U(1)_A symmetry at finite temperature with DWF

Guido Cossu

High Energy Accelerator Organization (KEK)

Joint Institute for Computational Fundamental Science (JICFuS)

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Outline symmetry Finite temperature Methods & Final Results Work in progress thoughts

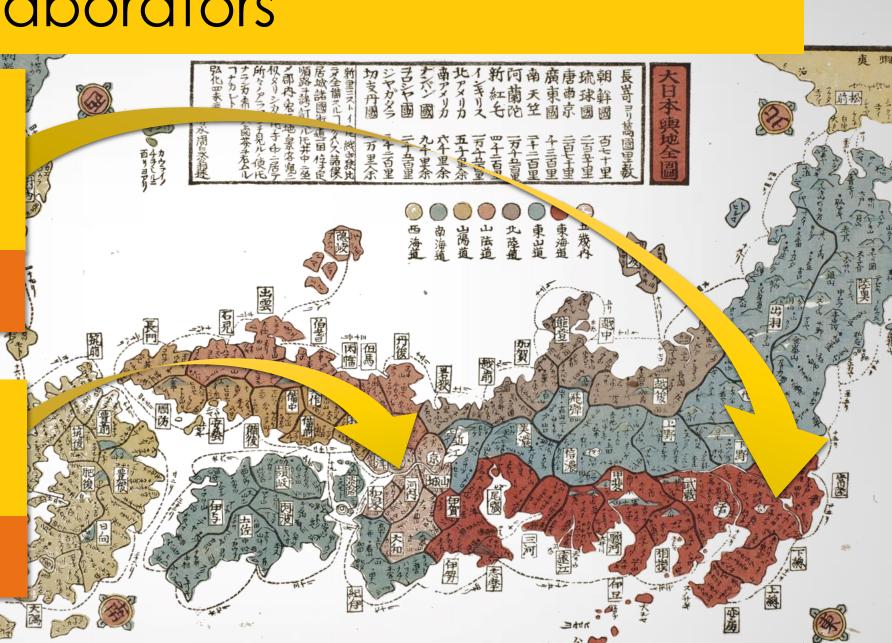
JLQCD collaborators

Shoji Hashimoto Jun-ichi Noaki Takashi Kaneko

KEK

Hidenori Fukaya Akio Tomiya

Osaka Un.





Low temperature – symmetries

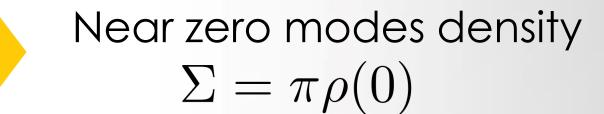
$$SU(2)_L \times SU(2)_R \times U(1)_V \times U(1)_A$$

Chiral condensate

Instantons

T<T_C

Dirac operator eigenmodes





$$\int \partial_{\mu} J_{\mu_5} \propto Q$$

High temperature – symmetries

$$SU(2)_L \times SU(2)_R \times U(1)_V \times U(1)_A$$

Current knowledge

Restoration of chiral symmetry at T_c

No condensate

Restoration at $T \to \infty$



Axial symmetry

High temperature – symmetries

$$SU(2)_L \times SU(2)_R \times U(1)_V \times U(1)_A$$

 $T \gtrsim T_c$?

No condensate

Axial symmetry?

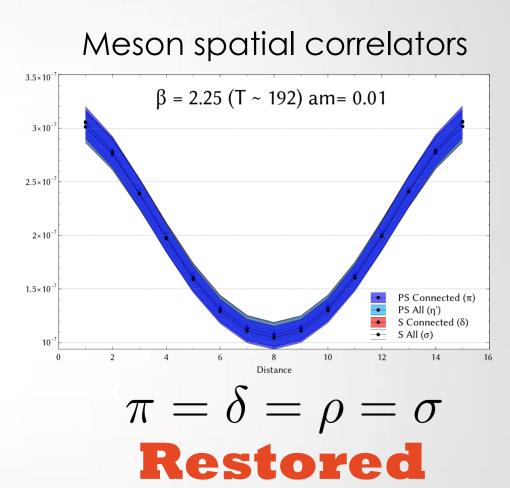


Recent literature - I

G. Cossu et al. (2013) for JLQCD Disconnected meson diagrams vanish at temperatures above T_c

Related: Gap in the Dirac spectrum

Aoki, Fukaya, Taniguchi (2012) Analytic calculation (Overlap) Dirac spectrum $\rho(\lambda) \sim c\lambda^3$ Implies **U(1)**_A anomaly invisible



Recent literature - II

Bazavov et al. (2012-13)
Domain wall, several volumes
Dirac spectrum, susceptibilities
NOT restored

Ohno et al., Sharma et al. (2012-13)
Overlap on HISQ configurations
Dirac spectrum
NOT restored

Brandt et al. (2013)
Wilson improved fermions
Screening masses
NOT restored

Our previous study

Exact chiral symmetry (Overlap)

