

UV Divergence of the Quasi-PDF Operator under the Lattice Regularization

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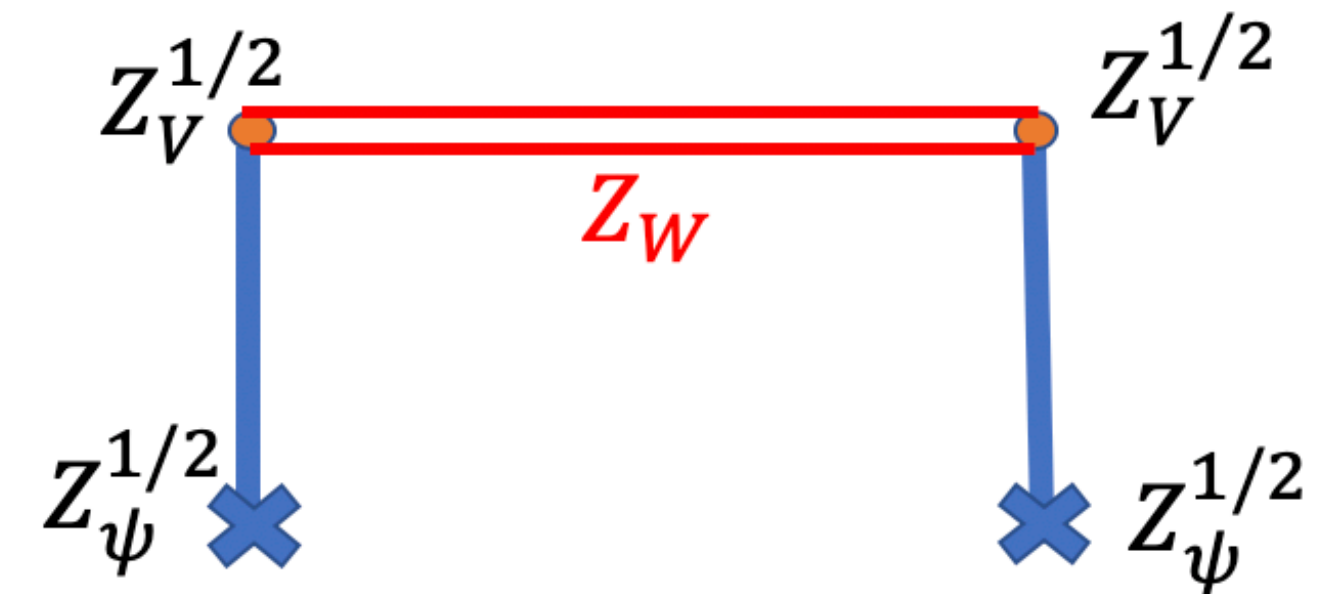
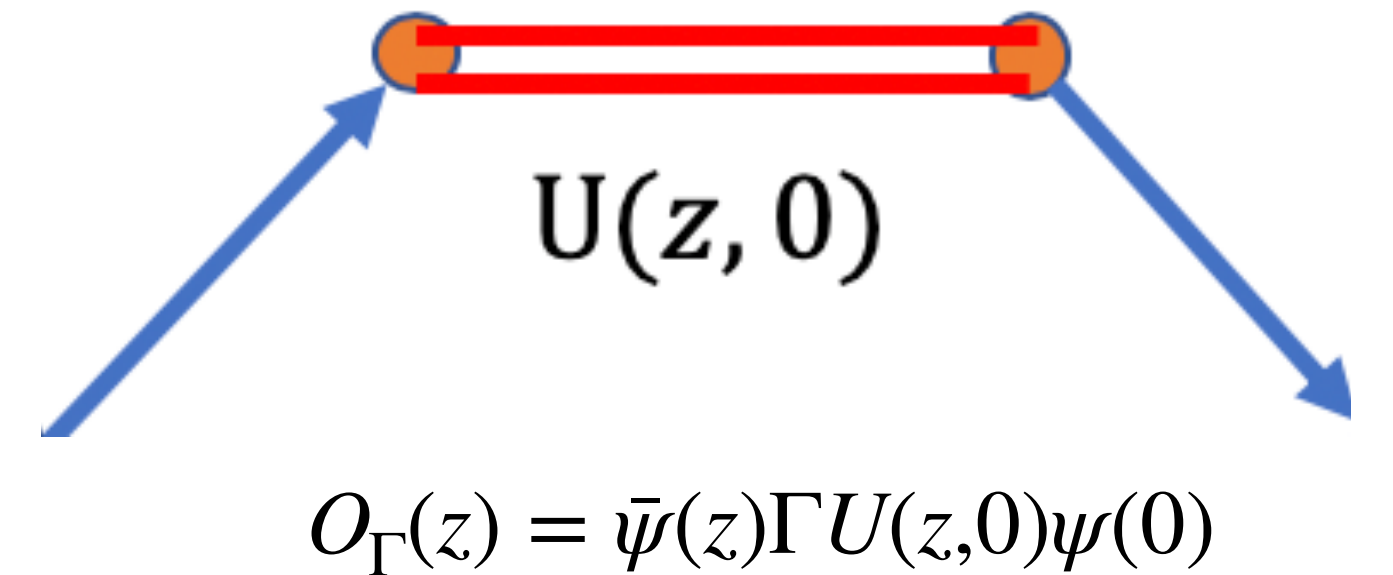
on behalf of Yi-Kai Huo and Peng Sun

Lattice Parton Collaboration

Quasi-PDF operator

Several renormalization procedures

- Naive counterterm $Z = Ce^{\frac{\delta m(a)}{a}z}$
- $\delta m(a)$ from $\langle O_{\gamma_z}(z) \rangle$,
- $\delta m(a)$ from $\langle O_{\gamma_t}(z, t) O_{\gamma_t}^\dagger(z, 0) \rangle$,
- $\delta m(a)$ from Wilson loop,
- $\delta m(a)$ from the regularization independent momentum subtraction (RI/MOM) scheme.
- Hybrid method based on RI/MOM



MILC configurations at 4 lattice spacings

β	L	L_t	C_{SW}	m_q	$a_{\text{phy}}(\text{fm})$
3.60	24	64	1.05088	-0.0695	0.120
3.78	32	96	1.04239	-0.05138	0.090
4.03	48	144	1.03493	-0.0398	0.060
4.20	64	192	1.03144	-0.0365	0.045

Accurate lattice spacing determination is in progress.

