

Identifying critical behavior in heavy-ion collision observables: A mean-field hadronic transport approach

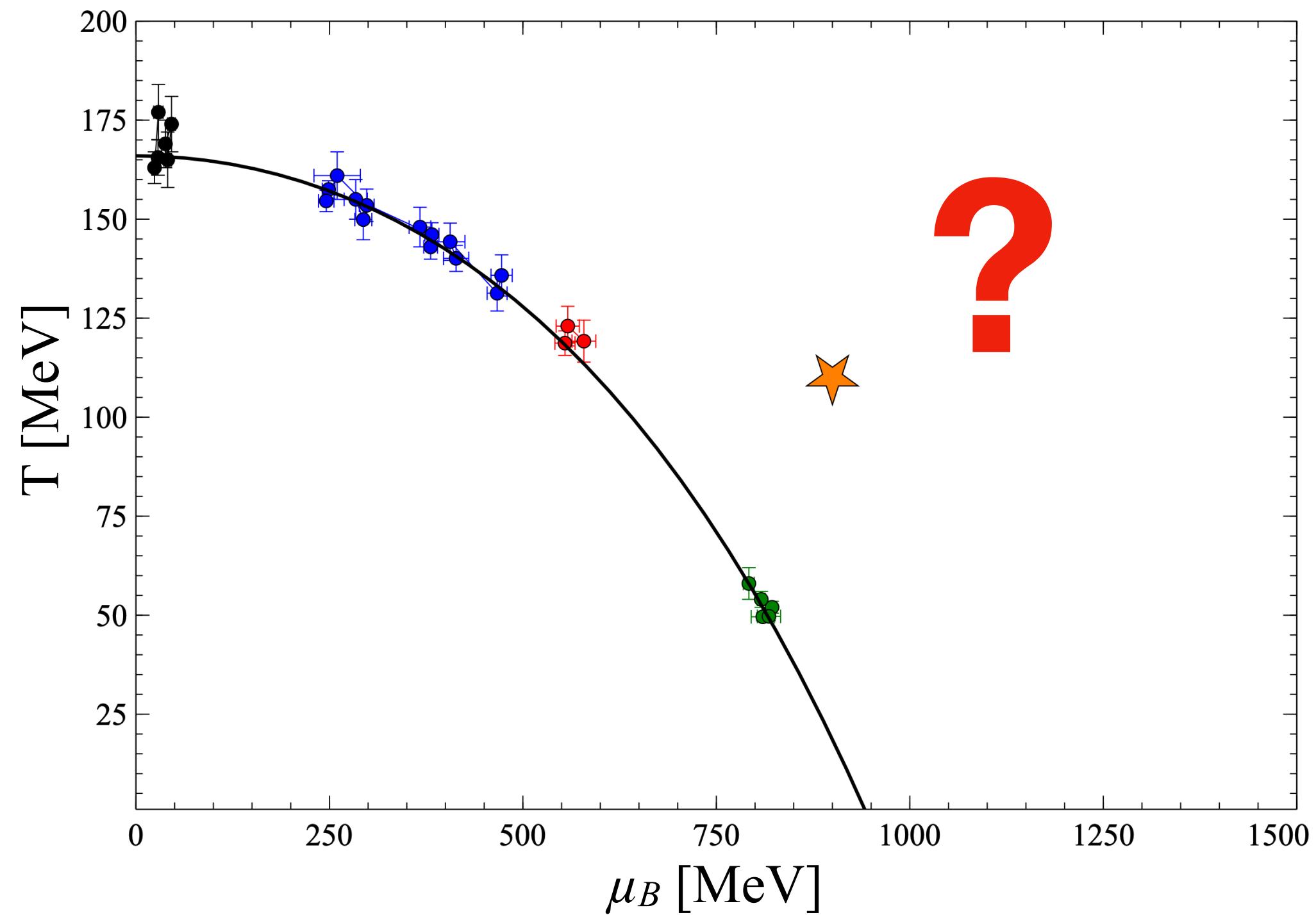
Agnieszka Wergieluk (UCLA/LBNL)

in collaboration with Volker Koch (LBNL)



May 16th, 2020
2020 Fall Meeting of the APS Division of Nuclear Physics

Motivation



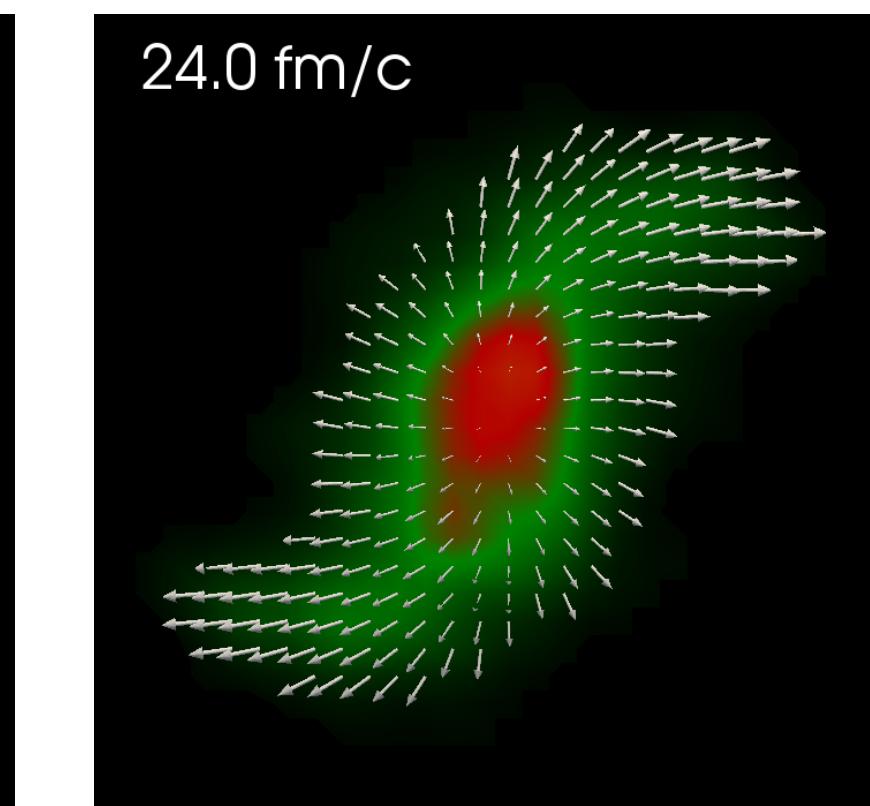
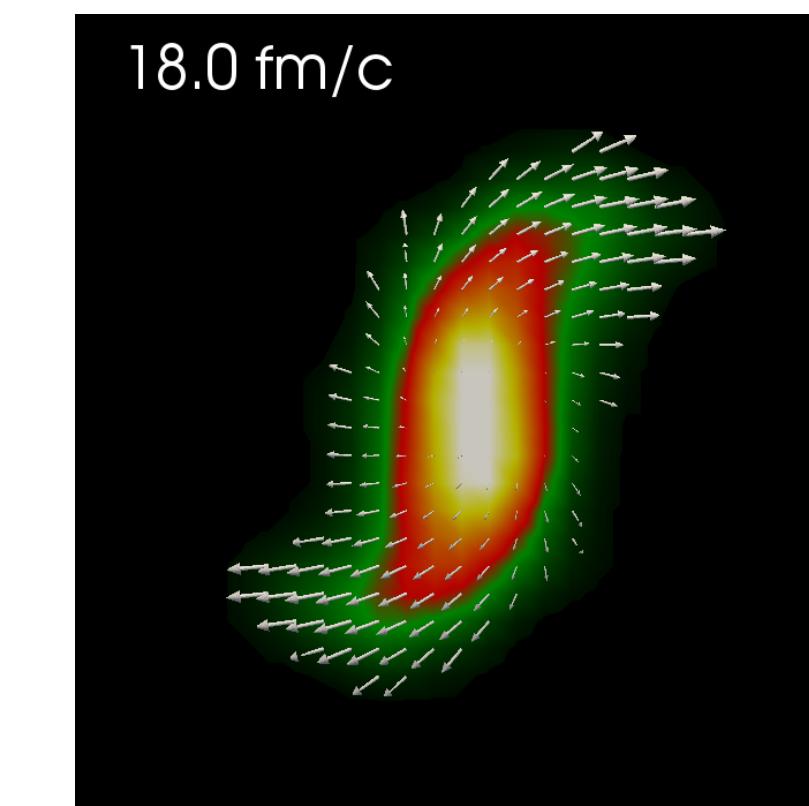
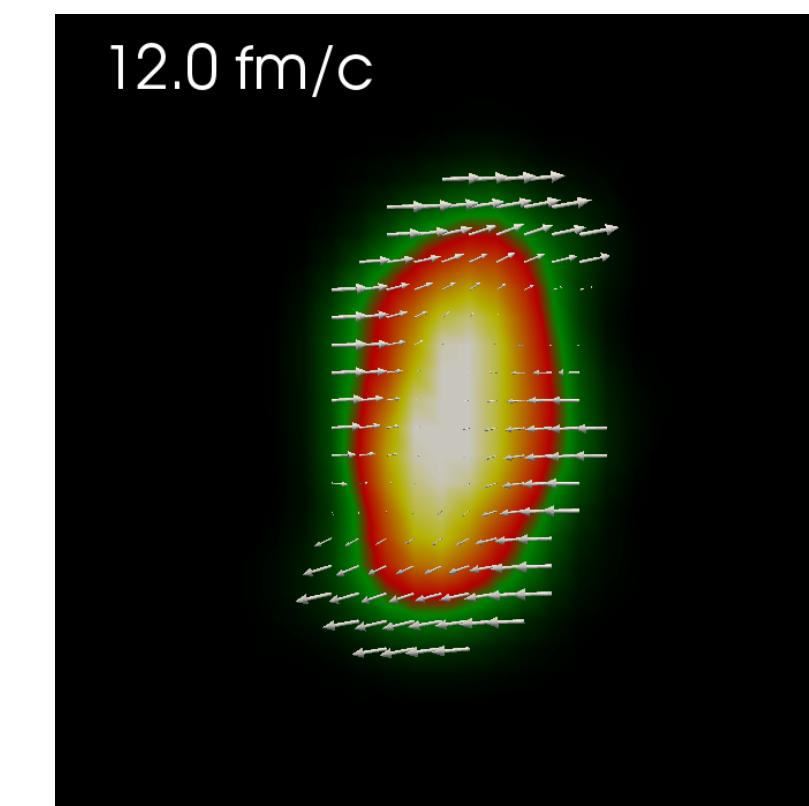
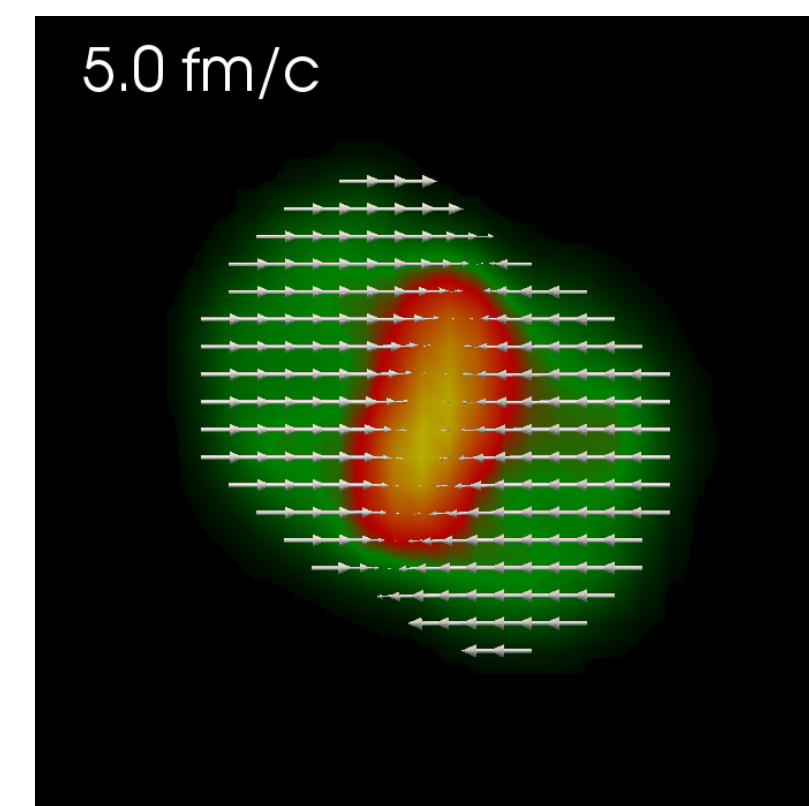
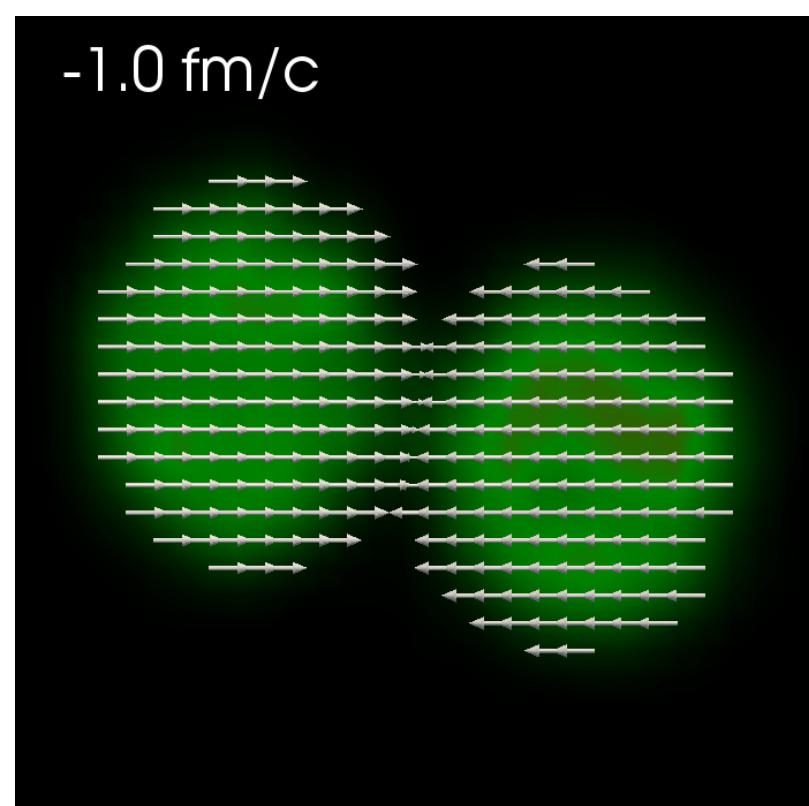
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How precisely can we measure it?

We try to answer by comparing simulations with data

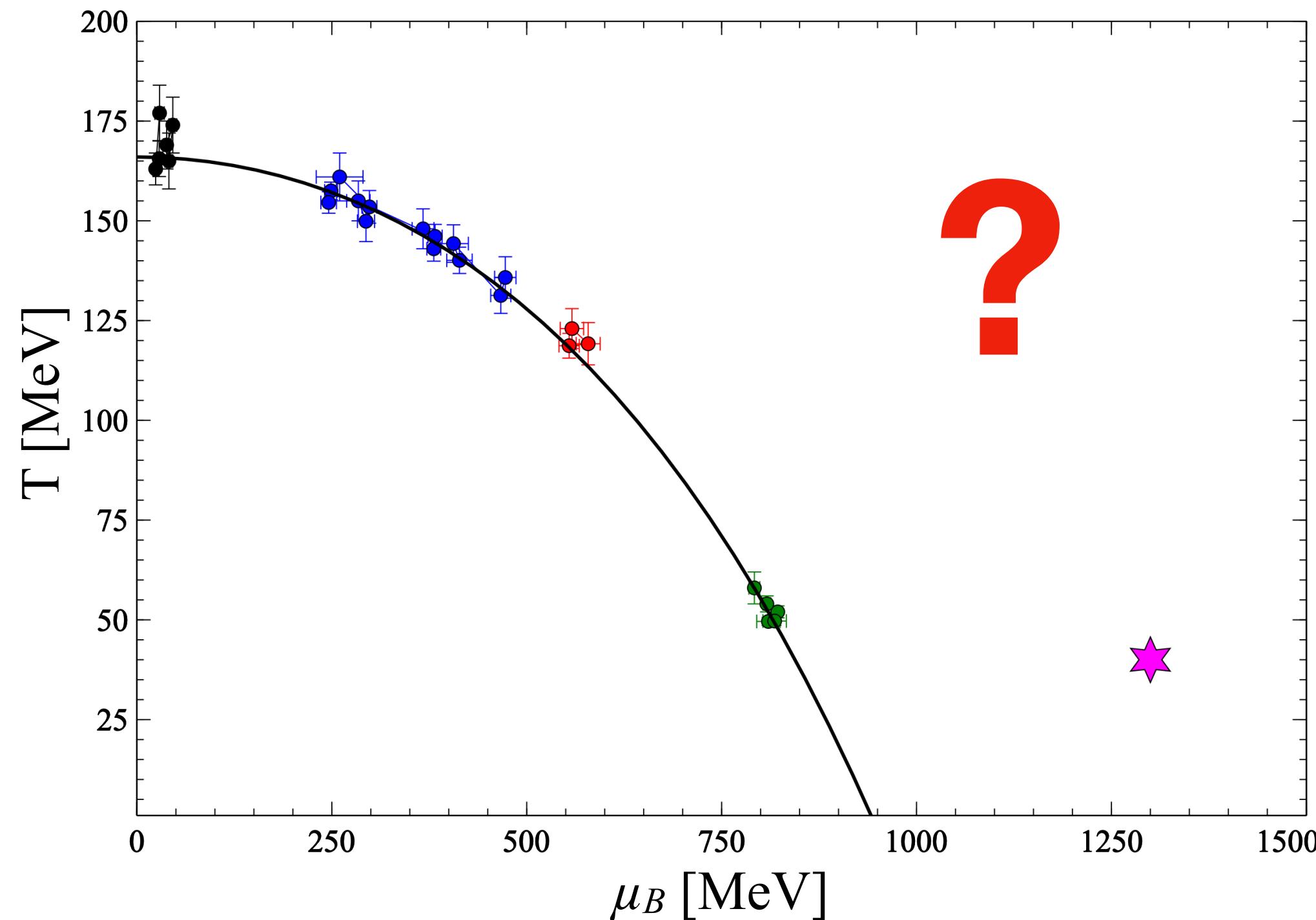
We need: thermodynamics (hydrodynamic evolution)
and particle dynamics (hadronic afterburner)

Hadronic phase is ever more important for low energies:
how is it affected by the QCD equation of state (EOS)?

H. Petersen, D. Oliinychenko, M. Mayer, J. Staudenmaier, S. Ryu,
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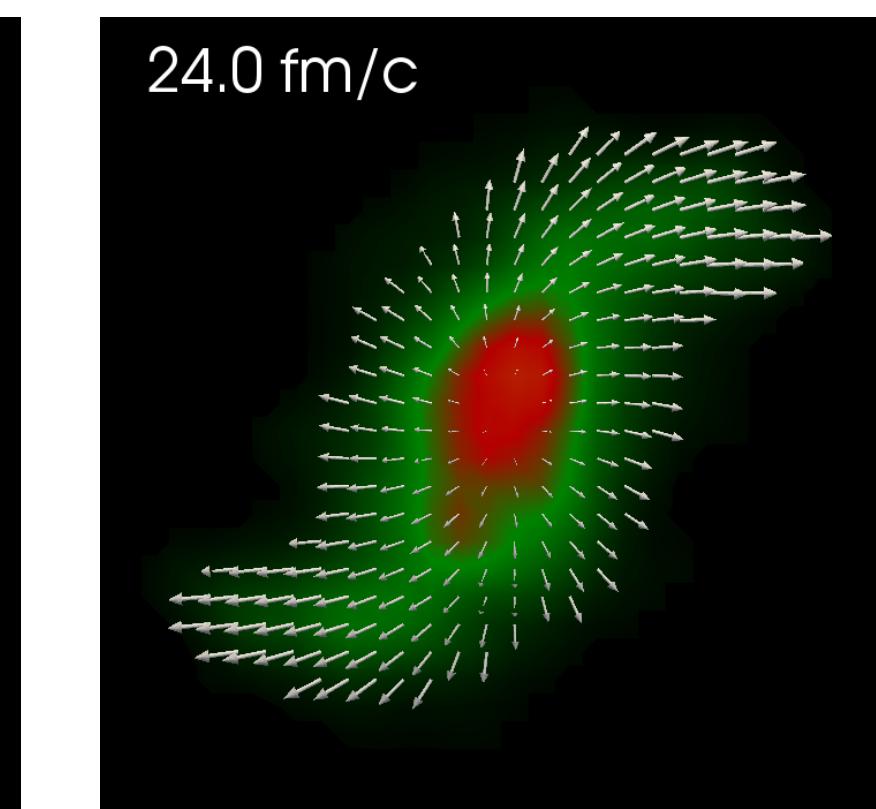
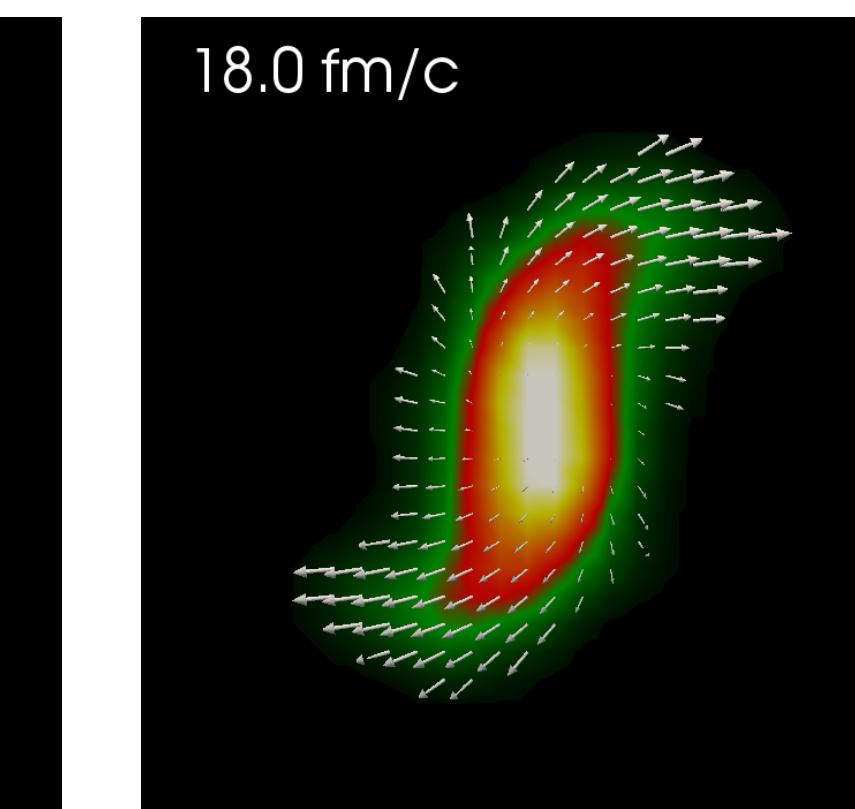
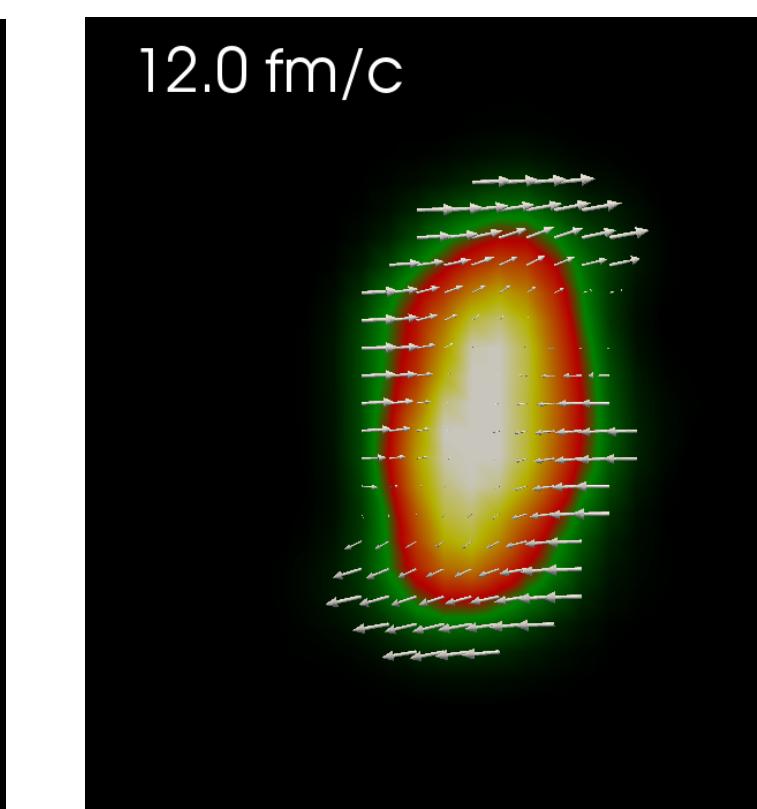
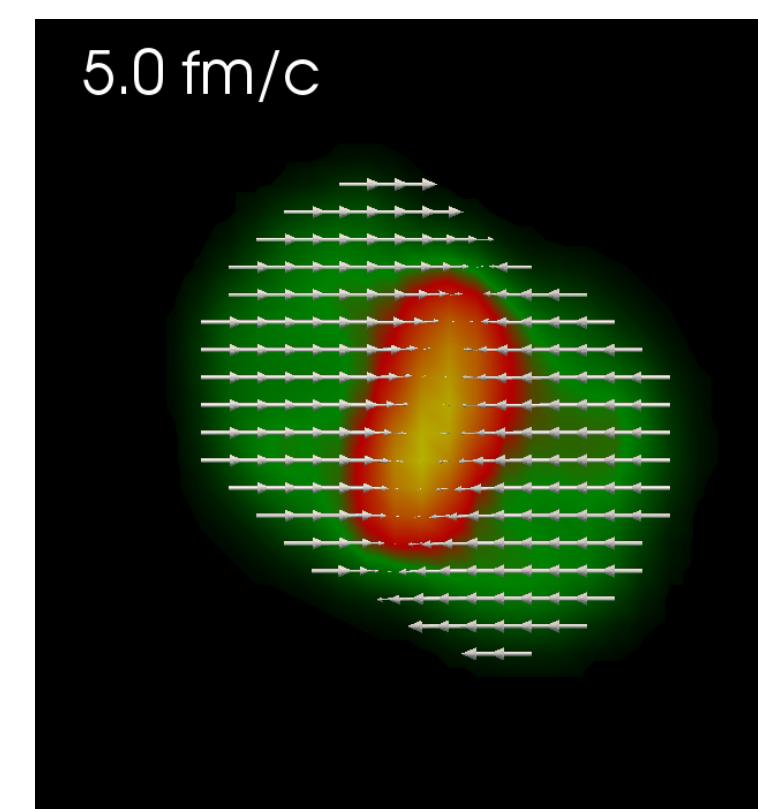
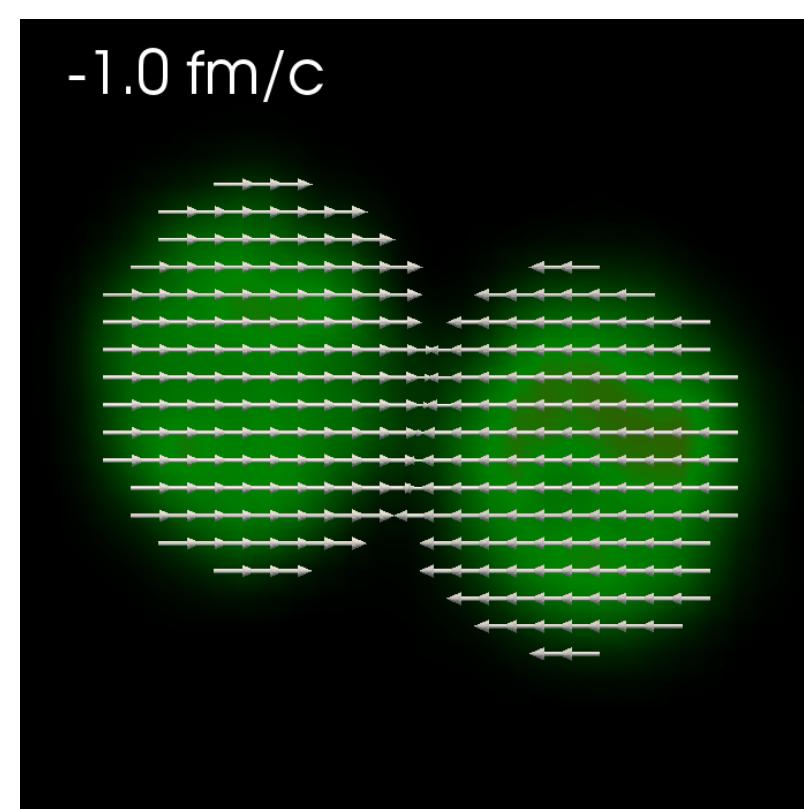
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- Flexible and systematic **parameterization** of the EOS of dense QCD matter
- Access to both thermodynamics and **single-particle equations of motion** for hadronic transport
- Relatively simple, **baryon-number-density—dependent interactions**: easy to use in hadronic transport (nucleons as the degrees of freedom, mean field level interactions only)

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We do **not** want to build a model that predicts the critical point.

We construct a consistent, flexible **parameterization** of critical behavior in hot and dense nuclear matter to use in simulations and compare the results with data.

Relativistic DF with vector-number-density–based interactions

G. Baym and S. A. Chin, “Landau Theory of Relativistic Fermi Liquids,” Nucl. Phys. A **262**, 527 (1976)

1) Postulate the energy density of the system:

$$\mathcal{E} = \mathcal{E}[f_{\mathbf{p}}] = \int \frac{d^3 p}{(2\pi)^3} \epsilon_{kin} f_{\mathbf{p}} + \sum_{i=1}^N C_i (j_\mu j^\mu)^{\frac{b_i}{2}-1} \left[j^0 j^0 - g^{00} \left(\frac{b_i - 1}{b_i} \right) j_\lambda j^\lambda \right]$$

← like Skyrme energy,
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$$\varepsilon_{\mathbf{p}} \equiv \frac{\delta \mathcal{E}[f_{\mathbf{p}}]}{\delta f_{\mathbf{p}}} \quad \frac{dx^i}{dt} \equiv - \frac{\partial \varepsilon_{\mathbf{p}}}{\partial p_i}, \quad \frac{dp^i}{dt} \equiv \frac{\partial \varepsilon_{\mathbf{p}}}{\partial x_i}$$

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3) Use the Boltzmann equation + EOMs to get conservation laws, the energy-stress tensor:

$$\frac{\partial f_{\mathbf{p}}}{\partial t} + (\nabla_{\mathbf{p}} \varepsilon_{\mathbf{p}}) \cdot (\nabla_{\mathbf{r}} f_{\mathbf{p}}) - (\nabla_{\mathbf{r}} \varepsilon_{\mathbf{p}}) \cdot (\nabla_{\mathbf{p}} f_{\mathbf{p}}) = 0 \quad \longrightarrow \quad \partial_\nu T^{\mu\nu} = 0$$

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4) Obtain pressure from the energy-stress tensor: $P = \frac{1}{3} \sum_i T^{ii} \Big|_{\text{rest frame}}$

← used for
parameterizing
the EOS

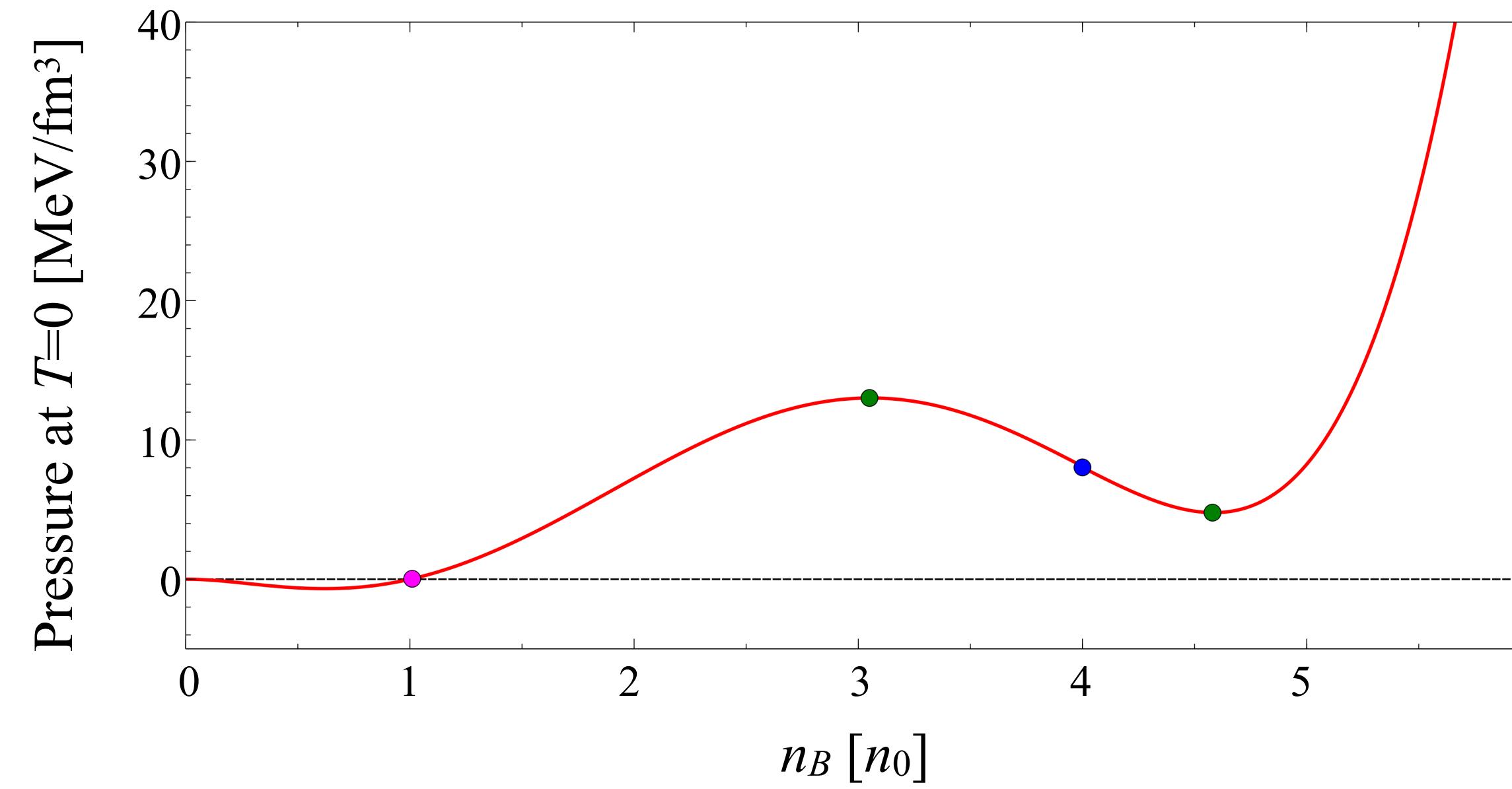
Relativistic DF with 2 phase transitions

Systems with two 1st order phase transitions: nuclear and “quark/hadron”

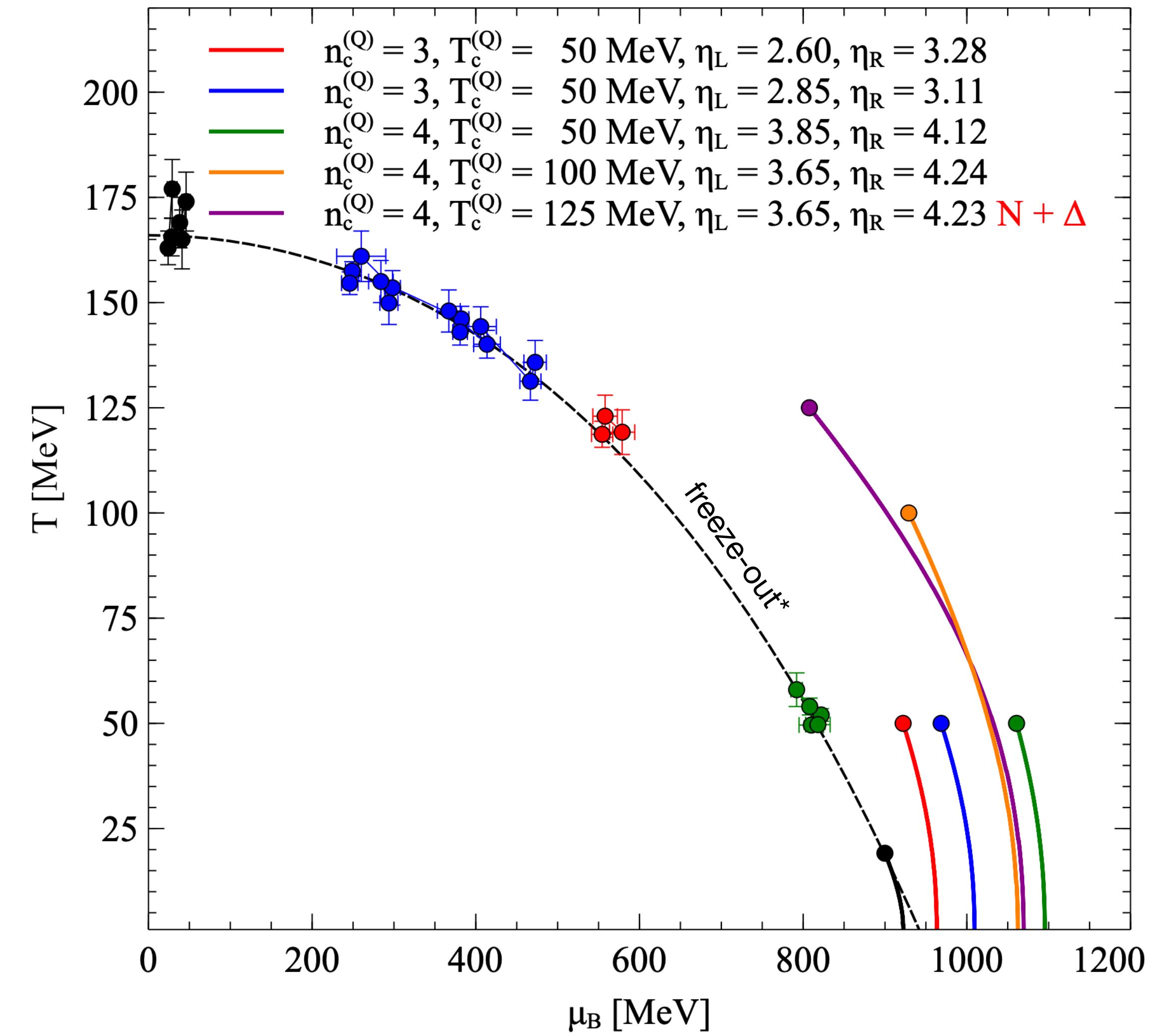
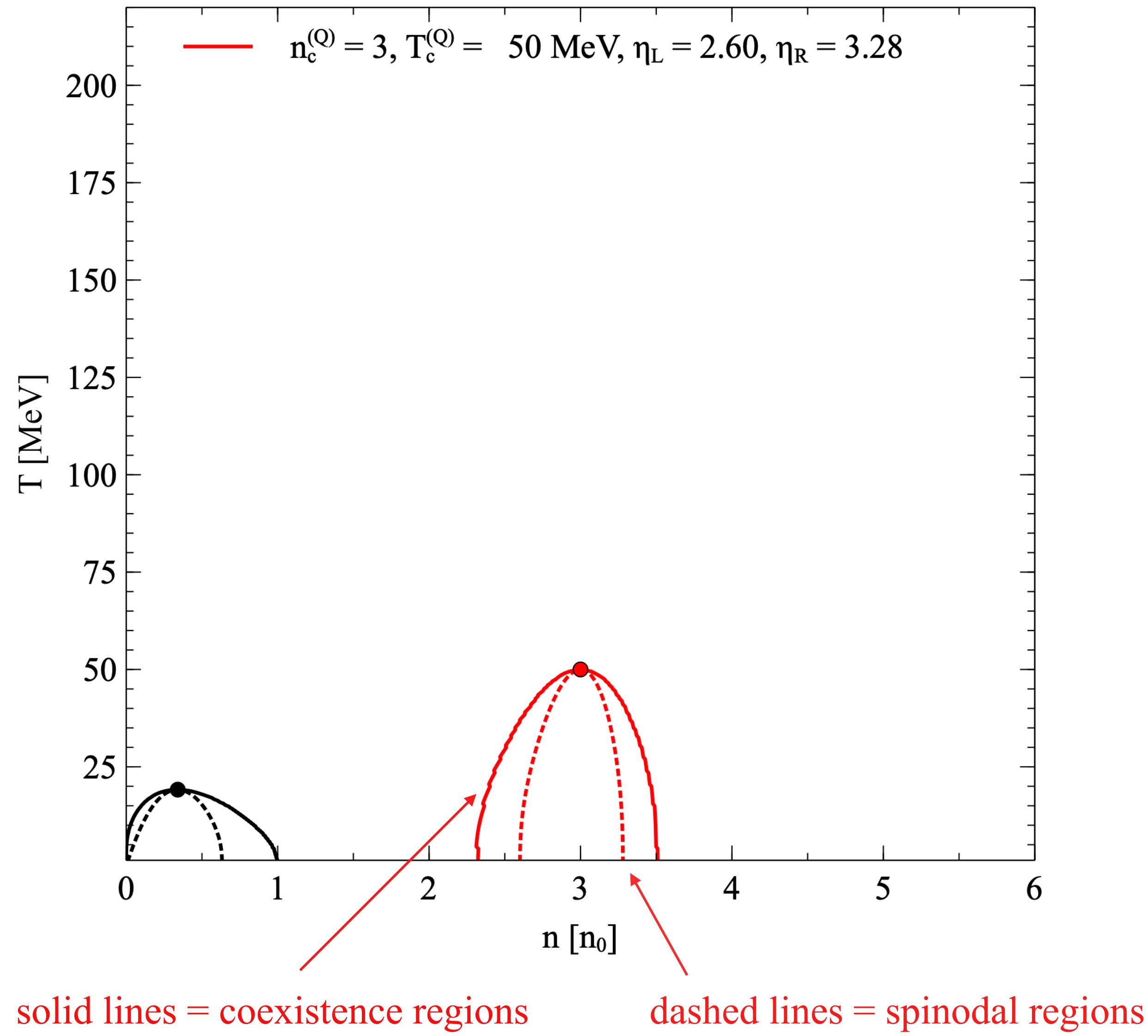
- Degrees of freedom: nucleons
- “Quark” matter coexists with dense nuclear matter, not vacuum
- 4 interactions terms = 8 parameters to fix: $\{C_1, C_2, C_3, C_4, b_1, b_2, b_3, b_4\}$

