# Experimental Measurements at the EIC

Proton

Neutron

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Lecture III

#### Questions from Wednesday

- 1. In a nutshell could you explain the concept of factorization one more time?
- 2. Why J M equation doesn't include the transverse spin contribution, but only the helicities? I naively think that it gives 3D spin contribution, which should include transverse contribution as well.
- 3. What can the EIC tell us about the helicity strange distribution?
- 4. Can parton helicities be further constrained by investigating GPDs via DVCS processes?
- 5. What type of eID and PID detectors are being proposed?

1. In a nutshell could you explain the concept of factorization one more time?

$$\sigma(ep o e'X) \propto fig(x,Q^2ig) \hat{\sigma}ig(\hat{s},\hat{t}\,,\hat{u}ig)$$

A factorization proof permits an observable, such as a cross-section or asymmetry, to be written as the multiplication (convolution) of two independent functions. In this case the parton PDF and the partonic cross-section.



$$p \xrightarrow{f_1^p} f_1^{p} \xrightarrow{f_1} f_1$$

$$\hat{\sigma}$$

$$p \xrightarrow{f_2} f_2^{p} \xrightarrow{f_2} f_2^{p}$$

$$\sigma(pp
ightarrow \pi X) \propto f_1ig(x,Q^2ig) f_2ig(x,Q^2ig) \hat{\sigma}ig(\hat{s},t,\hat{u}ig) D_\piig(z,Q^2ig)$$

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$$p \xrightarrow{f_1^p} f_1$$

$$f \xrightarrow{f_1} f_1$$

$$\hat{\sigma}$$

$$f_2$$

$$p \xrightarrow{f_2} f_2$$

$$\sigma(pp
ightarrow \pi X) \propto f_1ig(x,Q^2ig) f_2ig(x,Q^2ig) \hat{\sigma}ig(\hat{s},t,\hat{u}ig) D_\piig(z,Q^2ig)$$

If factorization holds for both processes then the PDFs should be universal. Same is true for FF.

2. Why J M equation doesn't include the transverse spin contribution, but only the helicities? I naively think that it gives 3D spin contribution, which should include transverse contribution as well.

At leading twist in a collinear framework there are three functions that completely define partonic kinematics Nucleon P



Number density of partons of flavor f with momentum fraction x inside a nucleon





Number density of longitudinally polarized partons inside longitudinally polarized nucleons (Helicity)





Number density of transversely polarized partons inside a transversely polarized nucleon (Transversity)



2. Why J M equation doesn't include the transverse spin contribution, but only the helicities? I naively think that it gives 3D spin contribution, which should include transverse contribution as well.

#### The Jaffe-Manohar Sum rule is formulated in the context of the helicity distributions.

R. L. Jaffe, A. Manohar, The G(1) problem: fact and fantasy on the spin of the proton, Nucl. Phys. B 337 (1990) 509–546.

$$rac{1}{2}=rac{1}{2}\Delta\Sigma(\mu)+1\Delta G(\mu)+L_{Q+G}(\mu)$$



Number density of longitudinally polarized partons inside longitudinally polarized nucleons (Helicity)





It is not possible to use J M formulation for the "transverse" spin sum rule.

# And here is your reward for all you hard work!

- 1. Built detectors to measure E and p of electron.
- 2. Correctly identified e- that took part in hard scattering interaction.
- 3. Spin sorted the electrons.
- 4. Measured the relative luminosities.
- 5. Measured the beam polarizations.
- 6. Defined my analysis bins to maximize purity and stability.

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- 7. Applied radiative and other kinematic corrections.
- 8. Extract  $g_{1}^{P}$  in my bins of x and  $Q^{2}$ !



#### Constraints on Helicity Distributions from Inclusive DIS



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#### How low can we go? Experiments never go low enough ...

Work by Pitonyak, Sievert and Kovchegov point to a new evolution in x that can *predict* low x gluon helicity contribution.

New constraints by JAMsmallx (2102.06159)



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JAMsmallx+EIC

 $\Delta u^+$ 

 $\Delta s^+$ 

0.2

0.1

0.0

 $Q^2 = 10 \,\mathrm{GeV^2}$ 

#### Constraints on OAM from Inclusive DIS



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## **Parity Violating Asymmetries**

- Inclusive DIS has contributions from both virtual photon and Z<sup>0</sup> exchange.
- Can isolate the contributions from γ\* Z<sup>0</sup> interference by looking at single spin asymmetries





Spin sort only by proton beam, integrate over e- beam.

