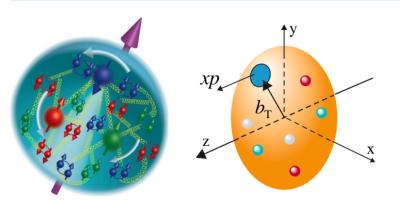
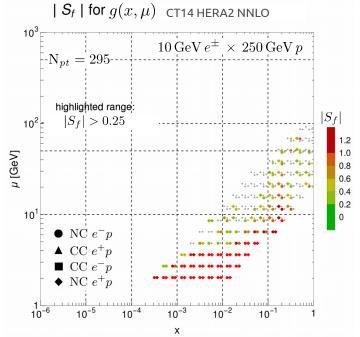
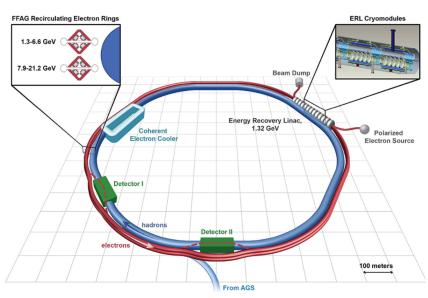
First-round studies of the EIC's PDF implications

Tim Hobbs, JLab EIC Center & CTEQ@SMU March 20th 2020















1st EIC Yellow Report Workshop; Temple Univ. 19-21 March 2020

Unraveling PDFs' flavor dependence is challenging; multiple EIC channels/processes needed

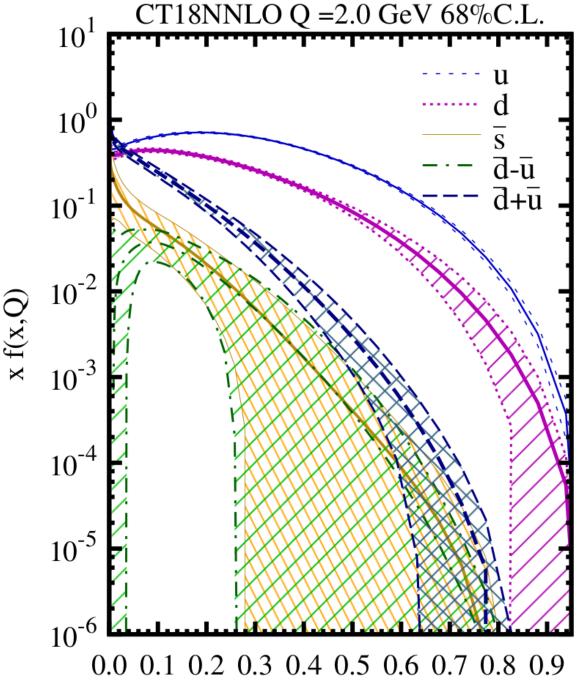


Figure: S. Dulat

note PDFs' different orders-of-mag.!

NC DIS: sensitivity to d-type quarks $\frac{1}{4}$ that of u-type

$$\sigma \propto \frac{4}{9}(u_+ + c_+) + \frac{1}{9}(d_+ + s_+ + b_+)$$

CC DIS: lower accuracy (1/10 lumi.)

high x (>0.1) inclusive DIS

- $\rightarrow u$ -quark dominates
- \rightarrow d-quark ½ of u, but harder to access in NC DIS (above)
- $\rightarrow \, \bar{d} + \bar{u} \, \, \sim \, \, {\rm few \; percent \; of } \, u$

...1% error on $u \rightarrow$ 50-100% error on $\bar{d} + \bar{u}$

- ightarrow for x ~0.1, $s pprox \bar{s} pprox \bar{d} \bar{u} < 0.1(\bar{d} + \bar{u})$
- \rightarrow at x>0.5, no separation for $\bar{u}, \bar{d}, \bar{s}$

we have a dedicated effort to explore PDF impacts of the EIC

Collaborators and consultants

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Southern Methodist

Sayipjamal Dulat

Xinjiang Univ.

C.-P. Yuan

Michigan State

Alberto Accardi Yulia Furletova Jefferson Lab

 the goal: use recentlydeveloped tools for PDF global analyses to examine the PDF pulls of EIC pseudodata

> 1803.02777 1806.07950 1904.00022 1907.00988 2001.07862



- needed for the Yellow Report Initiative: <u>quick</u>, <u>unambiguous</u> PDF impact metrics
 - → in turn, these can be incorporated into the YR workflow:

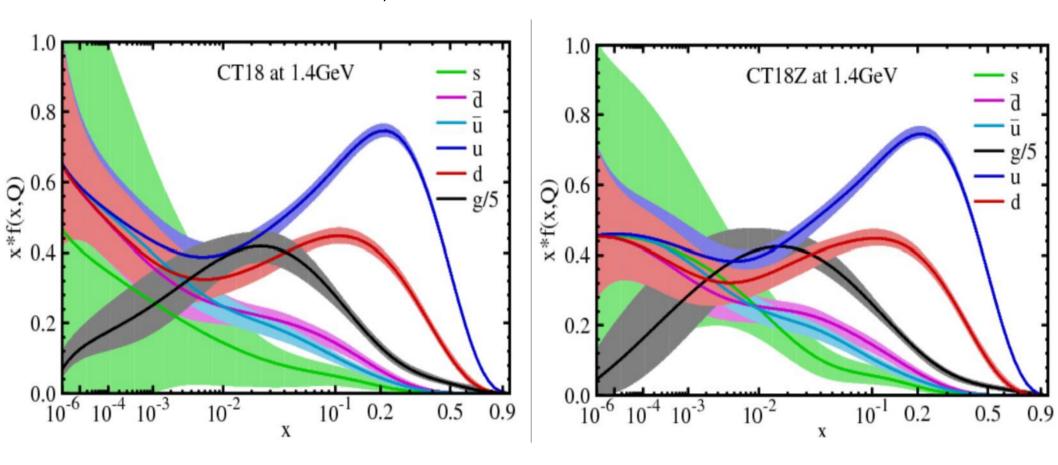
iteratively, machine design \rightarrow simulation \rightarrow physics

→ speed, simplicity, clarity can help ensure robust and timely convergence

CT18 parton distributions

PDF analyses are challenging! (theoretically, computationally, statistically, ...)