Pressure:

$$P = g \int \frac{d^3 p}{(2\pi)^3} T \ln \left[1 + e^{-\beta(\varepsilon_p - \mu_B)} \right] + \sum_{i=1}^{N=4} C_i \frac{b_i - 1}{b_i} n_B^{b_i}$$

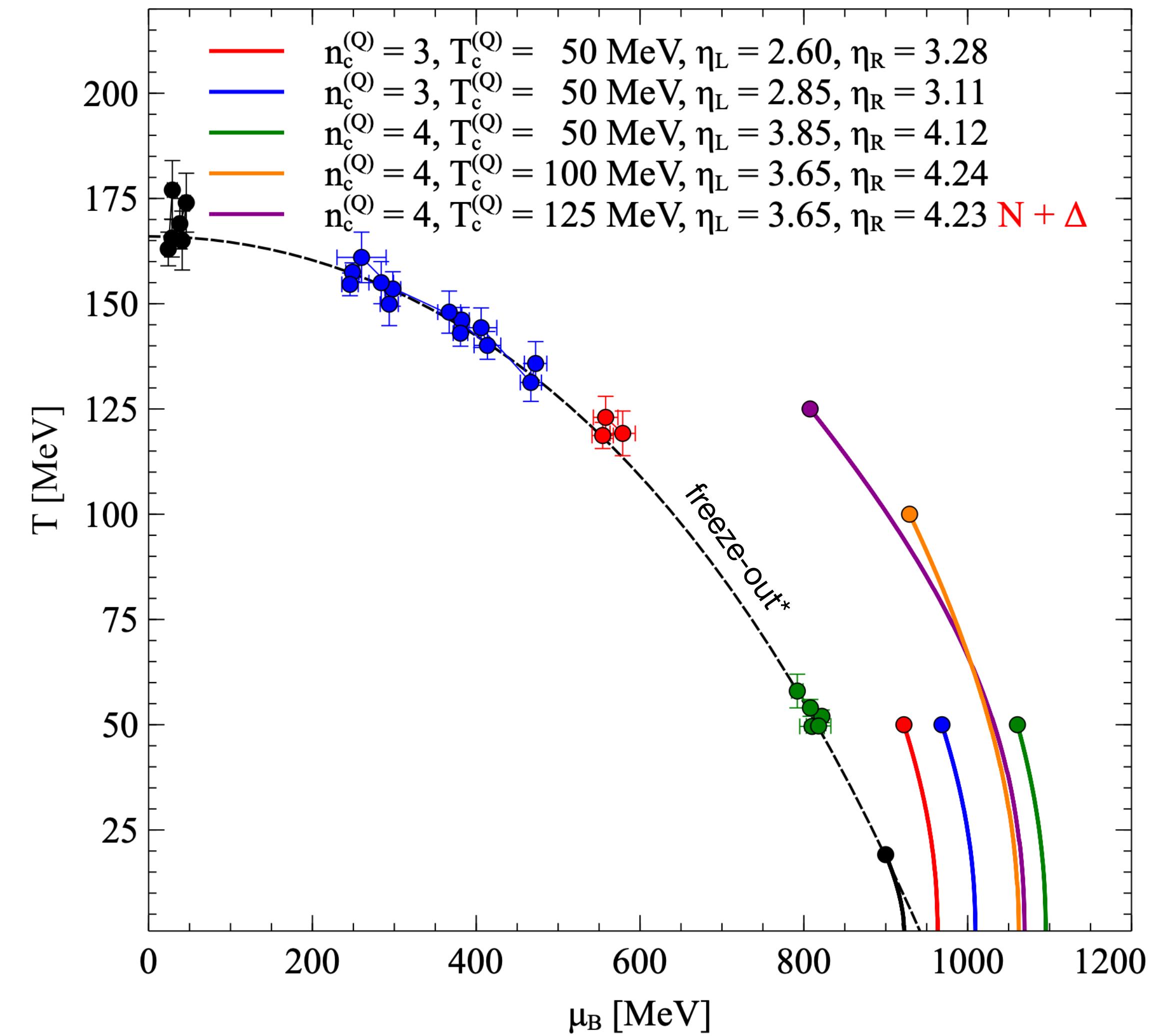
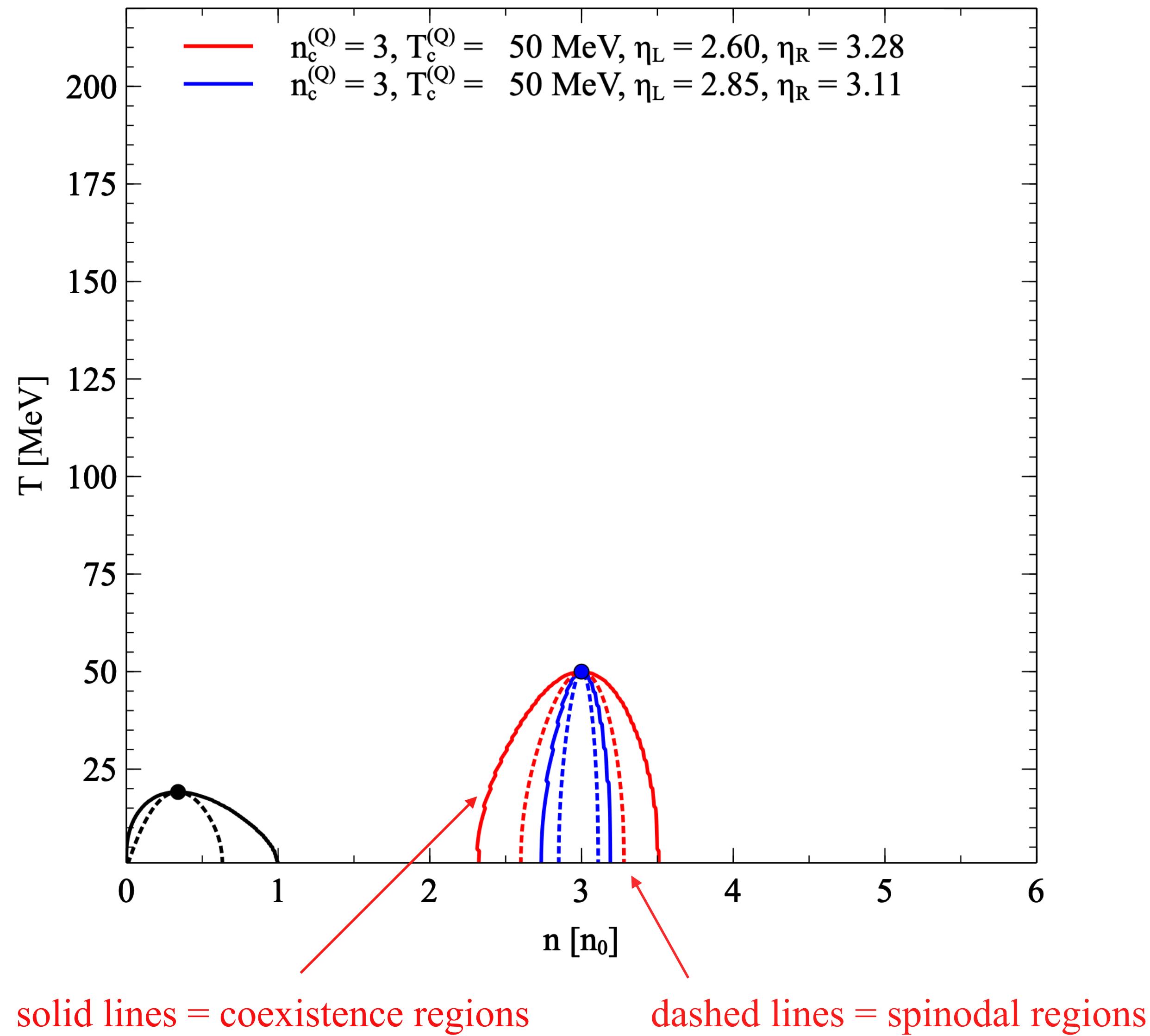


Phase diagram in the (T, n_B) and (T, μ_B) plane



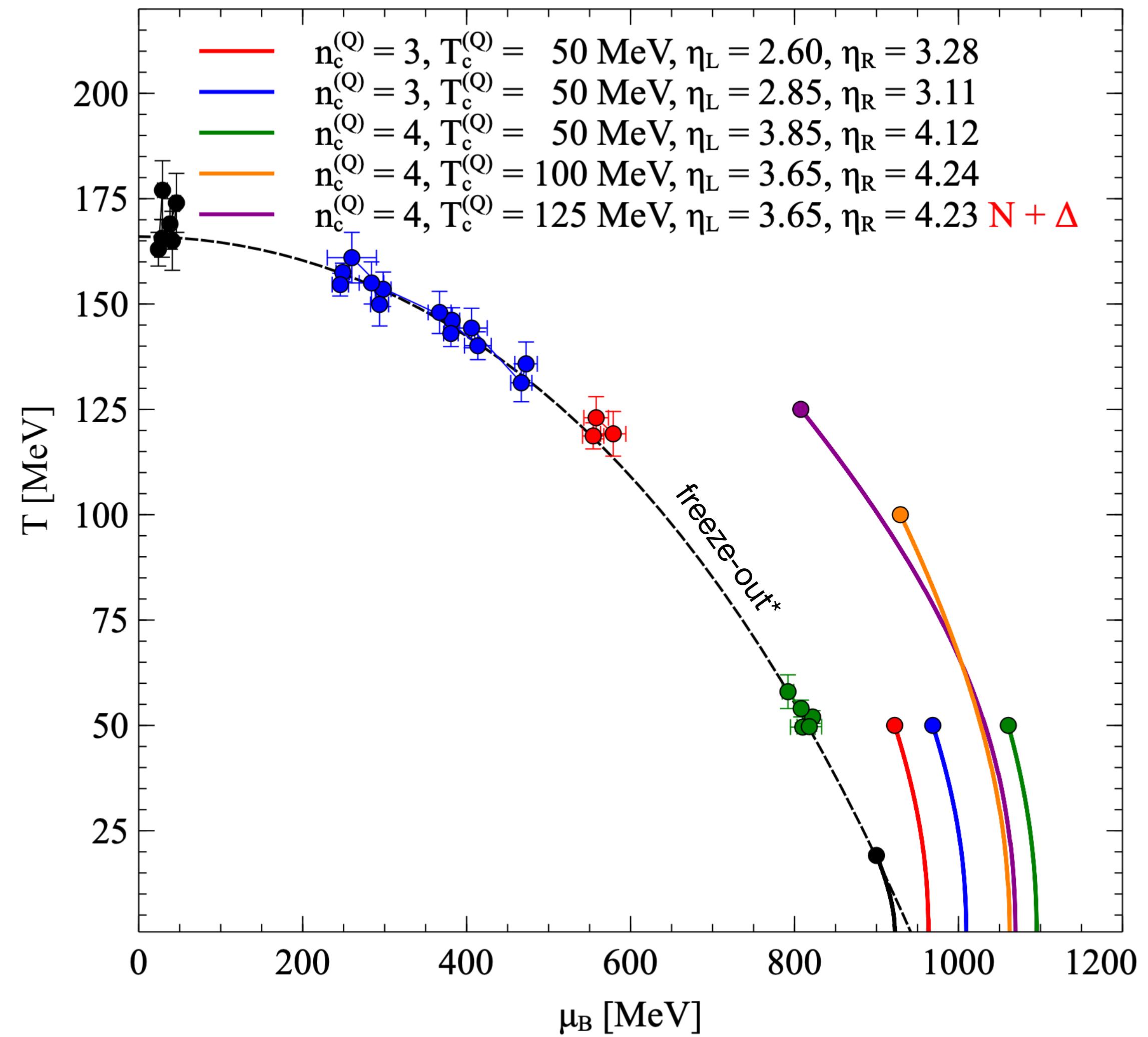
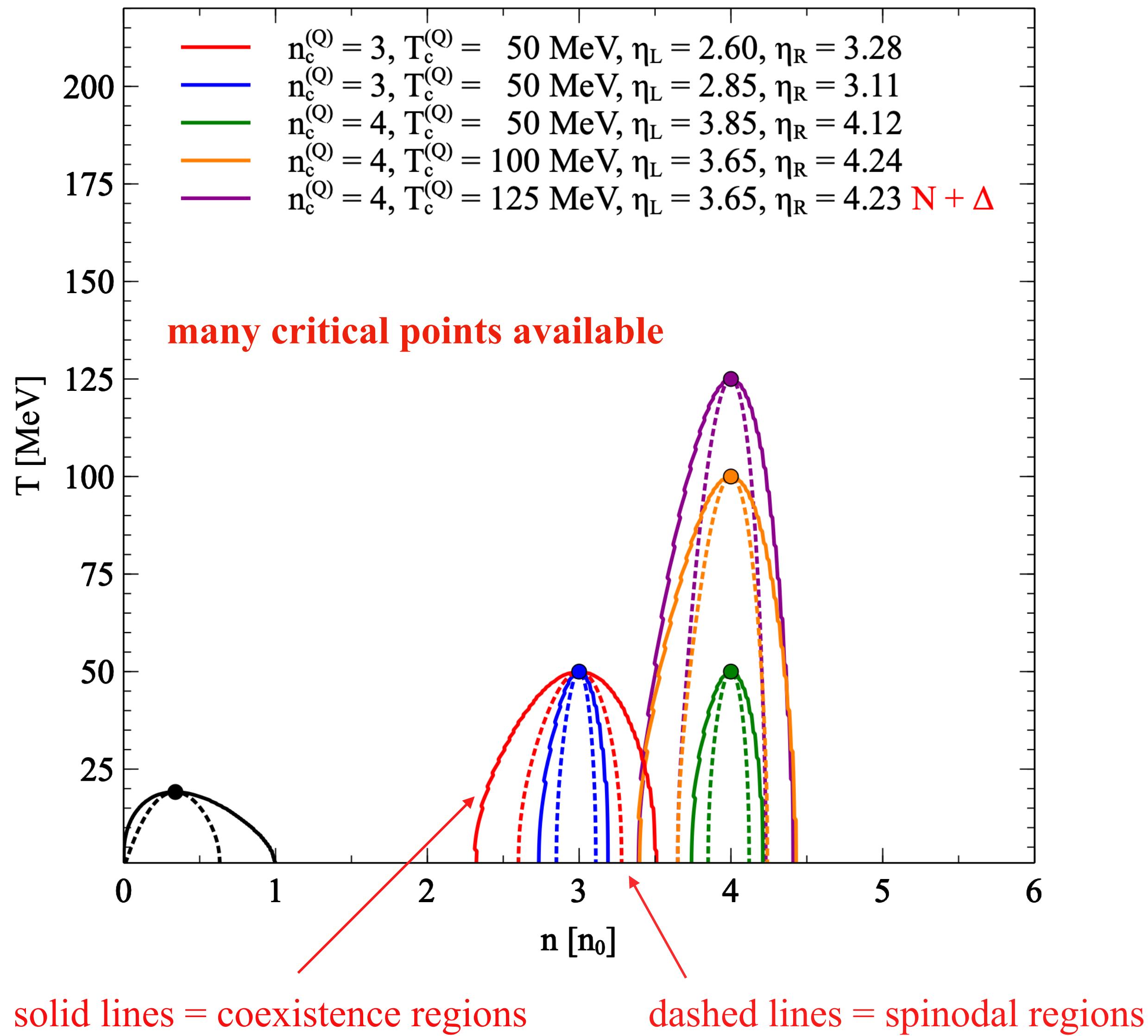
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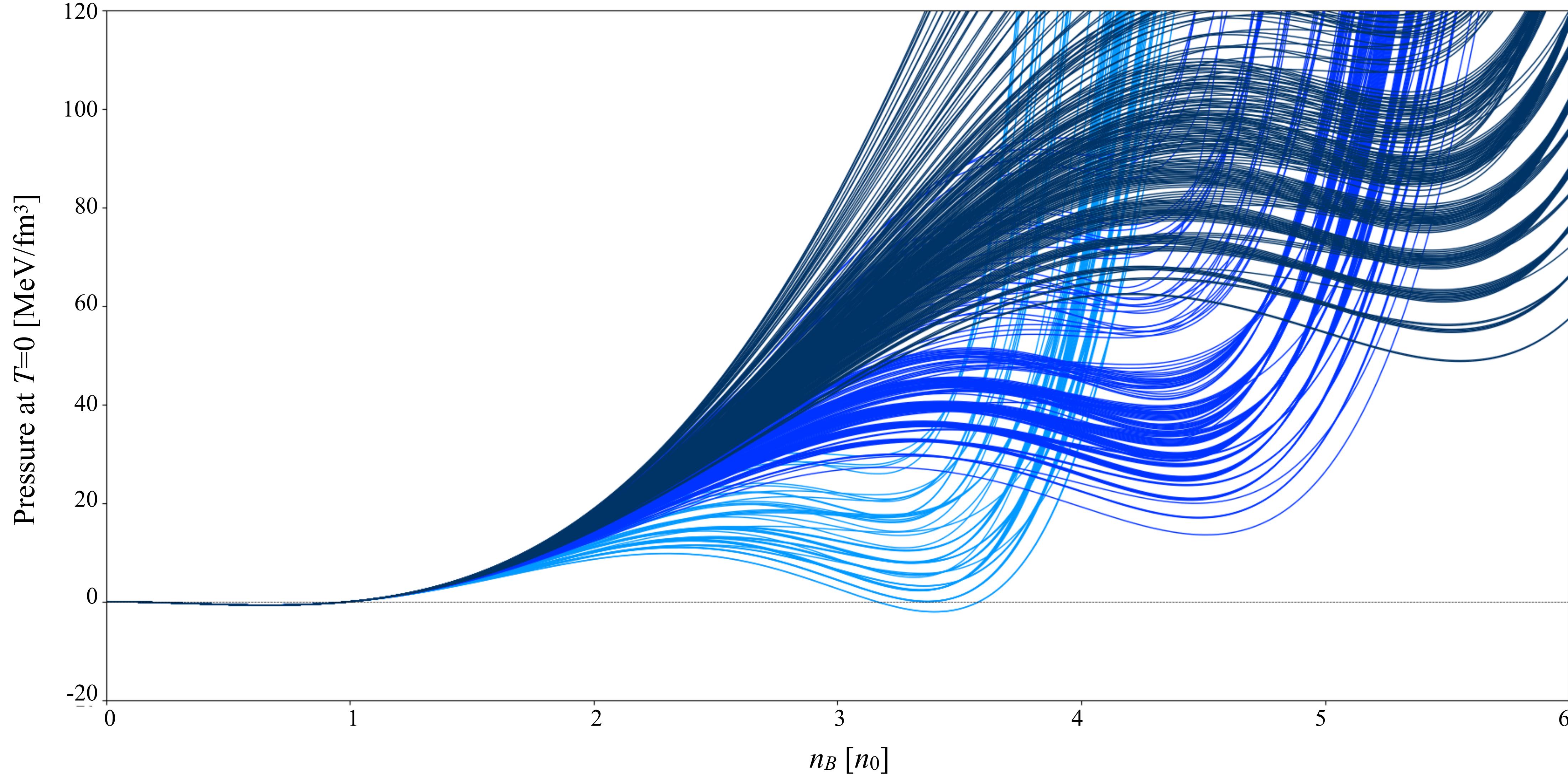
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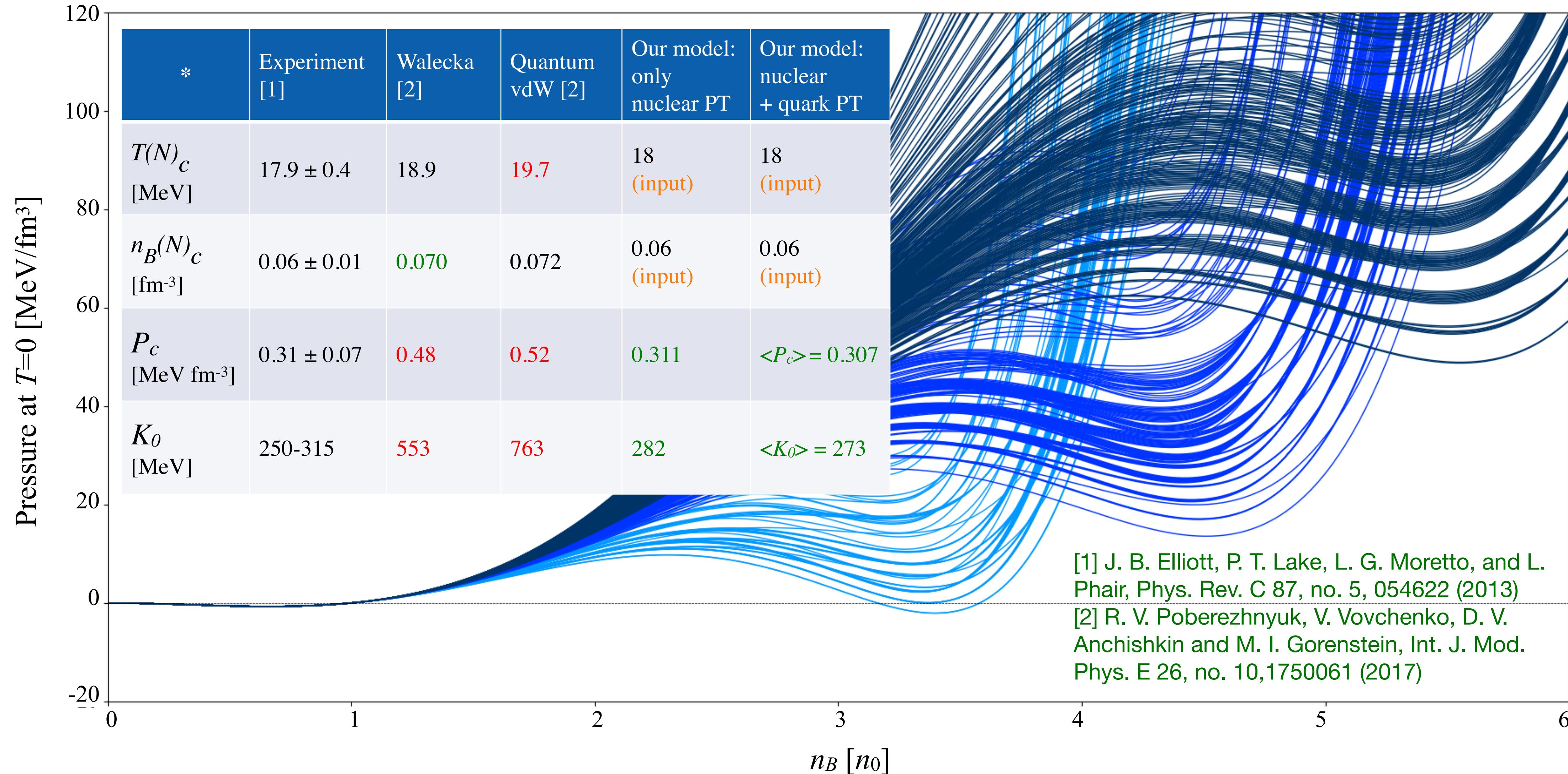
Relativistic DF with 2 phase transitions

Many equations of state to try!

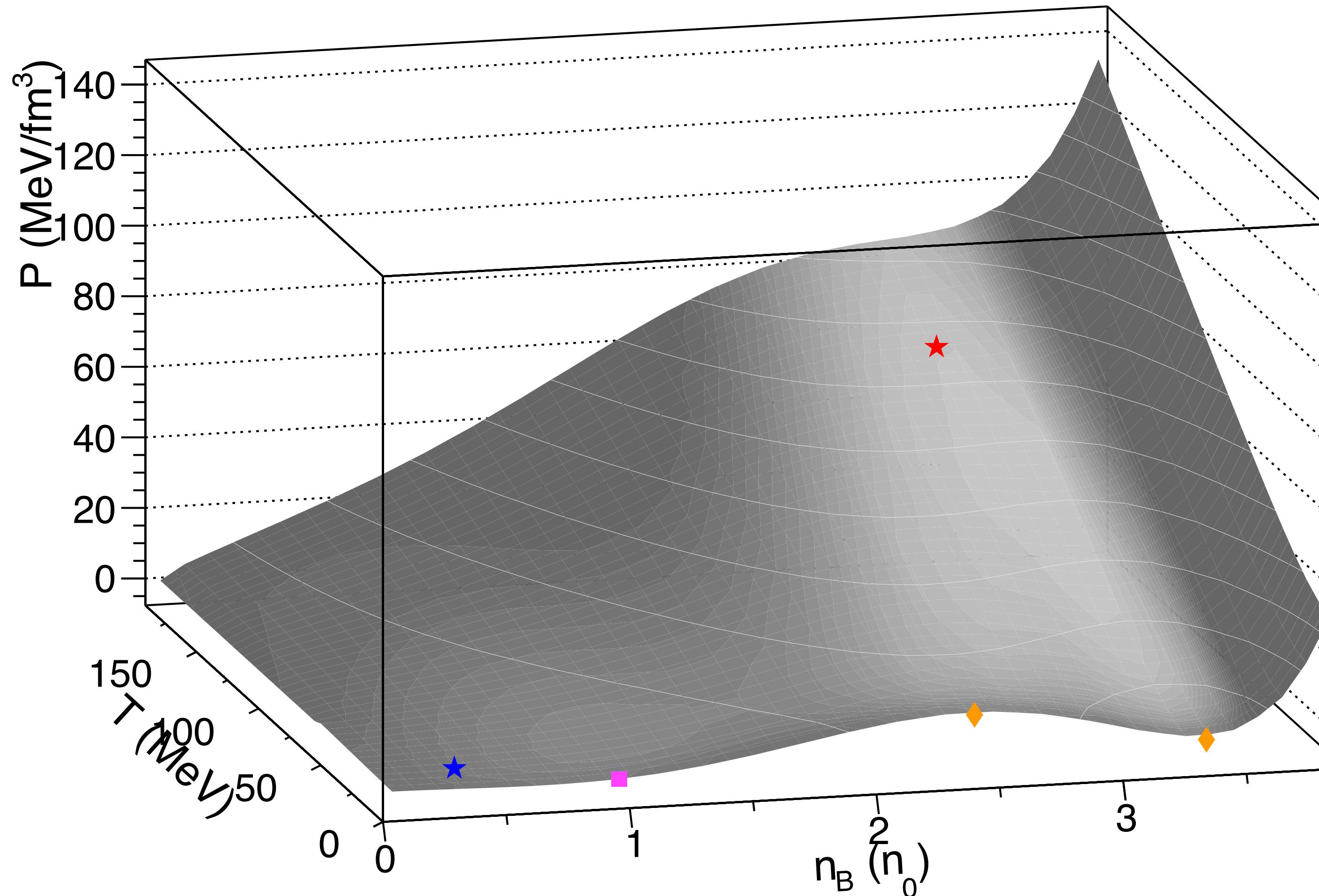


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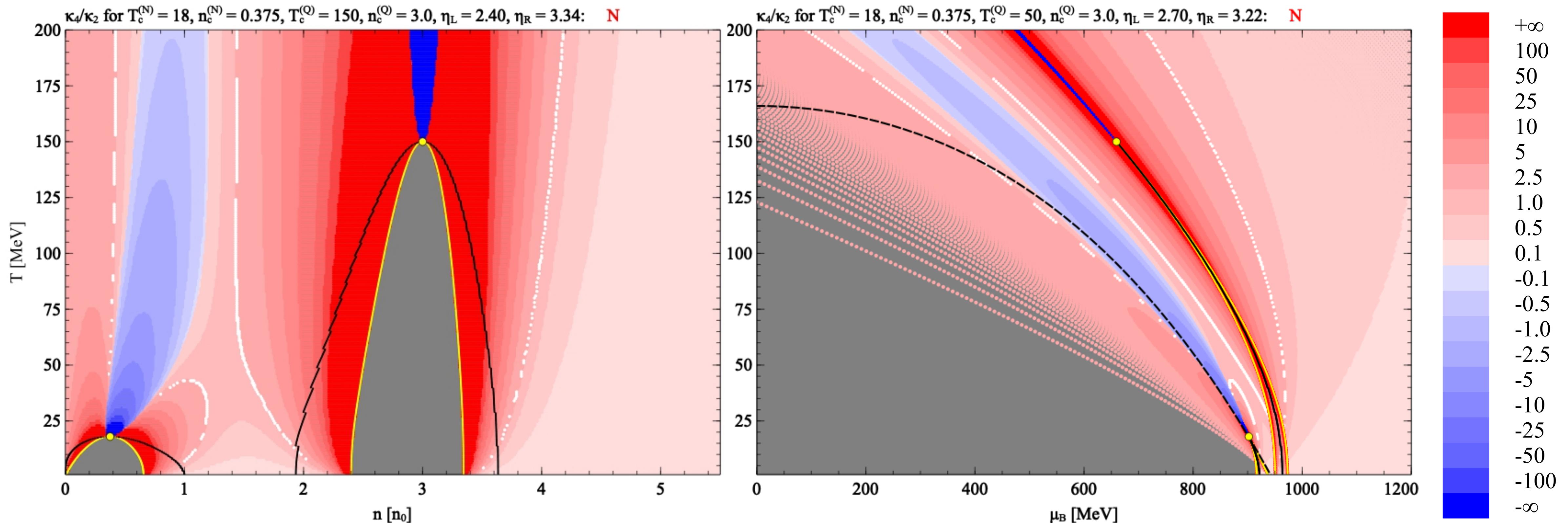


We choose one particular EOS to study



$$n_0 = 0.160 \text{ fm}^{-3}, E_B = -16.3 \text{ MeV}, T_{c(N)} = 18 \text{ MeV}, n_{c(N)} = 0.375n_0, T_{c(Q)} = 150 \text{ MeV}, n_{c(Q)} = 3n_0, \eta_L = 2.4n_0, \eta_R = 3.34n_0$$

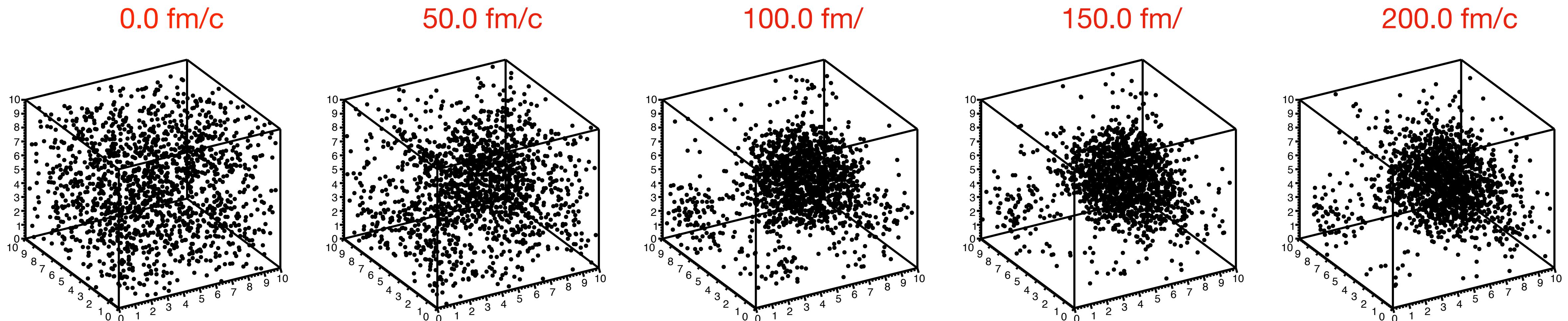
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SMASH* results

We simulate dense nuclear matter in a periodic box.
First, we are going to initialize the system inside the spinodal region:
super interesting and biggest effects expected!



$L = 10 \text{ fm}$, $n_B = 0.25n_0$, $T = 1 \text{ MeV}$, $N_{\text{test}} = 200$



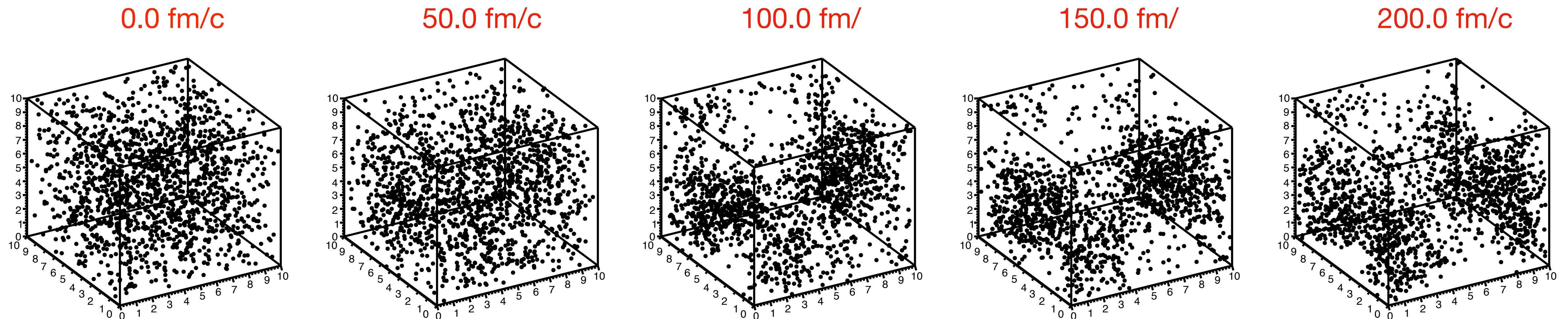
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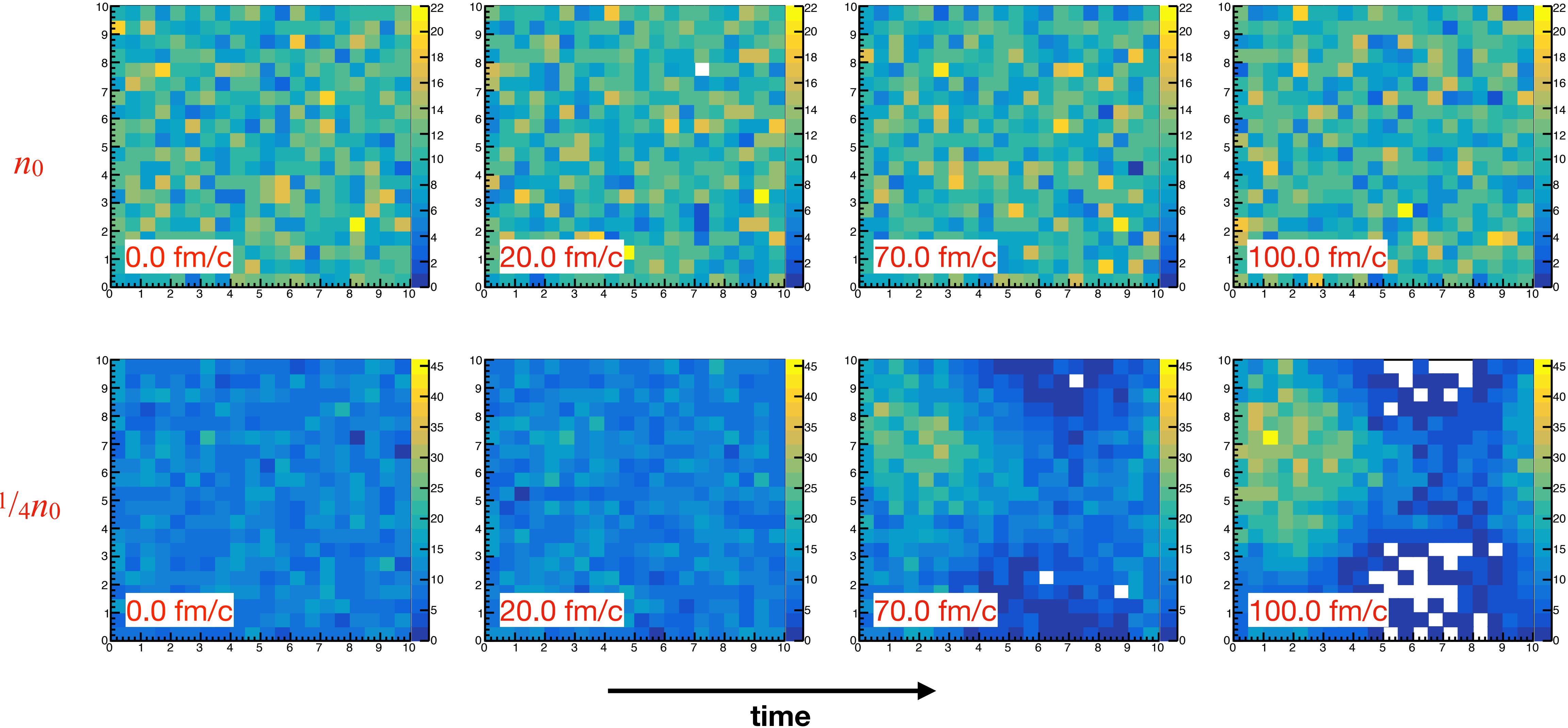
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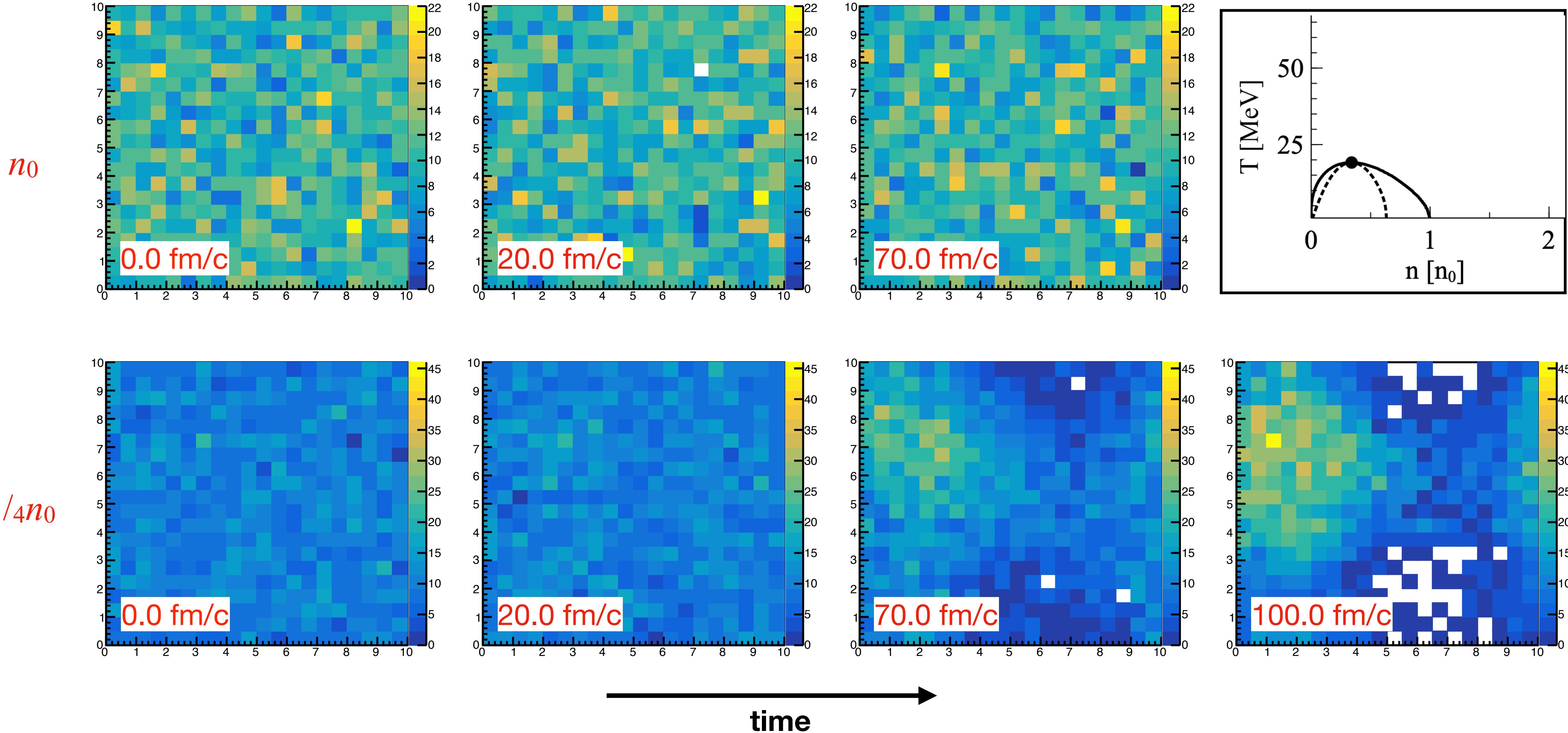
particle number projection onto the xy -plane



