

# Comparison of lattice definitions of the topological charge

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Presentation
$\mathbf{outline}$

- Introduction
- Results
- Conclusions

#### 1. Introduction

- Motivation
- Short overview of employed definitions
- 2. Results
  - Comparison at a single lattice spacing
  - Increase of correlation towards the continuum limit
  - Topological susceptibility
- 3. Conclusions



Presentation outline

Introduction

#### Motivation

- Definitions
- Index
- Spectral flow
- Spectral projectors
- Fermionic disc.
- Field theoretic
- Lattice setup
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# Motivation



- theoretical
  - $\star$  there are many definitions of the topological charge
  - $\star$  which of them should one use for different purposes?
    - $\diamond$   $\,$  to compute the topological susceptibility
    - $\diamond$  to sort configurations according to the topological charge
    - $\diamond$  for weighting results with the topological charge
  - $\star$  what are the pros and cons of different definitions?
  - $\star$  which definitions are theoretically clean and which are "suspicious"?
- practical
  - $\star~$  quite a lot of data for the topological charge from different definitions and several projects



# Definitions



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- index of the overlap Dirac operator [K.C., E. García Ramos, K. Jansen]
- spectral flow of the Hermitian Wilson-Dirac operator [U. Wenger]
- spectral projectors [K.C., E. García Ramos, K. Jansen]
- fermionic from disconnected loops [C. Michael, K. Ottnad, C. Urbach, F. Zimmermann]
- field theoretic with HYP smearing [U. Wenger, F. Zimmermann]
- field theoretic with APE smearing [A. Dromard, M. Wagner, F. Zimmermann]
- field theoretic with cooling [A. Dromard, M. Wagner]
- field theoretic with gradient flow [U. Wenger]



Index of the overlap Dirac operator



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- Chirally symmetric fermionic discretizations allow exact zero modes of the Dirac operator.
- The famous Atiyah-Singer index theorem [M. Atiyah, I.M. Singer, Ann. Math. 93, 139 (1971) 168] relates the topological charge to the number of zero modes of the Dirac operator:

 $Q = n_- - n_+$ 

- This is quite remarkable, because a property of gauge fields is linked to a fermionic observable.
- uniqueness: dependence on the *s* parameter of the kernel warning: locality [K. C., V. Drach, E. García Ramos, G. Herdoiza, K. Jansen, Nucl.Phys. **B**869 (2013) 131-163, arXiv:1211.1605 [hep-lat]]
- pros: theoretically clean, integer-valued, no renormalization
- cons: cost, cost, cost...



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# Spectral flow of the Hermitian Wilson-Dirac operator



- closely related and actually equivalent to the index
- $\mathbf{uniqueness}$ : dependence on the s parameter of the kernel
- pros: theoretically clean, integer-valued, no renormalization, one computation leads to the whole *s*-dependence of the index
- cons: cost (although still cheaper than index), non-trivial to analyze data



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#### Spectral projectors

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# Spectral projectors

- another fermionic definition, introduced in: [L. Giusti, M. Lüscher, 2008], [M. Lüscher, F. Palombi, 2010]
  - $\mathbb{P}_M$  projector to the subspace of eigenmodes of  $D^{\dagger}D$  with eigenvalues below  $M^2$ , evaluated stochastically
  - $Q = \operatorname{Tr} \{\gamma_5 \mathbb{P}_M\}$  for chirally symmetric fermions
- spectral projectors are then also equivalent to the index (=stochastic way of counting the zero modes)
- for non-chirally symmetric fermions it still gives a clean definition, although chirality of modes is no longer  $\pm 1 \longrightarrow \pm 1 + \mathcal{O}(a^2)$
- in practice, one evaluates the observable:

$$\mathcal{C} = \frac{1}{N} \sum_{k=1}^{N} \left( \mathbb{P}_M \eta_k, \gamma_5 \mathbb{P}_M \eta_k \right),$$

- uniqueness: dependence on the M of  $\mathbb{P}_M$
- pros: theoretically clean, still rather cheap
- cons: stochastic ingredient, non-integer,  $Z_S/Z_P$  needed



Fermionic from disconnected loops



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#### Fermionic disc.

