Physical and cut-off effects of heavy charm-like sea quarks

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Results

Motivation

Charm effects

- Estimate physical effects of the charm quark in QCD
- \blacktriangleright At large M effective theory in powers of 1/M
- $\blacktriangleright~M_{\rm c}\simeq 12 M_{\rm s}$: cut-off effects from charm can be large
- ► Here: study these effects for N_f = 2 O(a) improved Wilson fermions at a mass below and close to charm
- Can $1/M^2$ effects be measured?



Results

Effective field theory

Expansion in $(E/M)^n$: EFT for $E \ll M$

- Only virtual effects of quark with mass M No states with explicit heavy quark (the HQET part)
- Effective Lagrangian (here for $N_{\rm f} \rightarrow N_{\rm f} 1$)

$$\begin{aligned} \mathscr{L}_{\text{QCD}}^{(N_{\text{f}}-1)} &= \mathscr{L}_{\text{QCD}}^{(N_{\text{f}}-1)}(\psi_{\text{light}}, \bar{\psi}_{\text{light}}, A_{\mu}; \tilde{g}_{0}(M), m(M)) \\ &+ \frac{1}{M} \mathscr{L}_{\text{Pauli}} + \frac{1}{M^{2}} \mathscr{L}_{6} \\ \mathscr{L}_{\text{Pauli}} &= \frac{g^{2l}(M)}{M} \bar{\psi}_{\text{light}} \sigma_{\mu\nu} F_{\mu\nu} m_{\text{light}} \psi_{\text{light}} \\ &+ \mathsf{NP} \times \bar{\psi}_{\text{light}} \sigma_{\mu\nu} F_{\mu\nu} \psi_{\text{light}} \end{aligned}$$

NP expected to be: ${
m NP}=M^{-\gamma}\,,\;\gamma>0$ ($\gamma\geq 1$?)

• Coupling $\tilde{g}_0^2(M)$ or $\bar{g}^2(M)$ drops out for ratios $R(M) = \frac{t_0}{w_0^2}, \ \frac{r_1}{r_0}, \ldots$



Pick observables with a strong dependence on $N_{\rm f}$ Wilson flow: t_0 [Lüscher, arXiv:1006.4518] and w_0 [Borsanyi et al., arXiv:1203.4469]

$$t_0$$
 : $\mathcal{E}(t_0) = 0.3$, $\mathcal{E}(t) = t^2 \langle E(x,t) \rangle$
 w_0 : $w_0^2 \mathcal{E}'(w_0^2) = 0.3$

Static force: r_0 [Sommer, hep-lat/9310022] and r_1 [Bernard et al., hep-lat/0002028]

$$r^2 F(r)|_{r=r_c} = c, \quad r_0 \equiv r_{1.65}$$

RGI mass, fixed in $\,\mathrm{MeV}$ using F_K

$$M \equiv M_{\rm RGI} = \frac{M}{\overline{m}(\mu)} \frac{Z_{\rm A}(1 + \tilde{b}_{\rm A} am)}{Z_{\rm P}(\mu)(1 + \tilde{b}_{\rm P} am)} m \,, \qquad m \equiv$$



 $m_{\rm PCAC}$

Ensembles		

β	$a [\mathrm{fm}]$	$T \times L^3$	$M \; [\mathrm{MeV}]$	kMDU
5.3	0.0658(10)	64×32^3	200(4)	1
5.3		64×32^3	410(8)	2
5.5	0.0486(7)	120×32^3	200(4)	8
5.5		120×32^3	407(7)	8
5.5		96×48^3	780(14)	4
5.7	0.038	192×48^3	389(8)	4
5.7		192×48^3	745(16)	4

(a from $F_{
m K}$ [ALPHA, arXiv:1205.5380] and PT for eta=5.7)



Autocorrelation times



Virotta, arXiv:1009:5228]

Ratios		

Global fits of cut-off effects

Global fits to ratios $R=t_0/w_0^2\,,\;r_1/r_0\,,\;r_0^2/t_0$

continued lines, a^2 effects at M = 0 fixed from [Sommer, arXiv:1401.3270]: $s \approx 2$ for $R = t_0/w_0^2$, $s \approx 15$ for $R = r_0^2/t_0$

$$R = R(M) + s\frac{a^2}{8t_0} \left(1 + k_1M + k_2M^2\right)$$

dashed lines:

$$R = R(M) + k_0 \frac{a^2}{8t_0} + k_1 \frac{a^2}{8t_0} M$$

(continuum values for $N_{\rm f}=$ 0, $N_{\rm f}=$ 2 at physical point from M. Bruno)



$R = t_0 / w_0^2$		



Cancellation between cut-off effects, at a = 0.049 fm, M = 0.8 GeV: term without M: 3.3%; with M: -2.4%





$R = r_1 / r_0$		





$R = r_0^2 / t_0$		



(coefficient of M^2 effects $k_2 = 0$)



Introduction

Simulations

Results

Conclusions and outlook

The M - dependence



Conclusions and outlook

Decoupling rescaled	to $N_{\rm f}=1 \rightarrow 0$
relative effects ·	

$$\frac{1}{N_{\rm f}} \frac{\mathcal{O}(M) - \mathcal{O}(\infty)}{\mathcal{O}(\infty)} \qquad N_{\rm f} = 2$$

R	$M \rightarrow$	$M_{\rm c} = 1.6 {\rm GeV}$	$0.8{ m GeV}$	$0.4{ m GeV}$	$0.2{ m GeV}$	0
$\sqrt{t_0}/w_0$		0.14 - 0.3%	0.62(19)%	1.23(12)%	2.6(2)%	5.4%
r_{1}/r_{0}		0 - 1%	0.3(1.1)%	1.8(5)%	2.9(8)%	≈4%
$r_0/\sqrt{t_0}$		0 - 1%	0.1(7)%	0.7(6)%	1.7(6)%	3%

Range at $M_{\rm c}$ from two estimates:

- scaling with $1/M^2$ (behavior for large M)

- scaling with 1/M (observed between $M = 0.4 \,\text{GeV}$ and 0.8 GeV)

 $(M_{\rm c} \text{ taken from [Rolf and Sint, hep-ph/0209255]})$



Conclusions and outlook

Relevance for decoupling of charm in QCD

Our numbers provide a rough estimate for charm effects in low energy observables in 2 + 1 + 1 simulations.

Put differently: tiny effects are being missed in 2+1 simulations (at low energies).

Low energy: up to r_1^{-1} was investigated.

No qualitative difference between decoupling $2 \rightarrow 0$ and decoupling $2+1+1 \rightarrow 2+1$ is expected

Pauli term $\psi_{\text{light}} \sigma_{\mu\nu} F_{\mu\nu} \psi_{\text{light}}$ does not appear in PT (chiral symmetry) and is therefore nonperturbatively suppressed: $M^{-1-\gamma}, \gamma > 0$