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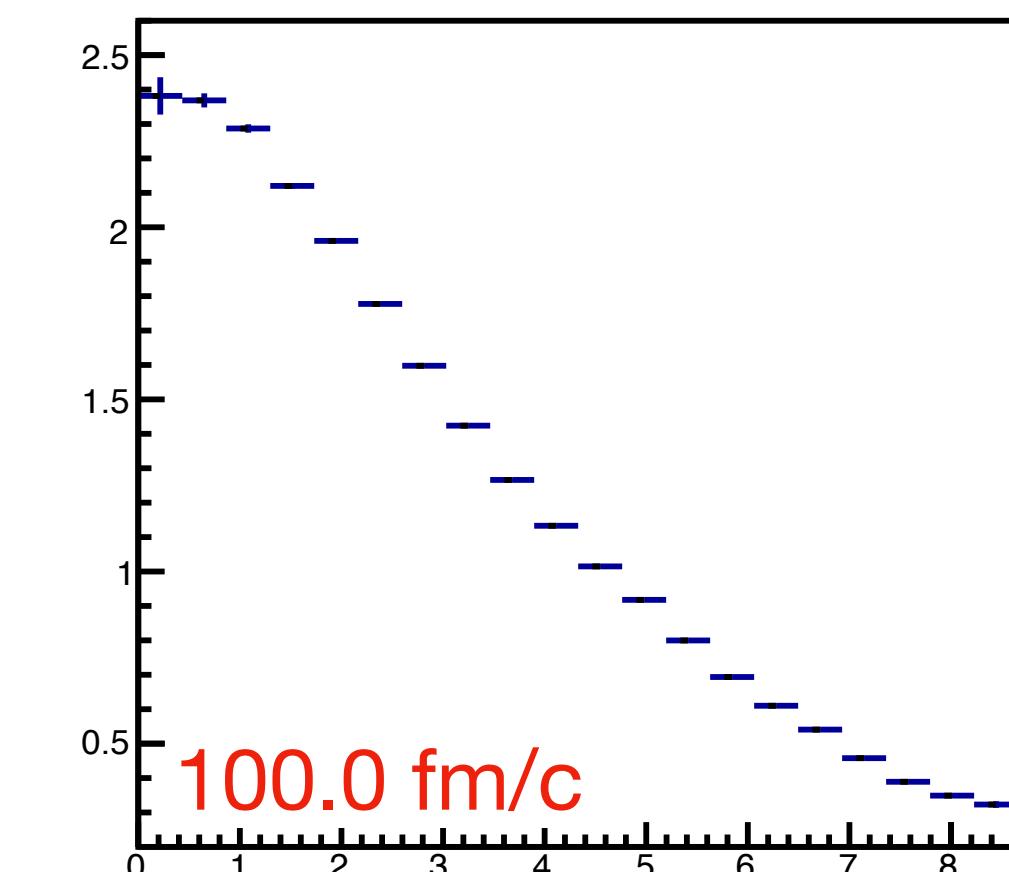
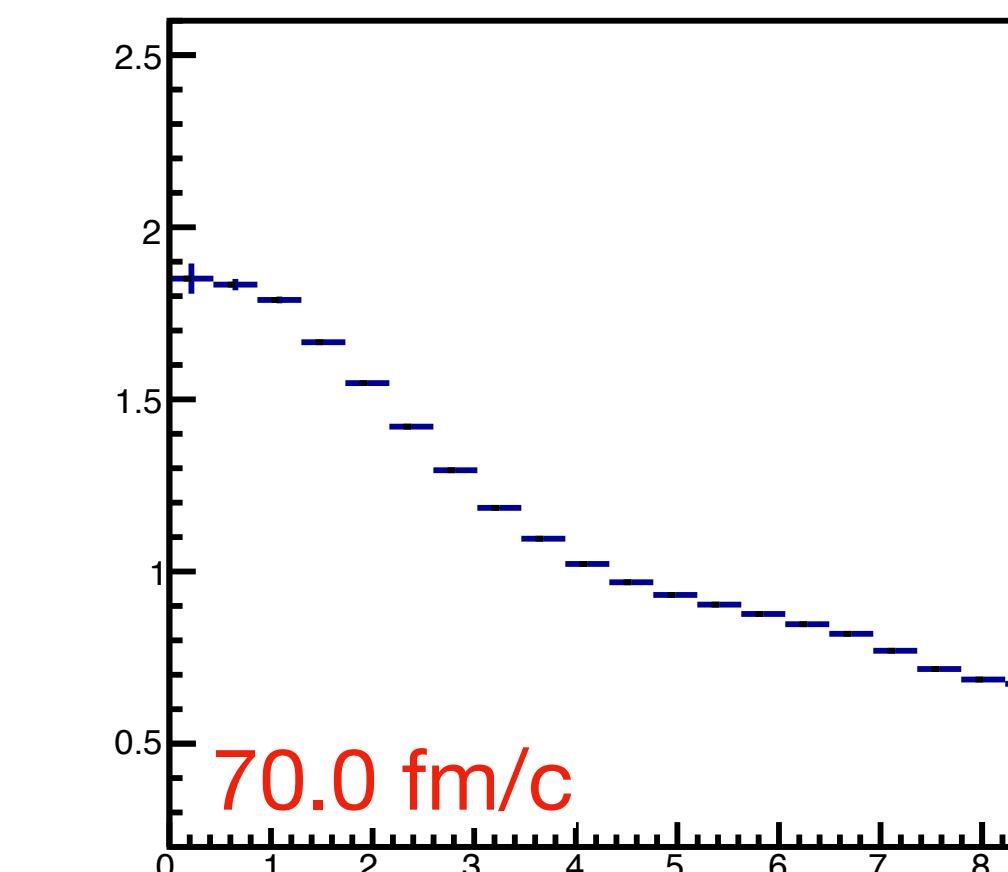
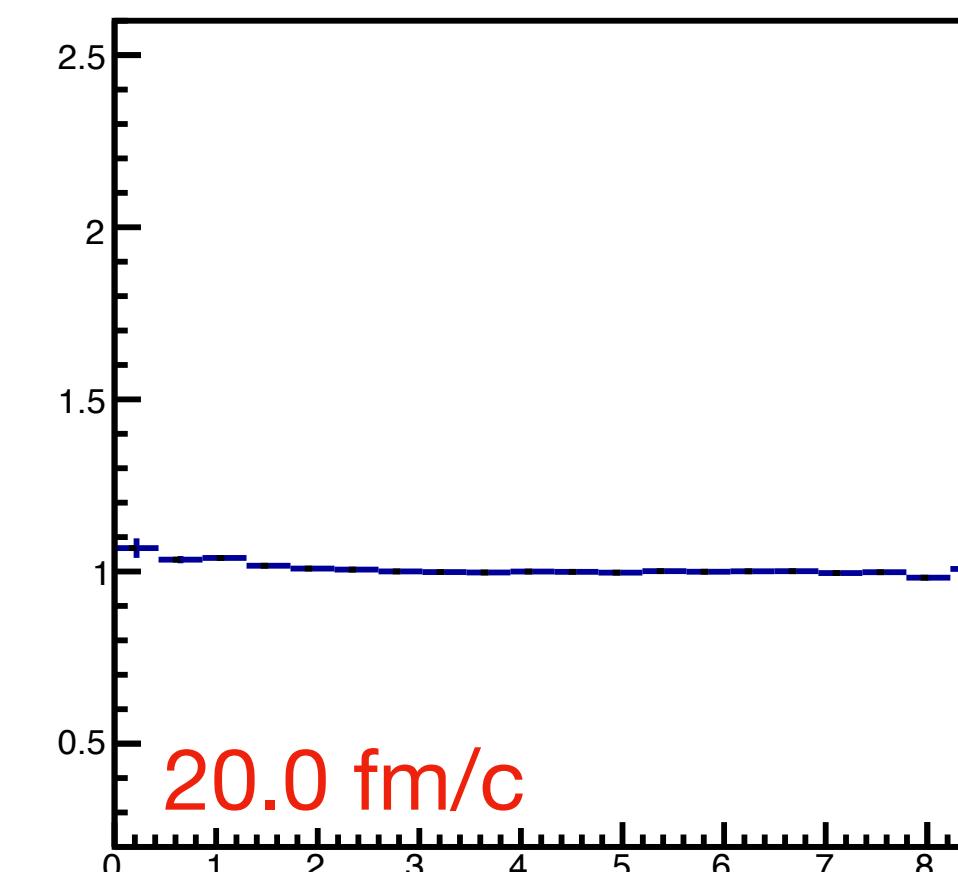
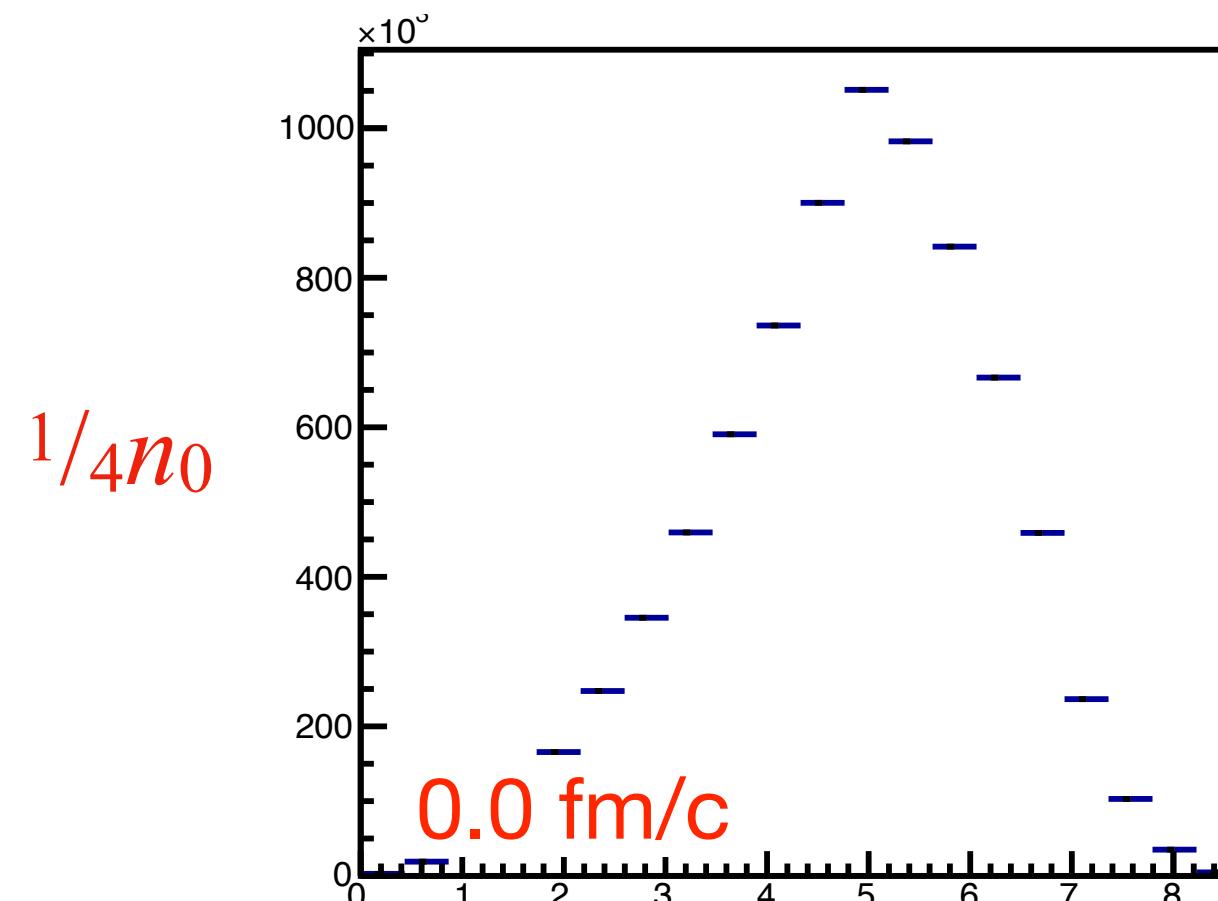
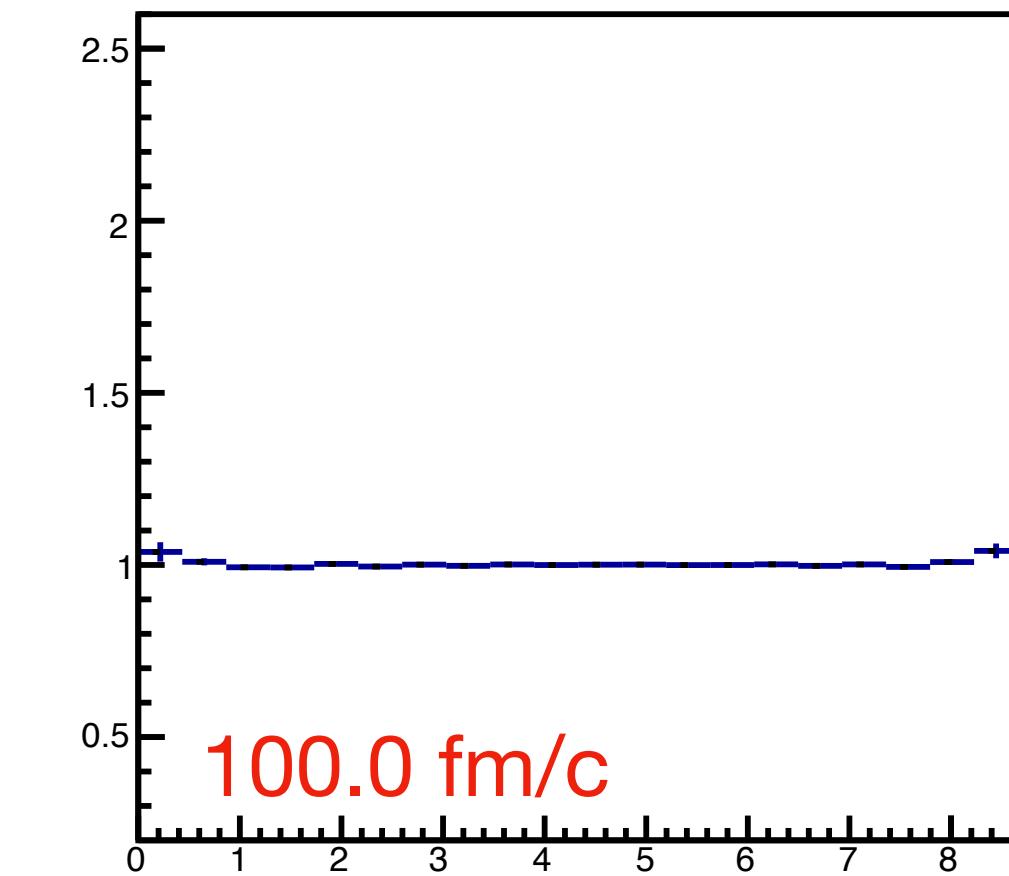
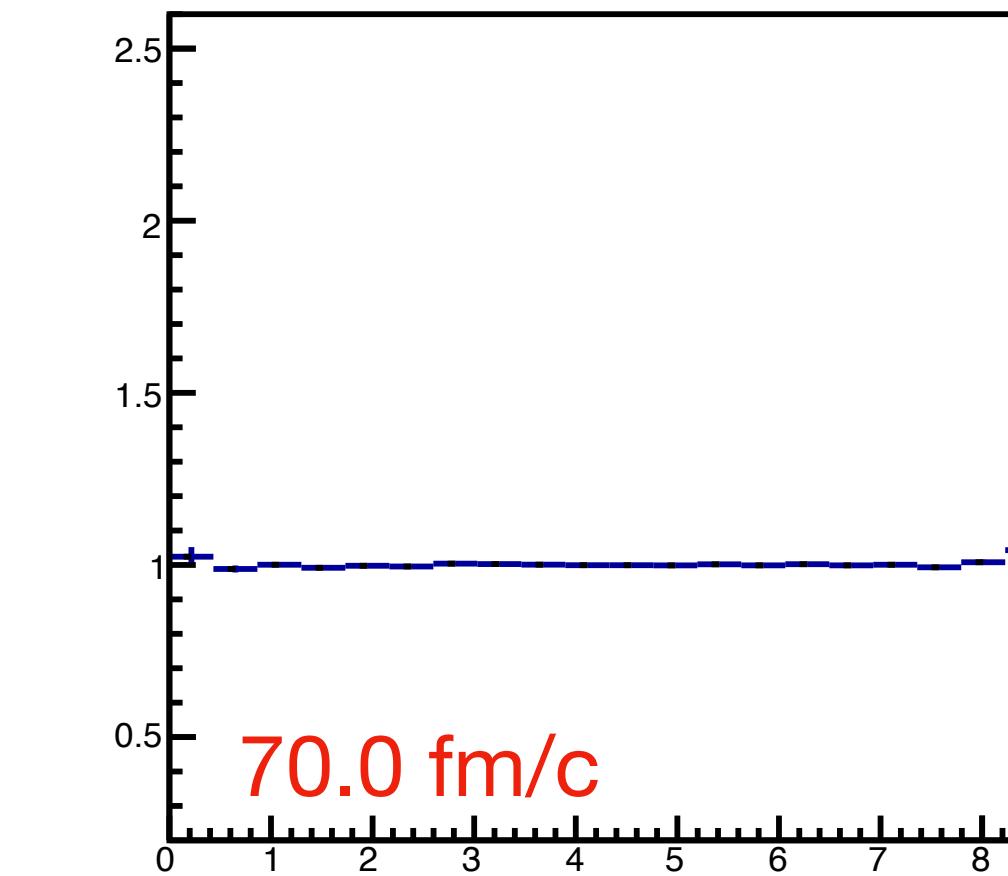
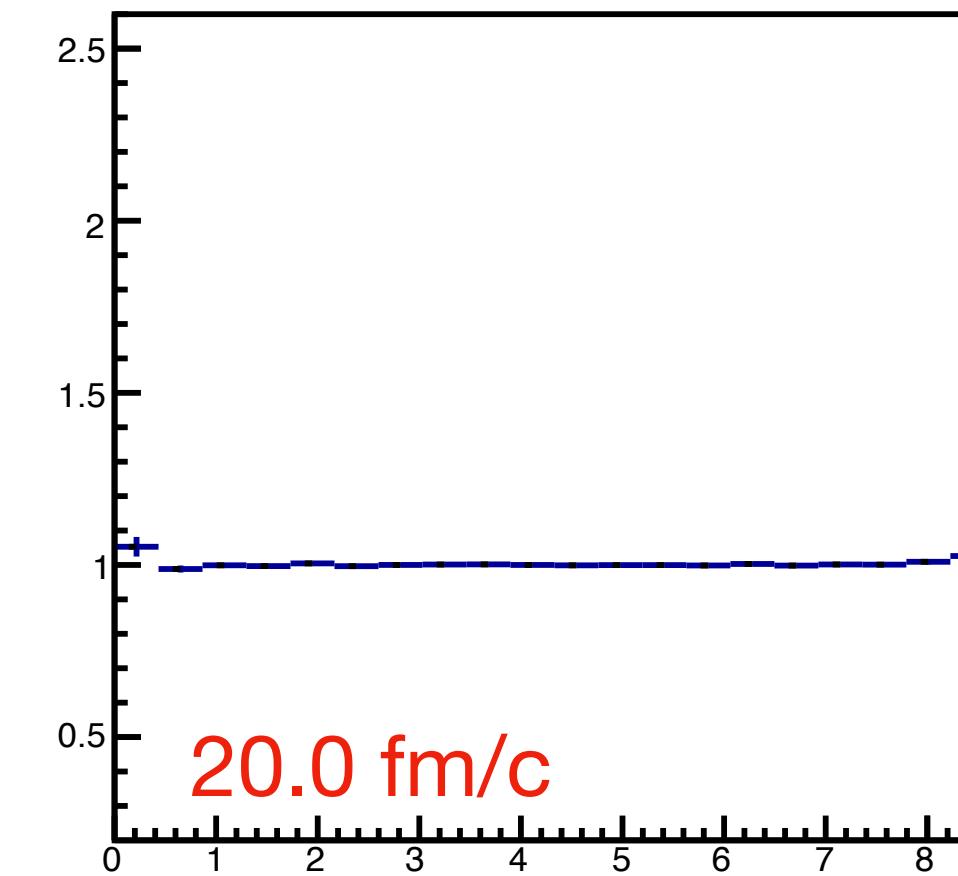
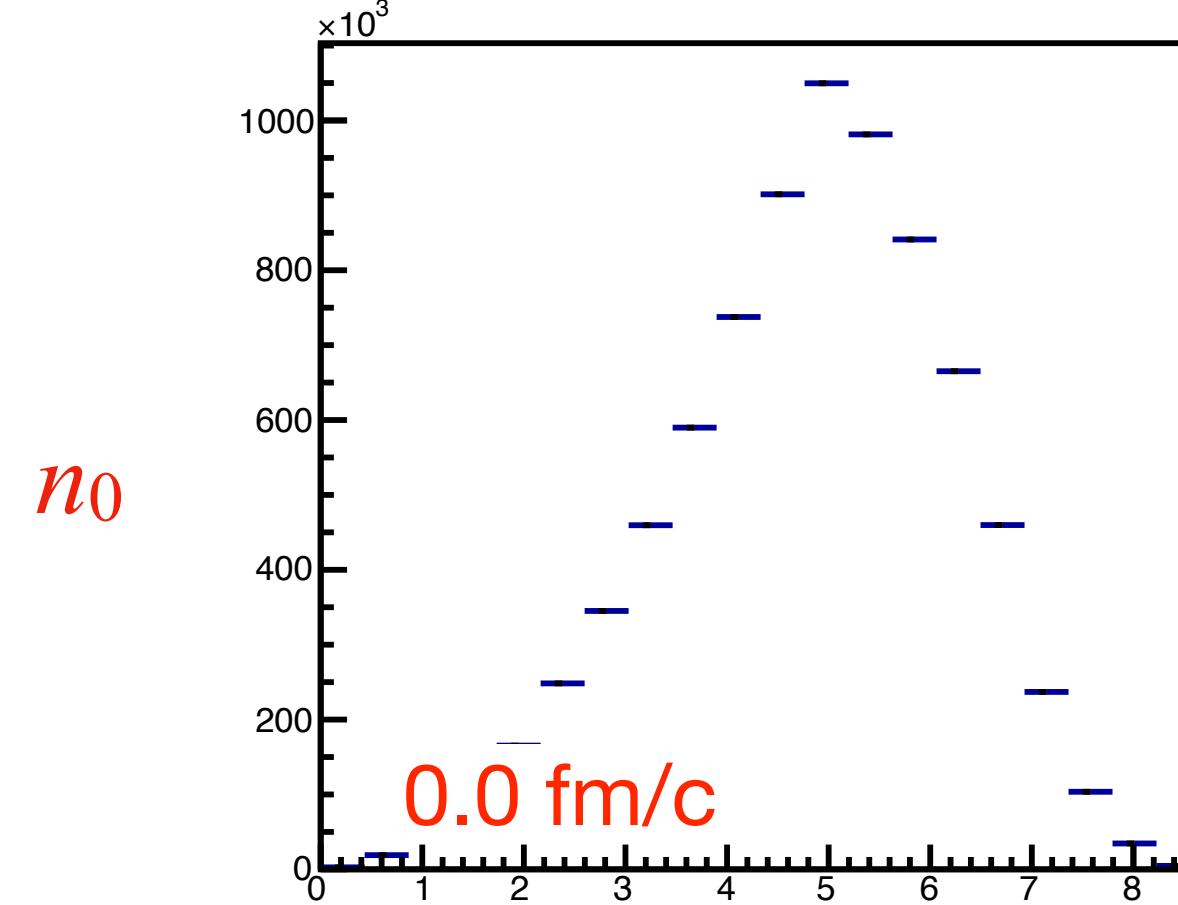
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SMASH results: periodic box

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pair separation distribution (scaled for $t > 0$)



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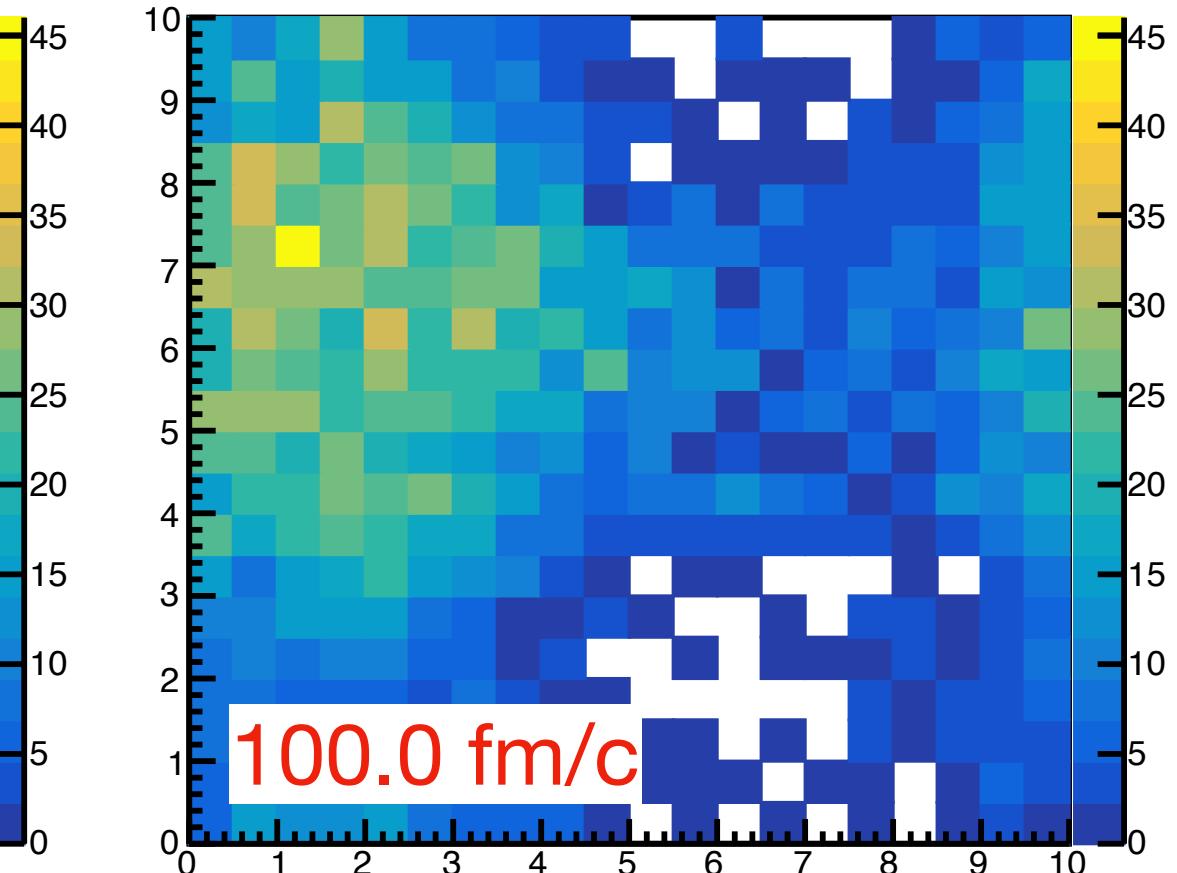
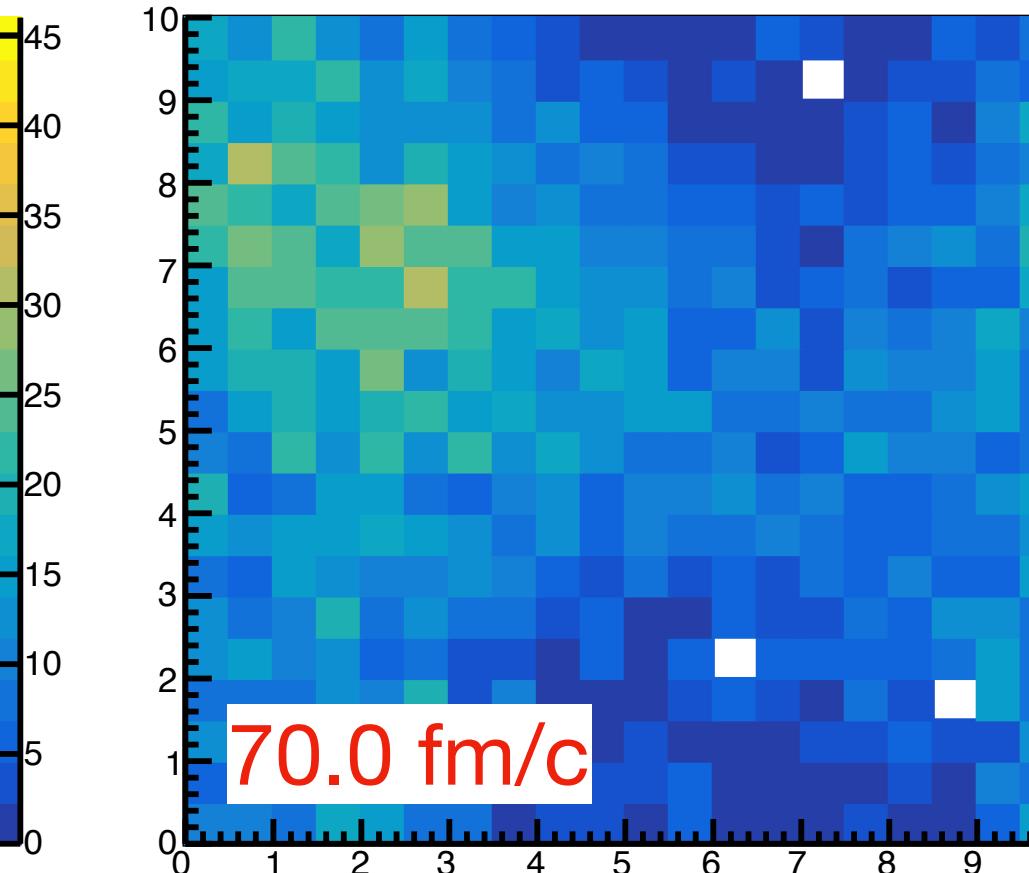
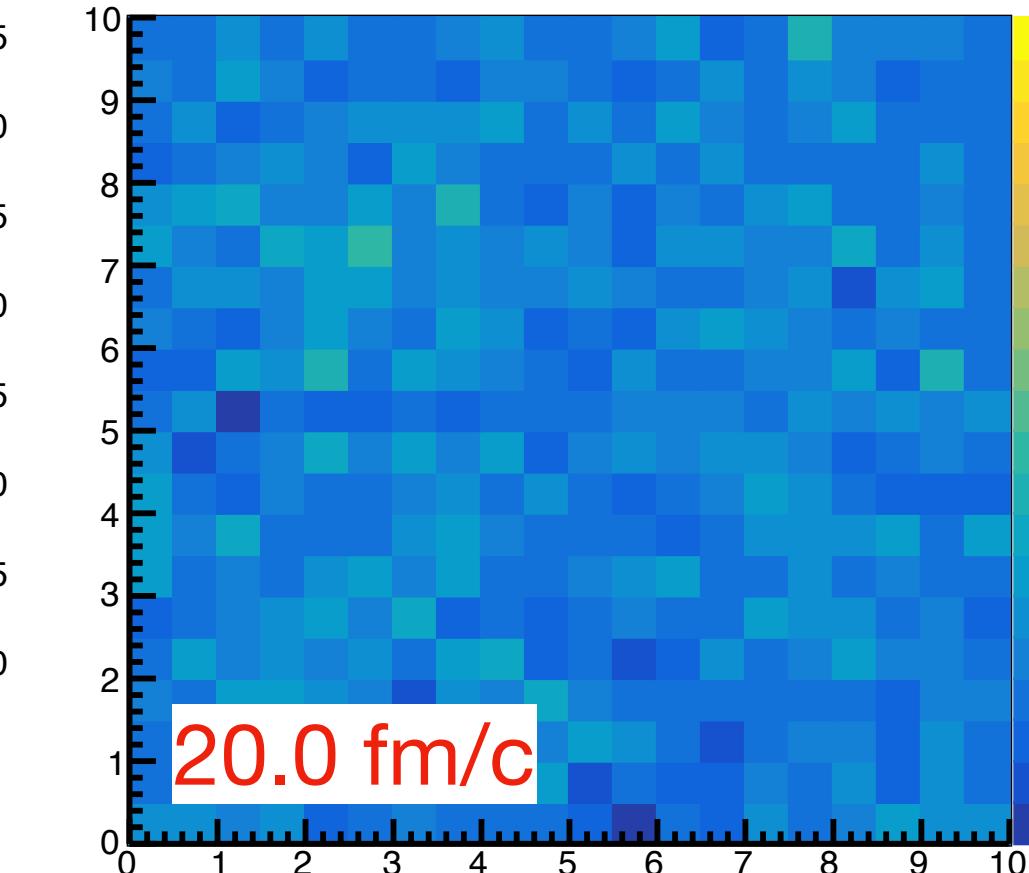
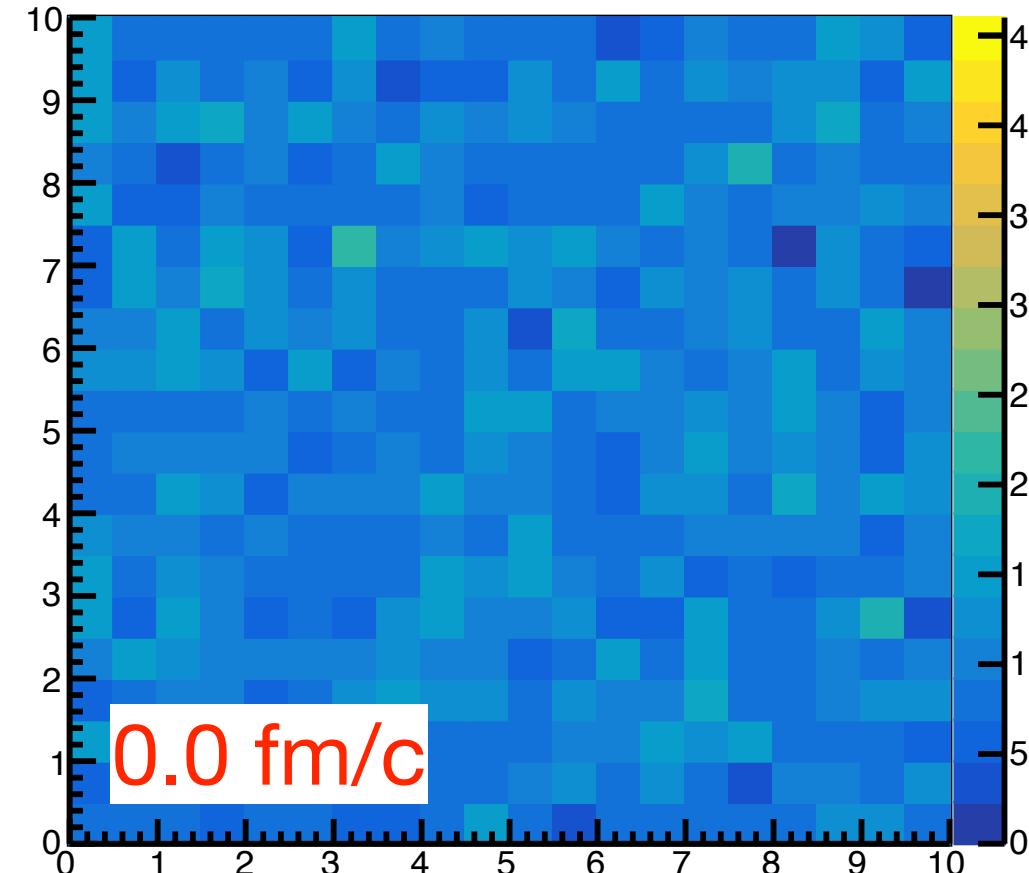
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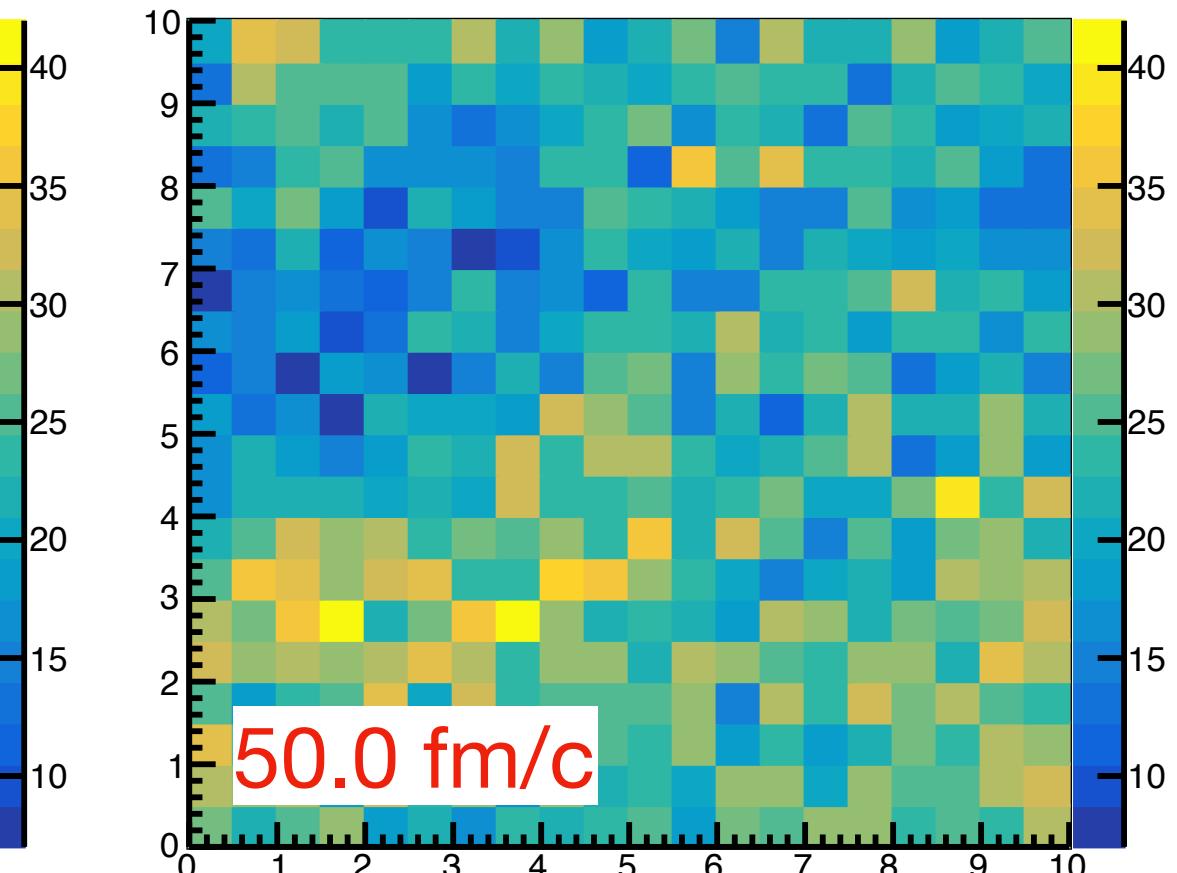
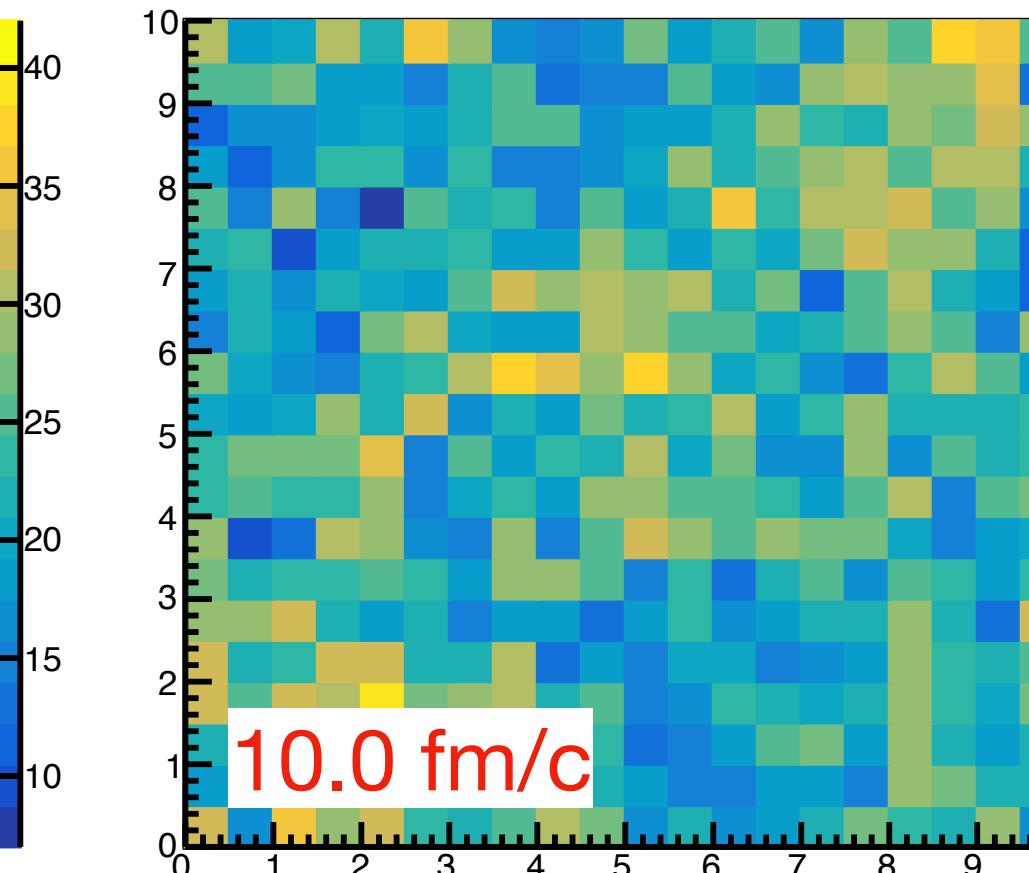
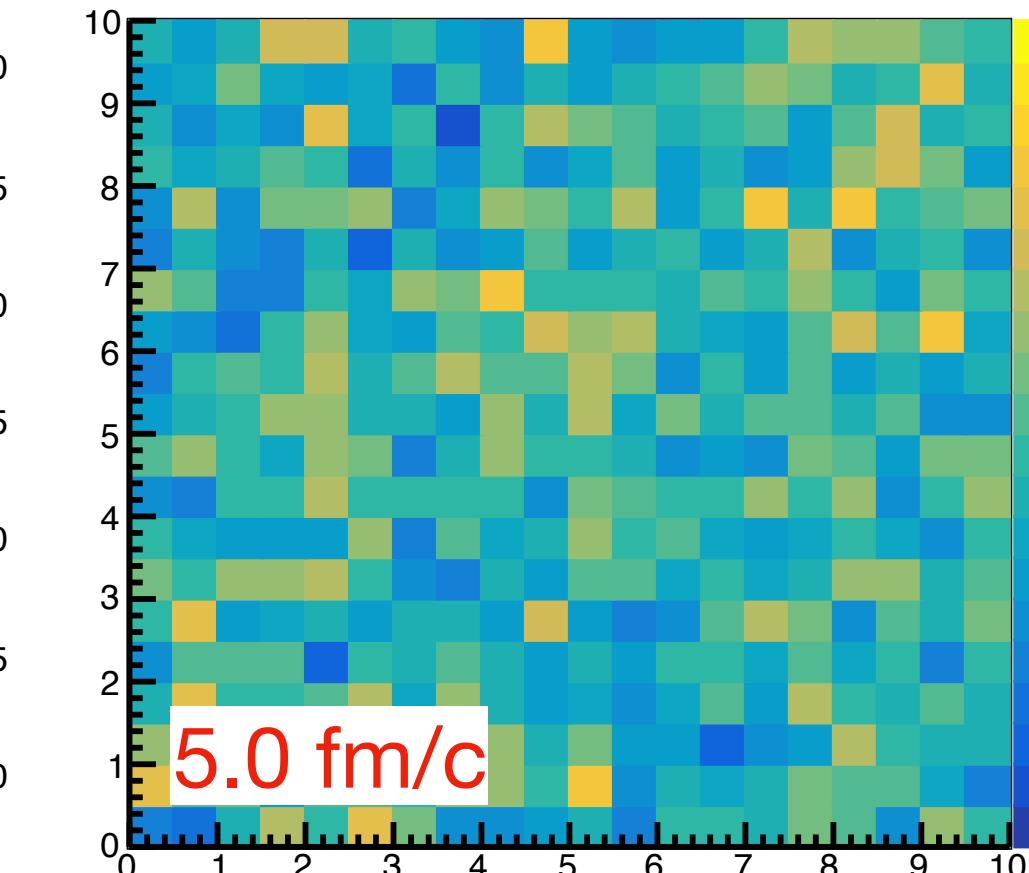
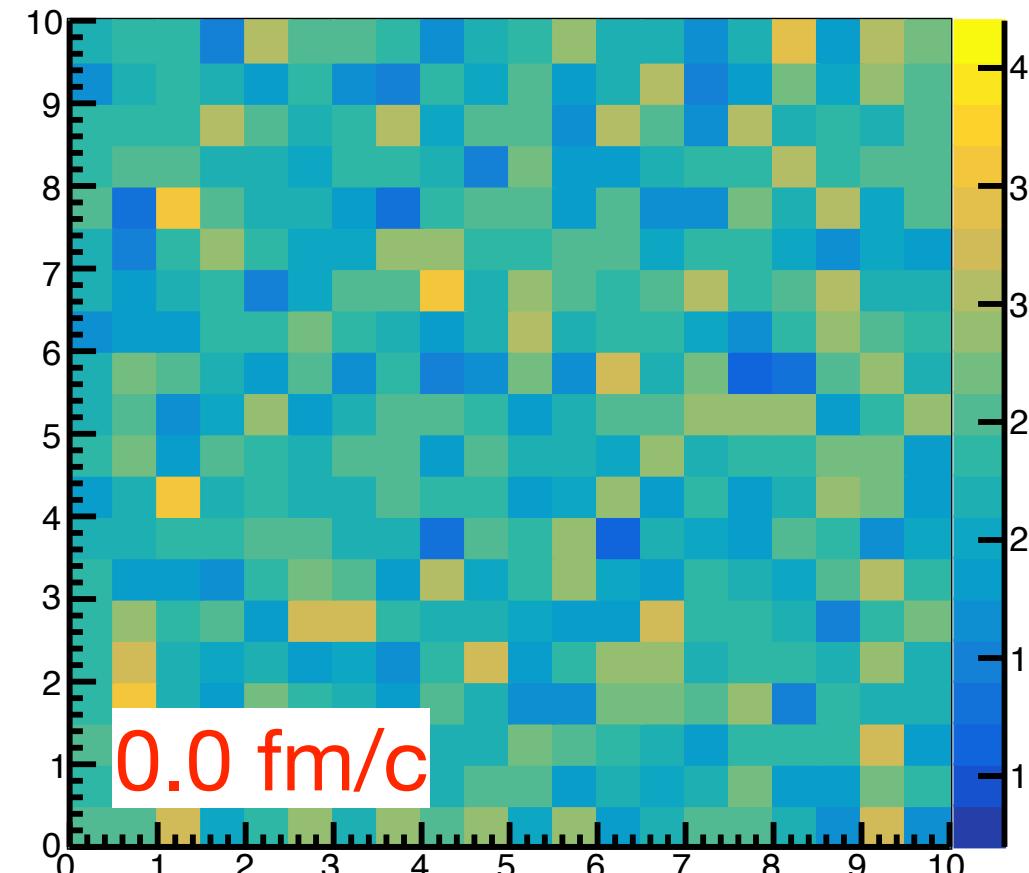
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$1/4n_0$



