

Four Lectures on non-perturbative QCD, hadron structure and parton physics

Xiangdong Ji

University of Maryland

Outline of Lectures

- Lecture 1: What is the fundamental reason for building EIC? Or, why QCD is so difficult?
- Lecture 2: An amateur's introduction to lattice QCD
- Lecture 3: What is parton physics? (at EIC)
- Lecture 4: Large momentum effective theory for calculating partons from first principles

Lecture 1: What is the
fundamental reason to
build EIC? Or
why QCD is so difficult?

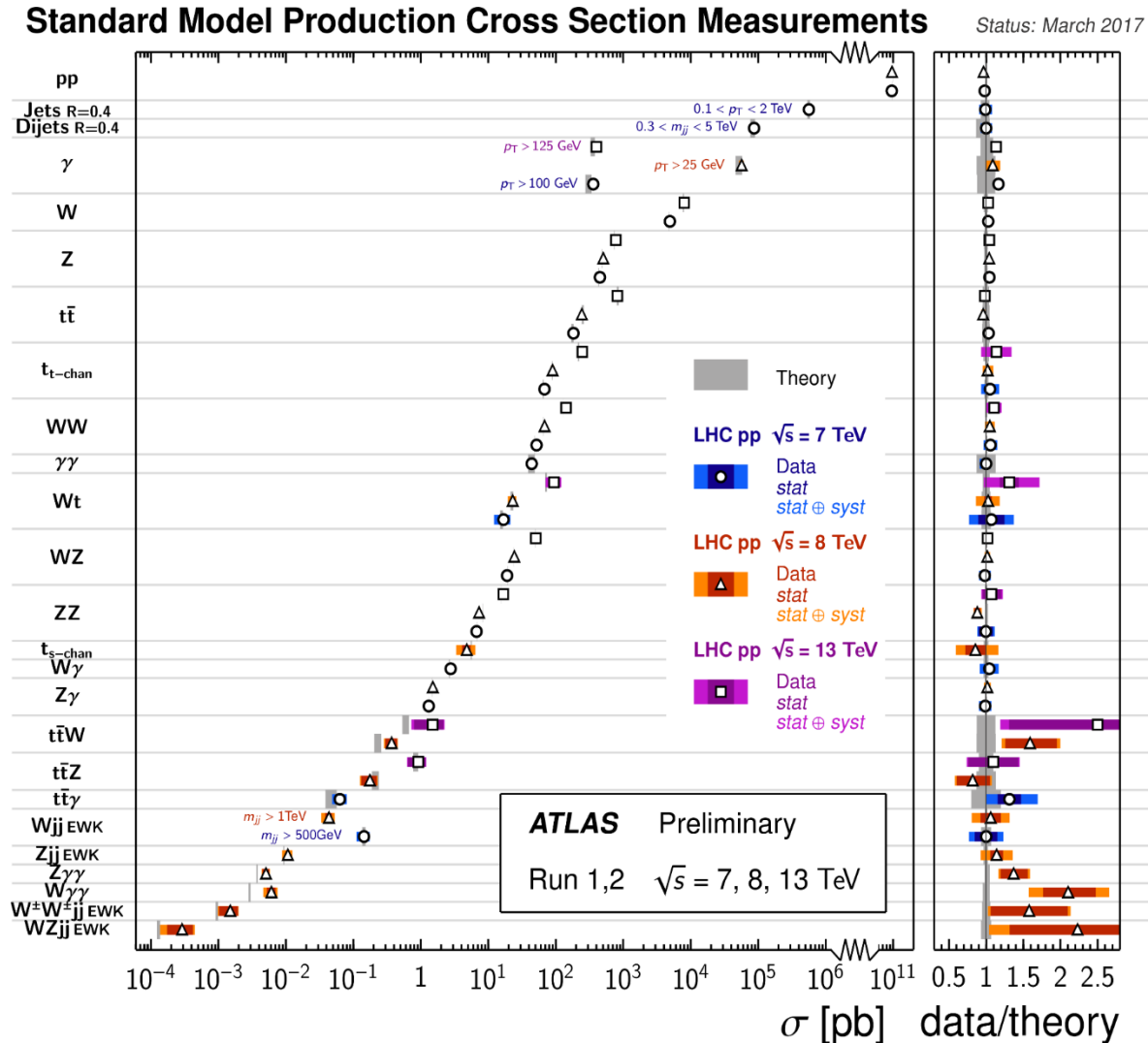
Standard model successes:

- The standard model itself has been hugely successful in explaining many physics phenomena

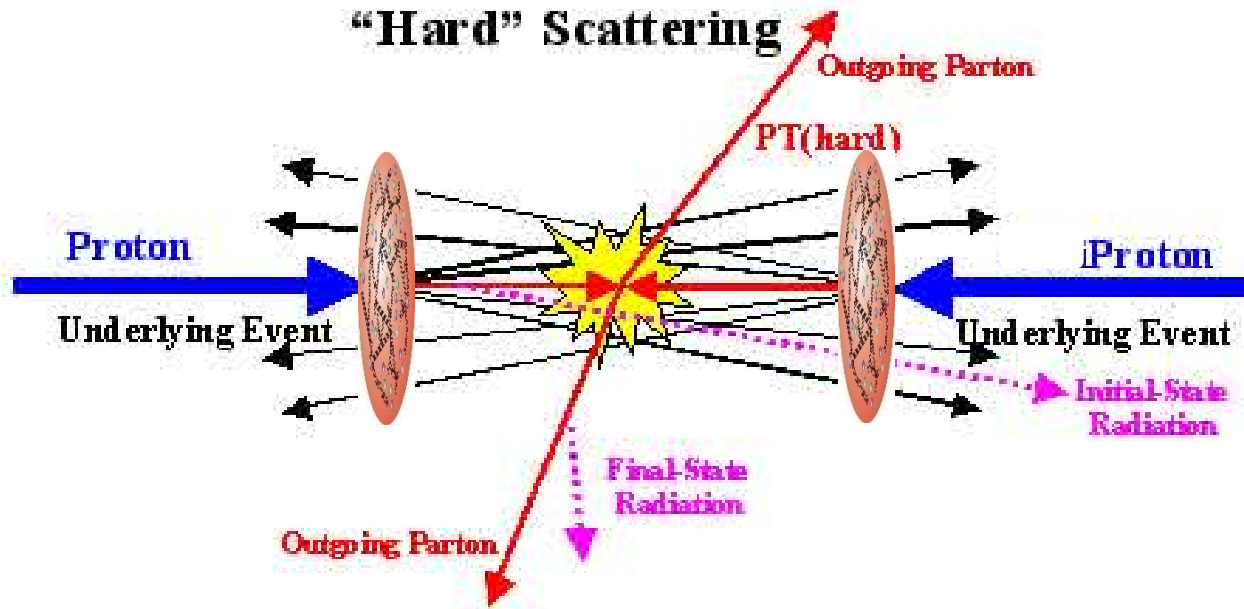
- Electroweak processes

- High-energy QCD processes

Perturbation theory works! (LHC)



