





Three-dimensional structure at the EIC (I)

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Outline

- Lecture 1 : Spatial distributions
- Lecture 2 : Parton distributions
- Lecture 3 : Wigner distributions
- Tutorial

Non-relativistic picture





3D structure is fundamental to understand physical properties

i.e. thermal, electrical, mechanical, ...



fullerene

nanotube

graphene

Example: X-ray diffraction



Diffraction pattern



 $\propto |A_{\rm scatt}|^2$

Scattered amplitude

$$\begin{split} A_{\rm scatt} \propto F(\vec{q}) &= \int {\rm d}^3 r \, e^{i \vec{q} \cdot \vec{r}} \, \rho(\vec{r}) \qquad \quad \vec{q} = \vec{k} - \vec{k}' \\ {}_{\rm Form \ factor} \qquad \qquad \begin{array}{c} {}_{\rm Scatterer} \\ {}_{\rm distribution} \end{array} \end{split}$$

Spatial distributions - Principles



Relativistic treatment

in Born approximation



$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} \Big/ \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} \Big|_{\mathrm{pointlike}} = \ [F(Q^2)]^2 \qquad {}^{\mathrm{Spin-0}}_{\mathrm{target}}$$

 $= \left\{ \left[G_E(Q^2) \right]^2 + \frac{\tau}{\epsilon} \left[G_M(Q^2) \right]^2 \right\} \frac{1}{1+\tau}$

Spin-1/2 target

Electric form factor Magnetic form factor

$$\begin{aligned} Q^2 &= -\Delta^2 \\ \tau &= Q^2/4M_N^2 \\ \epsilon &= (1+2(1+\tau)\tan^2\frac{\theta_e}{2})^{-1} \end{aligned}$$

Nucleon electromagnetic form factors

Proton

Neutron



[Alexandrou et al., PRD100 (2019) 014509]

Expectation value

$$\langle \psi | j^{\mu}(\vec{r}) | \psi \rangle = \int \frac{\mathrm{d}^{3}P}{(2\pi)^{3}} \frac{\mathrm{d}^{3}\Delta}{(2\pi)^{3}} \,\tilde{\psi}^{*}(\vec{P} + \frac{\vec{\Delta}}{2}) \tilde{\psi}(\vec{P} - \frac{\vec{\Delta}}{2}) \,\langle \vec{P} + \frac{\vec{\Delta}}{2} | j^{\mu}(\vec{r}) | \vec{P} - \frac{\vec{\Delta}}{2} \rangle$$

Normalization $\langle \vec{p}' | \vec{p} \rangle = (2\pi)^3 \, \delta^{(3)} (\vec{p}' - \vec{p})$

Wave packet $\tilde{\psi}($

 $ilde{\psi}(ec{p}) = \langle ec{p} \, | \psi
angle$

Probabilistic interpretation $\tilde{\psi}(\vec{P} \pm \frac{\vec{\Delta}}{2}) \approx \tilde{\psi}(\vec{P})$

$$\begin{split} \langle \psi | j^{\mu}(\vec{r}) | \psi \rangle \approx \int \frac{\mathrm{d}^{3}P}{(2\pi)^{3}} \, | \tilde{\psi}(\vec{P}) |^{2} \underbrace{\mathcal{J}^{\mu}_{\vec{P}}(\vec{r})}_{= \int \frac{\mathrm{d}^{3}\Delta}{(2\pi)^{3}} \, e^{-i\vec{\Delta}\cdot\vec{r}} \langle \vec{P} + \frac{\vec{\Delta}}{2} | j^{\mu}(\vec{0}) | \vec{P} - \frac{\vec{\Delta}}{2} \rangle \quad \text{Internal} \\ \text{distribution} \end{split}$$



 $\label{eq:Restframe} \begin{array}{ll} {\rm Rest\,frame} & |\vec{P}|=0 & \Rightarrow & P^0\approx M \end{array}$ Infinite momentum frame $\ |\vec{P}|\gg M & \Rightarrow & P^0\gg M \end{array}$

Phase-space approach

$$\langle \psi | j^{\mu}(\vec{r}) | \psi \rangle = \int \frac{\mathrm{d}^{3}P}{(2\pi)^{3}} \,\mathrm{d}^{3}R \,\rho_{\psi}(\vec{R},\vec{P}) \,\langle j^{\mu} \rangle_{\vec{R},\vec{P}}(\vec{r})$$

