## The Kaon Semileptonic Form Factor from Domain Wall QCD at the Physical Point

# Andreas Jüttner, Robert Mawhinney, David Murphy, Francesco Sanfilippo for the RBC/UKQCD Collaboration

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(3) Preliminary Analysis



3 Preliminary Analysis

- $K \to \pi l \nu$  decay gives most precise direct experimental constraint on CKM matrix element  $V_{us}$
- Measurements of decay rate constrain [M. Antonelli et al., Eur.Phys.J. C69 (2010)]

$$f_{+}^{K\pi}(0)|V_{us}| = 0.2163(5)$$

- Requires lattice calculation of non-perturbative form factor  $f_{+}^{K\pi}(0)$ 
  - Largest source of uncertainty
- Precision test of Standard Model unitarity

$$\delta_{CKM} = 1 - |V_{ud}|^2 - |V_{us}|^2 - |V_{ub}|^2$$

- Goal: ~ 0.3% lat. error on  $f_+^{K\pi}(0)$
- With recent physical point MILC result error from lattice calculation is approaching experimental error [A. Bazavov et al., Phys.Rev.Lett. 112 (2014)]

### $K \to \pi l \nu$ Decay in the Standard Model

#### Standard Model Results:

$$\Gamma_{K \to \pi l \nu} = C_K^2 \frac{G_F^2 |V_{us}|^2 m_K^5}{128\pi^3} \left| f_+^{K\pi}(\mathbf{0}) \right|^2 I_{Kl} S_{EW} \left( 1 + \delta_{SU(2)} + \delta_{EM} \right)$$
(1)  
$$\langle \pi(p_\pi) | V_\mu | K(p_K) \rangle = f_+^{K\pi}(q^2) \left( p_K + p_\pi \right)_\mu + f_-^{K\pi}(q^2) \left( p_K - p_\pi \right)_\mu$$
(2)

$$\langle \pi(p_{\pi})|S|K(p_K)\rangle|_{q^2=0} = \frac{m_K^2 - m_{\pi}^2}{m_s - m_u} f_+^{K\pi}(\mathbf{0})$$
 (3)

- Lattice:  $f_+^{K\pi}(0)$
- Theory:  $\delta_{SU(2)}$ ,  $\delta_{EM}$ ,  $S_{EW}$  (SU(2)-breaking, electromagnetic, and short-distance electroweak corrections)
- Experiment:  $I_{Kl}$  (phase-space integral)
- $q = p_K p_\pi$  (momentum transfer)
- $V_{\mu} = \overline{s} \gamma_{\mu} u$  (flavor-changing vector current)
- $S = \overline{s}u$  (flavor-changing scalar density)



3 Preliminary Analysis

- Focus on calculation for new RBC/UKQCD physical pion mass ensembles:
  - $48^3 \times 96 \times 24$  (48I):  $a^{-1} = 1.729(4)$  GeV
  - $64^3 \times 128 \times 16$  (64I):  $a^{-1} = 2.358(7)$  GeV
- Lattice setup:
  - ▶  $N_f = 2 + 1$  Möbius Domain Wall fermions
  - Iwasaki gauge action
  - All-mode averaging
  - Twisted boundary conditions
  - Coulomb gauge-fixed wall source/sink for  $K_{l3}$  3-point correlation functions



- $-\ell$ : light quark(s)
- s: strange quark
  - $\vec{p}$ : twist momentum

- In DWF simulations we have
  - Exactly conserved, point-split 5-d vector current  $\mathscr{V}^a_{\mu}(x)$
  - Local 4-d vector current  $V^a_{\mu}(x)$
- Latter is computationally cheaper, but needs to be renormalized
- We calculate  $Z_V$  from the pion and kaon electromagnetic form factors
- $\bullet\,$  Observe  $\sim 0.5\%$  dependence on flavor content of current

Ensemble	$Z_V^{\pi}$	$Z_V^K$
48I	0.71094(20)	0.71407(9)
64I	0.74295(18)	0.74517(5)

