

The electroweak transition and the equation of state in the $SU(2)$ -Higgs-model

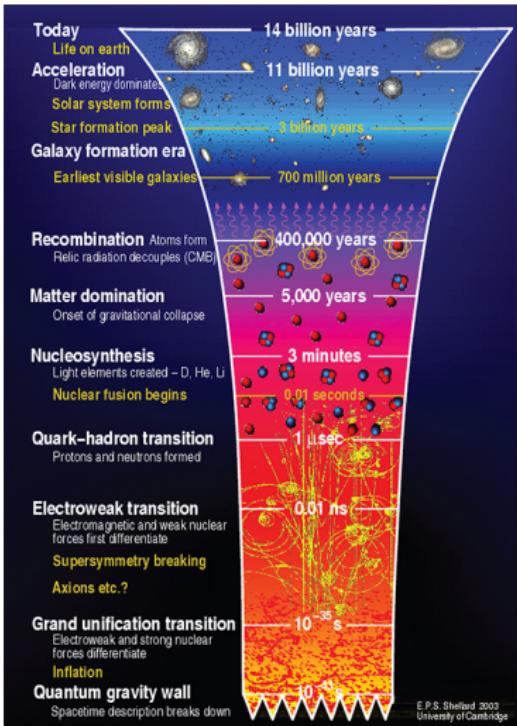
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Introduction

- ▶ with EoS the temporal development of the Universe can be described
- ▶ until 2012 calculation impossible, due to unknown m_H
- ▶ here first attempts to calculate the EoS in $SU(2)$ -Higgs-model



The $SU(2)$ -Higgs-model

Algorithms and observables

Line of constant physics

Equation of state

Work in progress

The transition

- ▶ in 1996 by Kajantie, Laine, Rummukainen and Shaposhnikov:
- ▶ for $m_H \lesssim m_W$ first order phase transition
- ▶ for $m_H \gtrsim m_W$ cross over
- ▶ in 1998 by Aoki, Fodor, Csikor . . . :
- ▶ critical endpoint (66.5 ± 1.4) GeV
- ▶ physical:
 - ▶ $m_W = 80.385$ GeV
 - ▶ $m_H = 125.9$ GeV
 - ▶ $m_H > m_W \longrightarrow$ cross over

The $SU(2)$ -Higgs-model

- ▶ electroweak interaction: $SU(2) \times U(1)$
- ▶ $SU(2)$ -Higgs-model $U(1)$ degrees of freedom are integrated out
- ▶ does not describe: strong interaction, fermions, $U(1)$ -gauge-fields

The action

The continuum action:

$$S_K = \int d^4x \frac{1}{4} \text{tr} F_{\mu\nu} F^{\mu\nu} + D_\mu \phi D^\mu \phi + m_K \phi_\mu \phi^\mu + \lambda_K (\phi_\mu \phi^\mu)^2$$

discretisation on the lattice:

$$\int d^4x D_\mu \phi D^\mu \phi \longrightarrow \sum_{x \in \Lambda} \sum_{\mu=1}^4 (U_{x\mu} \phi_{x+\hat{\mu}} - \phi_x)^\dagger (U_{x\mu} \phi_{x+\hat{\mu}} - \phi_x)$$

$$\int d^4x m_K^2 \phi_\mu \phi^\mu + \lambda_K (\phi_\mu \phi^\mu)^2 \longrightarrow \sum_{x \in \Lambda} m_K^2 \phi_x^\dagger \phi_x + \lambda_K (\phi_x^\dagger \phi_x)^2$$

Reparametrisation I

$$S_S = \sum_{x \in \Lambda} \sum_{\mu=1}^4 (U_{x\mu} \phi_{x+\hat{\mu}} - \phi_x)^\dagger (U_{x\mu} \phi_{x+\hat{\mu}} - \phi_x) + m_K^2 \phi_x^\dagger \phi_x + \lambda_K (\phi_x^\dagger \phi_x)^2.$$

$$\begin{aligned}\lambda_K &= \frac{\lambda}{\kappa} \\ m_K^2 &= \frac{1 - 2\lambda}{\kappa} - 2 \\ \phi_x &= \sqrt{\kappa} \Phi_x.\end{aligned}$$

$$S_S = \sum_{x \in \Lambda} \Phi_x^\dagger \Phi_x + \lambda (\Phi_x^\dagger \Phi_x - 1)^2 - 2\kappa \sum_{\mu=1}^4 \Re (\Phi_x^\dagger U_{x\mu} \Phi_{x+\hat{\mu}}).$$

