

Testing Composite Higgs models on the lattice

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Lattice Strong Dynamics collaboration

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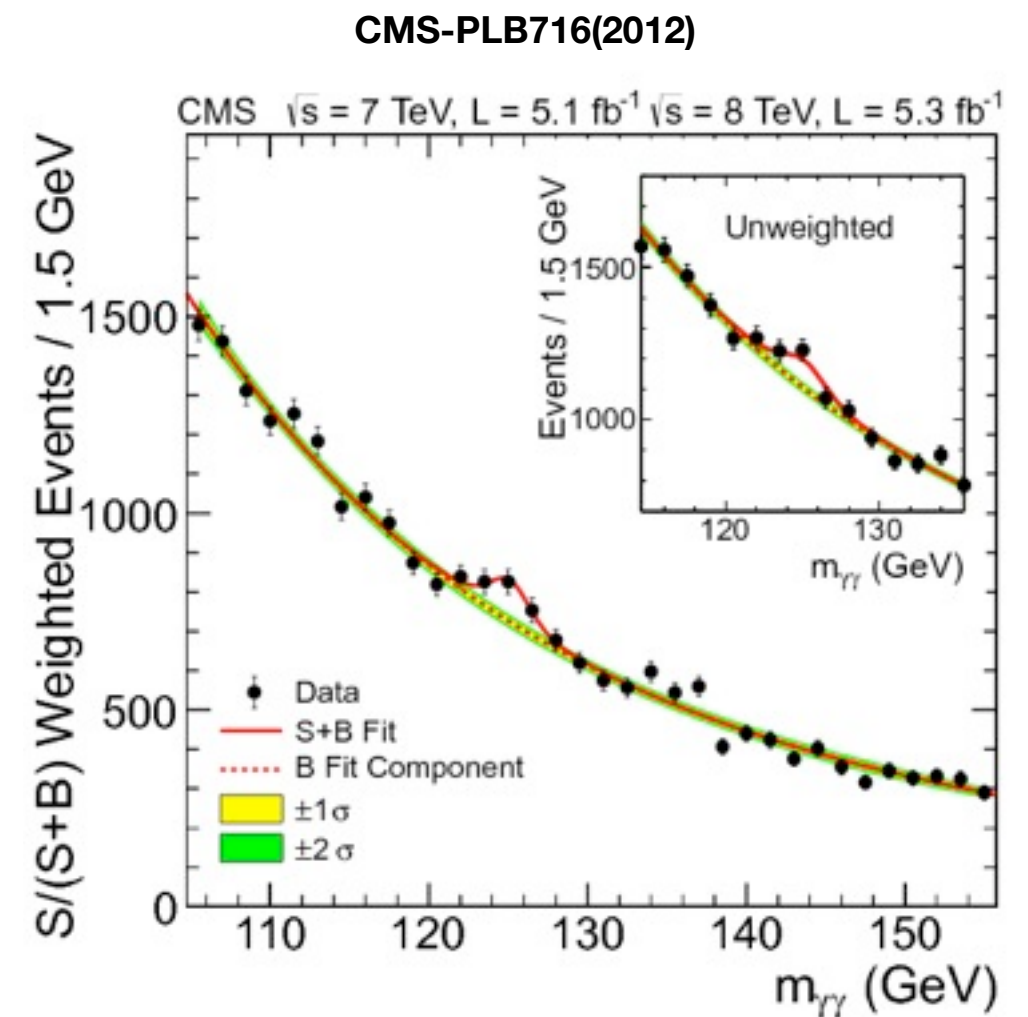
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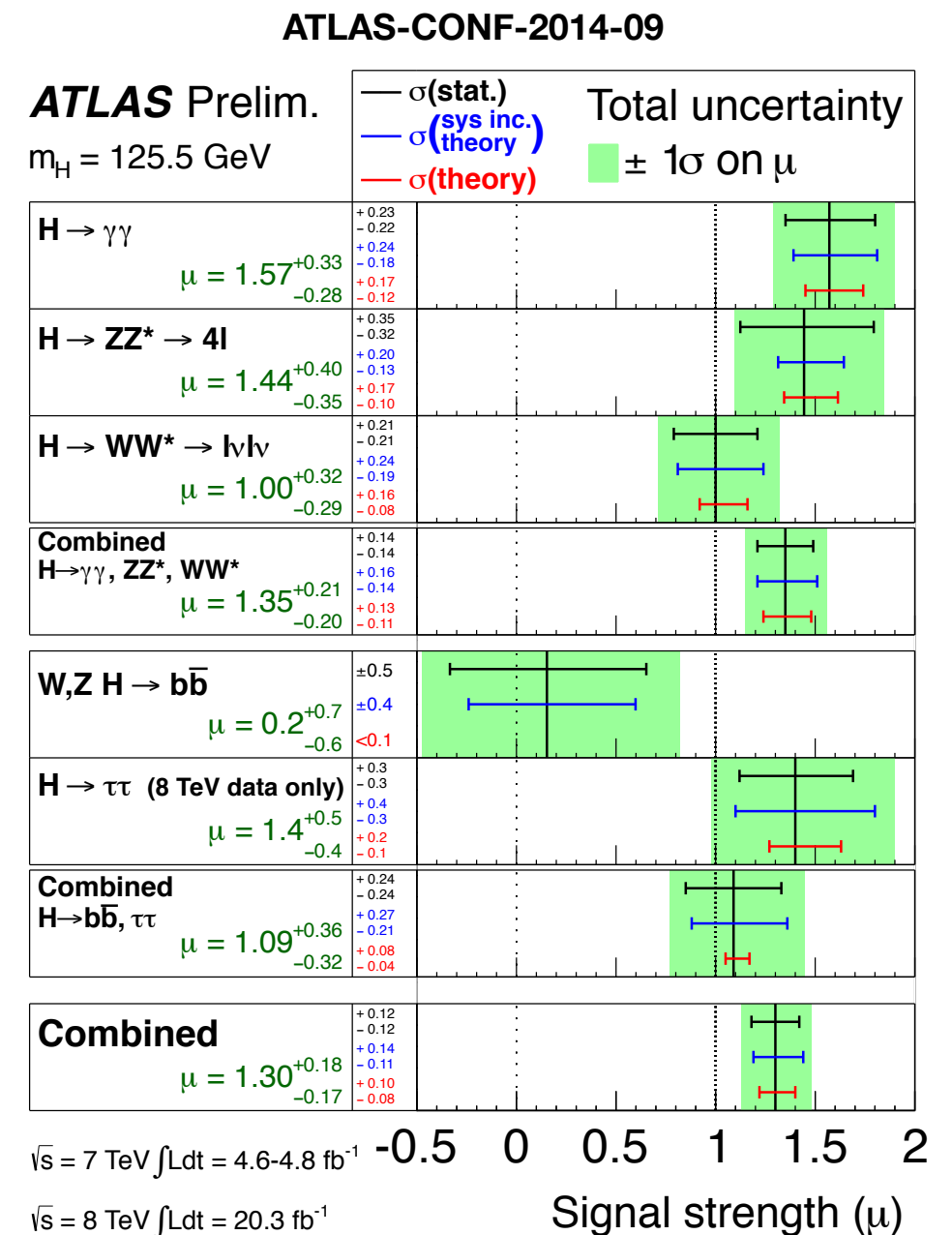
The SM Higgs boson