• Must agree in  $a^2 \to 0$  limit

### All-Mode Averaging [Blum et al., Phys.Rev. D88 (2013)]

- Most expensive part of calculation is light quark propagator inversion
- Idea is to decompose lattice observable  $\mathcal{O}(S_F)$  into
  - Cheap part, calculated with inexact fermion propagators  $(S'_F)$ :

$$\mathcal{O}' = \mathcal{O}(S'_F)$$

• Expensive part, calculated with exact fermion propagators  $(S_F)$ :

$$\Delta \mathcal{O} = \mathcal{O}(S_F) - \mathcal{O}(S'_F)$$

• AMA estimate is

$$\mathcal{O}_{AMA} = \frac{1}{N_g} \sum_{g \in G} \mathcal{O}' \circ g + \Delta \mathcal{O},$$

where G is the group of lattice translations

- For  $K_{l3}$  analysis:
  - Compute  $S'_F$  for all source/sink separations
  - Compute  $S_F$  for 7(8) source/sink separations on 48I(64I) ensemble
  - ▶ Form AMA observables, averaging  $\mathcal{O}'$  over all time translations

All-Mode Averaging [Blum et al., Phys.Rev. D88 (2013)]

• Exact and AMA correlators with a  $\pi - \pi$  separation of 20 on 48I ensemble:



D. Murphy (Columbia U.)



3 Preliminary Analysis

### Measurement Strategy

- Form factor and vector current renormalization constants extracted using ratio method [P.A. Boyle et al., JHEP 0705 (2007)]
  - $K_{l3}$  form factors:

$$R_{1}^{(\mu)}(\vec{p}_{K},\vec{p}_{\pi}) = 4\sqrt{E_{\pi}E_{K}}\sqrt{\frac{\mathcal{C}_{K\pi}^{(\mu)}(t,t_{\mathrm{snk}};\vec{p}_{K},\vec{p}_{\pi})\mathcal{C}_{\pi K}^{(\mu)}(t,t_{\mathrm{snk}};\vec{p}_{\pi},\vec{p}_{K})}{\tilde{\mathcal{C}}_{\pi}(t_{\mathrm{snk}};\vec{p}_{\pi})\tilde{\mathcal{C}}_{K}(t_{\mathrm{snk}};\vec{p}_{K})}} \qquad (4)$$

$$R_{2}^{(\mu)}(\vec{p}_{K},\vec{p}_{\pi}) = 2\sqrt{E_{\pi}E_{K}}\sqrt{\frac{\mathcal{C}_{K\pi}^{(\mu)}(t,t_{\mathrm{snk}};\vec{p}_{K},\vec{p}_{\pi})\mathcal{C}_{\pi K}^{(\mu)}(t,t_{\mathrm{snk}};\vec{p}_{\pi},\vec{p}_{K})}{\mathcal{C}_{\pi\pi}^{(0)}(t_{\mathrm{snk}};\vec{0})\mathcal{C}_{KK}(t_{\mathrm{snk}};\vec{0})}} \qquad (5)$$

$$R_{\alpha}^{(\mu)} \stackrel{t,t_{\mathrm{snk}}\gg1}{=} (p_{K}+p_{\pi})^{\mu}f_{+}^{K\pi}(q^{2}) + (p_{K}-p_{\pi})^{\mu}f_{-}^{K\pi}(q^{2})} \qquad (6)$$

 $\blacktriangleright$   $Z_V$ :

$$Z_V^{\pi} \stackrel{t \gg 1}{=} \frac{\tilde{\mathcal{C}}_{\pi}(t_{\mathrm{snk}})}{\mathcal{C}_{\pi\pi}(t_{\mathrm{src}}, t, t_{\mathrm{snk}})}, \quad Z_V^K \stackrel{t \gg 1}{=} \frac{\tilde{\mathcal{C}}_K(t_{\mathrm{snk}})}{\mathcal{C}_{KK}(t_{\mathrm{src}}, t, t_{\mathrm{snk}})}$$
(7)

• We have also explored global fit approach:

- Fit directly to correlators
- ▶ Reduce statistical error using maximal set of simultaneous constraints

#### Summary of Lattice Data

Lattice	$m_{\pi}$ (MeV)	$f_+^{K\pi}(0)$	Stat. error
24I	678	0.9992(1)	0.01%
24I	563	0.9956(4)	0.04%
24I	422	0.9870(9)	0.09%
24I	334	0.9760(43)	0.4%
24I	334	0.9858(28)	0.3%
<b>48I</b> (PRELIMINARY)	139	0.9727(25)	0.3%
32ID	248	0.9771(21)	0.2%
32ID	171	0.9710(45)	0.5%
32I	398	0.9904(17)	0.2%
32I	349	0.9845(23)	0.2%
32I	295	0.9826(35)	0.4%
<b>64I</b> (PRELIMINARY)	139	0.9701(22)	0.2%

