



Exploring the QCD phase diagram with fluctuations of conserved charges

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BNL-Bi-CCNU Collaboration:

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Definitions

Expansion of the pressure:

$$\frac{p}{T^4} = \sum_{i,j,k=0}^{\infty} \frac{1}{i!j!k!} \chi_{ijk,0}^{BQS} \left(\frac{\mu_B}{T}\right)^i \left(\frac{\mu_Q}{T}\right)^j \left(\frac{\mu_S}{T}\right)^k$$

$X = B, Q, S$: conserved charges

Lattice

$$\chi_n^X = \frac{\partial^n [p/T^4]}{\partial (\mu_X/T)^n} \Big|_{\mu_X=0}$$

generalized susceptibilities

\Rightarrow only at $\mu_X = 0$!

Experiment

$$\begin{aligned} VT^3 \chi_2^X &= \langle (\delta N_X)^2 \rangle \\ VT^3 \chi_4^X &= \langle (\delta N_X)^4 \rangle - 3 \langle (\delta N_X)^2 \rangle^2 \\ VT^3 \chi_6^X &= \langle (\delta N_X)^6 \rangle - 15 \langle (\delta N_X)^4 \rangle \langle (\delta N_X)^2 \rangle \\ &\quad + 30 \langle (\delta N_X)^2 \rangle^3 \end{aligned}$$

cumulants of net-charge fluctuations

$$\delta N_X \equiv N_X - \langle N_X \rangle$$

\Rightarrow only at freeze-out $(\mu_f(\sqrt{s}), T_f(\sqrt{s}))$!

Motivations

Explore the QCD phase diagram

- Analyze higher order cumulants and test universal scaling behavior
→ Explore the QCD phase diagram

This talk

Analyze freeze-out conditions

- Match various cumulant ratios of measured fluctuations to QCD
→ determine freeze-out parameter

BNL-Bielefeld, PRL 109 (2012) 192302.

Identify the relevant degrees of freedom

- Compare (lattice) QCD fluctuations to various hadronic/quasiparticle models:
 - deconfinement vs. chiral transition (melting of open strange/charm hadrons)
 - evidence for experimentally not yet observed hadrons



See talk by
H.-T. Ding

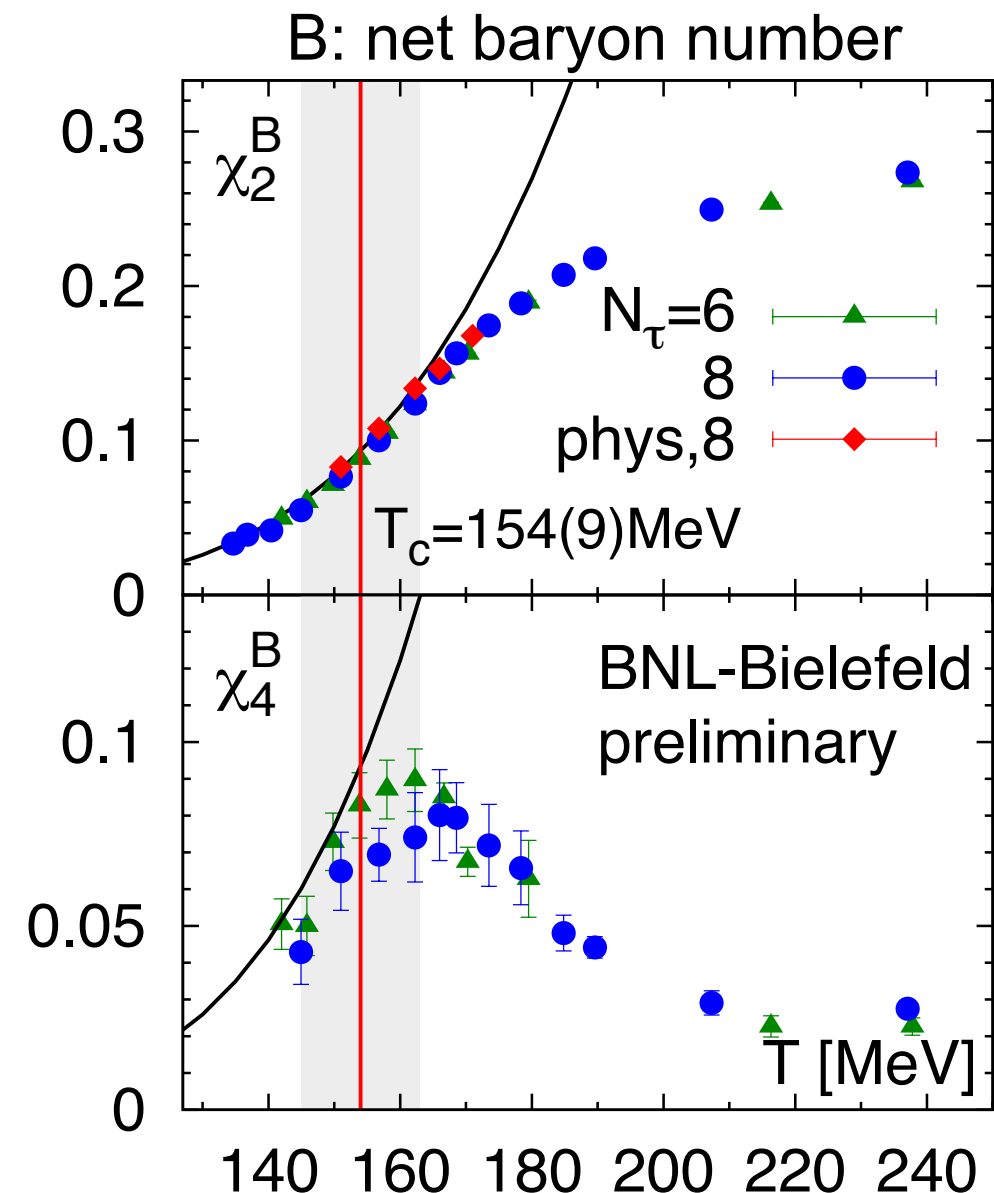
BNL-Bielefeld, PRL 111 (2013) 082301;

BNL-Bielefeld-CCNU, arXiv: 1404.4043, arXiv: 1404.6511.

The lattice setup

Lattice parameters:

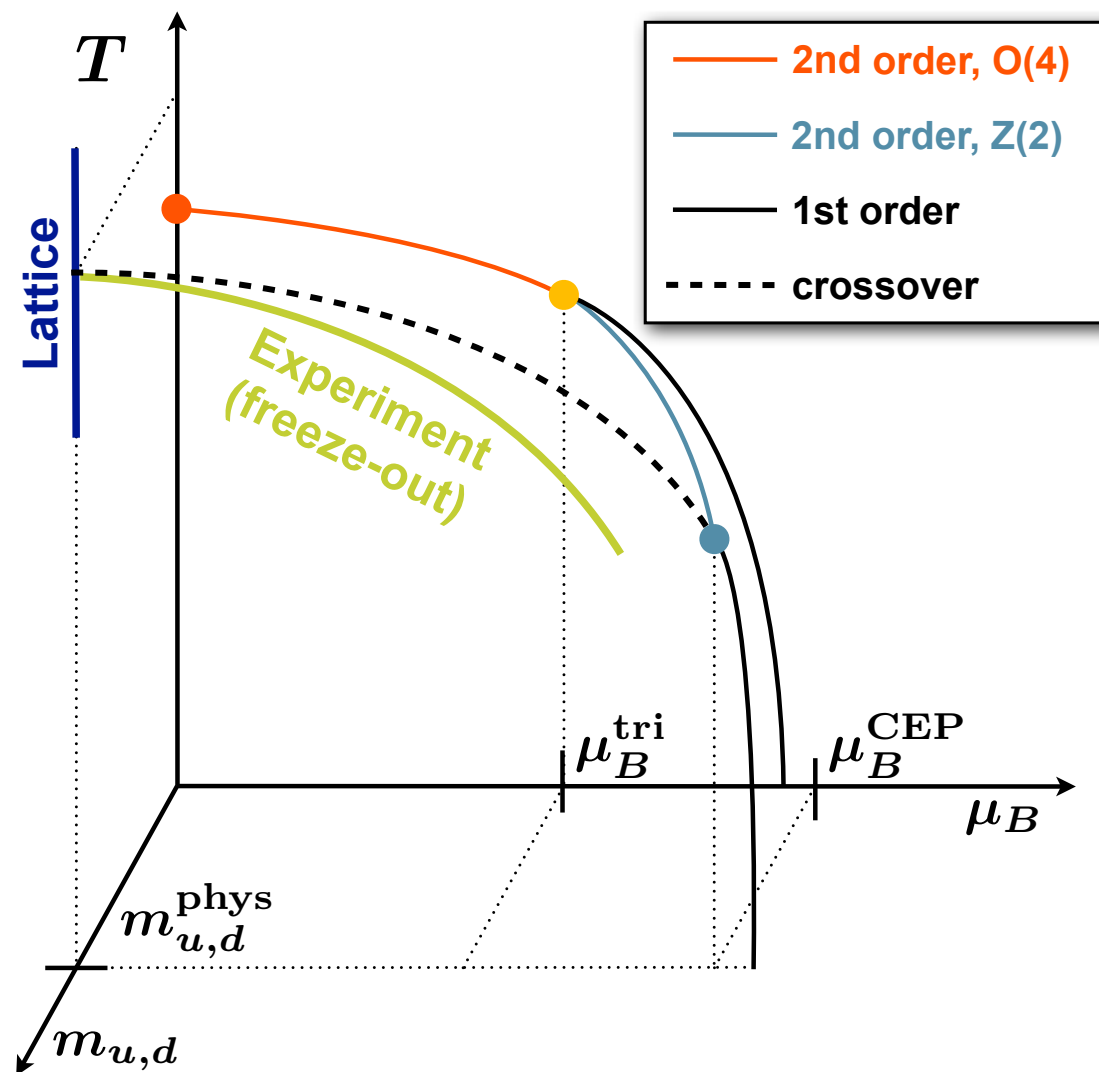
- (2+1)-flavor of highly improved staggered fermions (HISQ-action)
- a set of different lattice spacings ($N_\tau = 6, 8, 12$)
- two different pion masses: $m_\pi = 140, 160 \text{ MeV}$
- high statistics: $(10 - 16) \times 10^3$ configurations



⇒ statistical and systematical errors are under control

⇒ In general: find good agreement with HRG model for $T < 155 \text{ MeV}$

QCD critical behavior



assume scaling hypothesis
for the free energy:

$$f = f_s(t, h) + \text{regular}$$

with

$$\lambda f_s(t, h) = f_s(\lambda^{a_t} t, \lambda^{a_h} h)$$

\Rightarrow in the chiral limit ($h=0$):

$$f \sim A_{\pm} |t|^{2-\alpha} + \text{regular}$$

critical exponent:

α

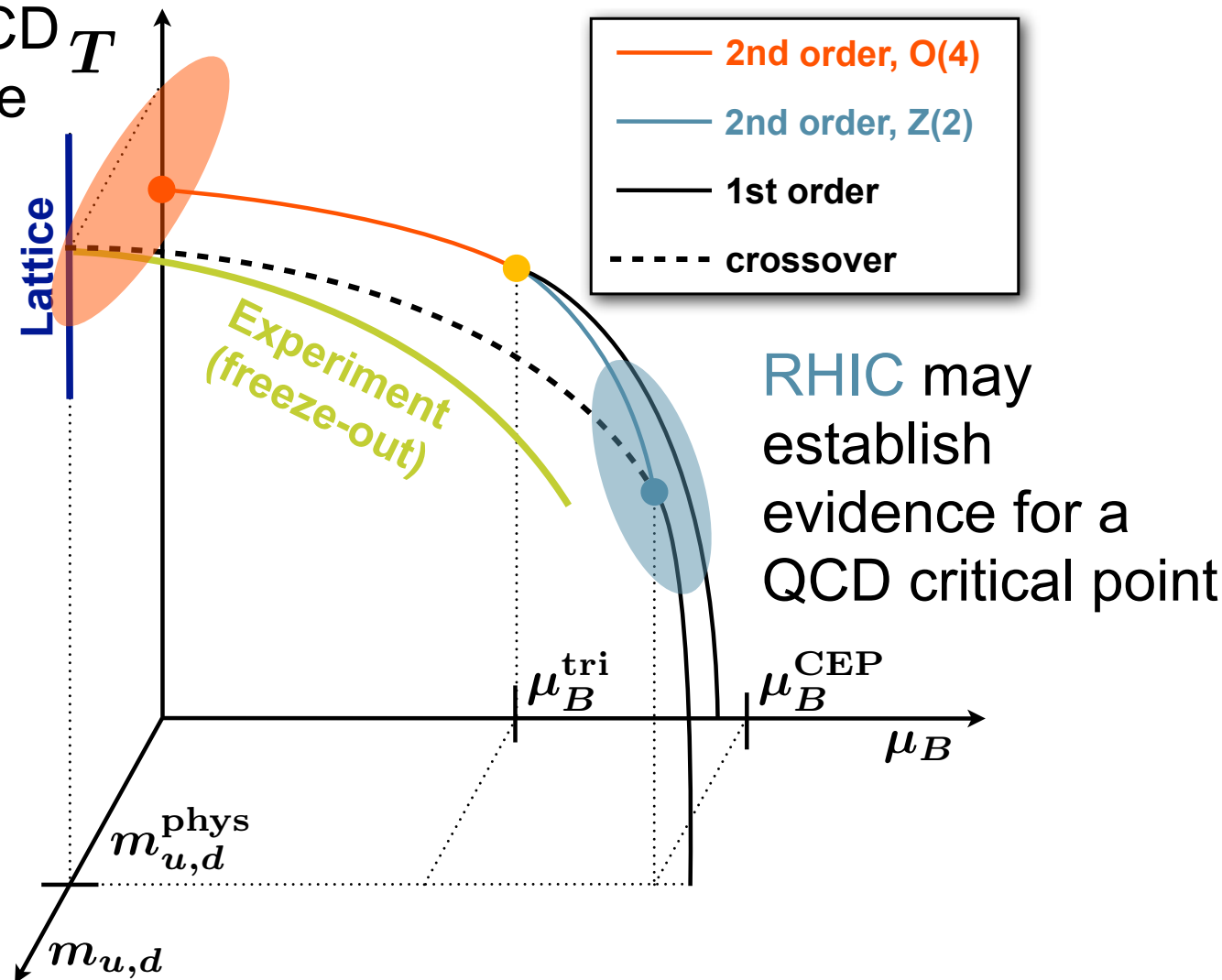
$$\text{O(4)} \quad -0.213$$

$$\text{Z(2)} \quad +0.107$$

\Rightarrow at $\mu_B = 0$, 4th order
cumulants develop a
cusp, 6th order
cumulants diverge

QCD critical behavior

LHC may establish contact with the QCD chiral phase transition



⇒ analyze universal scaling behavior

assume scaling hypothesis for the free energy:

$$f = f_s(t, h) + \text{regular}$$

with

$$\lambda f_s(t, h) = f_s(\lambda^{a_t} t, \lambda^{a_h} h)$$

⇒ in the chiral limit ($h=0$):

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critical exponent:

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$$\text{O}(4) \quad -0.213$$

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⇒ at $\mu_B = 0$, 4th order cumulants develop a cusp, 6th order cumulants diverge