CT18 main analysis \rightarrow arxiv: 1912.10053.



 have developed fast methods to guide global fits toward highest-impact data & calculations

[here, Hessian profiling with (I) ePump and (II) PDFSense]

→ can be quickly applied to EIC pseudodata to supplement fitting

high-energy EIC pseudodata

- reach in center-of-mass energy, $20 \leq \sqrt{s} \leq 140\,\mathrm{GeV}$
 - → luminosities 2-3 decades greater than at HERA
 - → á la HERA, the combination of precision & kinematic coverage provide constraining 'lever arm' on QCD evolution
 - \rightarrow QCD evolution: (high x, low Q) \leftrightarrow (low x, high Q)
- as a generic scenario, we consider here the simulated impact of a machine with: $10\,\mathrm{GeV}\,e^\pm\,\mathrm{on}\,250\,\mathrm{GeV}\,p\quad(\sqrt{s}=100\,\mathrm{GeV})$

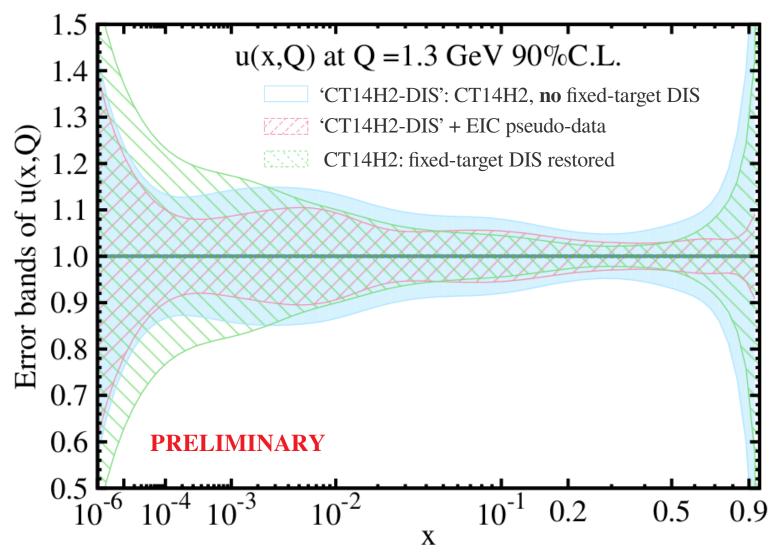
$$\mathcal{L} = 100 \,\text{fb}^{-1} \, e^{-} \,\text{pseudodata}$$

$$\mathcal{L} = 10 \,\text{fb}^{-1} \, e^{+} \,\text{pseudodata}$$
NC/C

→ generated based on CT14 HERA2 NNLO PDF fit

I) Hessian profiling [ePump] for EIC impacts on PDF errors

ePump: Schmidt, Pumplin, and Yuan; PRD98 (2018) no.9, 094005



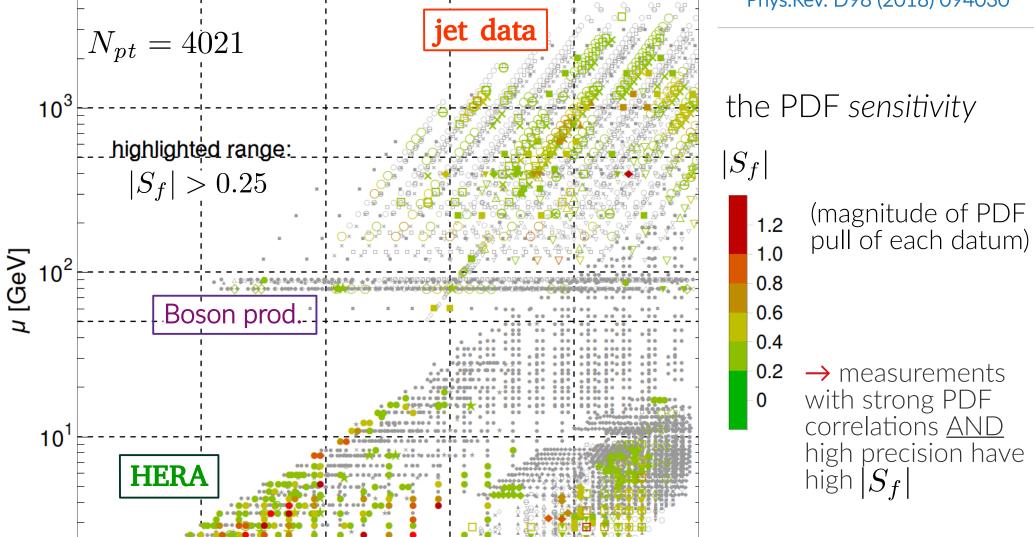
- EIC pseudodata supersede fixed-target DIS information in CT fits
- reweighting strongly depends on parametrizations; other ambiguities

→ complementary approaches welcome!