- Higgs boson discovery in Run I at LHC
➡ $m_H = 125.5 \text{ GeV}$
- Higgs boson decays to ➡ Gauge bosons and leptons+quarks
- Signal strength measurements ➡ coupling constant measurements
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compositeness of the Higgs boson

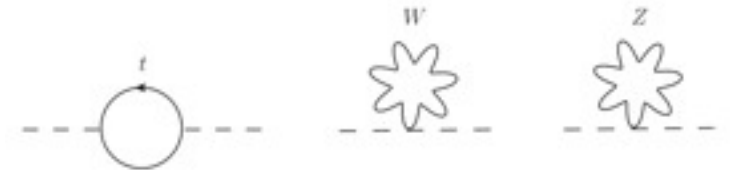
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Composite Higgs models

- the Higgs boson is not an **elementary** scalar particle ➡ **composite** bound state of new strong dynamics
- Technicolor** Higgs:
 - ✓ the new sector breaks the EW symmetry through a technifermion condensate
 - ✓ the Higgs is identified with the lightest scalar excitation of the condensate
 - ✓ can be light due to interactions with SM particles (obtained from ETC dynamics)

[Foadi et al., PRD87(095001)] [Di Chiara et al., arxiv:1405.7154]



- Walking Technicolor** Higgs:
 - ✓ walking coupling and large anomalous mass dimension $\gamma \approx 1$
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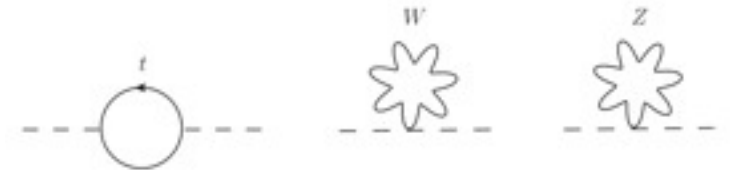
[Yamawaki et al., PRL56(1986)] [Bando et al., PLB178(1986)]

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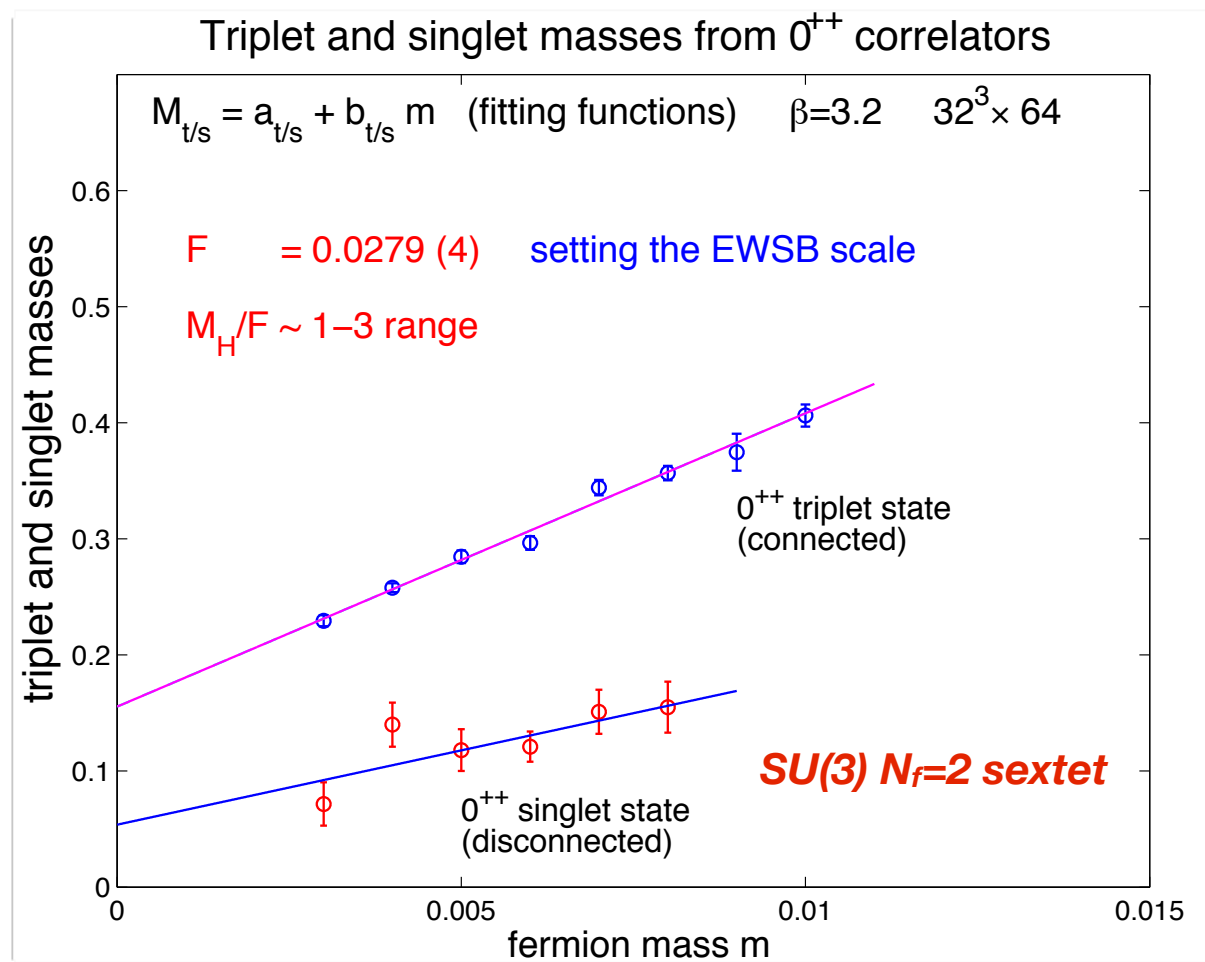
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Lattice simulations