Hard scattering theory

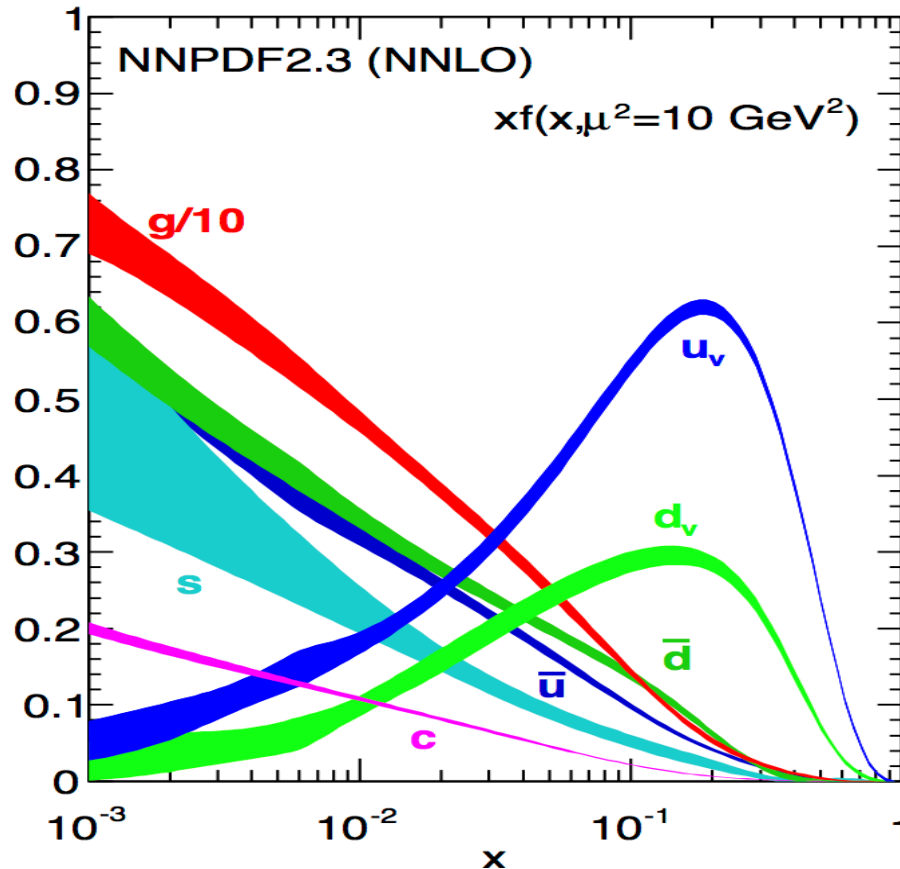


- **Factorization theorems:** The scattering cross sections are factorized in terms of PDFs and parton x-section.

$$\sigma = \int dx_a dx_b f_{a/A}(x_a) f_{b/B}(x_b) \hat{\sigma}$$

Phenomenological PDFs

- Use experimental data (~50 yrs) to extract PDFs



J. Gao, et al,
Phys. Rept. 742
(2018) 1-121

Structure of the proton

- To calculate PDFs from QCD, we need to understand the structure of the proton, **a quantum mechanical bound state**.
- Our experiences with bound states in QM are limited to mainly non-relativistic systems:
 - Atomic systems:** electrons + Coulomb int
 - Nuclear systems:** protons + neutrons

A lot of experience has been accumulated on non-relativistic many-body problems.

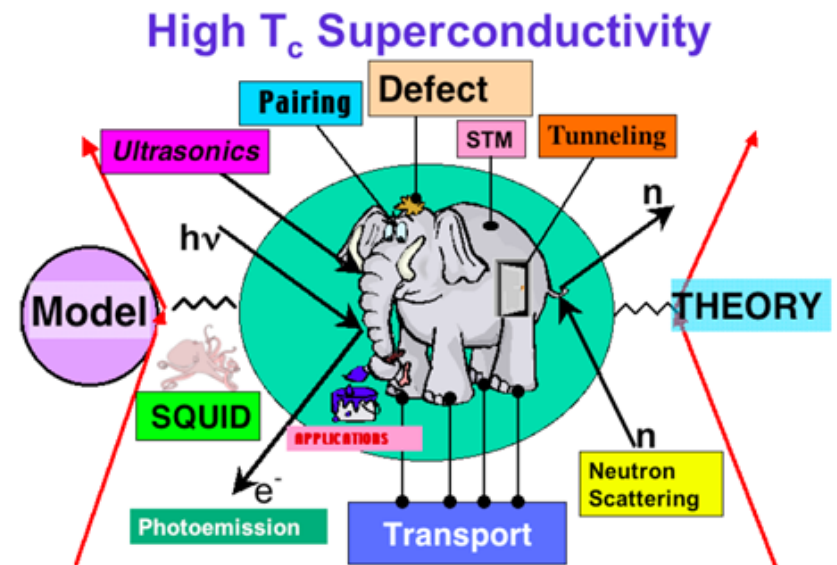
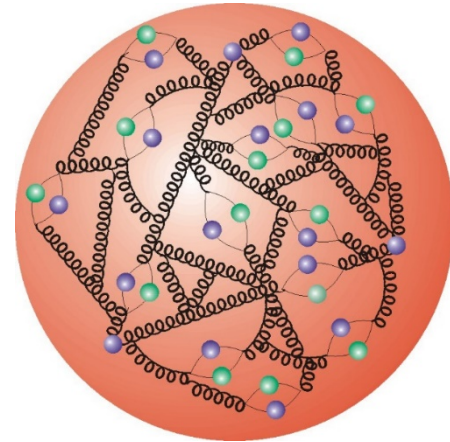
QCD bound states

- However, it is an unprecedented challenge to understand how QCD works at low energy, where theory becomes non-perturbative:

- Guts of Strong Interactions!