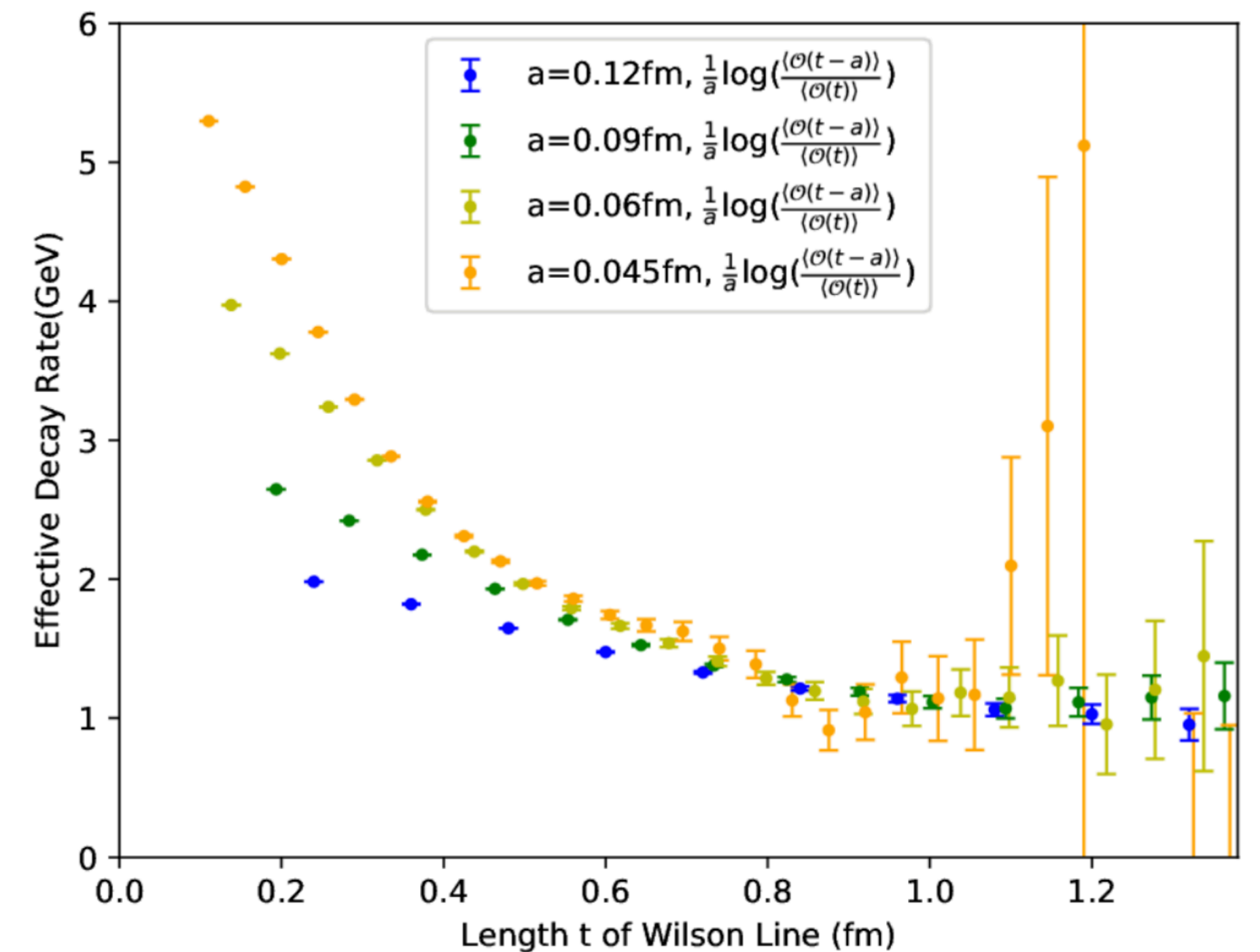
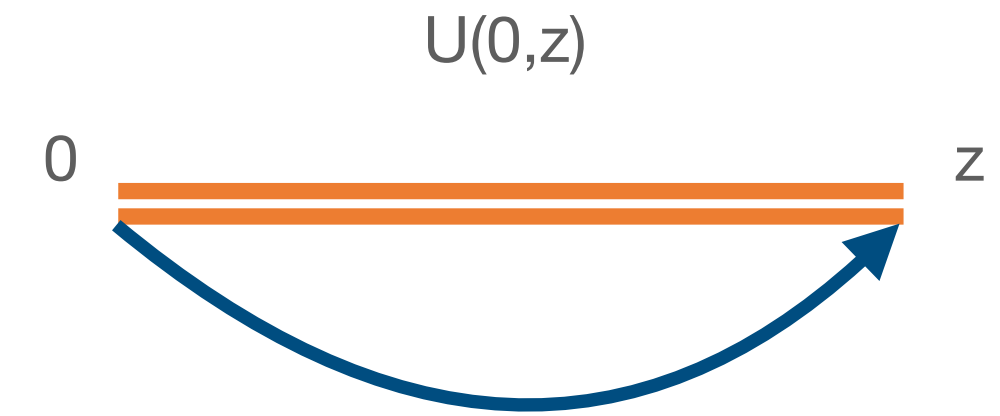
- $\delta m(a)$ from $\langle O_{\gamma_z}(z) \rangle$

- Vacuum expectation value of $O_{\gamma_z}(z)$

- $$\frac{1}{a} \log \left(\frac{\langle O_{\gamma_z}(z-a) \rangle}{\langle O_{\gamma_z}(z) \rangle} \right) = \frac{\delta m^{VEV}}{a} + f(z) + \mathcal{O}(a);$$

- Above ratio seems to converge in the continuum limit, and then $\delta m^{VEV} \sim 0$;

- It is understandable since the possible paths in the quark propagator increase exponentially with $1/a$, and then they would cancel the linear divergence.



- $\delta m(a)$ from $\langle O_{\gamma_t}(z, t) O_{\gamma_t}^\dagger(z, 0) \rangle$ and Wilson Loop

- $M(t, z) \equiv \langle O_{\gamma_t}(z) O_{\gamma_t}^\dagger(z) \rangle$

- $\frac{1}{a} \log\left(\frac{M(t, z - a)}{M(t, z)}\right) \Big|_{t=0.18\text{fm}} = 2\left(\frac{\delta m^{\text{VEV2}}}{a} + m_0(z)\right)$