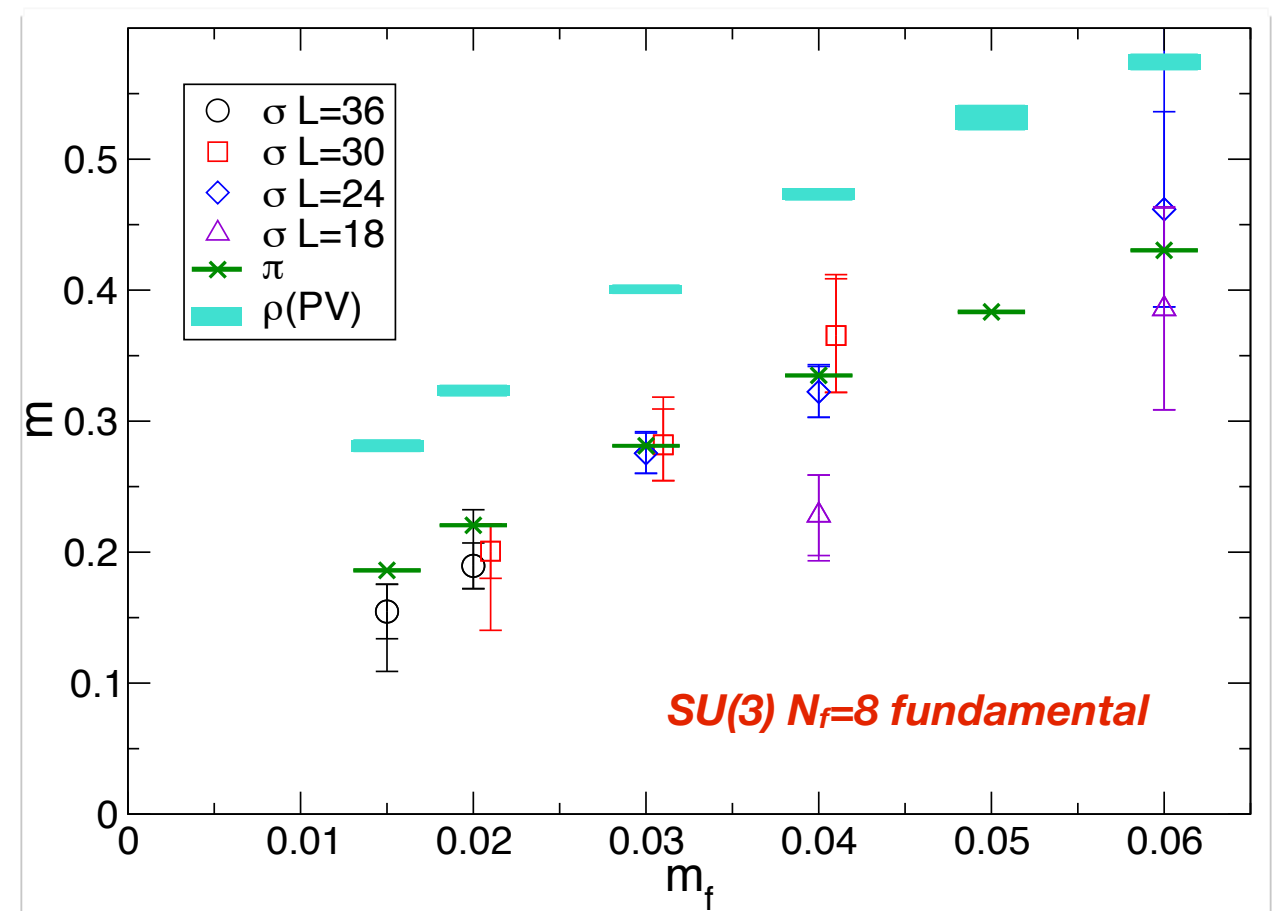
Composite Higgs models

LH Coll. PoS(2013)062



[see talk by R.Wong 2C for updates]

LatKMI Coll. PRD89(111502)R



[see talk by K.-i.Nagai 9C for updates]

Goldstone Higgs dynamics

- a composite Higgs can emerge as a **scalar** pseudo-NGB from breaking of a global symmetry in a new strong sector at the TeV scale
- phenomenologically interesting and well studied models are based on the breaking pattern **$SU(4) \rightarrow Sp(4)$** (equivalently $SO(6) \rightarrow SO(5)$) [Galloway et al., JHEP10(2010)89]
- a realization of this framework has been studied recently, showing how $SU(4)$ can break to $Sp(4)$ using a **$SU(2)$ gauge theory with 2 flavors in the fundamental representation** [Cacciapaglia&Sannino, JHEP04(2014)111]
 - ✓ $SU(4)/Sp(4)$ coset gives 5 NGBs \Rightarrow 3 pseudoscalars and 2 scalars
 - ✓ different choices of the quark condensate can be used when embedding the strong sector with the EW sector
 - ✓ interplay between the **2 scalars NGBs and the lightest excitation of the condensate** depending on the vacuum alignment give different scenarios [talk by A.Hietanen 2C]