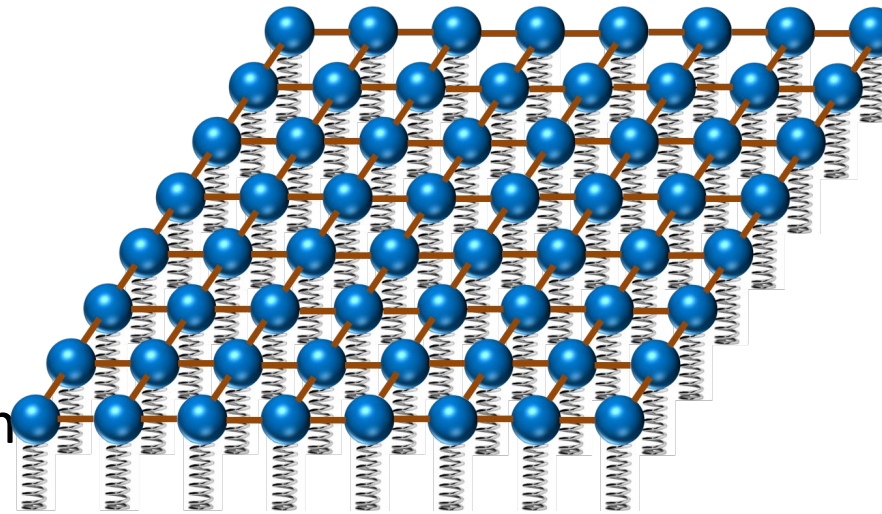
- Similar in nature to Condensed Matter Physics: the Lagrangian is known, but the solution is hard

High T_c , Hall effects, strongly coupled electron systems, etc



QCD (& SM) is a quantum field theory

- QFT = QM + Relativity
- The fundamental degrees of freedom are **quantized fields (quarks + gluons)**
- The fields fill in all space and time. Every point of space and time has a few quantum mechanical degrees of freedom (DOF), which are interacting with their neighboring DOFs (local QFT).
- Extremely complicated (infinite DOF) systems.....
Infinity is a problem



What is a particle?

- A particle is a quantized wave in a quantum field.



“Physical particles in QFT”



Two “easy” cases

- **Non-relativistic limit:**

Masses of the particles are much larger than the kinetic and potential energies.

Atomic and condensed matter physics are a non-relativistic limit of QED.

Nuclei physics is the non-relativistic limit of QCD.

Particle number is conserved. (For neutral systems, IR photon decouples)

QFT \rightarrow Hamiltonian dynamics

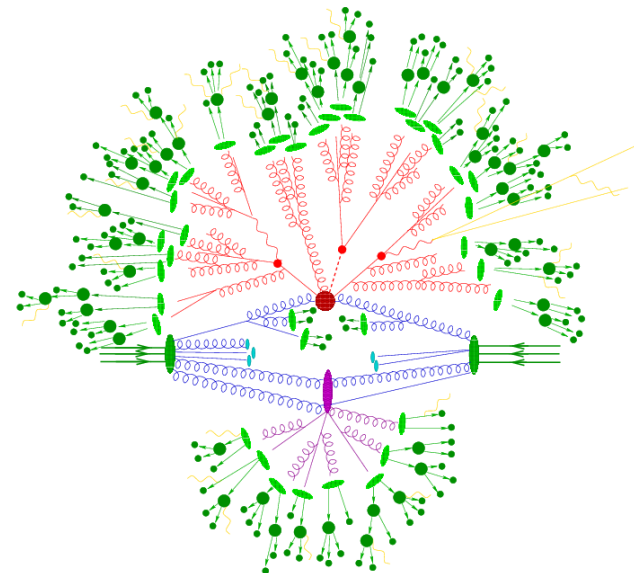
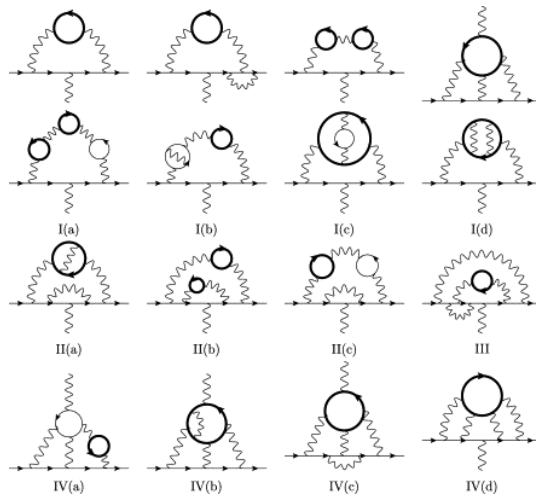
Two “easy” cases