$|S_f|$ for $g(x,\mu)$, CT14 HERA2 NNLO

II) visualizing impacts with PDFSense

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 10^{-2}

Χ

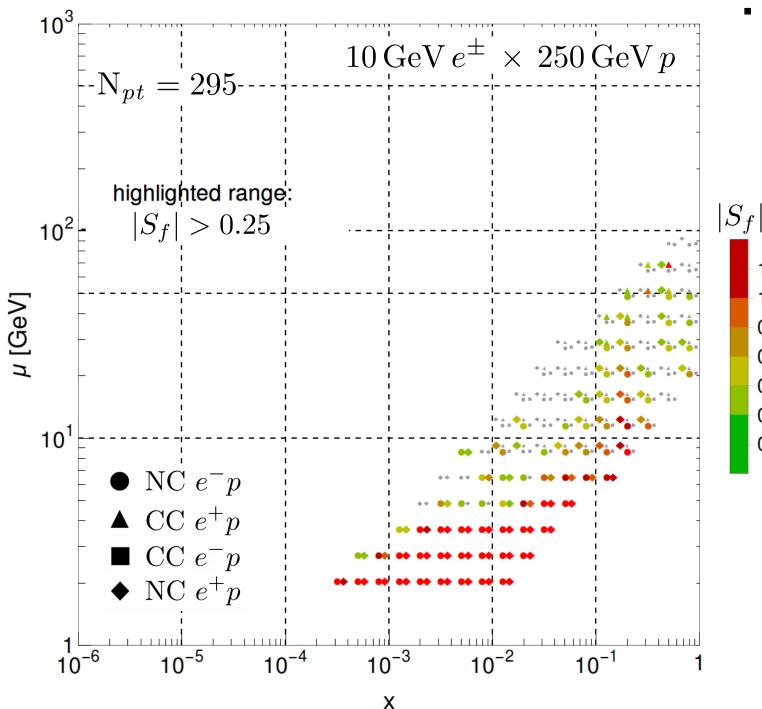
 10^{-4}

fixed-target

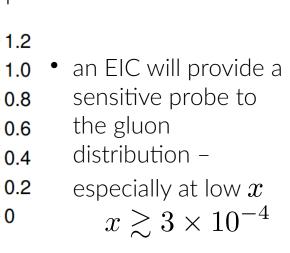
10⁻¹

 used to identify high-impact data for CT18

$\mid \mathcal{S}_{\it f} \mid$ for $g(x,\mu)$ CT14 HERA2 NNLO



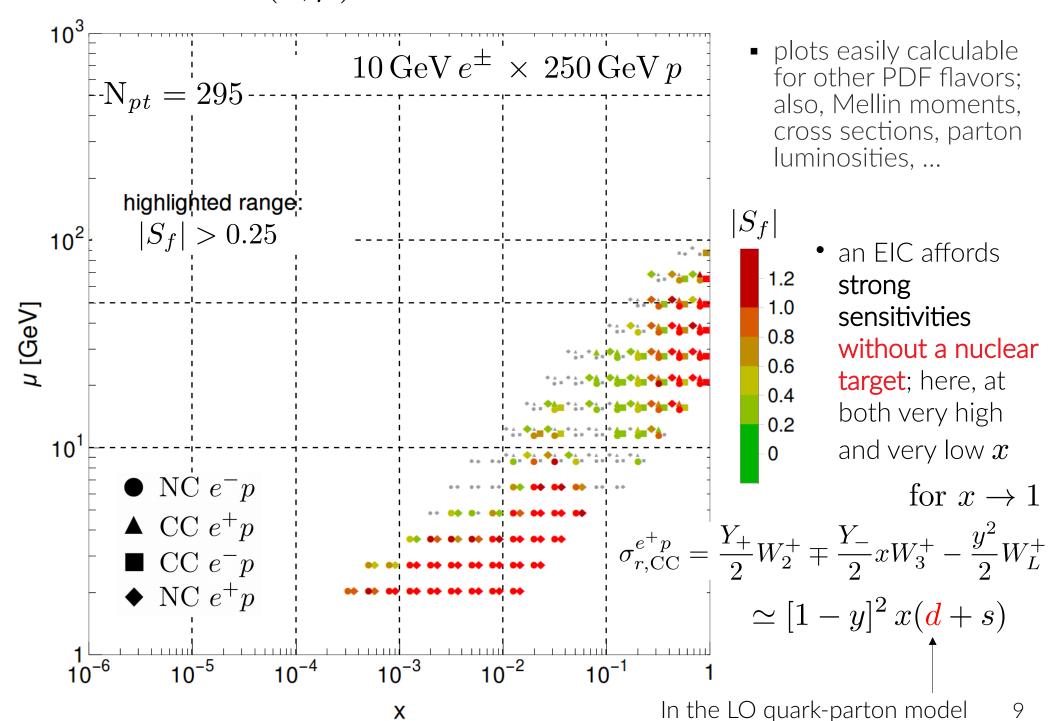
 PDFSense quickly visualizes the pulls of the EIC pseudodata, comparing on the same scale as for the candidate CT18 exps.



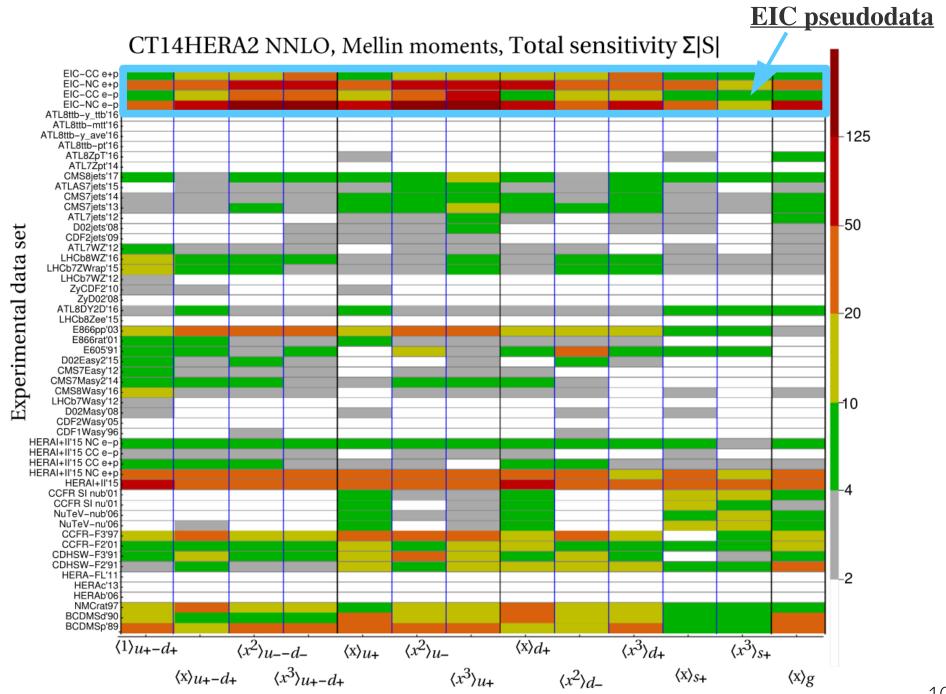
• these constraints arise from high statistics neutral current data on $\sigma_{n,NC}^{e^{\pm}p}$