topology fixed
Only 163x8 volume
Mass dependence
No continuum limit



Generalized Domain Wall

$$D^{4}(m) = \frac{1+m}{2} + \frac{1-m}{2}\gamma_{5}\operatorname{sgn}(H)$$

Play with the sign function

Möbius Kernel

$$H_M = \gamma_5 \frac{bD_W}{2 + cD_W}$$

Function approximation
Transfer matrix in 5D

- Hyperbolic tangent
- Rational approximation

Reduced residual mass

$$b=2 c=1 Scaled Shamir, m_{res} \sim 10^{-4}$$

Status of simulations



Symanzik + smeared DWF

Multipurpose code, HMC & measurements Available on request, soon online

Optimized for BlueGene/Q

Webpage: http://suchix.kek.jp/guido_cossu/

Collected data

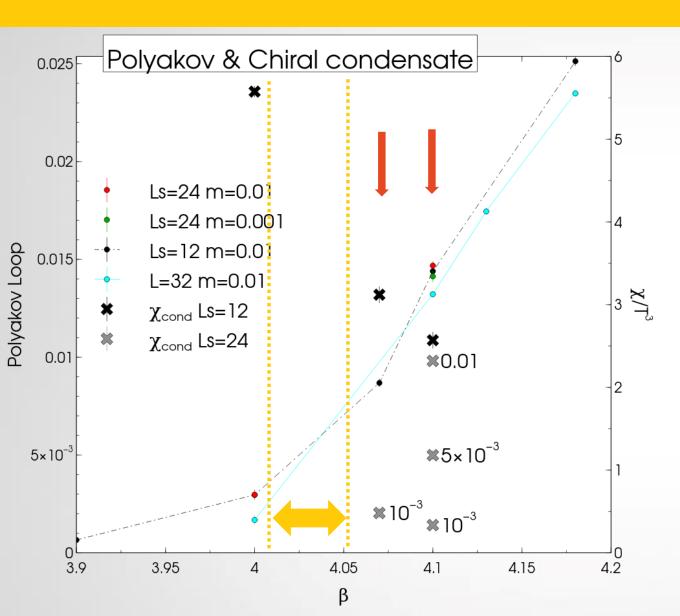
- 2 volumes
- 3 masses
- 5 temperatures
- Topology changes
- $-N_{t}=8, N_{t}=12$

Full analysis in progress





Phase transition



Today:

T ~ 184, 200 MeV (red arrows)

Phase transition at ~180 MeV

2 volumes Mass dependence

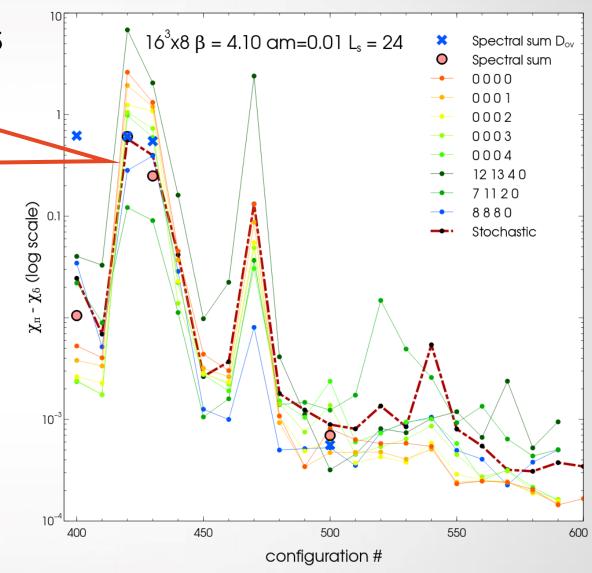
N_t=12 analysis not ready

Delta $\Delta = \chi_{\pi} - \chi_{\delta}$ $\chi_{X} = \int \langle X(0)X(x) \rangle$

Using local source is dangerous

Stochastic measurement is in nice accordance with the spectral sum

$$\Delta = \int \frac{2m^2 \rho(\lambda, m)}{(\lambda^2 + m^2)^2}$$



Source of the signal

Discrete spectral sum

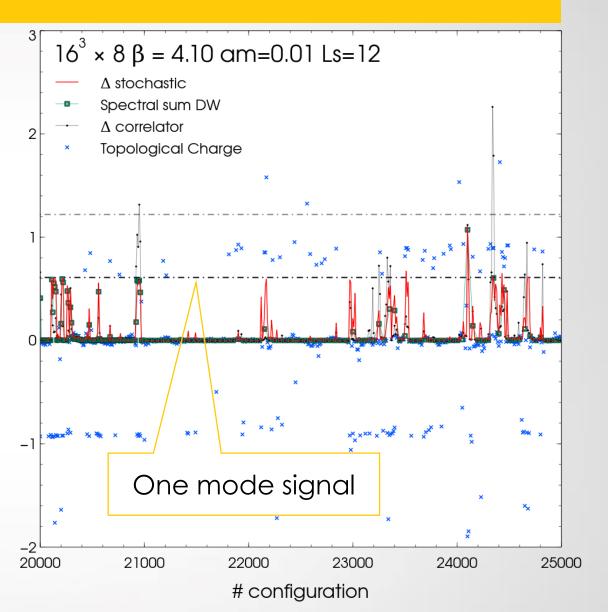
Zero modes

$$\Delta = \frac{2N_0}{Vm^2} + \sum_{\lambda \neq 0} \frac{2m^2}{V(\lambda^2 + m^2)^2}$$