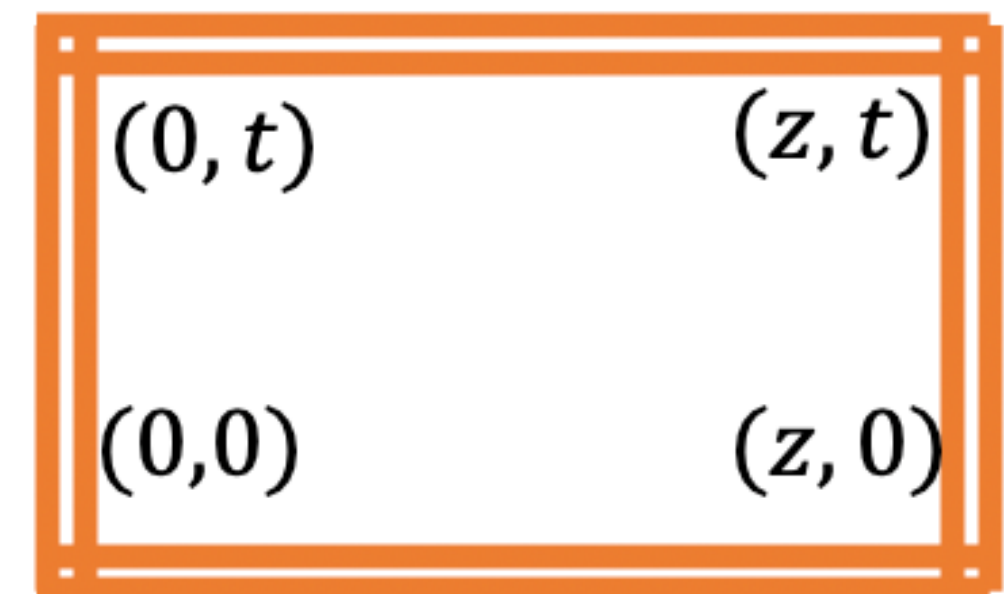
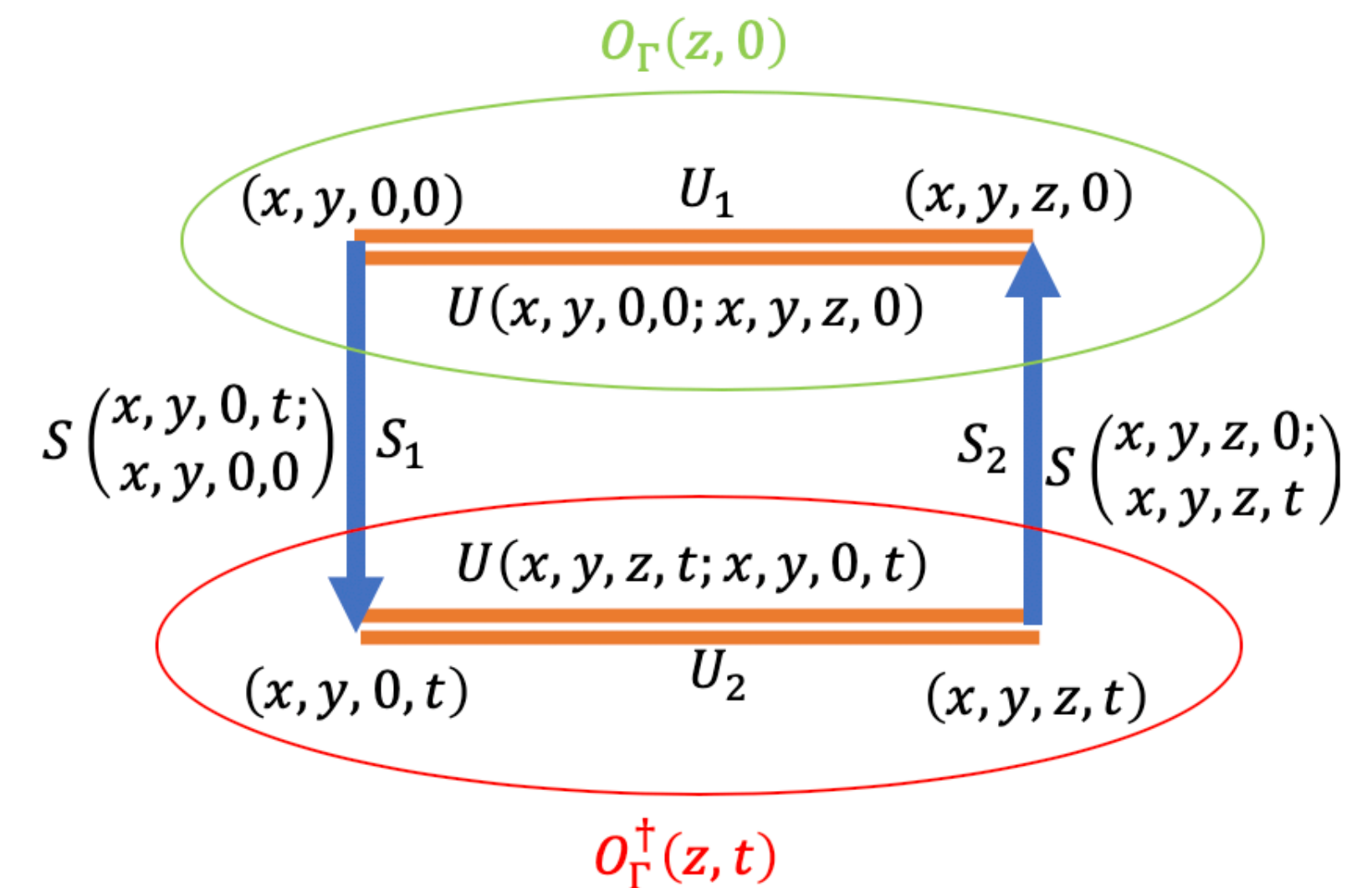
- **Wilson loop**

$$U(t, z) = \langle U(0, t; z, t) U(z, t; z, 0) U(z, 0; 0, 0) U(0, 0; 0, t) \rangle$$

- $\frac{1}{2} \log\left(\frac{U(t - a, z)}{U(t, z)}\right) \Big|_{t \rightarrow \infty} = \frac{c_0}{z} + 2\left(\frac{\delta m^{\text{Loop}}}{a} + m_0\right) + \sigma z$