Nucleon Wigner
distribution
$$\rho_{\psi}(\vec{R},\vec{P}) = \int d^{3}z \, e^{-i\vec{P}\cdot\vec{z}} \, \psi^{*}(\vec{R}-\frac{\vec{z}}{2})\psi(\vec{R}+\frac{\vec{z}}{2}) \qquad \qquad \psi(\vec{r}) = \int \frac{d^{3}p}{(2\pi)^{3}} \, e^{-i\vec{p}\cdot\vec{r}} \, \tilde{\psi}(\vec{p}) \\ = \int \frac{d^{3}q}{(2\pi)^{3}} \, e^{-i\vec{q}\cdot\vec{R}} \, \tilde{\psi}^{*}(\vec{P}+\frac{\vec{q}}{2})\tilde{\psi}(\vec{P}-\frac{\vec{q}}{2})$$

Quasi-probabilistic interpretation

$$\int d^{3}R \,\rho_{\psi}(\vec{R},\vec{P}) = |\tilde{\psi}(\vec{P})|^{2}$$
$$\int \frac{d^{3}P}{(2\pi)^{3}} \,\rho_{\psi}(\vec{R},\vec{P}) = |\psi(\vec{R})|^{2}$$



Internal distribution (for a state localized in phase-space)

$$\langle j^{\mu} \rangle_{\vec{R},\vec{P}}(\vec{r}) = \int \frac{\mathrm{d}^{3}\Delta}{(2\pi)^{3}} e^{-i\vec{\Delta}\cdot(\vec{r}-\vec{R})} \langle \vec{P} + \frac{\vec{\Delta}}{2} | j^{\mu}(0) | \vec{P} - \frac{\vec{\Delta}}{2} \rangle$$

<u>NB</u>: $\mathcal{J}^{\mu}_{\vec{P}}(\vec{r}) = \langle j^{\mu} \rangle_{\vec{0},\vec{P}}(\vec{r})$



Nucleon charge distributions (3D Breit frame)

Proton

Nuclei



Nucleon charge distributions (2D elastic frame)

0.4

0.2

0.2



0.8

0.6

0.4



Neutron



Proton-pion fluctuation



Nucleon charge distributions (2D elastic frame)

0.2

0.4

0.8

0.6



-0.6

 $\sim G_M(Q^2)$

Instead of charge current, we can look at four-momentum current



Four-momentum operator

$$P^{\mu}_{a} = \int \mathrm{d}^{3}r \, T^{0\mu}_{a}(r) \qquad a = q, g$$

Energy distribution (3D Breit frame)

$$\langle T^{00} \rangle_{\vec{0},\vec{0}}(\vec{r}) = \int \frac{\mathrm{d}^3 \Delta}{(2\pi)^3} e^{-i\vec{\Delta}\cdot\vec{r}} \langle \frac{\vec{\Delta}}{2} | T^{00}(0) | - \frac{\vec{\Delta}}{2} \rangle$$



Pressure distributions (3D Breit frame)

Isotropic pressure

Pressure anisotropy



Mechanical equilibrium



von Laue relation





$$\gamma = \int \mathrm{d}r \, s(r)$$





Angular momentum distributions

Orbital vs intrinsic

$$L^{i}(\vec{r}) = \epsilon^{ijk} r^{j} \langle T^{0k} \rangle_{\vec{0},\vec{0}}(\vec{r})$$

$$S^{i}(\vec{r}) = \langle \overline{\psi} \gamma^{i} \gamma_{5} \psi \rangle_{\vec{0},\vec{0}}(\vec{r})$$



Kinetic vs Belinfante

 $J^i(\vec{r}) = L^i(\vec{r}) + S^i(\vec{r})$

$$J_{\text{Bel}}^{i}(\vec{r}) = \epsilon^{ijk} r^{j} \langle \frac{1}{2} (T^{0k} + T^{k0}) \rangle_{\vec{0},\vec{0}}(\vec{r})$$



Large-*N_c* bag model

Graviton exchange



Waaaaaayyy too weak in practice



Some references

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- Polyakov, Schweitzer, IJMPA33 (2018) 26, 1830025
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