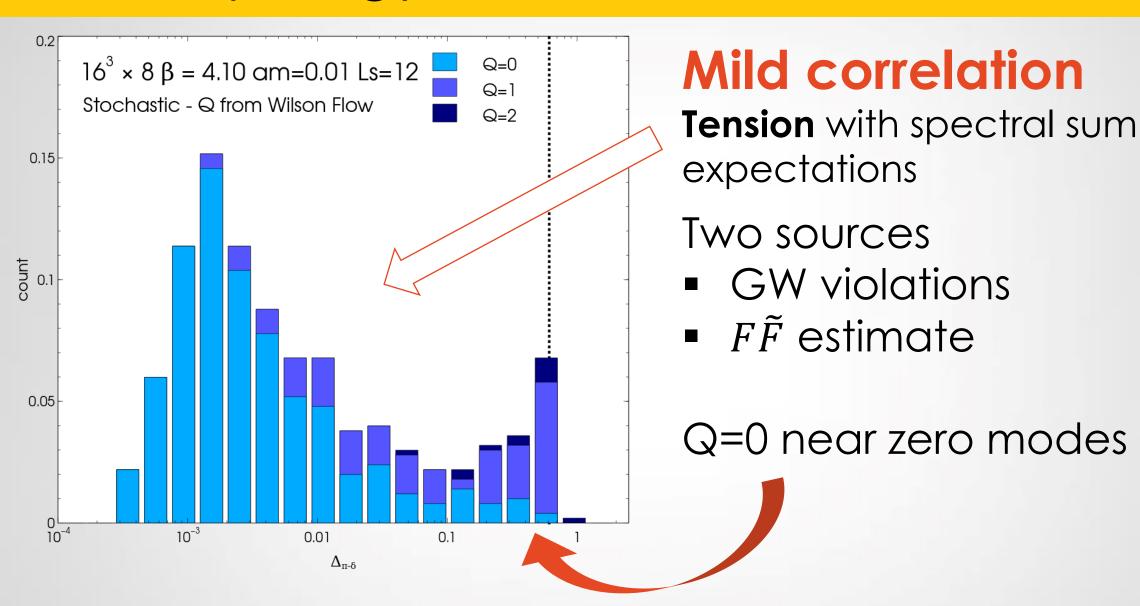
Peaks dominate the signal

76%

Fluctuations of **3 orders of magnitude**



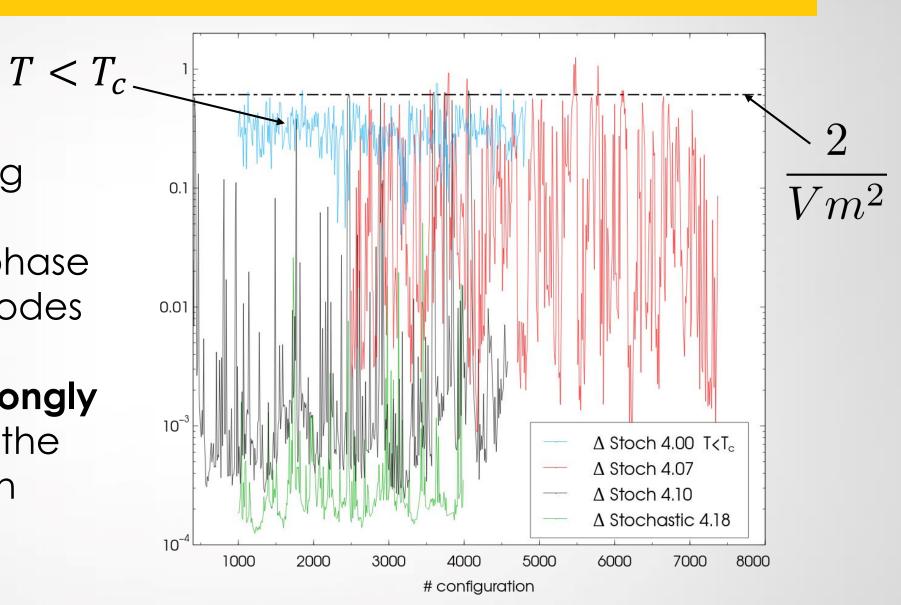
Δ - Topology correlation



Temperature dependence

Broad picture arising at this stage:

- Just above the phase transition zero modes dominate
- Then they are strongly suppressed and the signal goes down