Reparametrisation II

$$S_S = \sum_{x \in \Lambda} \Phi_x^\dagger \Phi_x + \lambda (\Phi_x^\dagger \Phi_x - 1)^2 - 2\kappa \sum_{\mu=1}^4 \Re(\Phi_x^\dagger U_{x\mu} \Phi_{x+\hat{\mu}}).$$

writing the scalar field as matrix:

$$\Phi = \begin{pmatrix} \Phi_1 \\ \Phi_2 \end{pmatrix} \longrightarrow \begin{pmatrix} \Phi_2^* & \Phi_1 \\ -\Phi_1^* & \Phi_2 \end{pmatrix} = \varphi$$

the full lattice action:

$$S[U, \varphi] = \beta \sum_{pl} \left(1 - \frac{1}{2} \text{tr } U_{pl} \right) + \sum_x \left(\frac{1}{2} \text{tr}(\varphi_x^\dagger \varphi_x) + \lambda \left(\frac{1}{2} \text{tr}(\varphi_x^\dagger \varphi_x) - 1 \right)^2 - \kappa \sum_{\mu=1}^4 \text{tr}(\varphi_x^\dagger U_{x\mu} \varphi_{x+\mu}) \right).$$

Used Algorithms

For simulations on GPUs:

- ▶ Heatbath-Algorithm
- ▶ Overrelaxation-Algorithm

For comparison on CPUs:

- ▶ HMC

Successfull comparison with

- ▶ Simulating the electroweak phase transition in the $SU(2)$ higgs model
- ▶ by Z. Fodor, J. Hein, K. Jansen, A. Jaster and I. Montvay
- ▶ published 1995 in Nucl. Phys. B439

Simple observables

$$S[U, \varphi] = \beta \sum_{pl} \left(1 - \frac{1}{2} \operatorname{tr} U_{pl} \right) + \sum_x \left(\frac{1}{2} \operatorname{tr}(\varphi_x^\dagger \varphi_x) + \lambda \left(\frac{1}{2} \operatorname{tr}(\varphi_x^\dagger \varphi_x) - 1 \right)^2 - \kappa \sum_{\mu=1}^4 \operatorname{tr}(\varphi_x^\dagger U_{x\mu} \varphi_{x+\mu}) \right)$$

The observables:

$$R_x = \det \varphi_x = \frac{1}{2} \operatorname{tr} (\varphi_x^\dagger \varphi_x) = \rho_x^2$$

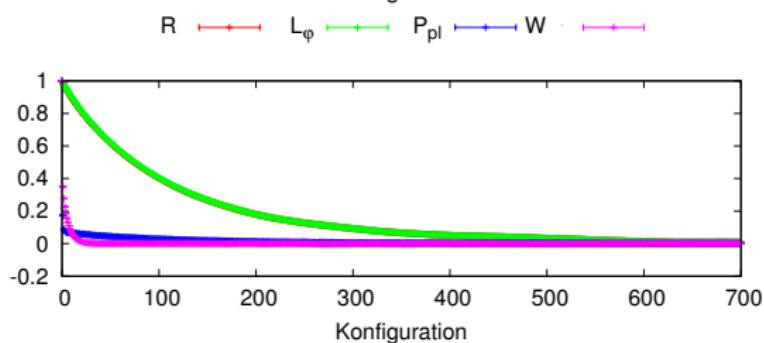
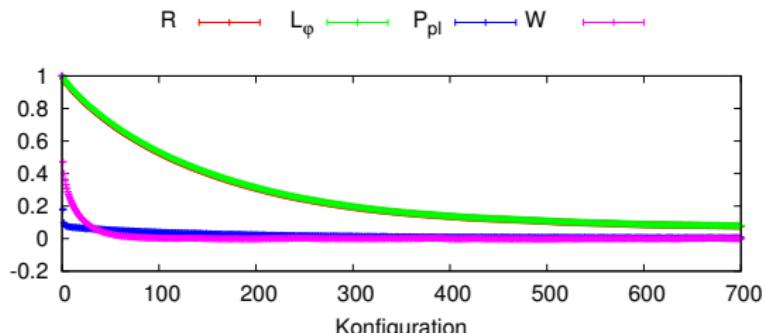
$$L_{\varphi, x\mu} = \frac{1}{2} \operatorname{tr} (\varphi_x^\dagger U_{x\mu} \varphi_{x+\hat{\mu}})$$