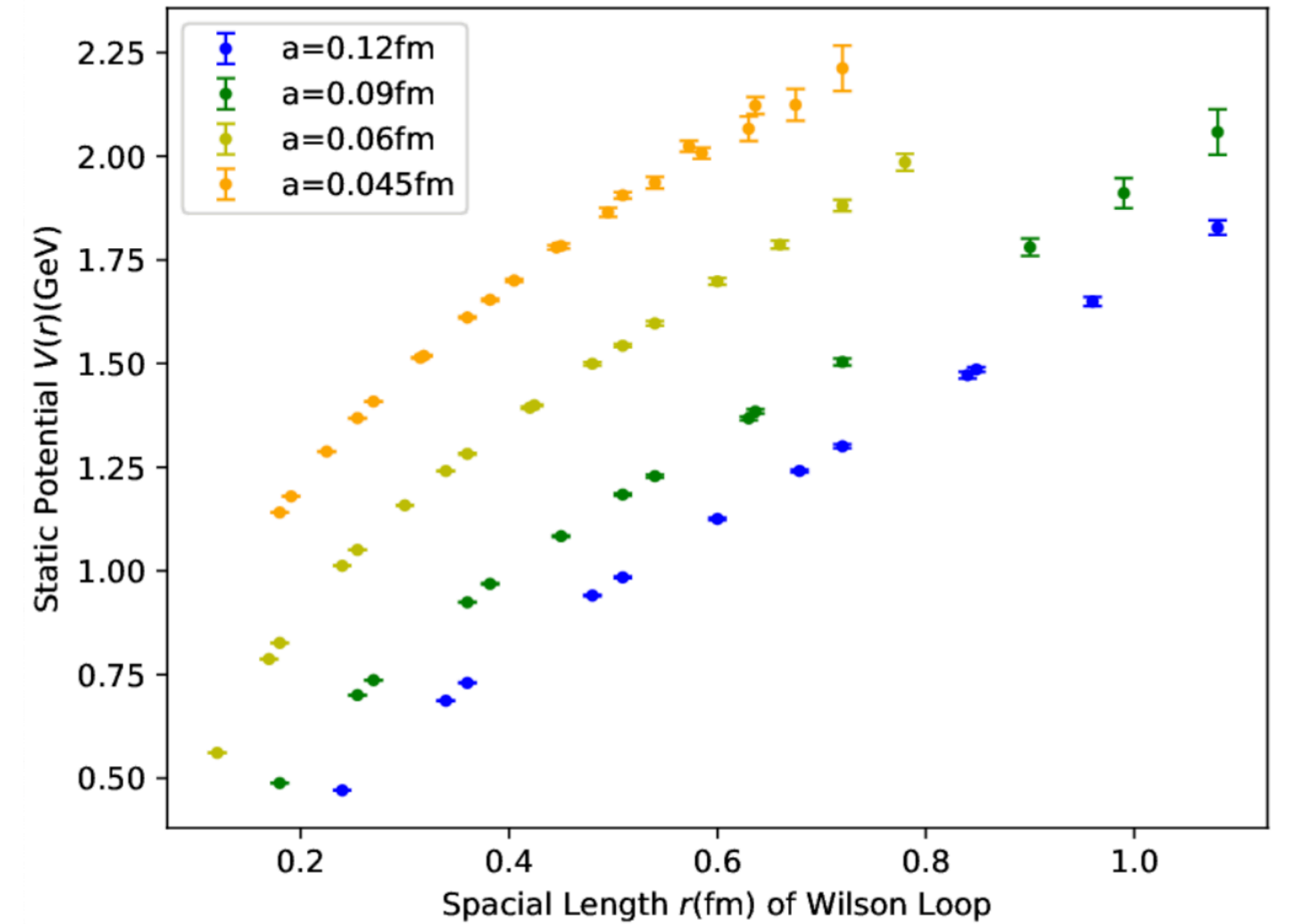
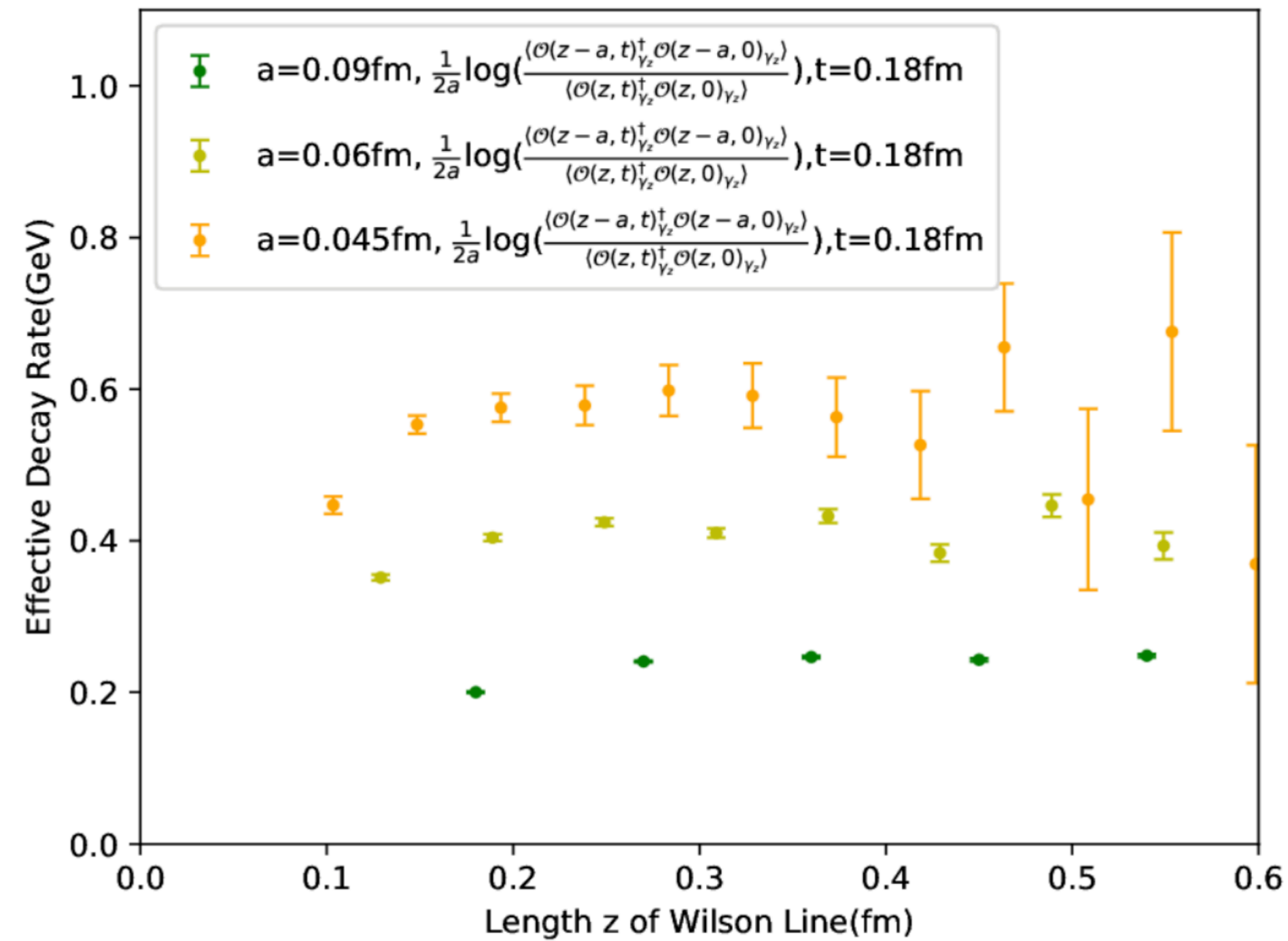
- Wilson loop uses infinite heavy quark line instead of the light quark propagator.

Definitions



- $\delta m(a)$ from $\langle O_{\gamma_t}(z, t) O_{\gamma_t}^\dagger(z, 0) \rangle$ and Wilson Loop

Results



$$\tilde{V}(z) = 2\left(\frac{\delta m^{VEV2}}{a} + m_0(z)\right)$$

$$\delta m^{VEV2} \sim 0.15(2)$$

The values in these two cases are consistent with each other.

$$V(z) = \frac{c_0}{z} + 2\left(\frac{\delta m^{Loop}}{a} + m_0\right) + \sigma z$$

$$\delta m^{Loop} \sim 0.15(1)$$

- $\delta m(a)$ from RI/MOM

- $$Z(z, a, \mu_R) = \frac{\langle p | O_{\gamma_t}(z) | p \rangle}{\langle p | O_{\gamma_t}(z) | p \rangle_{tree}} \Big|_{p^2 = -\mu_R^2, p_z = 0}$$

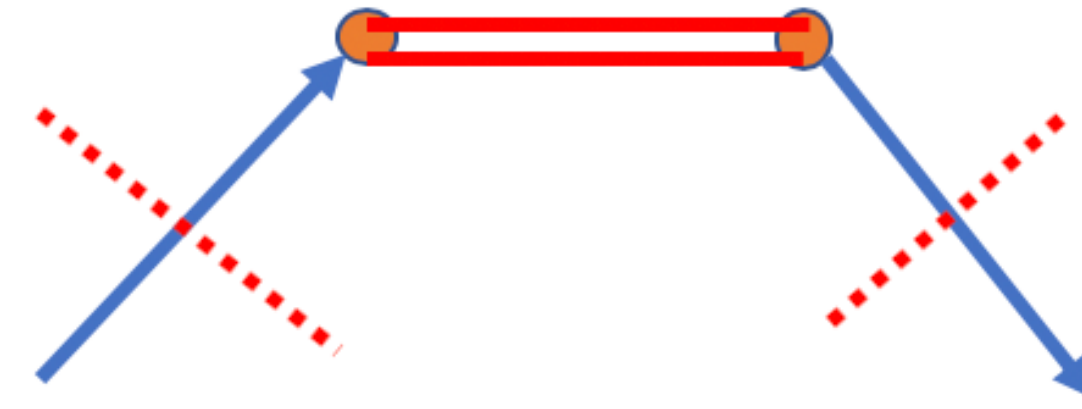
- $$\frac{1}{a} \log\left(\frac{Z(z, a)}{Z(z - a, a)}\right) = \frac{\delta m^{RI/MOM}}{a} + f(z);$$