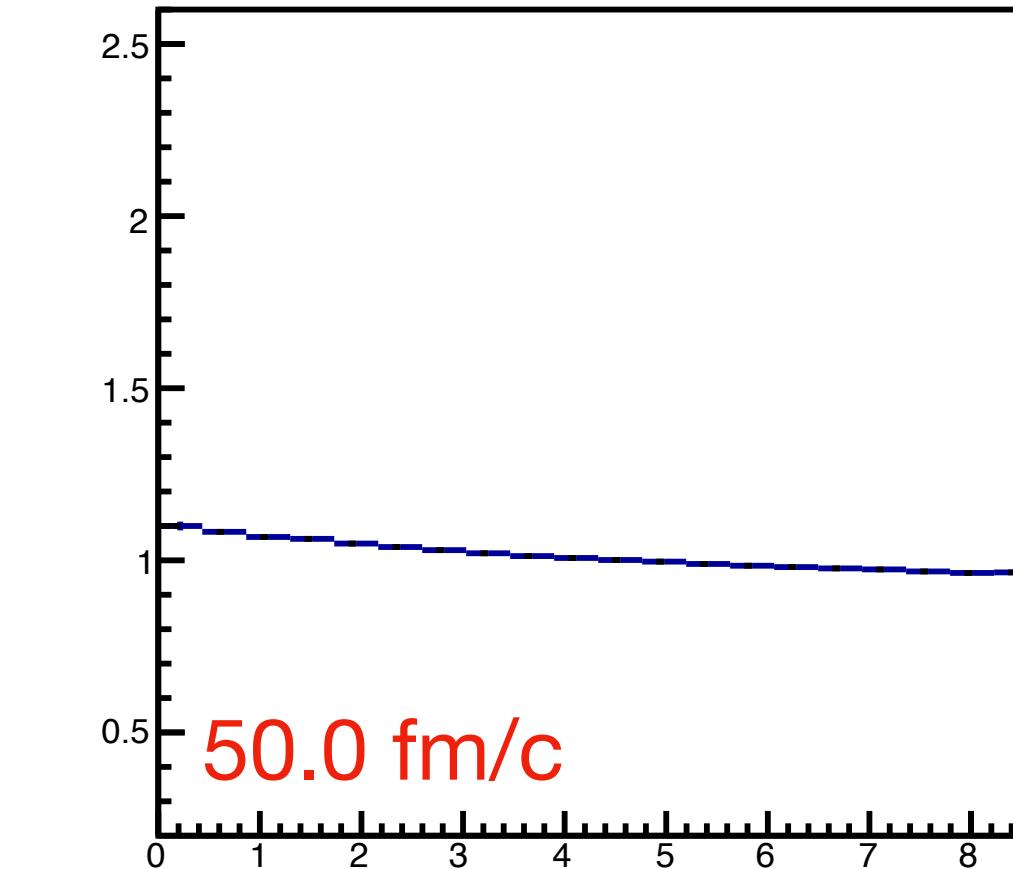
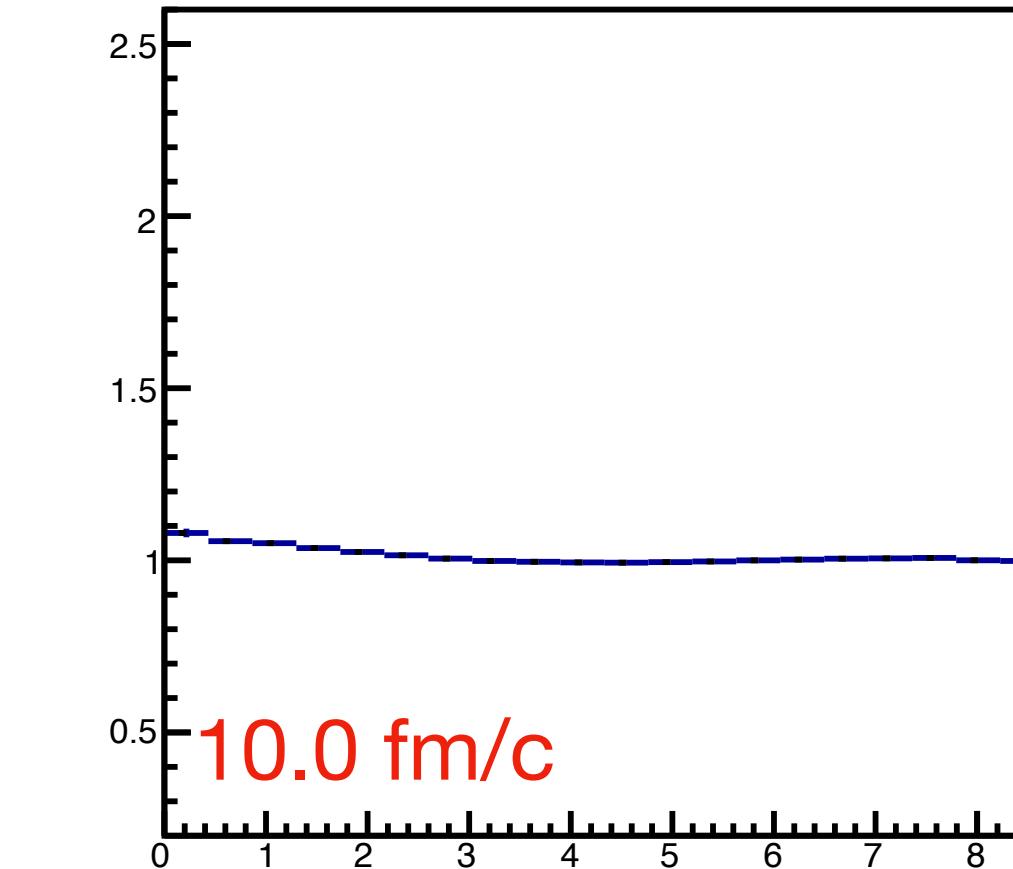
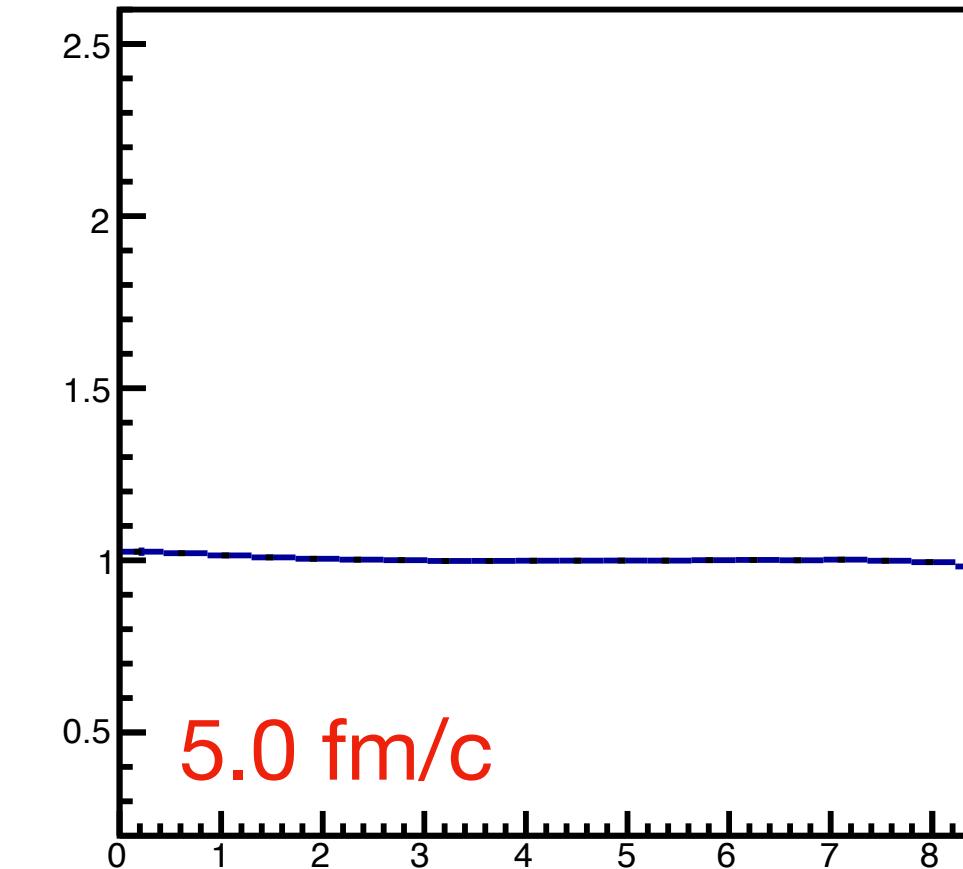
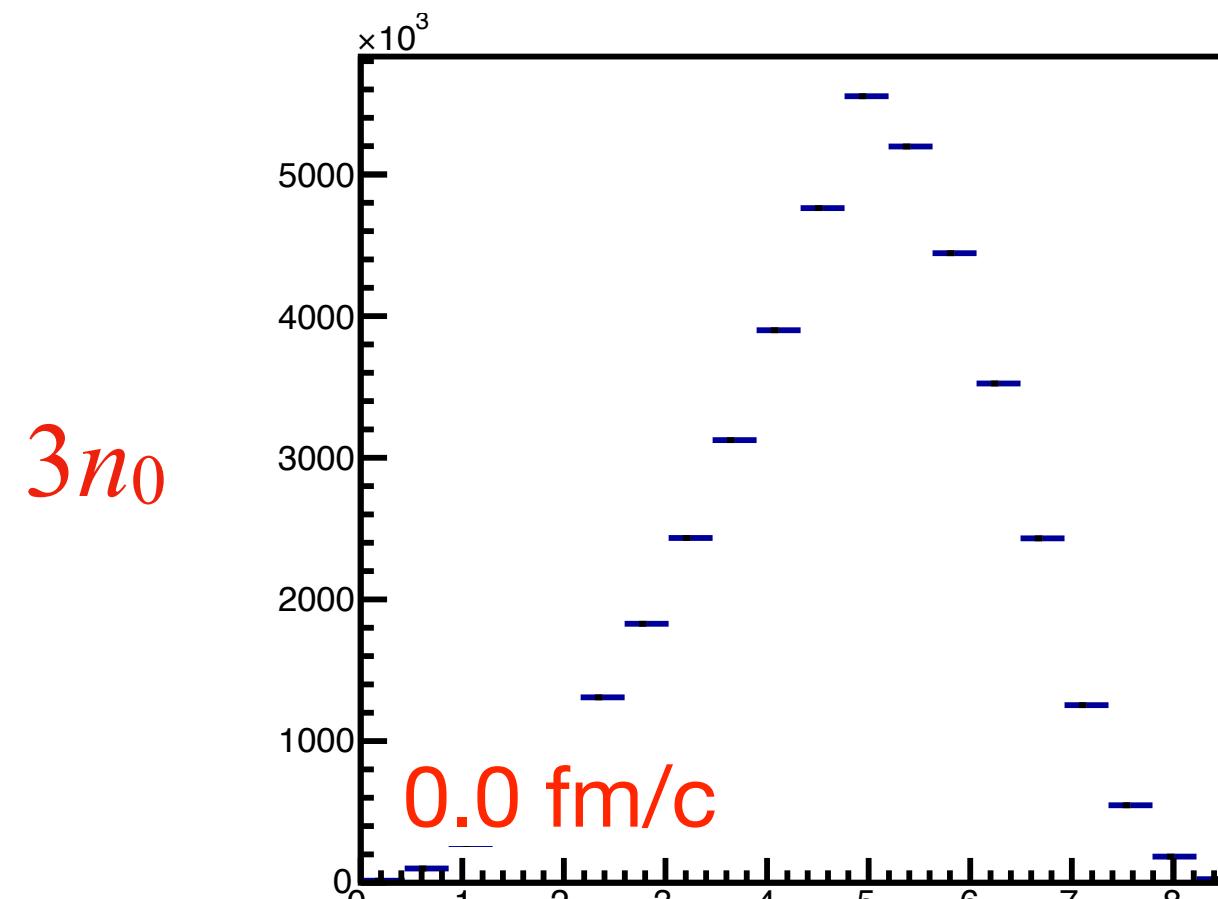
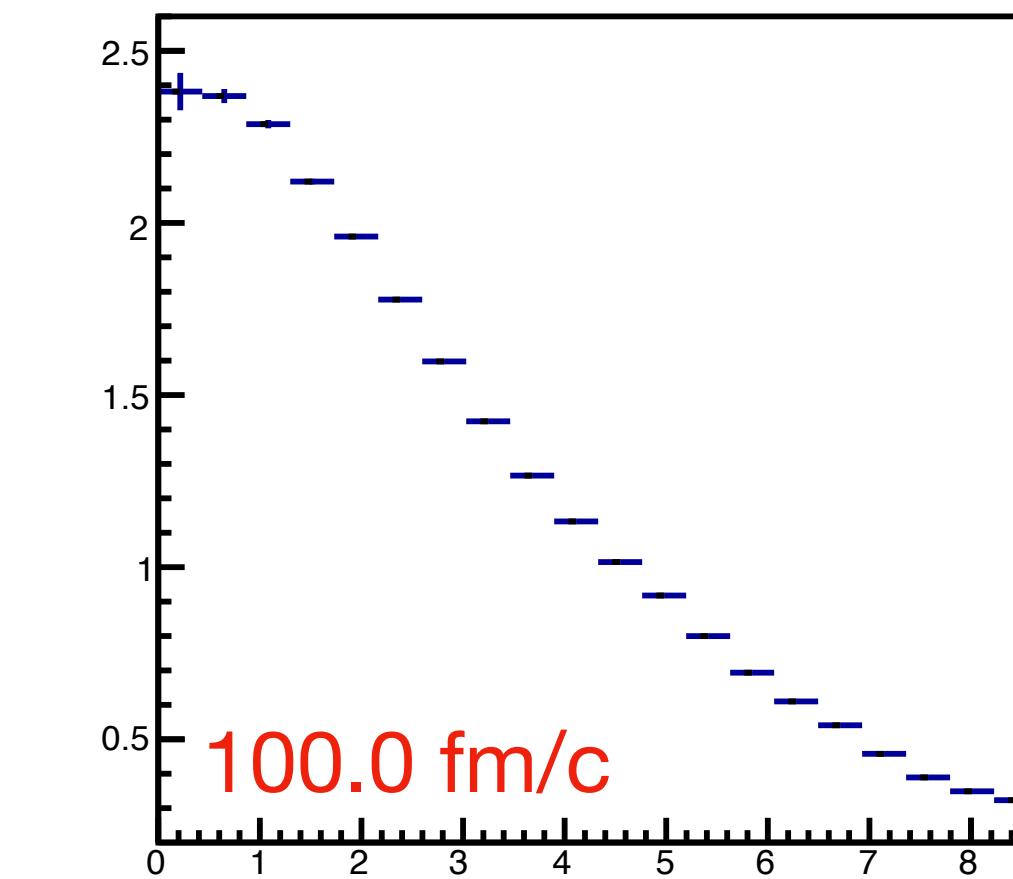
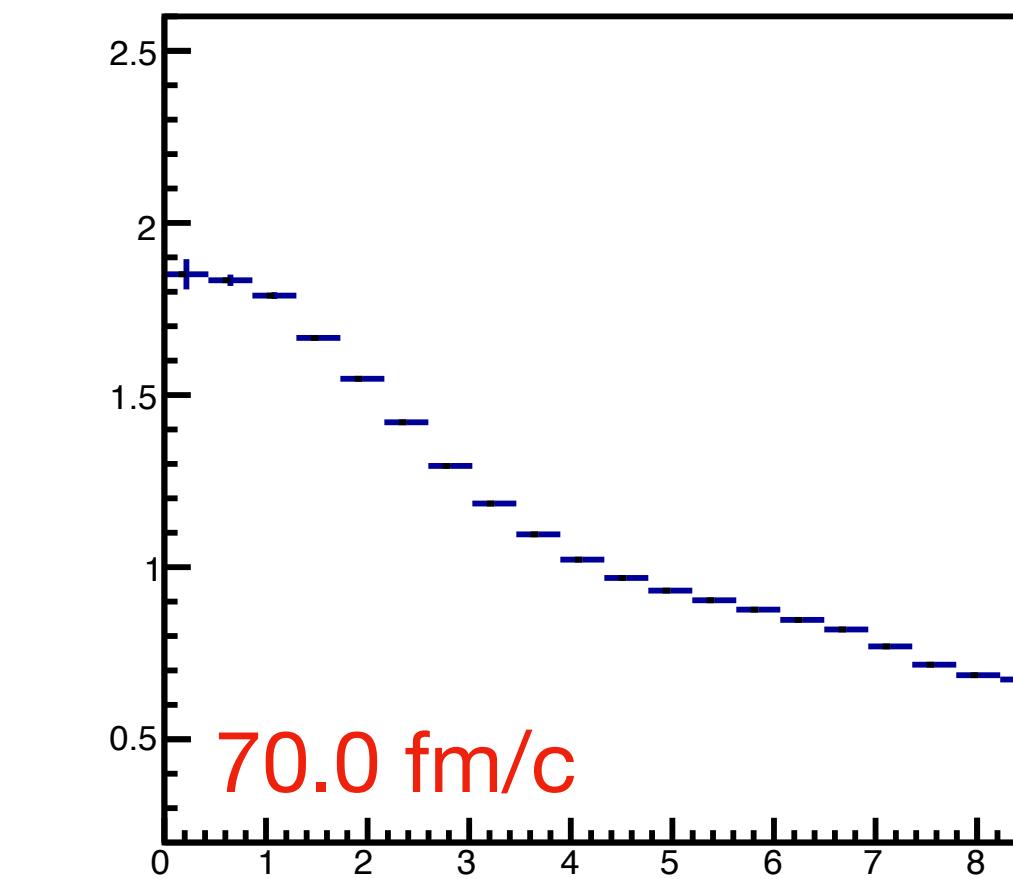
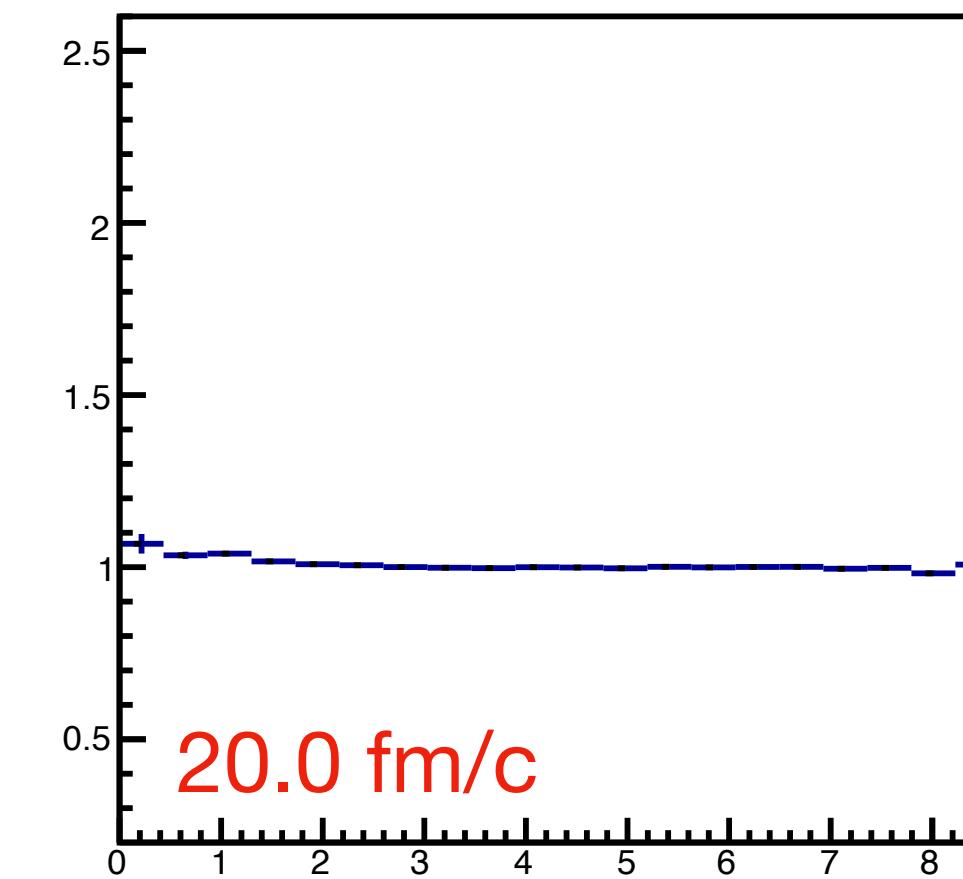
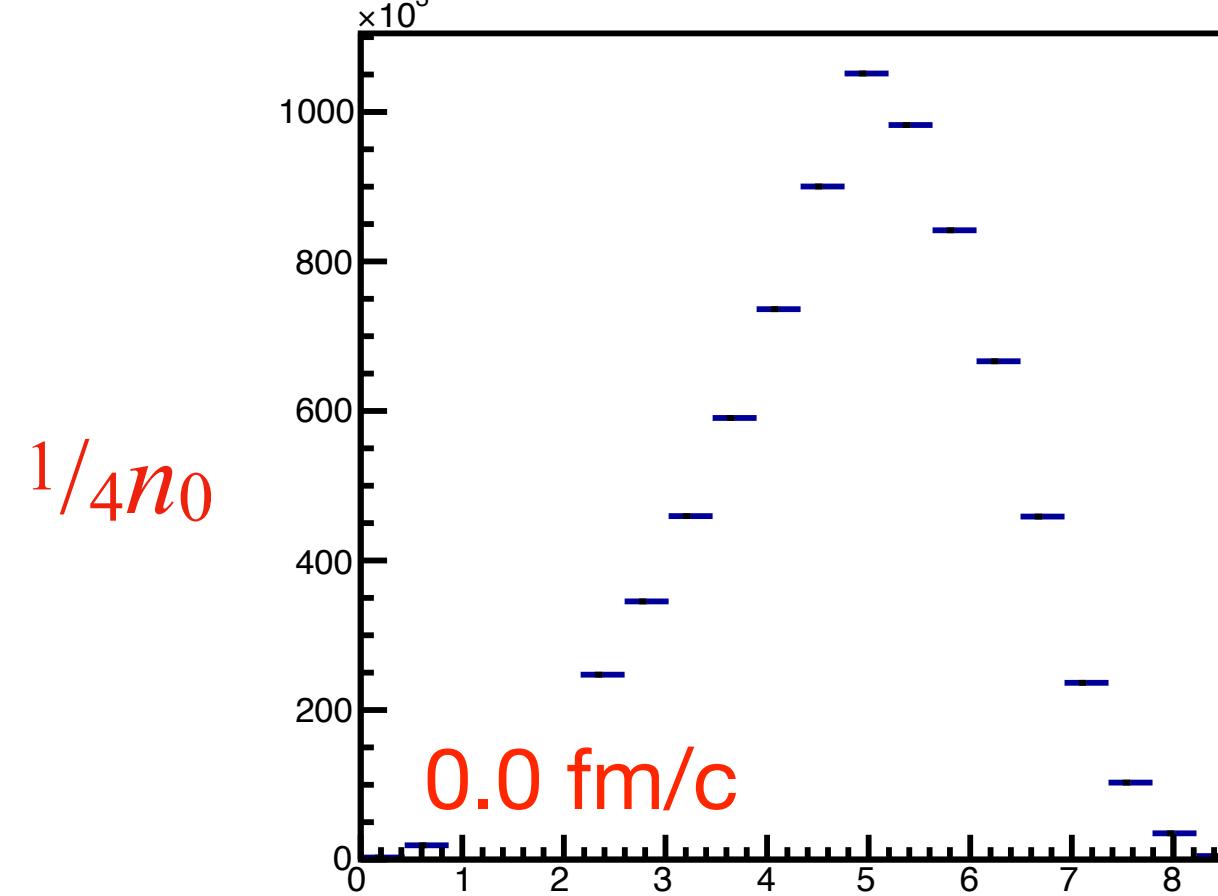
$3n_0$



SMASH results: periodic box

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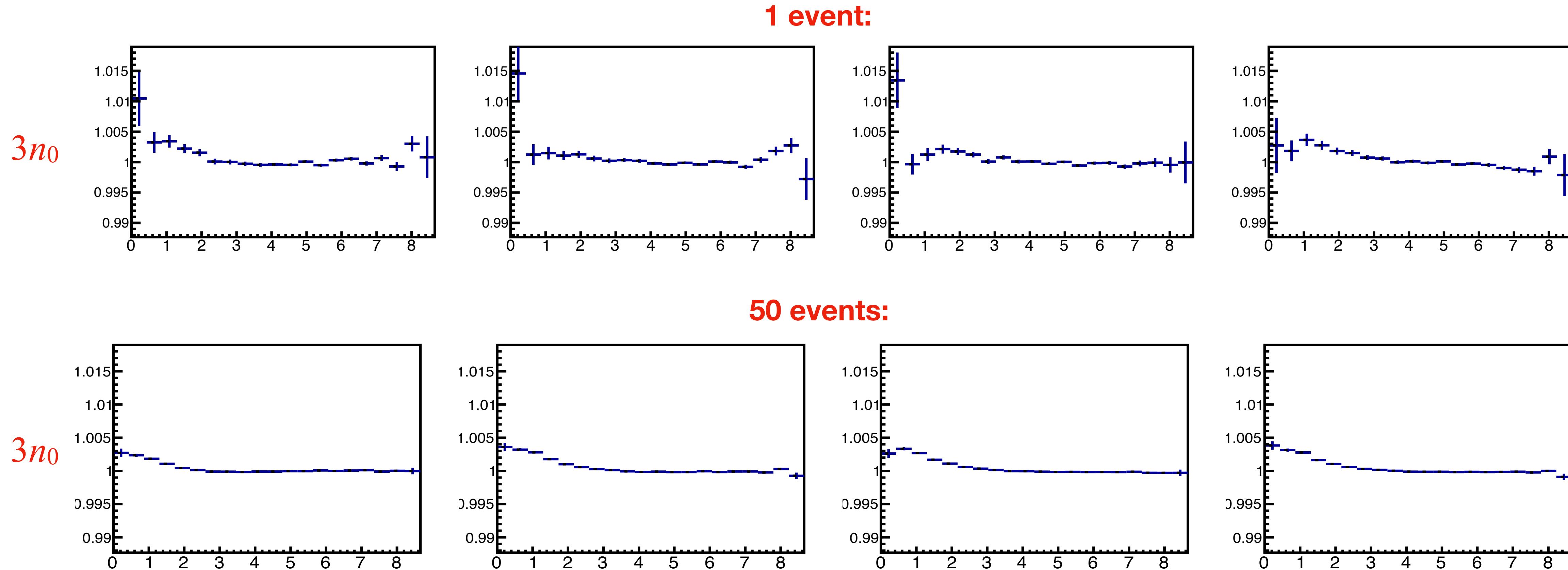
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SMASH results: periodic box averaged over events

$L = 10 \text{ fm}$ $T = 150 \text{ MeV}$

pair separation distribution (scaled for $t > 0$)



Averaging over many events needed especially for correlations at $T=T_c$

How can we *see* the phase separation for many events?

We can divide our box into cells (bins) and histogram the particle number distribution

SMASH results: periodic box, averaged over 250 events

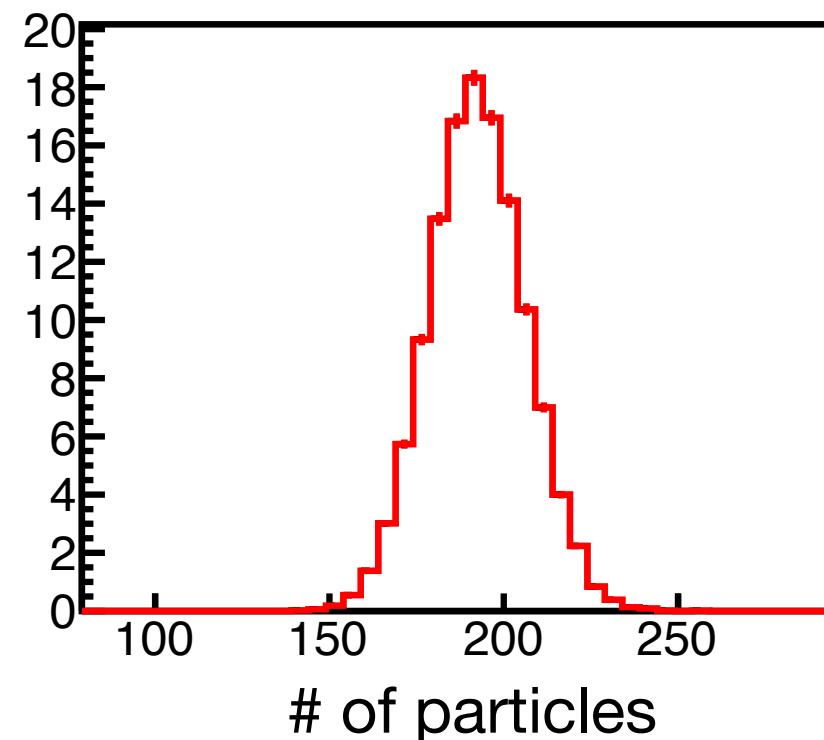
$L = 10 \text{ fm}$

$T = 1 \text{ MeV}$

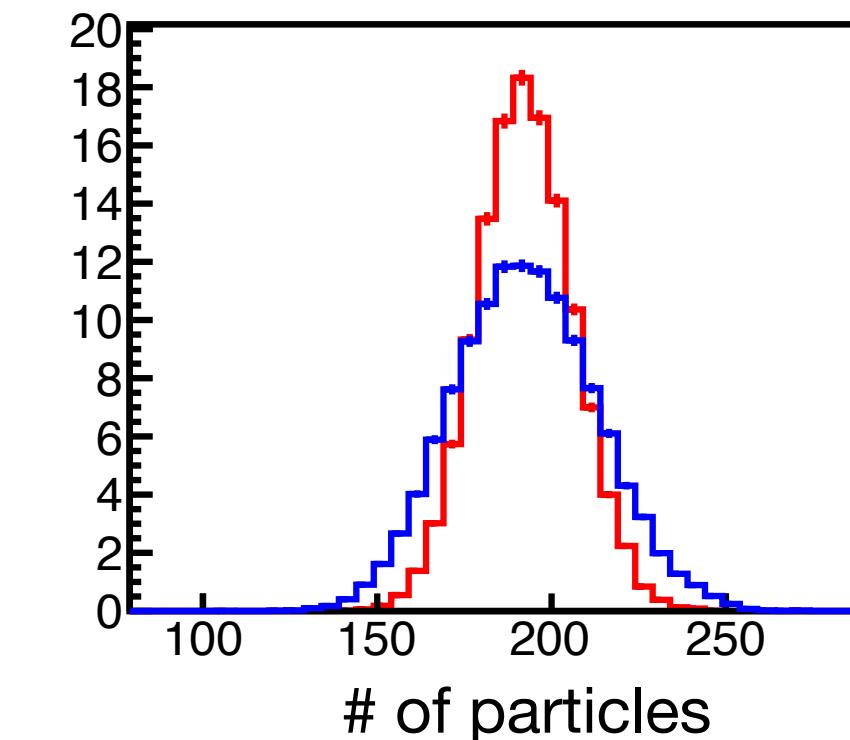
$n_B = 3n_0$

Cell (bin) length $L = 2 \text{ fm}$

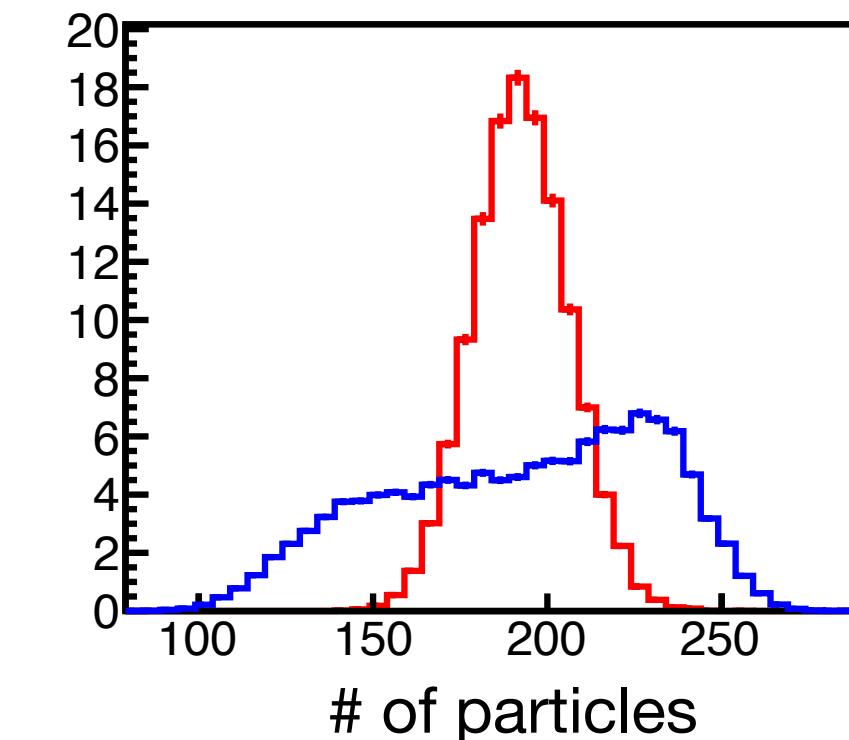
The distribution becomes bimodal as the system separates!