Let's increase volume – m=0.01

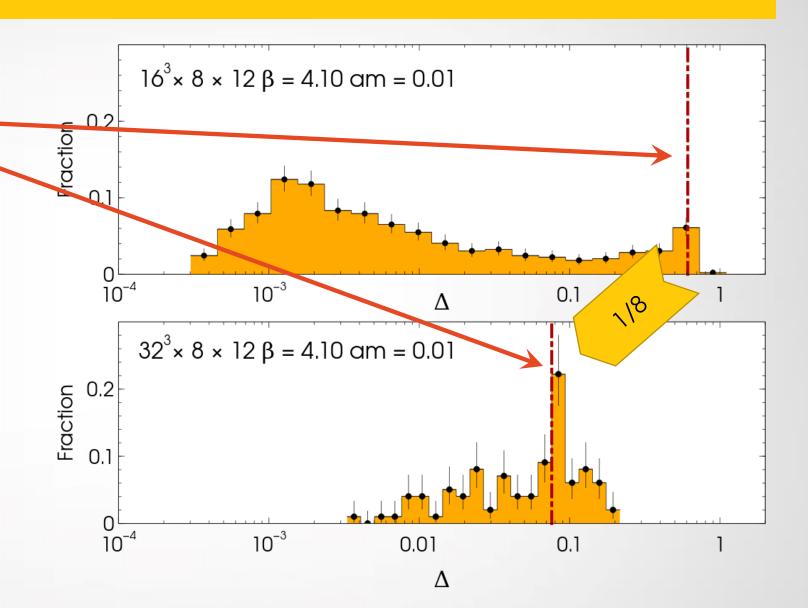
Zero mode contribution

suppressed ~1/V

As expected from spectral sum

Bulk contribution increases

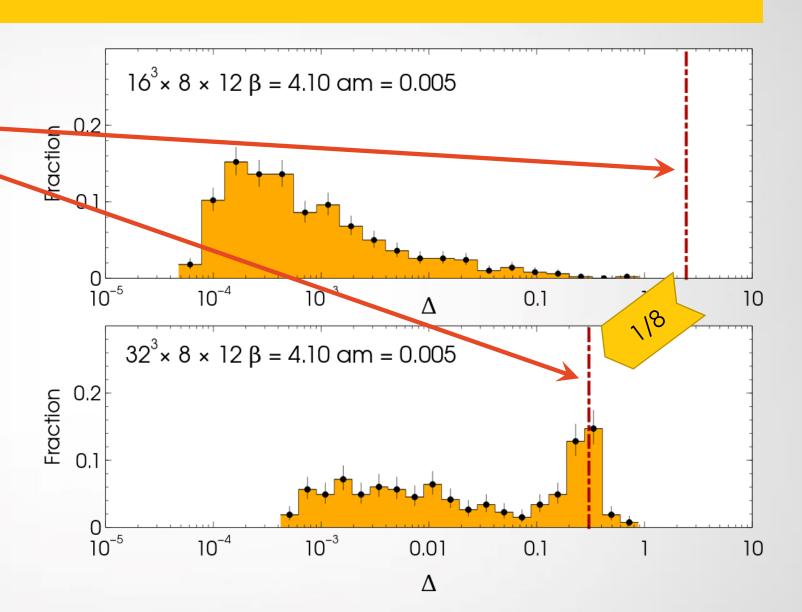
Decrease the mass?



Let's increase volume – m=0.005

Zero mode
contribution
suppressed ~1/V
As expected from
spectral sum

Bulk contribution increases



Volume&mass dependence

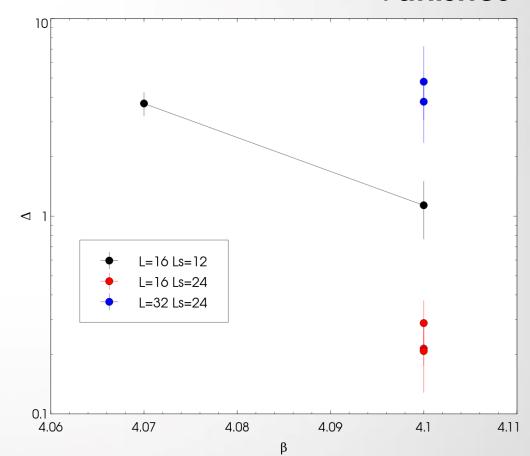
$$\chi_t = \lim_{V \to \infty} \frac{\langle Q^2 \rangle}{V} = \text{const.} \to \frac{N_0}{V} \to 0$$

Zero modes contribution vanishes

Conclusion: signal from the bulk part, near zero modes

Let's cut all configurations with Q>0

Signal constant with the mass



Is everything all right? – I

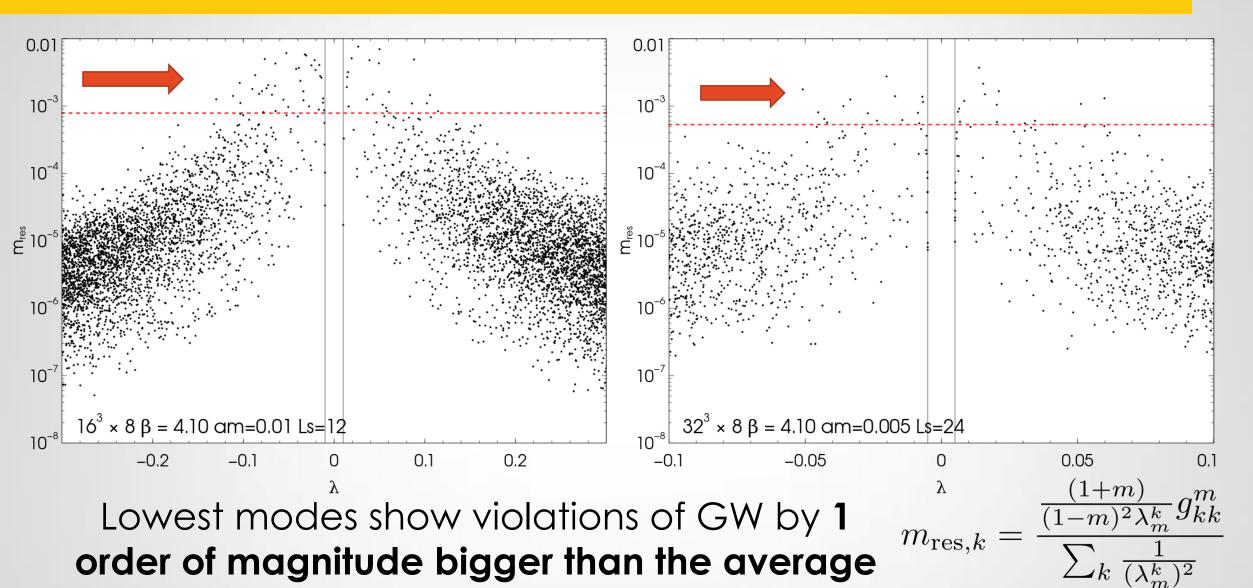
From the Ginsparg-Wilson relation we can measure the amount of violation for each mode, g_{kk}

$$\frac{1}{1-m^2} [(\gamma_5 - H_m)(H_m - m\gamma_5) + (H_m - m\gamma_5)(\gamma_5 - H_m)] = 0$$

$$g_{kk}^m = \psi_m^k {}^{\dagger} \gamma_5 \psi_m^k - \frac{(\lambda_m^k)^2 + m}{\lambda_m^k (1+m)}$$