8

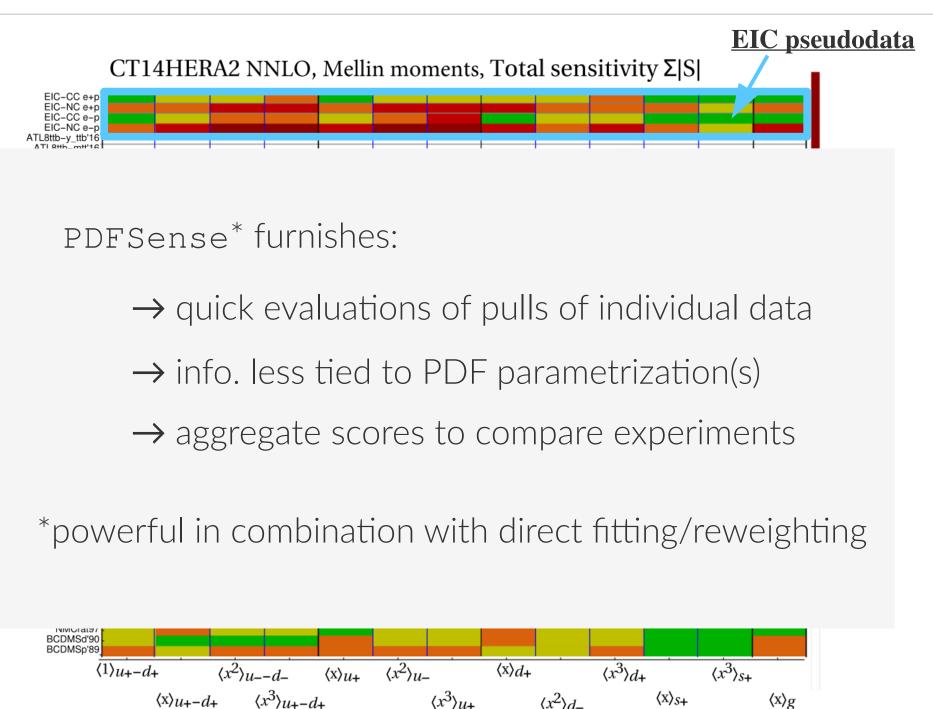
$|S_f|$ for $d(x,\mu)$ CT14 HERA2 NNLO



sensitivities can be aggregated for direct comparisons of exps



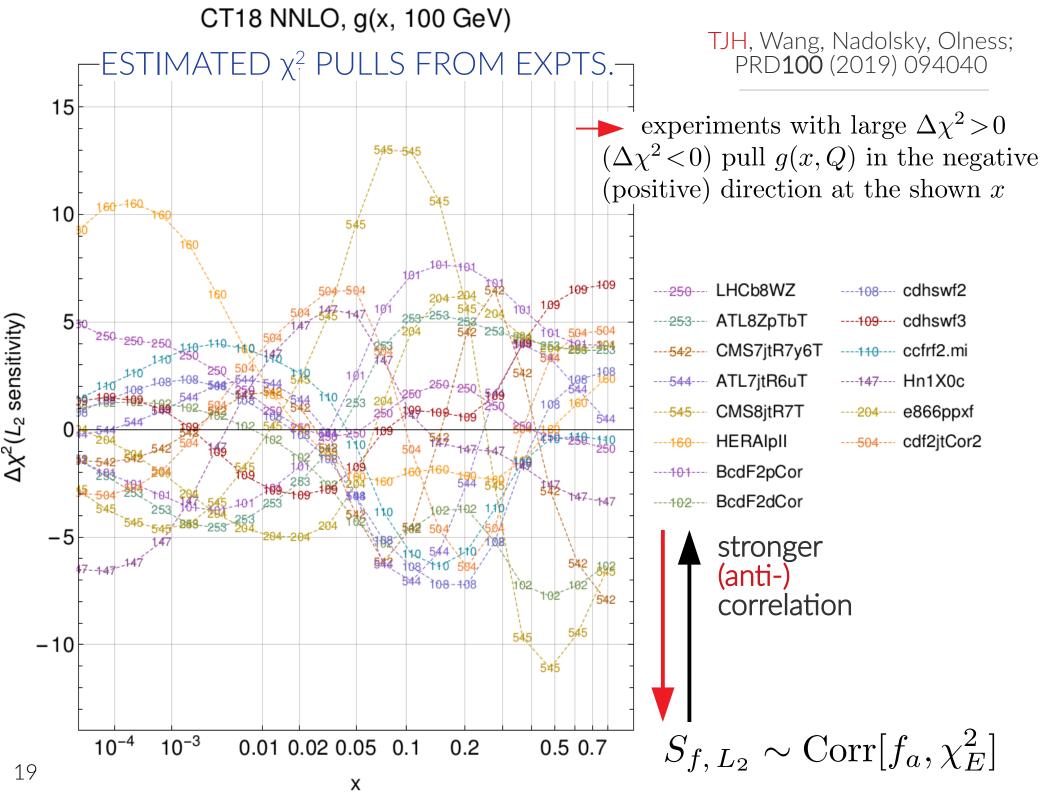
sensitivities can be aggregated for direct comparisons of exps



conclusions, recommendations

- the Yellow Initiative has an aggressive timeline; necessitates quick action
 - → need: an efficient, effective workflow to connect EIC design scenarios to physics output
 - → in conjunction with fitting, fast methods here can inform/integrate into this workflow
 - → we need guidelines on EIC operation as initial input
 - performance in energy, lumi, systs (baseline ↔ optimistic scenarios)
 - → identify 'lumi-hungry' vs. syst-limited measurements; explore error scenarios
- inclusive DIS will constrain: glue (scaling violations); d-PDF (high-lumi CC); lattice-calculable quantities; ...
 - → not sufficient alone for flavor sep.; need input from SIDIS/HQ production
 - ... IRG activity must coordinate with other WGs!

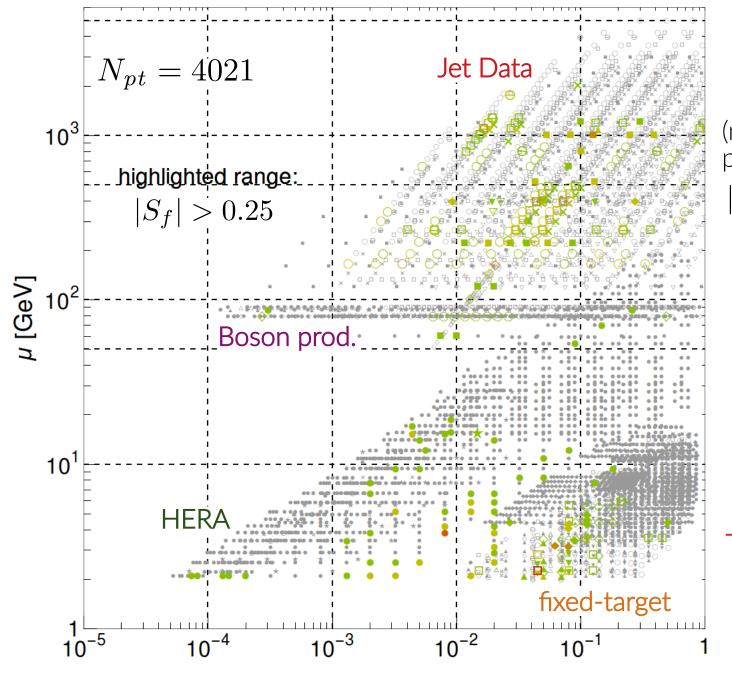
supplementary material ———





B.-T. Wang, TJH, S. Doyle, J. Gao, T.-J. Hou, P. M. Nadolsky, F. I. Olness

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X

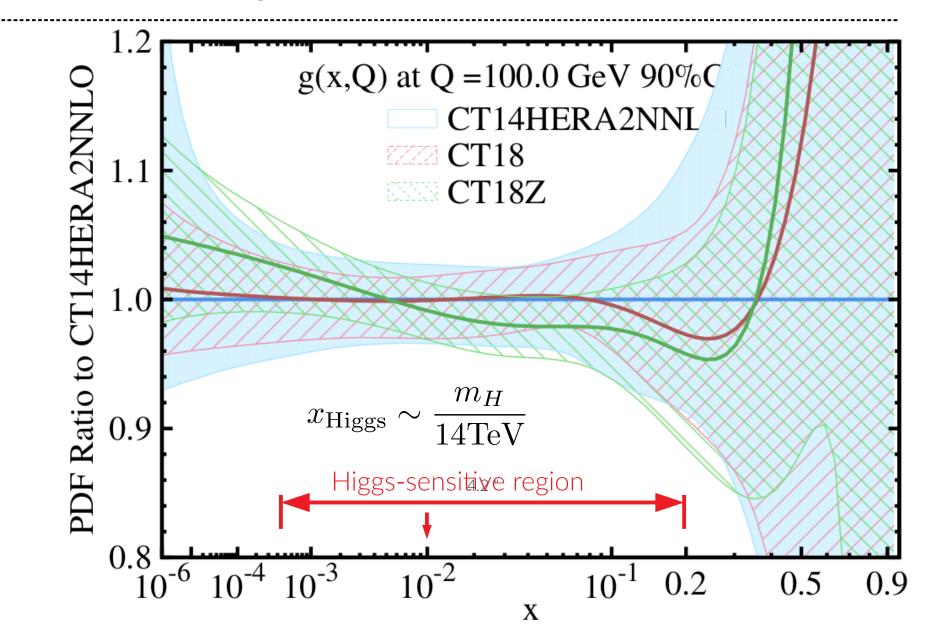
(magnitude of PDF pull of each datum)

 $|S_f|$



after the aggregated HERA data, inclusive jet production – greatest total sensitivity!