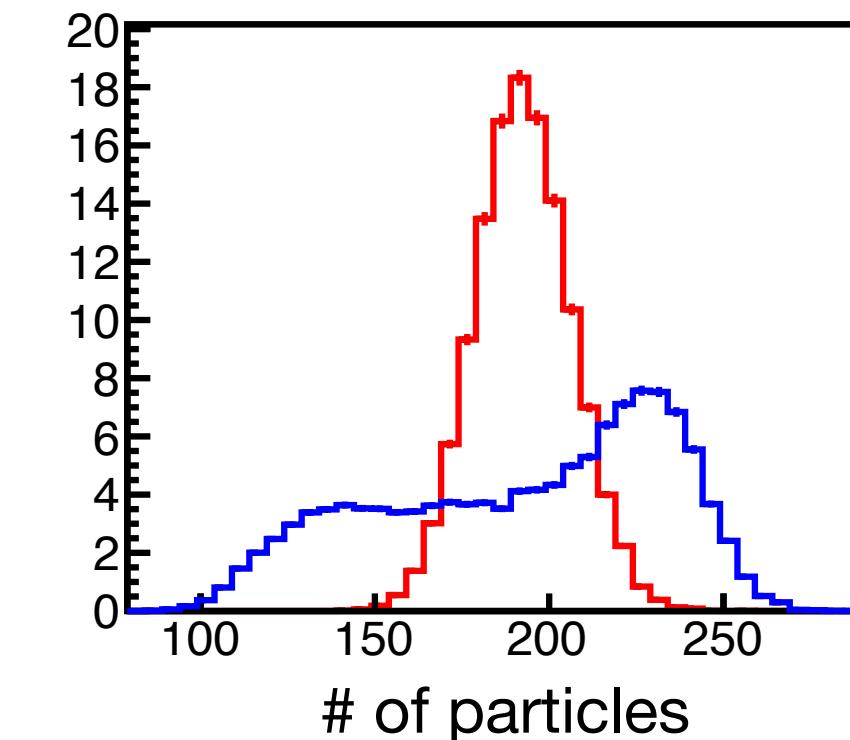
0.0 fm/c



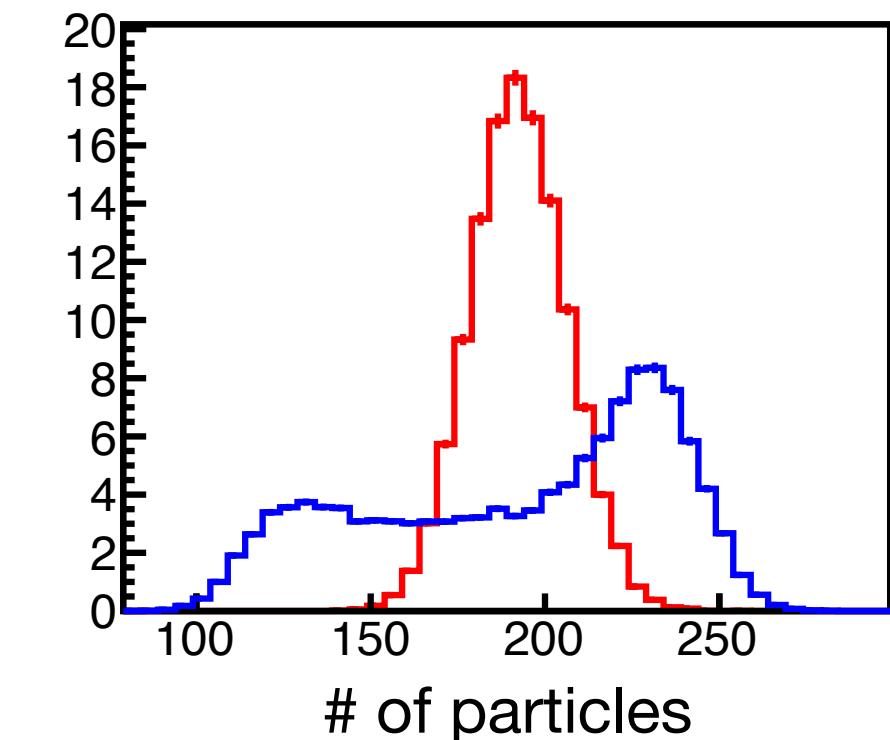
5.0 fm/c



15.0 fm/c



20.0 fm/c



30.0 fm/c

SMASH results: periodic box, averaged over 250 events

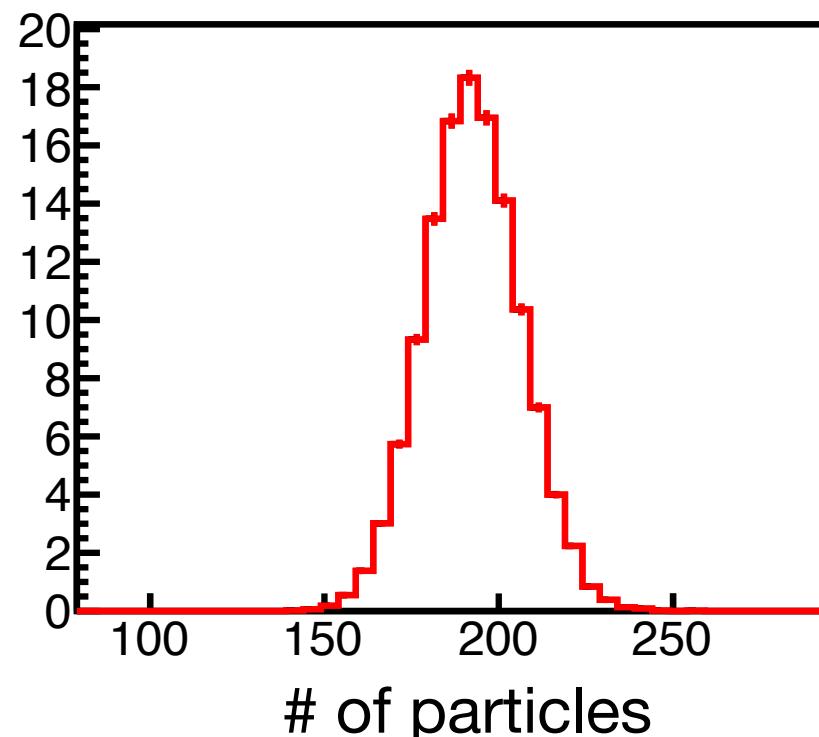
$L = 10 \text{ fm}$

$T = 1 \text{ MeV}$

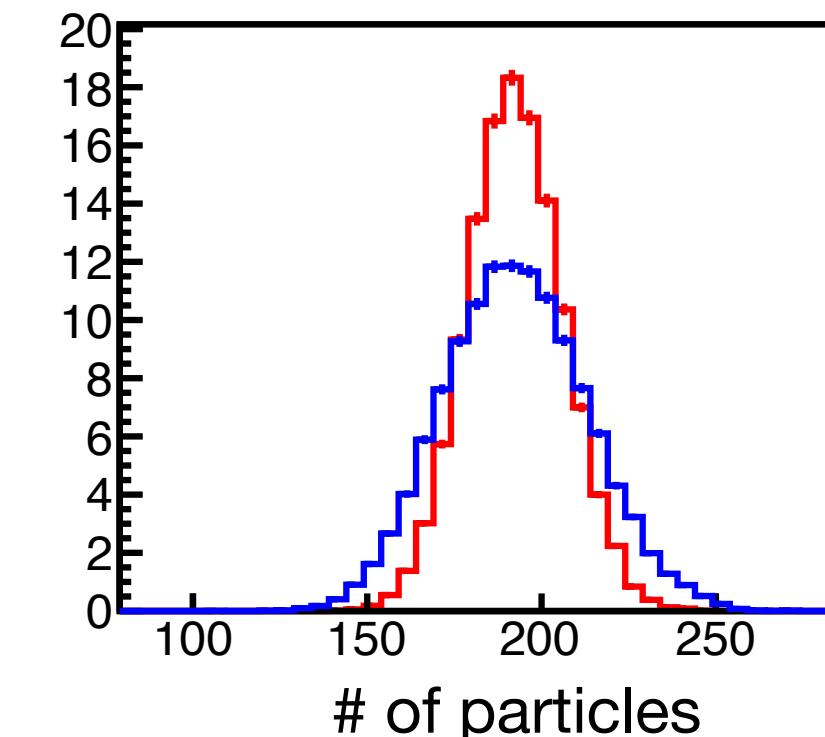
$n_B = 3n_0$

Cell (bin) length $L = 2 \text{ fm}$

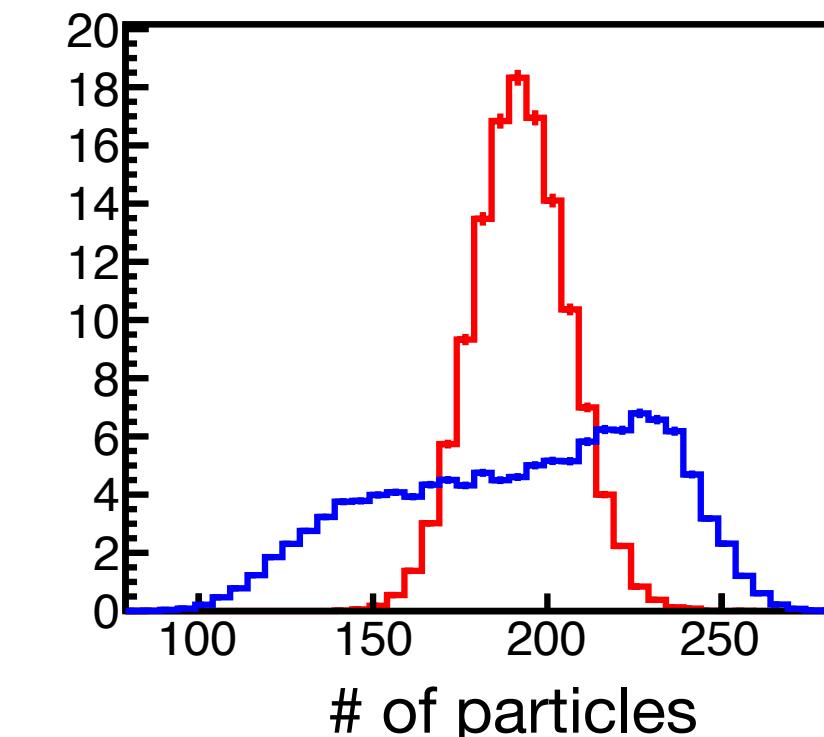
The distribution becomes bimodal as the system separates!



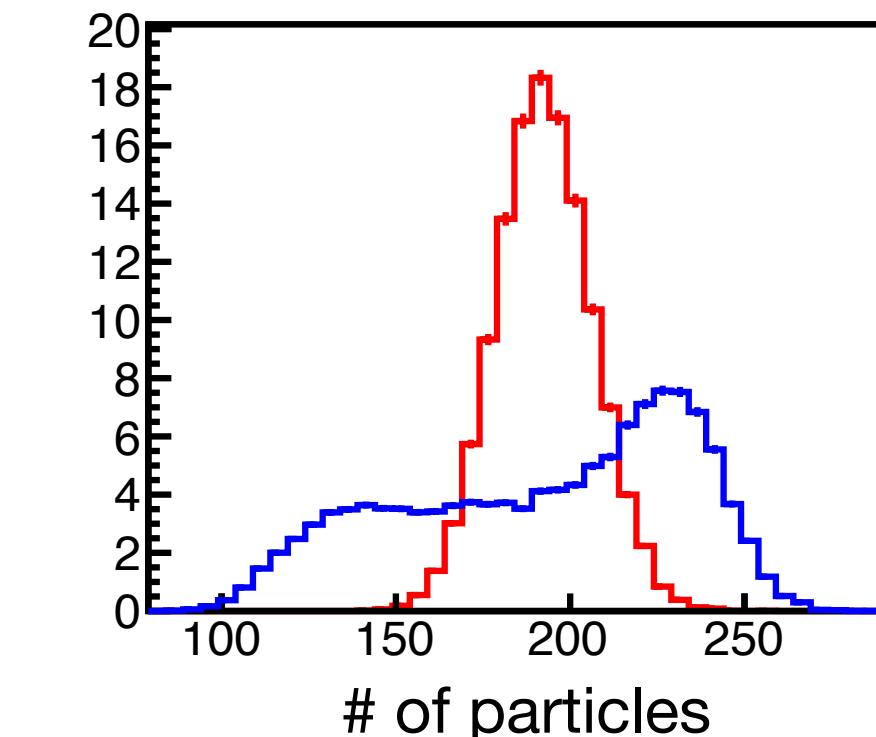
0.0 fm/c



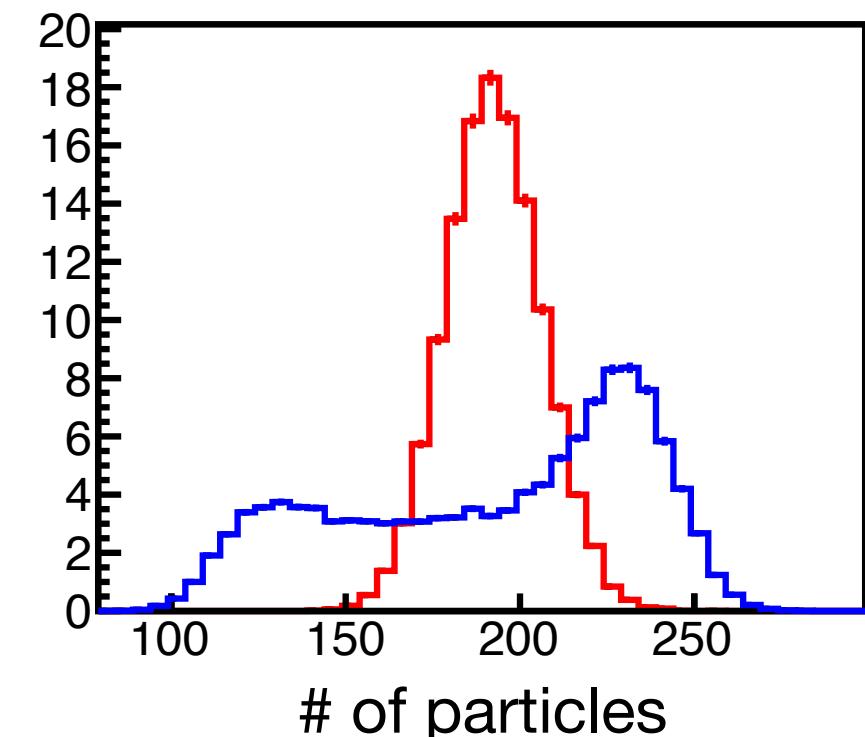
5.0 fm/c



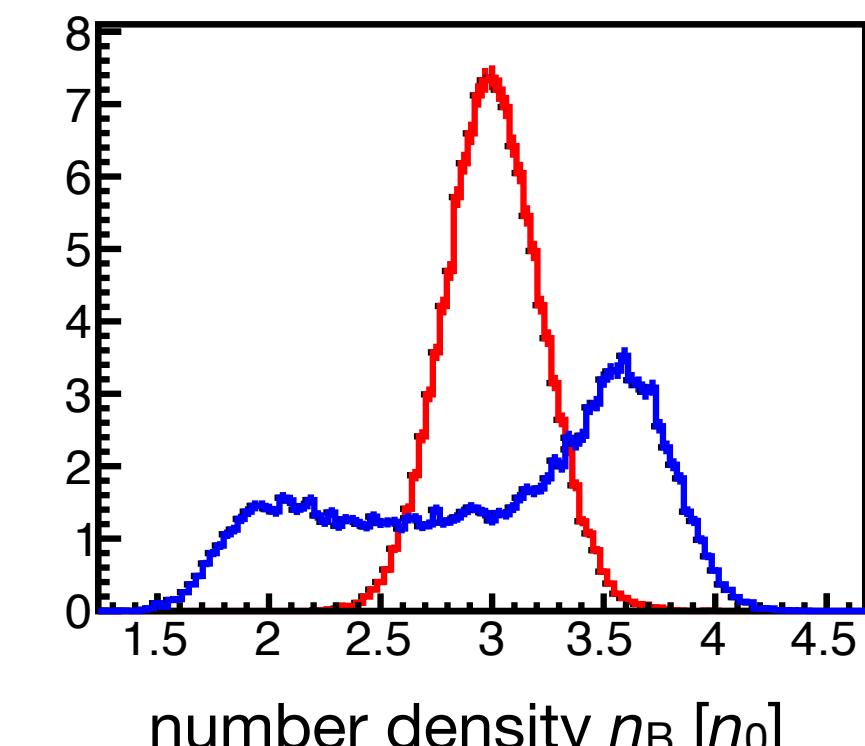
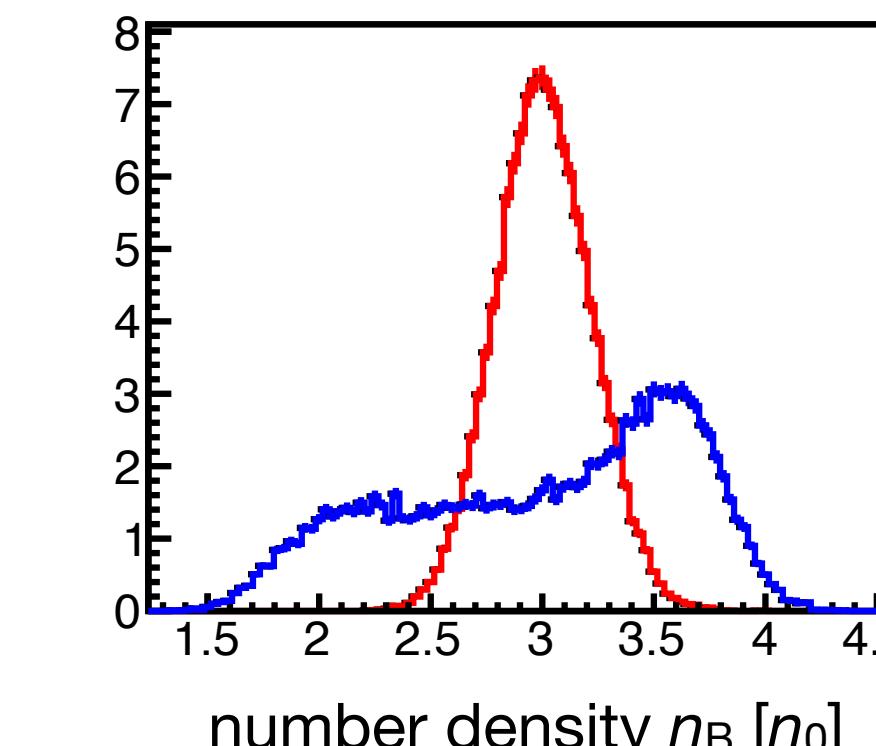
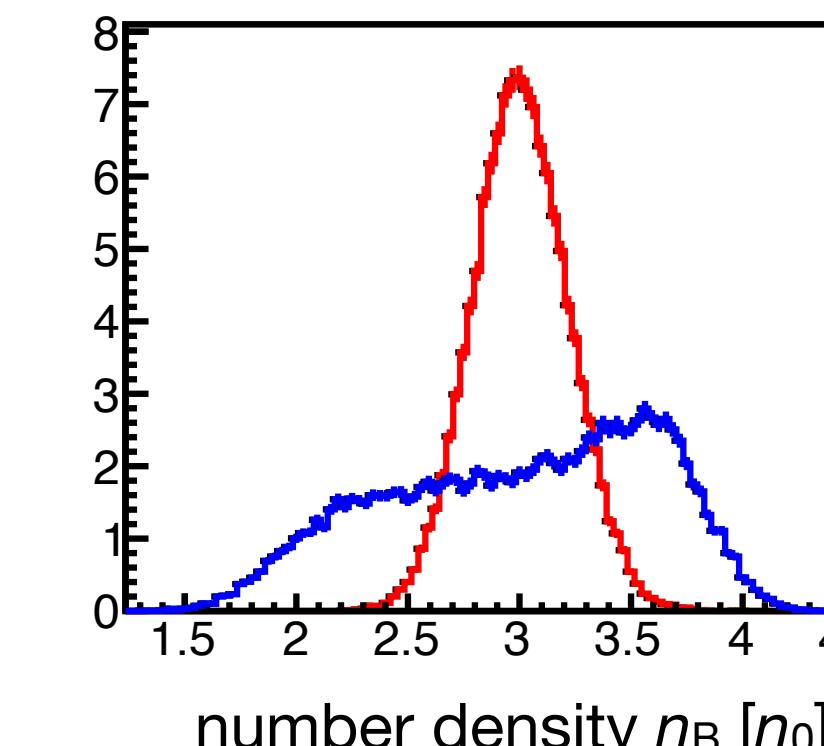
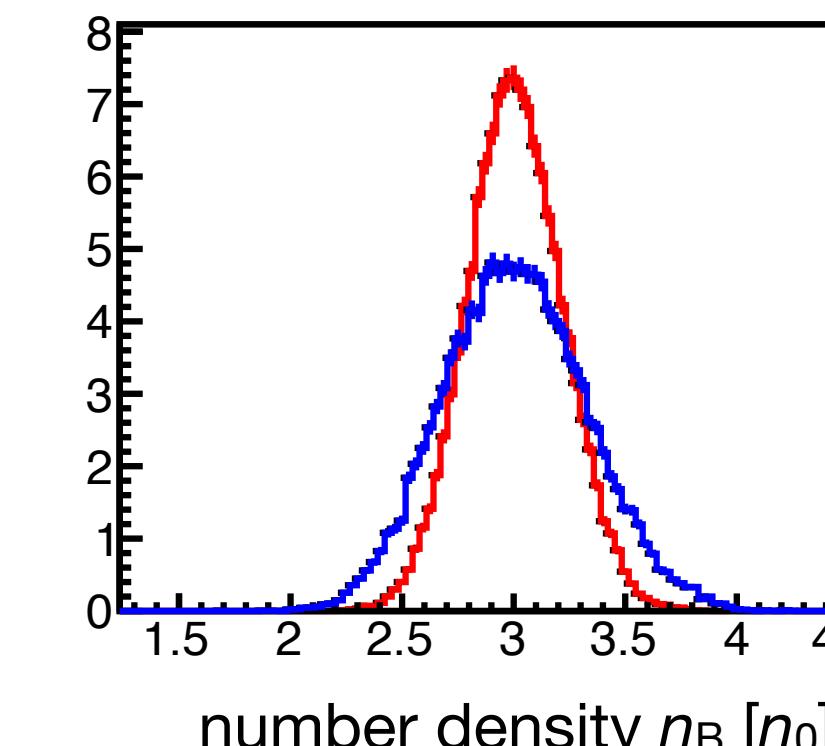
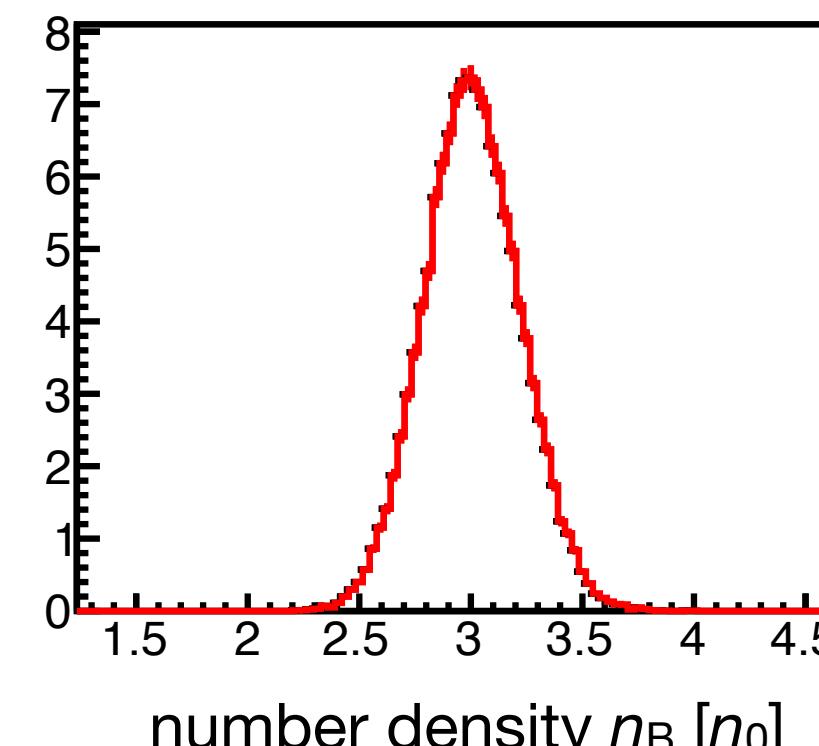
15.0 fm/c



20.0 fm/c



30.0 fm/c



SMASH results: periodic box, averaged over 250 events

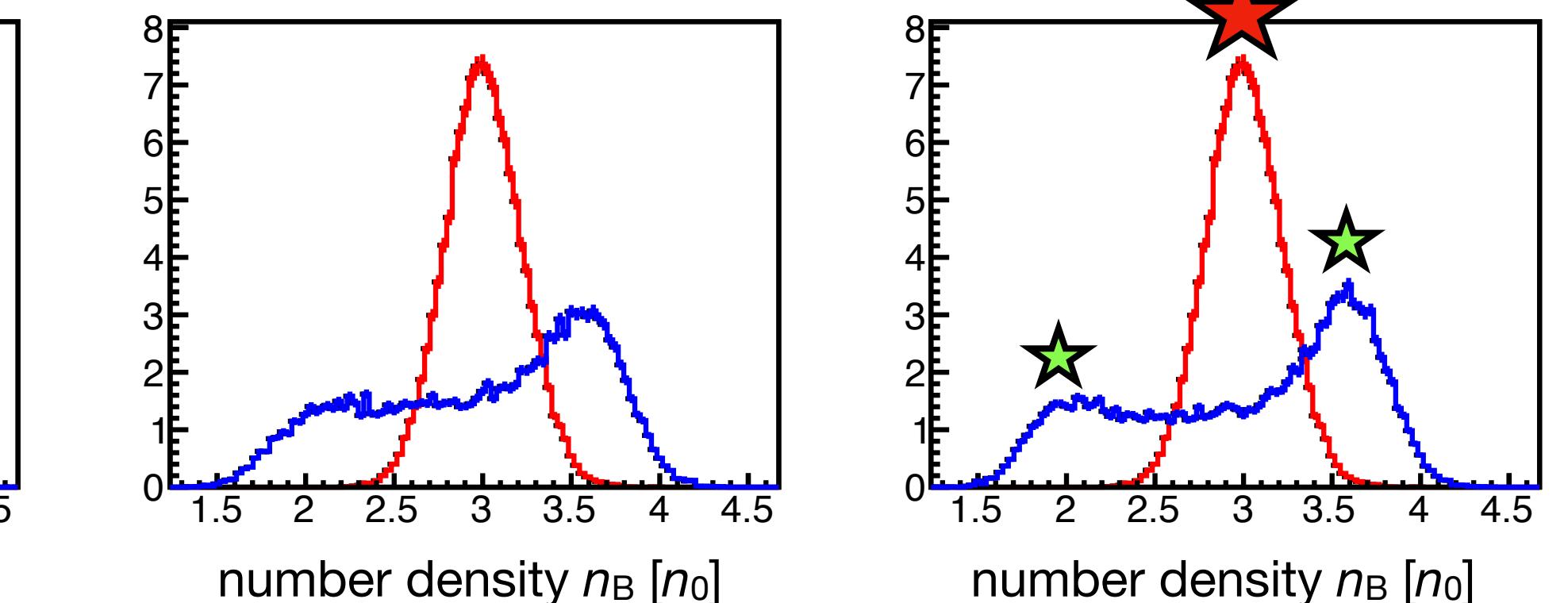
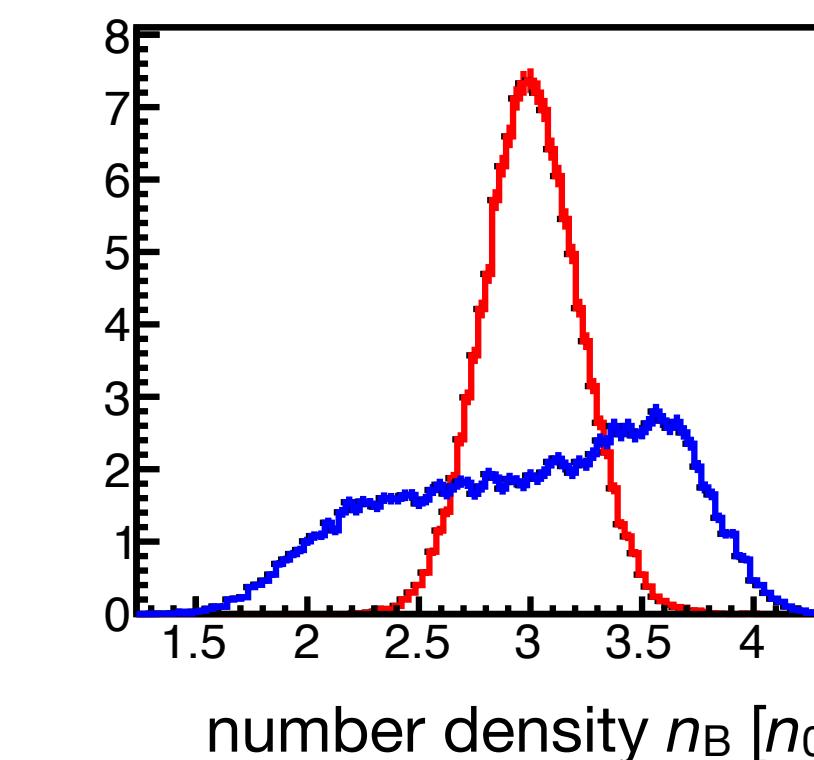
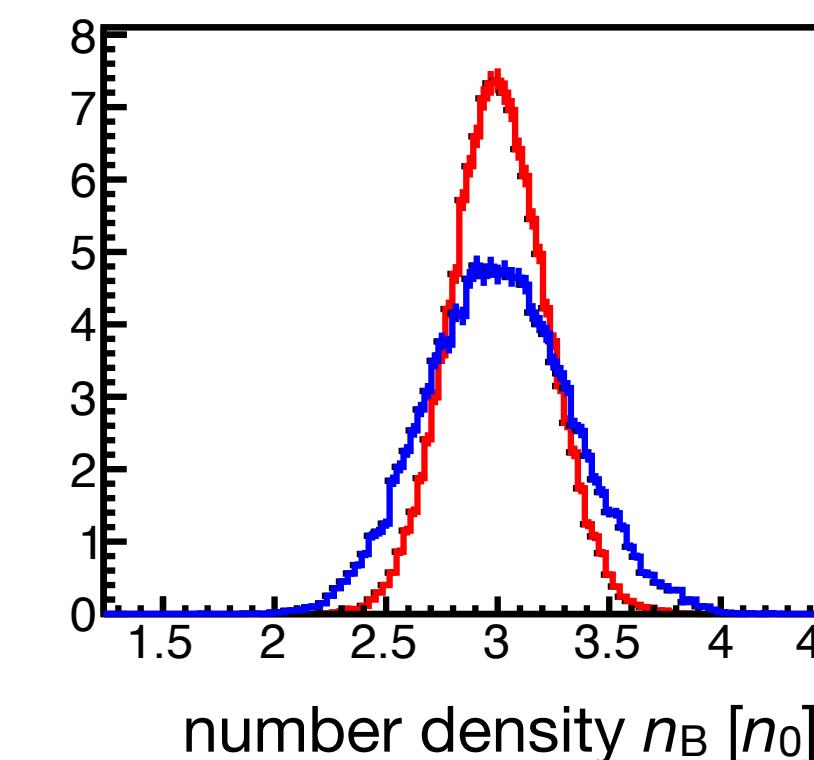
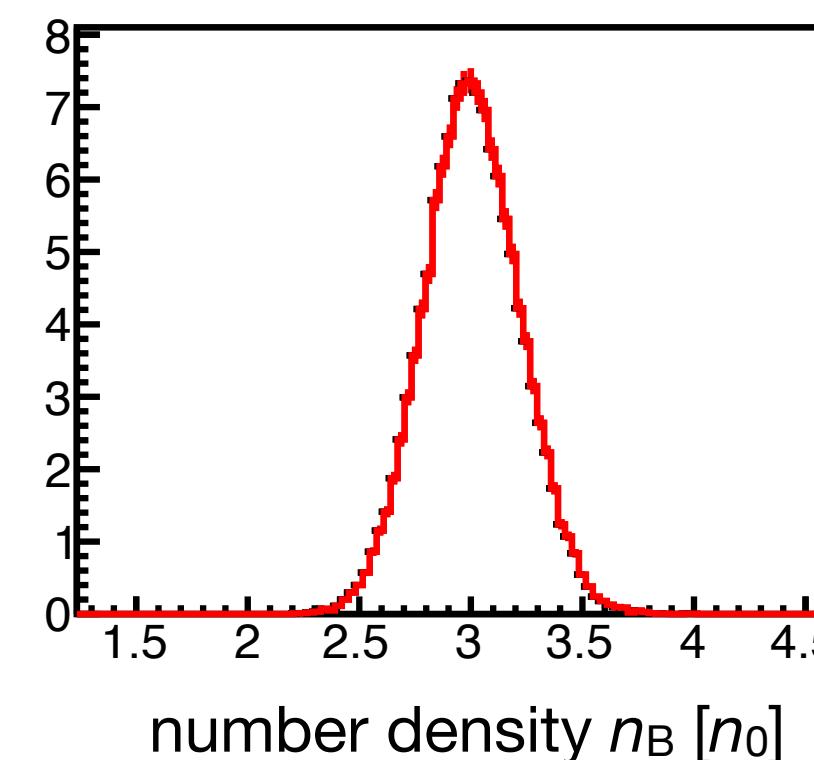
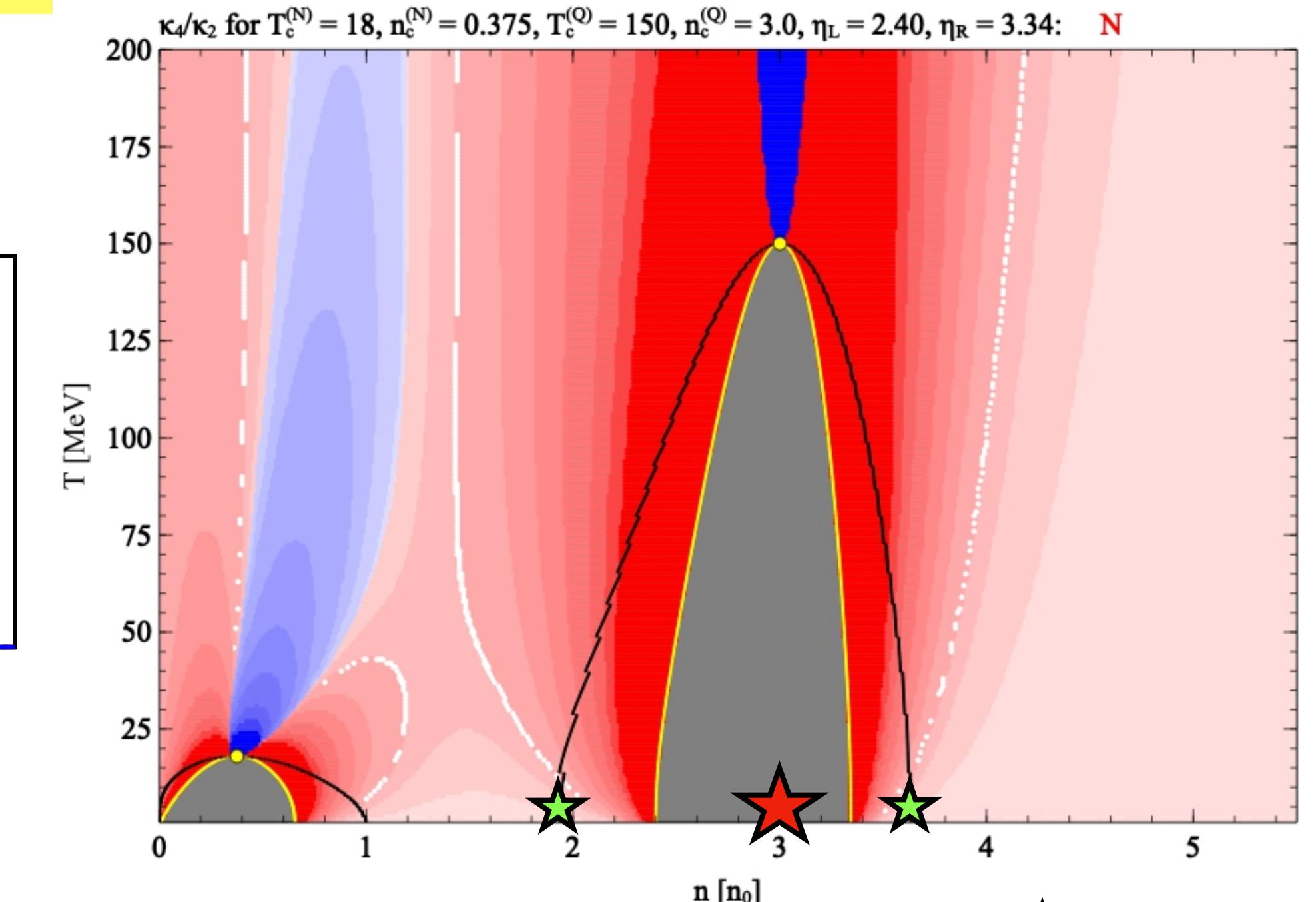
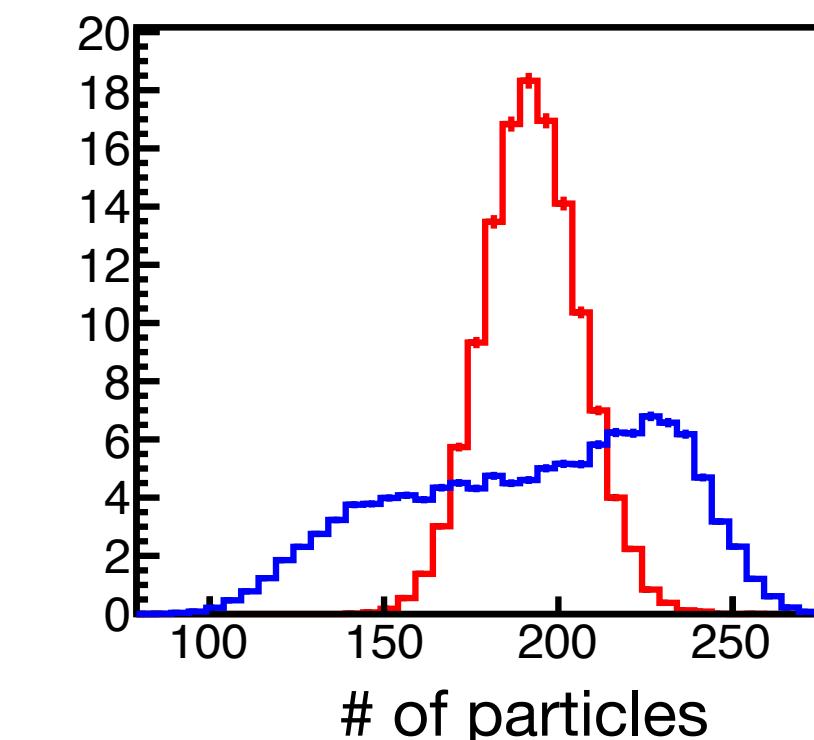
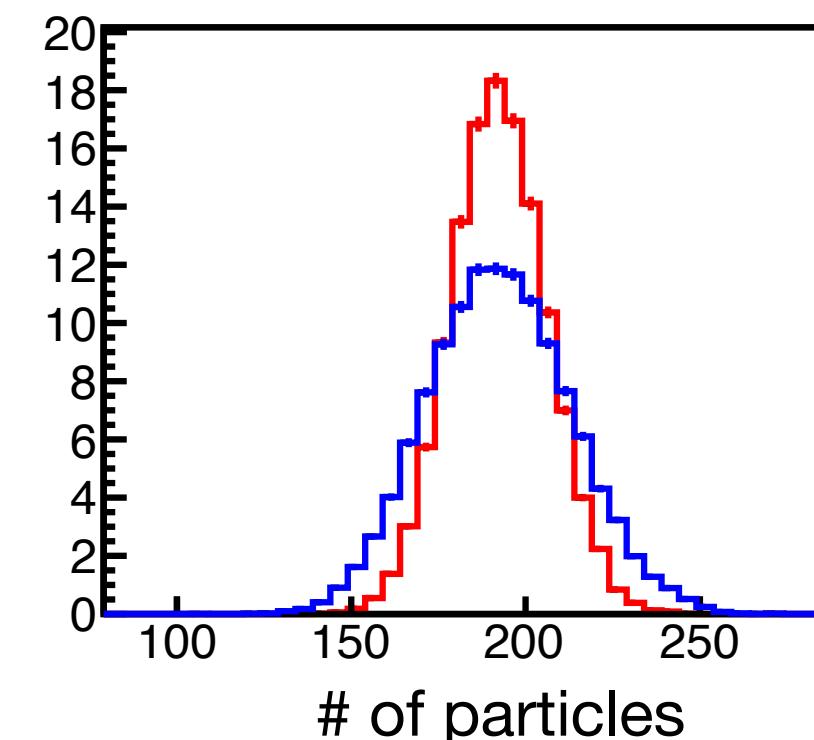
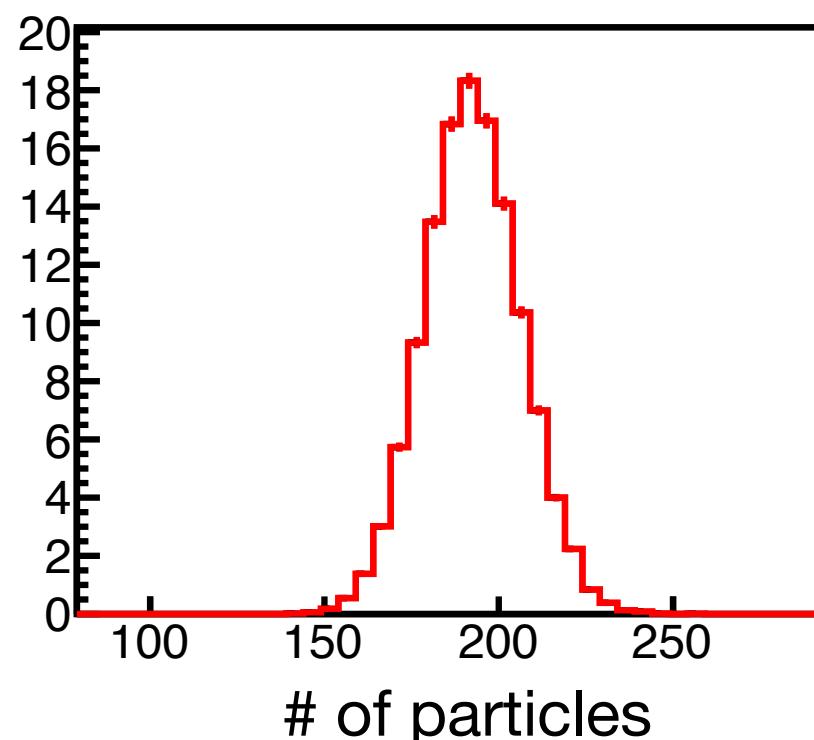
$L = 10 \text{ fm}$

$T = 1 \text{ MeV}$

$n_B = 3n_0$

Cell (bin) length $L = 2 \text{ fm}$

The distribution becomes bimodal as the system separates!