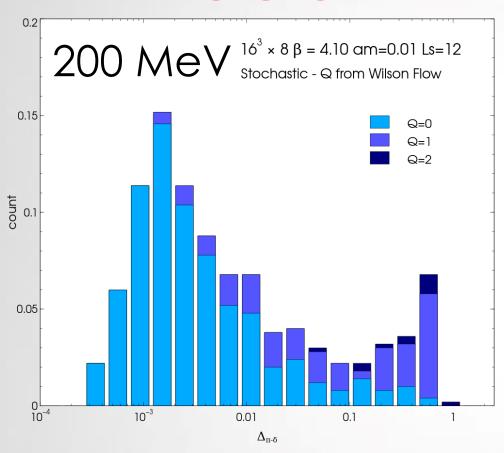
$$m_{\text{res}} = \frac{\sum_k \frac{(1+m)}{(1-m)^2 \lambda_m^k} g_{kk}^m}{\sum_k \frac{1}{(\lambda_m^k)^2}}$$

Is everything all right? – II

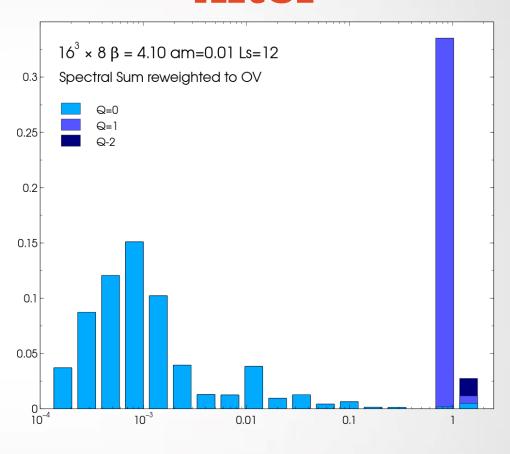


Reweight it! (DWF to Overlap)

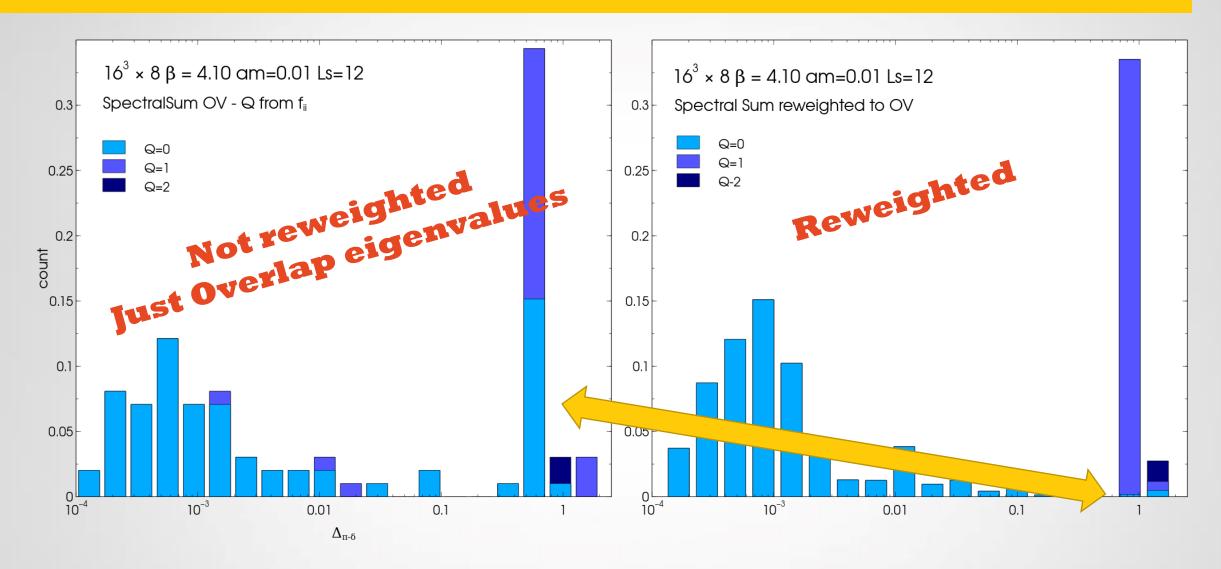
Before



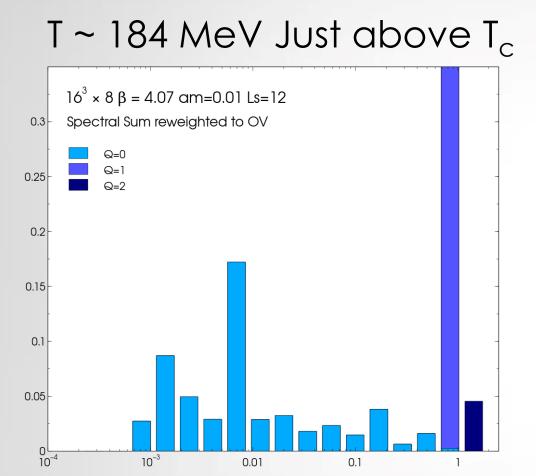
After

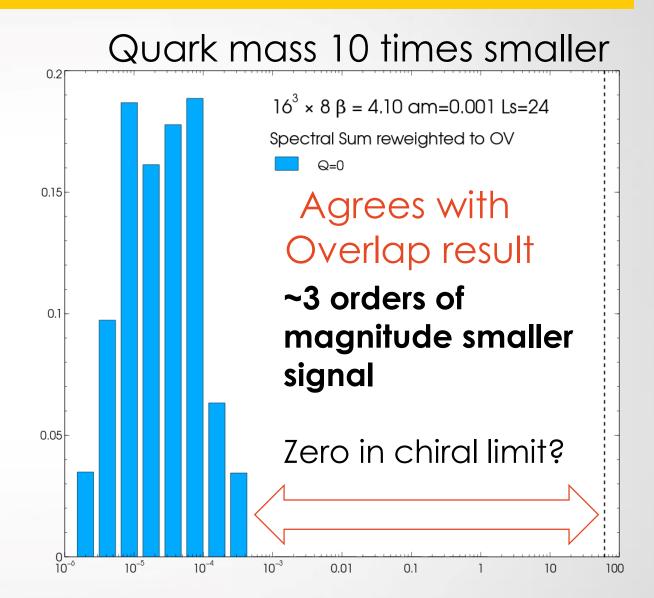


Reweighting alters the final answer!