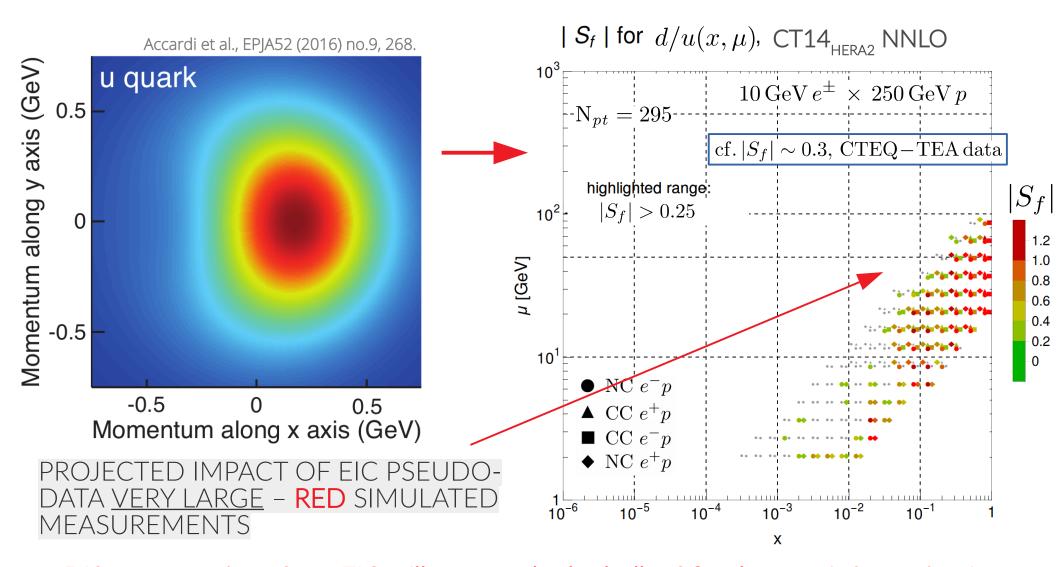
▶ large correlations for E866, BCDMS, CCFR, CMS WASY, Z p_T and ttbar production, but smaller numbers of highly-sensitive points



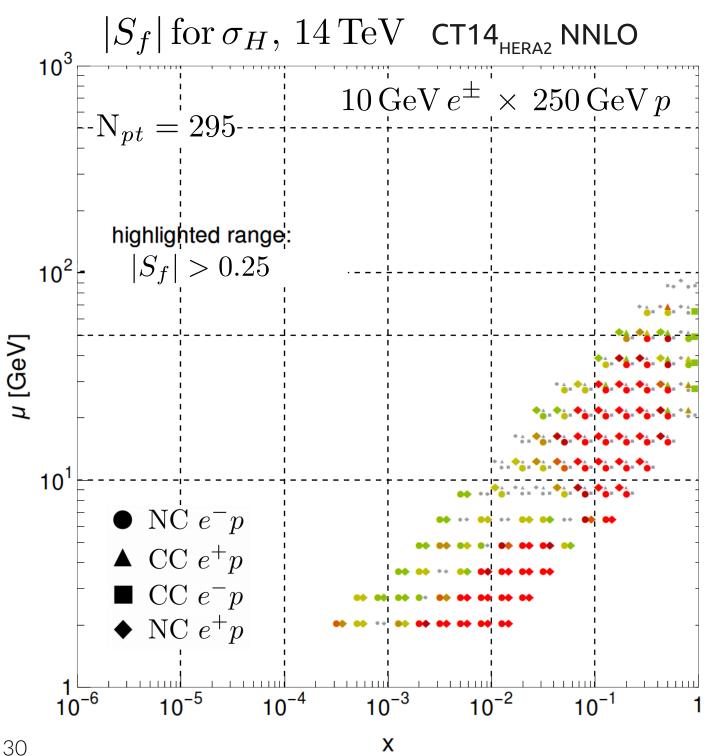
• while LHC Run-1 data drive important PDF improvements, including for the gluon at high-, low-x, the effect is relatively incremental

the EIC tomography program will deliver high-precision DIS

 by measuring the nucleon's multi-dimensional wave function with high precision, the EIC will hugely constrain proton collinear structure



 DIS cross sections from EIC will supercede the bulk of fixed-target information in contemporary QCD fits; provide an 'anchor-point' to resolve systematic PDF tensions



strong predicted impact on the Higgs sector

 $|S_f|$

1.2

1.0

8.0

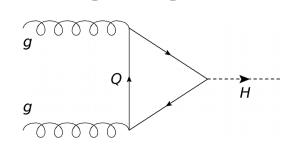
0.6

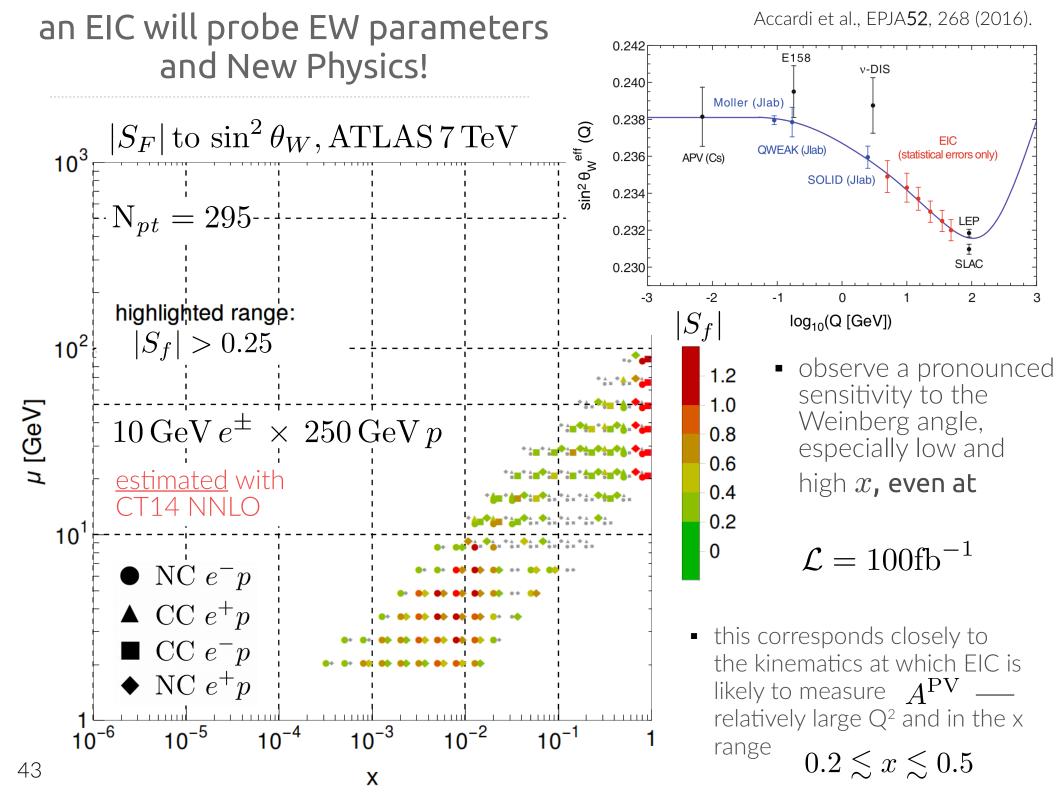
0.4

0.2

the impact of an EIC upon the theoretical predictions for inclusive Higgs production arises from a very broad region of the kinematical space it can access