$$P_{Pl} = 1 - \frac{1}{2} \operatorname{tr} U_{pl}$$

$$Q_x = (\rho_x^2 - 1)^2$$

$$S_x = 6\beta P_{Pl} + R_x + \lambda Q_x - 8\kappa L_{\varphi, x\mu}$$

Autocorrelation time without and with overrelaxation



Correlators

Used correlators for the higgs boson mass:

$$R_x = \det \varphi_x$$

$$L_{\varphi,x\mu} = \frac{1}{2} \text{tr} (\varphi_x^\dagger U_{x\mu} \varphi_{x+\hat{\mu}})$$

Used correlators for the W boson mass:

$$W_{xrk}^{(n)} = \frac{1}{2} \text{tr} (\sigma_r \varphi_x^\dagger U_{xk} U_{x+\hat{k},k} U_{x+2\hat{k},k} \dots U_{x+(n-1)\hat{k},k} \varphi_{x+n\hat{k}})$$

(Bunk, Ilgenfritz, Kripfganz, Schiller: The finite-temprature phase transition in lattice $SU(2)$ Higgs theory at weak couplings, Nucl. Phys. B403 (1993))

Renormalized coupling

The renormalized coupling

$$g_R^2 = \frac{16}{3}\pi A$$

can be calculated from the static potential

$$V(R) = C - \frac{A}{R}e^{-MR} + DG(M, R)$$

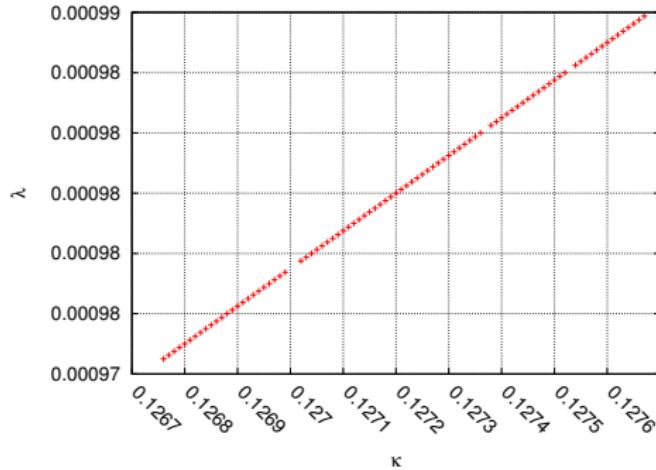
- ▶ A, C, D, M are fit parameter (M is screening mass)
- ▶ $G(M, R)$ are lattice corrections (1986 by Langguth, Montvay, Weisz)

The static potential can be calculated from the Wilson-Loops

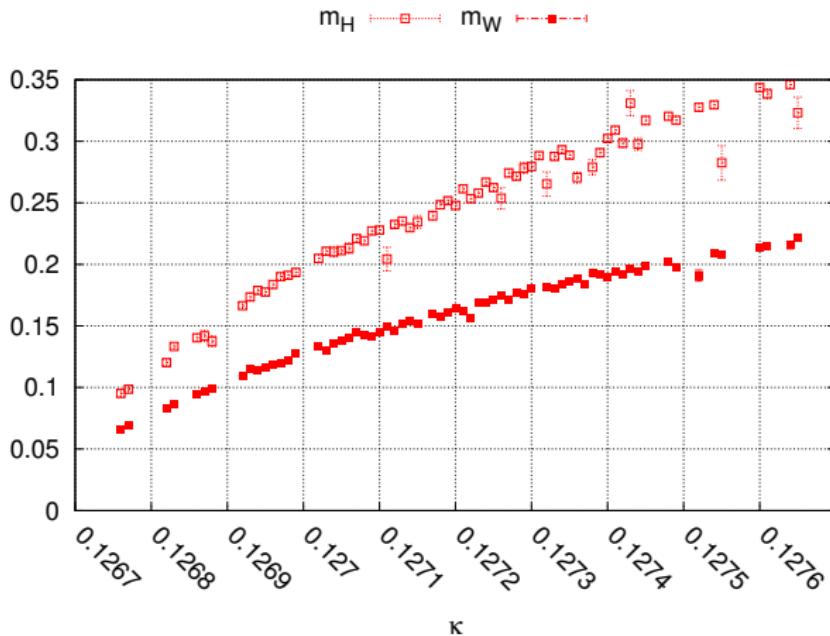
$$V(R) = -\lim_{T \rightarrow \infty} \frac{1}{T} \ln \mathcal{W}(R, T)$$