QCD critical behavior

matching scaling fields to QCD at $\mu_B = 0$:

$$h = \frac{m_q}{h_0} \quad t = \frac{1}{t_0} \left(\left(\frac{T - T_c}{T_c} \right) + \kappa \left(\frac{\mu_B}{T} \right)^2 \right)$$

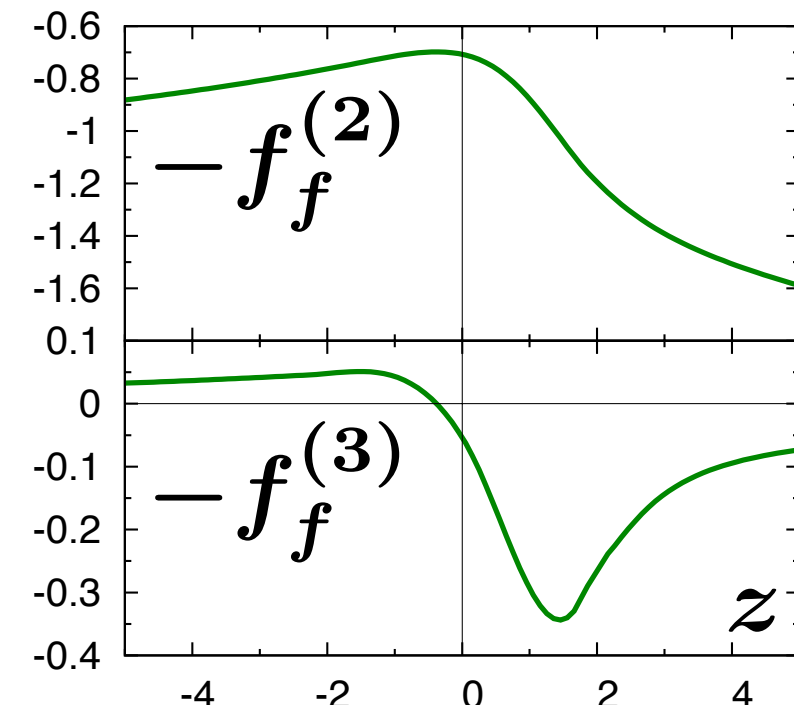
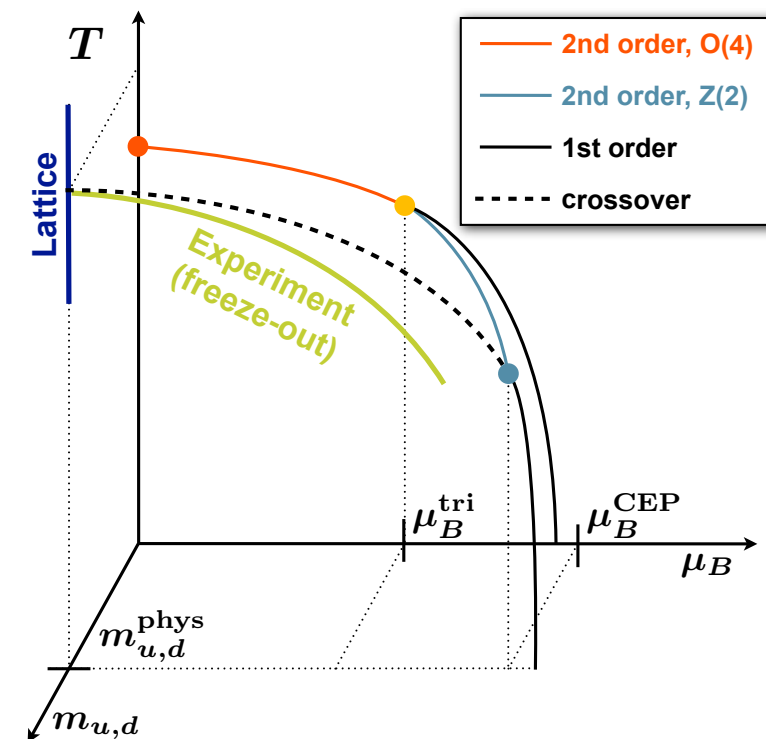
controlled by non-universal normalization constants t_0, h_0, κ

$$\frac{p}{T^4} = -h^{(2-\alpha)/\beta\delta} \underbrace{f_f(t/h^{1/\beta\delta})}_{\text{(universal scaling function)}} - \underbrace{f_r(V, T, \vec{\mu})}_{\text{(regular part)}}$$

f_f is function of scaling variable: $z = t/h^{1/\beta\delta}$

\Rightarrow critical behavior of cumulants (at $\mu_B = 0$):

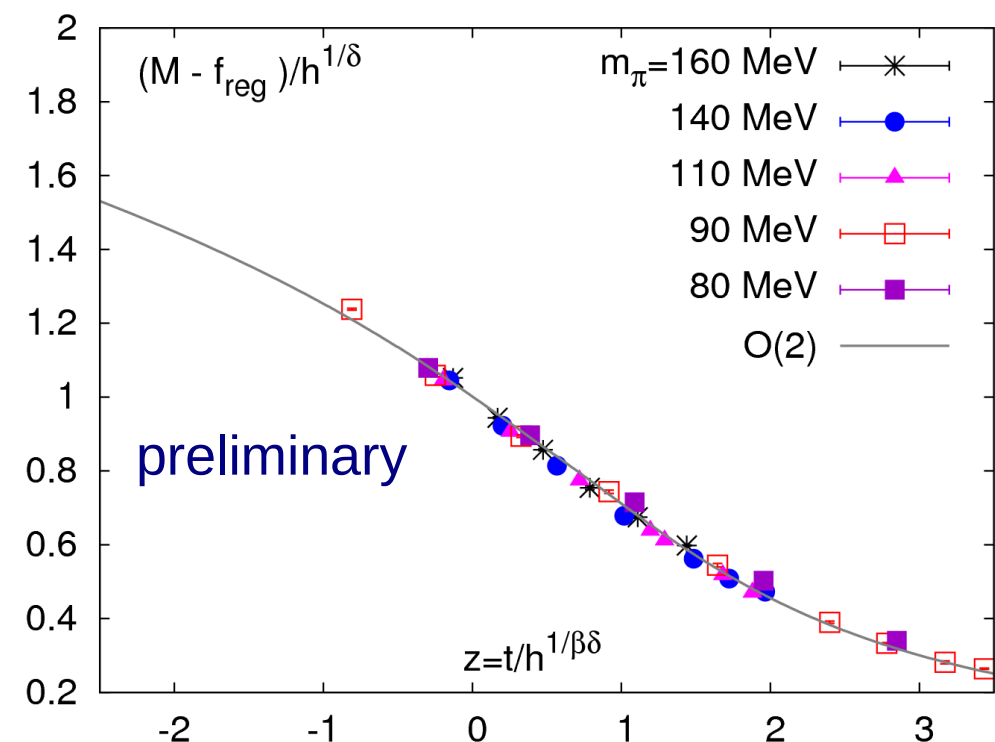
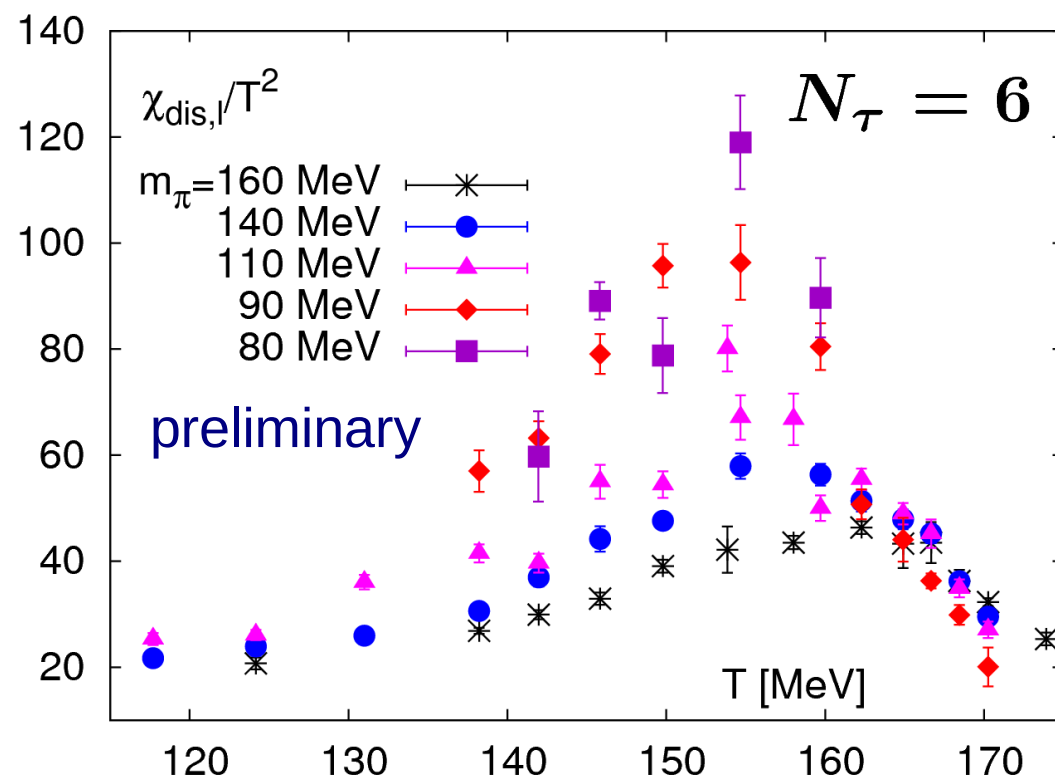
$$\chi_B^{(n)} \sim m_q^{(2-\alpha-n/2)/\beta\delta} f_f^{(n/2)}(t/h^{1/\beta\delta})$$