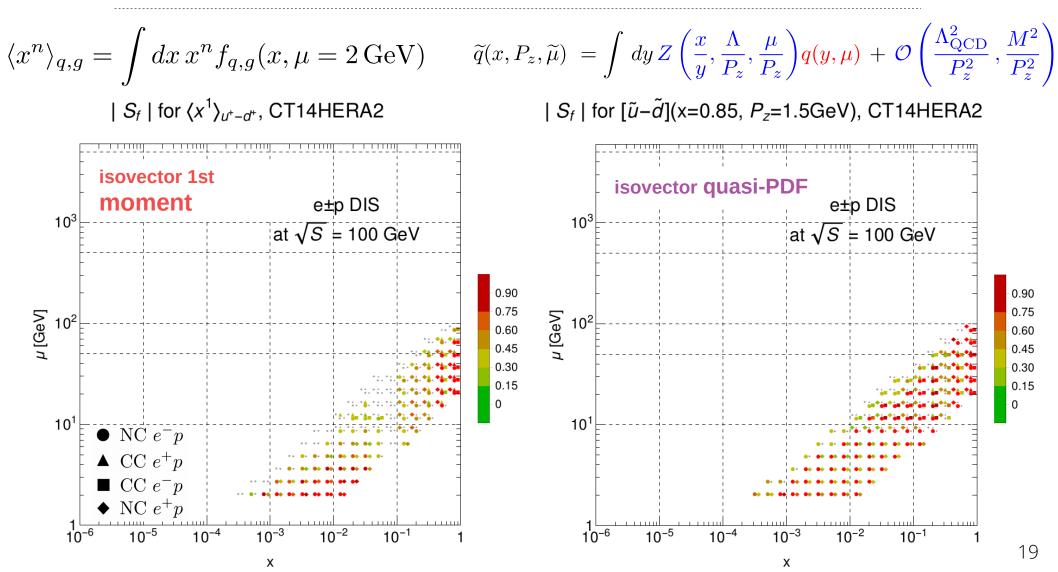
impact rather closely tied to that of the integrated gluon PDF:



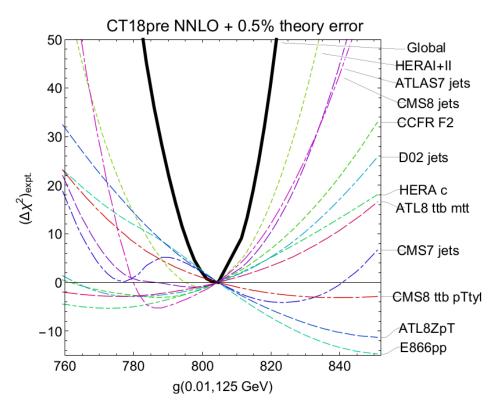


An EIC would drive lattice phenomenology

- A high-luminosity lepton-hadron collider will impose very tight constraints on many lattice observables; below, the isovector first moment and qPDF; this is crucial for benchmarking!
- Many of the experiments most sensitive to PDF Mellin moments and qPDFs involve nuclear targets → eA data from EIC would sharpen knowledge of nuclear corrections



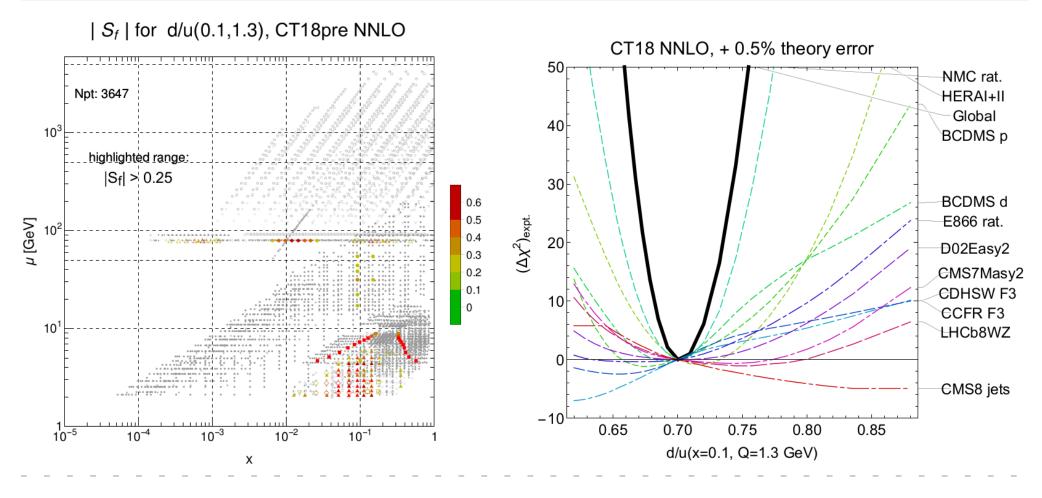
...for the gluon PDF in the Higgs region, $g(0.01,m_H)$



$g(x=0.01, \mu=125 \text{ GeV})$		
PDFSense		LM scan
CT14HERA2	CT18pre	CT18pre
HERAI+II'15	HERAI+II'15	HERAI+II'15
CMS8jets'17	CMS8jets'17	CMS8jets'17
CMS7jets'14	CMS7jets'14	ATL8ZpT'16
ATLAS7jets'15	E866pp'03	E866pp'03
E866pp'03	ATLAS7jets'15	ATLAS7jets'15
BCDMSd'90	BCDMSd'90	CCFR-F2'01
CCFR-F3'97	BCDMSp'89	D02jets'08
D02jets'08	D02jets'08	HERAc'13
NMCrat'97	NMCrat'97	NuTeV-nub'06
BCDMSp'89	CDHSW-F2'91	CCFR-F3'97

• PDFSense identifies the most sensitive experiments with high confidence and in accord with other methods such as the LM scans. It works the best when the uncertainties are nearly Gaussian, and experimental constraints agree among themselves [arXiv:1803.02777]

PDFSense predictions can be validated against actual fits



- PDFSense successfully predicts the highest impact data sets before fitting, as shown in this illustration for the large x PDF ratio $\,d/u\,$
- Lagrange Multiplier scans provide an independent test of which datasets most drive the global fit in connection with specific PDFs

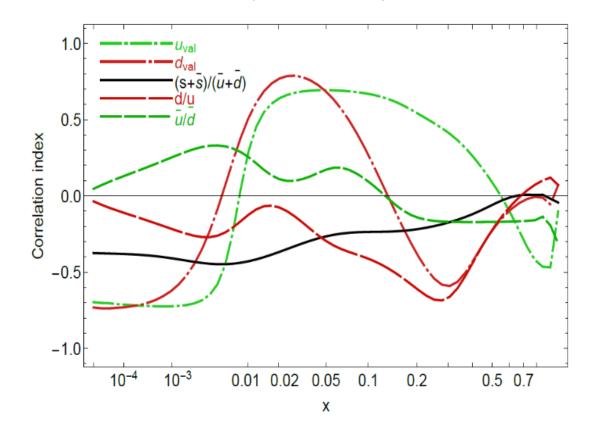
HERA and fixed-target (BCDMS, NMC) data are dominant!