- Field theoretic Lattice setup
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• another fermionic definition, given by Chris Michael:

$$N_f Q = m_q \sum \bar{\psi} \gamma_5 \psi = \sum \frac{m_q \gamma_5}{D + m_q},$$

in the limit as  $m_q \to 0$ .

- allows for a Q computation as a by-product of evaluation of disconnected loops
- uniqueness: yes
- pros: cheap if treated as a by-product
- cons: unclear to what extent it is valid, stochastic ingredient, non-integer,  $Z_S/Z_P$  needed



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# Field theoretic

- a very natural definition
- in the continuum:

$$Q = \frac{1}{32\pi^2} \int d^4x \,\epsilon_{\mu\nu\rho\sigma} \mathrm{tr}[F_{\mu\nu}(x)F_{\rho\sigma}]$$

- on the lattice one has to choose some discretization
- renormalization:

 $q_R[U] = \operatorname{round}(Zq_{\operatorname{bare}}[U]),$ 

Z is the (non-zero) solution of:

$$\min \sum_{U} \left( Zq_{\text{bare}} - \text{round}(Zq_{\text{bare}}[U]) \right)^2$$

• smoothing of gauge fields needed

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# Field theoretic



- smoothing:
  - ★ cooling an iterative minimization of the lattice action, eliminates rough topological fluctuations while keeping large instantons unchanged and decreases renormalization effects by smoothing out the UV noise
    - [B. Berg, 1981], [Y. Iwasaki et al., 1983], [M. Teper, 1985],
    - [E. Ilgenfritz et al., 1986]
  - ★ HYP/APE smearing [M. Albanese et al., 1987], [A. Hasenfratz,
    F. Knechtli, 2001]
  - $\star$  gradient flow (GF) [M. Lüscher, 2010]
- GF is very important from the point of view of validity of the field theoretic approach
- uniqueness: discretization of F, level of smoothing
- pros: very cheap (but: cost of smoothing), theor. clean if GF used for smoothing and no renorm. then [M. Lüscher, P. Weisz, 2011]
- cons: HYP/APE or cooling a bit *ad hoc* (require renorm.)



# Lattice setup



Ensemble	eta	lattice	$a\mu_l$	$\mu_R \; [\text{MeV}]$	$\kappa_c$	$L \; [\mathrm{fm}]$	$m_{\pi}L$
b40.16	3.90	$16^3 \times 32$	0.004	21	0.160856	1.4	2.5
c30.20	4.05	$20^3 \times 40$	0.003	19	0.157010	1.3	2.4
d20.24	4.20	$24^3 \times 48$	0.002	15	0.154073	1.3	2.4
e17.32	4.35	$32^3 \times 64$	0.00175	16	0.151740	1.5	2.4
Ensemble	eta	lattice	$a\mu_l$	$\mu_{l,R}$ [MeV]	$\kappa_c$	L [fm]	$m_{\pi}L$
A30.32	1.90	$32^3 \times 64$	0.0030	13	0.163272	2.8	4.0
A40.32	1.90	$32^3 \times 64$	0.0040	17	0.163270	2.8	4.5
A50.32	1.90	$32^3 \times 64$	0.0050	22	0.163267	2.8	5.1
A60.24	1.90	$24^3 \times 48$	0.0060	26	0.163265	2.1	4.2
A80.24	1.90	$24^3 \times 48$	0.0080	35	0.163260	2.1	4.8
B25.32	1.95	$32^3 \times 64$	0.0025	13	0.161240	2.5	3.4
B35.32	1.95	$32^3 \times 64$	0.0035	18	0.161240	2.5	4.0
B55.32	1.95	$32^3 \times 64$	0.0055	28	0.161236	2.5	5.0
B75.32	1.95	$32^3 \times 64$	0.0075	38	0.161232	2.5	5.8
B85.24	1.95	$24^3 \times 48$	0.0085	45	0.161231	1.9	4.7
D20.48	2.10	$48^3 \times 96$	0.0020	12	0.156357	2.9	3.9
D30.48	2.10	$48^3 \times 96$	0.0030	19	0.156355	2.9	4.7
D45.32	2.10	$32^3 \times 64$	0.0045	29	0.156315	1.9	3.9