Temperature and mass dependence







Instanton gas – hints?

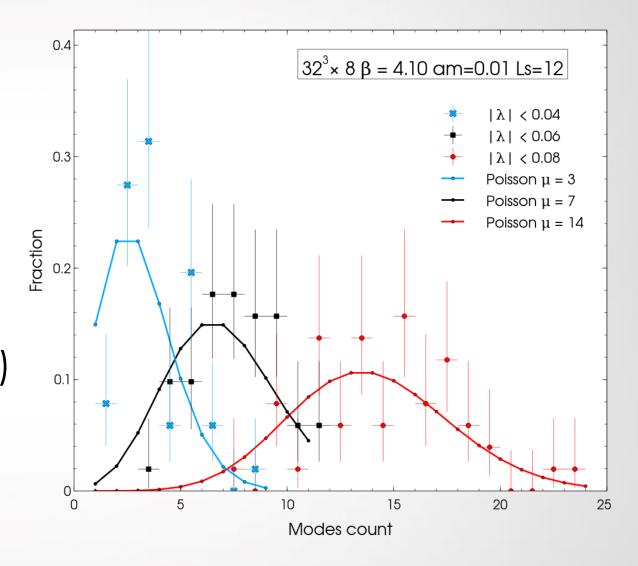
Results not yet conclusive (analysis running right now)

If the large volume signal is not coming from lattice artifacts

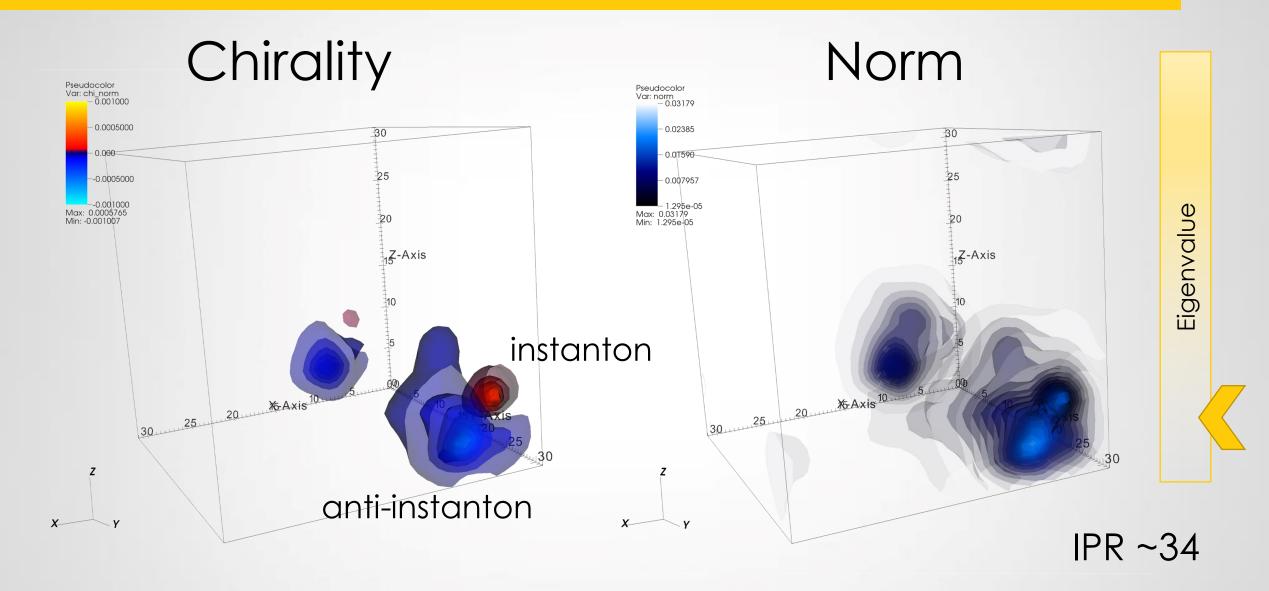
Near zero modes are responsible for breaking U(1)

What are they?

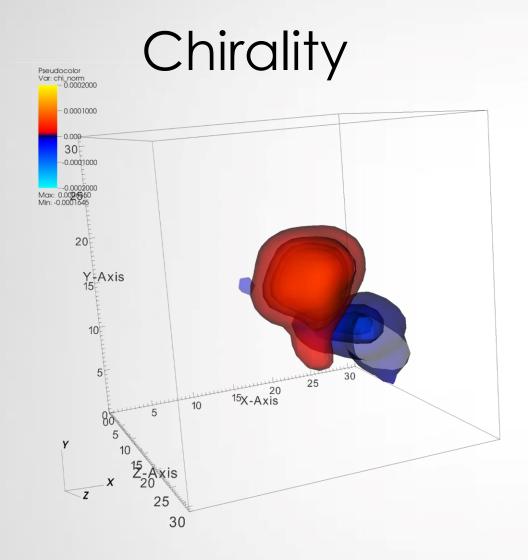
Poisson distributed?