Higgs and dark matter candidates

- interest in this special theory is not limited to Higgs compositeness:
 - models for **composite DM** have been recently studied based on this framework, without connection with EW symmetry breaking [\[Buckley&Neil, PRD87\(043510\)\]](#)
 - a model for **dark nucleosynthesis** based on the same strong sector now exists [\[Detmold et al., arxiv:1406.2276,arxiv:1406.4116\]](#) [talk by W.Detmold 8C]
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study the isosinglet scalar channel on the lattice

Lattice simulations with FUEL

- FUEL (Framework for Unified Evolution of Lattices) it's a lightweight and flexible wrapper for level 3 USQCD libraries with focus on lattice generation for BSM theories [\[https://github.com/jcosborn/qhmc\]](https://github.com/jcosborn/qhmc) [See J. Osborn's talk in 1F]
- capabilities to do arbitrary number of colors, dimensions and flavors
- Staggered fermion formulations, as well as Wilson ones are included
- use HMC to generate $O(10^3)$ trajectories for SU(2) with Wilson plaquette gauge action and 2 Wilson fermions
- one coupling $\beta=2.2$, one volume $32^3 \times 64$, six bare fermion masses $m_0=\{-0.68,-0.70,-0.72,-0.735,-0.75,-0.755\}$

Glueball spectroscopy: operators

- eigenstates of the Hamiltonian are classified according to the irreducible representations of the cubic group

$$\{A_1(1), A_2(1), E(2), T_1(3), T_2(3)\}$$

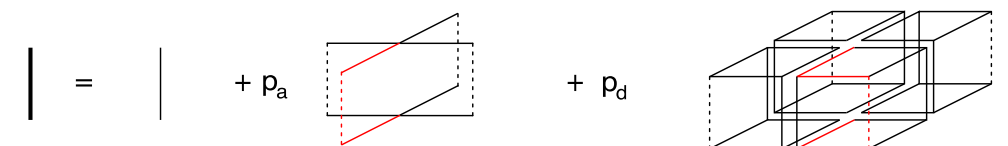
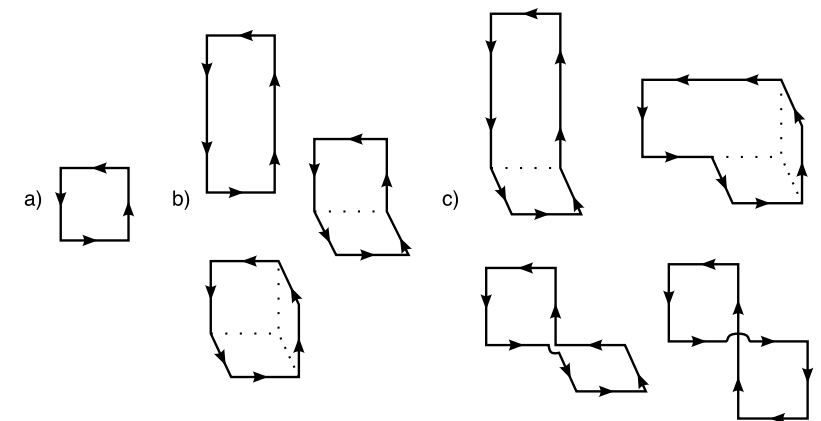
J	A_1	A_2	E	T_1	T_2
0	1	0	0	0	0
1	0	0	0	1	0
2	0	0	1	0	1
3	0	1	0	1	1
4	1	0	1	1	1

- suitable gauge-invariant operators must be constructed that respect the symmetries

$$\mathcal{O}_G(t) = \frac{1}{L^3} \sum_{x \in L^3} \text{Tr} \left(\prod_{l \in \mathcal{W}(x)} U_l \right) \quad \mathcal{O}_G^{(R)}(t) = \sum_{\alpha=1}^{24} a_{\alpha}^{(R)} \mathcal{R}_{\alpha} [\mathcal{O}_G(t)]$$

- vacuum contributions must be subtracted in the scalar case

$$\mathcal{O}^{(A_1)}(t) - \langle 0 | \mathcal{O}^{(A_1)} | 0 \rangle$$



- improved operators are obtained by blocking and smearing algorithms

Lucini, Rago, ER
JHEP08(2010)