Line of constant physics

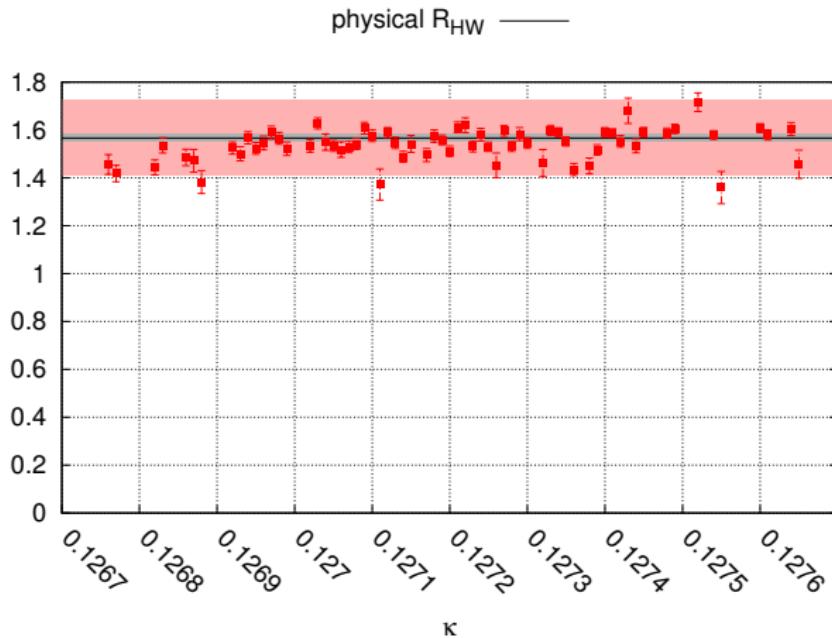
- ▶ Lattice: 32^4
- ▶ β is constant
- ▶ λ depends linear on κ