- Fit $Z(z, a)$ with the form

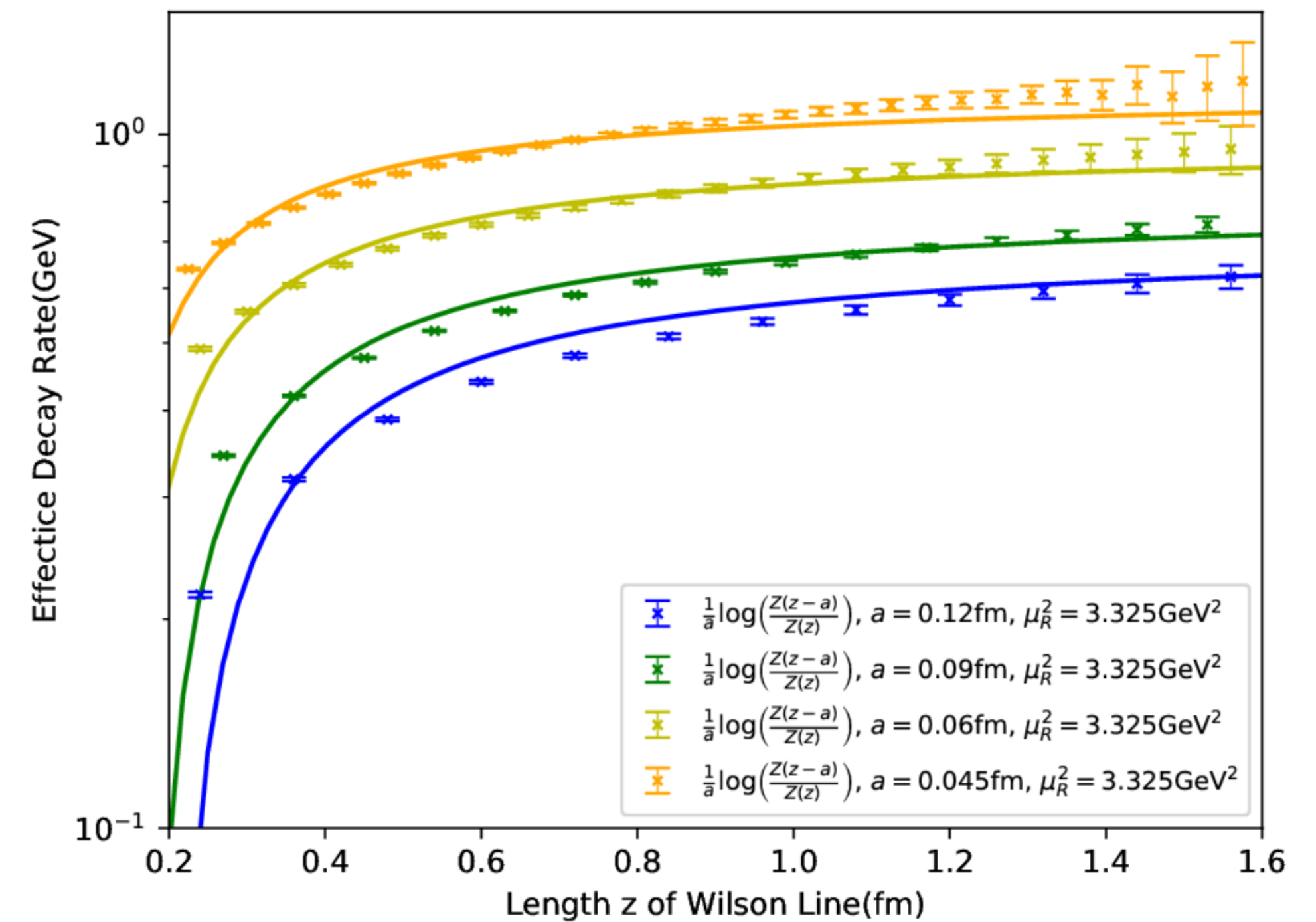
$$Z(z, a) = C_0(a) e^{-\frac{\alpha_s(a)}{\alpha_s(a=0.12\text{fm})} \left\{ \left(\frac{\delta m^{RI}}{a} + m_0^{RI} \right) z + c_1 \log\left(\frac{z}{z_0}\right) \right\}}$$

in the range $z \in [0.3, 2]\text{fm}$ and

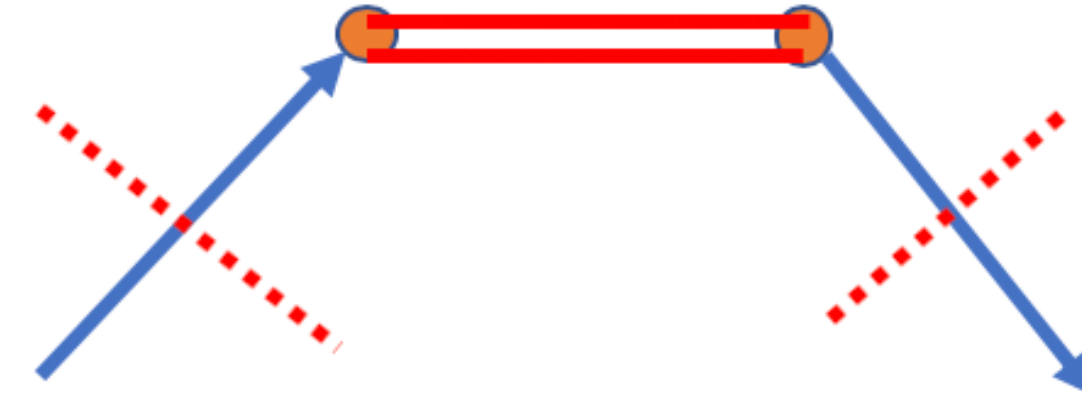
$a \in [0.04, 0.12]\text{fm}$ gives $\chi^2 \sim 130$ and $\delta m \sim 0.23(1)$.



