Leptonic B and D decay constants with 2+1 flavor asqtad fermions

Ethan T. Neil (University of Colorado/RIKEN BNL) Lattice 2014 for the Fermilab/MILC Collaboration

Motivation

• Leptonic heavy meson decays are sensitive to both strong and weak physics:



- Accurate f_H (from lattice) is crucial for precise CKM matrix elements!
- Aside from determining CKM, decay constants are needed for rare leptonic decays - whether mediated by standard model or by new physics



(from talk by J. Serrano (LHCb), La Thuille 2014)

$B_{s/d} \rightarrow \mu^+ \mu^-$

CMS PAS BPH-13-007, LHCb-CONF-2013-012

- CMS (25 fb⁻¹) and LHCb (3 fb⁻¹) both found evidence for the very rare decay $B_s \rightarrow \mu^+ \mu^-$, in agreement with SM
- Combining CMS and LHCb: first observation of $B_s \rightarrow \mu^+ \mu^-$

 $BR(B_{S}^{0} \to \mu^{+}\mu^{-}) = (2.9 \pm 0.7) \times 10^{-9} \quad BR(B^{0} \to \mu^{+}\mu^{-}) = (3.6^{+1.6}_{-1.4}) \times 10^{-10}$



Simulation details

- MILC asqtad ensembles (see right), clover heavy quarks w/ Fermilab interpretation
- Open discs show the ensembles used in last iteration of this analysis, filled circles are new; circle area ~ statistics
- New analysis has more lattice spacings, more stats, lower accessible quark masses. 4 time-sources per config.



Two-point correlators

 Decay constant comes from two-point function of pseudoscalar (PS) source, axial-vector sink. Calculate PS-PS correlator as well to fix normalization.

$$C_{ij}(t) = \sum_{n=0}^{N_X} \left[A_{i,n} A_{j,n} \left(e^{-E_n t} + e^{-E_n (N_t - t)} \right) - (-1)^t A'_{i,n} A'_{j,n} \left(e^{-E'_n t} + e^{-E'_n (N_t - t)} \right) \right]$$

- Our procedure: fit uniformly down to small t (t_{min}=4) for all correlators, using as many excited states as necessary.
- N_x=3 (4+4 states) for most ensembles, with 1 more state pair used on finest lattices
- This approach is relatively simple (no tuning fit ranges) and doesn't throw away statistics, but as a multi-exponential fit is highly prone to numerical instability!

Multi-exponential fits and instability

- Without additional constraints, multi-exponential fits suffer from ordering ambiguity: who says "E₀" is the smallest energy? N! minima in probability space!
- Variable re-mapping can force a particular ordering, e.g.

 $E_0, \log(E_1 - E_0), \log(E_2 - E_1), \dots$

- Still numerically challenging, many flat directions even if we lift the spurious minima
- Getting the ground state right tends to stabilize the rest of the fit...



Two-stage constrained fits

- Lots of different methods out there for dealing with this problem; this isn't a review talk so I'll just focus on my approach
- Constrained curve fitting (Lepage et al, arXiv:hep-lat/0110175) eliminates ^{0.56}
 flat directions by use of Gaussian 0.54
 "prior" constraints on fit parameters 0.53
- "Priors" tend to be set in a datadriven way in many analyses, e.g. by looking at a subset of the data.
 So why not use the same correlator to set up the full fits?

"Two-stage" constrained fits



1) Fit ground state in plateau region

2) Use fit mean values and errors (scaled up) to set "priors"

3) Fit full correlator with N_x excited states

Stability for replication methods (jackknife/bootstrap)!









"Two-stage" features and plans

- Very stable, insensitive to detailed choice of plateau or factor used to scale errors up. (We picked one plateau on one correlator, then rescaled in terms of r₁/a and got stable results almost everywhere.)
- Method write-up to be included in finished paper; code to be publicly released
- Fitter is implemented in Python, with the following features:
 - Bayesian constrained fits or unconstrained, interface to standard Python optimizers, MINUIT, and MCMC packages
 - Symbolic manipulation of model functions (no need to input derivatives by hand)
 - On-the-fly compilation to C or Fortran for quick fitting useful for bootstrap and jackknife loops
 - Black-box implementation of two-stage fit procedure with "plateau finder" (user controllable)

Renormalization and tuning

 "Mostly non-perturbative" renormalization of results: Z_V determined non-perturbatively, leftover piece computed in lattice PT.

$$\phi_Q = \sqrt{2} Z_{A_{Qq}^4} A_{A_{Qq}^4} = \sqrt{2} \left(\rho_{A_{Qq}^4} \sqrt{Z_{V_{qq}^4} Z_{V_{QQ}^4}} \right) A_{A_{Qq}^4}$$

 Simulation values of heavy-quark mass parameter к do not exactly match, so we have to re-tune:

$$\phi_Q \to \phi_Q + \Delta \phi_Q = \phi_Q + \left(\frac{d\phi_Q}{d\kappa}\right) (\kappa_{\rm sim} - \kappa_{\rm tune})$$

• Tuning factor calculated non-perturbatively from runs at various κ on same ensembles; numerically small (mistunings are O(10⁻³).)

Chiral/continuum extrapolation

• Fit to rooted staggered chiPT (rSxPT) to extrapolate to continuum and physical quark masses:

 $\phi = \phi^{0} [1 + (\text{chiral logs}) + (\text{NLO analytic}) + (\text{NNLO analytic}) + (\text{LQ discretization}) + (\text{HQ discretization})]$

- Terms included for taste-breaking effects, finite-volume corrections in chiral logs, hyperfine/LQ flavor splitting of heavy-light mesons (details in arXiv: 1112.3051.) Hyperfine splittings found to be important for good fits!
- Fits can be used to split out and quantify each discretization error. Remaining systematic effects (mainly from exact form of chiral fit and scale-setting) are estimated by comparison with variant chiral fits.
- NNLO terms found necessary to fit points with valence mass near strange.

Warning: results include an unknown blinding normalization!





Conclusion

- Correlator fits and central chiral fits are done, unblinding after systematic error estimation is in place (explicit fit parameters and variant chiral fits)
- Details on "two-stage" fitting method and Python fitting code to be released with paper
- Competitive results for decay-constant ratios (*preliminary*: 0.5% stats+discretization for f_{Ds}/f_D, 0.8% for f_{Bs}/f_B)
- Combination of f_B/f_D with MILC HISQ results for D decay constants planned, may allow for more precise f_B estimate