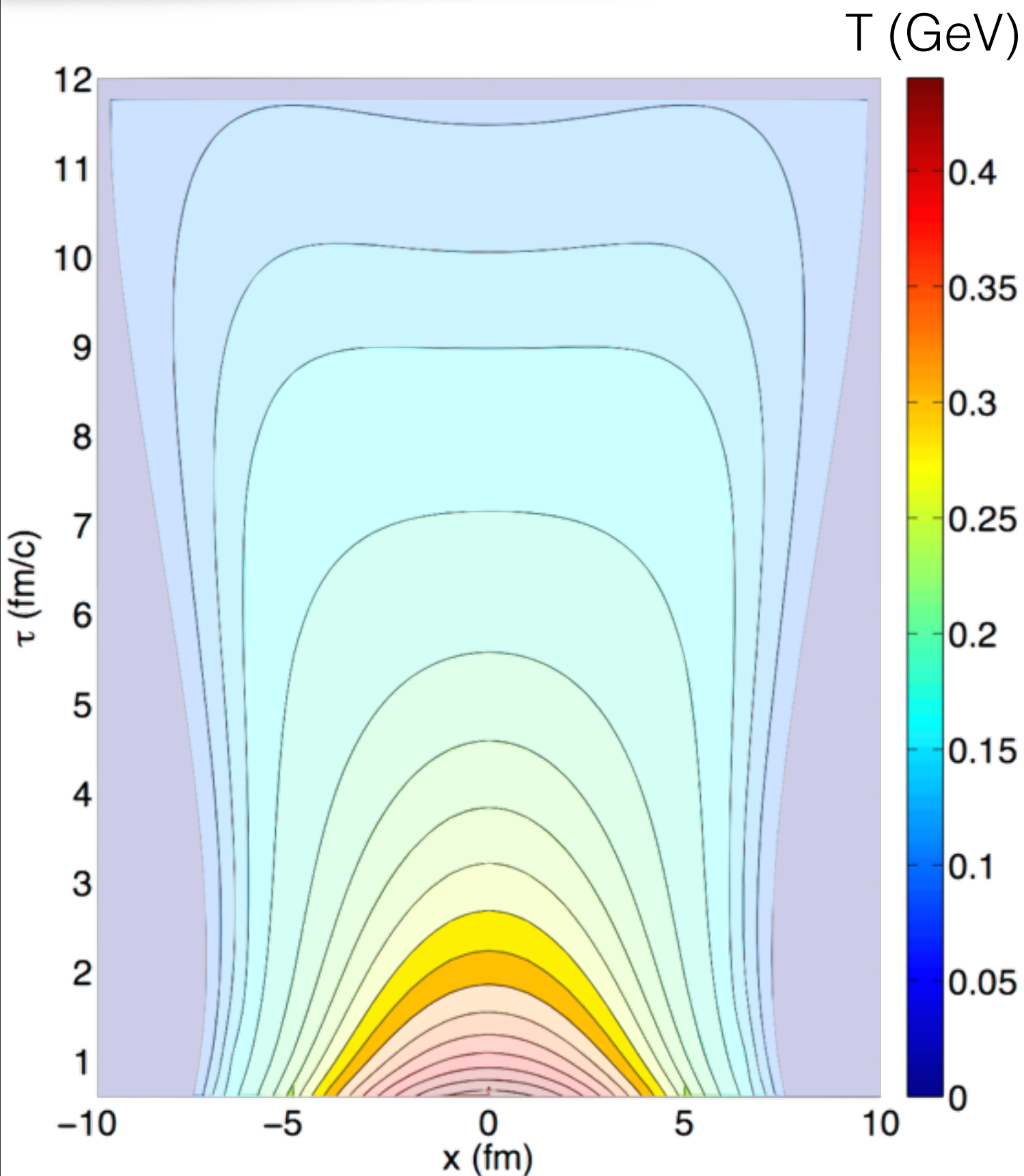
Slicing the
hydrodynamic
medium



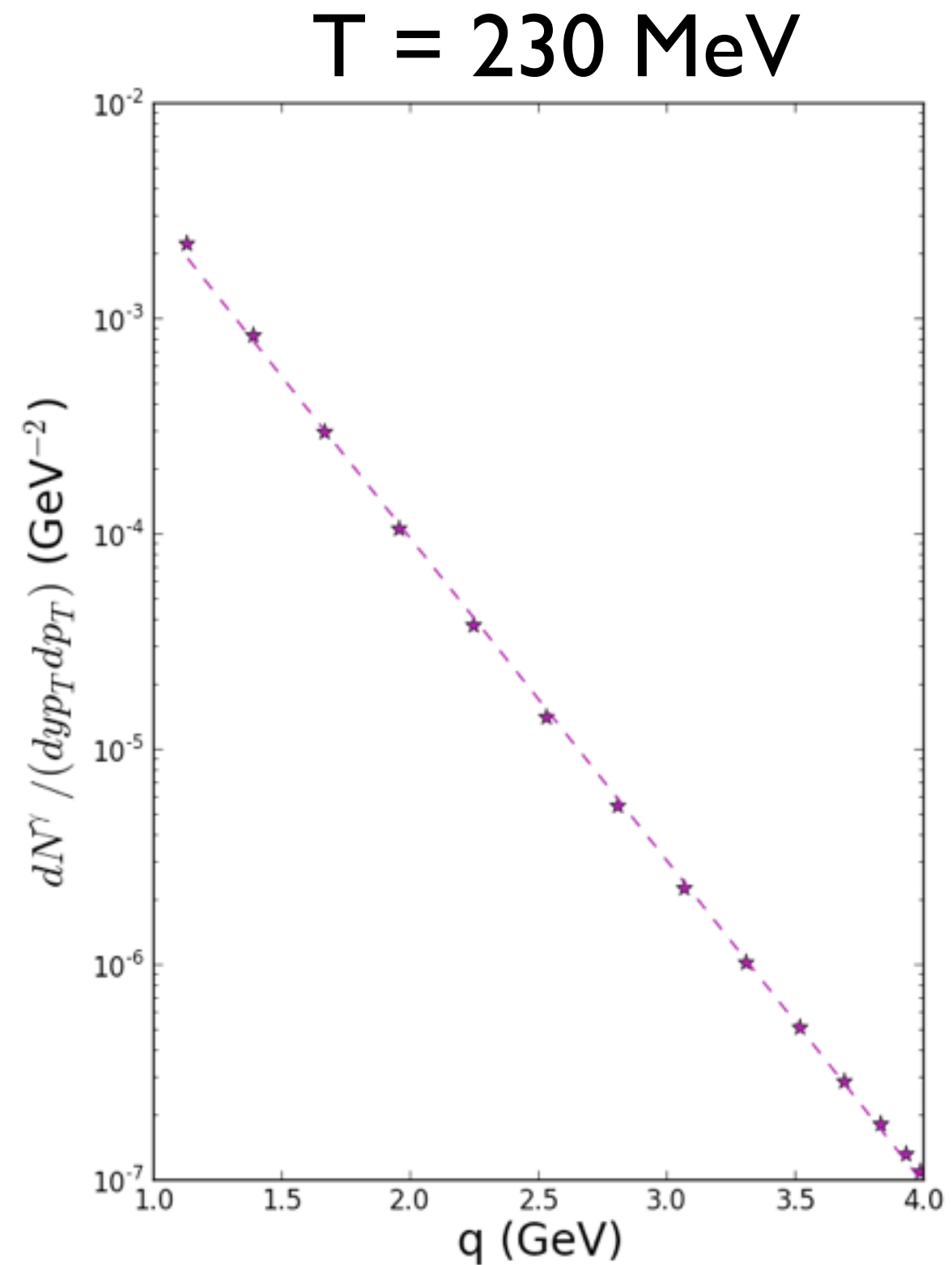
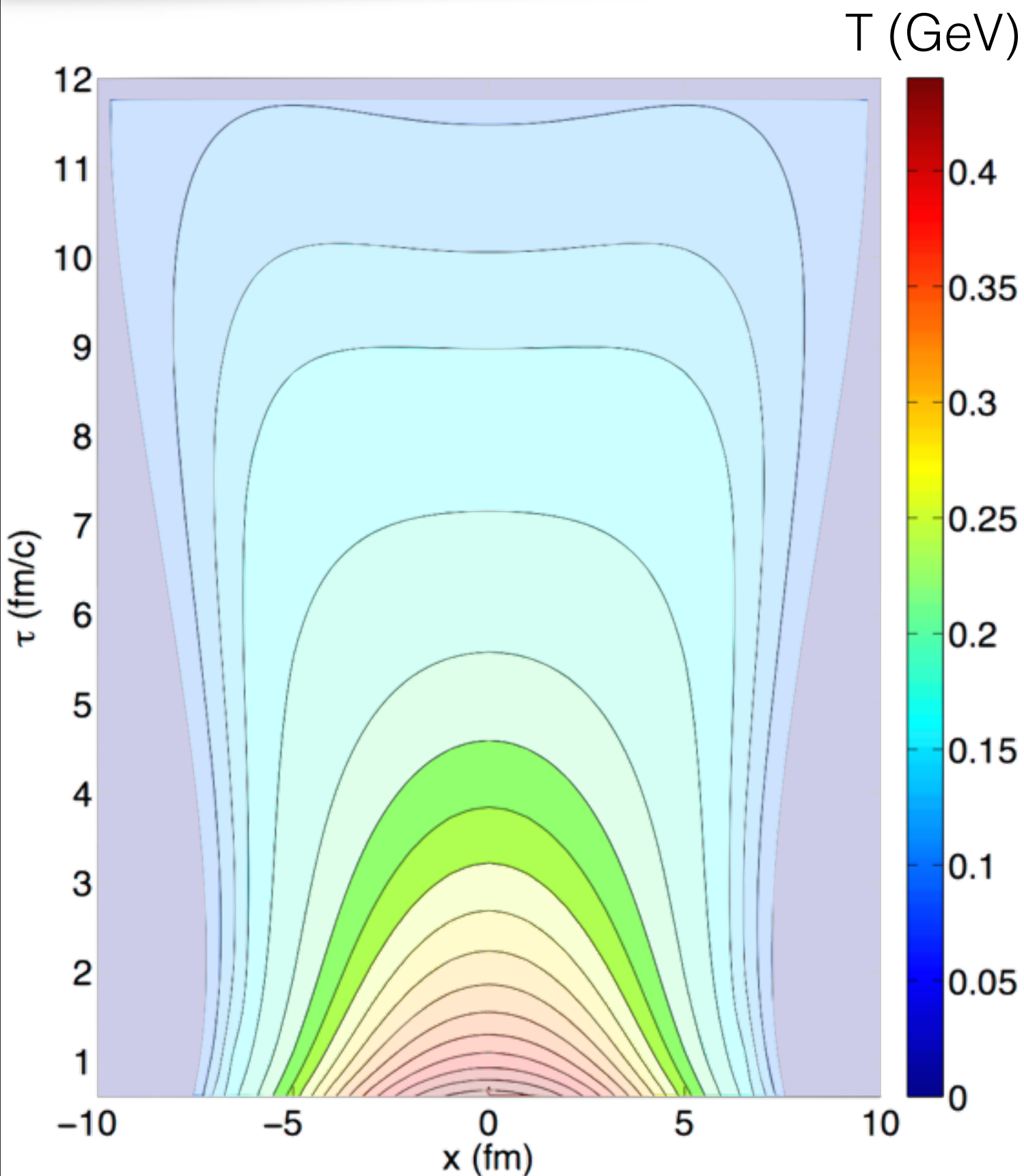
Slope of Photon Spectrum