Simple fit



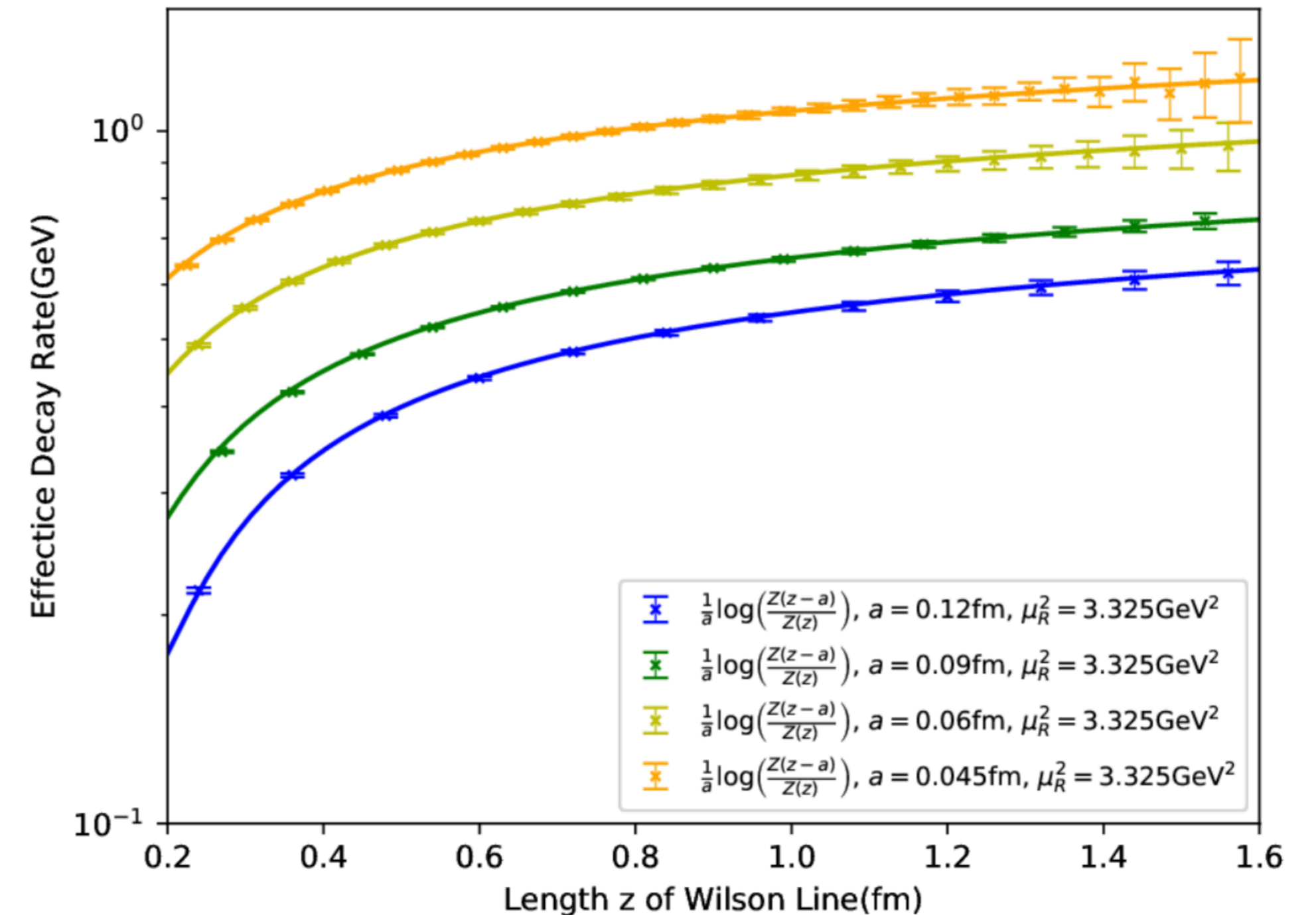
- $\delta m(a)$ from RI/MOM



Modified fit

$$Z(z, a) = C_0(a) e^{-\frac{\alpha_s(a)}{\alpha_s(a=0.12\text{fm})}(1+c_2\log(\frac{z}{z_0}))} \left\{ \left(\frac{\delta m^{RI}}{a} + m_0^{RI} \right) z + c_1 \log(\frac{z}{z_0}) \right\}$$

- The $1 + c_2 \log(\frac{z}{z_0})$ factor would come from the resummation on the $1/z$ scale dependence, and make the linear divergence to be stronger at larger z .
- Fit $Z(z, a)$ with above form in the range $z \in [0.3, 2]\text{fm}$ and $a \in [0.04, 0.12]\text{fm}$ gives $\chi^2 \sim 1$ and $\delta m \sim 0.16(1)$ with $z_0 = 0.3\text{fm}$.

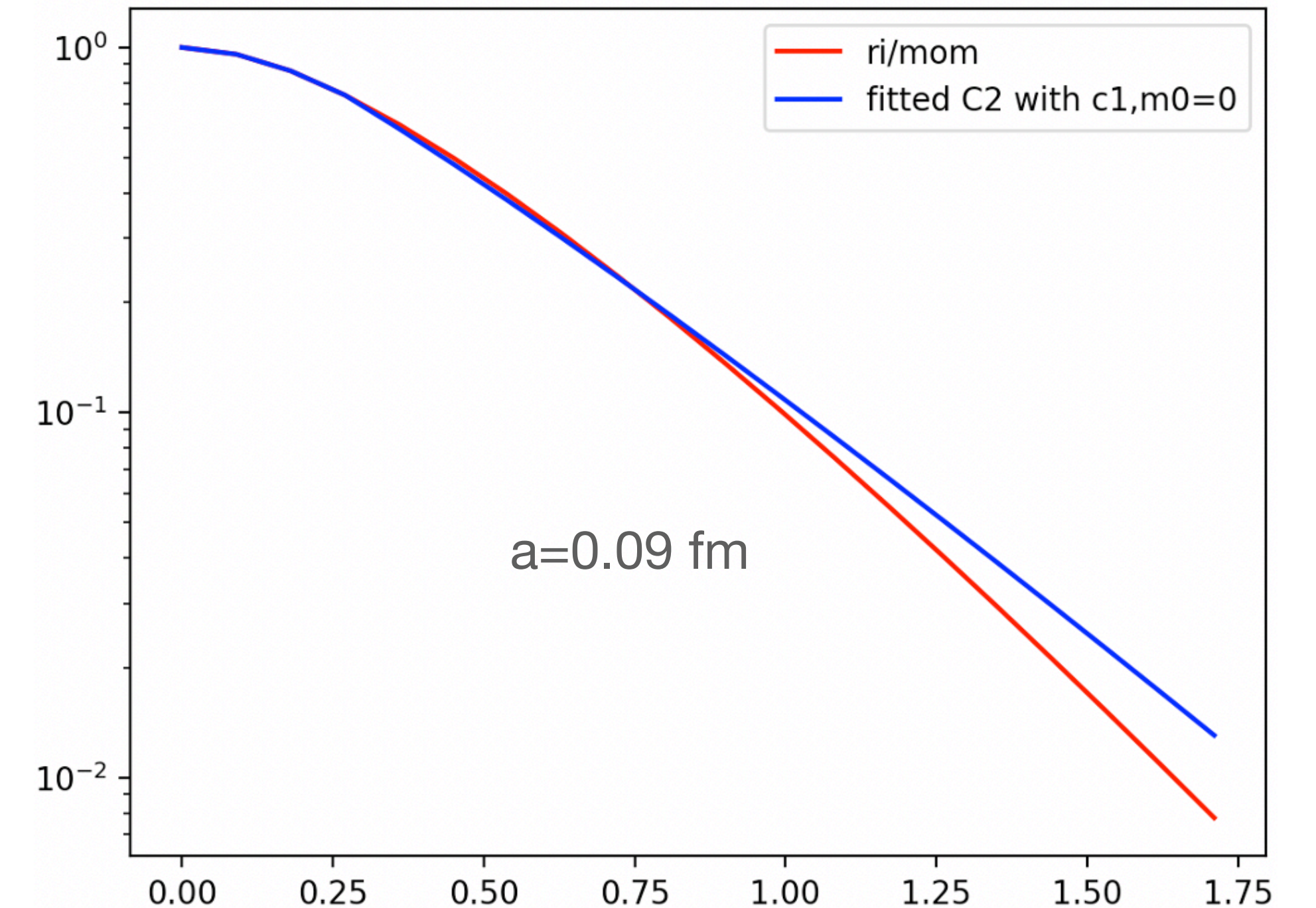


- Hybrid method based on RI/MOM
- Separate the z -range into different parts and treat them differently
- At small distance $0 \leq z \leq z_0$ ($z_0 \sim 0.3\text{fm}$), use RI/MOM renormalization

