A study of the anomalous magnetic moment of the muon computed from the Adler function

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The Adler function

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Adler function

The Adler function is defined as¹

$$\frac{D(q^2)}{q^2} = \frac{3\pi}{\alpha} \frac{d}{dq^2} \Delta \alpha_{had}(q^2),$$

and it can be measured in $e^+e^- \rightarrow hadrons$. The Adler function is related to the vacuum polarization by

$$D(q^2) = 12\pi^2 q^2 \frac{d\Pi(q^2)}{d(q^2)},$$

and if we use $D(q^2)$ to determine $\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$ we avoid the extrapolation to $\Pi(q^2 = 0)$. Steps of the analysis:

- $\ \, {\rm 0} \ \, {\rm determine} \ \, \Pi(q^2) \ \, {\rm on} \ \, {\rm the} \ \, {\rm lattice} \ \ \,$
- 2 take the numerical derivative of $\Pi(q^2)\,\rightarrow\,D(q^2)$
- ${igsin}$ in order to describe the q^2 -dependence apply a fit to $D(q^2)$
- determine a_{μ}^{HLO} from this fit
- ¹Adler, Phys. Rev. **D** 10, 3714, 1974

CLS-Ensembles

In our study we use O(a)-improved Wilson-fermions with $N_f = 2$ with partially twisted boundary conditions. The strange and charm quarks are partially quenched.

Label	V	β	$a[fm]^*$	$m_{\pi}[MeV]$	$m_{\pi}L$	N_{cnfg}	N_{meas}
A3	64×32^3	5.20	0.079	473	6.0	251	1004
A4	64×32^3	5.20	0.079	363	4.7	400	1600
A5	64×32^3	5.20	0.079	312	4.0	251	1004
B6	96×48^3	5.20	0.079	267	5.1	306	1224
E5	64×32^3	5.30	0.063	456	4.7	1000	4000
F6	96×48^3	5.30	0.063	325	5.0	300	1200
F7	$96 imes 48^3$	5.30	0.063	277	4.2	250	1000
G8	128×64^3	5.30	0.063	193	4.0	205	820
N5	96×48^3	5.50	0.050	430	5.2	347	1392
N6	96×48^3	5.50	0.050	340	4.1	559	2236
07	128×64^3	5.50	0.050	261	4.4	138	552

* [arXiv:1110.6365]

Step 1: Determination of $\Pi(q^2)$

The vacuum polarization tensor can be computed by

$$\Pi_{\mu\nu}(q^2) = \int d^4x e^{iqx} \left\langle J_{\mu}(x) J_{\nu}(0) \right\rangle.$$

From Euclidean invariance and current conservation one finds

$$\Pi_{\mu\nu}(q^2) = \left(g_{\mu\nu}q^2 - q_{\mu}q_{\nu}\right)\Pi(q^2).$$





Fit to $\Pi(q^2)$

Fit an ansatz to $\Pi(q^2)$, and compute the derivative of the fit function. We use the Padé ansatz

$$\Pi_{fit}(q^2) = \Pi(0) - q^2 \left(\frac{a_1}{q^2 + b_1} + \frac{a_2}{q^2 + b_2}\right),$$
$$\frac{d}{dq^2} \Pi_{fit}(q^2) = -\frac{a_1 b_1}{(b_1 + q^2)^2} - \frac{a_2 b_2}{(b_2 + q^2)^2}.$$

[arXiv:1112.2894, arXiv:1205.3695]

a

Numerical derivative of $\Pi(q^2)$

We use linear fits with varying ranges to approximate the derivative of $\Pi(q^2)$.

Step 2: Procedures for the numerical derivative 1

Procedure |

• at each q^2 perform a linear fit

$$\Pi_{fit}^{[l]}(q^2) = a_l + b_l q^2,$$

- \bullet repeat these fits for several fit ranges $\epsilon \in [0.1, 1.0] {\rm GeV}^2$,
- search for a region in ϵ where variations in b_l are small.



