Update on staggered Wilson fermions

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Introduction

Will discuss the current situation for staggered versions of Wilson fermions ("staggered Wilson fermions") regarding

- Theoretical viability
- Computational efficiency vs usual Wilson fermions
- Usability

As illustration of usability, will discuss construction of pseudoscalar meson operators for 2-flavor staggered Wilson fermions.

What are staggered Wilson fermions?

Usual Wilson fermion:

Naïve fermion (16 species) + Wilson term

 \rightarrow 1 physical species, 15 doublers

Staggered Wilson fermion (the idea):

Staggered fermion (4 species) + "staggered Wilson term"

 \rightarrow 1 or 2 physical species, 3 or 2 doublers

Origin of staggered Wilson fermions

- 2-flavor staggered version of overlap fermions arose from spectral flow approach to the index of the staggered Dirac operator in [DA, PRL (2010), PLB (2011)]. Has 2-flavor staggered Wilson fermion as kernel.
- 1-flavor version proposed later by C. Hoelbling [PLB (2011)]
- The staggered Wilson terms in both cases are combinations of the "flavored mass terms" of Golterman & Smit [NPB (1984)] which lift the degeneracy of the 4 staggered fermion flavors.

Motivation: Staggered more efficient than Wilson

Consider the momentum space Brillioun zone in 2D



Wilson

Staggered



+ spinor index α = 1,2

- 1 physical species
- 3 doublers (junk)

- 2 physical species
- no junk, no waste
- smaller Dirac matrix!

Staggered formalism is more efficient, but what to do about the extra species?

Possibilities:

- Live with the 4 species ('tastes') for each quark.
 -- The current approach for Asqtad and HISQ.
- Use staggered fermion as a description of 4 quarks (flavors)
 -- Investigated by Golterman & Smit (1984).
 - -- Problems arise when the mass degeneracy of the 4 species is lifted.
- Add a `Wilson term' to reduce the number of physical species [DA, C. Hoelbling].

Staggered Wilson fermion formulation

Fermion action $\overline{\chi}(D_{sW} + m)\chi$ where

 $D_{sW} = D_{st} + \frac{1}{a}M$ staggered Wilson Dirac operator

 $D_{st} = \eta_{\mu} \nabla_{\mu}$ usual staggered Dirac operator

M = flavored mass terms [Golterman-Smit]

In the spin-flavor interpretation of the free field momentum space rep have

$$\hat{M}(q) = \mathbf{1} \otimes \mathbf{M} + O(a^2 q^2)$$

Zero-eigenvector of $\mathbf{M} \rightarrow$ massless physical fermion species Nonzero-eigenvector of $\mathbf{M} \rightarrow$ heavy species, mass ~ 1/a (doublers)

Some details of the flavored mass terms

$$M = M^{(0)} + M^{(2)} + M^{(4)}$$
 $M^{(n)} = n$ -link term

In the free field momentum space spin-flavor interpretation have

$$\hat{M}(q) = r_0(\mathbf{1} \otimes \mathbf{1}) + ir_{\mu\nu}c_{\mu}(aq)c_{\nu}(aq)(\mathbf{1} \otimes \xi_{\mu}\xi_{\nu}) + r_5c_5(aq)(\mathbf{1} \otimes \xi_5)$$
$$c_{\mu}(aq) \equiv \cos(aq_{\mu}) , \quad c_5(aq) \equiv \Pi_{\mu=1,\dots4}\cos(aq_{\mu})$$
$$\xi_{\mu}: \text{ Dirac matrices on flavor space}$$

$$\Rightarrow \hat{M}(q) = \mathbf{1} \otimes (r_0 \mathbf{1} + i r_{\mu\nu} \xi_{\mu} \xi_{\nu} + r_5 \xi_5) + O(a^2 q^2)$$

Lifts the degeneracy of the 4 staggered fermion flavors.

Can get zero-eigenvector (\rightarrow massless fermion) by tuning r_0 .

2-flavor version of staggered Wilson fermion

$$M = M^{(0)} + M^{(4)}$$
, $M^{(2)}$ is absent

$$\hat{M}(q) = \mathbf{1} \otimes (r_0 \mathbf{1} + r_5 \xi_5) + O(a^2 q^2)$$

→ Set $r_0 = -r_5 \equiv r$, then

$$\hat{M}(q) = r\mathbf{1} \otimes (\mathbf{1} - \xi_5) + O(a^2 q^2)$$

The additional content of the second second

Chirality aspect of staggered Wilson fermions

Recall: staggered fermion has an exact *flavored* chiral symmetry

$$\varepsilon D_{st} = -D_{st}\varepsilon$$

$$\varepsilon \chi(x) = (-1)^{x_1 + \dots + x_4} \chi(x)$$
 $\varepsilon \cong \gamma_5 \otimes \xi_5$

Crucial fact:

 ε acts as *unflavored* $\gamma_5 \otimes \mathbf{1}$ on the physical species of the staggered Wilson fermion.

This is because $\xi_5 = \mathbf{1}$ on the zero-eigenspace of **M**.