- **Weak coupling limit**

In the weak coupling limit, one can use perturbation theory developed by Feynman et al.

High-loop calculations for g-2

Parton Shower



A remark about the non-relativistic limit

- Not all cases of the non-relativistic limit are easy.
- Heavy quarkonium systems involve non-perturbative potentials which cannot be easily calculated.
- Some level of perturbative expansion in the interactions is needed. (multiple photon effects are perturbative)

High-precision spectroscopy in H-atom

Why strong QCD is so difficult?

- **Strongly coupled: Similar to NR electron systems**
 - Non-perturbative approximation methods must be devised beyond fermi liquid theory.
 - Ab initio numerical simulations
 - Effective degrees of freedom?
- **Extra difficulty: Relativity**
 1. Center-of-mass and internal motions coupled
 2. The QCD vacuum (ground state)
 3. Number of particles not fixed (going to infinity)

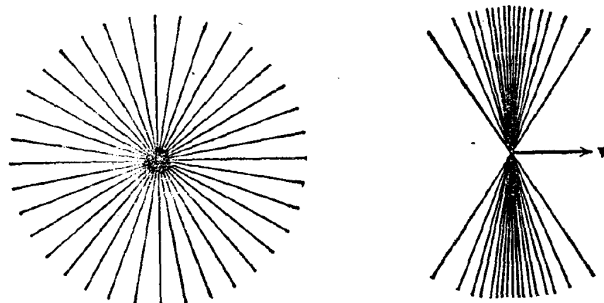
Relativity: internal states are frame-dependent

- The center-of-mass motion is part of the physics: the bound state has definite total momentum
- Because the boost operator is dynamical, the internal states are different at different momenta!