Karsch, Engels, PRD 85 (2012) 094506

QCD critical behavior

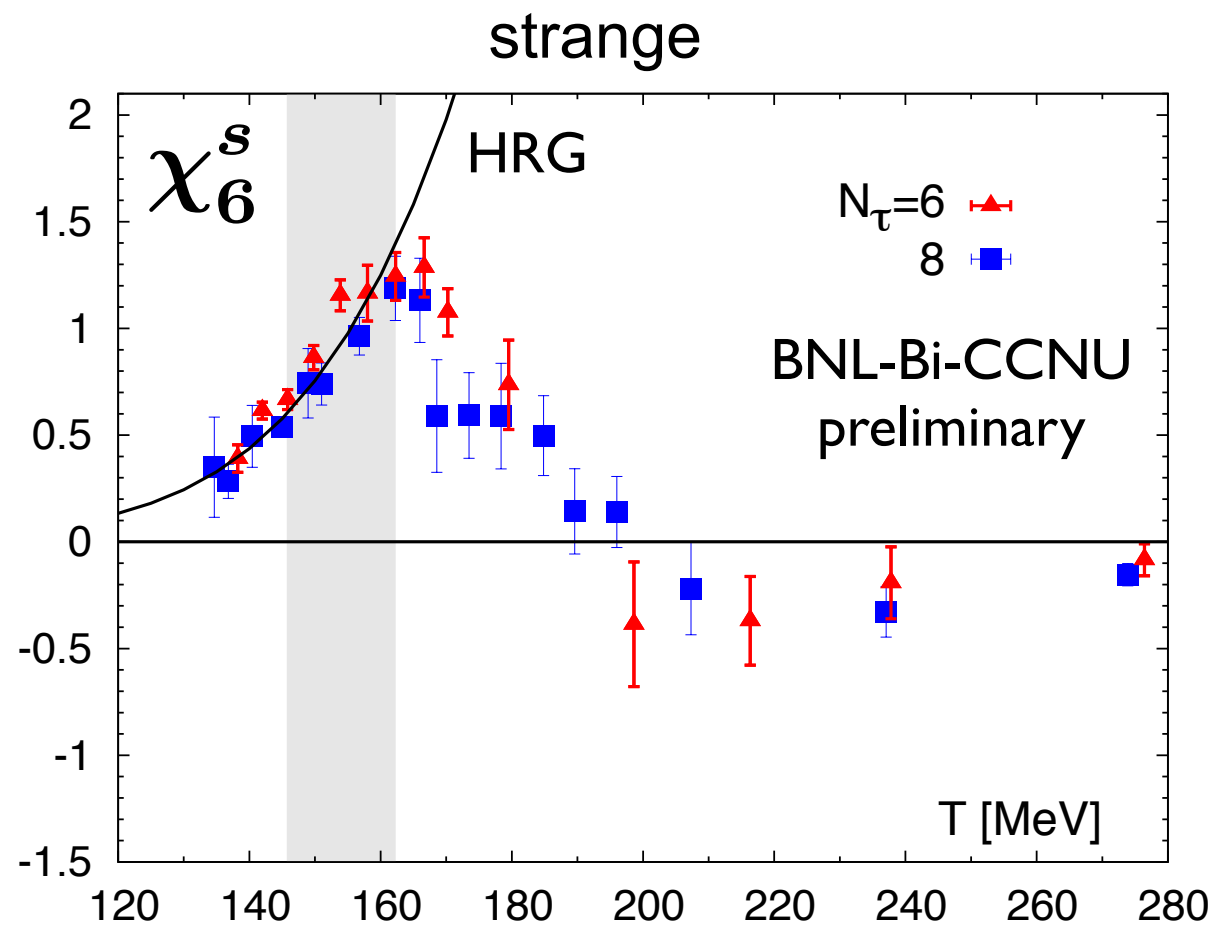
evidence for universal scaling behavior with HISQ from chiral condensate and chiral susceptibility (H.-T. Ding et al. Lattice 2013)