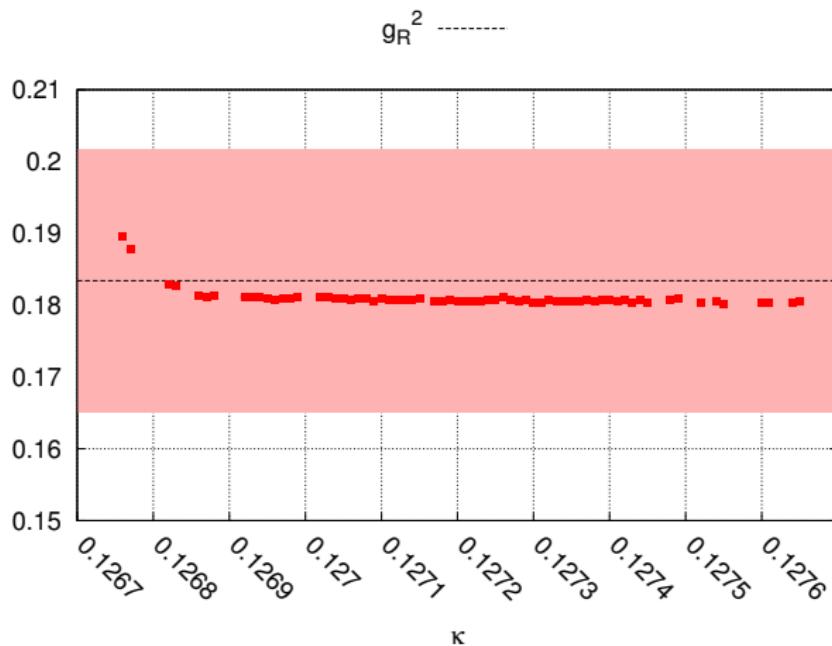
Masses



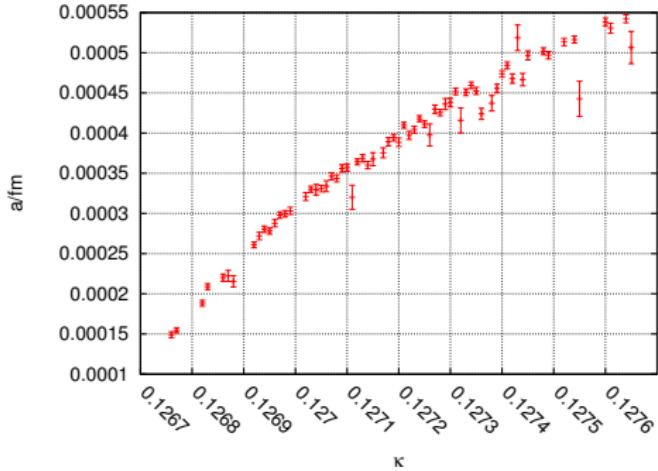
$$R_{HW} = \frac{m_H}{m_W}$$



g_R^2

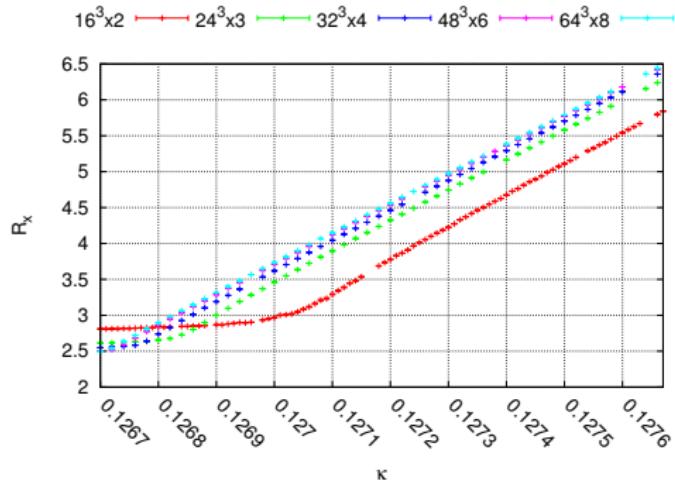


The scale



Electroweak scale is 0.0008 fm (246 GeV)

The transition



$$R_x = \det \varphi_x = \frac{1}{2} \text{tr} (\varphi_x^\dagger \varphi_x) = \rho_x^2$$

Equation of state

- ▶ Interaction measure:

$$\begin{aligned} I &= -\frac{N_t^3}{N_s^3} m_H \left(\frac{\partial \ln Z}{\partial \beta} \frac{\partial \beta}{\partial \kappa} \frac{d\kappa}{dm_H} + \frac{\partial \ln Z}{\partial \lambda} \frac{\partial \lambda}{\partial \kappa} \frac{d\kappa}{dm_H} + \frac{\partial \ln Z}{\partial \kappa} \frac{d\kappa}{dm_H} \right) \\ &= -N_t^4 m_H \left(-6\langle P_{\text{PI}} \rangle \frac{\partial \beta}{\partial \kappa} \frac{\partial \kappa}{\partial m_H} - \langle Q \rangle \frac{\partial \lambda}{\partial \kappa} \frac{\partial \kappa}{\partial m_H} + 8\langle L_\varphi \rangle \frac{\partial \kappa}{\partial m_H} \right) \end{aligned}$$

- ▶ Pressure:

$$\begin{aligned} \frac{p}{T^4} &= N_t^4 \int_{(\kappa_0, \beta_0, \lambda_0)}^{(\kappa, \beta, \lambda)} d(\kappa, \beta, \lambda) \left(\frac{1}{N_t N_s^3} \begin{pmatrix} \frac{\partial \ln Z}{\partial \kappa} \\ \frac{\partial \ln Z}{\partial \beta} \\ \frac{\partial \ln Z}{\partial \lambda} \end{pmatrix} - \frac{1}{N_{s0}^3 N_{t0}} \begin{pmatrix} \frac{\partial \ln Z_0}{\partial \kappa} \\ \frac{\partial \ln Z_0}{\partial \beta} \\ \frac{\partial \ln Z_0}{\partial \lambda} \end{pmatrix} \right) \\ &= N_t^4 \sum -6\Delta\beta\langle P_{\text{PI}} \rangle - \Delta\lambda\langle Q \rangle + 8\Delta\kappa\langle L_\varphi \rangle \end{aligned}$$