SMASH results: periodic box, averaged over 250 events

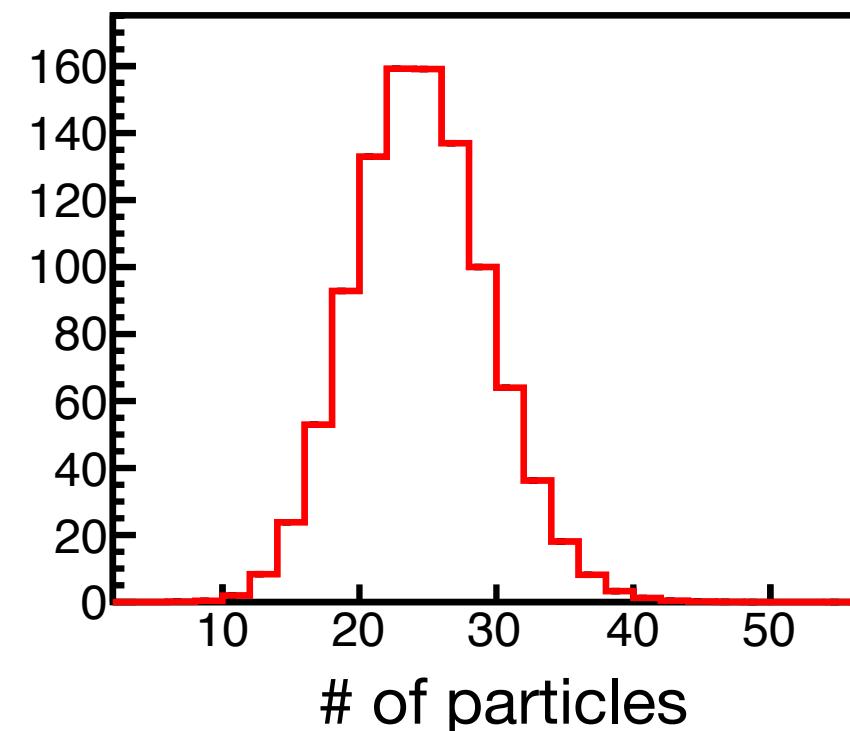
$L = 10 \text{ fm}$

$T = 1 \text{ MeV}$

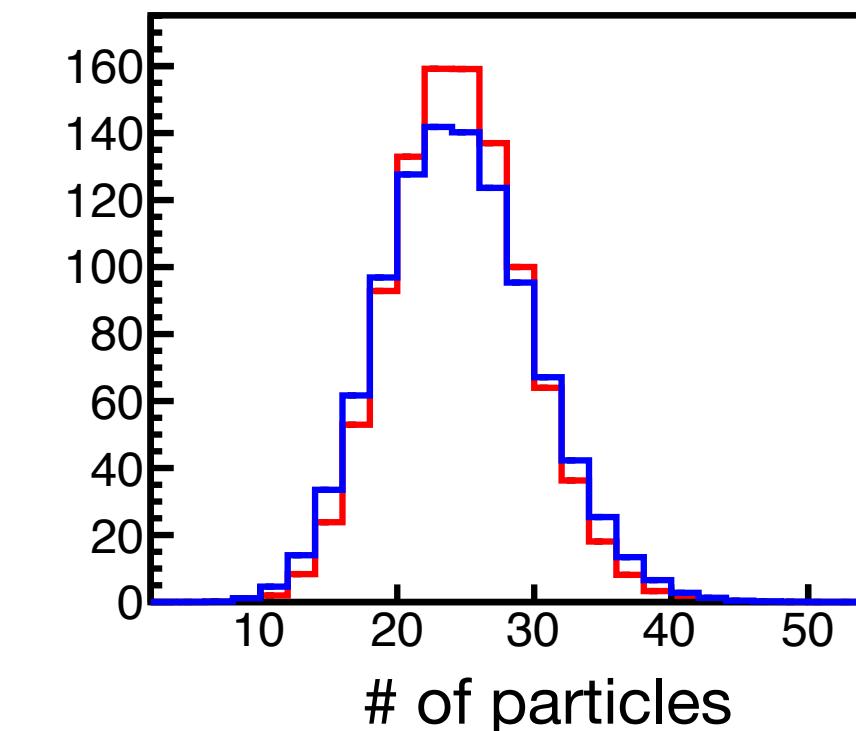
$n_B = 3n_0$

Cell (bin) length $L = 1 \text{ fm}$

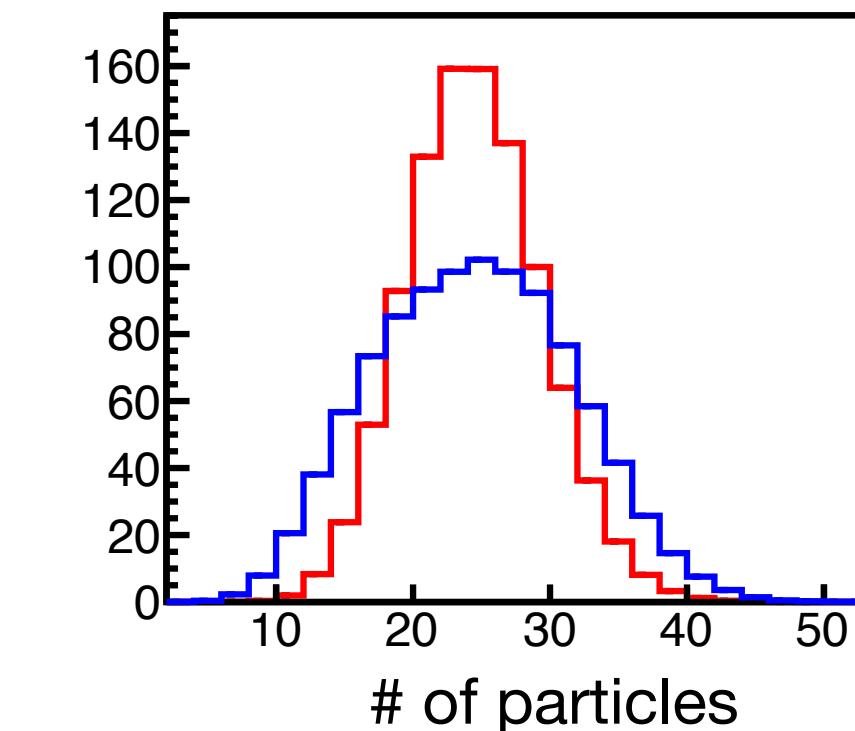
The degree of bimodality is sensitive to binning



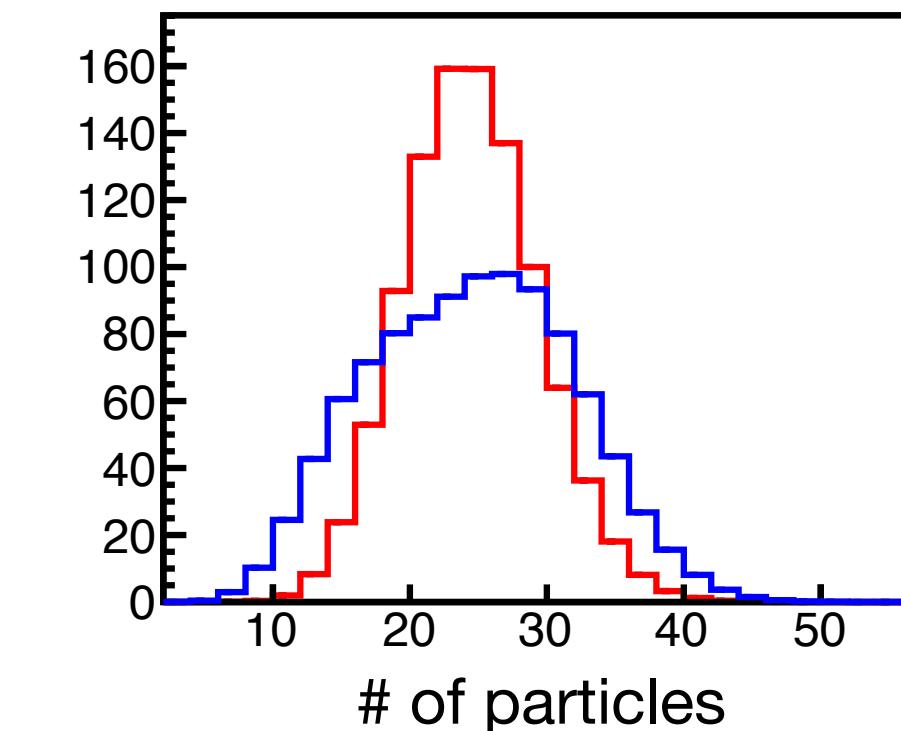
0.0 fm/c



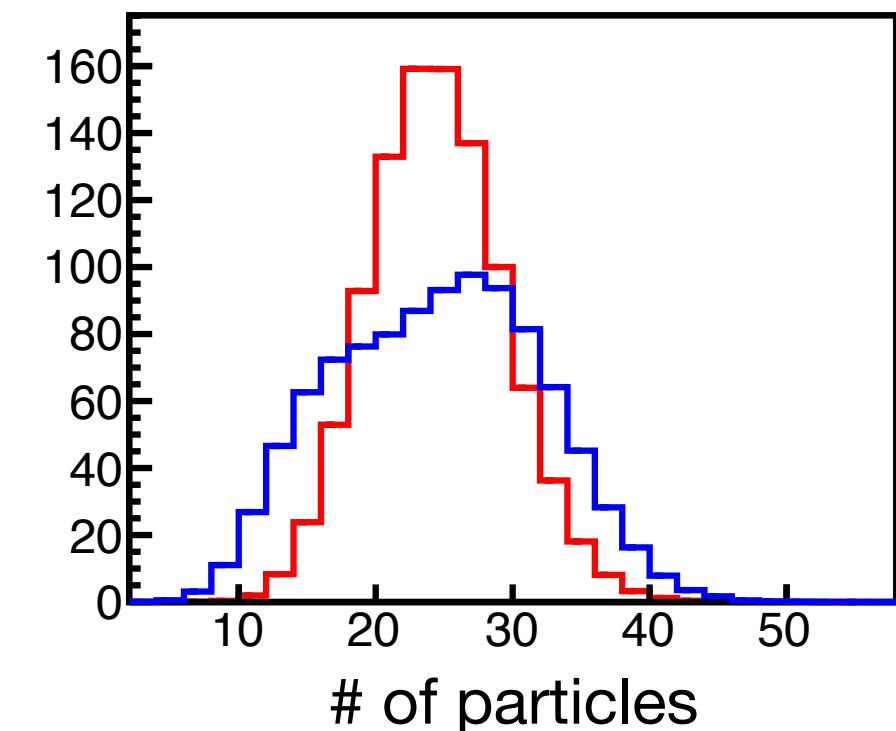
5.0 fm/c



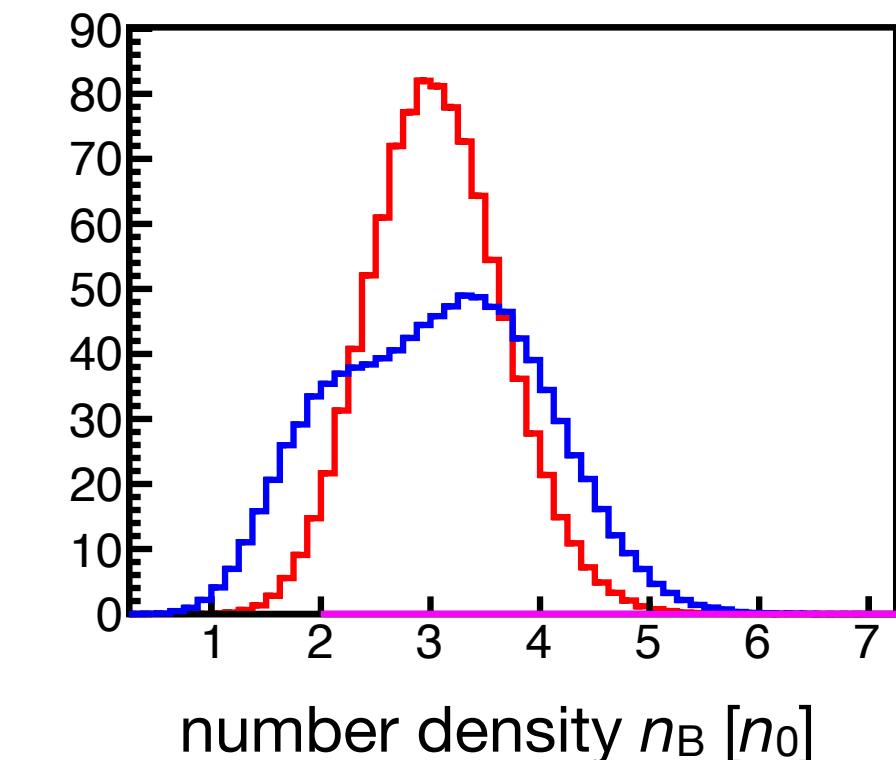
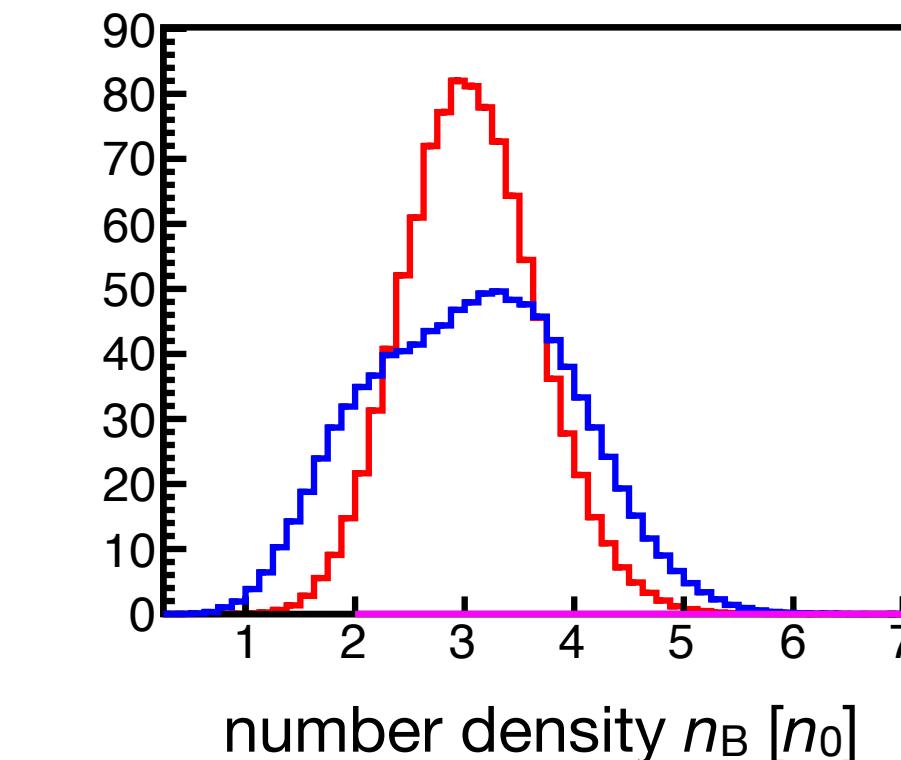
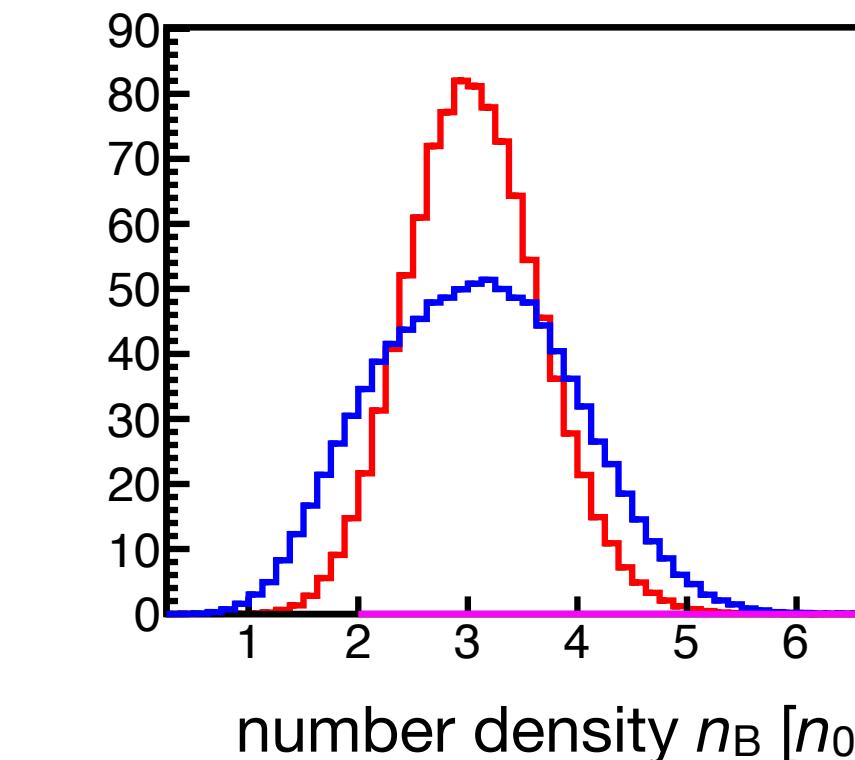
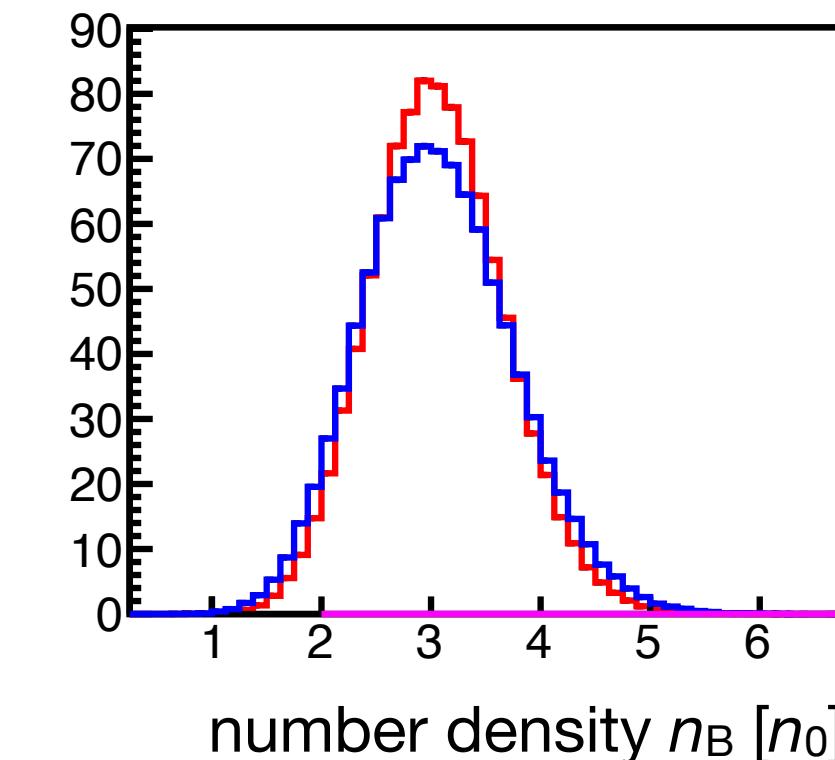
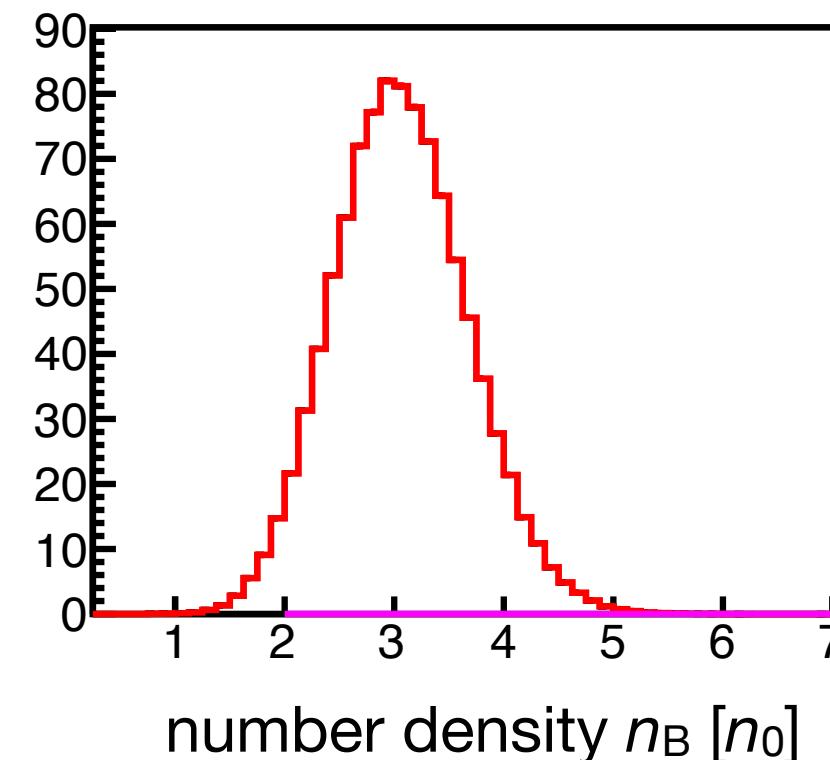
15.0 fm/c



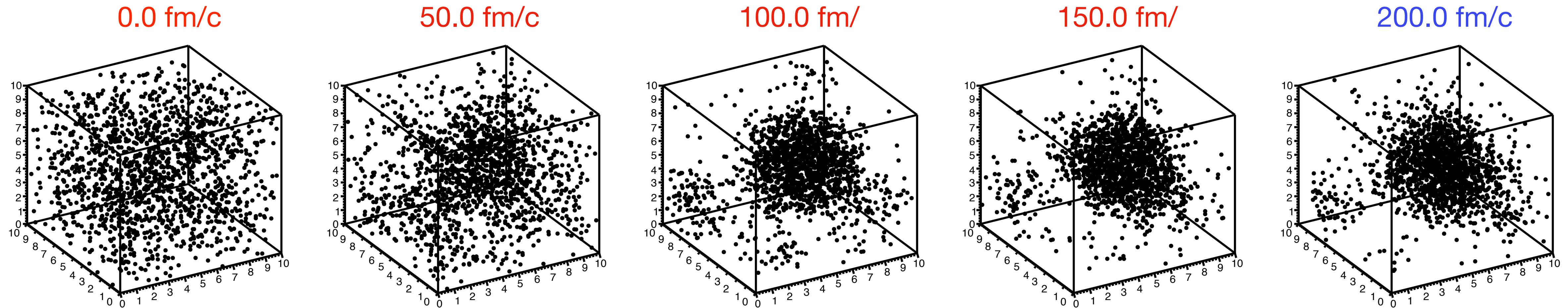
20.0 fm/c



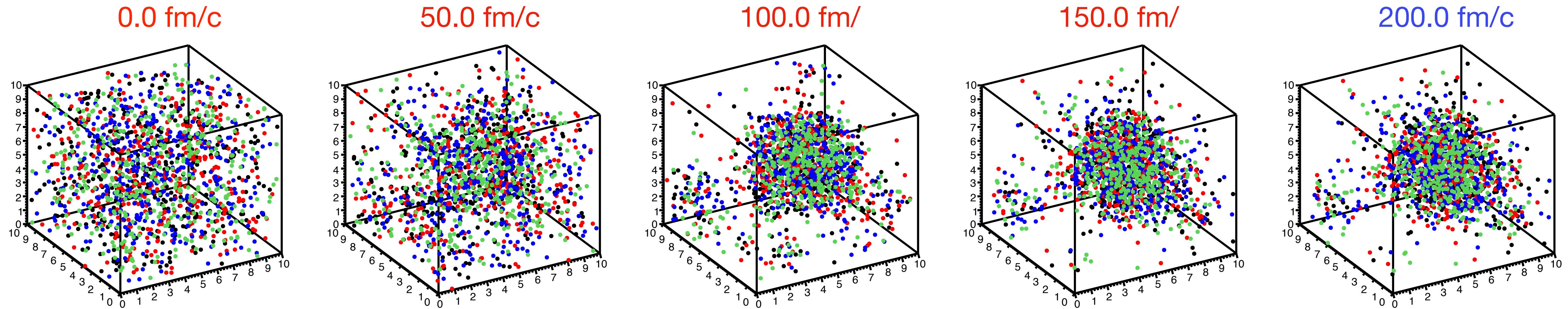
30.0 fm/c



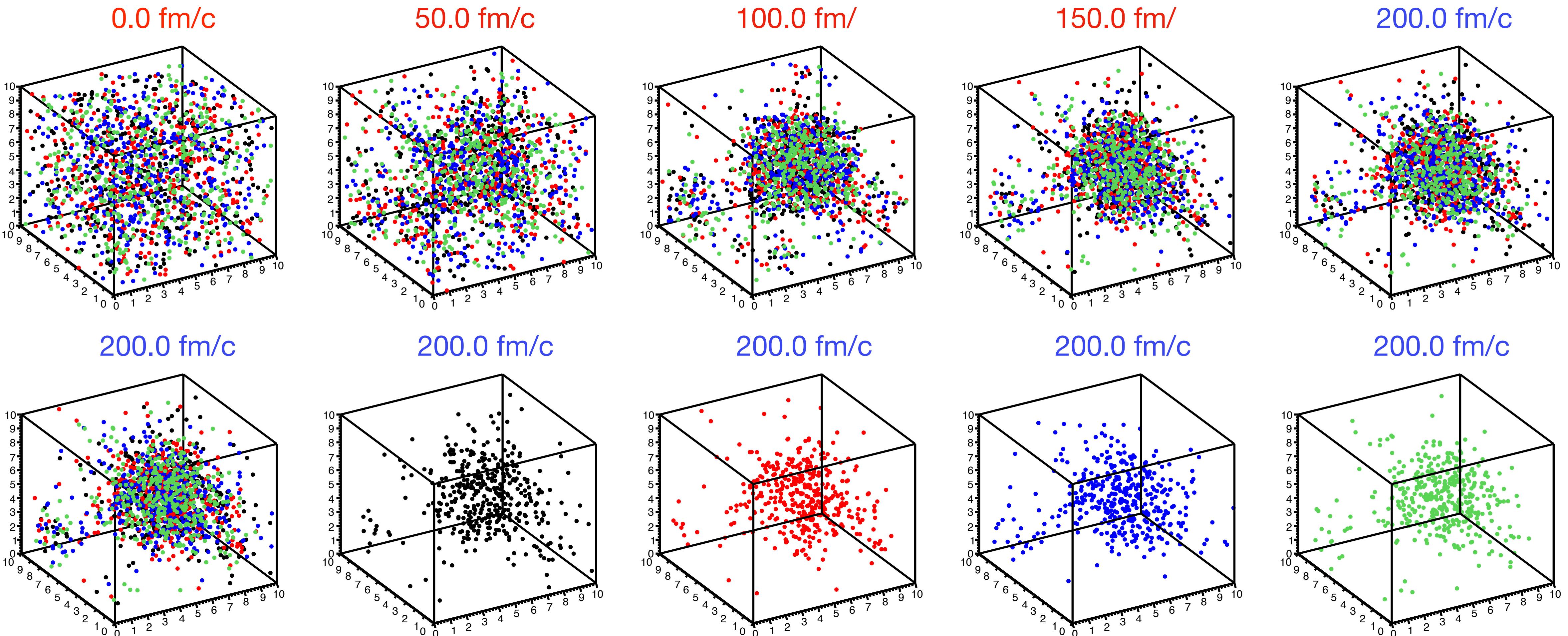
SMASH results: periodic box with “parallel ensembles”



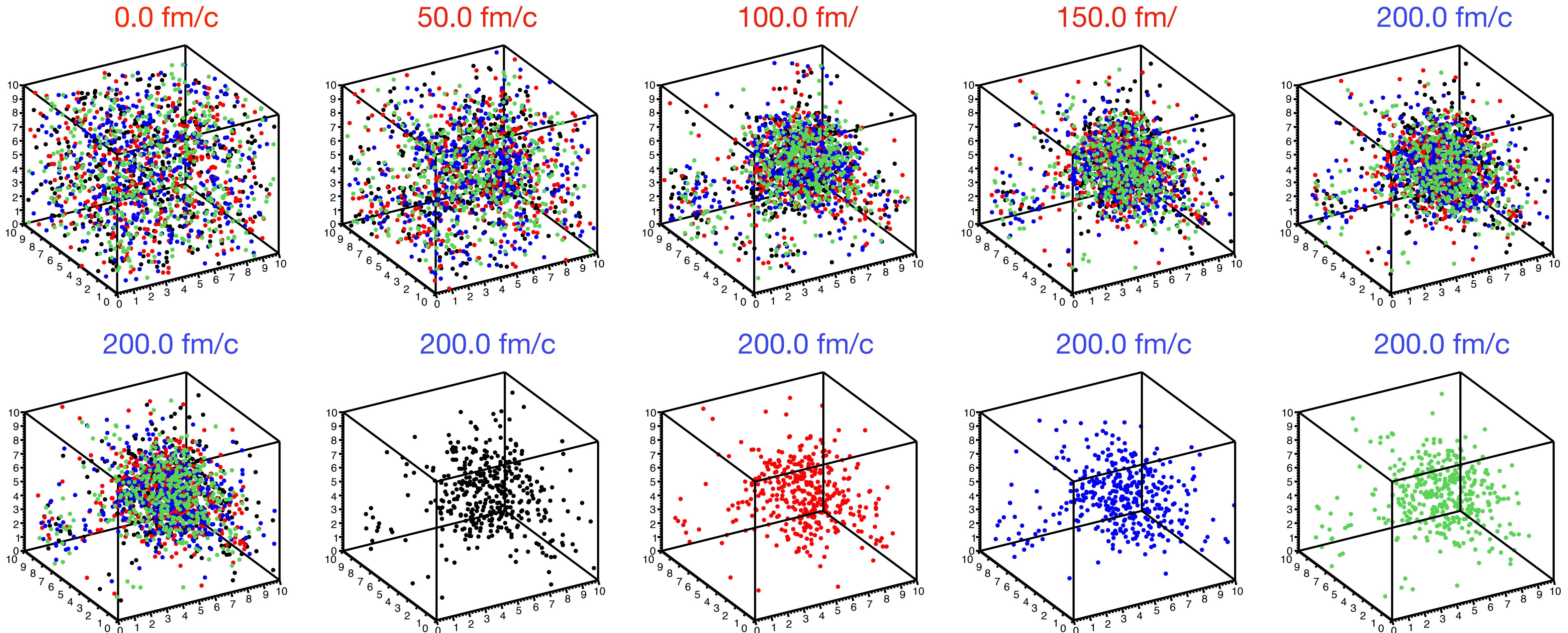
SMASH results: periodic box with “parallel ensembles”



SMASH results: periodic box with “parallel ensembles”



SMASH results: periodic box with “parallel ensembles”



We introduce “parallel ensembles” *post factum*:
We take # test particles $N_T = \#$ “parallel ensembles” →
and create “parallel boxes” at the analysis stage

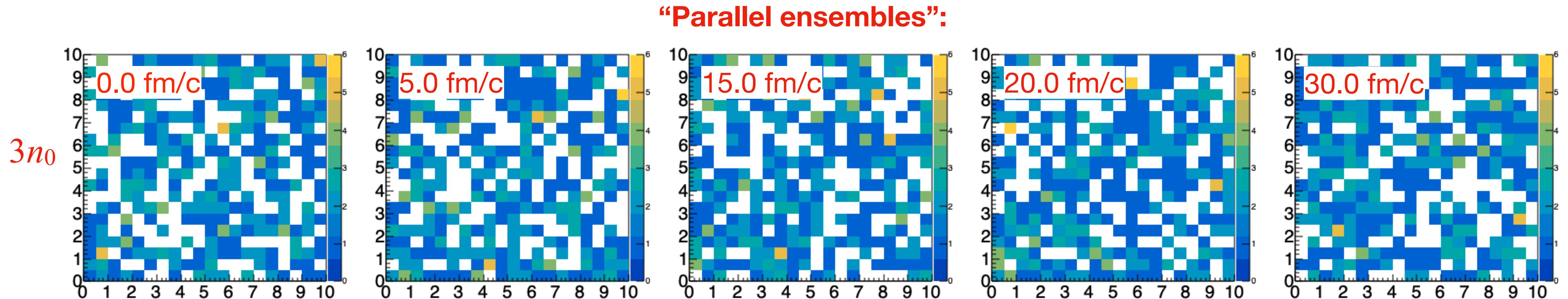
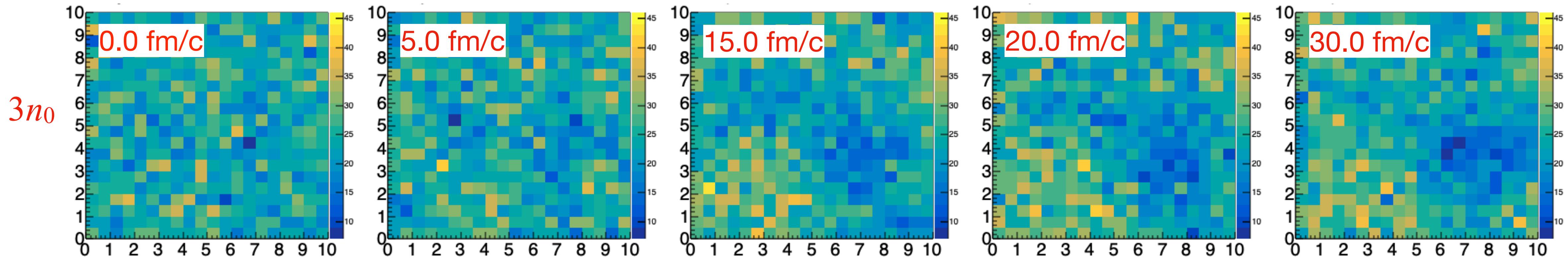
At $n_B = 3n_0$, this means 480 baryons
instead of 24,000 test particles ($N_T = 50$).
This is more like experiment!

SMASH results: periodic box with “parallel ensembles”

$L = 10 \text{ fm}$ $T = 1 \text{ MeV}$

particle number projection onto the xy -plane

$\mathbf{NT} = 20$

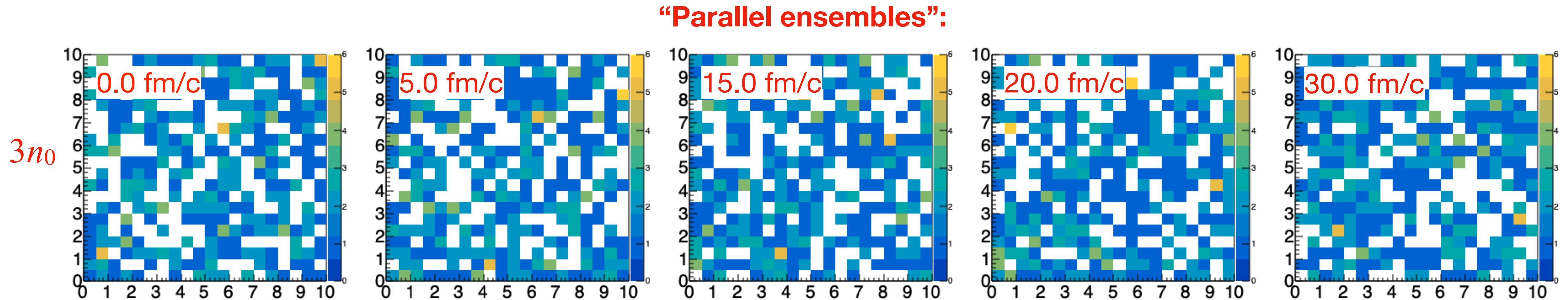
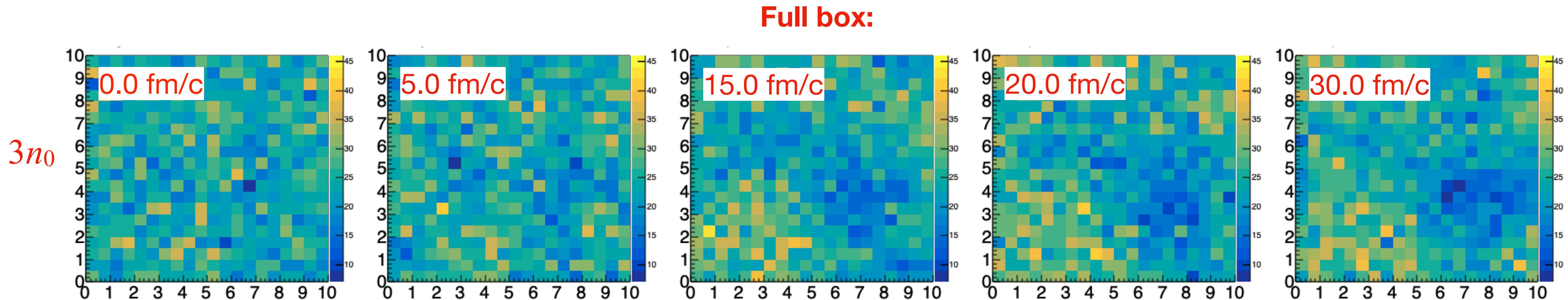


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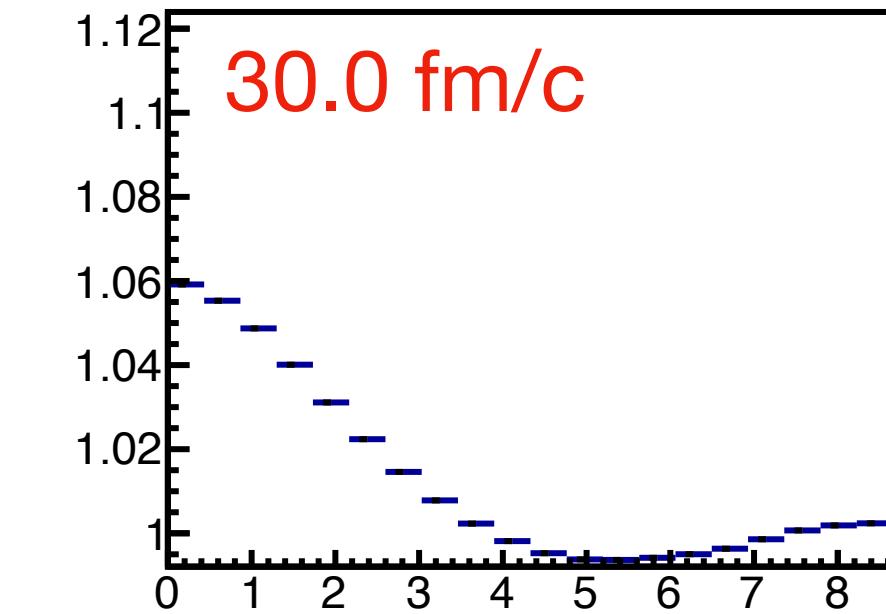
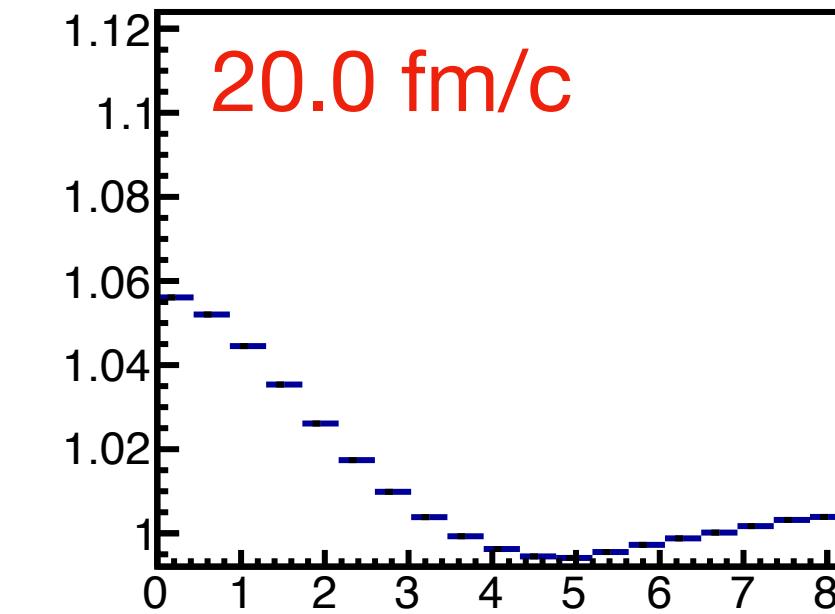
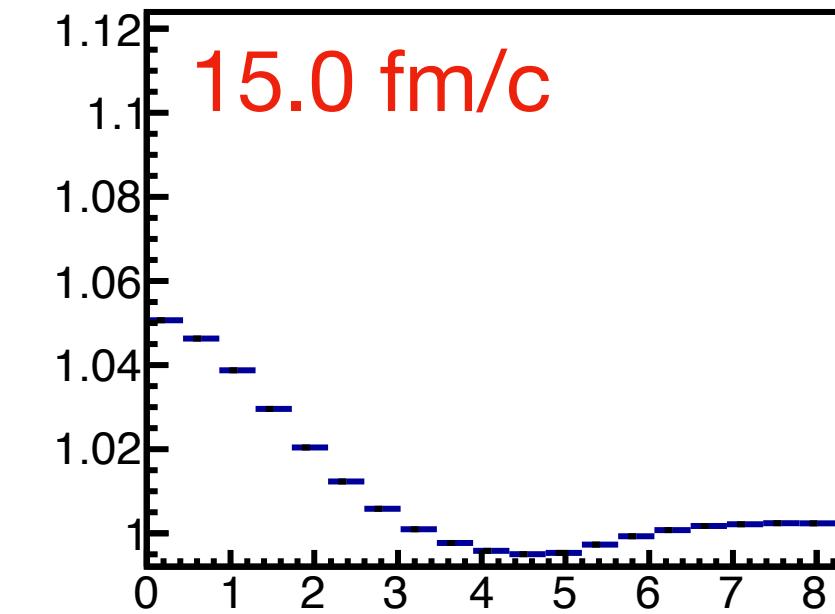
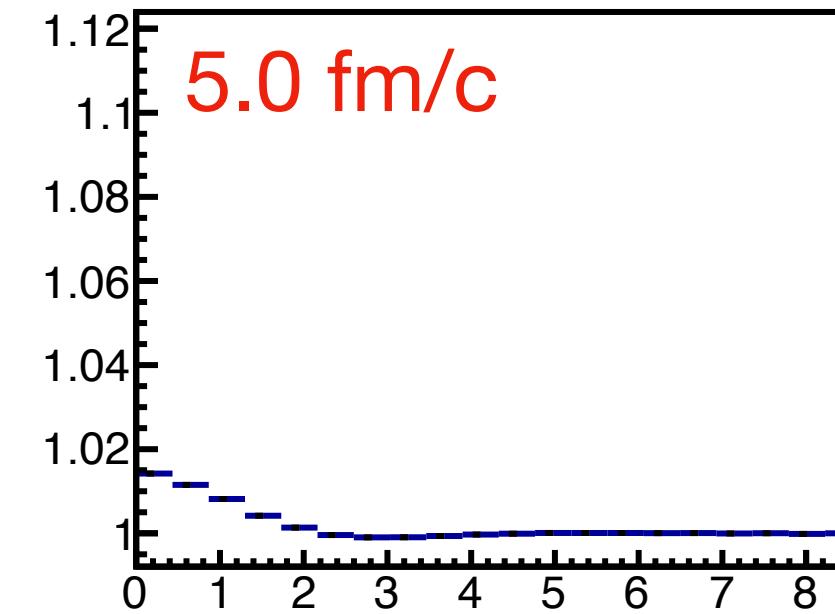
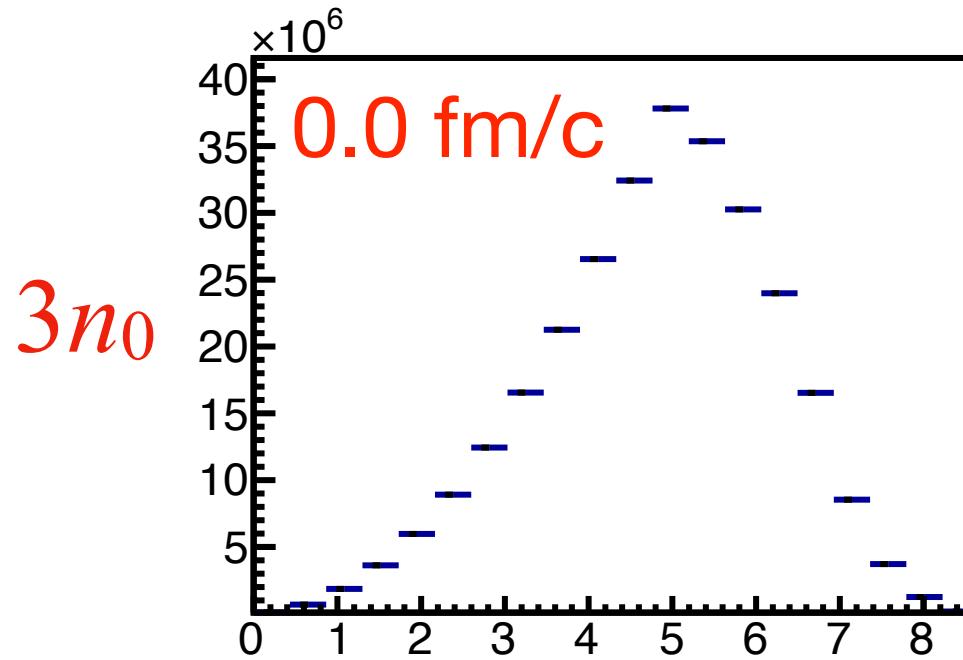
Is information about phase separation preserved in “parallel ensembles”?