$$M = m_s \langle \bar{\psi} \psi \rangle / T^4$$

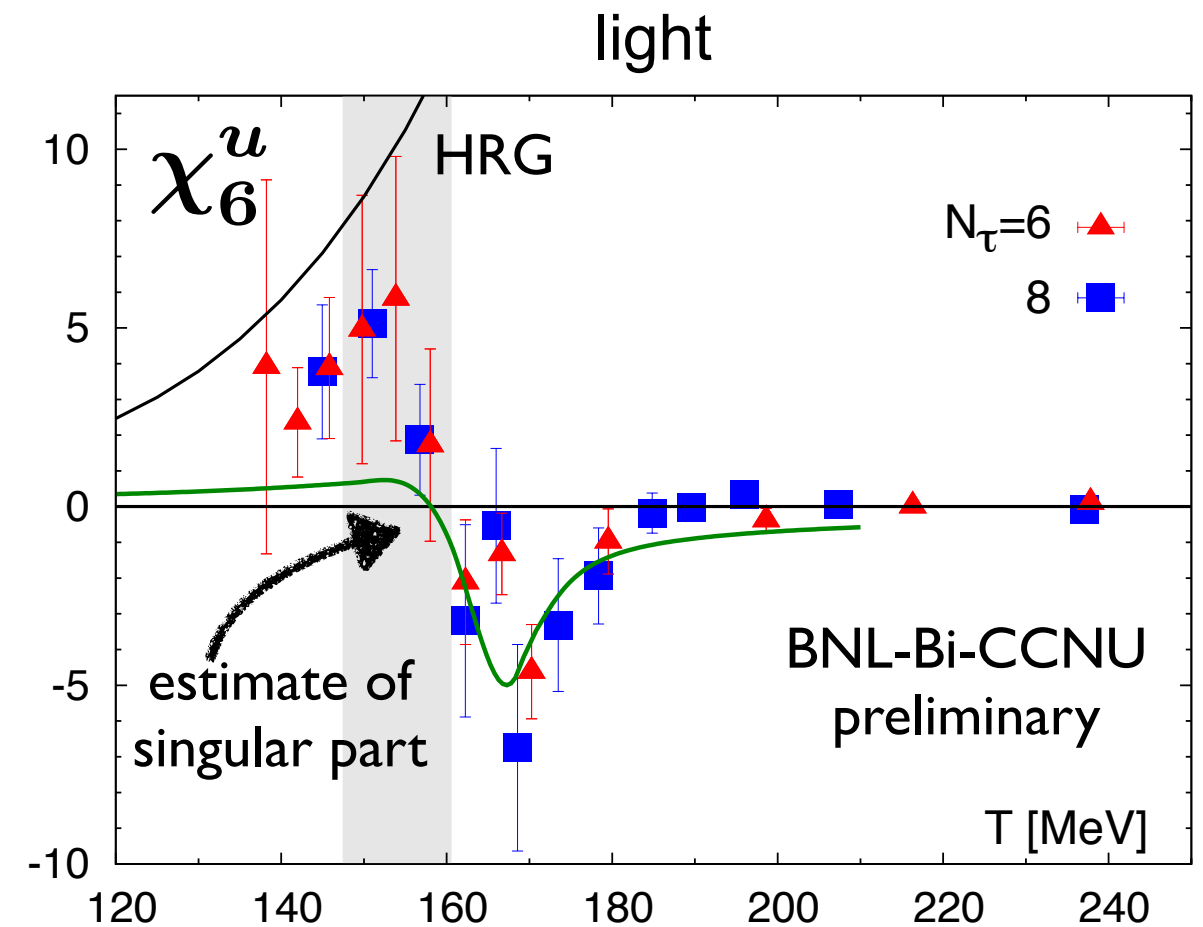
⇒ scaling region extends to physical pion mass

QCD critical behavior



⇒ no evidence for typical $O(4)$ singular structure

⇒ regular contribution dominates



⇒ clear evidence for typical $O(4)$ singular structure

⇒ regular and singular contribution

QCD critical behavior

some universal numbers:

width of the transition
region (as seen by χ_6^B):

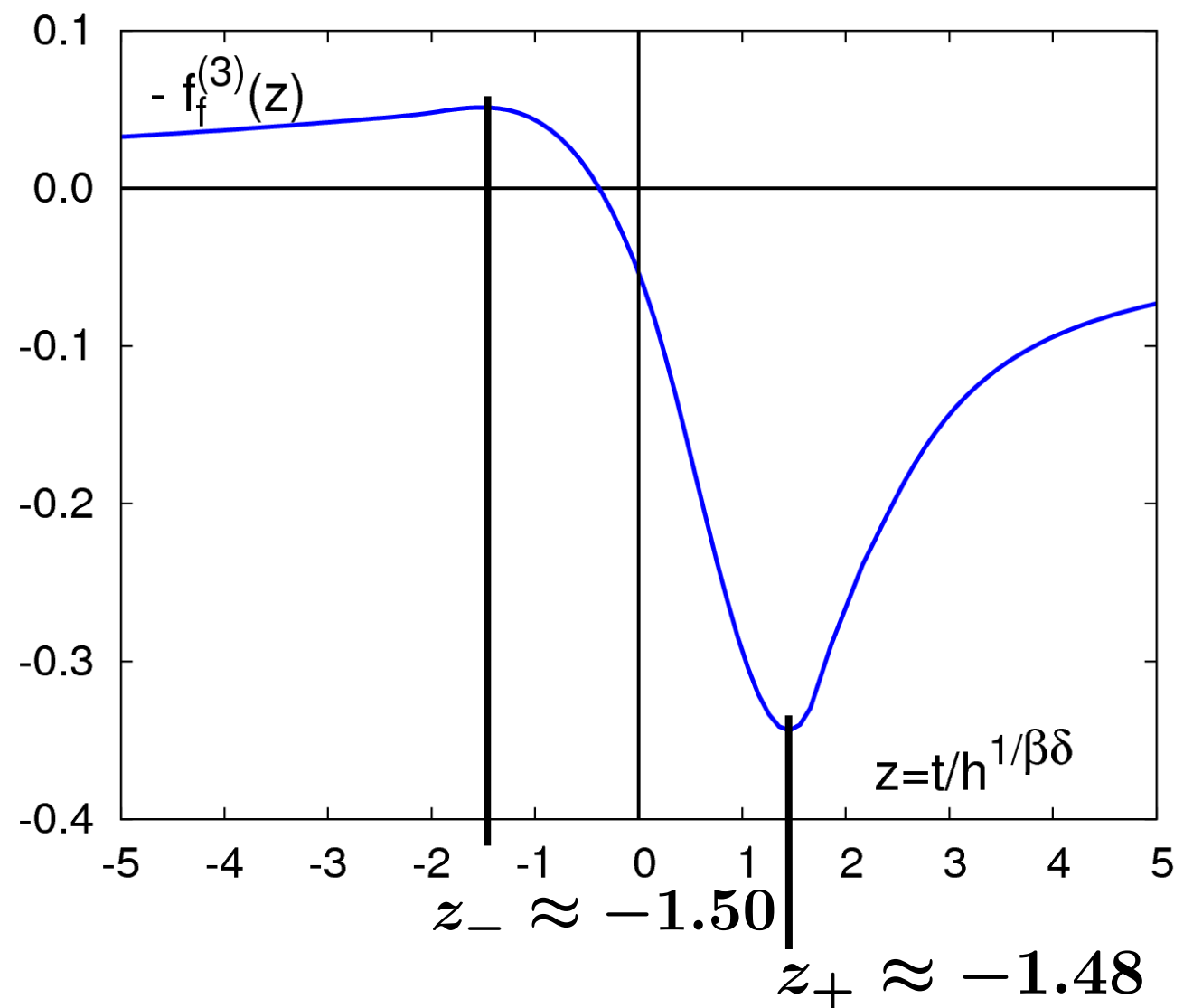
$$\Delta z = z_+ - z_- \approx 3$$

\Rightarrow

$$T_+ - T_- = \frac{1}{\Delta z} \frac{t_0 T_c}{h_0^{1/\beta\delta}} \left(\frac{m_l}{m_s} \right)^{1/\beta\delta}$$

at the physical point:

$$T_+ - T_- \approx 0.2 T_c$$



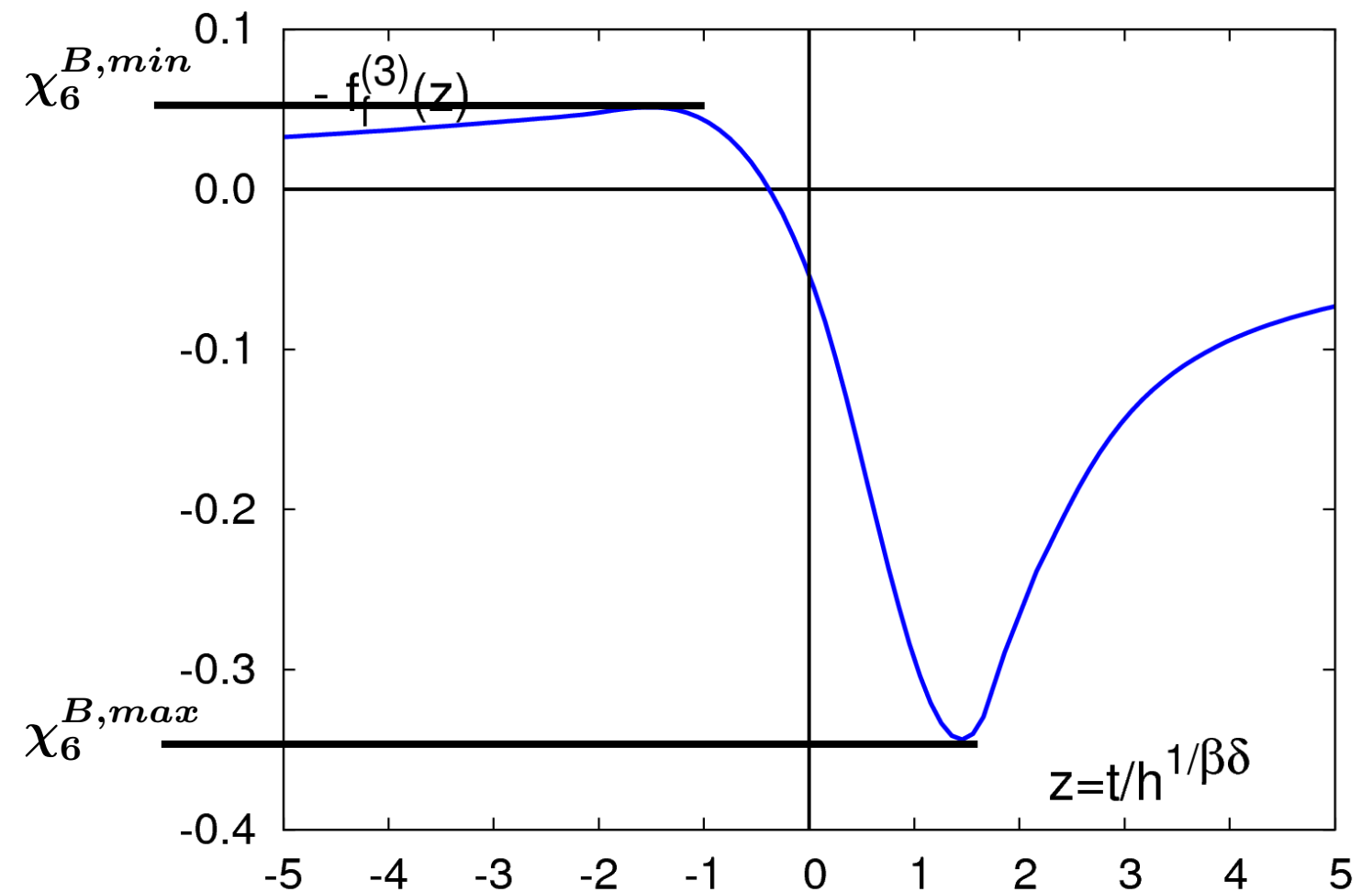
QCD critical behavior

some universal numbers:

ratio of minimum to maximum:

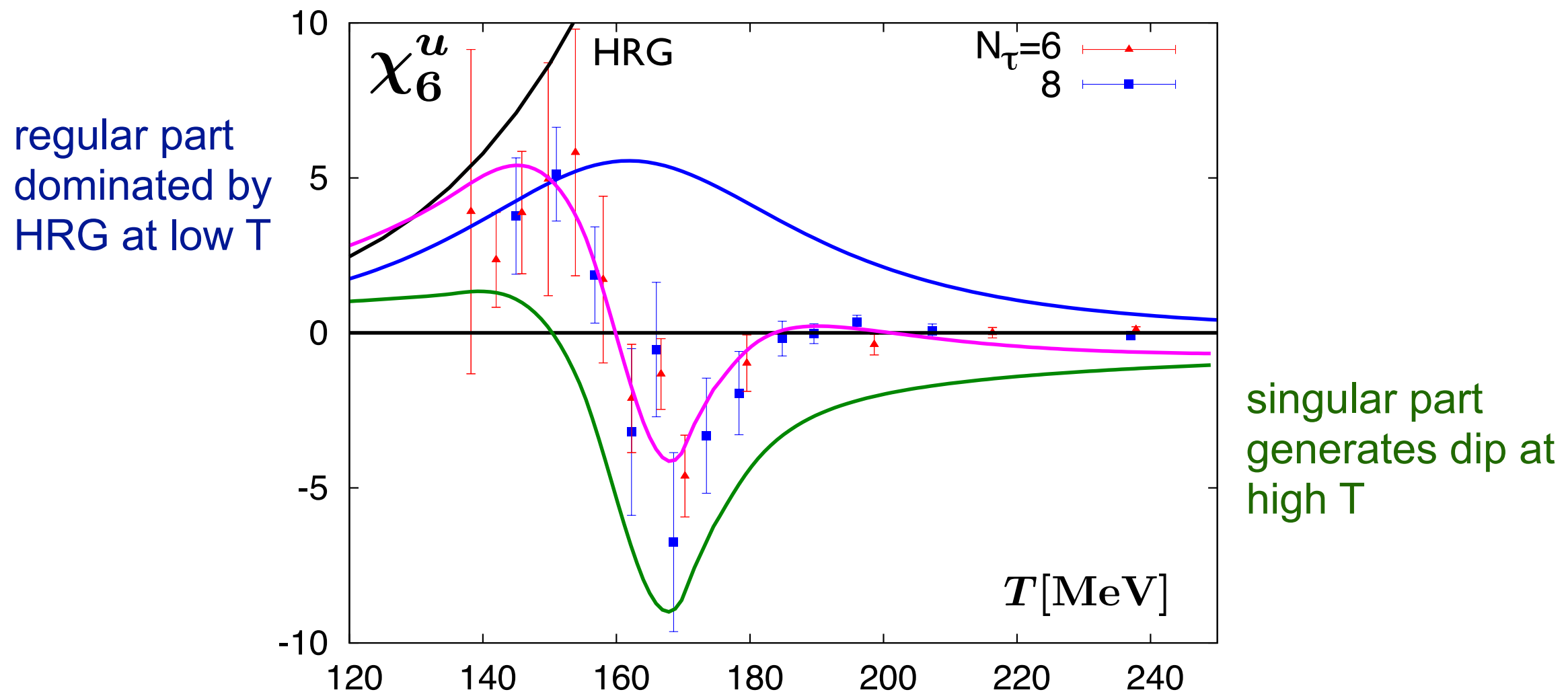
$$\chi_6^{B,min} / \chi_6^{B,max} \approx -6.7$$

\Rightarrow depth of minimum at high T
fixes maximal singular
contribution at high T



QCD critical behavior

on the interplay of regular and singular contributions (so far a guess and not a fit)