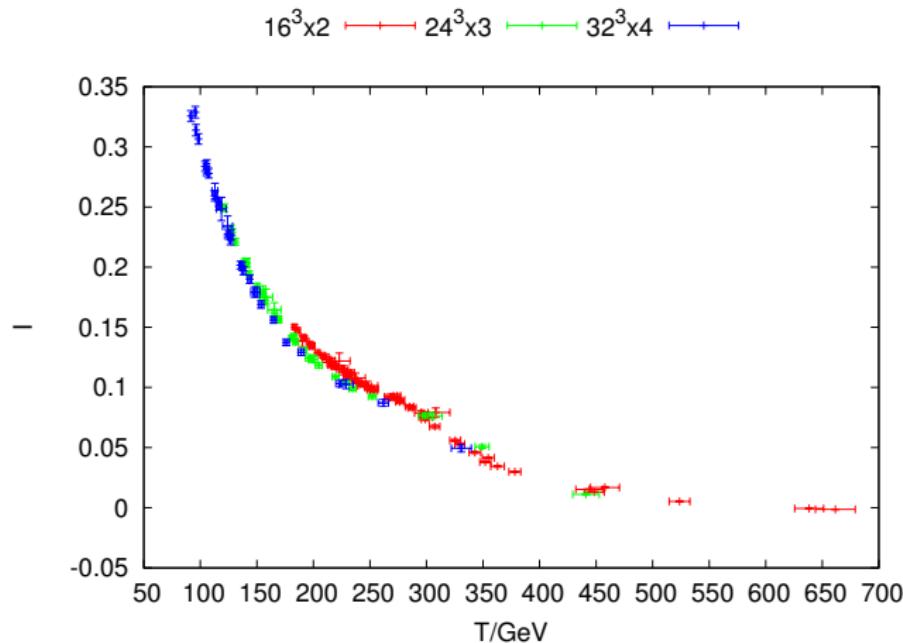
- ▶ Energy density:

$$\frac{\epsilon}{T^4} = I + \frac{3p}{T^4}$$

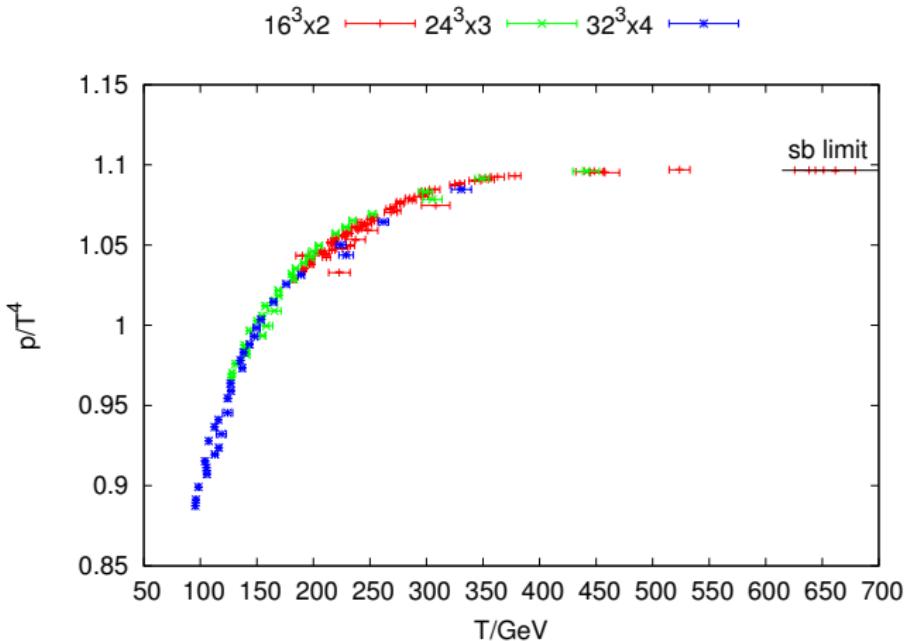
- ▶ Entropy density:

$$\frac{s}{T^3} = \frac{\epsilon}{T^4} + \frac{p}{T^4}$$

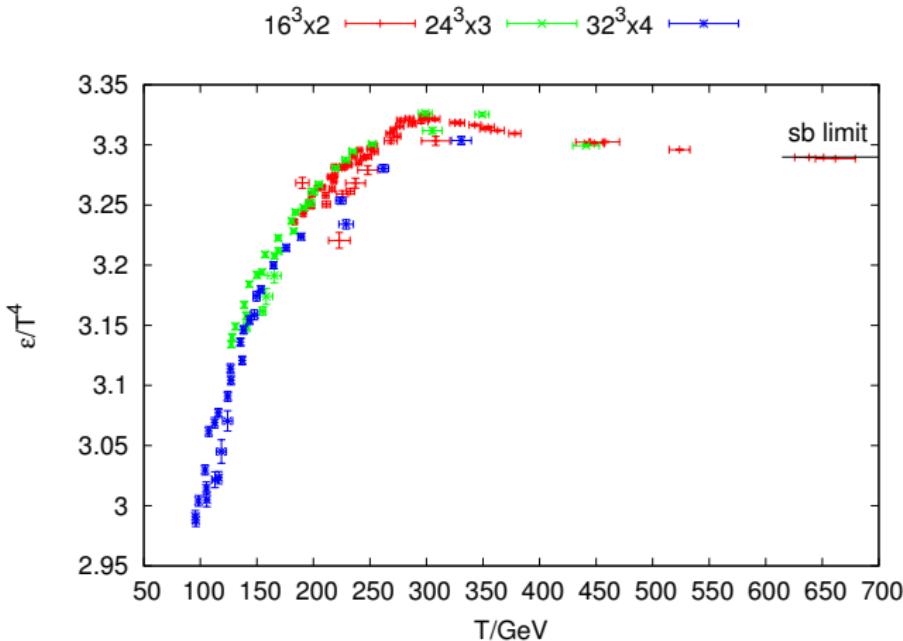
Interaction measure I



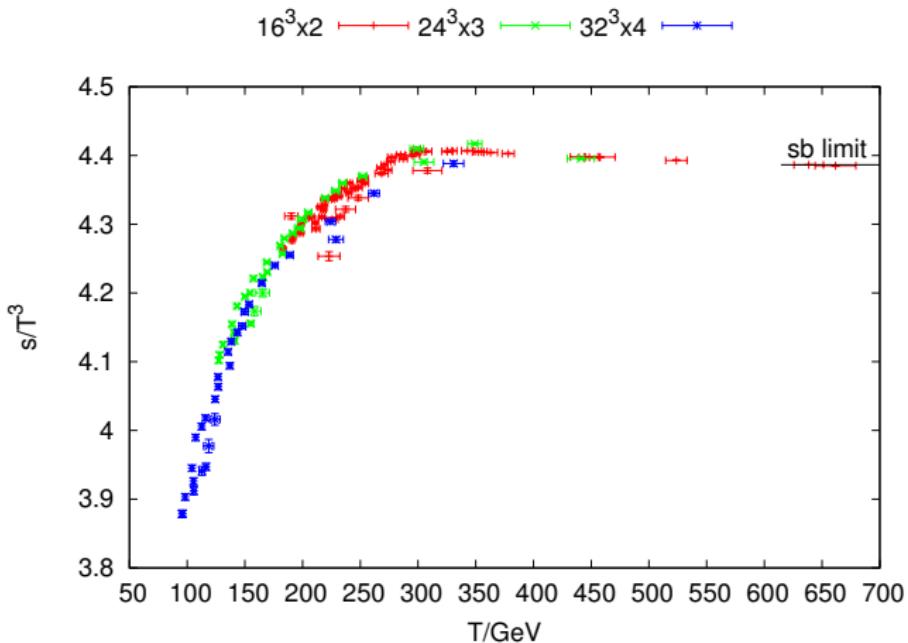
Pressure p



Energy density ϵ



Entropy density s



Work in progress

- ▶ Larger lattices will hopefully make it possible to get to smaller masses.
- ▶ First results from tuning on 48^4 lattices with GPUs.
- ▶ Blue Gene/Q program is ready for larger lattices.