### Histograms – b40.16





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### Histograms – b40.16





Field theoretic definition

#### 10 HYP iterations

40 HYP iterations

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### Histograms – b40.16





Field theoretic definition with cooling

5 steps with tol. 5%

10 steps with tol. 1%

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# Cooling vs. APE/HYP smearing



Field theoretic definition

#### cooling

#### APE/HYP smearing

Europeq.

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ÂM

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	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21
1	1.000(0)	0.889(12)	0.950(7)	0.877(15)	0.951(4)	0.929(9)	0.924(10)	0.585(51)	0.473(45)	0.864(16)	0.827(21)	0.814(20)	0.180(59)	0.907(12)	0.889(15)	0.909(10)	0.838(19)	0.830(20)	0.810(21)	0.889(12)	0.848(18)
2	0.889(12)	1.000(0)	0.903(10)	0.881(11)	0.901(10)	0.893(10)	0.888(11)	0.570(42)	0.489(45)	0.849(16)	0.807(22)	0.792(22)	0.208(40)	0.892(11)	0.874(13)	0.987(1)	0.850(16)	0.836(18)	0.795(21)	1.000(0)	0.867(14)
3	0.950(7)	0.903(10)	1.000(0)	0.898(10)	0.990(2)	0.937(6)	0.933(6)	0.553(43)	0.461(48)	0.864(16)	0.822(22)	0.806(23)	0.168(44)	0.924(8)	0.906(9)	0.923(9)	0.847(16)	0.837(18)	0.808(20)	0.903(10)	0.863(14)
4	0.877(15)	0.881(11)	0.898(10)	1.000(0)	0.907(9)	0.855(13)	0.847(13)	0.499(46)	0.395(50)	0.775(25)	0.733(31)	0.715(33)	0.184(41)	0.834(15)	0.811(17)	0.875(12)	0.770(21)	0.752(24)	0.725(25)	0.881(11)	0.782(21)
5	0.951(4)	0.901(10)	0.990(2)	0.907(9)	1.000(0)	0.935(6)	0.928(6)	0.549(42)	0.456(51)	0.865(15)	0.823(21)	0.808(22)	0.173(43)	0.919(8)	0.901(9)	0.920(8)	0.845(15)	0.836(16)	0.806(20)	0.901(10)	0.859(14)
6	0.929(9)	0.893(10)	0.937(6)	0.855(13)	0.935(6)	1.000(0)	0.994(0)	0.561(46)	0.455(44)	0.916(9)	0.867(16)	0.853(17)	0.155(45)	0.967(3)	0.944(5)	0.928(7)	0.882(12)	0.880(12)	0.834(15)	0.893(10)	0.898(10)
7	0.924(10)	0.888(11)	0.933(6)	0.847(13)	0.928(6)	0.994(0)	1.000(0)	0.561(48)	0.454(47)	0.923(8)	0.875(16)	0.860(17)	0.154(41)	0.972(2)	0.951(5)	0.926(8)	0.892(11)	0.890(11)	0.842(15)	0.888(11)	0.906(10)
8	0.585(51)	0.570(42)	0.553(43)	0.499(46)	0.549(42)	0.561(46)	0.561(48)	1.000(0)	0.593(34)	0.707(46)	0.716(44)	0.716(47)	0.093(39)	0.619(40)	0.608(41)	0.581(42)	0.619(38)	0.613(40)	0.649(36)	0.570(42)	0.613(40)
9	0.473(45)	0.489(45)	0.461(48)	0.395(50)	0.456(51)	0.455(44)	0.454(47)	0.593(34)	1.000(0)	0.545(44)	0.558(42)	0.566(38)	0.053(45)	0.516(47)	0.502(52)	0.501(46)	0.519(45)	0.510(44)	0.553(39)	0.489(45)	0.520(45)