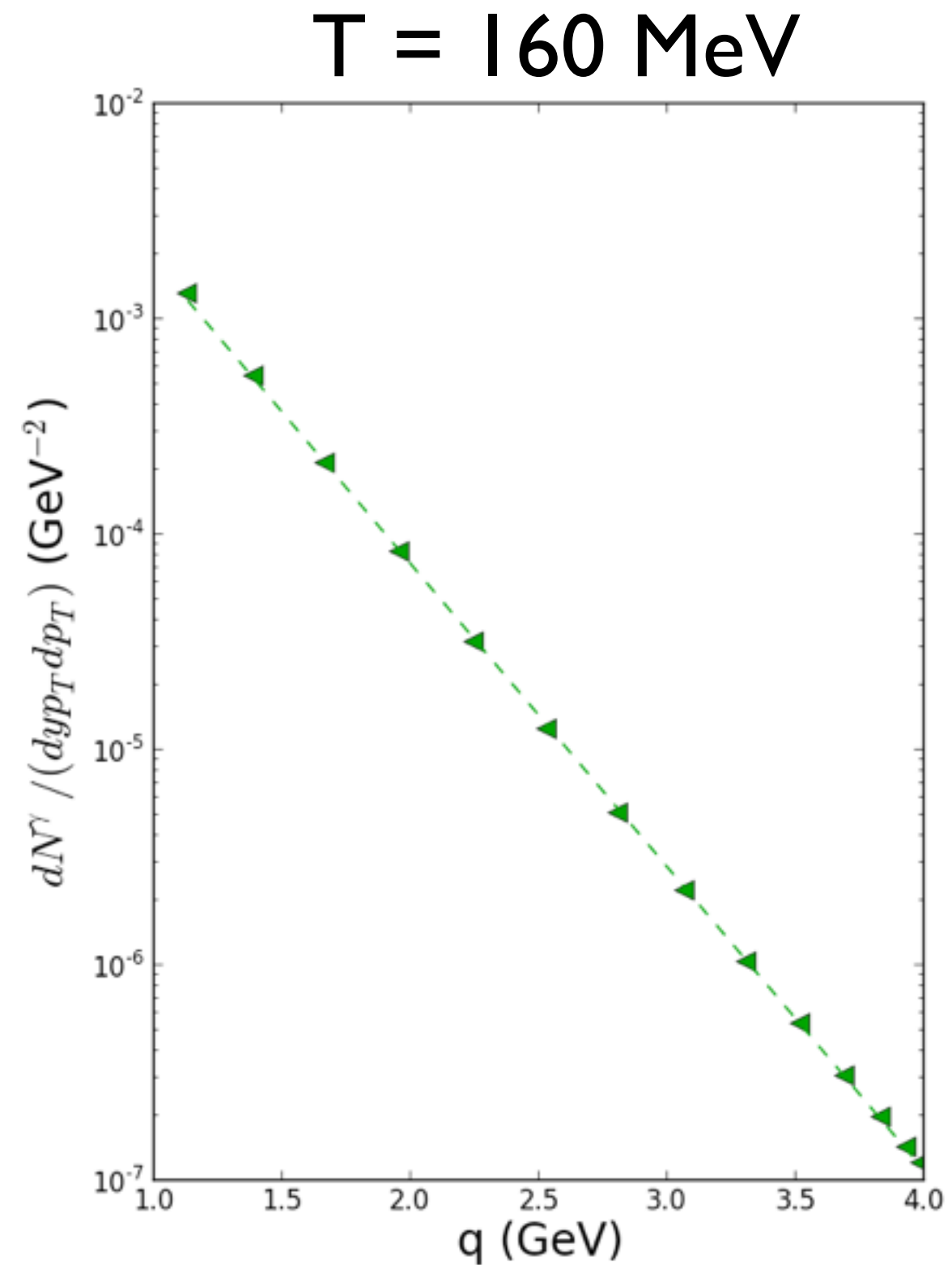
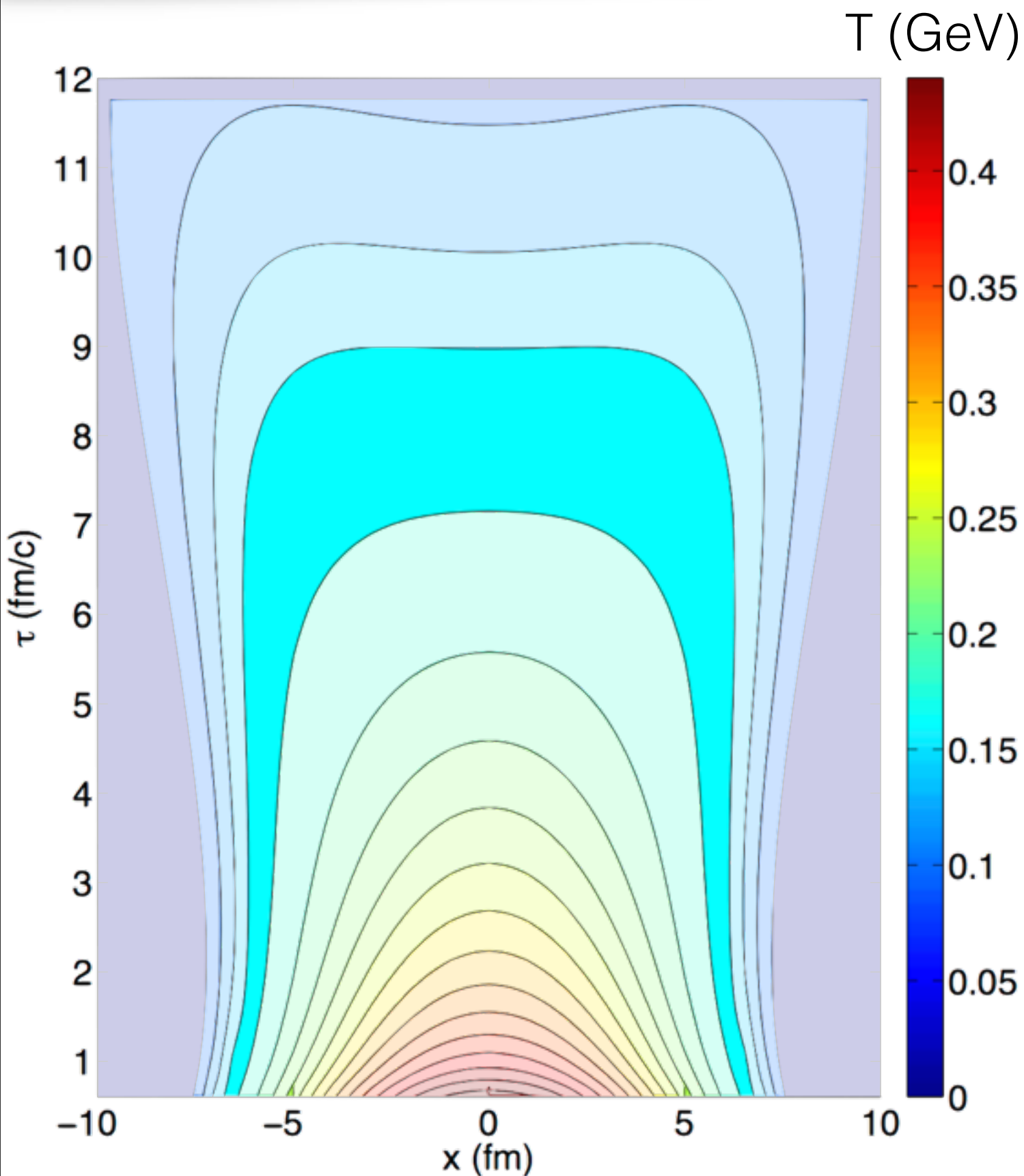
Slope of Photon Spectrum