Glueball spectroscopy: variational analysis

- basis of operators \rightarrow

$$\{\mathcal{O}_1(t), \dots, \mathcal{O}_n(t)\}$$

- matrix of correlators \rightarrow

$$C_{ij}(t) = \sum_{\tau} \langle 0 | \mathcal{O}_i^{\dagger}(\tau + t) \mathcal{O}_j(\tau) | 0 \rangle$$

- generalised eigenvalue problem \rightarrow

$$C_{ij}(t) v_j^{\alpha} = \lambda^{\alpha} v_i^{\alpha}$$

- ground state correlator fit ($\alpha=0$) \rightarrow

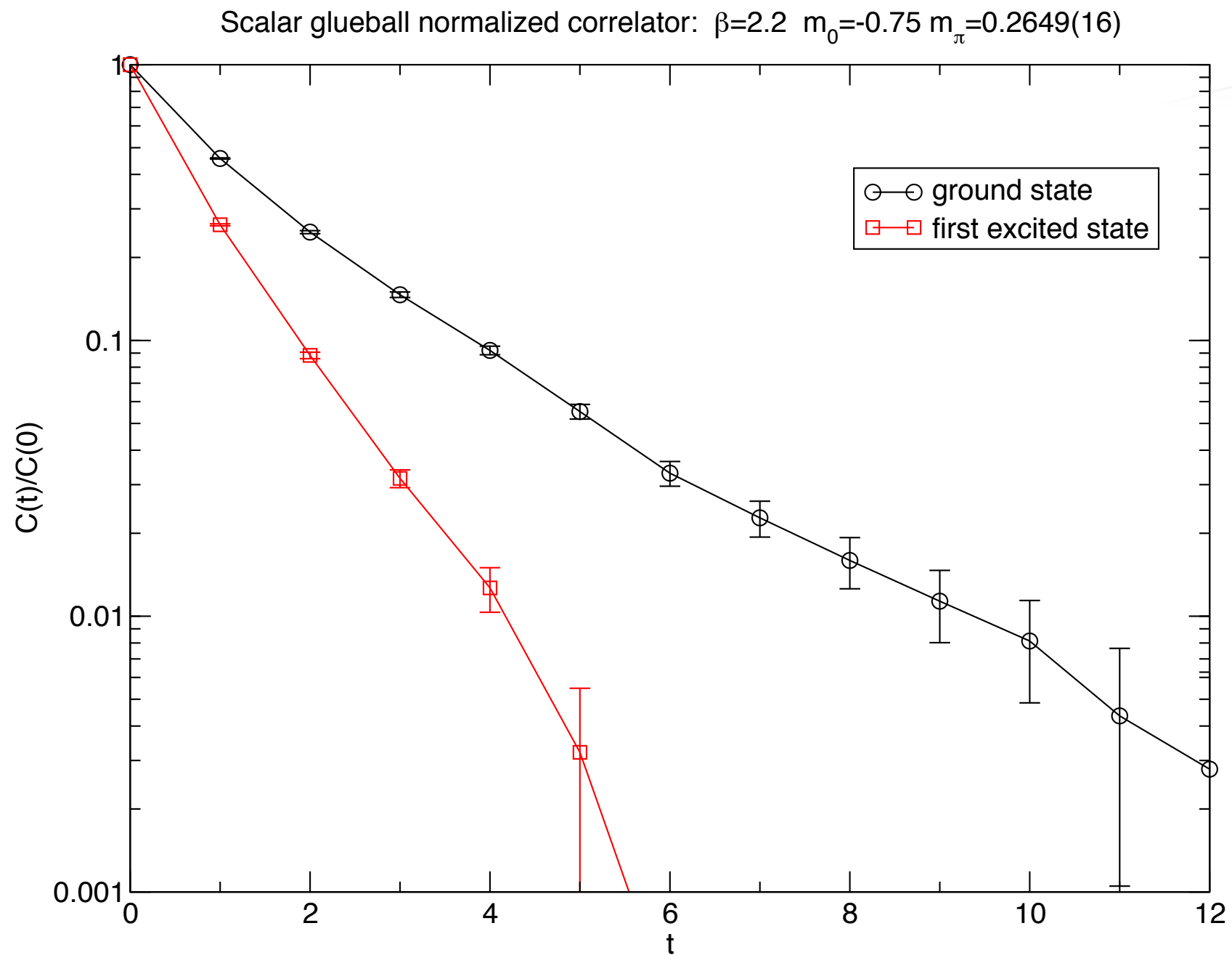
$$\Phi_{\alpha}(t) = \sum_{i=1}^n v_i^{\alpha} \mathcal{O}_i(t)$$

$$\langle \Phi_{\alpha}^{\dagger}(t) \Phi_{\alpha}(0) \rangle = |c_{\alpha}|^2 \left(e^{-m_{\alpha} t} + e^{-m_{\alpha} (T-t)} \right)$$