## Parity Violating Asymmetries

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- Can isolate the contributions from  $y^*$  Z<sup>0</sup> • interference by looking at single spin asymmetries





Spin sort only by proton beam, integrate over e- beam.

> Structure functions from virtual photon exchange.



$$p$$

$$(X:\Lambda \to n\pi^0)$$

$$g_1^{\text{proton, }\gamma Z} \approx \frac{1}{9} (\Delta u + \Delta \bar{u} + \Delta d + \Delta \bar{d} + \Delta s + \Delta \bar{s} + \Delta c + \Delta \bar{c})$$
$$g_5^{\text{proton, }\gamma Z} = \frac{1}{3} (\Delta u_V + \Delta c - \Delta \bar{c}) + \frac{1}{6} (\Delta d_V + \Delta s - \Delta \bar{s})$$

## Parity Violating Asymmetries

$$\begin{split} A_{\rm PV}^{\rm hadron} &= \frac{\sigma^{(+)} - \sigma^{(-)}}{\sigma^{(+)} + \sigma^{(-)}} \\ &= \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} [g_V^e \frac{g_5^{\gamma Z}}{F_1^{\gamma}} + g_A^e \frac{Y_-}{Y_+} \frac{g_1^{\gamma Z}}{F_1^{\gamma}}]. \end{split}$$



 $10^{-1}$ 

 $x_{\min}$ 

Constraints depend on source of axial coupling constant.



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$$F_1^{\text{proton, }\gamma Z} \approx \frac{1}{9} (u + \bar{u} + d + \bar{d} + s + \bar{s} + c + \bar{c}),$$
  
$$F_3^{\text{proton, }\gamma Z} = \frac{2}{3} (u_V + c - \bar{c}) + \frac{1}{3} (d_V + s - \bar{s}),$$



 $10^{-2}$ 

 $10^{-1}$ 

r



Potential to constraint strange contribution depends on electron ID systematics.

#### slido



## How can we separate quark and anti-quark helicity distributions?

(i) Start presenting to display the poll results on this slide.

#### Semi-inclusive Deep Inelastic Scattering (SIDIS)

- Identify and reconstruct scattered electron as in inclusive DIS and use it to determine kinematics
- In addition to e- tag final state hadrons.
- Hadrons serve as quark flavor filter. For example:
  - $\pi^+ \rightarrow$ u+dbar  $\pi^- \rightarrow$ d+ubar
  - $K+ \rightarrow u+sbar$   $K- \rightarrow s+ubar$
- Requires knowledge about fragmentation functions  $D(z, Q^2)$  and this introduces additional uncertainties.
- Z is the momentum fraction of the scattered quark carried by the hadron. The correlation between the hadron and the fragmenting parton's flavor increases as z → 1.
- Low z hadrons are more likely to originate from the proton remnant, rather than the hard interaction.



#### Where do the hadrons live?



#### Hadron PID → velocity

- 1. Time-of-Flight (TOF)
- Direct measurement of of particle velocity
- Required tracking to measures the momentum p
- Extract  $\beta$  from scintillator/silicon detector plus start time detector
- Mass =  $\beta \gamma / p$
- 2. Velocity dependent interactions with materials





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• mRICH - modular Ring Imaging Cerenkov Detector. Aerogel based with Frensel lens. Used for both electron ID and PID.





- mRICH modular Ring Imaging Cerenkov Detector. Aerogel based with Frensel lens.
- mRICH eID out to ~ 2GeV
- mRICH pi/K PID out to ~7GeV





Table 8.6 Requested PID momentum coverage for 3σ pion/kaon separation. (Yellow Report)

Pseudorapidity Range	Momentum Range
$-3.5 < \eta < -1.0$	$\leq$ 7 GeV/c
$-1.0 < \eta < 0.5$	$\leq 10  \text{GeV}/c$
$0.5 < \eta < 1.0$	$\leq 15  \text{GeV}/c$
$1.0 < \eta < 1.5$	$\leq$ 30 GeV/c
$1.5 < \eta < 2.5$	$\leq 50  \text{GeV}/c$
$2.5 < \eta < 3.0$	$\leq$ 30 GeV/c
$3.0 < \eta < 3.5$	$\leq 20  \mathrm{GeV}/c$