- I: Iwasaki gauge action, ID: Iwasaki+DSDR gauge action
- $f_{+}^{K\pi}(0)$  extracted from vector current
- Unphysical ensembles analyzed in [P.A. Boyle et al., JHEP 1308 (2013)]

### Physical Point Extrapolation (PRELIMINARY)

- Physical point measurements still contain small errors from mistuning of input quark masses, cutoff effects, etc.
- Can use heavy data to correct to physical point:
  - SU(3)-breaking ansatz [P.A. Boyle et al., JHEP 1308 (2013)]:

$$f_{+}^{K\pi}(0; m_{\pi}^{2}, m_{K}^{2}) = A + \frac{(m_{K}^{2} - m_{\pi}^{2})^{2}}{m_{K}^{2}} \left(A_{0} + (m_{K}^{2} + m_{\pi}^{2})A_{1}\right)$$
(8)

#### Analytic ansätze:

 $f_{+}^{K\pi}(0;m_l,m_s) = f_{+}^{K\pi}(0;m_l^{\text{phys}},m_s^{\text{phys}}) + \frac{\partial f_{+}^{K\pi}}{\partial m_l} \Delta m_l + \frac{\partial f_{+}^{K\pi}}{\partial m_s} \Delta m_s + \frac{\partial^2 f_{+}^{K\pi}}{\partial m_l^2} \Delta m_l^2 + \frac{\partial^2 f_{+}^{K\pi}}{\partial m_l \partial m_s} \Delta m_l \Delta m_s + \frac{\partial^2 f_{+}^{K\pi}}{\partial m_s^2} \Delta m_s^2$ (9)  $f_{+}^{K\pi}(0;m_l,m_s^{\text{sea}},m_s^{\text{val}}) = f_{+}^{K\pi}(0;m_l^{\text{phys}},m_s^{\text{phys}},m_s^{\text{phys}}) + \frac{\partial f_{+}^{K\pi}}{\partial m_l} \Delta m_l + \frac{\partial f_{+}^{K\pi}}{\partial m_s^{\text{sea}}} \Delta m_s^{\text{sea}} + \frac{\partial f_{+}^{K\pi}}{\partial m_s^{\text{val}}} \Delta m_s^{\text{val}}$ (10)

- $\Delta m_q \equiv m_q m_q^{\text{phys}}$
- Physical quark masses  $(m_q^{\text{phys}})$  from SU(2) ChPT global fit analysis [R. Arthur et al., Phys.Rev. D87 (2013)]
- Can also add  $a^2$  scaling term

## Physical Point Extrapolation (PRELIMINARY)



(a) Quadratic analytic ansatz



- SU(3) ansatz:  $m_{\pi}^{\text{phys}} = 135.0 \text{ MeV}, m_{K}^{\text{phys}} = 495.7 \text{ MeV}$
- Estimate model error from difference in fits
- Preliminary results suggest systematic error from physical point extrapolation is  $\lesssim 0.1\%$  effect

- We have generated two new physical pion mass lattices, and computed  $Kl_3$  3-point functions on each
- Physical point measurements have reduced extrapolation error to sub-statistical
- Improved precision means deeper understanding of systematics necessary
  - Mass,  $a^2$  dependence of  $Z_V$
- Ongoing work:
  - ▶ Finalize fit systematics
  - Incorporate data from scalar current
  - Explicit continuum extrapolation?

# Thank you!

### Twisted Boundary Conditions

• Momentum states in periodic, cubic lattice with side length L and spacing a are

$$p_i = \frac{2\pi k_i}{L}, \ k_i \in \mathbb{Z}$$

• On our physical point ensembles, smallest allowed momentum states are

$$48^3: \frac{2\pi}{48a} \approx 226 \,\text{MeV}, \quad 64^3: \frac{2\pi}{64a} \approx 232 \,\text{MeV}$$

• 
$$q^2 = 0$$
 requires

$$|\vec{p}_{\pi}| = \frac{m_K^2 - m_{\pi}^2}{2m_K} \approx 229 \,\mathrm{MeV}$$

• Avoid momentum interpolation via twisted boundary conditions

$$\psi(x + L\hat{e}_i) = e^{i\theta_i}\psi(x) \implies p_i = \frac{2\pi k_i}{L} + \frac{\theta_i}{L}$$

• Further reduction in statistical noise from moving pion by twisting in all 3 spatial directions and averaging spatial components of vector correlator