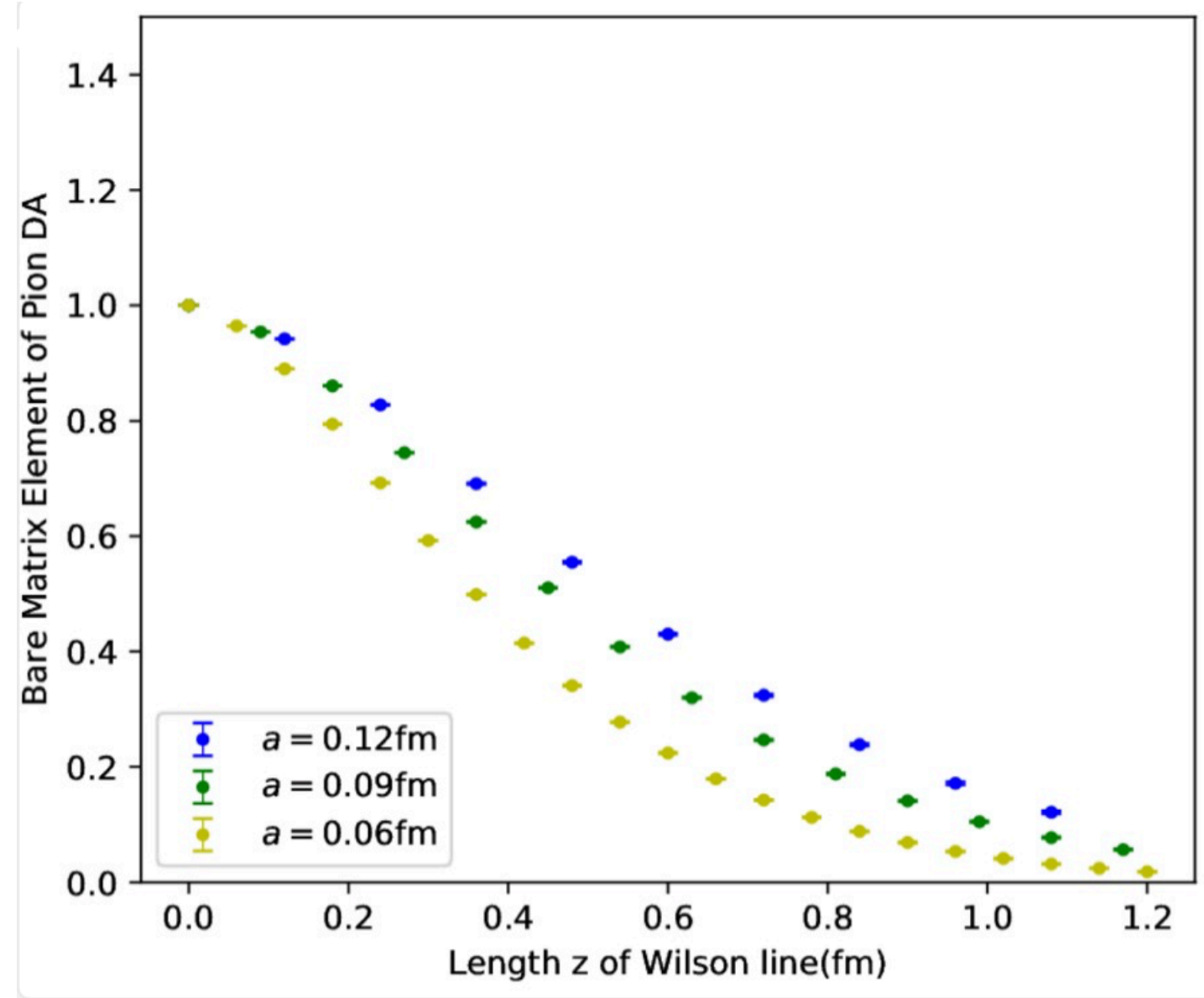
$$Z(z < z_0, a) = Z^{RI}(z, a)$$

- At larger distance $z_0 \leq z$, use modified mass renormalization

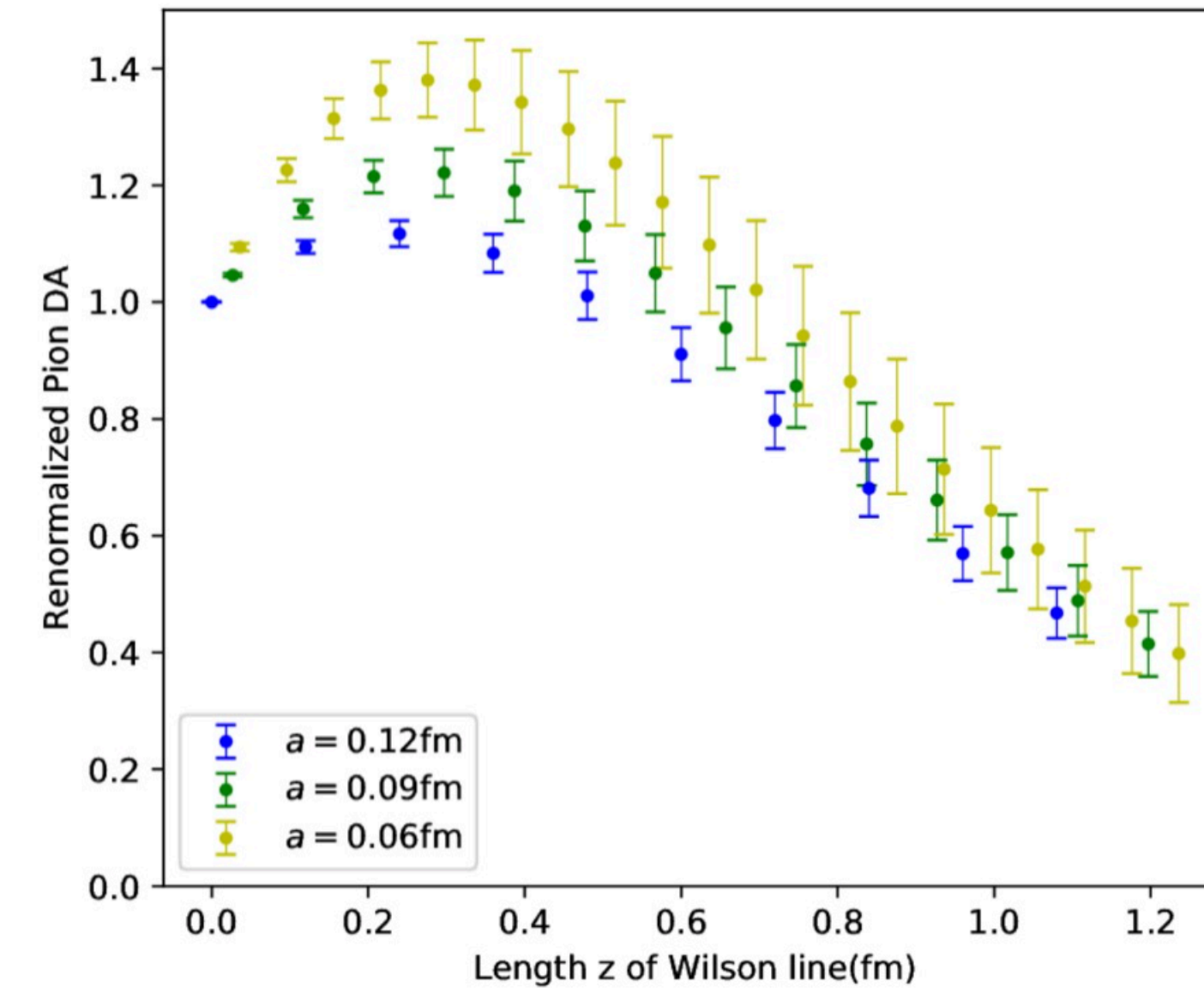
$$Z(z > z_0, a) = Z(z_0, a) e^{-\frac{\alpha_s(a)}{\alpha_s(a=0.12\text{fm})}(1+c_2\log(\frac{z}{z_0}))\frac{\delta m^{RI}}{a}(z-z_0)}$$