SMASH results: periodic box with “parallel ensembles”

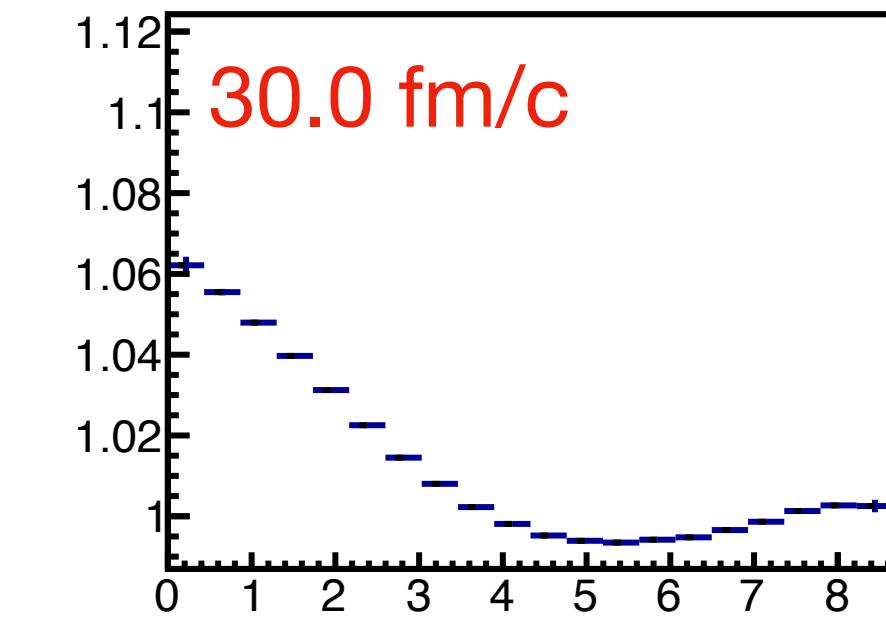
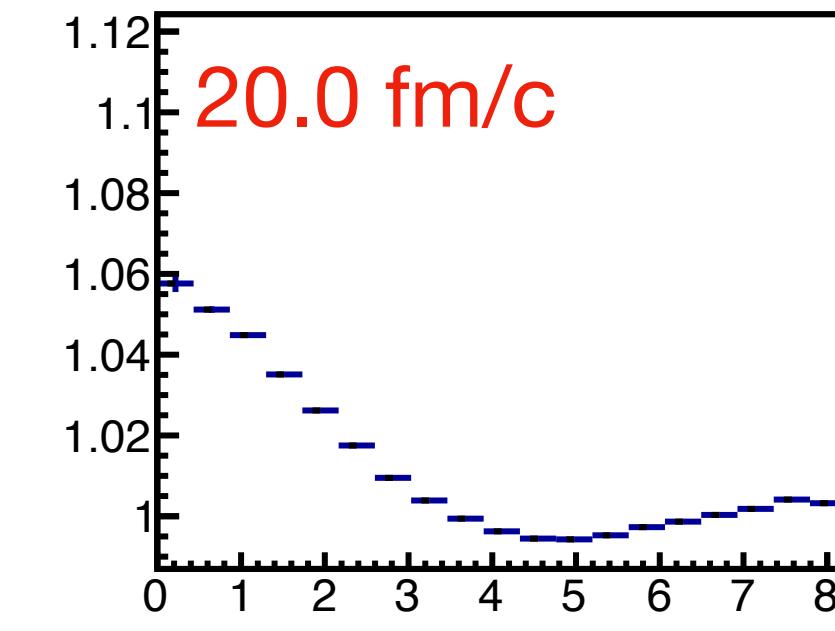
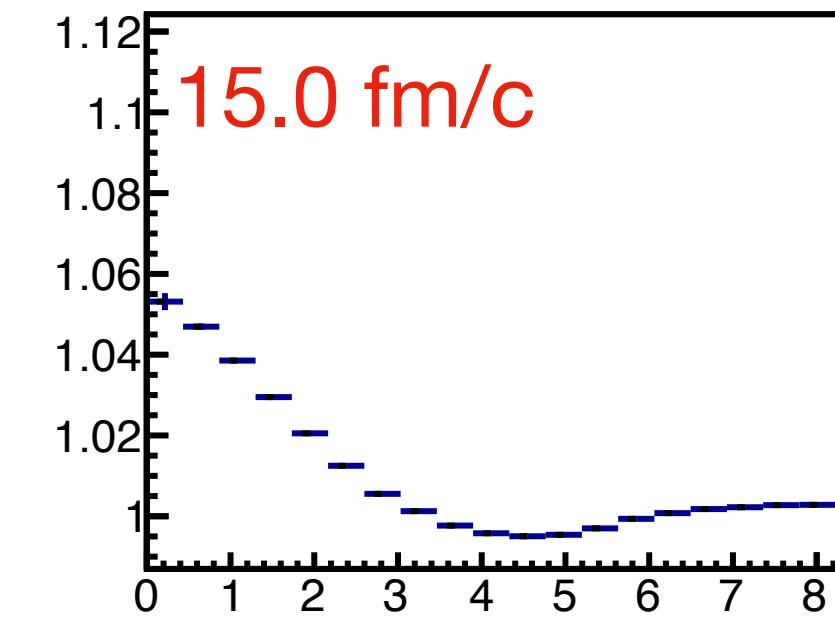
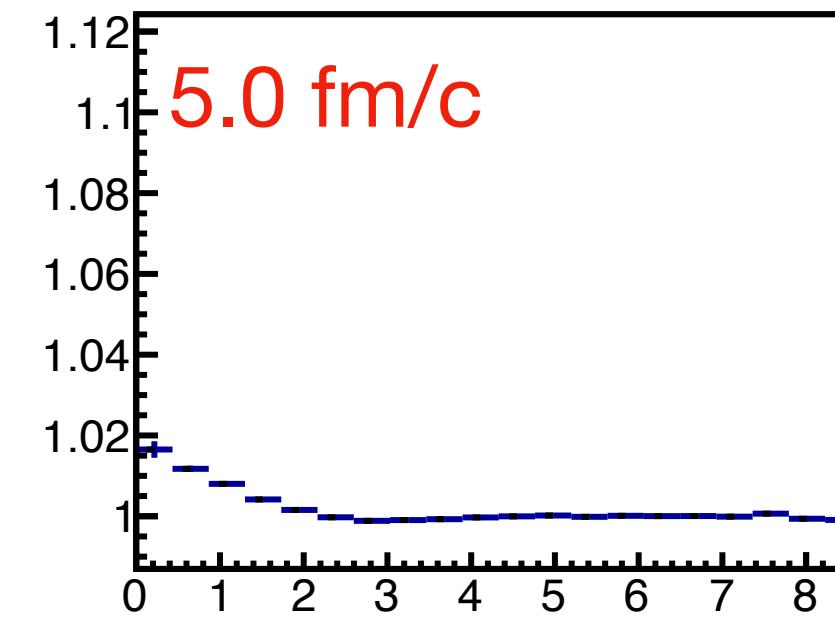
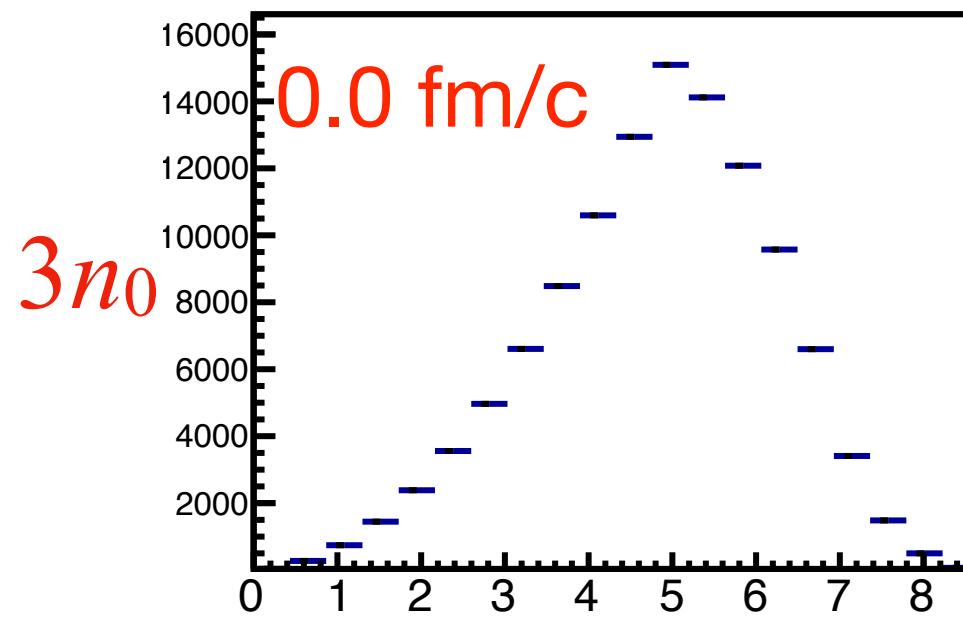
$L = 10 \text{ fm}$ $T = 1 \text{ MeV}$

pair separation distribution (scaled for $t > 0$)

$N_T = N_{\text{“parallel”}} = 50$



Full box:



$\sqrt{(\vec{r}_i - \vec{r}_j)^2} [\text{fm}]$

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Is information about phase separation preserved in “parallel ensembles”? Yes!

SMASH results: periodic box with “parallel ensembles”

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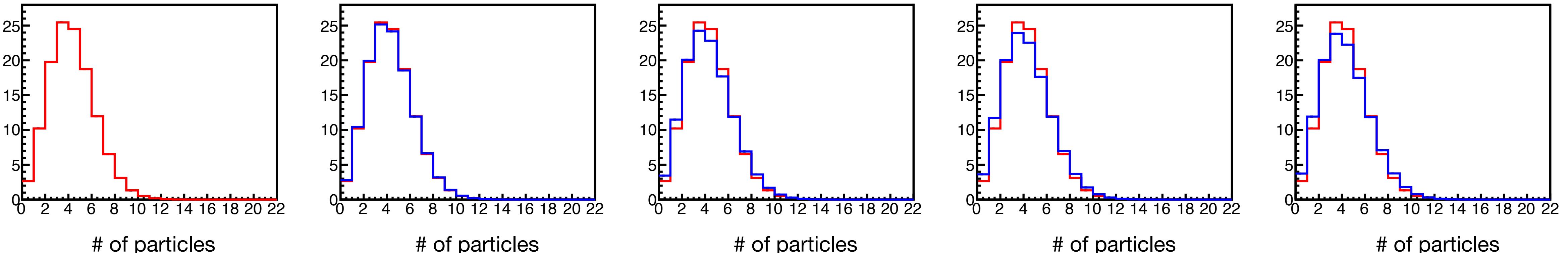
$n_B = 3n_0$

Cell (bin) length $L = 2 \text{ fm}$

$N_T = N_{\text{“parallel”}} = 50$

But bimodality is gone!

Same physics, same phase transition, but small particle numbers affect what the distribution looks like



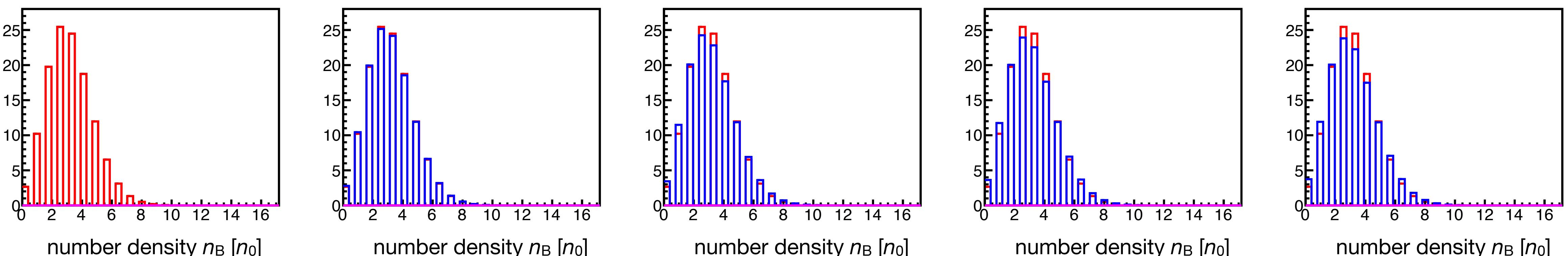
0.0 fm/c

5.0 fm/c

15.0 fm/c

20.0 fm/c

30.0 fm/c



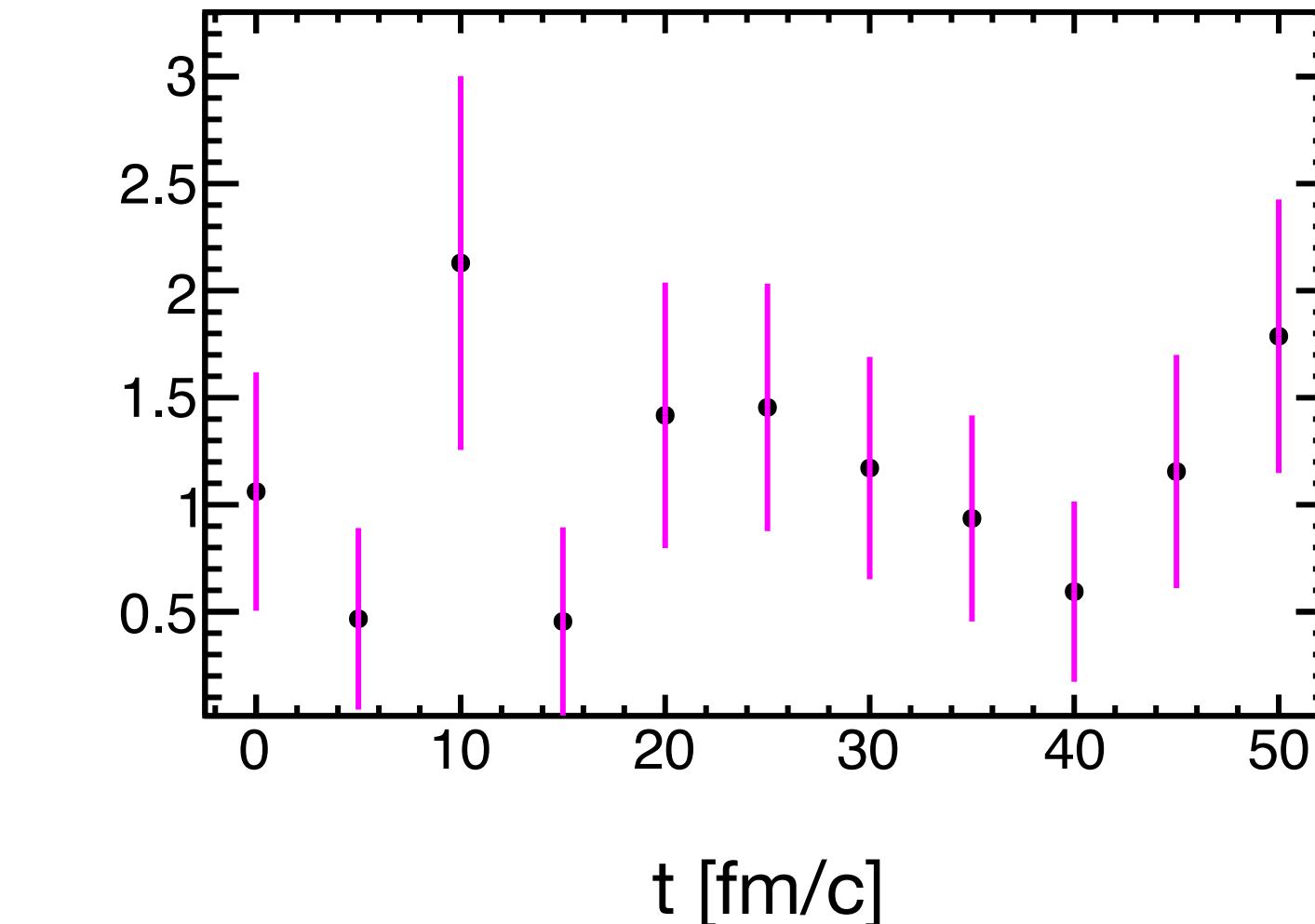
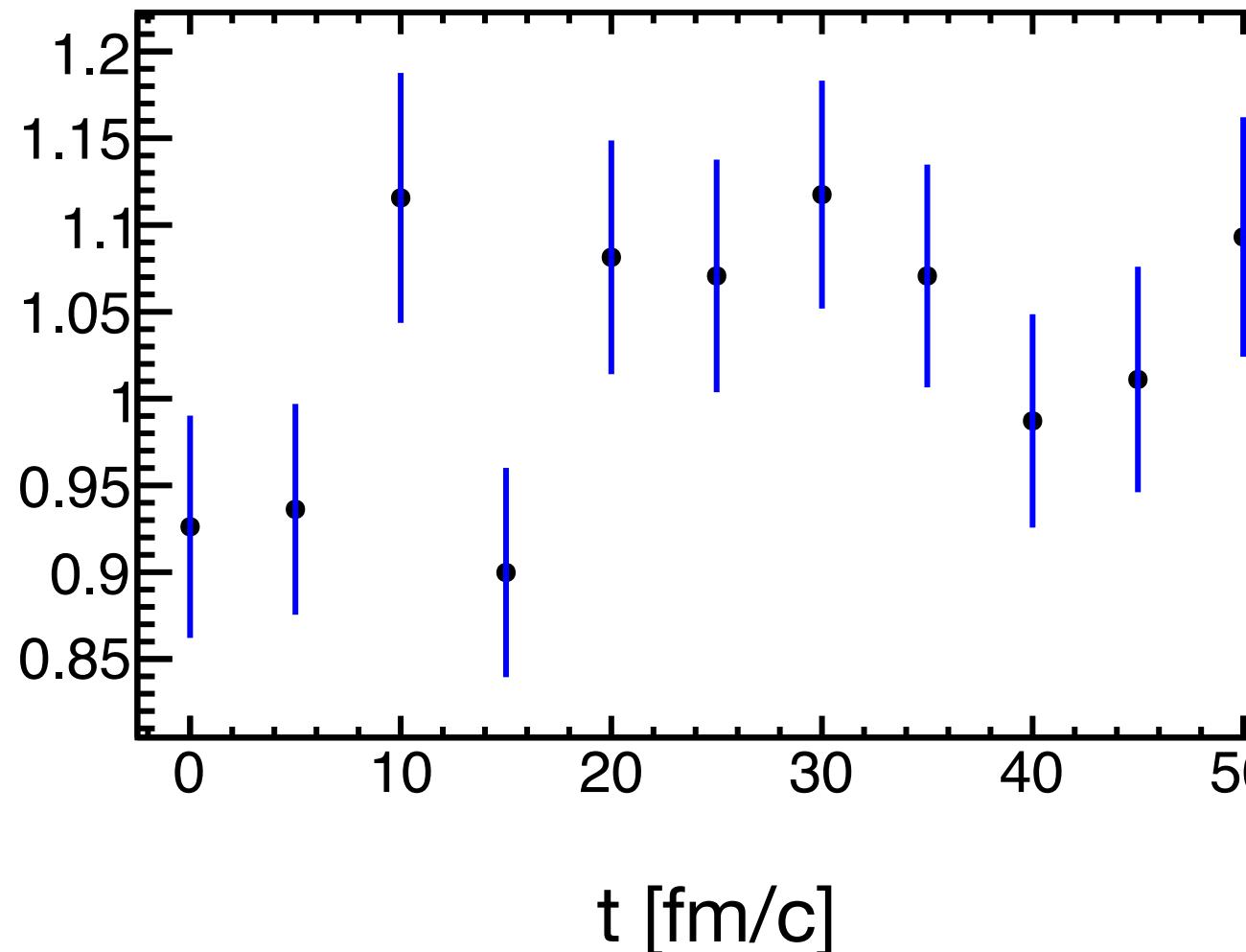
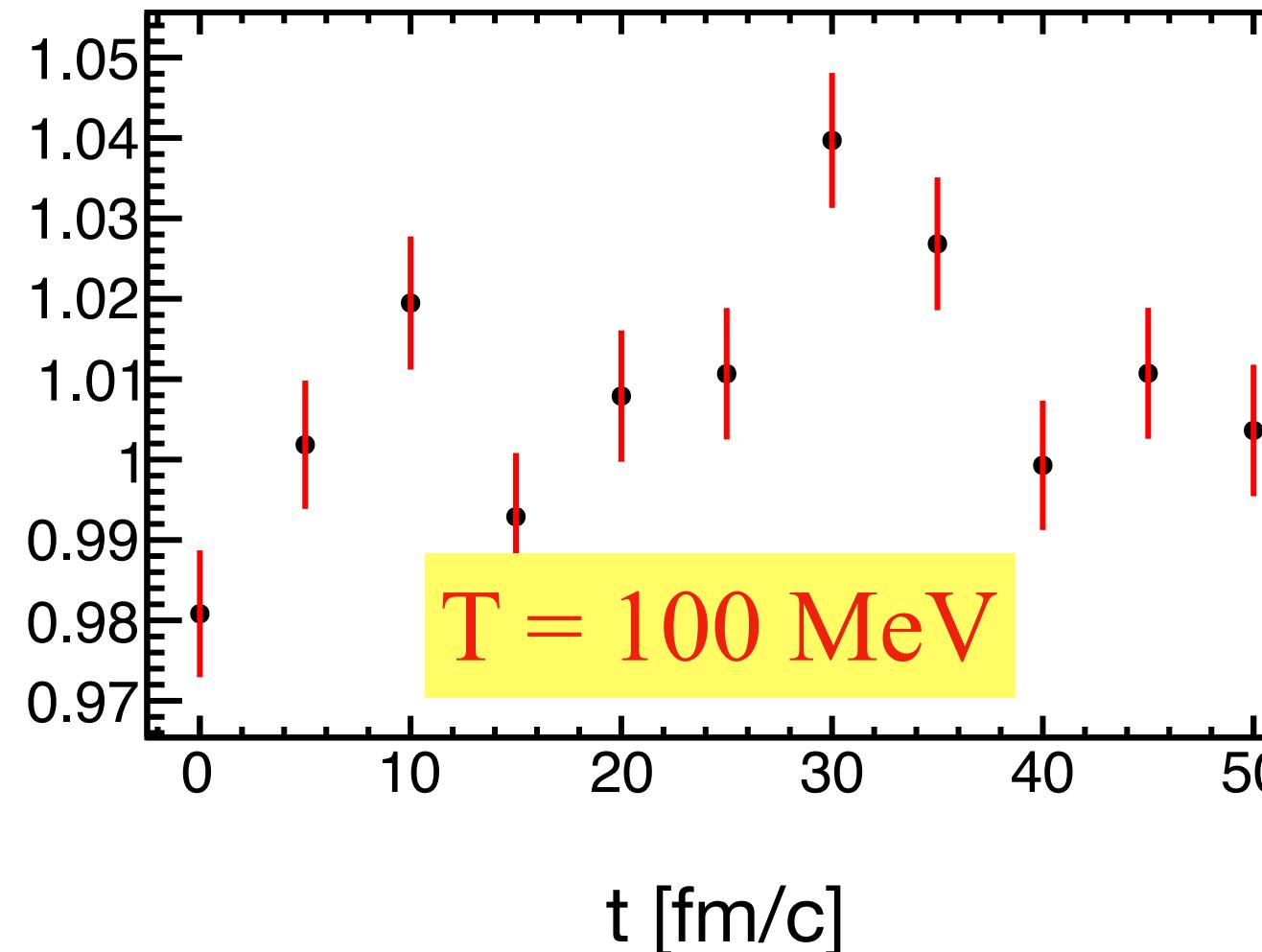
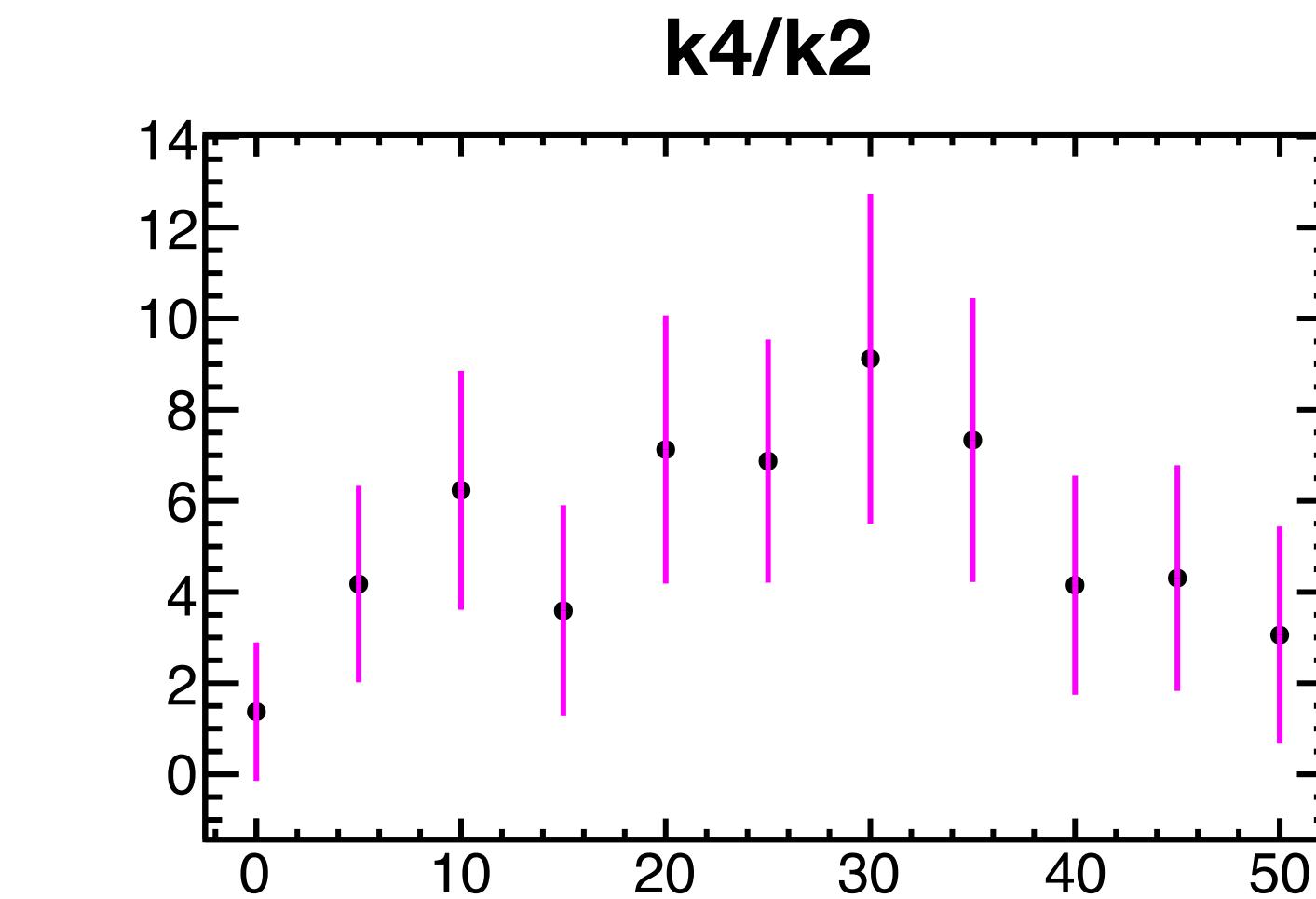
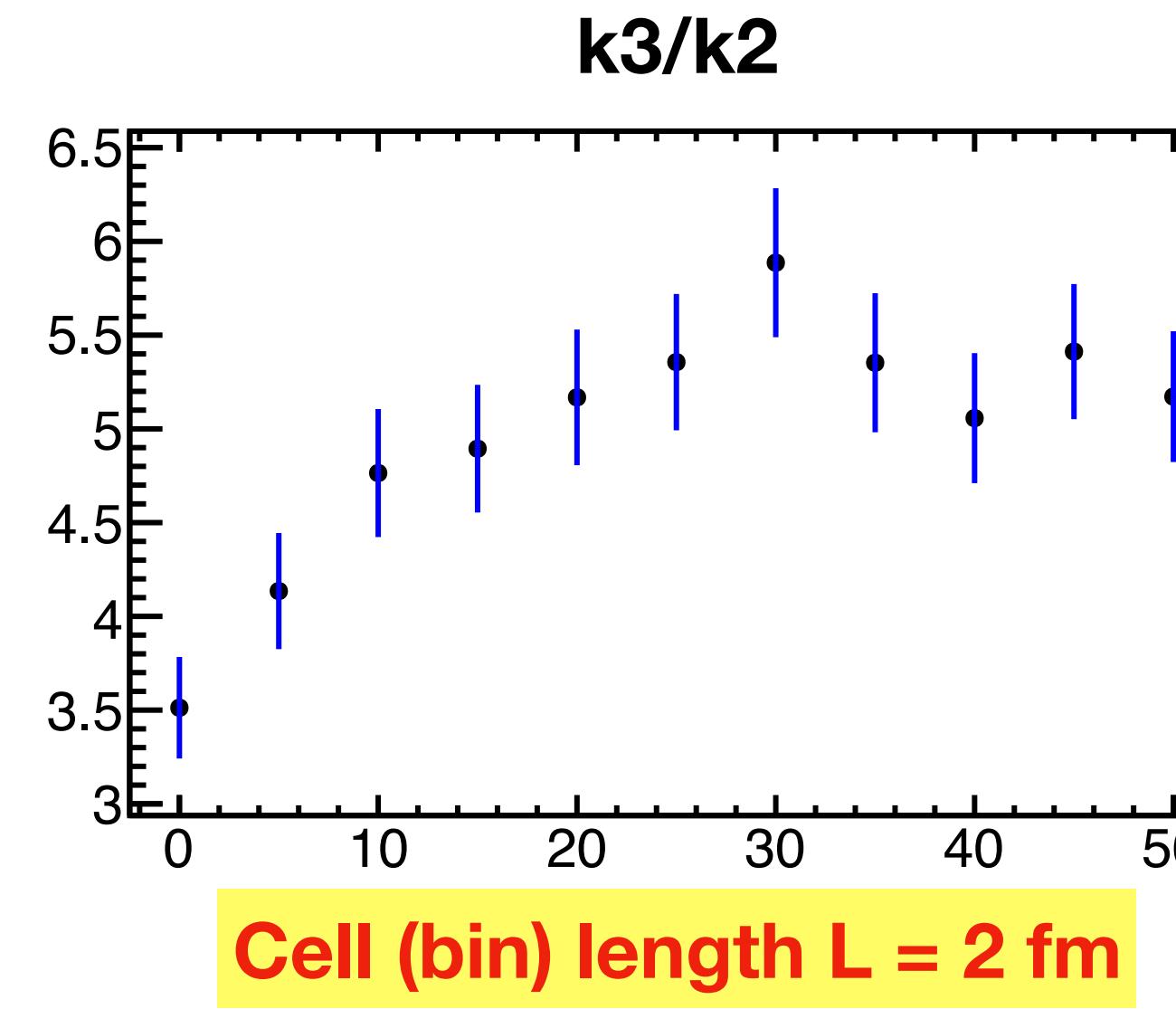
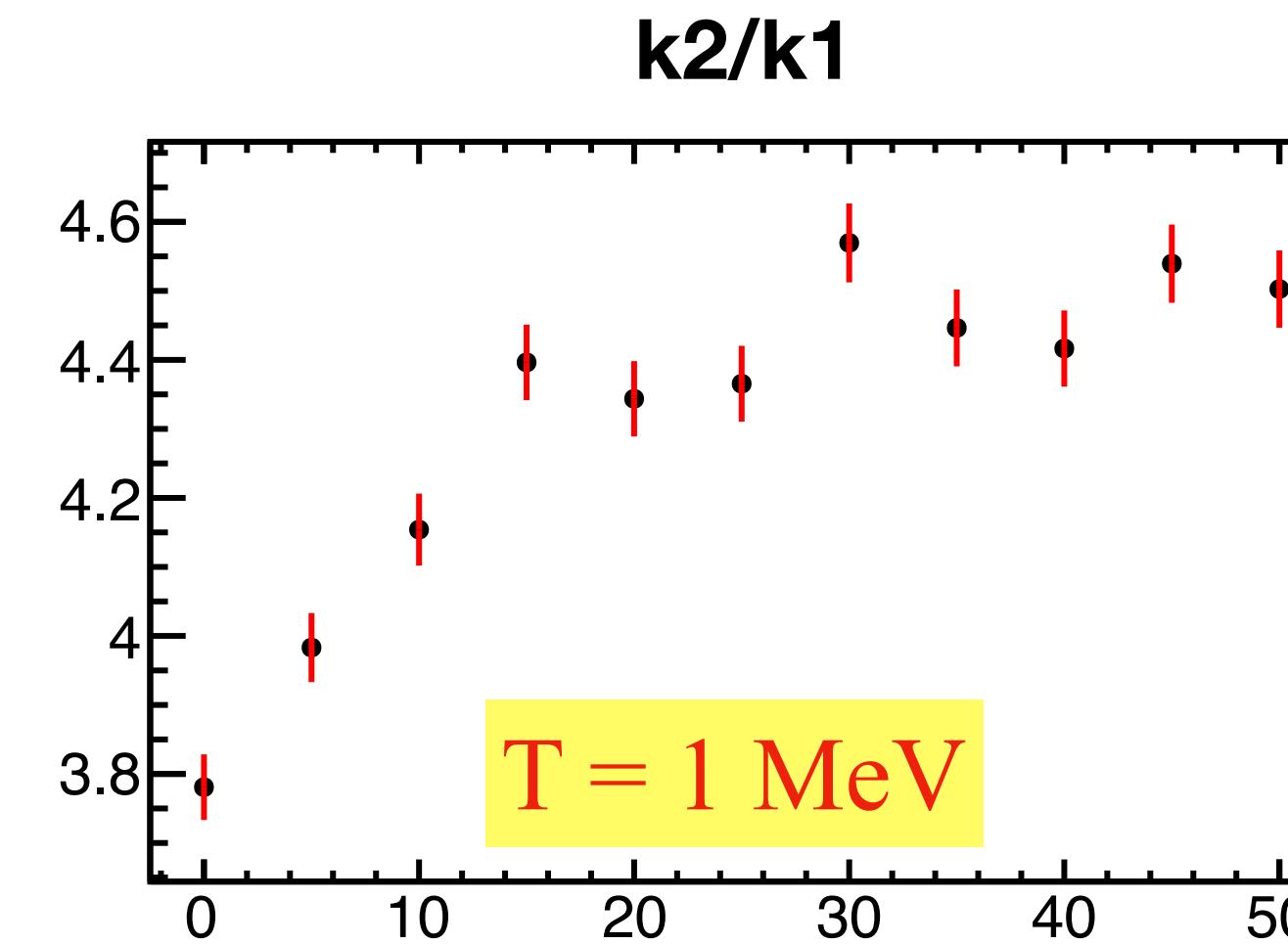
number density $n_B [n_0]$

SMASH results: periodic box with “parallel ensembles”

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Cumulants of the baryon number