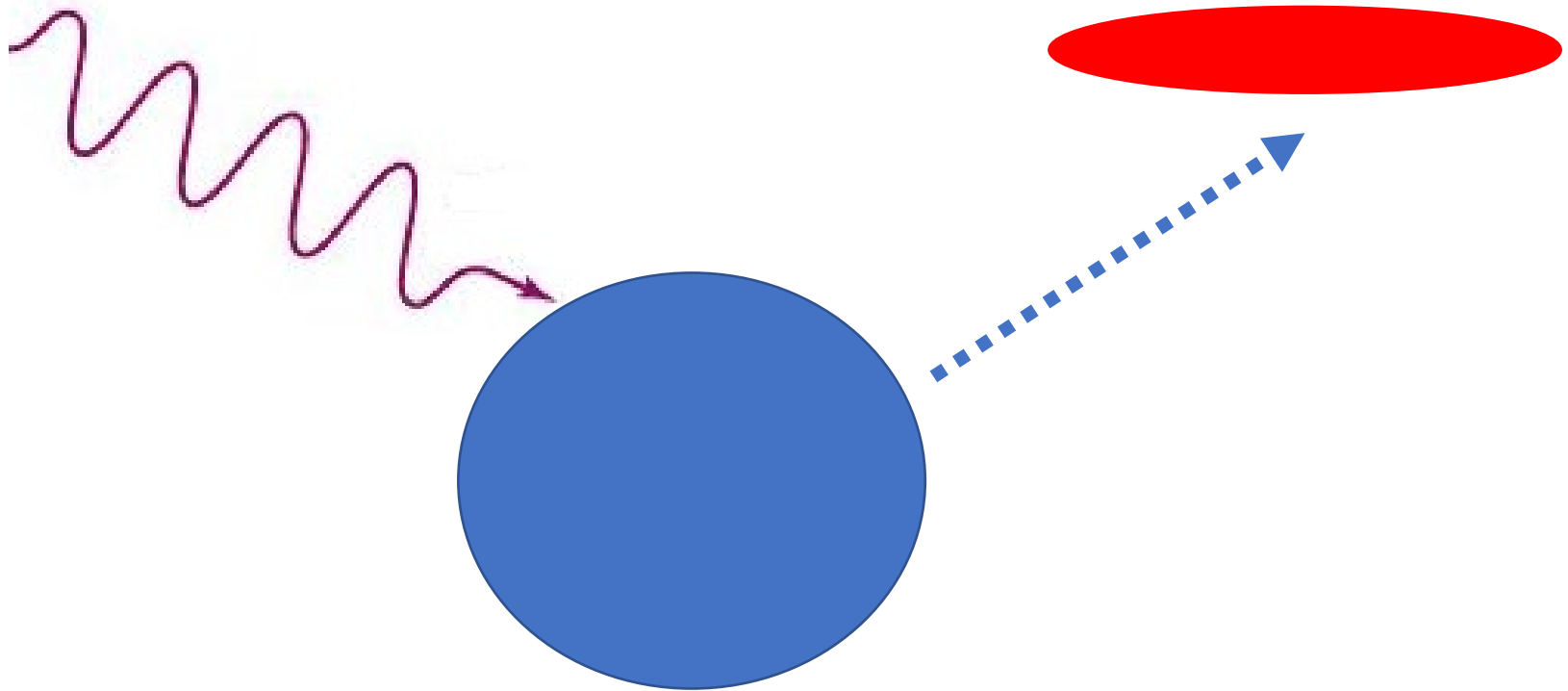
$$|p'\rangle = U(L)|p\rangle$$

where $|p'\rangle$ is different from $|p\rangle$ dynamically!

- The electromagnetic fields of a moving charge depends on its velocity or $\beta = v/c$



Elastic scattering: form factors
hadron states (wave functions) depending
on the external momenta



The QCD vacuum (ground state)

- Hadron systems are built upon the QCD vacuum which in itself is **extraordinary** complex
 - Similar to a strongly-interacting fermi sea in Condensed Matter Systems, where Landau's fermi liquid theory breaks down!
- And the hadron physics phenomena occur as **complex excitations** of this vacuum.
- It will be difficult to understand the excitations without understanding the ground state.

Understanding the water waves



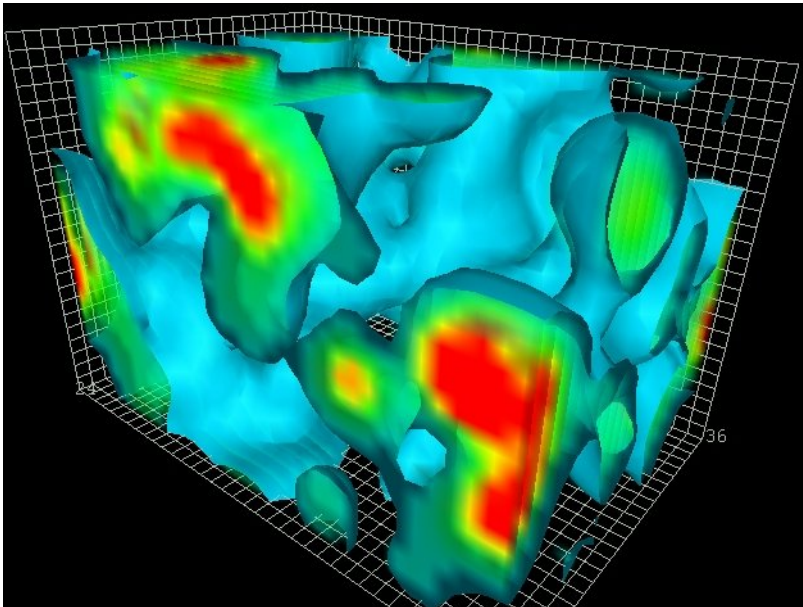
**Hadron physics that we
try to understand!**



**QCD vacuum that we
don't observe**

What does QCD vacuum look like?