...as a follow-on to Alesandro's EW-focused overview:

important PDF correlations for the ATLAS extraction of $\sin^2 heta_W$

Example: $\sin^2 \theta_{weak} \equiv s2w$ measured by ATLAS 8 TeV

Correlation, $\sin \theta_w$ (ATLAS 8 TeV CB) and f(x,Q) at Q=81.45 GeV 2018/11/11, PRELIMINARY, CT14 NNLO



Strongest correlations of s2w with u_{val} , d_{val} at $0.005 \lesssim x \lesssim 0.2$

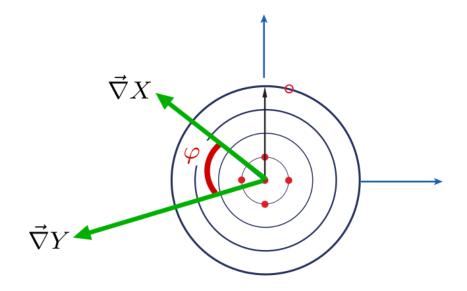
weak correlations with \bar{u} , \bar{d} , \bar{s} , g

 u_{val} , d_{val} changed between CT10 and CT14 [1506.07433, Sec. 2B]

It is instructive to explore the data pulls on $u_{val},\,d_{val}$

rather than the costly LM scans, we can examine a "cheaper" measure which yields comparable information

the L2 sensitivity

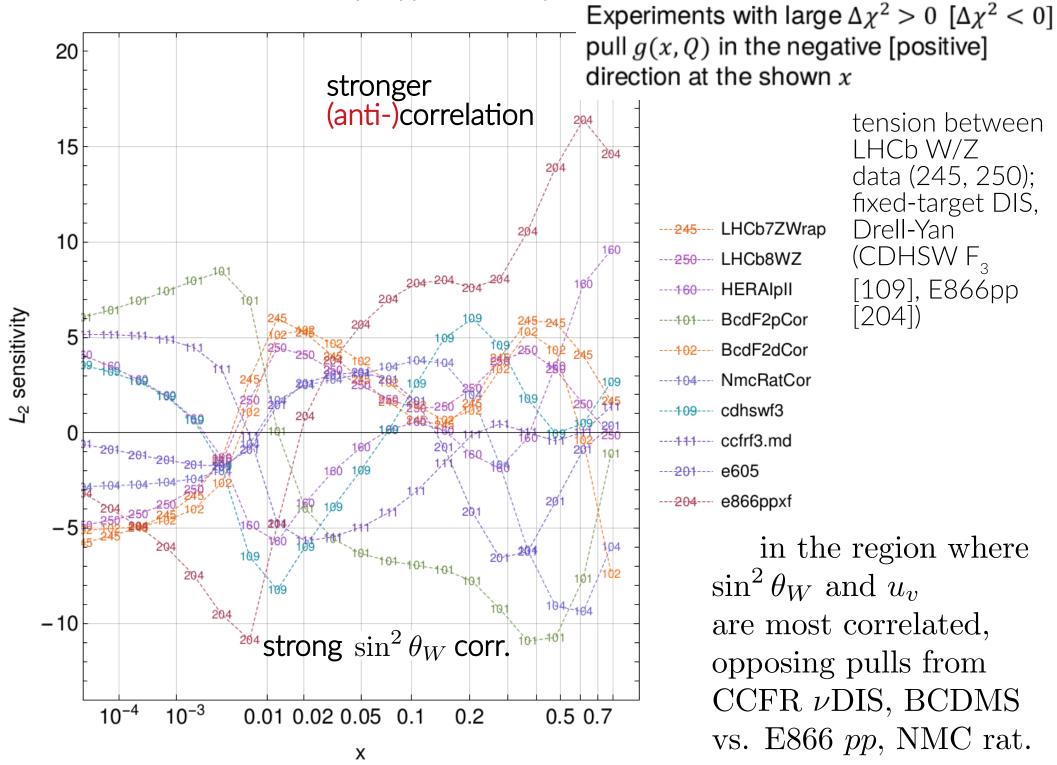


 L_2 sensitivity. Take $X=f_a(x_i,Q_i)$ or $\sigma(f)$; $Y=\chi_E^2$ for experiment E. Find $\Delta Y(\vec{z}_{m,X})$ for the displacement $|\vec{z}_{m,X}|=1$ along the direction $|\vec{\nabla}X|/|\vec{\nabla}X|$ (corresponding to $\Delta\chi_{tot}^2=T^2$ and $X(\vec{z})=X(0)+\Delta X$):

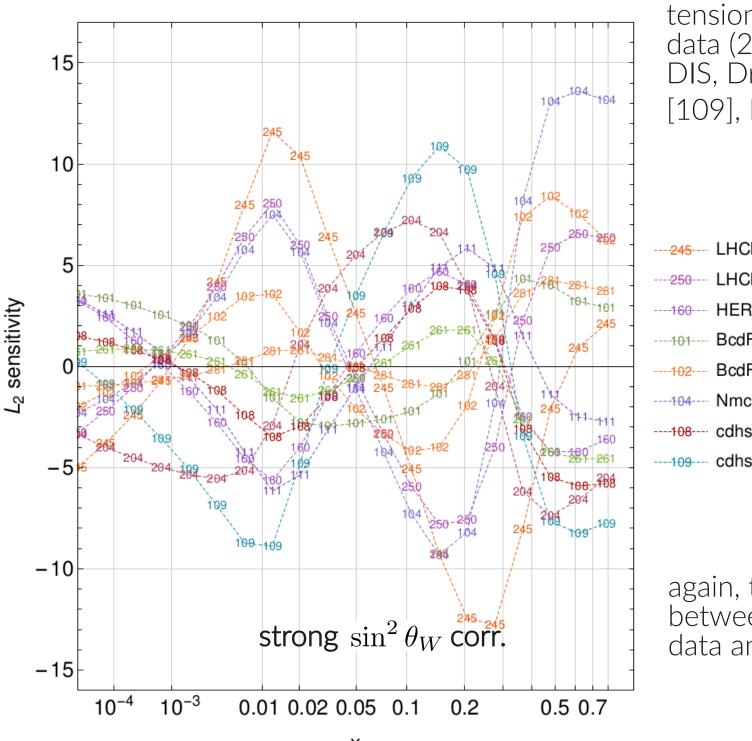
$$S_{f,L_2} \equiv \Delta Y(ec{z}_{m,X}) = ec{
abla} Y \cdot ec{z}_{m,X} = ec{
abla} Y \cdot rac{ec{
abla} X}{|
abla X|}$$
 or, $\sim \operatorname{Corr}[f_a,\chi_E^2]$ $= \Delta Y \, \cos arphi$

...extent to which total χ^2_F of specific expts. correlates with x-dep. of PDFs

CT18 NNLO, $u_V(x,Q)(x, 100 \text{ GeV})$



CT18 NNLO, $d_V(x,Q)(x, 100 \text{ GeV})$

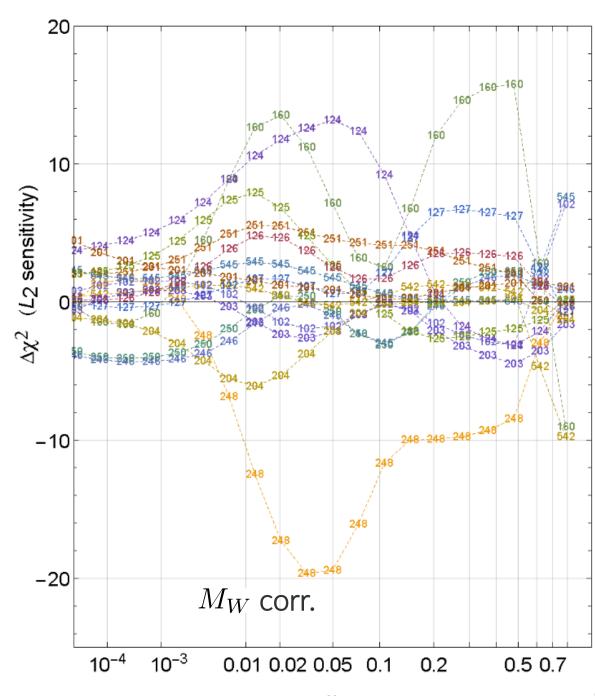


tension between LHCb W/Z data (245, 250); fixed-target DIS, Drell-Yan (CDHSW F₃ [109], E866pp [204])

