

## Lattice simulations of $\mathbf{G}_{\mathbf{2}} \mathbf{- Q C D}$ at finite density $\mathbf{I}$

## $32^{\text {nd }}$ International Symposium on Lattice Field Theory

Lattice 2014


## Lorenz von Smekal



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- $\mathbf{G}_{2}$ Gauge Theory
- Phase Diagram
A. Maas, L.v.S., B. Wellegehausen \& A. Wipf, Phys. Rev. D 86 (2012) 111901(R)
- G2-QCD Spectroscopy \& Baryon Density
B. Wellegehausen, A. Maas, A. Wipf \& L.v.S., Phys. Rev. D 89 (2014) 056007
- Summary and outlook


## Phase Diagram



## QCD-like Theories

- compare lattice simulations with functional methods and effective models where there's no sign problem
- apply to ultracold fermi gases exploit analogies and more experimental data

- strongly correlated fermions in 2+1 dimensions electronic properties of graphene
[see Dominik Smith's talk on Thu, Applications beyond QCD]


## Fermion-Sign Problem

sign problem:

$$
\left(\operatorname{Det} D\left(\mu_{f}\right)\right)^{*}=\operatorname{Det} D\left(-\mu_{f}\right)
$$

- except if:
(a) two degenerate flavors with isospin chemical potential
fermion determinant $\rightsquigarrow \operatorname{Det}\left(D\left(\mu_{I}\right) D\left(-\mu_{I}\right)\right)$

$$
\beta=2
$$

## QCD at finite isospin density

(b) anti-unitary symmetry $\quad T D(\mu) T^{-1}=D(\mu)^{*} \quad T^{2}= \pm 1$
fermion color representation:
(i) pseudo-real $\quad T^{2}=1$
two-color QCD
$\beta=1$
(ii) real $\quad T^{2}=-1 \quad$ adjoint $\mathbf{Q C D}$, or $\mathbf{G}_{2}-\mathbf{Q C D}$
$\beta=4$

## Two-Color QCD

- Polyakov-Quark-Meson-Diquark model phase diagram:
- Lattice simulations:



Strodthoff \& L.v.S., PLB 731 (2014) 350

Can we describe the two-color world with the 3d effective lattice theory for heavy quarks? [cf. Philipp Scior's talk]

Cotter, Giudice, Hands \& Skullerud,
PRD 87 (2013) 034507

## $\mathbf{G}_{2}$ Gauge Theory

- smallest exceptional Lie group subgroup of SO(7)
- rank $=2$ (as SU(3)), dimension $=14$

7 colors, 14 gluons
fund. reps.: $\{7\}=(1,0),\{14\}=(0,1)(=$ adjoint $)$

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- simple \& simply connected, no center yet (as $\operatorname{SU}(3)$ ), $1^{\text {st }}$ order deconfinement finite $T$ phase transition in pure gauge theory (also in chiral condensate)


Holland, Minkowski, Pepe \& Wiese, Nucl. Phys. B 668 (2003) 207 Pepe \& Wiese, NPB 768 (2007) 21
Danzer, Gattringer, Maas, JHEP 01 (2009) 024
Wellegehausen, Wipf \& Wozar, PRD 83 (2011) 114502

## $\mathbf{G}_{2}$ Gauge Theory

- smallest exceptional Lie group subgroup of SO(7)
- rank = 2 (as SU(3)), dimension = 14

7 colors, 14 gluons
fund. reps.: $\{7\}=(1,0),\{14\}=(0,1)(=$ adjoint $)$

- simple \& simply connected, no center yet (as SU(3)), $1^{\text {st }}$ order deconfinement finite $T$ phase transition in pure gauge theory (also in chiral condensate)
- all reps. real

Dirac operator $\mathbf{D}(\mu)$ has antiunitary symmetry $\mathbf{S}$, with $\mathrm{S}^{2}=-1$ (symplectic, $\beta=4$ ) and extended $\operatorname{SU}\left(2 \mathrm{~N}_{\mathrm{f}}\right)$ flavor symmerty


Holland, Minkowski, Pepe \& Wiese, Nucl. Phys. B 668 (2003) 207
Pepe \& Wiese, NPB 768 (2007) 21
Danzer, Gattringer, Maas, JHEP 01 (2009) 024
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- no sign problem
real and positive for single flavor: $\mathrm{SU}(2) \rightarrow \mathbf{U}(1)_{\mathrm{B}}$ 2 Goldstone bosons: scalar (anti)diquarks


## $\mathbf{G}_{2}$ Gauge Theory

## - breaks down to QCD

Higgs

$$
G_{2} \longrightarrow S U(3)
$$

coset:

$$
\begin{aligned}
& G_{2} / S U(3) \sim S O(7) / S O(6) \sim S^{6} \\
&\{7\} \rightarrow\{3\} \oplus\{\overline{3}\} \oplus\{1\} \\
&\{14\} \rightarrow\{3\} \oplus\{\overline{3}\} \oplus\{8\}
\end{aligned}
$$