10	0.864(16)	0.849(16)	0.864(16)	0.775(25)	0.865(15)	0.916(9)	0.923(8)	0.707(46)	0.545(44)	1.000(0)	0.976(3)	0.962(4)	0.152(56)	0.971(3)	0.967(4)	0.904(10)	0.981(3)	0.974(4)	0.954(6)	0.849(16)	0.988(1)
11	0.827(21)	0.807(22)	0.822(22)	0.733(31)	0.823(21)	0.867(16)	0.875(16)	0.716(44)	0.558(42)	0.976(3)	1.000(0)	0.992(1)	0.133(58)	0.934(9)	0.938(8)	0.859(16)	0.982(3)	0.966(6)	0.978(4)	0.807(22)	0.973(5)
12	0.814(20)	0.792(22)	0.806(23)	0.715(33)	0.808(22)	0.853(17)	0.860(17)	0.716(47)	0.566(38)	0.962(4)	0.992(1)	1.000(0)	0.139(57)	0.918(10)	0.922(9)	0.843(18)	0.964(6)	0.952(7)	0.983(3)	0.792(22)	0.955(7)
13	0.180(59)	0.208(40)	0.168(44)	0.184(41)	0.173(43)	0.155(45)	0.154(41)	0.093(39)	0.053(45)	0.152(56)	0.133(58)	0.139(57)	1.000(0)	0.168(45)	0.166(50)	0.200(42)	0.149(45)	0.175(47)	0.128(45)	0.208(40)	0.161(45)
14	0.907(12)	0.892(11)	0.924(8)	0.834(15)	0.919(8)	0.967(3)	0.972(2)	0.619(40)	0.516(47)	0.971(3)	0.934(9)	0.918(10)	0.168(45)	1.000(0)	0.987(1)	0.935(7)	0.943(6)	0.938(8)	0.903(11)	0.892(11)	0.954(5)
15	0.889(15)	0.874(13)	0.906(9)	0.811(17)	0.901(9)	0.944(5)	0.951(5)	0.608(41)	0.502(52)	0.967(4)	0.938(8)	0.922(9)	0.166(50)	0.987(1)	1.000(0)	0.919(9)	0.946(6)	0.941(7)	0.907(11)	0.874(13)	0.955(5)
16	0.909(10)	0.987(1)	0.923(9)	0.875(12)	0.920(8)	0.928(7)	0.926(8)	0.581(42)	0.501(46)	0.904(10)	0.859(16)	0.843(18)	0.200(42)	0.935(7)	0.919(9)	1.000(0)	0.897(11)	0.887(13)	0.840(18)	0.987(1)	0.915(9)
17	0.838(19)	0.850(16)	0.847(16)	0.770(21)	0.845(15)	0.882(12)	0.892(11)	0.619(38)	0.519(45)	0.981(3)	0.982(3)	0.964(6)	0.149(45)	0.943(6)	0.946(6)	0.897(11)	1.000(0)	0.979(3)	0.955(6)	0.850(16)	0.994(0)
18	0.830(20)	0.836(18)	0.837(18)	0.752(24)	0.836(16)	0.880(12)	0.890(11)	0.613(40)	0.510(44)	0.974(4)	0.966(6)	0.952(7)	0.175(47)	0.938(8)	0.941(7)	0.887(13)	0.979(3)	1.000(0)	0.939(8)	0.836(18)	0.979(4)
19	0.810(21)	0.795(21)	0.808(20)	0.725(25)	0.806(20)	0.834(15)	0.842(15)	0.649(36)	0.553(39)	0.954(6)	0.978(4)	0.983(3)	0.128(45)	0.903(11)	0.907(11)	0.840(18)	0.955(6)	0.939(8)	1.000(0)	0.795(21)	0.942(7)
20	0.889(12)	1.000(0)	0.903(10)	0.881(11)	0.901(10)	0.893(10)	0.888(11)	0.570(42)	0.489(45)	0.849(16)	0.807(22)	0.792(22)	0.208(40)	0.892(11)	0.874(13)	0.987(1)	0.850(16)	0.836(18)	0.795(21)	1.000(0)	0.867(14)
$^{21}$	0.848(18)	0.867(14)	0.863(14)	0.782(21)	0.859(14)	0.898(10)	0.906(10)	0.613(40)	0.520(45)	0.988(1)	0.973(5)	0.955(7)	0.161(45)	0.954(5)	0.955(5)	0.915(9)	0.994(0)	0.979(4)	0.942(7)	0.867(14)	1.000(0)