- In semi-classical approximation, the vacuum is filled with interacting instantons/anti-instantons, a particular classical gluon configuration.



- The BPST instanton is an essentially non-perturbative classical solution of the Yang–Mills field equations
- $\vec{iE} = \vec{B}$, or $\vec{iE} = -\vec{B}$
- Instanton has zero energy density and zero angular momentum density

Vacuum properties: effective scalar field or gluon condensate

- Vacuum scalar field

$$\phi = \langle 0|F^2|0\rangle \sim \text{density of instantons}$$

The vacuum has a scalar field density which generates a dimension, Λ^4 , which sets the strong interaction scale parameter.

- The real QCD physics is independent of this scale parameter, unless there is a new scale introduced, such as quark masses.
- This scalar field is similar to the Higgs field in electroweak theory.

Vacuum properties: chiral symmetry & spontaneous breaking

- When N_f massless quarks q_i are introduced, there is a chiral symmetry

$$U_L(N_f) \times U_R(N_f) = U_V(1) \times U_A(1) \times SU_L(N_f) \times SU_R(N_f)$$

acting on the spaces of $q_{iL}(x)$ and $q_{iR}(x)$

- In the QCD vacuum, the symmetry is broken by instantons

$$U_L(N_f) \times U_R(N_f) \rightarrow U_V(1) \times SU_V(N_f)$$

($U_A(1)$ symmetry is broken explicitly by instantons)

Smoking gun for chiral symmetry breaking

- Non-zero chiral condensate in the instanton vacuum

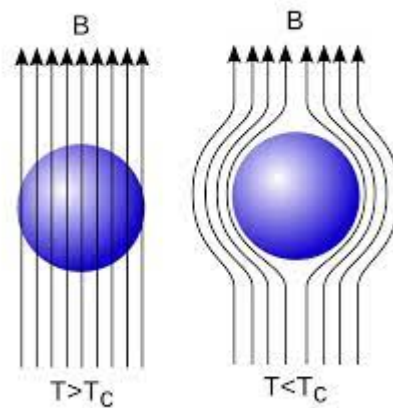
$$\langle 0 | \bar{q}_L q_R | 0 \rangle \neq 0 \quad (\text{density of instantons})$$