$N_T = N_{\text{“parallel”}} = 50$

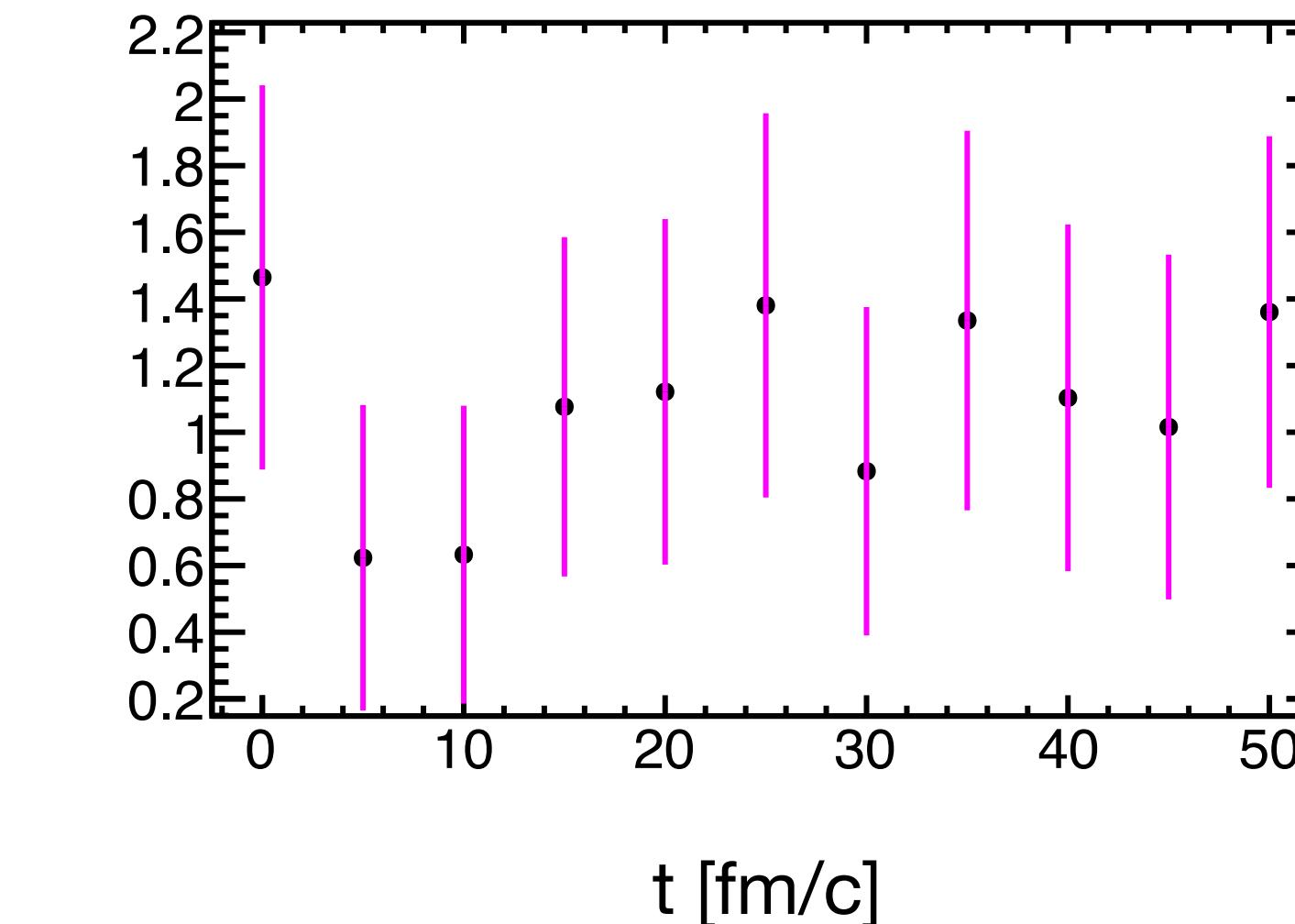
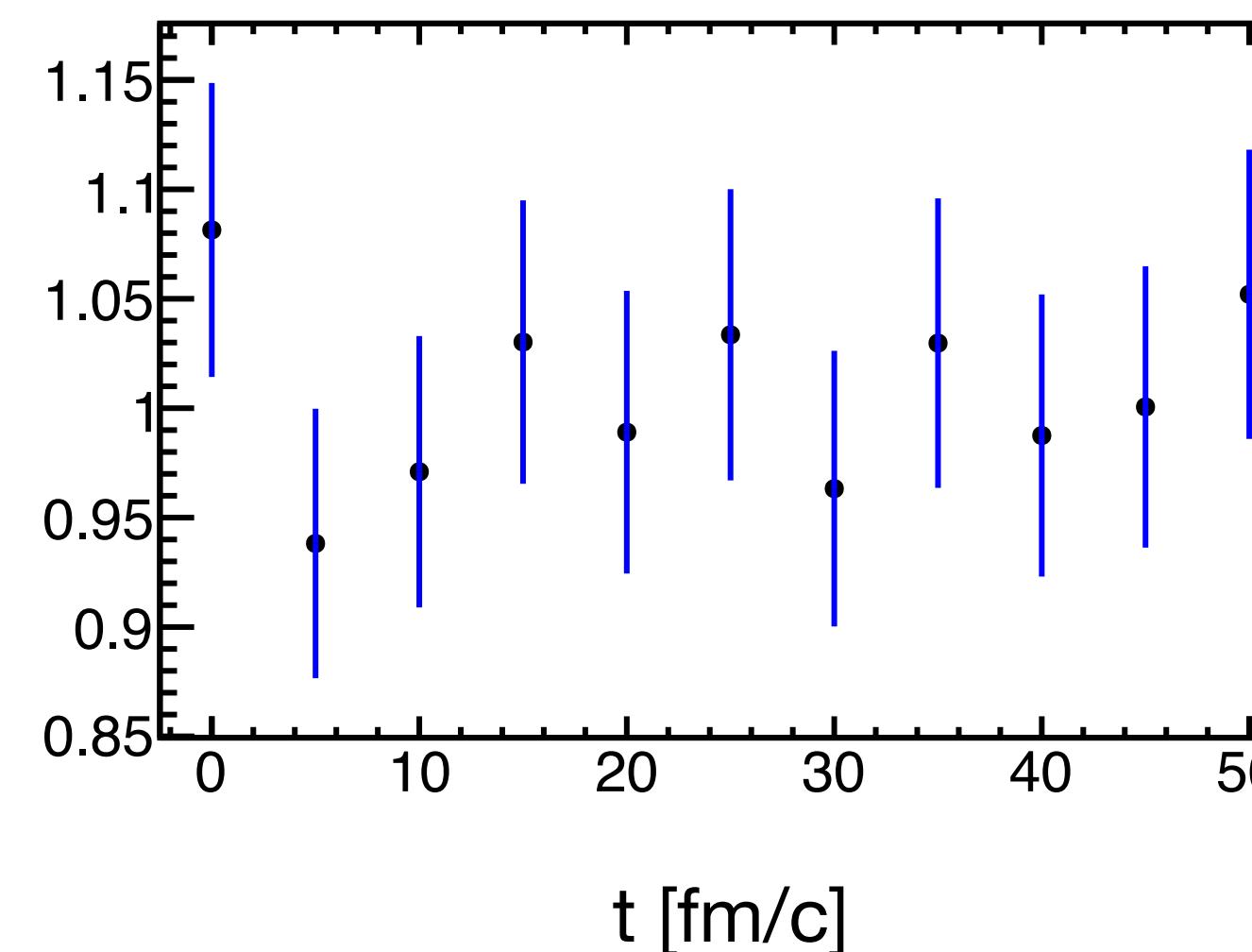
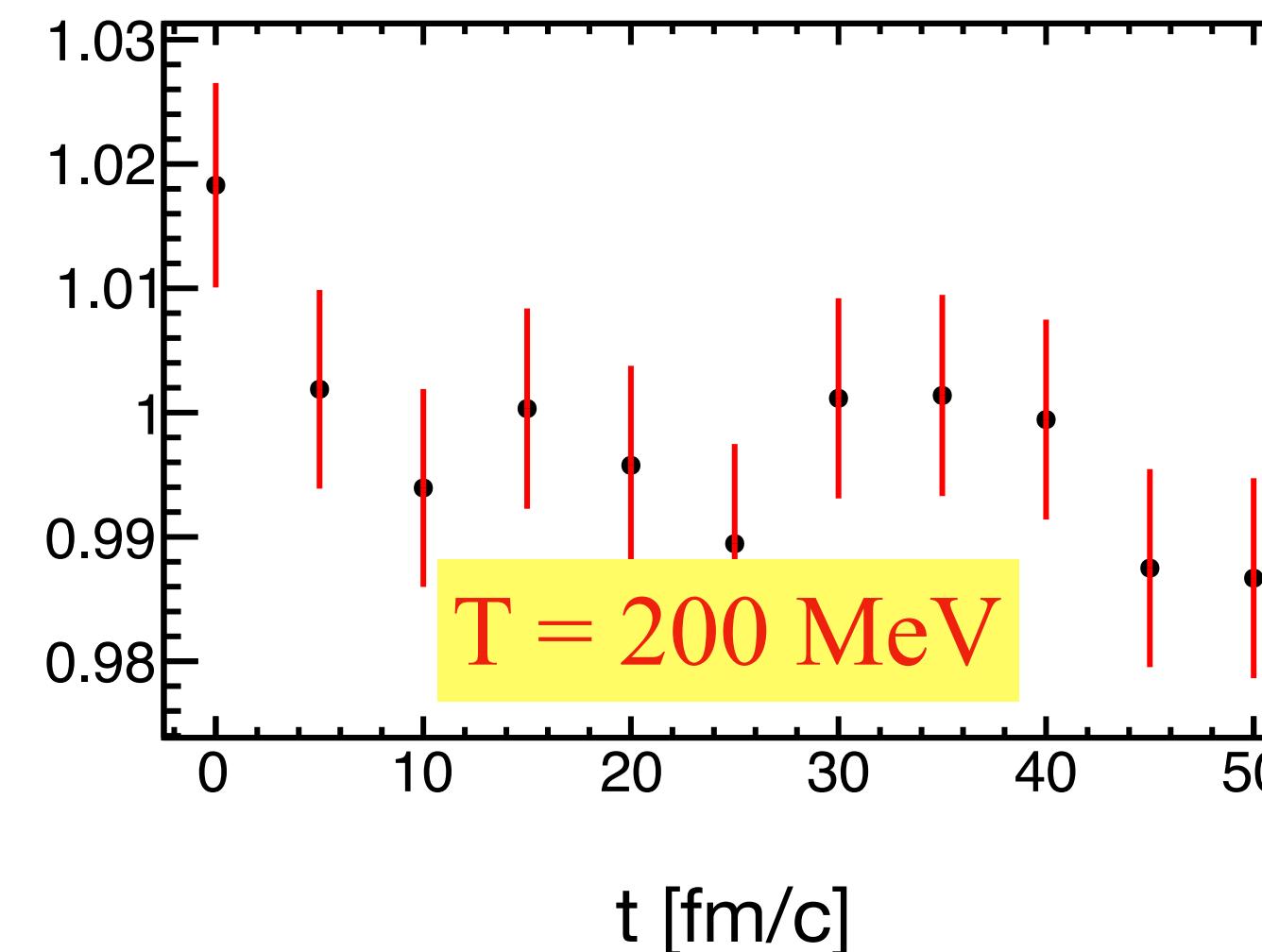
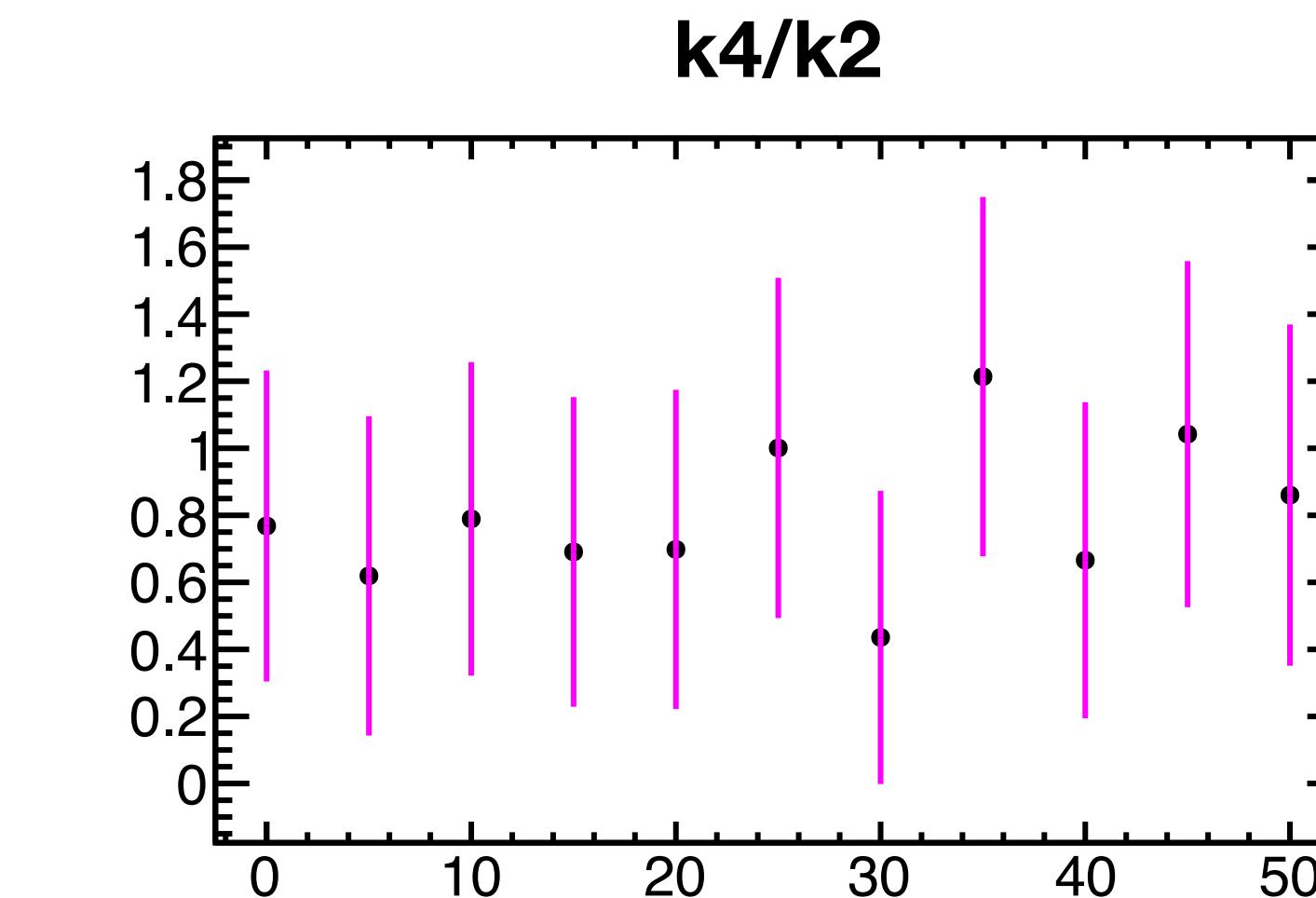
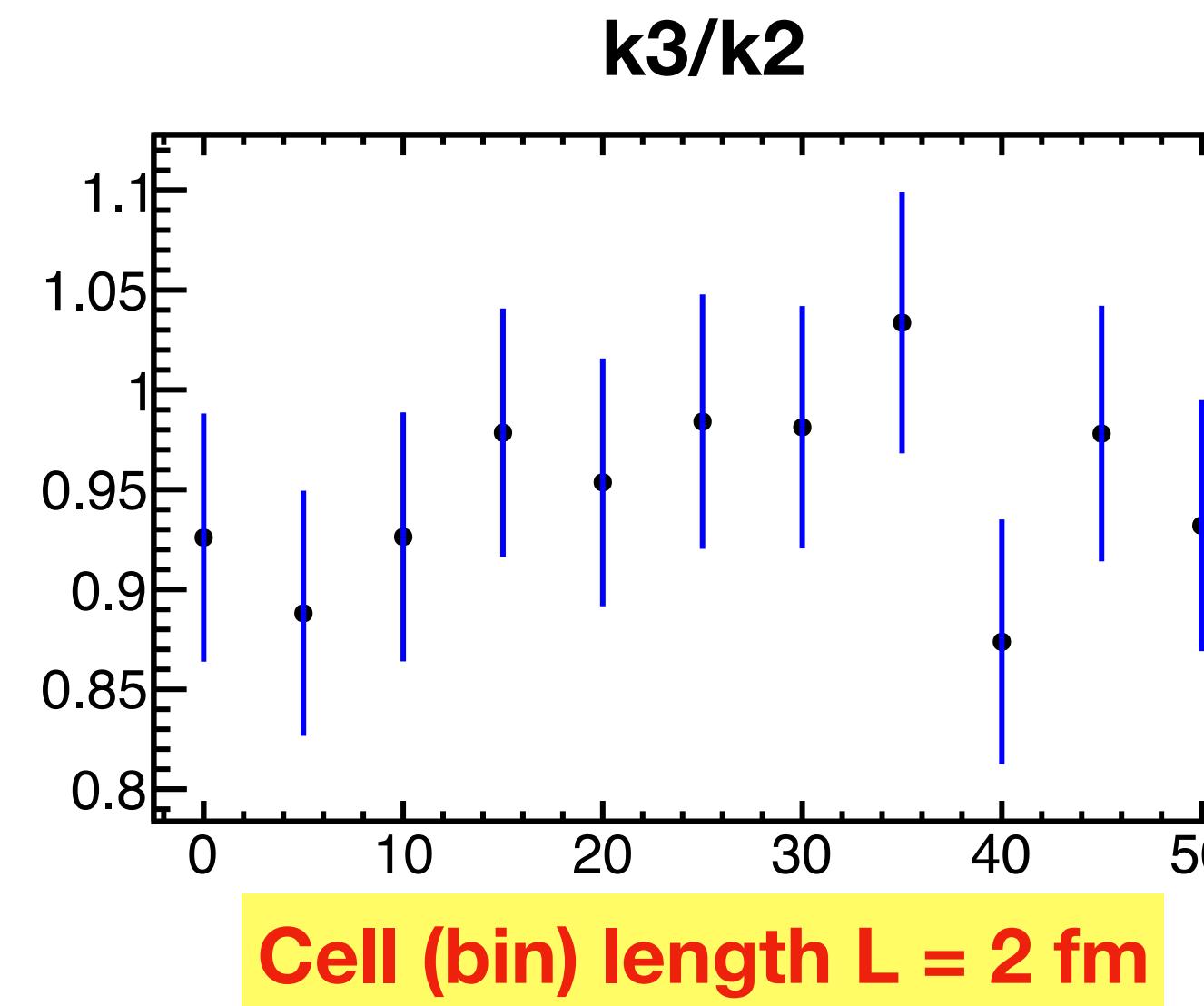
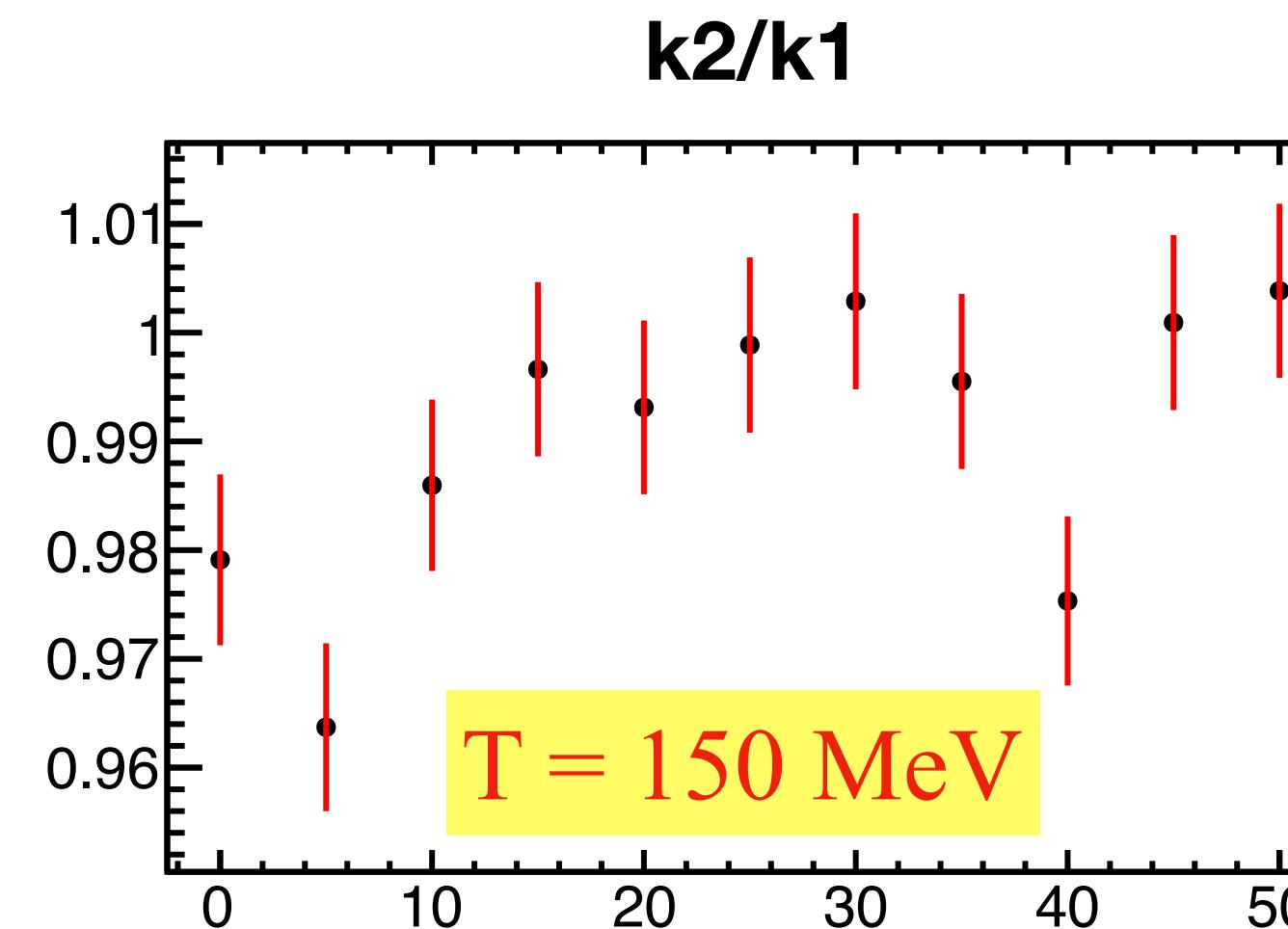


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0.0 fm/c

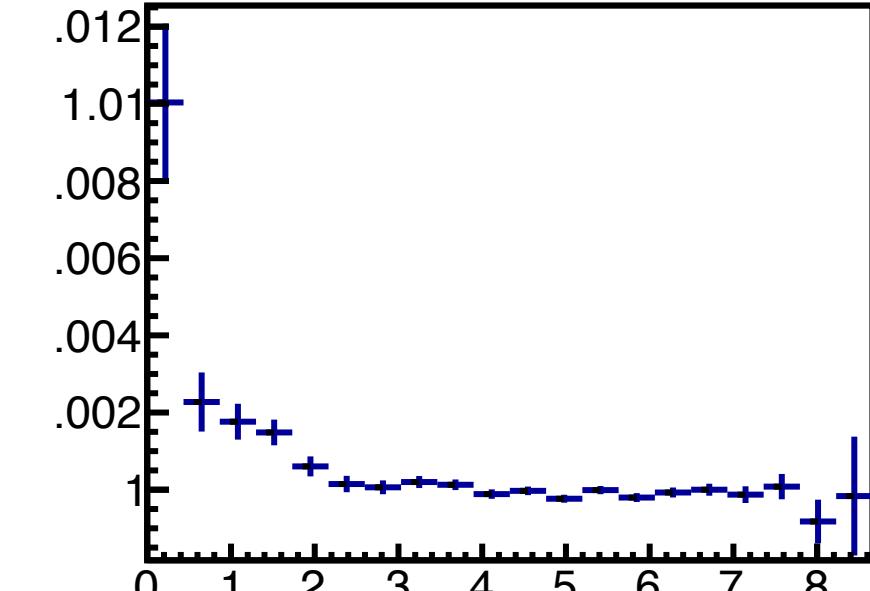
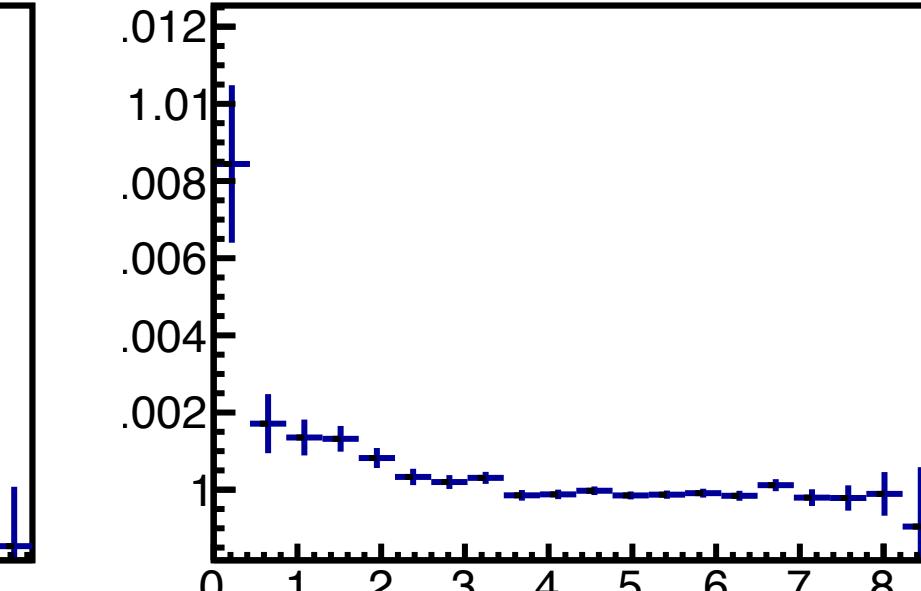
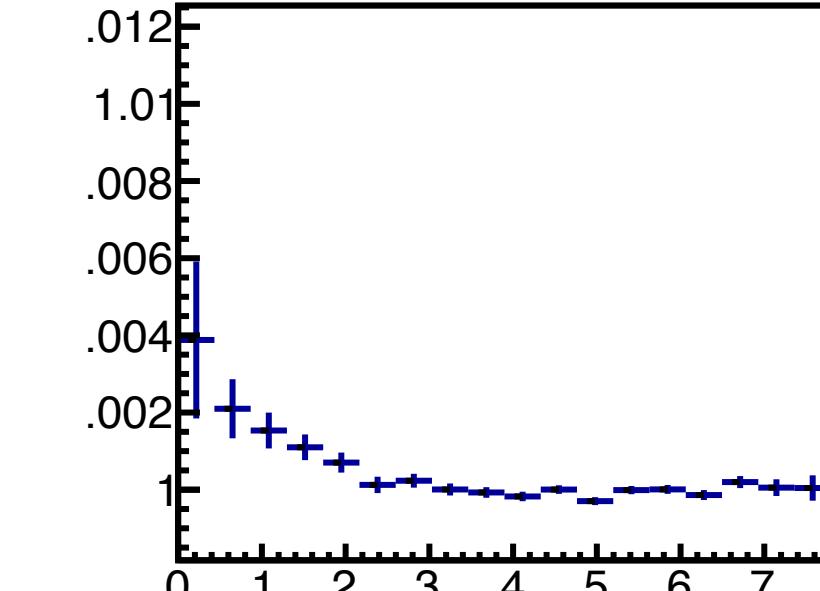
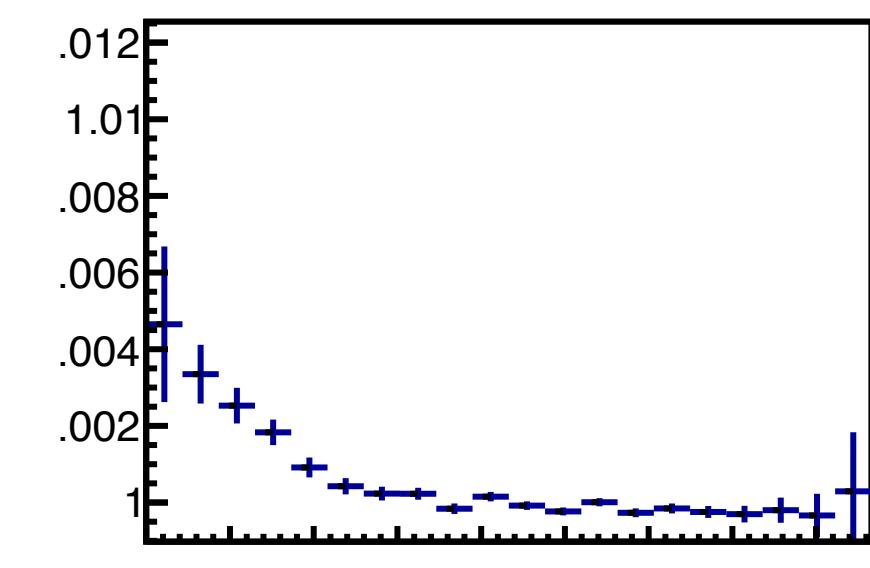
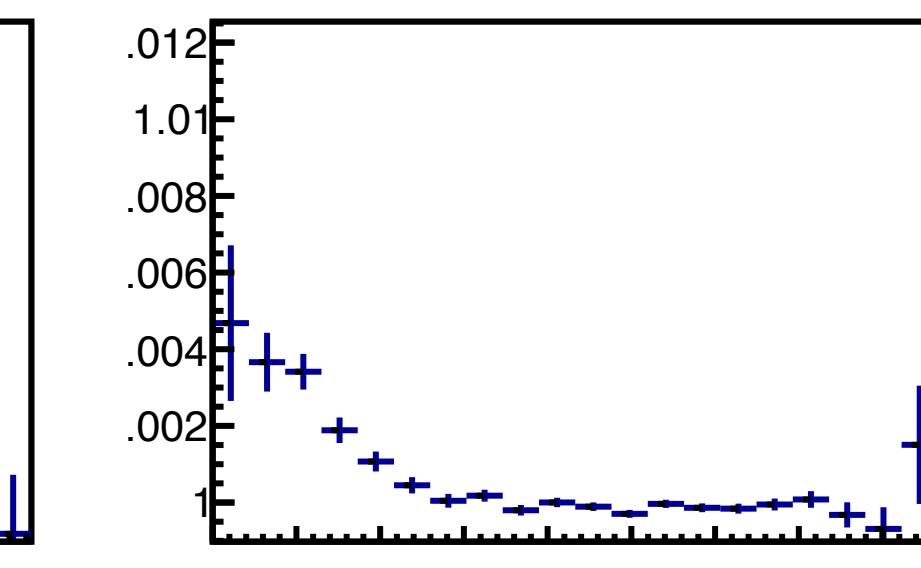
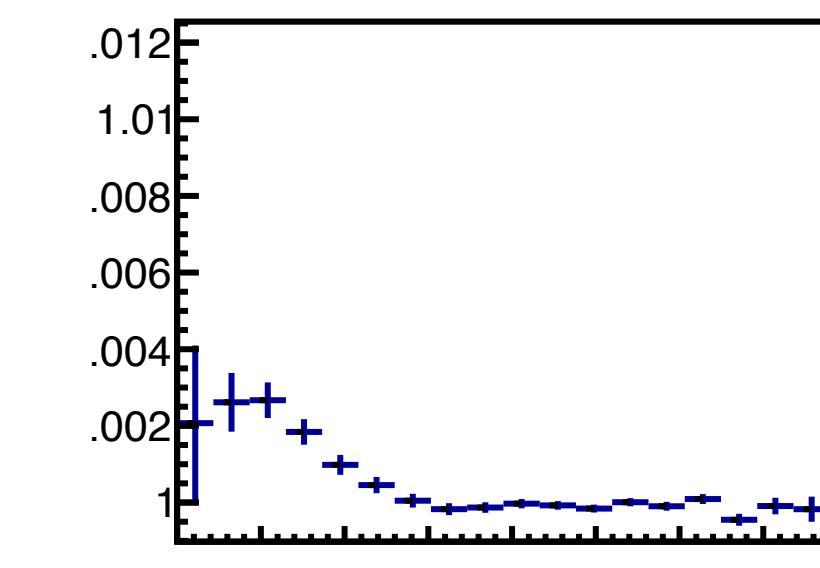
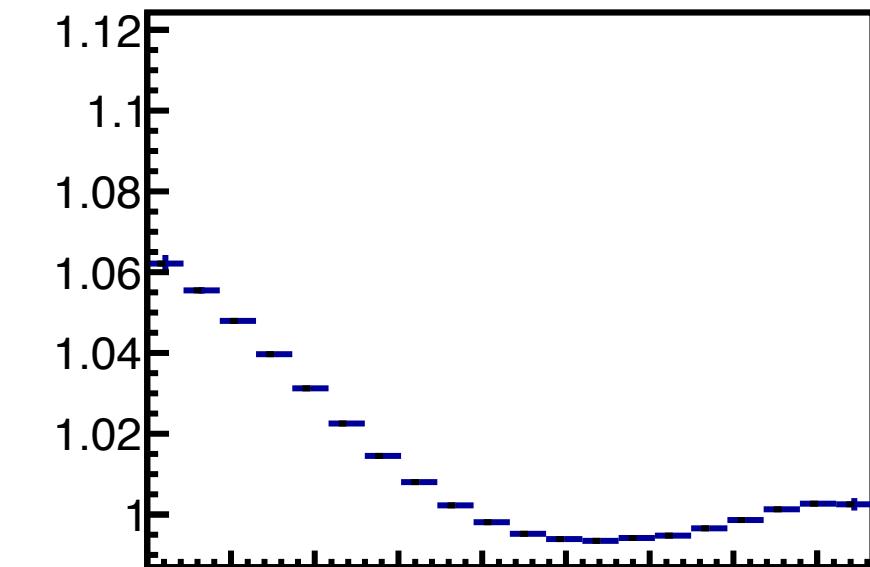
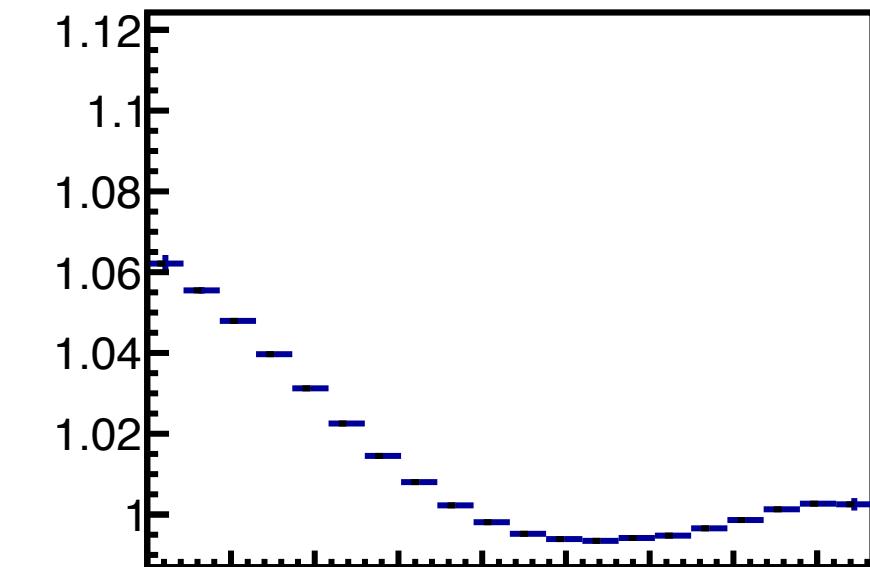
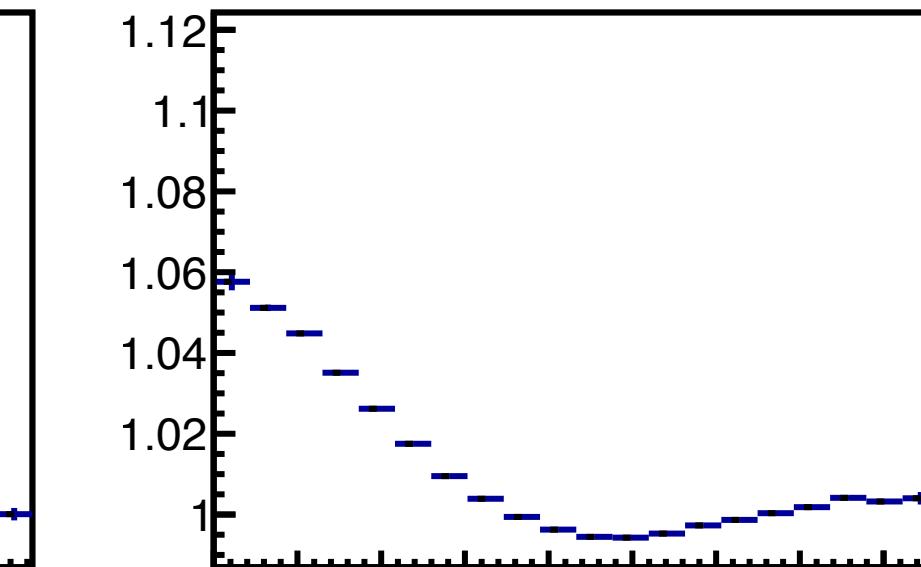
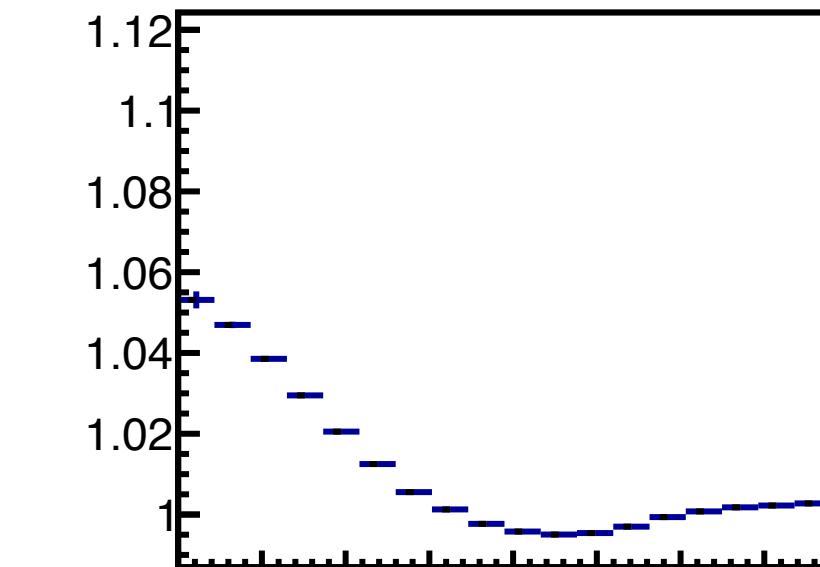
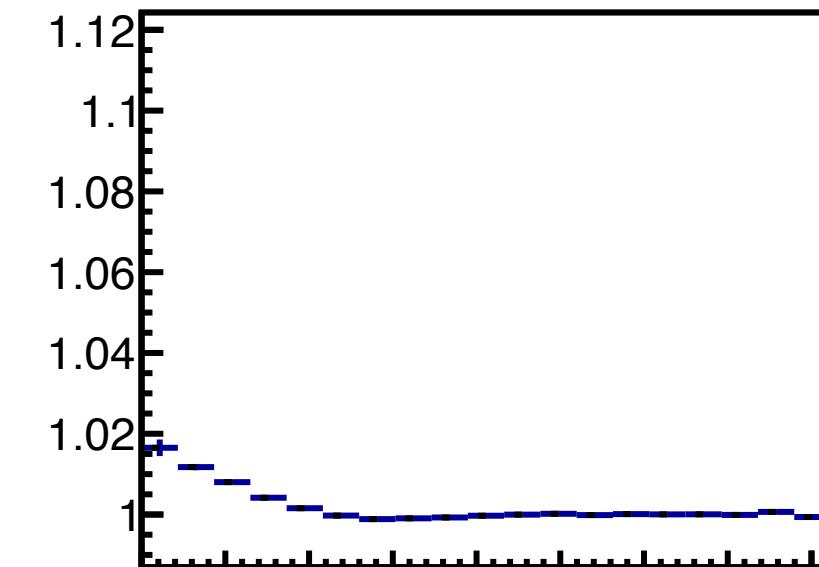
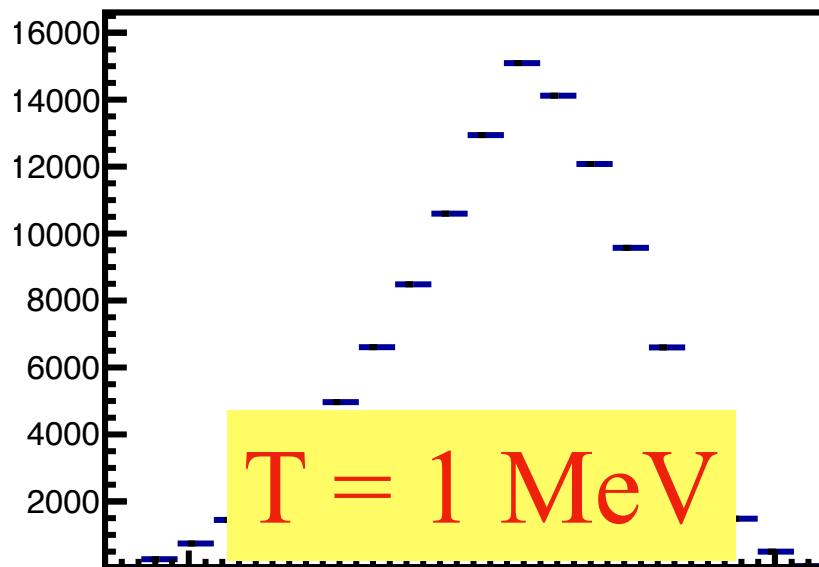
5.0 fm/c

15.0 fm/c

20.0 fm/c

30.0 fm/c

$3n_0$



$\sqrt{(\vec{r}_i - \vec{r}_j)^2} [\text{fm}]$

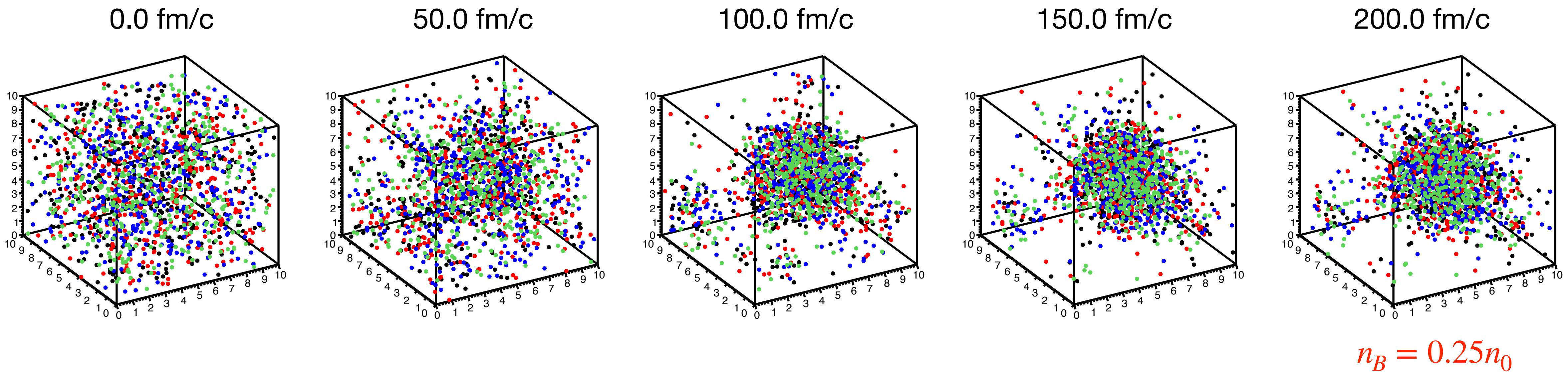
Summary

- Flexible parameterization of the dense QCD EOS including the known properties of nuclear matter
- Comprehensive approach to studying thermodynamics and dynamical evolution of heavy-ion collisions
- Promising initial results from implementation in SMASH

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Thank you for your attention



This material is based upon work supported by the U.S. Department of Energy, Office of Science, Office of Nuclear Physics, under contract number DE-AC02-05CH11231.