Step 2: Procedures for the numerical derivative 2

Procedure II

 ${\, \bullet \,}$ at each q^2 we fit the two functions

$$\Pi_{fit}^{[l]}(q^2) = a_l + b_l \ln(q^2),$$

$$\Pi_{fit}^{[q]}(q^2) = a_q + b_q \ln(q^2) + c_q \left(\ln(q^2)\right)^2,$$

- \bullet repeat these fits for several fit ranges $\epsilon \in [0.1, 1.0] {\rm GeV}^2$,
- apply cuts to the fits, such as removing fits with a large curvature c_q ,
- from the fits that survive pick the result, where the coefficients b_l and b_q are similar.





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Results for the Adler function on F7: Procedure |



Results for the Adler function on F7: Procedure II



Comparison of the different methods on F7



Comparison of the different methods on F7



Comparison with the mixed representation method



For the mixed representation method see [arXiv:1306.2532], and the talk by Anthony Francis.



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Step 3: Combined fits

To determine the Adler function at the physical point and in the continuum we apply combined fits of the type:

$$D(q^2) = A(q^2)(1 + B(a,q) + C(m_{\pi},q^2))$$

• The momentum dependence is described by the derivative of Padé-approximants

$$A_{[12]}(q^2) = q^2 \left(\frac{a_1 b_1}{(b_1 + q^2)^2} + \frac{a_2 b_2}{(b_2 + q^2)^2} \right),$$

$$A_{[22]}(q^2) = q^2 \left(\frac{a_1 b_1}{(b_1 + q^2)^2} + \frac{a_2 b_2}{(b_2 + q^2)^2} + a_0 \right).$$

• The lattice spacing dependence is given by

$$B_1(a,q) = c_1(aq) + c_2(4\pi f_{\rm K}a), \quad B_2(a,q) = c_1(aq)^2 + c_2(4\pi f_{\rm K}a)^2.$$

• The light quark mass dependence is

$$C(m_{\pi}, q^2) = d_1 \frac{m_{\pi}^2 - (m_{\pi}^{exp})^2}{d_2 + q^2}$$

Results for $m_{\pi} \approx 270$ MeV for P[2,2], $O(a^2)$, $N_f = 2$



Results for $\beta = 5.30$, a = 0.063 fm, $N_f = 2$



Results for $\beta = 5.30$, a = 0.063 fm for P[2,2], $\overline{O(a^2)}$, $N_f = 2$



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Step 4: Determination of a_{μ}^{HLO}

We can use the results from the combined fits to determine a_{μ}^{HLO} :



where we insert the coefficients of the combined chiral and continuum fits of the Adler function to determine the q^2 -behaviour of $\hat{\Pi}(q^2)$:

$$\begin{split} \hat{\Pi}(q^2) \, \to \, \hat{\Pi}_{12}(q^2) &= -q^2 \left(\frac{a_1}{b_1 + q^2} + \frac{a_2}{b_2 + q^2} \right) \text{ or } \\ &\to \, \hat{\Pi}_{22}(q^2) = -q^2 \left(\frac{a_1}{b_1 + q^2} + \frac{a_2}{b_2 + q^2} + a_0 \right). \end{split}$$

Preliminary results for a_{μ}^{HLO} from combined fits, $N_f=2$





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Summary and Outlook

- From the vacuum polarization we can obtain the Adler function with different methods, which agree within errors for a large range of q^2 , and are similar to the mixed representation method, cf. talk by Anthony Francis on Monday (1D).
- We presented a method to extract the Adler function in the continuum and at the physical point using a combined fit for the light quark sector.
- From the extrapolated Adler function we can extract the hadronic contribution to the anomalous magnetic moment of the muon.
- We plan to extend the analysis to the already available strange and charm data.
- E. Shintani is currently investigating AMA for the vacuum polarization.
- In the future we will also investigate methods which make use of the moments of $\Pi(q^2)$ to compute a_{μ}^{HLO} [arXiv:1403.1778, arXiv:1406.4671].

Thank you for your attention.

Backup

Preliminary results for a_{μ}^{HLO} for P[2,2], $N_f = 2$



Low q^2 -region of $\Pi(q^2)$