1 = index nonSmear s = 0.4, 2 = index nonSmear s = 0, 3 = index HYP1 s = 0, 4 = SF HYP1 s = 0.75, 5 = SF HYP1 s = 0, 6 = SF HYP5 s = 0.5, 7 = SF HYP5 s = 0, 8 = spec.proj. nonSmear, 9 = fermionic (from disc. loops), 10 = GF flow time  $t_0$ , 11 = GF flow time  $2t_0$ , 12 = GF flow time  $3t_0$ , 13 = field theor. nonSmear, 14 = field. theor. HYP10, 15 = field. theor. HYP40, 16 = field. theor. APE10, 17 = field. theor. APE30, 18 = field theor.cool. tol. 5%, 19 = field theor. cool.tol. 1%, 20 = field theor. cool.10 steps, 21 = field theor.cool. 30 steps



## Plot of correlations





#### Results

Histograms Cool. vs.

APE/HYP

Pair comparisons

#### Correlations

Topo. susc.

Conclusions

# method 2



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## Correlation towards the continuum limit





0.18

0.2









# Topological susceptibility



used only pion masses  $m_{\pi} \leq 400 \text{ MeV}$ 



spectral projectors:  $r_0 \Sigma^{1/3} = 0.646(59)$ 

Free-level 
$$\chi$$
PT fit:  $r_0^4 \chi = \frac{r_0^3 \Sigma \cdot r_0 \mu_R}{2}$ 

compare to direct determination

 $[\mathrm{K.C.,\ E.\ Garcı́a\ Ramos,\ K.\ Jansen,\ JHEP\ 10(2013)175}]$ 

$$r_0 \Sigma_{cont,N_f=2+1+1}^{1/3} = 0.680(29)$$

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fermionic:  $r_0 \Sigma^{1/3} = 0.644(20)$ 



# Topological susceptibility





spectral projectors:  $r_0 \Sigma^{1/3} = 0.612(57)$ 

Tree-level 
$$\chi$$
PT fit:  $r_0^4 \chi = \frac{r_0^3 \Sigma \cdot r_0 \mu_R}{2}$ 

compare to direct determination

 $[\mathrm{K.C.,\ E.\ Garcı́a\ Ramos,\ K.\ Jansen,\ JHEP\ 10(2013)175}]$ 

$$r_0 \Sigma_{cont,N_f=2+1+1}^{1/3} = 0.680(29)$$

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fermionic:  $r_0 \Sigma^{1/3} = 0.622(19)$ 



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Cleanness

Conclusions

# My subjective ranking of theoretical cleanness

From the cleanest to the problematic ones:

- index of the overlap Dirac operator spectral flow field theoretic with gradient flow spectral projectors
- 2. field theoretic with smearing
- 3. field theoretic with cooling
- 4. fermionic from disconnected loops

The theoretical progress of recent years (spectral projectors, gradient flow) makes the non-clean definitions rather unattractive – there is basically no need any more to use them for reasons of (relatively small) computational intensity.





# Conclusions



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Cleanness

 $\operatorname{Conclusions}$ 

- There are many definitions of the topological charge on the lattice
- No clear answer which definition to use
- We observe high correlations between different definitions...
- ...and these correlations seem to increase towards the continuum limit
- Perhaps the best compromise between theoretical cleanness and computational cost: field theoretic with gradient flow

[M. Bruno, S. Schaefer, R. Sommer, arXiv:1406.5363 [hep-lat]]

Thank you for attention!

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