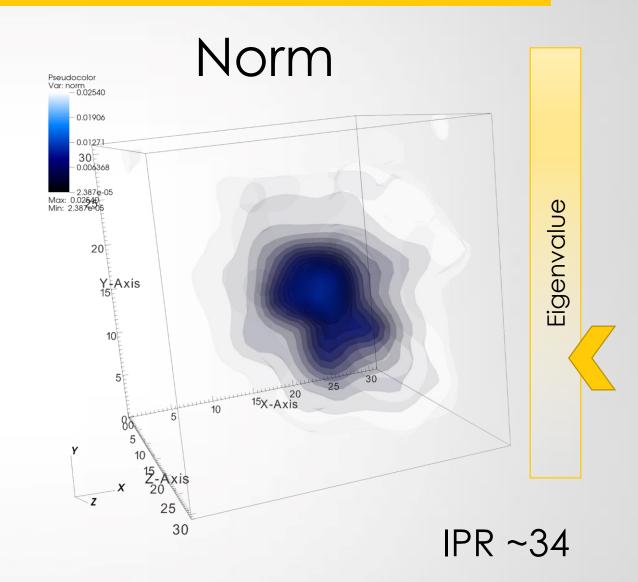


Fun with 3D – put your glasses on

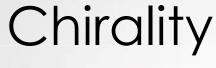


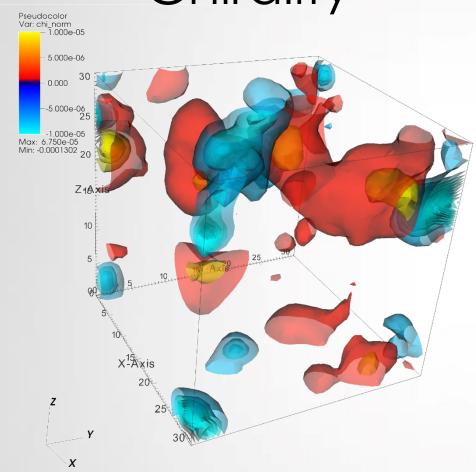
Fun with 3D – put your glasses on





Fun with 3D – put your glasses on





Norm Pseudocolor Var: norm - 0.01219 0.009168 0.006149 0.00313125 Eigenvalue - 0.0001132 Max: 0.01219 Min: 0.0001132 20 ZAAxis

IPR ~5

Summary – one more slide...

DWF volume & mass dependence suggests that near zero modes are the source of U(1) breaking

Lattice artifacts can spoil the signal

Exact chiral symmetry results differ from DWF

DWF lowest modes look like an instanton weakly interacting gas

Are we finished?

The talk is over the work is not!

Collected data yet to analyze

- Reweighting
- Continuum limit
- Chiral limit

Lattice artifacts?

Gas of instanton pairs, dyons?

Correlation with Polyakov loop?

U(1) restoration above critical

temperature

is still an open question.

Thanks!

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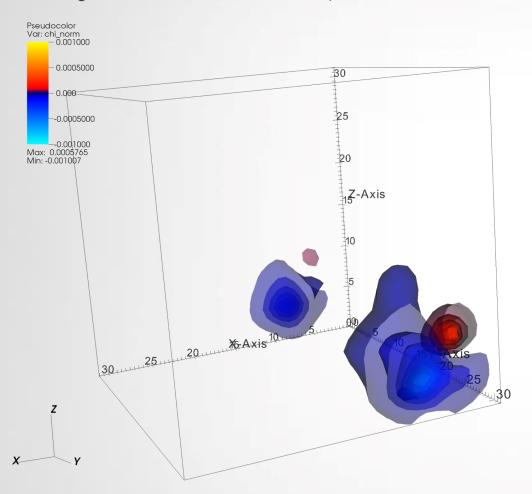




