

TWO-B ARYON SYSTEMS WITH TWISTED BOUNDARY CONDITIONS



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In collaboration with:

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TWISTED BOUNDARY CONDITIONS



A GENERAL TWO-BARYON FINITE-VOLUME QUANTIZATION CONDITION



Arbitrary CM momentum, twisted boundary conditions, arbitrary masses and relativistic kinematics. *Briceno, ZD and Luu, Phys. Rev. D* 88, 034502.

Briceno, ZD and Luu, Phys. Rev. D 88, 034502. Briceno, ZD, Luu and Savage, Phys. Rev. D 89, 074509.

See R. Briceno, Phys. Rev. D 89, 074507 for generalization to arbitrary spin.



Kim, Sachrajda, Sharpe, et al, Nucl.Phys. B727 (2005) 218-243.





TWO-NUCLEON CORRELATION FUNCTION $T \to \infty, \ a \to 0$ \sim V \overline{O} + VV σ +VV σ' A' $(\mathcal{M}_{\infty}) V$ (A) $V\left(\mathcal{M}_{\infty}\right)$ (\mathcal{M}_{∞}) (A')VVV+Luscher, Nucl.Phys. B354, 531-578 (1991). $\det\left(\delta \mathcal{G}^V + \mathcal{M}^{-1}\right) = 0$ Briceno, ZD and Luu, Luscher, Commun.Math.Phys., 105:153 Phys. Rev. D 88, 034502. (1986). Non-diagonal in A function of Parametrized by phase shifts Diagonal in J both L and J basis and mixing angles energy and volume basis



FINITE-VOLUME FUNCTION

$$\begin{split} \left[\delta \mathcal{G}^{V} \right]_{JM_{J},LS;J'M'_{J},L'S'} &= i\eta \frac{k^{*}}{8\pi E^{*}} \delta_{SS'} \left[\delta_{JJ'} \delta_{M_{J}M'_{J}} \delta_{LL'} + i \sum_{l,m} \frac{(4\pi)^{3/2}}{k^{*l+1}} c_{lm}^{\mathbf{d},\phi_{1},\phi_{2}}(k^{*2};L) \\ &\times \sum_{M_{L},M'_{L},M_{S}} \langle JM_{J} | LM_{L}, SM_{S} \rangle \langle L'M'_{L}, SM_{S} | J'M'_{J} \rangle \int d\Omega \ Y^{*}_{LM_{L}} Y^{*}_{lm} Y_{L'M'_{L}} \right] \end{split}$$

Ishizuka, PoS, LAT2009, 119 (2009).

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ON-SHELL KINEMATICS AND FINITE-VOLUME FUNCTION



Volume and its shape, momentum of the CM, and the boundary conditions

Li and Liu, Phys.Lett. B587, 100 (2004). Detmold W and Savage M J 2004 Nucl.Phys. A743 170–193. Feng X, Li X and Liu C 2004 Phys.Rev. D70 014505.

$$\psi_{\text{Lab}}(\mathbf{x}_1 + L\mathbf{n}_1, \mathbf{x}_2 + L\mathbf{n}_2) = e^{i\phi_1 \cdot \mathbf{n}_1 + i\phi_2 \cdot \mathbf{n}_2} \psi_{\text{Lab}}(\mathbf{x}_1, \mathbf{x}_2)$$
$$\mathbf{n}_1, \mathbf{n}_2 \in \mathbb{Z}^3$$

 $\mathbf{r} = \frac{1}{L} \,\hat{\gamma}^{-1} \left[2\pi (\mathbf{n} - \alpha \mathbf{d}) - (\alpha - \frac{1}{2})(\phi_1 + \phi_2) + \frac{1}{2}(\phi_1 - \phi_2) \right]$

Briceno, ZD, Luu and Savage, Phys. Rev. D 89, 074509.

$$\alpha = \frac{1}{2} \left(1 + \frac{m_1^2 - m_2^2}{E^{*2}} \right)$$

$$\det\left(\delta \mathcal{G}^V + \mathcal{M}^{-1}\right) = 0$$

ON-SHELL KINEMATICS AND FINITE-VOLUME FUNCTION

r² Depends on:

Volume and its shape, momentum of the CM, and the boundary conditions

Li and Liu, Phys.Lett. B587, 100 (2004). Detmold W and Savage M J 2004 Nucl.Phys. A743 170–193. Feng X, Li X and Liu C 2004 Phys.Rev. D70 014505.

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Briceno, ZD, Luu and Savage, Phys. Rev. D 89, 074509.