Wellegehausen, Wipf \& Wozar, PRD 83 (2011) 114502

## $\mathbf{G}_{2}$ Gauge Theory

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Higgs
$G_{2} \longrightarrow S U(3)$
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G_{2} / S U(3) \sim S O(7) / S O(6) \sim S^{6}
$$

$\begin{aligned}\{7\} & \rightarrow\{3\} \oplus\{\overline{3}\} \oplus\{1\}^{<} \text {massive Higgs } \\ \{14\} & \rightarrow\{3\} \oplus\{\overline{3}\} \rightarrow\{8\}^{\text {\& }} \text { _gluons }\end{aligned}$

## $\mathbf{G}_{2}$ Gauge Theory

## - breaks down to QCD

Higgs

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$$

| $\{7\}$ | $\rightarrow\{3\} \oplus\{\overline{3}\}$ |
| ---: | :--- |
| $\{14\}$ | $\rightarrow\{1\}^{\text {massive Higgs }}$ |
| heavy gauge | $\{3\} \oplus\{\overline{3}\}$ |

- $\mathrm{G}_{2}$ glueball sectrum, Casimir scaling \& string breaking

Pepe \& Wiese, NPB 768 (2007) 21
Wellegehausen et al., PRD 83 (2011) 016001
Lacroix et al., PRD 87 (2013) 054025

- vortices, monopoles, instantons...

Greensite et al., PRD 75 (2007) 034501
Di Giacomo et al., JHEP 10 (2008) 096
Ilgenfritz \& Maas, PRD 86 (2012) 114508


Wellegehausen, Wipf \& Wozar, PRD 83 (2011) 114502


## $\mathbf{G}_{2}-\mathbf{Q C D}$ at Finite Density


$U(1)_{B}$ breaks spontaneously at $\mu_{B}=m_{d_{0}^{+}}=m_{\pi}$

- diquark condensation as in $\mathrm{QC}_{2} \mathrm{D}$



Bjoern Wellegehausen, PhD thesis, Jena 2012

## $\mathbf{G}_{2}$ Gauge Theory at Finite Density

- but has fermionic baryons also (as adjoint QCD, in principle)
- finite baryon density (bosonic and fermionic)

Polyakov loop

quark condensate

baryon density

$G_{2}$ nuclear matter?

Maas, LvS, Wellegehausen \& Wipf, Phys. Rev. D 86 (2012) 111901R

## G $_{2}$-QCD Phase Diagram

## - 1 flavor dynamcial Wilson



Maas, LvS, Wellegehausen \& Wipf, Phys. Rev. D 86 (2012) 111901R

## $\mathbf{G}_{2}$ Spectroscopy

$$
\begin{aligned}
\{7\} \otimes\{7\} & =\{1\} \oplus\{7\} \oplus\{14\} \oplus\{27\} \\
\{7\} \otimes\{7\} \otimes\{7\} & =\{1\} \oplus 4 \cdot\{7\} \oplus 2 \cdot\{14\} \oplus \ldots \\
\{14\} \otimes\{14\} & =\{1\} \oplus\{14\} \oplus\{27\} \oplus \ldots, \\
\{14\} \otimes\{14\} \otimes\{14\} & =\{1\} \oplus\{7\} \oplus 5 \cdot\{14\} \oplus \ldots, \\
\{7\} \otimes\{14\} \otimes\{14\} & =\{1\} \oplus \ldots
\end{aligned}
$$

mesons (baryon number 0)

| Name | $\mathcal{O}$ | $T$ | J | P | C |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\pi$ | $\bar{u} \gamma_{5} d$ | SASS | 0 | - | + |
| $\eta$ | $\bar{u} \gamma_{5} u$ | SASS | 0 | - | + |
| $a$ | $\bar{u} d$ | SASS | 0 | + | + |
| $f$ | $\bar{u} u$ | SASS | 0 | + | + |
| $\rho$ | $\bar{u} \gamma_{\mu} d$ | SSSA | 1 | - | + |
| $\omega$ | $\bar{u} \gamma_{\mu} u$ | SSSA | 1 | - | + |
| $b$ | $\bar{u} \gamma_{5} \gamma_{\mu} d$ | SSSA | 1 | + | + |
| $h$ | $\bar{u} \gamma_{5} \gamma_{\mu} u$ | SSSA | 1 | + | + |

diquarks (baryon number 2)