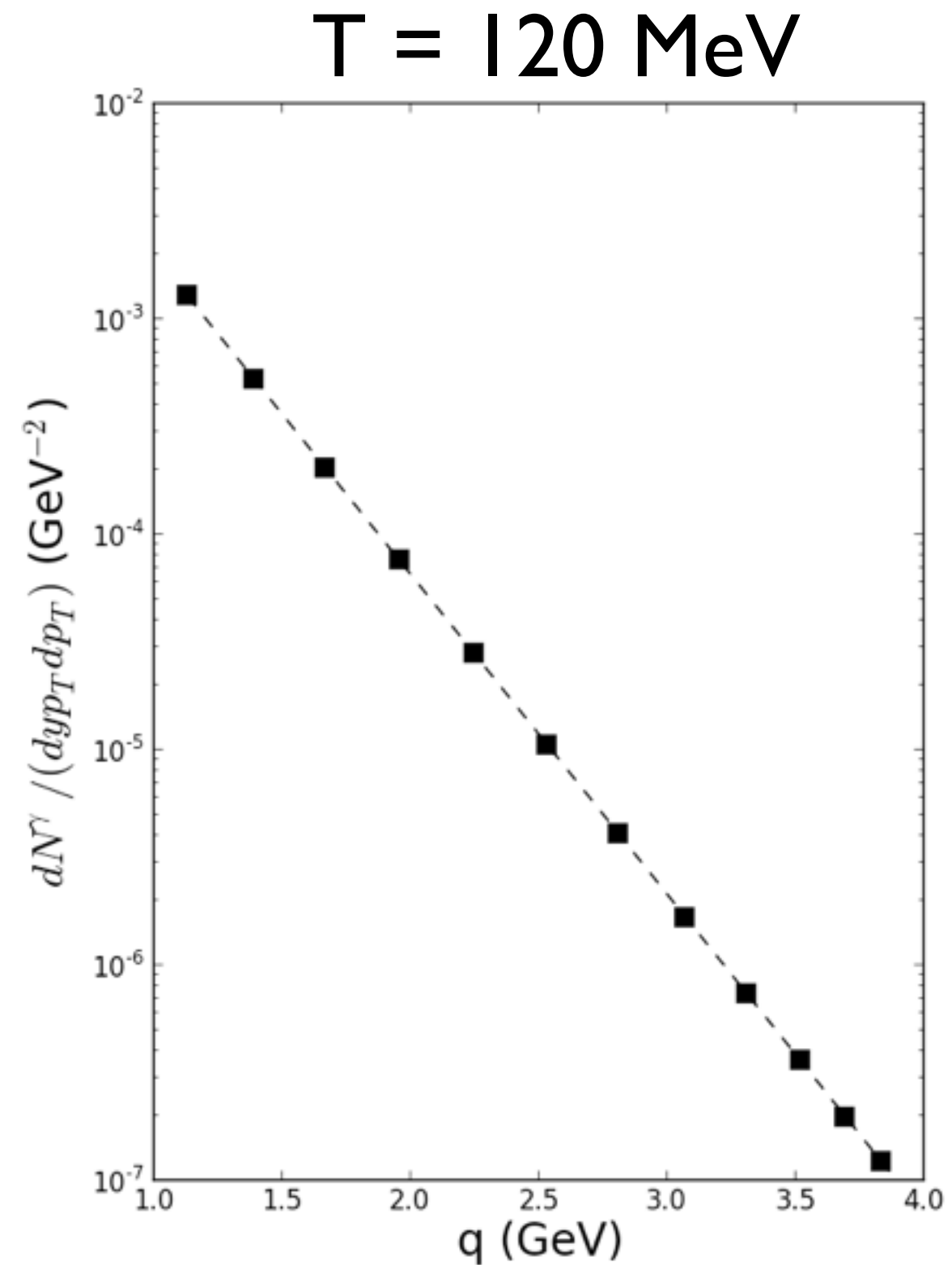
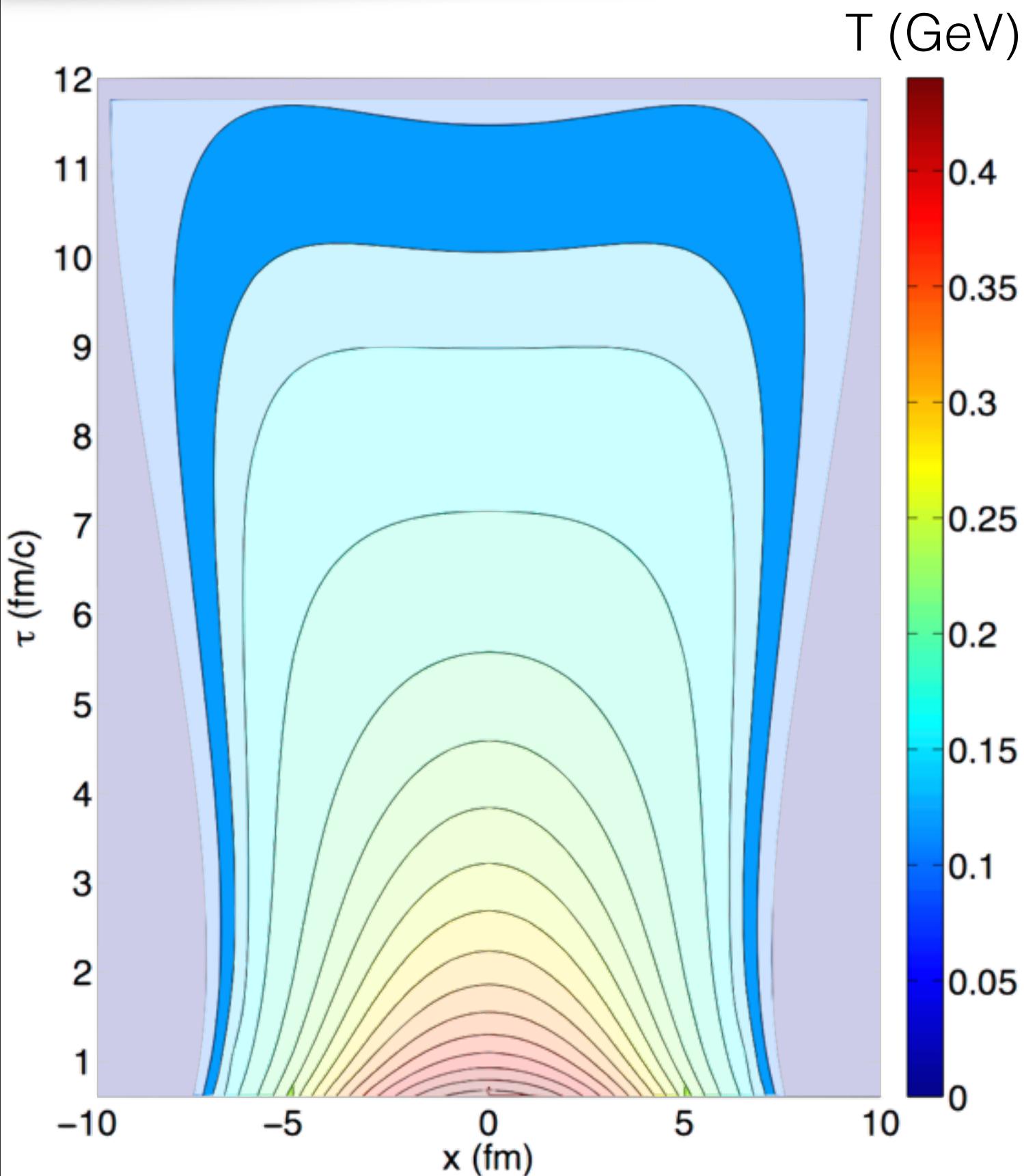
Slope of Photon Spectrum