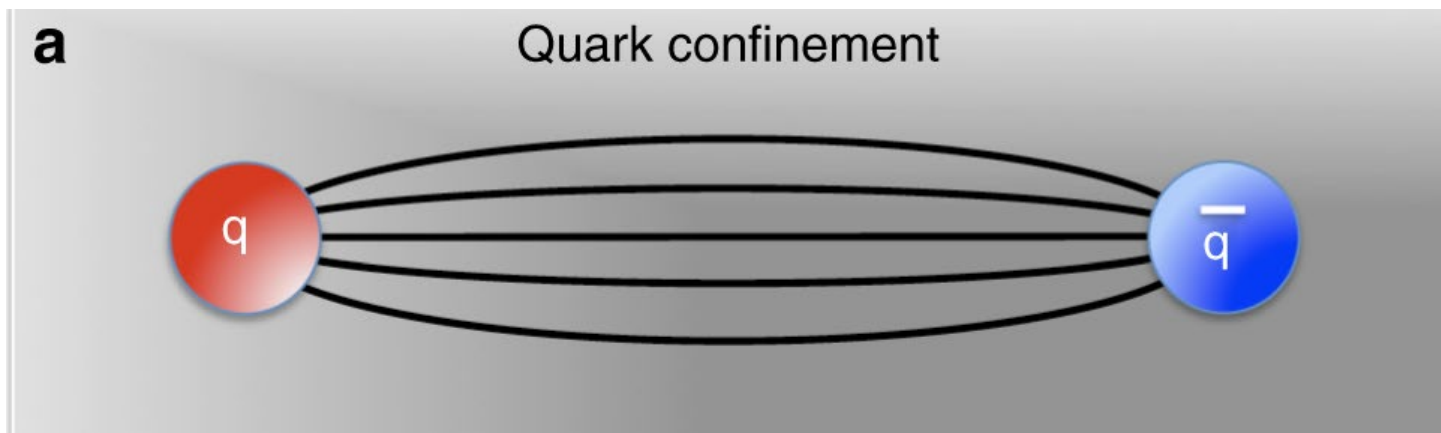
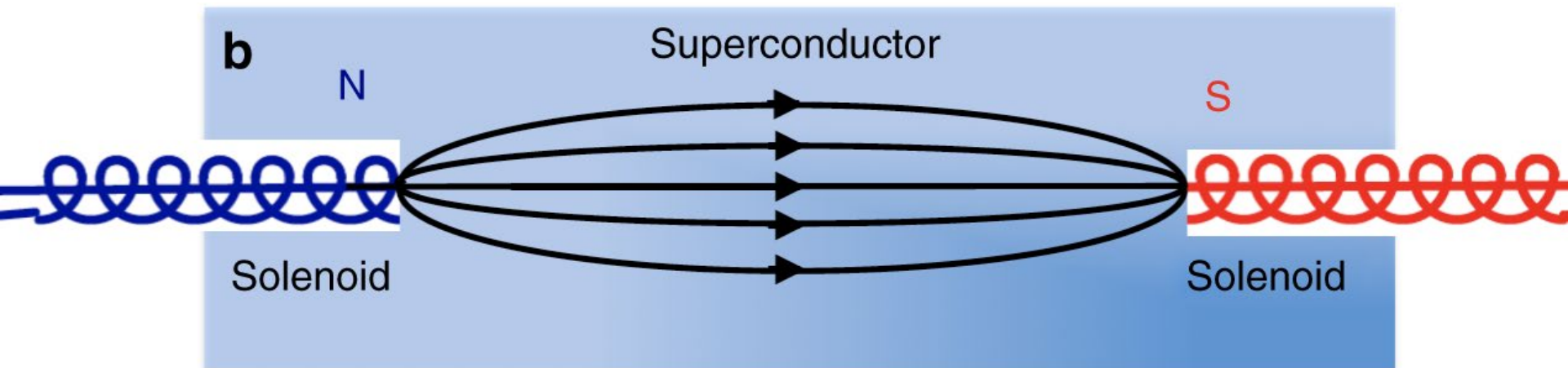
- Appearance of massless particles: Goldstone bosons as **vacuum excitations**.
- When quarks have small masses, we have pseudo-Goldstone bosons π, K which has mass-squared proportional to quark mass

$$m_\pi^2 \sim m_q$$

Color confinement

- An isolated colored charge has infinite energy in the QCD vacuum. Thus, colored states are not part of the QCD spectrum.
- The QCD vacuum expels the color electric field, just like the superconductor expels the magnetic field





Hadrons

- Is it possible to understand the structure of hadrons without understanding the QCD vacuum?
- Is it possible that there are effective degrees of freedom using which one can describe the hadrons?
 - Basis for model building (quark models, bag models)
 - QFT-> simple few body problem.
- **In hadron physics, a universal effective description of hadrons has not been found**
 - Existing ones are partially effective in limited domains.
 - **We are forced to start from scratch**

More on water-wave analogy

- We know the basic interactions between water molecules
 - but we don't know how the state of water is formed, or how to calculate the properties of water.
 - how are the wave excitations formed on the top of it?

Low-energy effective theory: Navier-Stokes equation

- To understand the waves, we just need to solve Navier-Stokes equation
- Turbulences?

Ab initio calculations

- The only first principles approach to solve QCD is lattice method, first proposed by Ken Wilson in 1974
- Formulated in Euclidean space-time, allowing calculating Euclidean correlation functions.
- The vacuum properties can be calculated.
- The certain properties of hadrons can be calculated through the Euclidean correlation functions.
- Methods of quantum fields, not particles!

Understanding proton structure in non-perturbative QCD

- Analytically solving QCD is a long shot (AdS/CFT is hardly a controlled approximation)
- Calculations using lattice QCD have made important progress, but is difficult to produce penetrating insight.
- What are the deep insights about the structure of the proton in nonperturbative QCD?
- This the most important reason for EIC!

To theorists:

- Don't tell us how good your predictions match with experimental data (a lot of them roughly do!)
- Tell us what are the systematic errors of your theory predictions or calculations in non-perturbative QCD (QFT)!
- QCD sum rules, Bethe-Salpeter equations, AdS/CFT, etc are not yet controlled approximations.