- Renormalized pion distribution amplitude (DA) matrix elements in the rest frame



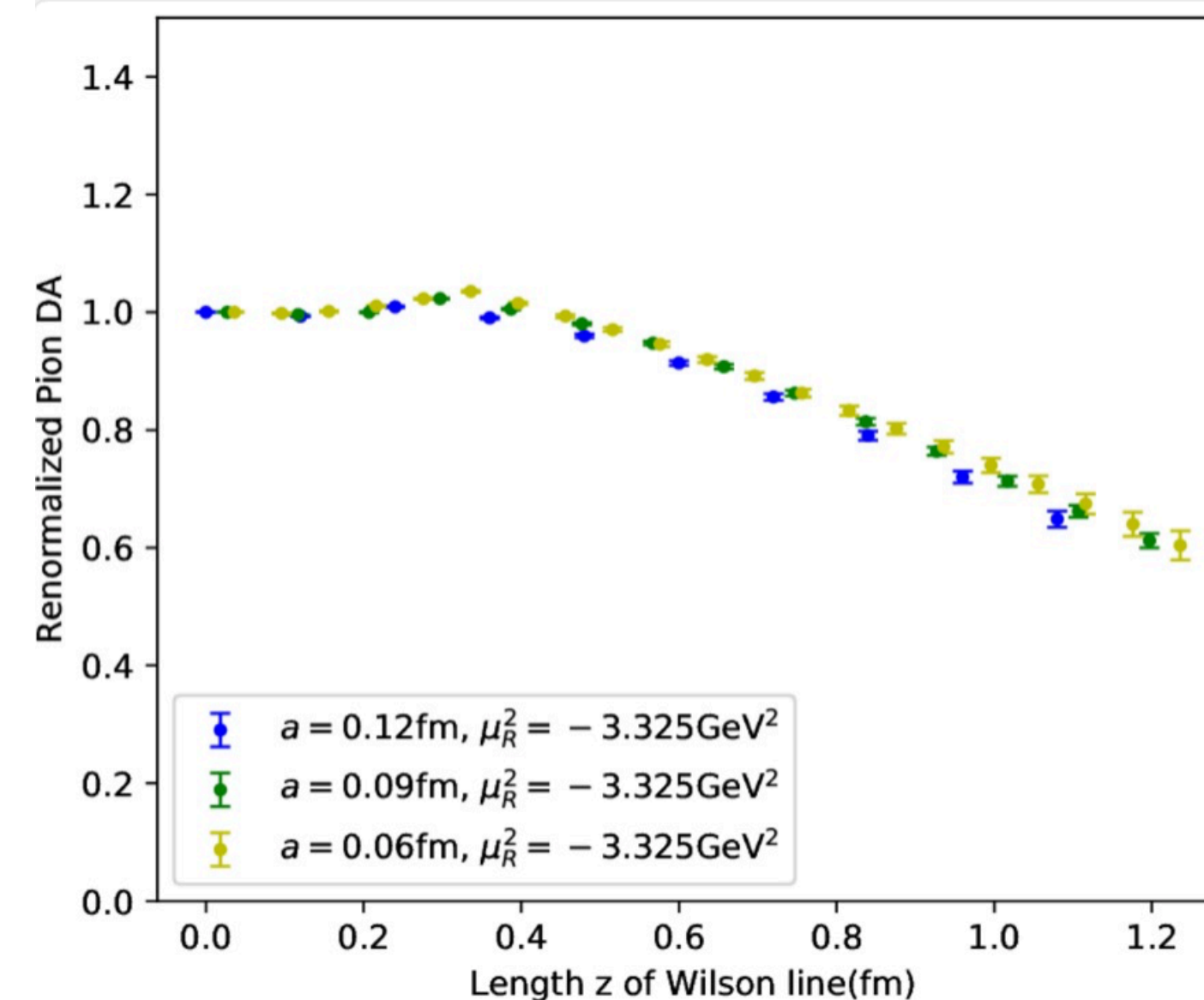
Bare DA with $\delta m = 0$



Renormalized DA with $\delta m = 0.15(2)$

$$Z \propto e^{\frac{\alpha_s(a)}{\alpha_s(a=0.12\text{fm})} \frac{\delta m^{Loop}}{a} z}$$

- Pion DA matrix element with $P_z = 0$;
- Only the hybrid scheme with modified mass renormalization can remove the linear divergence in the bare pion DA matrix element.



Renormalized DA with hybrid scheme

$$Z \propto e^{\frac{\alpha_s(a)}{\alpha_s(a=0.12\text{fm})} (1+c_2 \log(\frac{z}{z_0})) \frac{\delta m^{RI}}{a} (z-z_0)}$$

Summary

- We calculated the lattice spacing dependence of kinds of matrix elements with Wilson links.
- A $\log(z)$ dependence of the linear divergence term exists in both the RI/MOM quark matrix element and hadron ones, but not $\langle O_{\gamma_z}(z) \rangle$, $\langle O_{\gamma_t}(z, t) O_{\gamma_t}^\dagger(z, 0) \rangle$ or Wilson loop.
- Hybrid renormalization scheme with the modified mass renormalization,
$$Z(z > z_0, a) = Z(z_0, a) e^{-\frac{\alpha_S(a)}{\alpha_S(a=0.12\text{fm})} (1 + c_2 \log(\frac{z}{z_0})) \frac{\delta m^{RI}}{a} (z - z_0)},$$
 can be a proper choice to remove the linear divergence in the bare hadron matrix element of $O_{\gamma_z}(z)$.