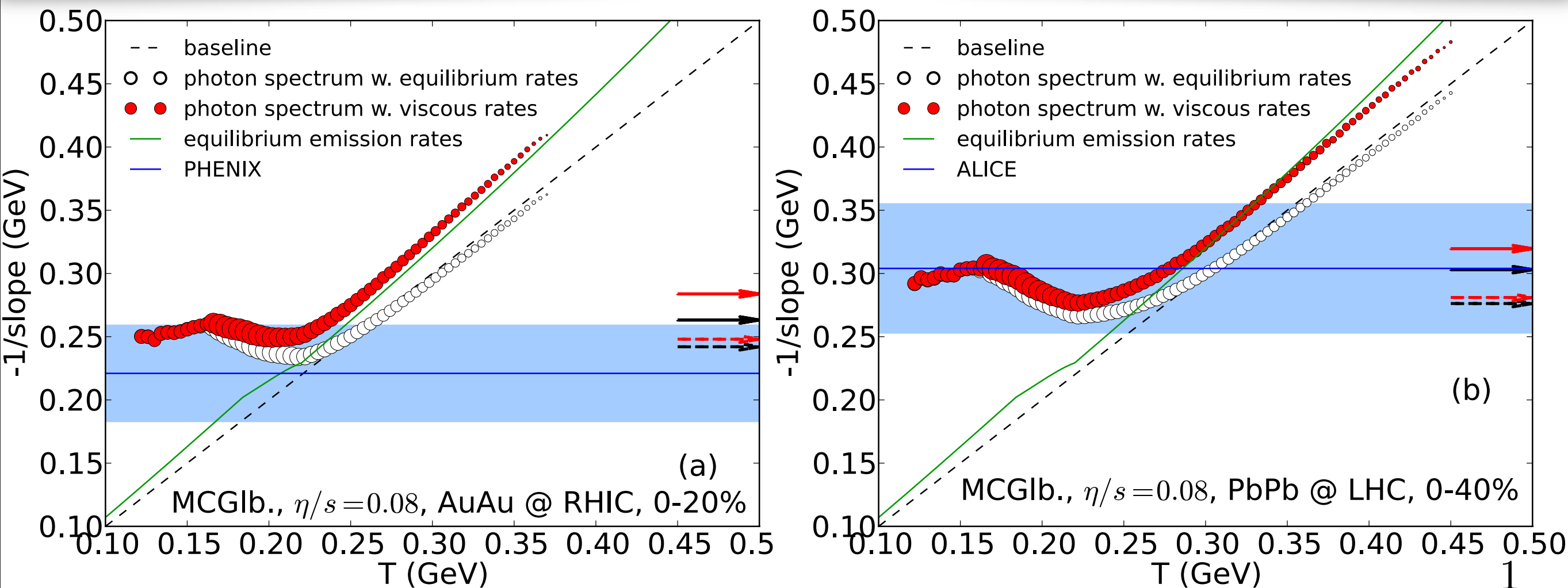
Slope of Photon Spectrum



Slope of Photon Spectrum



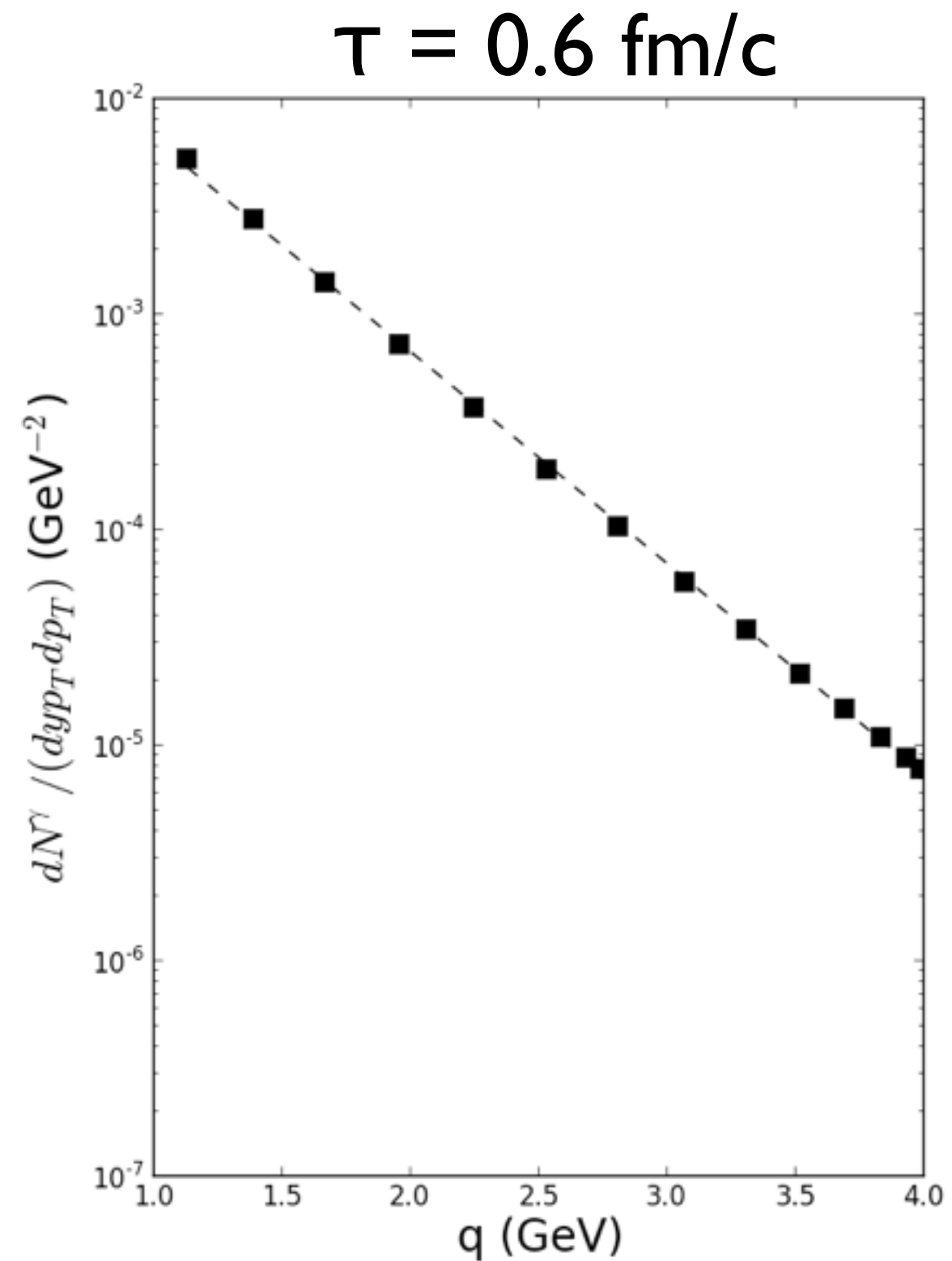
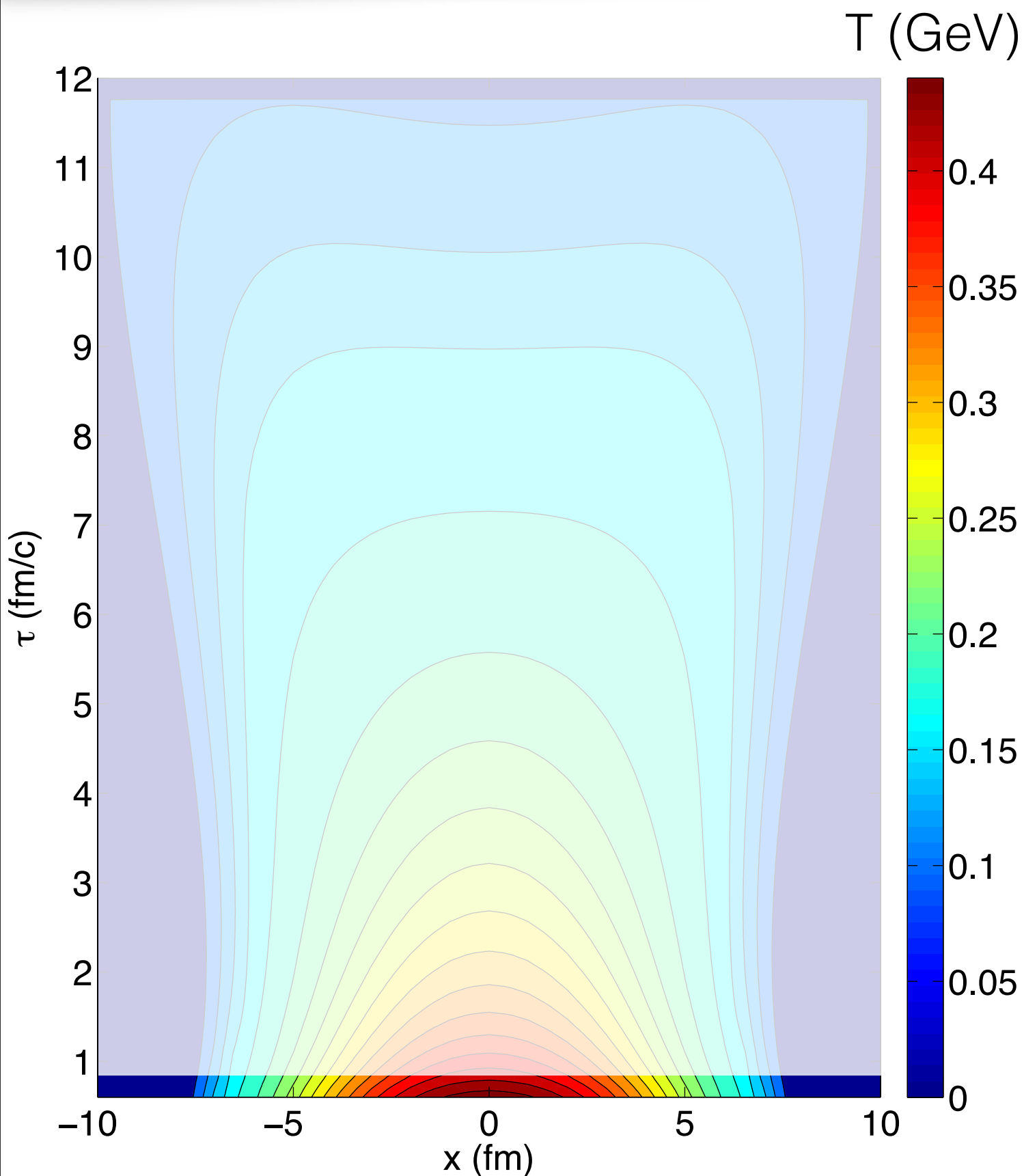
Fitted T_{eff} vs. True Temperature



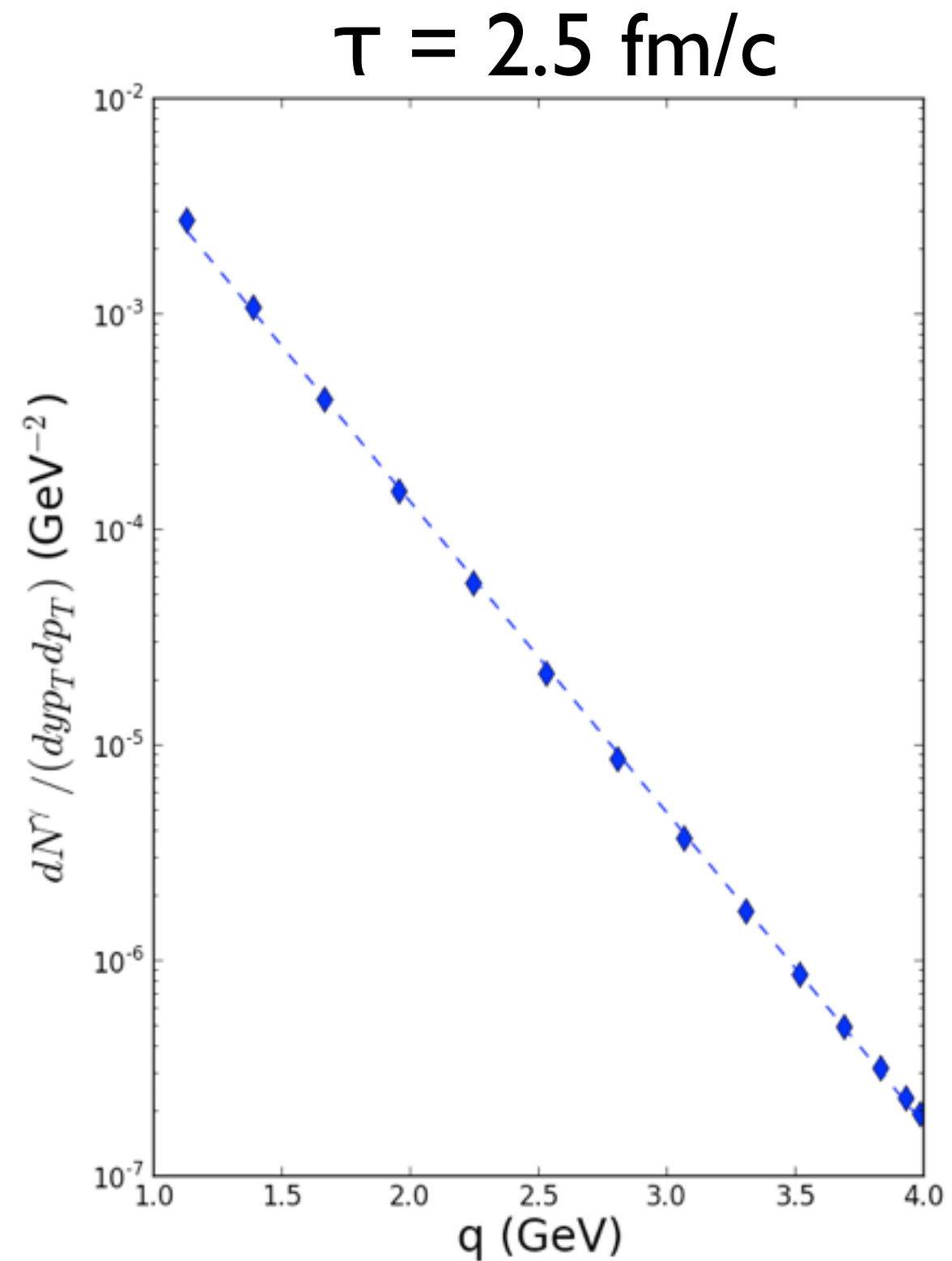
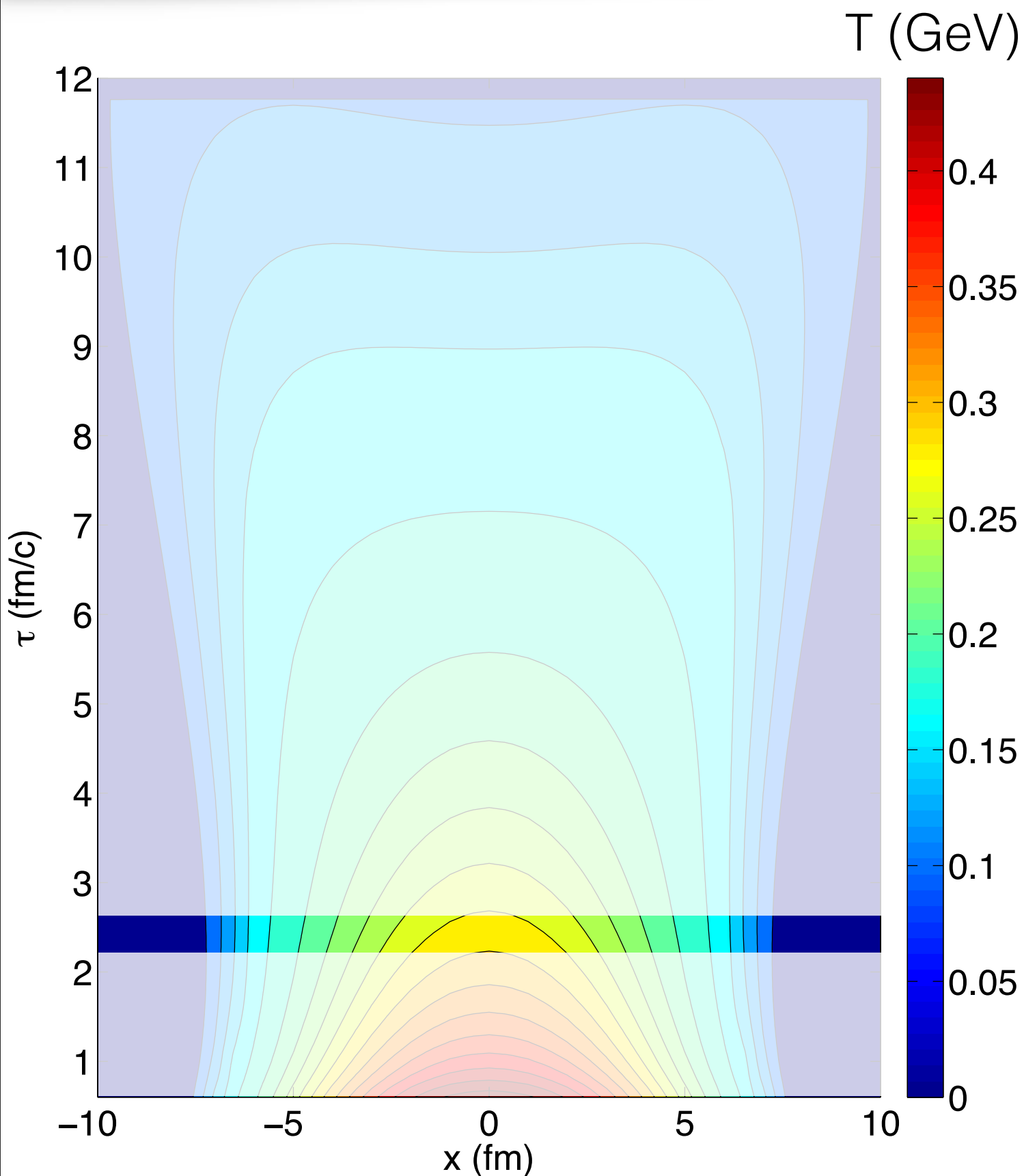
- Photon emission rates $\propto \exp(-E/T) \log(E/T)$, $T_{\text{eff}} > T$
- **All** photons with $T < 250 \text{ MeV}$ at RHIC and $< 300 \text{ MeV}$ at LHC carries T_{eff} within the experimental fitted region
- About **50-60%** of photons are emitted from $T = 165 \sim 250 \text{ MeV}$, they are strongly blue shifted by radial flow