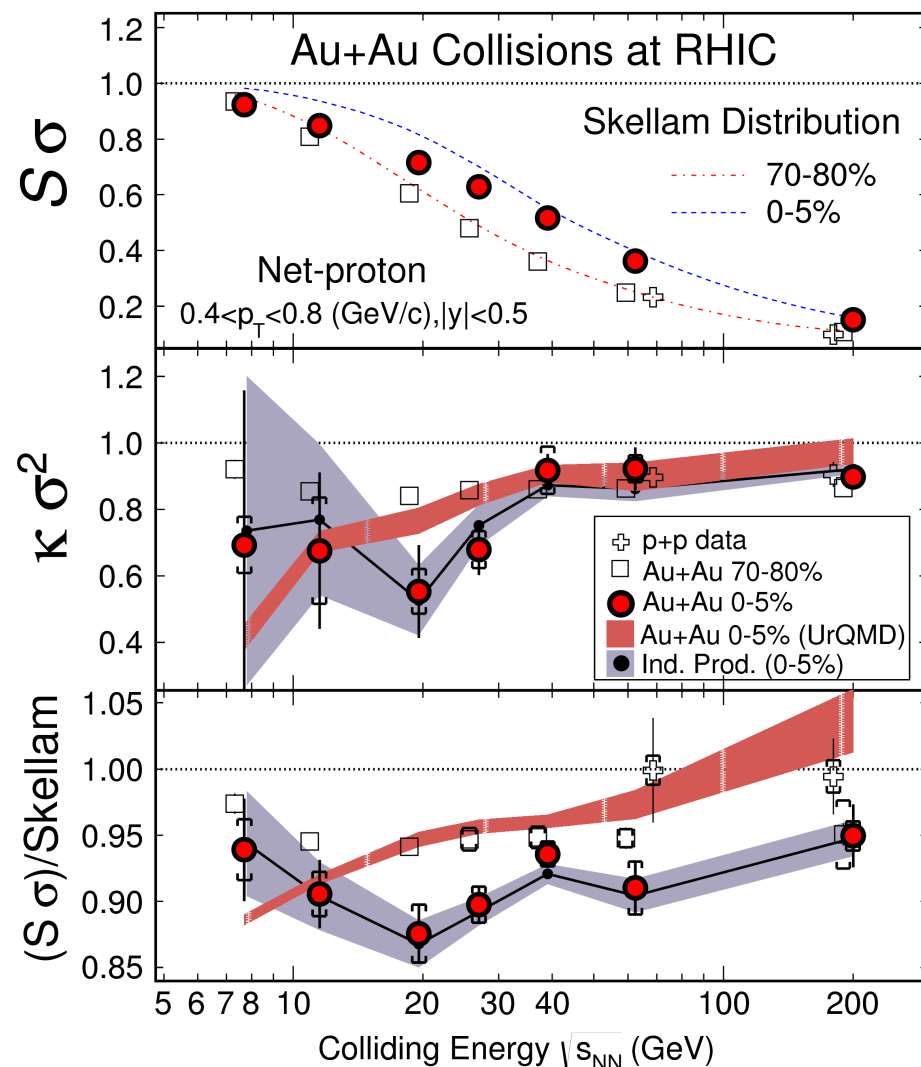


total = singular + regular

⇒ approximate agreement with HRG and sensibility to $O(4)$ scaling are not in disagreement

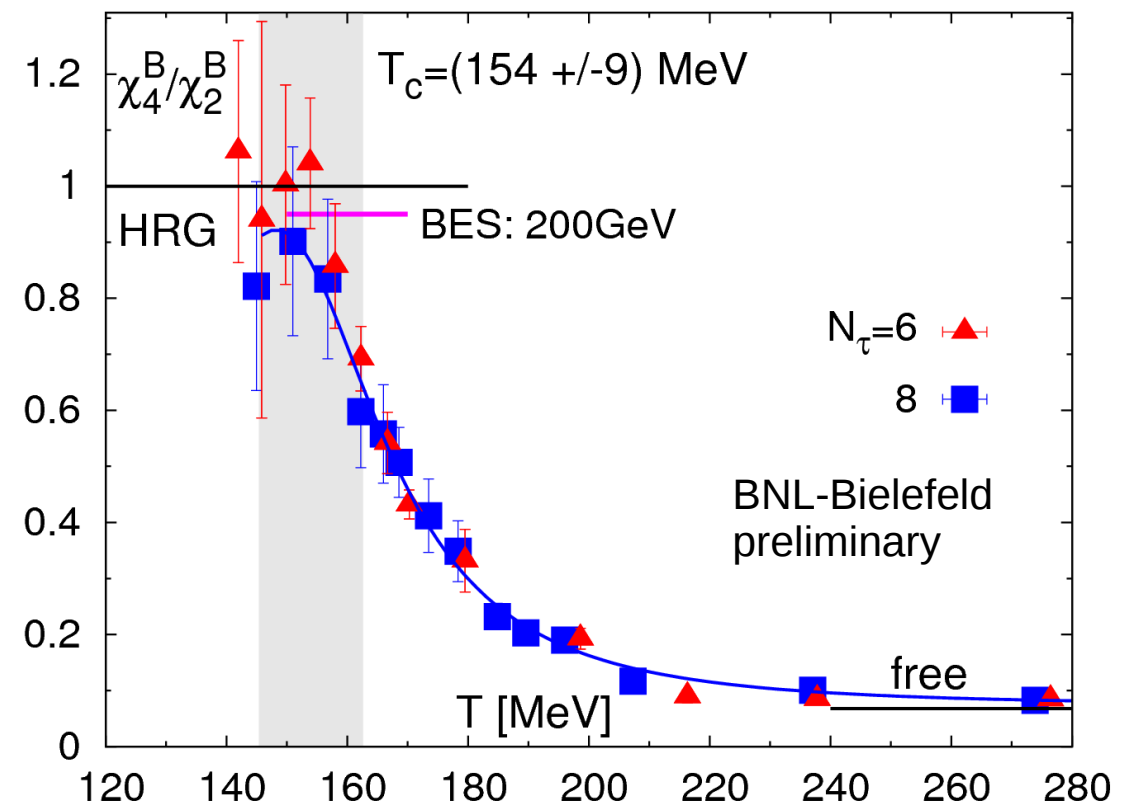
Critical point search

proton number fluctuations



STAR, arXiv: 1309.5681

relative strength of the NLO correction to the pressure is controlled by χ_4^B / χ_2^B



$$T < T_c : \quad 0.8 \leq \chi_4^B / \chi_2^B \leq 1.0$$

$$\Rightarrow \frac{\mathcal{O}(\mu_B^4) \text{ contribution to pressure}}{\mathcal{O}(\mu_B^2) \text{ contribution to pressure}} < 1$$

for $\mu_B/T \lesssim 3.5$

Critical point search

$$\frac{p}{T^4} = \frac{1}{VT^3} \ln Z(V, T, \mu_B) = \sum_{n, \text{even}} \frac{1}{n!} \chi_n^B \left(\frac{\mu_B}{T} \right)^n$$

\Rightarrow consider radius of convergence

$$\left(\frac{\mu_B}{T} \right)_{crit} = \lim_{n \rightarrow \infty} \sqrt{\left| \frac{(n+2)! \chi_n^B}{n! \chi_{n+2}^B} \right|}$$

