Suppression of excited-state effects in lattice determination of nucleon electromagnetic form factors

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New York, NY, USA June 22–28, 2014

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 - why do *ep* scattering and muonic hydrogen results disagree so badly?

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- Here: explore how to suppress excited-state effects

Form Factors

• eN scattering cross section parameterized in terms of Sachs form factors G_E , G_M via Rosenbluth formula

$$\left(\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega}\right) \propto \left[\frac{\mathcal{G}_{\mathcal{E}}^2 + \tau \, \mathcal{G}_{\mathcal{M}}^2}{1 + \tau} + 2\tau \, \mathcal{G}_{\mathcal{M}}^2 \tan^2\left(\frac{\theta}{2}\right)\right]\,, \quad \tau = \frac{Q^2}{4m_N^2}$$

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 Matrix element of vector current between nucleon states decomposed in terms of Dirac and Pauli form factors F₁, F₂ as

$$\langle N(p',s')|V_{\mu}|N(p,s)\rangle = \overline{u}(p',s')\left[\gamma_{\mu}F_{1}+i\frac{\sigma_{\mu\nu}q_{\nu}}{2m_{N}}F_{2}\right]u(p,s)$$

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• Relationship given by

Lattice Setup

• Form ratios

$$R_{\gamma_{\mu}}(\mathbf{q},t,t_{s}) = \frac{C_{3,\gamma_{\mu}}(\mathbf{q},t,t_{s})}{C_{2}(\mathbf{q},t_{s})} \sqrt{\frac{C_{2}(\mathbf{q},t_{s}-t)C_{2}(\mathbf{0},t)C_{2}(\mathbf{0},t_{s})}{C_{2}(\mathbf{0},t_{s}-t)C_{2}(\mathbf{q},t)C_{2}(\mathbf{q},t_{s})}}$$

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• Extract Sachs form factors as

$$\operatorname{Re}\left[R_{\gamma_0}\right] = \sqrt{\frac{m_N + E_{\mathbf{p}}}{2E_{\mathbf{p}}}} G_E(Q^2)$$

and

$$\operatorname{Re}\left[R_{\gamma_{i}}\right]_{i=1,2} = \epsilon_{ij}p_{j}\frac{1}{\sqrt{2E_{\mathbf{p}}(m_{N}+E_{\mathbf{p}})}}G_{M}(Q^{2})$$

$$G_X^{\text{eff}}(Q^2, t, t_s) = G_X(Q^2) + c_{X,1}(Q^2) \mathrm{e}^{-m_\pi t} + c_{X,2}(Q^2) \mathrm{e}^{-2m_\pi (t_s - t)} + \dots$$

- Plateau method: Identify plateaux in t
- Problem: Need large t_s, where signal-to-noise ratio is poor
- Observe systematic trend in t_s even for $t_s \sim 1.4$ fm

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Summation method:

$$S_X(Q^2, t_s) = \sum_{t=0}^{t_s} G_X^{ ext{eff}}(\mathbf{q}, t, t_s) o c + t_s G_X(Q^2) + \mathcal{O}\left(\mathrm{e}^{-m_\pi t_s}
ight)$$

- Advantage: Parametrically reduced excited state contamination $m_{\pi}t_s$ instead of $m_{\pi}t$
- Disadvantage: Increased statistical errors

$$G_X^{\text{eff}}(Q^2, t, t_s) = G_X(Q^2) + c_{X,1}(Q^2) \mathrm{e}^{-m_\pi t} + c_{X,2}(Q^2) \mathrm{e}^{-2m_\pi (t_s - t)} + \dots$$

- **Excited-state fits:** Explicitly fit $G_X(Q^2, t, t_s)$ to leading excited-state contributions
 - as a function of t, $t_s t$ at each t_s separately, or
 - as a function of t_s , t at all t_s simultaneously
- Advantage: Fully removes leading excited state contamination
- Disadvantage: Somewhat model-dependent, hard to assess trustworthiness of results

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Chiral behaviour



• Parameterize each Sachs form factor as a dipole

$$G_X(Q^2)=rac{G_X(0)}{\left(1+rac{Q^2}{M_X^2}
ight)^2}$$

where $G_E(0) = 1$, $G_M(0) = \mu$

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• Charge radii and anomalous magnetic moment given by

$$\frac{1}{M_E^2} = \frac{r_E^2}{12} = \frac{r_1^2}{12} + \frac{\kappa}{8m_N^2} \qquad \frac{1}{M_M^2} = \frac{r_M^2}{12} = \frac{r_1^2 + \kappa r_2^2}{12(1+\kappa)}$$

where $\kappa = \mu - 1$

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• Given $r_E \approx r_M$, can extract magnetic moment also using

$$\mu = \lim_{Q^2 \to 0} \frac{G_M(Q^2)}{G_E(Q^2)}$$

with a flat extrapolation









Chiral Extrapolation



G.M. von Hippel

Excited State Effects on Nucleon Form Factors













- ullet Systematic trend in ${\it G_E}$ plateau values persists to $t_s\sim 1.4$ fm
- Even with summation method, G_E systematically too high
- Considering only the largest values of t_s brings summation method closer to experiment at the expense of large statistical errors
- Excited-state fits indicate a possible reason:
 - with small gap m_{π} , approach to plateau is very slow
 - summed ratios still receive sizeable corrections

Thank you for your attention



- BACKUP -