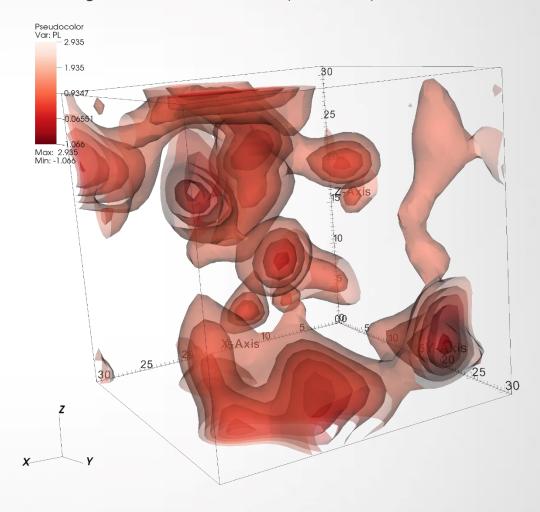


Lowest mode vs Polyakov Loop

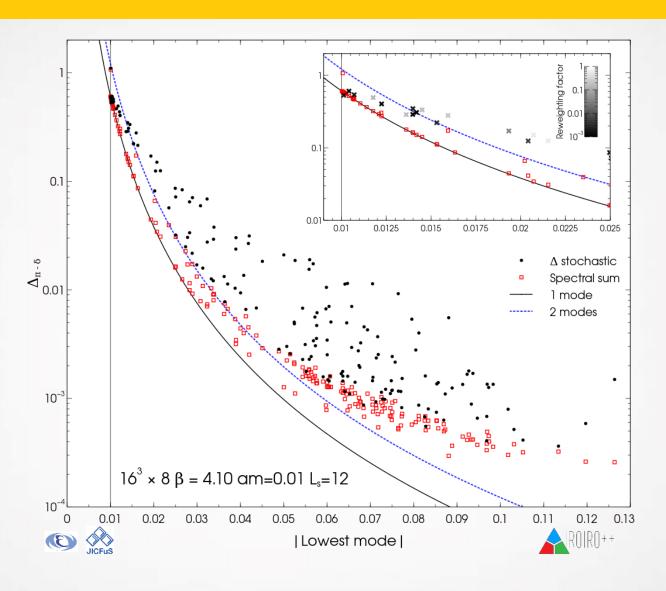
DB: eigenHDW_Conf2080_2_chirality_t0



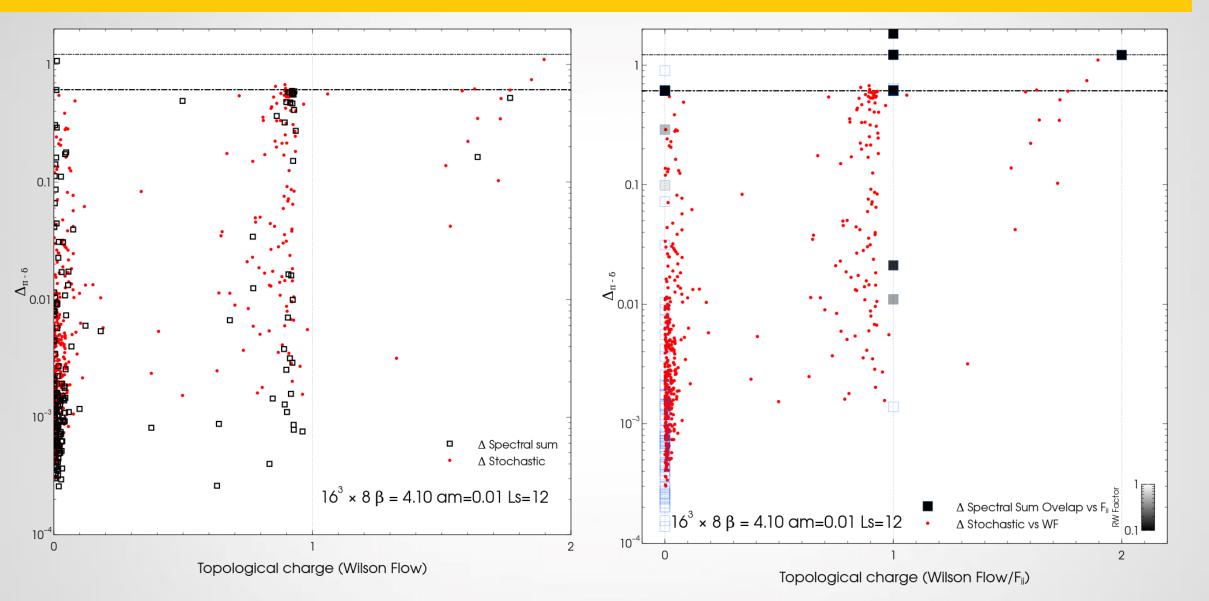
DB: eigenHDW_Conf2080_PolyakovLoop



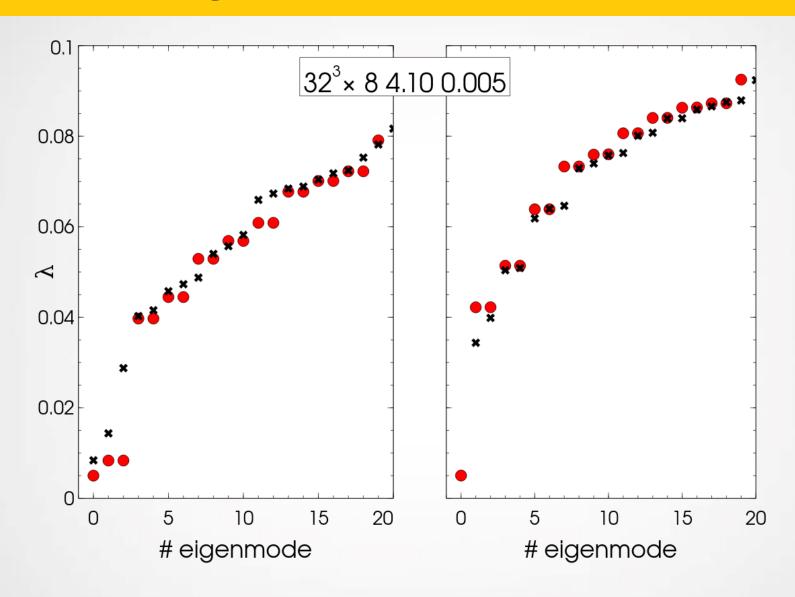
GW violations



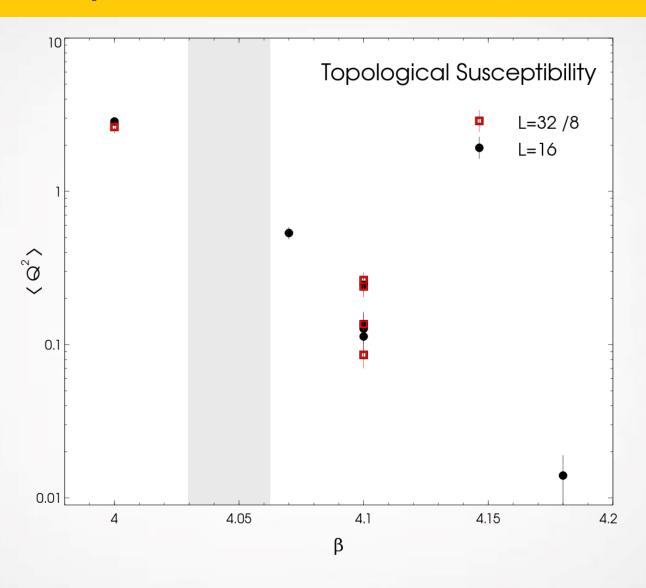
Many configurations violate GW



DW – OV eigenvalue mismatch



Susceptibility scales with volume



Let's increase volume

Zero mode contribution ~1/V -

- Bulk contribution increases