Rummukainen, Gottlieb, Nucl.Phys. B450 (1995) 397-436. Kim, Sachrajda, Sharpe, et al, Nucl.Phys. B727 (2005) 218-243.

Bedaque, Phys.Lett. B593 (2004) 82-88.

Agadijanov, Meissner and Rusetsky, JHEP01. 103 (2014).

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 $\det\left(\delta \mathcal{G}^V + \mathcal{M}^{-1}\right) = 0$

ZD, Savage, Phys.Rev. D84 (2012) 114502. Fu, Phys.Rev. D85 (2012) 014506. Leskovec, Prevolesk, Phys.Rev. D85 (2012) 114507.

THE SPECTRUM OF THE DEUTERON IN A FINITE VOLUME



1) Periodic boundary conditions and varying the CM momentum

See Briceno, ZD Luu and Savage, arXiv:1309.1556 (2013).

2) Varying the twisted boundary conditions with zero CM momentum

Briceno, ZD, Luu and Savage, Phys. Rev. D 89, 074509.

INPUT: SCATTERING PARAMETERS FROM PHENOMENOLOGICAL ANALYSES

S-1

 4π

 $\mathcal{M} =$

$$S_{(J=1)} = \begin{pmatrix} \cos \epsilon_1 & -\sin \epsilon_1 \\ \sin \epsilon_1 & \cos \epsilon_1 \end{pmatrix} \begin{pmatrix} e^{2i\delta_{1\alpha}} & 0 \\ 0 & e^{2i\delta_{1\beta}} \end{pmatrix} \begin{pmatrix} \cos \epsilon_1 & \sin \epsilon_1 \\ -\sin \epsilon_1 & \cos \epsilon_1 \end{pmatrix}$$

Biedenharn, Blatt, Phys.Rev. 86, 399 (1952). Biedenharn, Blatt, Phys.Rev. 93, 1387 (1954).

 ik^*



Wednesday, June 25, 2014

NN-OnLine

DEUTERON SPECTRUM TWISTED BOUNDARY CONDITIONS AND VOLUME IMPROVEMENT

 $\phi^p = -\phi^n \equiv \phi = (0, 0, 0) \text{ (PBCs)}, \ (\pi, \pi, \pi) \text{ (APBCs)}, \text{ and } (\frac{\pi}{2}, \frac{\pi}{2}, \frac{\pi}{2})(i - PBCs)$



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DEUTERON SPECTRUM TWISTED BOUNDARY CONDITIONS AND VOLUME IMPROVEMENT

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VOLUME-EFFECTS REDUCTION AN ANALYTICAL INVESTIGATION

The S-wave limit

$$k^{*} \cot \delta^{(^{3}S_{1})}|_{k^{*}=i\kappa} + \kappa = \sum_{\mathbf{n}\neq\mathbf{0}} e^{i(\alpha-\frac{1}{2})\mathbf{n}\cdot(\phi^{p}+\phi^{n})} e^{-i\frac{1}{2}\mathbf{n}\cdot(\phi^{p}-\phi^{n})} e^{i2\pi\alpha\mathbf{n}\cdot\mathbf{d}} \frac{e^{-|\hat{\gamma}\mathbf{n}|\kappa\mathbf{L}}}{|\hat{\gamma}\mathbf{n}|\mathbf{L}}$$
$$\kappa = \sqrt{MB_{d}} \approx 45.7 \text{ MeV}$$

VOLUME-EFFECTS REDUCTION AN ANALYTICAL INVESTIGATION

The S-wave limit

$$\begin{aligned} k^* \cot \delta^{(^3S_1)}|_{k^*=i\kappa} + \kappa &= \sum_{\mathbf{n}\neq\mathbf{0}} e^{i(\alpha-\frac{1}{2})\mathbf{n}\cdot(\phi^n+\phi^n)} e^{-i\frac{1}{2}\mathbf{n}\cdot(\phi^p-\phi^n)} e^{i2\pi\alpha\mathbf{n}\cdot\mathbf{d}} \frac{e^{-|\gamma\mathbf{n}|\kappa\mathbf{L}}}{|\gamma\mathbf{n}|\mathbf{L}} \\ &\quad \kappa &= \sqrt{MB_d} \approx 45.7 \text{ MeV} \end{aligned}$$