\Rightarrow basic quantities

$$\chi_n^B / \chi_{n+2}^B$$

=1 for HRG

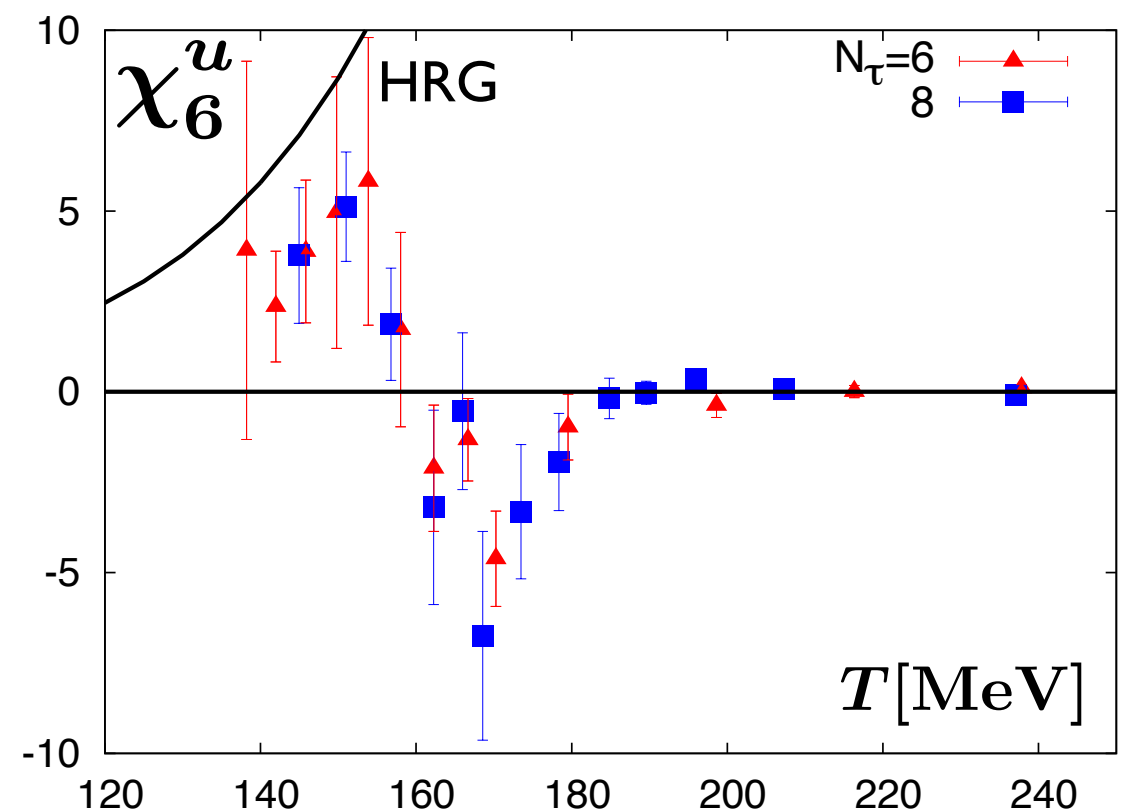
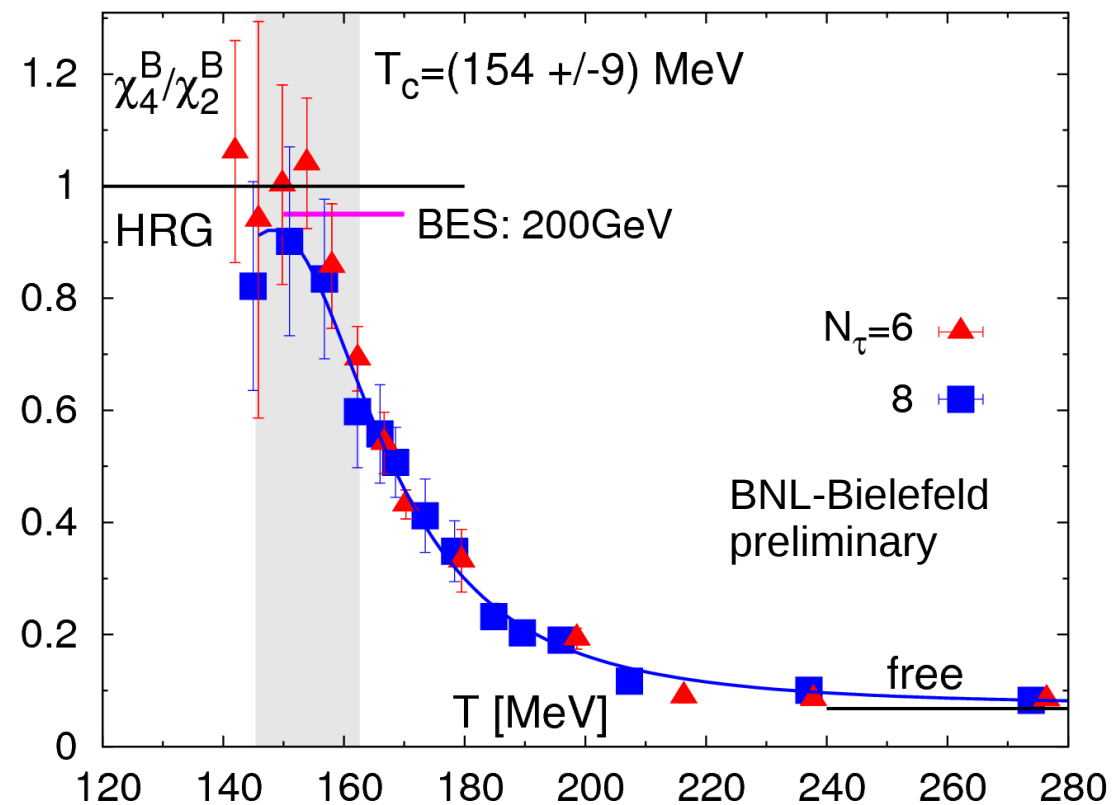
need to deviate from HRG like n^2
to obtain finite radius of convergence

\Rightarrow singularity on the real axis **only if**

$$\chi_n^B > 0 \text{ for all } n > n_0$$

Critical point search

we find so far no evidence for large enhancement over HRG for $T < T_c$ ($n \leq 6$)



\Rightarrow this suggests large μ_B^{crit}



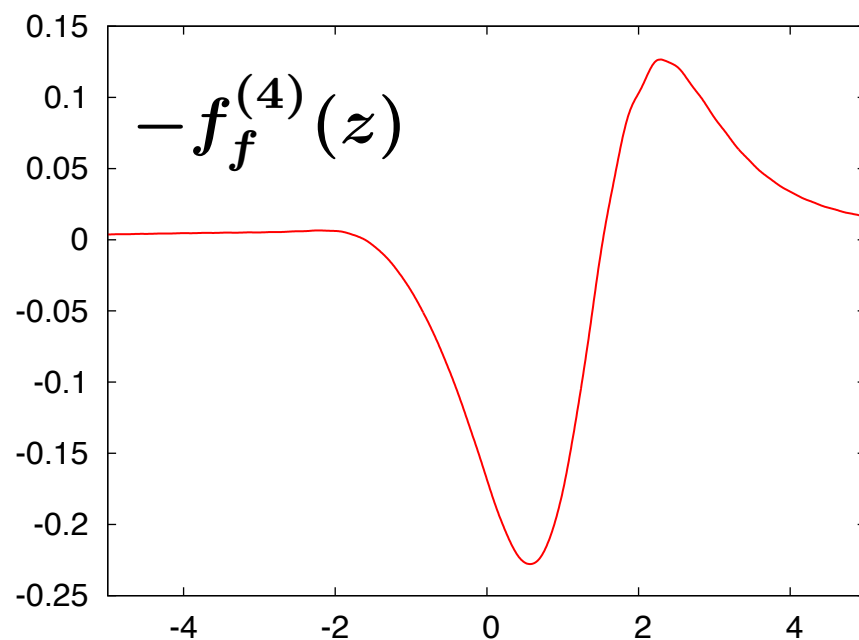
Summary

- Approximate agreement with HRG model calculations at freeze-out and sensitivity to $O(4)$ singular behavior are not inconsistent with each other
 - Higher order fluctuations need to deviate from HRG like n^2 in order to obtain a finite radius of convergence
-
- 6th order cumulants are sensitive to $O(4)$ scaling but will pick up only a small singular contribution at low T . This favors estimates for the location of a critical end point at large baryon chemical potentials

QCD critical behavior at $\mu_B > 0$

$$\begin{aligned} \mu_B > 0 : \quad \chi_{4,\mu}^B = & -3(2\kappa_B t_0^{-1})^2 h^{-\alpha/\Delta} f_f^{(2)}(z) \\ & -6(2\kappa_B t_0^{-1})^3 (\hat{\mu}_B^c)^2 h^{-(1+\alpha)/\Delta} f_f^{(3)}(z) \\ & -(2\kappa_B t_0^{-1} \hat{\mu}_B^c)^4 h^{-(2+\alpha)/\Delta} f_f^{(4)}(z) + \text{regular} \end{aligned}$$

dominates in the chiral
limit or if $\hat{\mu}_B^c > 0 \gtrsim 1$



\Rightarrow close to T_c :

$$\chi_4^B(\mu_B) < 0$$

B.Friman, FK, K.Redlich,V.Skokov,
Eur. Phys. J. C71, 1694 (2011)

\Rightarrow mapping of scaling variables non trivial

M. Stephanov, PRL 107 (2011) 052301