Hybrid renormalization for quasilight-front correlations in LaMET

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Quasi-light-front correlations

■ LaMET calculation begins with the spatial correlation [Ji, PRL 13' & SCPMA 14']

$$\tilde{h}(z, a, P^z) = \frac{1}{2N} \langle P | \mathcal{O}_{\Gamma}(z, 0) | P \rangle = \frac{1}{2N} \langle P | \bar{\psi}(z) \Gamma W(z, 0) \psi(0) | P \rangle$$

The defining operator renormalizes multiplicatively [Ji, JHZ, Zhao, PRL 18', Ishikawa et al, PRD 17', Green et al, PRL 18']

$$O_{\Gamma}(z,0) = e^{\delta m|z|} Z(a) O_{\Gamma,R}(z,0)$$

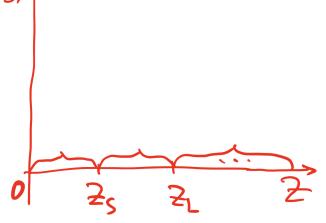
- Renormalization
 - Type 1: Determining δm , Z(a) separately
 - Type 2: Forming a ratio of matrix elements of the same operator [Stewart, Zhao, PRD 18', Radyushkin, PRD 17']

 [Braun, Vladimirov, JHZ, 19']

$$\begin{split} \tilde{h}_R(z,a,P^z) &= \tilde{h}(z,a,P^z)/Z_X(z,a) \\ Z_X(z,a) &= \frac{\langle p \,|\, O_\Gamma(z,0) \,|\, p \rangle\,|_{p^2=-\mu_R^2}}{\langle p \,|\, O_\Gamma(z,0) \,|\, p \rangle_{\text{tree}}}, \quad \tilde{h}(z,a,P^z=0), \quad |\, \langle 0 \,|\, O_\Gamma(z,0) \,|\, 0 \rangle \end{split}$$

Separate the z-range into different ways parts and treat them differently

 Avoid introducing extra nonperturbative effects at large z in the renormalization stage



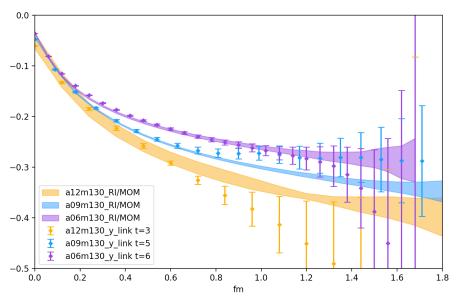
- At short distance $(0 \le |z| \le z_S)$, use the ratio in type 2 renormalization
- At long distance ($z_S \le |z| \le z_L$), use the type 1 mass renormalization
- ullet Beyond z_L , use physics-based extrapolation

Determination of renormalization factors

$$\delta m = \frac{m_{-1}}{a} + m_0, \quad Z(a)$$

- $oldsymbol{o} = m_0 \sim \mathcal{O}(\Lambda_{\rm QCD})$ depends on the matrix elements
- Asymptotically, one has

$$\tilde{h}(z, a, P^z) \sim e^{-\delta m|z|}$$



Plot from Jun Hua

Determination of renormalization factors

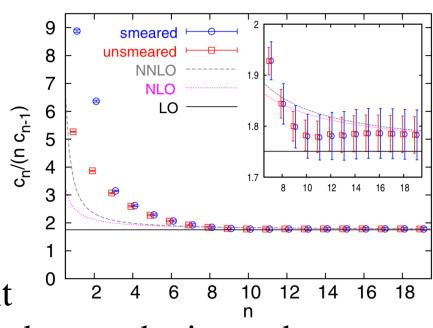
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■ In perturbation theory, m_{-1} starts from $\mathcal{O}(\alpha_s)$, can mix with m_0 due to renormalon effect

■ It has been argued that [Bauer, Bali & Pineda, PRL 12']

$$\delta m = \frac{1}{a} \sum_{n>0} c_n \alpha_s^{n+1} (1/a)$$

• m_0 is approximately constant $\frac{1}{2}$ $\frac{4}{6}$ $\frac{6}{8}$ $\frac{10}{n}$ $\frac{12}{14}$ when truncated at high enough perturbative orders



Determination of renormalization factors

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- In perturbation theory, m_{-1} starts from $\mathcal{O}(\alpha_s)$, can mix with m_0 due to renormalon effect
- Ideally, it would be chosen to match non-perturbatively to the continuum MSbar scheme condition

$$\delta m_{\overline{\rm MS}} = 0$$

• In practice, an optimal way could be to vary and identify the value that yields least momentum dependence

- Beyond z_L , we use a physics-based extrapolated form and may begin with $x^a(1-x)^b$ for the parton distribution
- Its Fourier transform

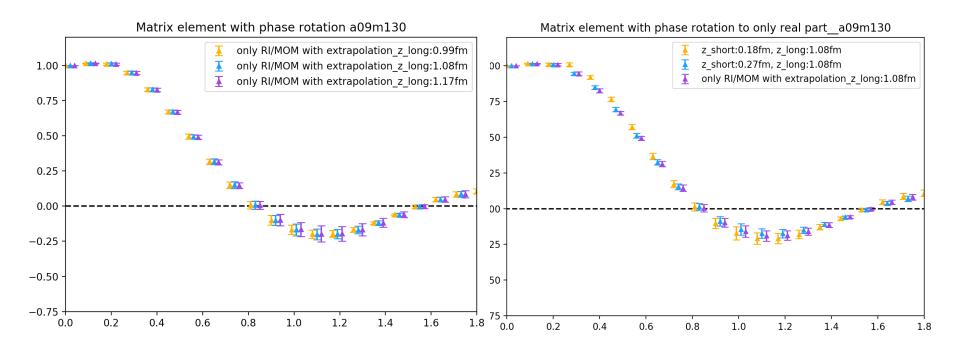
$$h(\lambda) = \int_0^1 dx \, e^{ix\lambda} x^a (1-x)^b \sim \frac{\Gamma(1+a)}{(-i|\lambda|)^{a+1}} + e^{i\lambda} \frac{\Gamma(1+b)}{(i|\lambda|)^{b+1}}$$
 at large λ

- Perturbative matching changes it to $\tilde{h}(\lambda) \sim |\lambda|^{-\alpha(a,b)+2\gamma}$, which also has a power decay behavior
- Extrapolation form can be chosen as

$$\frac{c_1}{(-i|\lambda|)^{d_1}} + e^{i\lambda} \frac{c_2}{(i|\lambda|)^{d_2}}$$

with unknown parameters being fixed from continuity requirement

- We need to investigate the dependence on the choice of z_s, z_L
- Example: pion distribution amplitude (see J. Hua's talk)



Plots from Jun Hua

Summary

 Hybrid renormalization offers a conceptually more viable renormalization scheme for LaMET calculations

- Further studies are required to establish robustness of this scheme
 - Optimal strategy to extract mass counterterms
 - Check independence of m_{-1} on various matrix elements used to compute it
 - \odot Determination of m_0
 - Applications to various parton quantities