$$\begin{aligned} \sum_{\mathbf{n}\neq\mathbf{0}} e^{-i\mathbf{n}\cdot\phi} \frac{e^{-|\mathbf{n}|\kappa\mathbf{L}|}}{|\mathbf{n}|\mathbf{L}} \\ \phi &= (0,0,0) \quad \mathcal{O}(e^{-\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{2}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{3}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-2\kappa\mathbf{L}}/\mathbf{L}) \\ \phi &= (\frac{\pi}{2}, \frac{\pi}{2}, \frac{\pi}{2}) \quad \mathcal{O}(e^{-\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{2}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{3}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-2\kappa\mathbf{L}}/\mathbf{L}) \\ \phi &= (\pi, \pi, \pi) \quad \mathcal{O}(e^{-\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{2}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{3}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-2\kappa\mathbf{L}}/\mathbf{L}) \\ A^{\text{verage of}} \quad \mathcal{O}(e^{-\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{2}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-\sqrt{3}\kappa\mathbf{L}}/\mathbf{L}) \quad \mathcal{O}(e^{-2\kappa\mathbf{L}}/\mathbf{L}) \end{aligned}$$

PARTIALLY TWISTED BOUNDARY CONDITIONS

NN sector

Bedaque and Chen, Phys.lett. B616, 208 (2005).

* No s-channel diagram with intermediate mesons containing a sea quark.

* The quantization condition valid up to exponential corrections in the masses of exchanged bosons.



Two-meson sector

* s-channel diagrams with intermediate mesons containing a sea quark exist.

* Quantization condition? Intricate cancellations in partially quenched QCD?



Agadijanov, Meissner and Rusetsky, JHEP01. 103 (2014).

SUMMARY AND CONCLUSIONS

A general quantization condition for the two-hadron systems in a finite volume with twisted boundary conditions is presented.



The quantization condition holds for partial twisting scenario in the two-nucleon sector and makes this approach computationally plausible.



Calculations with optimized TBCs make the extraction of binding energies feasible at smaller volumes.





Volume improvement in three (and higher)-particle bound states energies using optimized boundary conditions?

THANK YOU

BACKUP SLIDES

ON-SHELL KINEMATICS AND FINITE-VOLUME FUNCTION

$$\psi_{\text{Lab}}(\mathbf{x}_1 + L\mathbf{n}_1, \mathbf{x}_2 + L\mathbf{n}_2) = e^{i\phi_1 \cdot \mathbf{n}_1 + i\phi_2 \cdot \mathbf{n}_2} \psi_{\text{Lab}}(\mathbf{x}_1, \mathbf{x}_2)$$

(1) $\psi_{\text{Lab}}(x_1, x_2) = e^{-iEX^0 + i\mathbf{P}\cdot\mathbf{X}}\varphi_{\text{Lab}}(0, \mathbf{x}_1 - \mathbf{x}_2)$ $X = \alpha x_1 + (1 - \alpha)x_2$ $\alpha = \frac{1}{2}\left(1 + \frac{m_1^2 - m_2^2}{E^{*2}}\right)$

Luscher, Nucl.Phys. B354 (1991) 531-578. Rummukainen, Gottlieb, Nucl.Phys. B450 (1995) 397-436.

(2)
$$\varphi_{\text{Lab}}(0, \mathbf{x_1} - \mathbf{x_2}) = \varphi_{\text{CM}}(\hat{\gamma}(\mathbf{x_1} - \mathbf{x_2}))$$

 $\hat{\gamma}\mathbf{x} = \gamma \mathbf{x_{\parallel}} + \mathbf{x_{\perp}}$

Bour, et al, Phys.Rev. D84 (2011) 091503. ZD, Savage, Phys.Rev. D84 (2012) 114502.

(3)
$$e^{i\alpha \mathbf{P} \cdot (\mathbf{n}_1 - \mathbf{n}_2)L + i\mathbf{P} \cdot \mathbf{n}_2 L} \varphi_{\mathrm{CM}}(\mathbf{y}^* + \hat{\gamma}(\mathbf{n}_1 - \mathbf{n}_2)L) = e^{i\phi_1 \cdot \mathbf{n}_1 + i\phi_2 \cdot \mathbf{n}_2} \varphi_{\mathrm{CM}}(\mathbf{y}^*)$$

 $\mathbf{y}^* = \mathbf{x}_1^* - \mathbf{x}_2^*$

$$\mathbf{r} = \frac{1}{L} \,\hat{\gamma}^{-1} \left[2\pi (\mathbf{n} - \alpha \mathbf{d}) - (\alpha - \frac{1}{2})(\phi_1 + \phi_2) + \frac{1}{2}(\phi_1 - \phi_2) \right]$$

Briceno, ZD, Luu and Savage, Phys. Rev. D 89, 074509.

See also: R. Briceno, Phys. Rev. D 89, 074507.