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- DIRC Detection of internally reflected Cerenkov
- DIRC use material (like quartz) with index n > 2, allows for total internal reflection and propagation to light to end of material.
- DIRC allow for pi/K separation up to 6 GeV with very thin detector
- hpDIRC "high performance" DIRC provides ring focusing and therefore improved resolution





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- dRICH "Duel" RICH
- dRICH uses two different radiators to allow for full momentum coverage
- dRICH takes the most real estate



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Extends PID coverage to lower momentum range and complements the RICH-based PID coverage. In the barrel region, hpDIRC and TOF provide an overlapping coverage from ~200 MeV/c to 2 GeV/c.

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#### Implementation in ECCE detector (very similar in ATHENA)



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#### Constraints from SIDIS at an EIC



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#### **Charged Current Interactions**

- At high  $Q^2$  the virtual photon is replaced with  $W^{+/-}$
- Scattered electron is replaced by neutrino, which goes undetected
- Rely on hadronic recoil for reconstruction of kinematics via Jaquet-Blondel method

$$Q_{JB}^2 = \frac{p_T^2}{1 - y_{JB}}$$
  $y_{JB} = \frac{(E - p^z)}{2E}$   $x_{JB} = \frac{Q_{JB}^2}{sy_{JB}}$ 

- Here p<sub>T</sub><sup>2</sup> and (E-p<sup>z</sup>) are summed over all hadrons in the final state. Challenge is to collect entire final state and to optimize resolution.
- Advantage is there are no fragmentation functions



#### Where is the CC final state?

- Requires both hadronic and electromagnetic calorimetry and tracking detectors in the forward region.
- Even with far forward coverage, many particles will be lost down the beamline.





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#### Single Spin Asymmetry A<sub>1</sub><sup>W</sup>

$$A_L^{W^-,N} \equiv \frac{d^2 \Delta \sigma^{W^-,N}/dxdy}{d^2 \sigma^{W^-,N}/dxdy}$$

The spin dependent CC cross-section for scattering of a left-handed electron (Wexchange) off of a longitudinally polarized nucleon target with helicity  $\lambda = \pm 1$ .

Discus relation: gL = g4 - 2xg5

g1 and g5 related to sea-quark distributions!

Very clean, but statistically limited observable.

$$\begin{aligned} \frac{d^2 \Delta \sigma^{W^-,N}}{dxdy} &= \\ &= \frac{1}{2} \left[ \frac{d^2 \sigma^{W^-,N} (\lambda_N = -1)}{dxdy} - \frac{d^2 \sigma^{W^-,N} (\lambda_N = +1)}{dxdy} \right] \\ &= \frac{2\pi \alpha_{em}^2}{xyQ^2} \eta \left[ 2Y_- x g_1^{W^-,N} - Y_+ g_4^{W^-,N} + y^2 g_L^{W^-,N} \right] (1) \end{aligned}$$

$$g_1^{W^-,p}(x) = \Delta u(x) + \Delta \bar{d}(x) + \Delta c(x) + \Delta \bar{s}(x) ,$$
  

$$g_5^{W^-,p}(x) = -\Delta u(x) + \Delta \bar{d}(x) - \Delta c(x) + \Delta \bar{s}(x)$$

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#### Summary of Lectures I-III

#### 1. Deep dive into experimental analysis - use helicity distributions as example

- a. Introduction to inclusive scattering
- b. Longitudinal spin asymmetries (single and double spin) and structure functions
- c. Electron identification and kinematic reconstruction
- d. Polarization, relative luminosity and radiative corrections
- e. Constraints from inclusive EIC measurement on helicity PDFs

#### 2. Flavor separation and sea quark helicity distributions

- a. Introduction to semi-inclusive scattering
- b. Hadron kinematics and PID
- c. Constraints from SIDIS EIC measurements on helicity PDFs
- d. Quick introduction to charge-current interactions and JB reconstruction.

# Lecture IV - Broaden our scope into Transverse Momentum Distributions (TMDs) and Generalized Parton Distributions (GPDs)