$$T_{\text{eff}} = T \sqrt{\frac{1+v}{1-v}}$$

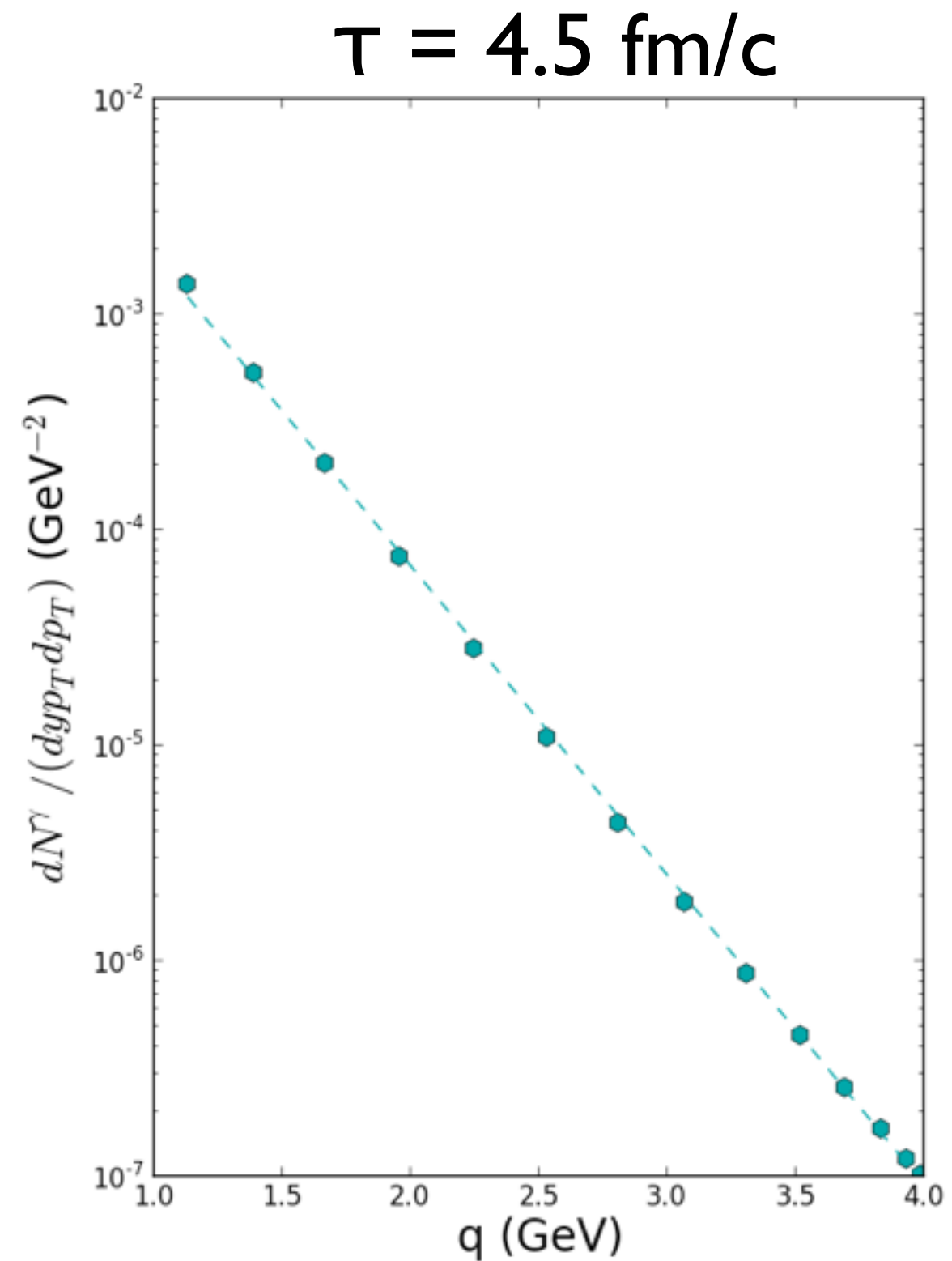
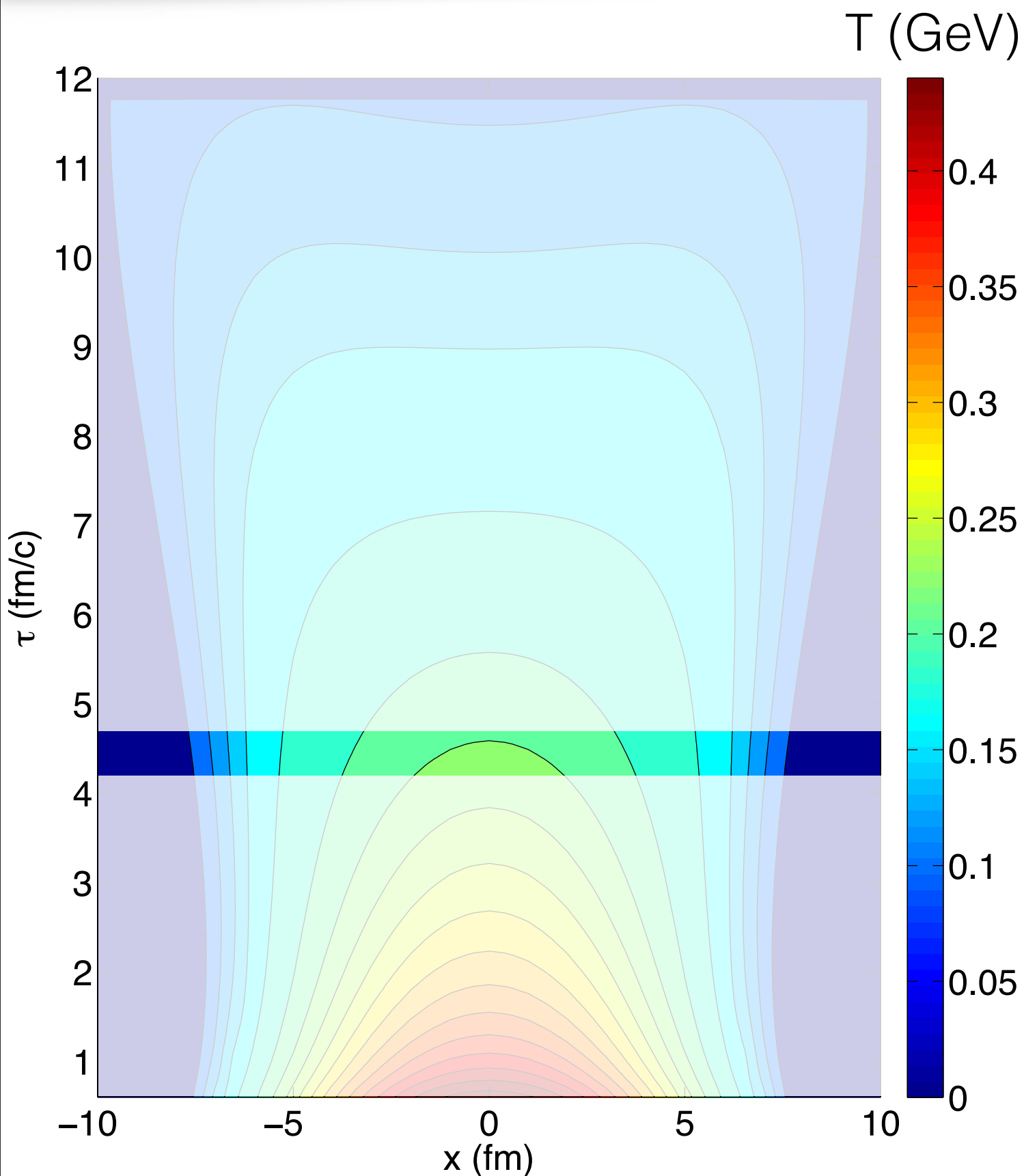
Slope of Photon Spectrum