Chirality aspect (continued)

 \rightarrow Use ε as the unflavored γ_5 for staggered Wilson fermions. Note:

$$D_{sW} = \varepsilon D_{sW} \varepsilon$$
 (ε -hermiticity)
 $\varepsilon^2 = 1$

These don't hold for the usual unflavored γ_5 in the staggered formalism!

Get staggered versions of domain wall and overlap fermions by $D_{\rm W} \rightarrow D_{\rm sW}$, $\gamma_5 \rightarrow \epsilon$ [DA, PLB (2011)]

Theoretical viability of staggered Wilson fermions?

- $M^{(2)}$ and $M^{(4)}$ break the staggered `shift' symmetries
- *M*⁽²⁾ also breaks the flavored lattice rotation symmetry of staggered fermions.
 - \rightarrow Questions:
 - Is the staggered Wilson spin-flavor structure preserved at the quantum level?
 - New counter-terms appear? Fine-tunings required?

Results on theoretical viability [DA, to appear]

- There are new fermionic counterterms, but their only effect on the physical species is a wavefunction renormalization.
 → No fine-tunings needed for these.
- For the <u>1-flavor versions</u> of staggered fermions, which break lattice rotation symmetry, a new gluonic counter-term arises from the fermion loop contributions to the gluonic 2-, 3- and 4-point functions.
 - \rightarrow Needs to be included in the bare action and fine-tuned.

<u>Conclusion</u>:

1-flavor versions of staggered Wilson fermions are problematic, but 2-flavor version is fine.

Will restrict considerations to 2-flavor staggered Wilson fermions henceforth.

Computational efficiency

- Previous study on quenched 16³ × 32 lattice with β = 6 found a speed-up factor of 4-6 for staggered Wilson compared to usual Wilson for inverting the Dirac operator on a source at fixed pion masses. [DA, D. Nogradi & C. Zielinski, Lattice 2013 proc.]
- New study on $20^3 \times 40$ lattice with $\beta = 6$ and $\beta = 6.14$ tests the dependence of the speed-up factor on physical volume and lattice spacing.

Results:

- Increased speed-up for smaller lattice spacing at fixed volume
- Speed-up mostly unchanged (slight increase) for larger volume at fixed lattice spacing.

Details in C. Zielinski's talk.

Usability of 2-flavor staggered Wilson fermions

- Need meson & baryon operators for the 2 flavors
- Situation is complicated by the fact that only a subset of the flavor symmetries are unbroken.
- However, meson & baryon operators for the 4 degenrate flavors described by usual staggered fermion are already known.
 [Golterman & Smit, Sharpe, ...]
- Can adapt these to the staggered Wilson case

 will illustrate for pseudoscalar mesons.

Pseudoscalar mesons with staggered Wilson

- For usual staggered fermions (4 flavors) have 16 pseudoscalar mesons
- Their operators $\overline{\chi}(\gamma_5 \otimes \xi_F) \chi$ form irreps of the flavored lattice rotation symmetry group as follows:

$$\xi_F = \mathbf{1}$$
 (scalar) $\xi_F = \xi_\mu$ (vector) $\xi_F = \xi_\mu \xi_\nu$ (tensor)

$$\xi_F = \xi_\mu \xi_5$$
 (pseudovector) $\xi_F = \xi_5$ (pseudoscalar)

Use same operators for staggered Wilson case. They continue to form irreps of the flavored rotation symmetry group, but their physical meaning changes...

Pseudoscalar mesons with staggered Wilson (cont'd)

$$\xi_5 = \begin{pmatrix} \mathbf{1} & \mathbf{0} \\ \mathbf{0} & -\mathbf{1} \end{pmatrix} \qquad \qquad \xi_F = \begin{pmatrix} \text{light-light} & \text{light-heavy} \\ \text{heavy-light} & \text{heavy-heavy} \end{pmatrix}$$

The light-light part of ξ_F determines the physical meaning of $\chi(\gamma_5 \otimes \xi_F)\chi$ for staggered Wilson fermions.

→ The tensor irrep $\xi_F = \xi_\mu \xi_\nu$ gives the 3 pions. They are degenerate since these pion operators belong to the same irrep.

E.g.

$$\xi_1 \xi_2 = \begin{pmatrix} \mathbf{0} & \mathbf{\sigma}_1 \\ \mathbf{\sigma}_1 & \mathbf{0} \end{pmatrix} \begin{pmatrix} \mathbf{0} & \mathbf{\sigma}_2 \\ \mathbf{\sigma}_2 & \mathbf{0} \end{pmatrix} = \begin{pmatrix} i\mathbf{\sigma}_3 & \mathbf{0} \\ \mathbf{0} & i\mathbf{\sigma}_3 \end{pmatrix}$$

 $\rightarrow -i\overline{\chi}(\gamma_5 \otimes \xi_1 \xi_2)\chi$ becomes operator for π_0 .

Conclusions

The situation for 2-flavor staggered Wilson fermions:

- Theoretically viable
- Computationally efficient
- Usable

Potentially advantageous for

- QCD thermodynamics
- High precision computation of the eta' mass.

C. Zielinski's talk:

Details for the computational efficiency, including spectrum of the staggered Wilson Dirac operator for a range of lattices.