- the effective mass plateau is used to determine the fitting window on the correlator

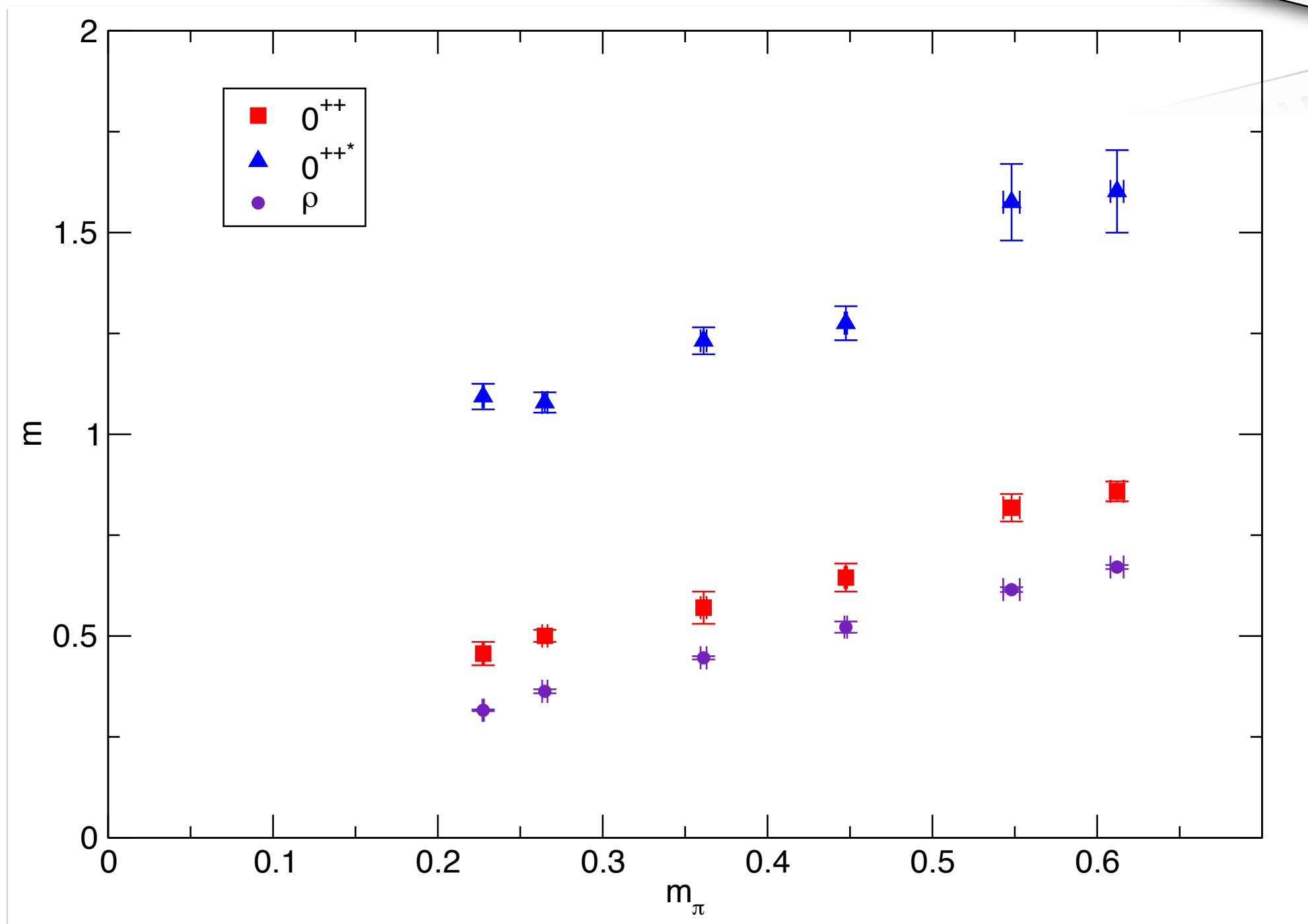
The scalar glueball

Preliminary



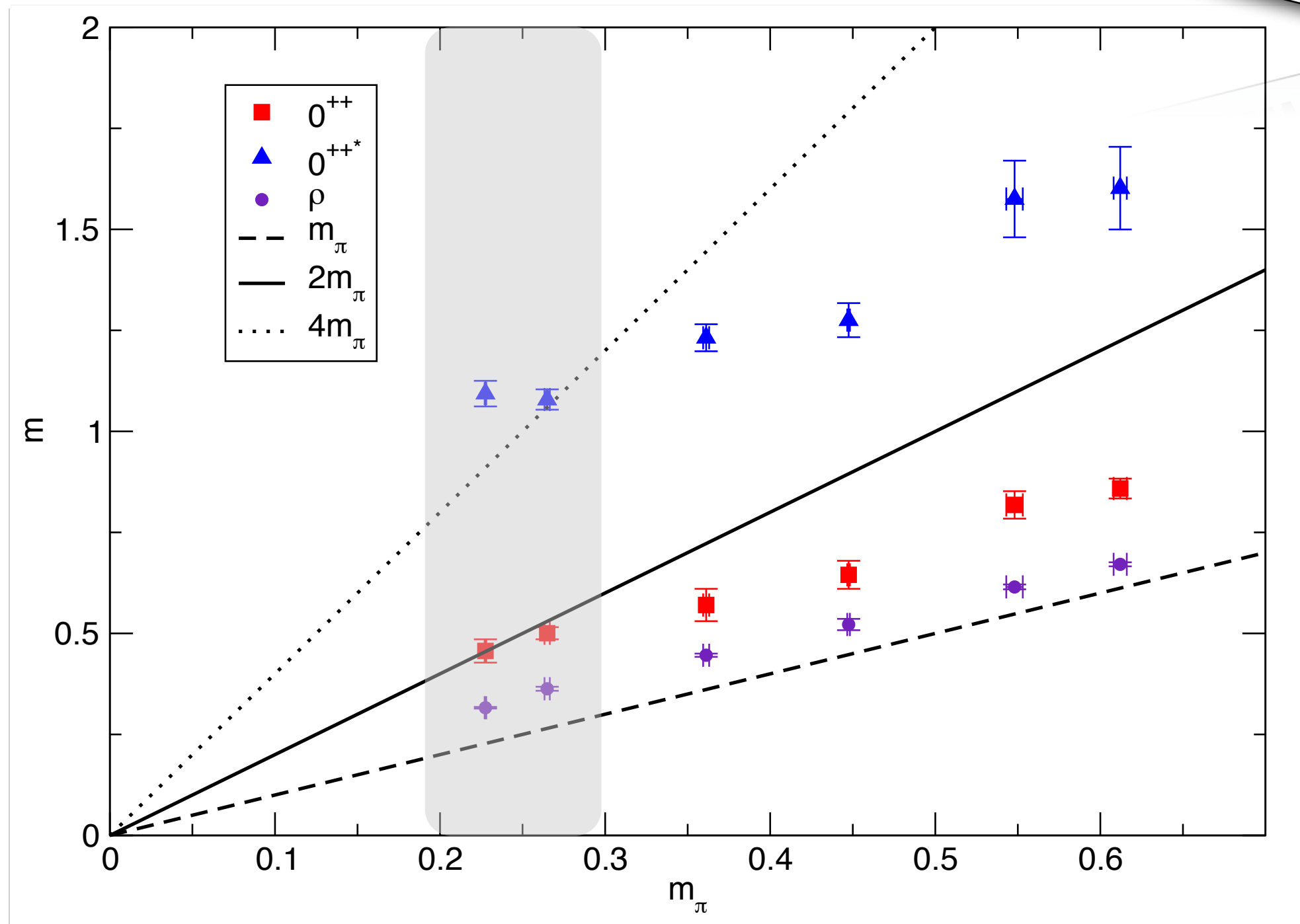
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Conclusions and future directions

- States in the **isosinglet scalar** channel have been investigated with **gluonic operators**, those coupling to glueballs in a pure gauge theory
- The ground and the first excited state are **heavier** than the isotriplet vector meson in the mass region explored
- Contributions due to **multi-pion** states become relevant in the light mass region
- Finite volume effects are being investigated for the lightest mass point on a $48^3 \times 96$ lattice
- A fermionic isosinglet scalar correlator, including the disconnected diagram, is being measured to address the mixing with the ground state

Thanks for your attention