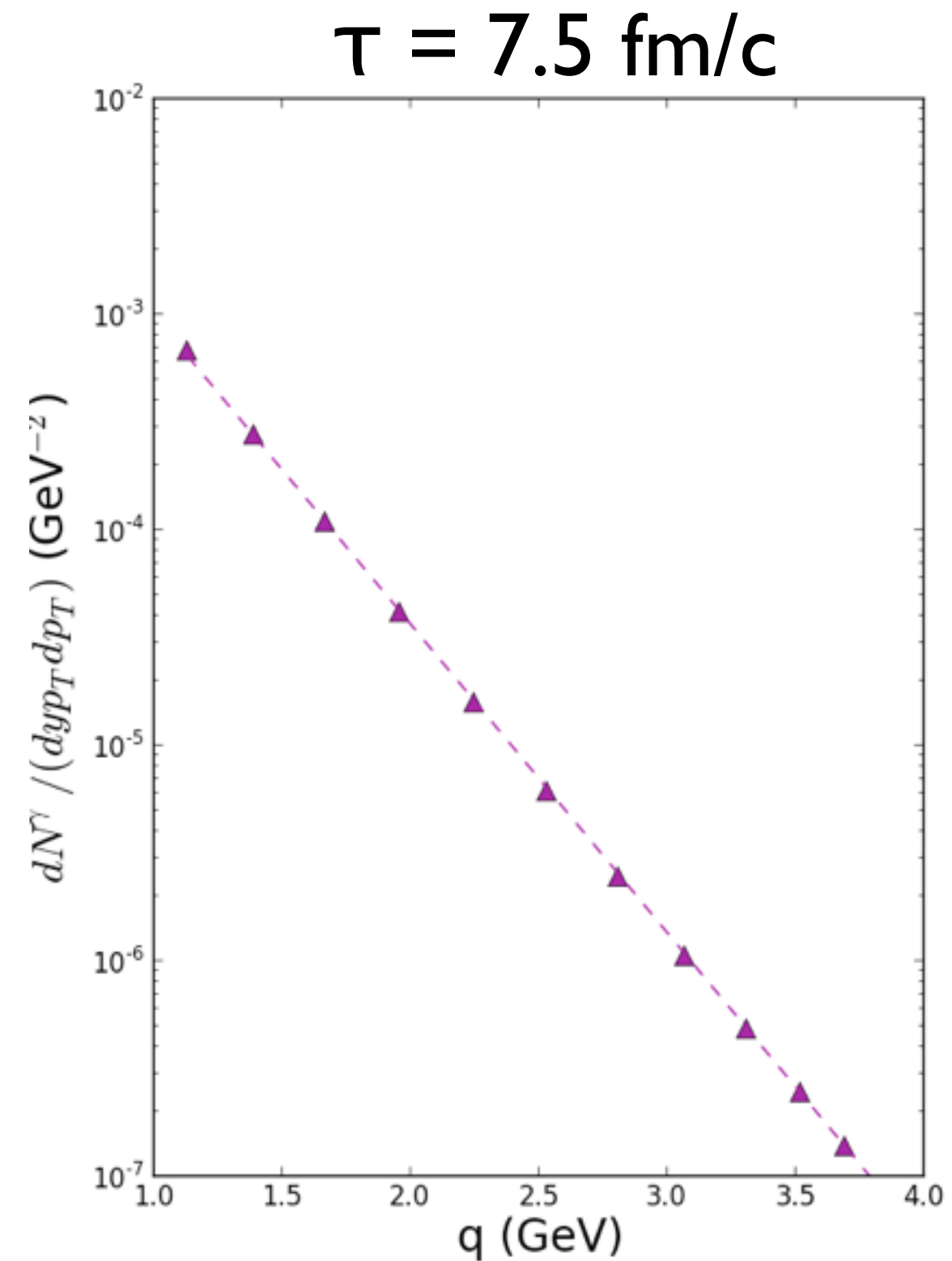
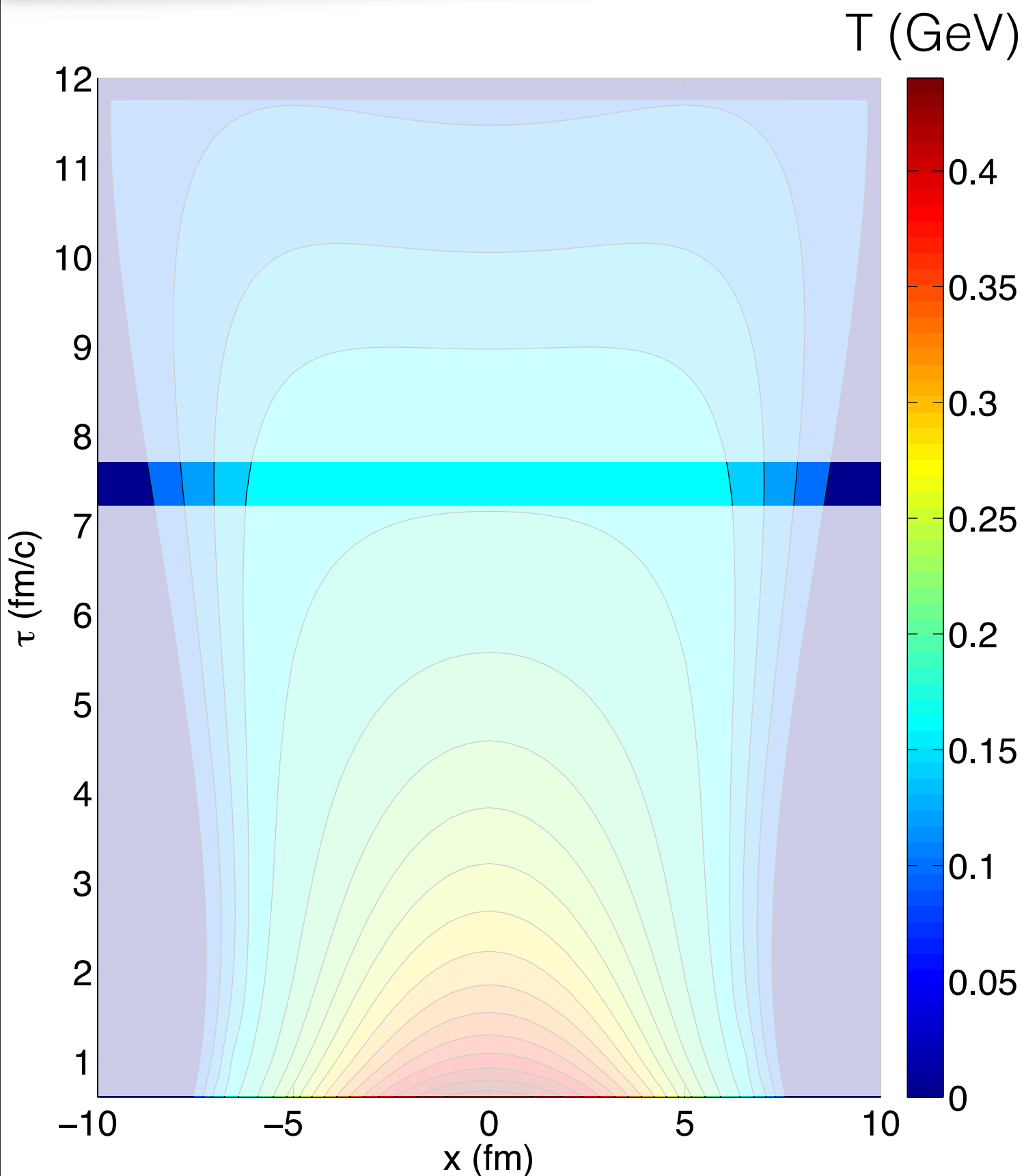
Slope of Photon Spectrum