| Name | $\mathcal{O}$ | $T$ | $J$ | $P$ | $C$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $d\left(0^{++}\right)$ | $\bar{u}^{\mathrm{C}} \gamma_{5} u+$ c.c. | SASS | 0 | + | + |
| $d\left(0^{+-}\right)$ | $\bar{u}^{\mathrm{C}} 5_{5} u-$ c.c. | SASS | 0 | + | - |
| $d\left(0^{-+}\right)$ | $\bar{u}^{\mathrm{C}} u+$ c.c. | SASS | 0 | - | + |
| $d\left(0^{--}\right)$ | $\bar{u}^{\mathrm{C}} u-$ c.c. | SASS | 0 | - | - |
| $d\left(1^{1+}\right)$ | $\bar{u}^{\mathrm{C}} \gamma_{\mu} d-\bar{d}^{\mathrm{C}} \gamma_{\mu} u+$ c.c. | SSSA | 1 | + | + |
| $d\left(1^{+-}\right)$ | $\bar{u}^{\mathrm{C}} \gamma_{\mu} d-\bar{d}^{\mathrm{C}} \gamma_{\mu} u-$ c.c. | SSSA | 1 | + | - |
| $d\left(1^{-+}\right)$ | $\bar{u}^{\mathrm{C}} \gamma_{5} \gamma_{\mu} d-\bar{d}^{\mathrm{C}} \gamma_{5} \gamma_{\mu} u+$ c.c. | SSSA | 1 | - | + |
| $d\left(1^{--}\right)$ | $\bar{u}^{\mathrm{C}} \gamma_{5} \gamma_{\mu} d-\bar{d}^{\mathrm{C}} \gamma_{5} \gamma_{\mu} u-$ c.c. | SSSA | 1 | - | - |

exotic particles (baryon number 1)

| Name | $\mathcal{O}$ | $T$ | J | P | C |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $N^{\prime}$ | $T^{a b c}\left(\bar{u}_{a} \gamma_{5} d_{b}\right) u_{c}$ | SAAA | $1 / 2$ | $\pm$ | $\pm$ |
| $\Delta^{\prime}$ | $T^{a b c}\left(\bar{u}_{a} \gamma_{\mu} u_{b}\right) u_{c}$ | SSAS | $3 / 2$ | $\pm$ | $\pm$ |
| Hybrid | $\epsilon_{a b c d e f g} u^{a} F_{\mu \nu}^{b c} F_{\mu \nu}^{d e} F_{\mu \nu}^{f g}$ | SSSS | $1 / 2$ | $\pm$ | $\pm$ |

baryons (baryon number 3)

| Name | $\mathcal{O}$ | $T$ | $J$ | P | C |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $N$ | $T^{a b c}\left(\bar{u}_{a}^{C} \gamma_{5} d_{b}\right) u_{c}$ | SAAA | $1 / 2$ | $\pm$ | $\pm$ |
| $\Delta$ | $T^{a b c}\left(\bar{u}_{a}^{C} \gamma_{\mu} u_{b}\right) u_{c}$ | SSAS | $3 / 2$ | $\pm$ | $\pm$ |

## $\mathbf{G}_{2}$ Spectroscopy

- $\boldsymbol{N}_{f}=1$ : real and positive for single flavor: $S U(2) \rightarrow U_{B}(1)$

2 Goldstone bosons: scalar (anti)diquarks


Wellegehausen, Maas, Wipf \& LvS, PRD 89 (2014) 056007

## $\mathbf{G}_{2}$ Spectroscopy



Wellegehausen, Maas, Wipf \& LvS, PRD 89 (2014) 056007

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## Finite Baryon Density

heavy ensemble


Wellegehausen, Maas, Wipf \& LvS, PRD 89 (2014) 056007

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## Finite Baryon Density



Wellegehausen, Maas, Wipf \& LvS, PRD 89 (2014) 056007

## Finite Baryon Density

heavy ensemble

light ensemble

$$
n_{q} a^{3}
$$



Wellegehausen, Maas, Wipf \& LvS, PRD 89 (2014) 056007

## Finite Baryon Density


light ensemble

$1^{\text {st }}$ order?
liquid-gas transition of G2 nuclear matter?

Wellegehausen, Maas, Wipf \& LvS, PRD 89 (2014) 056007

## Summary \& Outlook

- $\mathbf{G}_{2}$-QCD, a useful laboratory for finite density studies
- no sign problem, most QCD-like
- finite baryon density region of phase diagram with MC simulations
- refine functional methods and models
- test effective lattice theories for heavy quarks
- spectroscopy \& baryon density
- physics of bosonic baryons as in two-color QCD
- fermionic baryons dominate above $\mathrm{G}_{2}$-nuclear matter transition
- further clarify nature of cold and dense phases
[to be continued in Bjoern Wellegehausen's talk next]



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## Thank You for Your Attention!