---245--- LHCb7ZWrap ---111--- ccfrf3.md
---250--- LHCb8WZ ---204--- e866ppxf
---160--- HERAIpII ---261--- ZyCDF2
---101--- BcdF2pCor ---281--- d02Easy5
---102--- BcdF2dCor
---104--- NmcRatCor
---108--- cdhswf2
---109--- cdhswf3

again, tensions observed between, e.g., NMC ratio data and CDHSW, E866pp

CT18Z NNLO, s(x, 2 GeV)



L_2 sensitivity, strangeness: CT18Z

Most sensitive experiments

```
----246--- LHCb8Zeer ----124--- NuTvNuChXN
----248--- ATL7ZW.xF ----125--- NuTvNbChXN
----250--- LHCb8WZ ----126--- CcfrNuChXN
----251--- ATL8DY ----127--- CcfrNbChXN
----542--- CMS7jtR7y6T ----201--- e605
----545--- CMS8jtR7T ----203--- e866f
----160--- HERAIpII ----204--- e866ppxf
----102---- BcdF2dCor
```

A tension trend between DIS (HERA I+II, CCFR, NuTeV) and Drell-Yan (ATLAS 7 Z/W, LHCb W/Z, E866 pp, ...) experiments

pronounced effect of ATLAS 7 TeV Z/W data!

a brief statistical aside, i

 the CTEQ-TEA global analysis relies on the Hessian formalism for its error treatment

$$\chi_E^2(\vec{a}) = \sum_{i=1}^{N_{pt}} r_i^2(\vec{a}) + \sum_{i=1}^{N_{\lambda}} \overline{\lambda}_{\alpha}^2(\vec{a}) \qquad \blacksquare$$

nuisance parameters to handle correlated errors

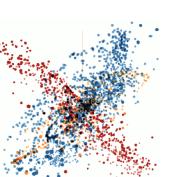
$$r_i(\vec{a}) = \frac{1}{s_i} \left(T_i(\vec{a}) - D_{i,sh}(\vec{a}) \right)$$

these result in systematic shifts to data central values:

$$D_i \to D_{i,sh}(\vec{a}) = D_i - \sum_{\alpha=1}^{N_{\lambda}} \beta_{i\alpha} \overline{\lambda}_{\alpha}(\vec{a})$$

ullet a 56-dimensional parametric basis $ec{a}$ is obtained by diagonalizing the Hessian

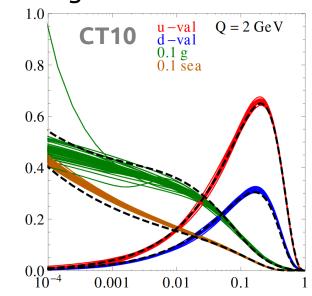
matrix H determined from χ^2 (following a 28-parameter fit)



use this basis to compute 56component "normalized" residuals :

$$\delta_{i,l}^{\pm} \equiv \left(r_i(\vec{a}_l^{\pm}) - r_i(\vec{a}_0) \right) / \langle r_0 \rangle_E$$

where
$$\langle r_0
angle_E \equiv \sqrt{rac{1}{N_{pt}} \sum_{i=1}^{N_{pt}} r_i^2(ec{a}_0)}$$

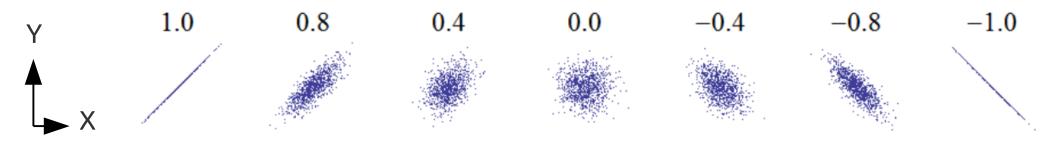


a brief statistical aside, ii

• ... but how does the behavior of these residuals relate to the fitted PDFs and their uncertainties?

for example, how does the PDF uncertainty (at specific x, μ) correlate with the residual associated with a theoretical prediction at the same x, μ ?

examine the Pearson correlation over the 56-member PDF error set between a PDF of given flavor and the residual



[X,Y] are exactly (anti-)correlated at the far (right) left above.

• we may then evaluate correlations between arbitrary PDF-derived quantities over the ensemble of error sets ([X,Y] may be PDFs, cross sections, residuals,...):

$$Corr[X, Y] = \frac{1}{4\Delta X \Delta Y} \sum_{j=1}^{N} (X_j^+ - X_j^-)(Y_j^+ - Y_j^-) \qquad \Delta X = \frac{1}{2} \sqrt{\sum_{j=1}^{N} (X_j^+ - X_j^-)^2}$$

Correlation C_f and sensitivity S_f

The relation of data point i on the PDF dependence of f can be estimated by:

• $C_f \equiv \text{Corr}[\rho_i(\vec{a})), f(\vec{a})] = \cos\varphi$

 $\vec{\rho}_i \equiv \vec{\nabla} r_i / \langle r_0 \rangle_E$ -- gradient of r_i normalized to the r.m.s. average residual in expt E;

$$\left(\vec{\nabla}r_i\right)_k = \left(r_i(\vec{a}_k^+) - r_i(\vec{a}_k^-)\right)/2$$

$$Corr[X, Y] = \frac{1}{4\Delta X \Delta Y} \sum_{j=1}^{N} (X_j^+ - X_j^-)(Y_j^+ - Y_j^-)$$

 C_f is **independent** of the experimental and PDF uncertainties. In the figures, take $|C_f| \gtrsim 0.7$ to indicate a large correlation.

•
$$S_f \equiv |\vec{\rho}_i| cos \varphi = C_f \frac{\Delta r_i}{\langle r_0 \rangle_E}$$
 -- projection of $\vec{\rho}_i(\vec{a})$ on $\vec{\nabla} f$

 S_f is proportional to $\cos\varphi$ and the ratio of the PDF uncertainty to the experimental uncertainty. We can sum $|S_f|$. In the figures, take $|S_f| > 0.25$ to be significant.

2nd aside: kinematical matchings

 residual-PDF correlations and sensitivities are evaluated at parton-level kinematics determined according to leading-order matchings with physical scales in measurements

deeply-inelastic scattering:

$$\mu_i \approx Q|_i, \ x_i \approx x_B|_i$$

 x_i : parton mom. fraction

 μ_i : factorization scale

hadron-hadron collisions:

$$AB \to CX$$

$$\mu_i \approx Q|_i, \ x_i^{\pm} \approx \left. \frac{Q}{\sqrt{s}} \exp(\pm y_C) \right|_i$$

single-inclusive jet production:

$$Q = 2p_{Tj}, \ y_C = y_j$$

$$tar{t}$$
 pair production:

$$t \bar{t}$$
 pair production: $Q = m_{t \bar{t}}, \ y_C = y_{t \bar{t}}$

etc...

$$d\sigma/dp_T^Z$$
 measurements:

$$d\sigma/dp_T^Z$$
 measurements: $Q=\sqrt{(p_T^Z)^2+(M_Z)^2},\ y_C=y_Z$