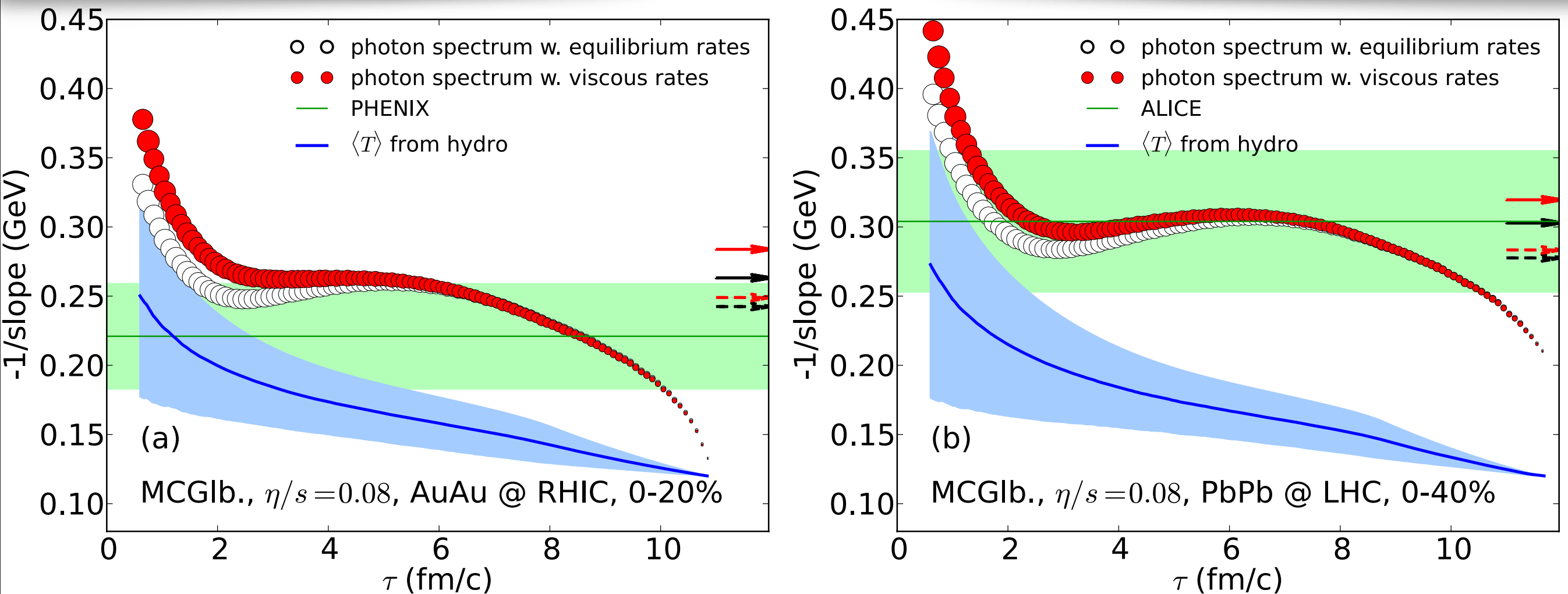
Slope of Photon Spectrum



Slope of Photon Spectrum

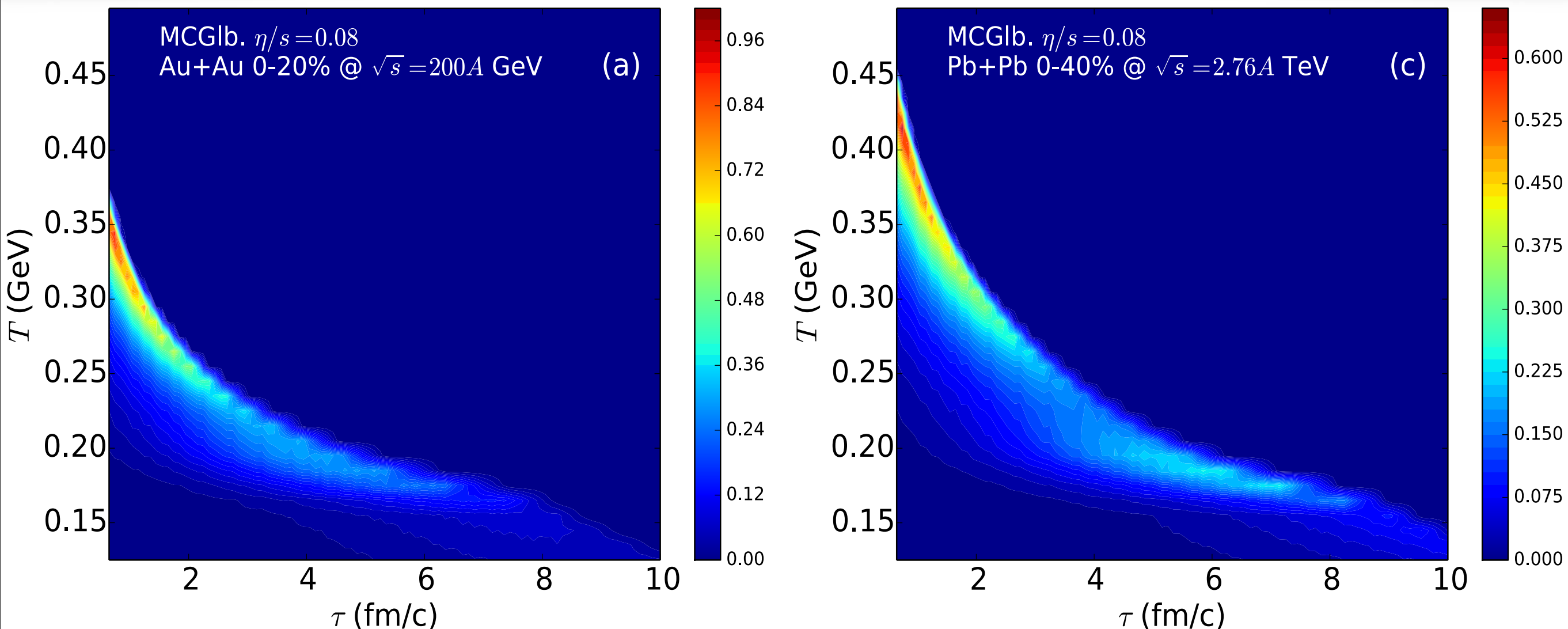


Fitted T_{eff} vs. Emission Time



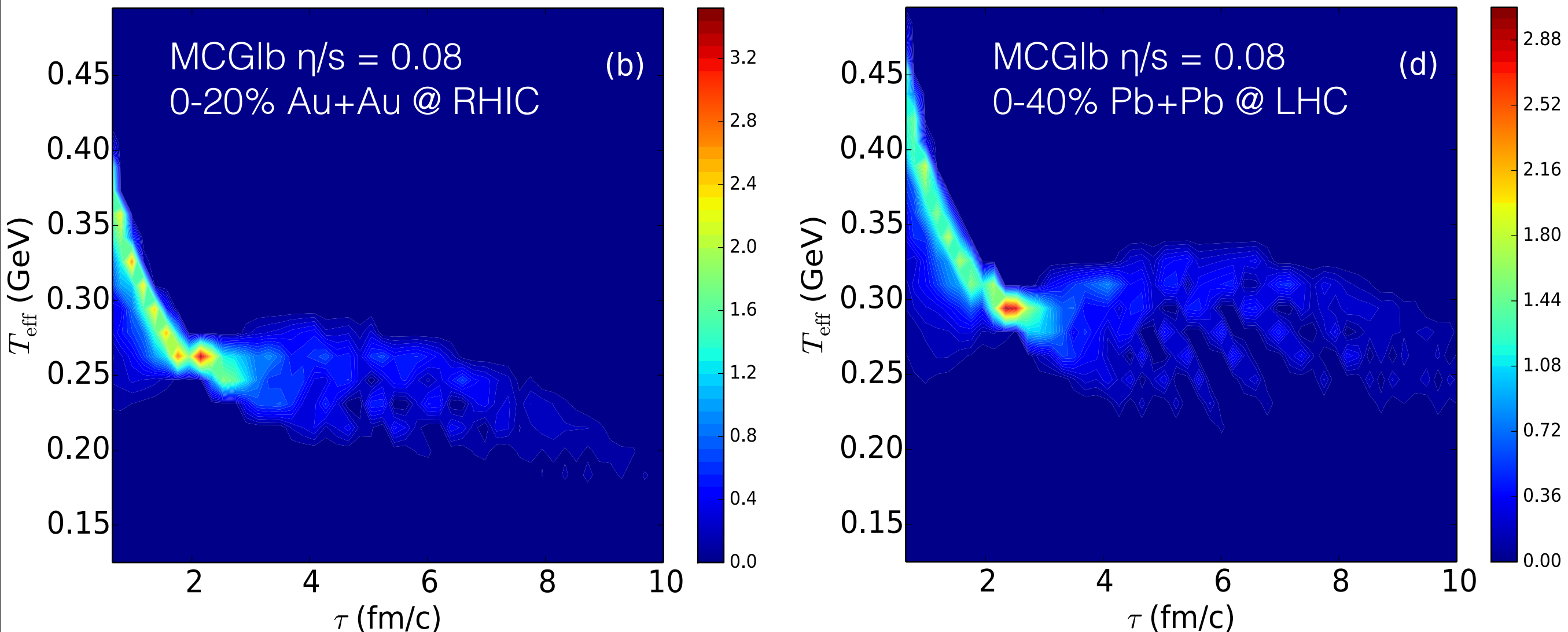
- About 25% of thermal photons are emitted in the first 2 fm/c
- After 2 fm/c, thermal photons are significantly blue shifted by radial flow
- Viscous corrections to the slope of photon spectra are stronger during the early part of the evolution

Mapping thermal photon emission



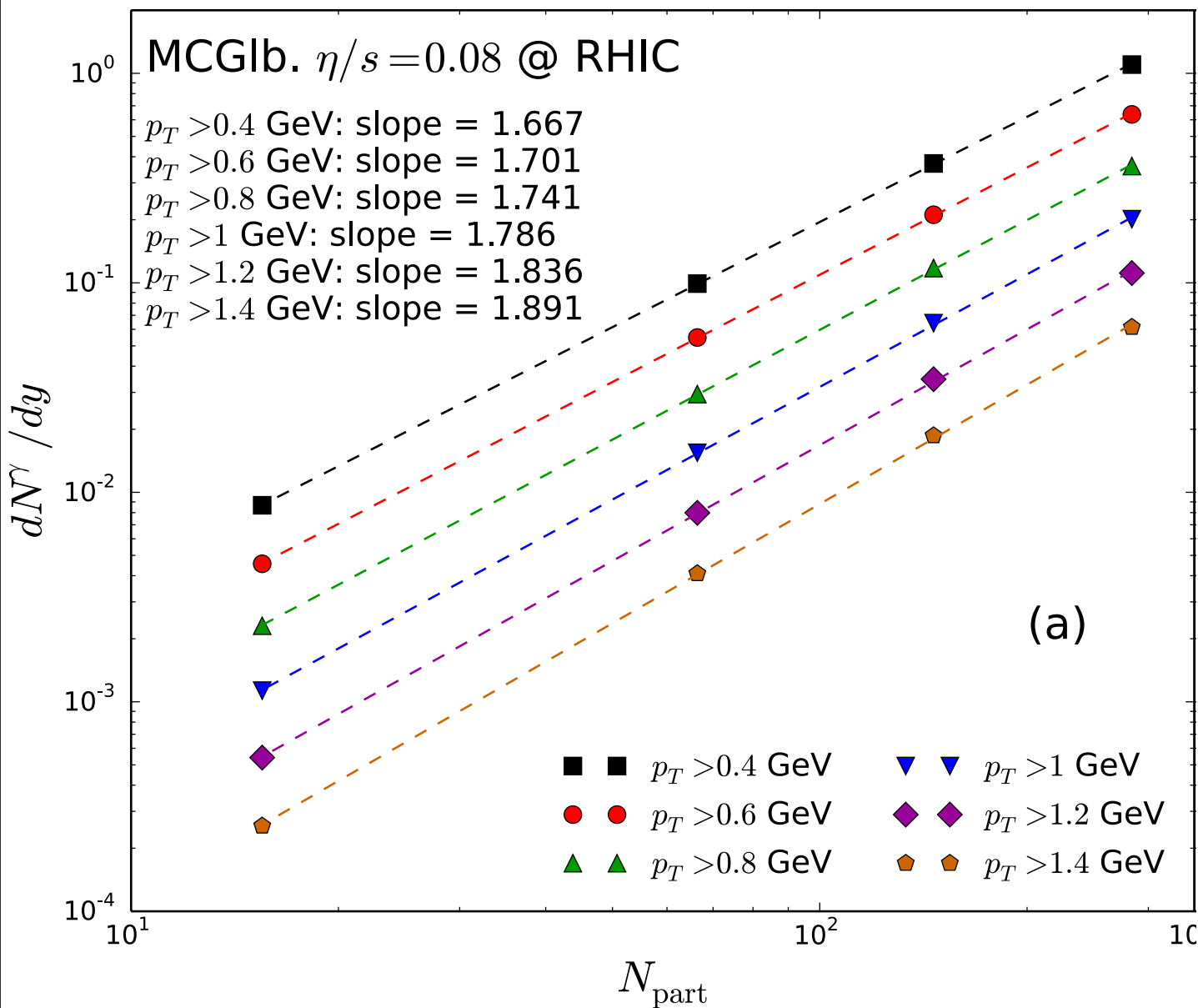
- By cutting hydro medium both in T and τ , we observe a **two-wave** thermal photon production
 - early time production — high rates at high temperatures
 - near transition region — growing of space-time volume

