



Error reduction with all-mode-averaging in Wilson fermion

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1. Introduction

Lattice QCD is powerful tool to deal with the strong interaction between quark and gluon, however, in order to give a reliable solution, we need to reduce the statistical fluctuation as much as possible. Recently idea of all-mode-averaging (AMA) [1,2] is stateof-the-art algorithm to reduce statistical error of correlation function in Monte-Carlo simulation, and it seems to be broadly applicable. In this poster we present a performance test of AMA using Wilson-Clover fermion.

2. All-mode-averaging (AMA)

The improved estimator is defined as

4. Test of covariance in SAP approximation

The covariance is easily checked by consistency test with gauge shift and source shift,

 $O[U](x + \mu, y + \mu) = O[U^g](x, y), \ U^g(x) = U(x - \mu)$

The precision of discrepancy should be below machine precision (~10⁻¹⁶ in double) if the translational symmetry is preserved. We checked same block shift has good consistency.

Table 1: 64 \times 32³ lattice with 4 \times 4³ SAP domain in ε ~0.01 residue

		Other block shift	(4,0,0,0)	
		Gauge shift	Source shift	diff.
	<psps>,</psps>	2.0810882153191448e-01	2.0814053466832835e-01	1.5e-4
	t=1			

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \ \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

where $O^{(appx)}$ is approximation whose cost is much smaller than O. g denotes the lattice transformation of the symmetry G. Here the translational invariance is employed. Using deflation method[2,3], the approximation is defined as the combination of deflation field and truncated solver as

$$\mathcal{O}^{(\text{appx})} = \mathcal{O}^{(\text{AMA})}[S^{(\text{all})}], S^{(\text{all})}(x,y) = \sum_{k,l}^{N_{\lambda}} \Lambda_{kl}\psi_k(x)\psi_l(y) + f_{\varepsilon}(D(x,y))$$

where two parameters, N_{λ} and ε , control the quality of approximation and computational cost.

The standard deviation of O^(1mp) expects to be

$$\begin{split} \frac{\sigma^{\mathrm{imp}}}{\sigma} &\simeq \sqrt{\frac{1}{N_G} + 2\Delta r + R^{\mathrm{corr}}}, \ R^{\mathrm{corr}} = \frac{1}{N_g^2} \sum_{g \neq g'} r_{gg'}, \ \Delta r = 1 - r \\ \text{where with } \Delta \mathrm{O} &= \mathrm{O} - <\!\mathrm{O}\!\!> \text{we have} \\ r &= \frac{\langle \Delta \mathcal{O} \Delta \mathcal{O}^{(\mathrm{appx})} \rangle}{\sigma \sigma^{(\mathrm{appx})}}, \ r_{gg'} = \frac{\langle \Delta \mathcal{O}^{(\mathrm{appx}),g} \Delta \mathcal{O}^{(\mathrm{appx}),g'} \rangle}{\sigma^{(\mathrm{appx}),g} \sigma^{(\mathrm{appx}),g'}}, \end{split}$$
these quantities denote the correlation between O and

<NN>, 1.9793954905488386e-07 | 1.9401235321009282e-07 | 2.0e-2 t=10

	Same block shift	(4,4,4,4)	
<psps>,</psps>	1.9670606430843440e-01	1.9670606430843440e-01	<1e-16
t=1			
<nn>,</nn>	3.5454273713838414e-07	3.5454273713838419e-07	1.4e-16
t=10			

5. Approximation with SAP+deflation

The left- and right-handed deflation [3] is given as, $P_L = 1 - \sum (D\phi_k)\phi_l^{\dagger} A_{kl}^{-1}, \quad P_R = 1 - \sum \phi_k (D\phi_l)^{\dagger} A_{kl}^{-1}, \quad A_{kl} = (\phi_k, D\phi_l),$ where decomposed matrix A is called as "little Dirac op.". The deflation procedure for solve Dx = b (D is Hermitian) is 1. $P_L Dx = P_L b \Rightarrow D P_R x = P_L b$, 2. $P_R x = D^{-1} P_L b = x - \sum_{k,l} \phi_k (D\phi_l, x) A_{kl}^{-1} = x - \sum_{k,l} \phi_k (\phi_l, b) A_{kl}^{-1}$ 3. $x = D^{-1}P_Lb + \sum \phi_k(\phi_l, b)A_{kl}^{-1}$ The deflation subspace is given from N_s fields generated from

O^(appx), and different g of O^(appx). For error reduction, we need to seach the approximation with small Δr and R^{corr} .

3. Approximation with SAP method

Schwartz alternative procedure (SAP) [2] is applied for both generation of deflation field and preconditioning. SAP is approximation to Wilson-Dirac kernel as decomposing lattice into domain Λ and Λ^* in which <u>Dirichlet BC</u> is utilized.

$$D_{w} \simeq \begin{pmatrix} D_{\Lambda} & \partial D_{\Lambda} \\ \partial D_{\Lambda^{*}} & D_{\Lambda^{*}} \end{pmatrix} \qquad M_{sap} = K \sum_{\nu=0}^{n_{cy}-1} (1 - DK)^{\nu}$$
$$K = R_{\Lambda}^{T} D_{\Lambda}^{-1} R_{\Lambda} + R_{\Lambda^{*}}^{T} D_{\Lambda^{*}}^{-1} R_{\Lambda^{*}} - R_{\Lambda^{*}}^{T} D_{\Lambda^{*}}^{-1} D_{\partial\Lambda^{*}} D_{\Lambda^{*}}^{-1} R_{\Lambda}$$
$$R_{\Lambda} = \begin{pmatrix} 1 & 0 \end{pmatrix}, \qquad R_{\Lambda^{*}} = \begin{pmatrix} 0 & 1 \end{pmatrix}$$

When we use SAP as an approximation $O^{(appx)}$ in AMA, translational invariance between domains is taken into account:

- 1. Shift inside domain: NOT translational invariance.
- 2. Shift other domian: NOT translational invariance.
- 3. Shift same domain and same local position: OK.

smoothing process. Deflation projection with random field in the SAP domain decomposition is given as

$$P = \sum_{\Lambda,\Lambda'}^{N_b} \sum_{kl}^{N_s} R_{\Lambda}^T \phi_k^{\Lambda} \phi_l^{\Lambda' \dagger} R_{\Lambda} (A_{kl}^{\Lambda\Lambda'})^{-1}, \quad A_{kl}^{\Lambda\Lambda'} = (\phi_k^{\Lambda}, R_{\Lambda}^T D R_{\Lambda'} \phi_l^{\Lambda'})$$

The quality of approximation can be controlled by N_s and SAP domain size in addition to stopping condition ε .

5. Performance test of AMA in SAP+deflation

In this test we use DDHMC library package and GCR algorithm in solver part. Here the mixed precision method is also applied. $N_f = 2$ Wilson-Clover fermion of CLS is used.

• 64x32³ lattice at *a*=0.06 fm in 451 MeV pion.



It is safe to assure covariance of approximation that source location is set to the next domain Λ and same local position.



Original source

location

References: [1] T. Blum, T. Izubuchi, E. Shintani, PRD88.094503 (2013), 1402.0244 [hep-lat]. [2] M. Luscher, Comp.Phys.Comm.156,209 (2004). [3] M.Luscher, JHEP 07,081(2007).