Mapping T_{eff}



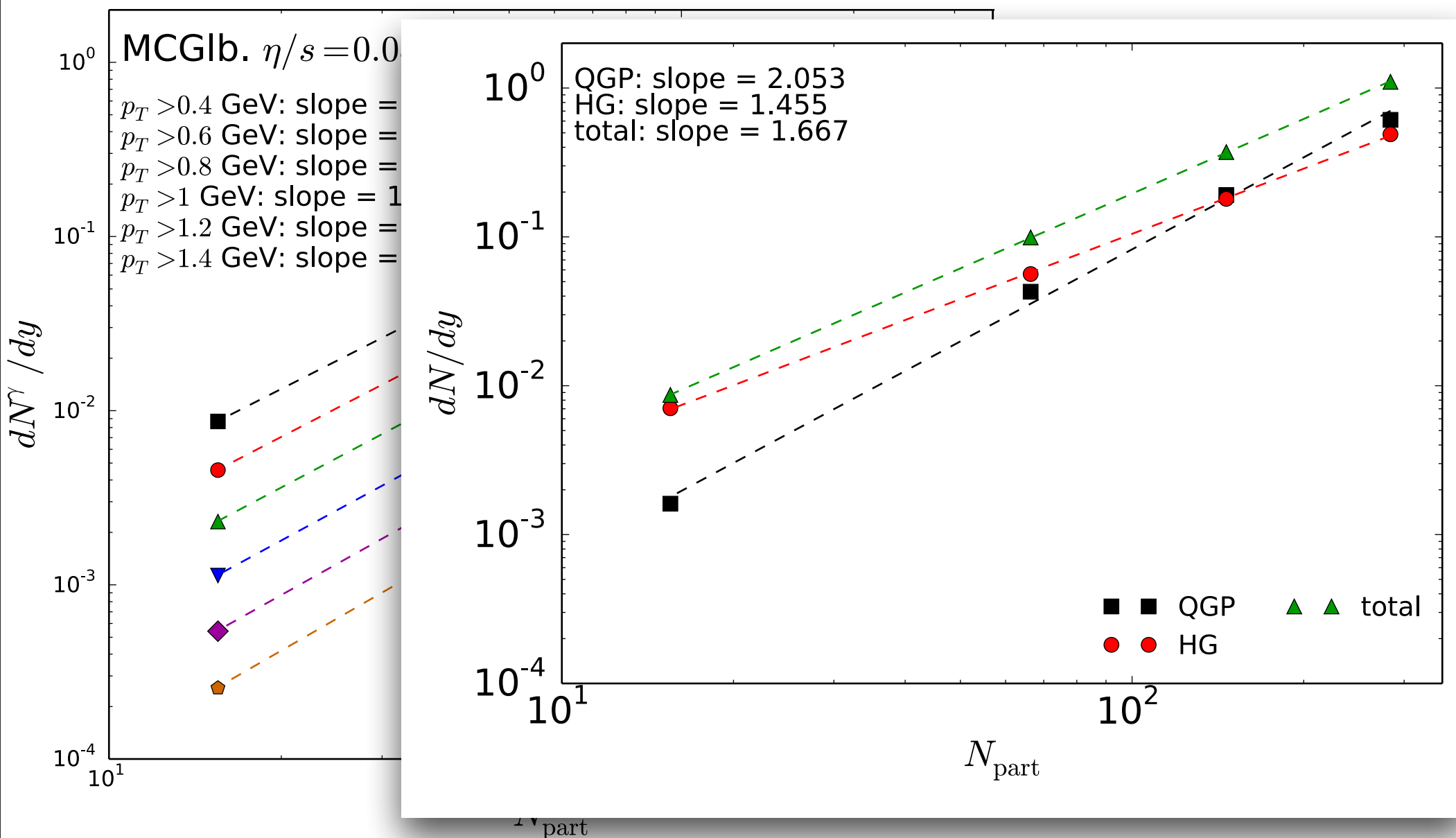
- Hydrodynamic radial flow strongly blue shifts the slopes of photon spectra
- Around 2 fm/c, it greatly shrinks the photon yield distribution in terms of the effective temperature compared to the real temperature

Centrality dependence of photon yield



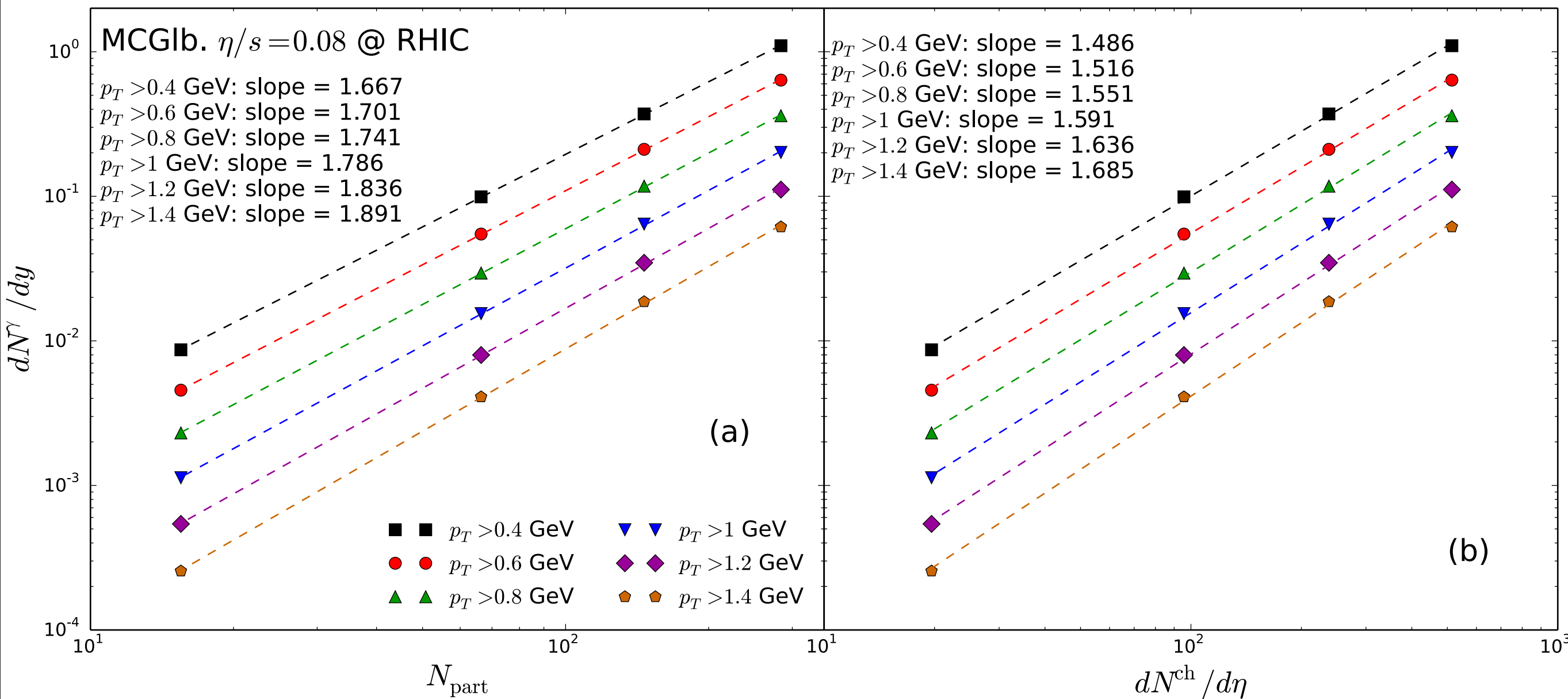
- Thermal photons from hydrodynamic medium qualitatively reproduce the centrality dependence of the direct excess photon yield at the top RHIC energy

Centrality dependence of photon yield



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Centrality dependence of photon yield



- Thermal photons from hydrodynamic medium qualitatively reproduce the centrality dependence of the direct excess photon yield at the top RHIC energy

dN^γ/dy vs. $dN^{\text{ch}}/d\eta$  **less model dependent comparison**

Conclusion

- Current theoretical calculation still **underestimates** the experimental measured direct photon spectra at both RHIC and LHC
- Thermal photon spectra are strongly **blue shifted** by hydrodynamic radial flow; their inverse slopes are **not** reflecting the true underlying temperature of the evolving fireball
- By cutting hydrodynamic space-time volume, we map out thermal photon emission as a function of temperature and emission time
- Thermal photon yields show a **stronger** centrality dependence than charged hadrons. The extracted exponent suggests a larger emission from late stage of the hydrodynamic evolution in the hadronic phase