

Polarisation Measurement and Spin Tracking at the ILC

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DESY

Beam Polarization and Polarimetry at the EIC, June 26, 2020

Introduction

Collision Data

Compton Polarimeters

Detector R&D

Spin Tracking

Conclusions

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Compton Polarimeters

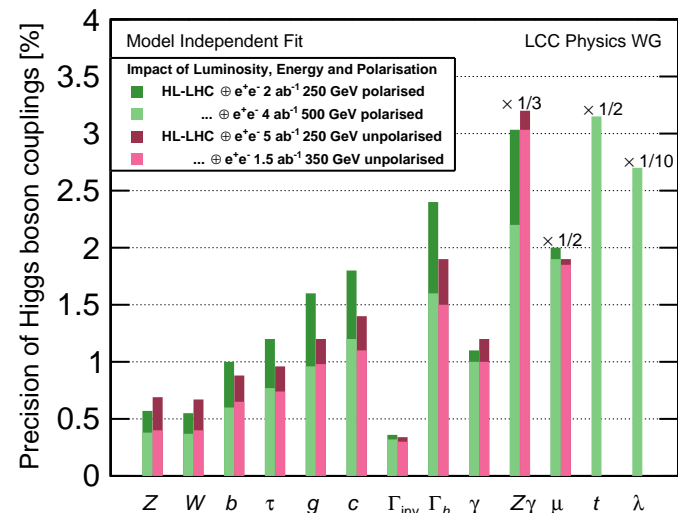
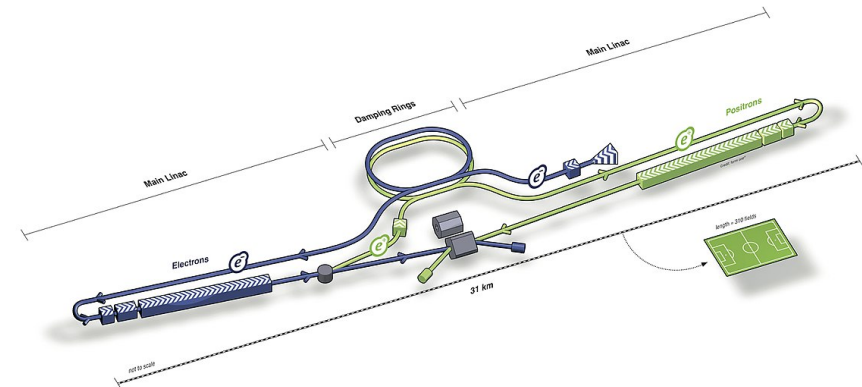
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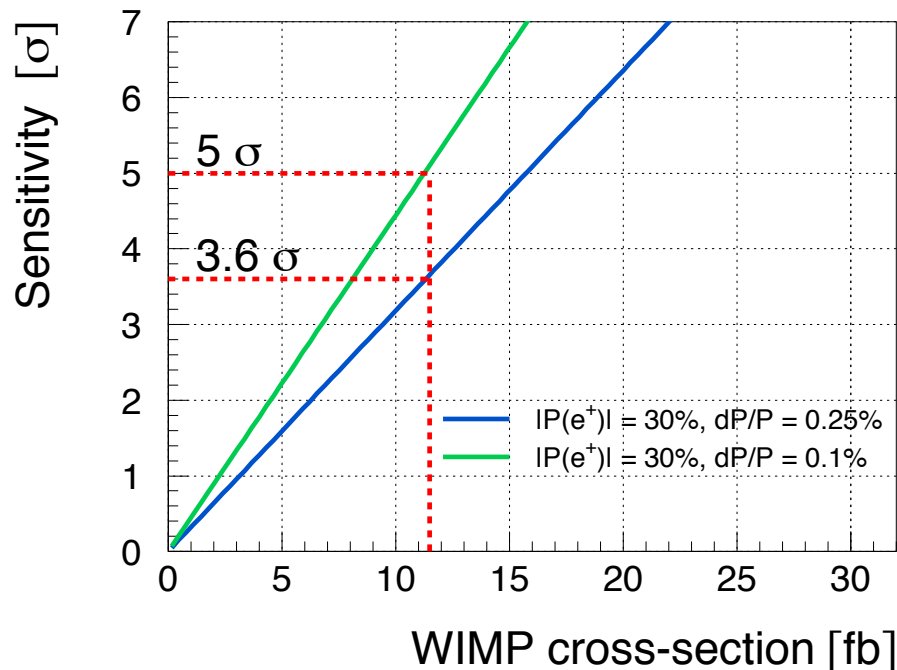
The International Linear Collider.

- $e^+ e^-$ “Higgs factory” with $\sqrt{s} = 250$ GeV, upgradable to up to 1 TeV
- *both* beams polarised:
 - $|P(e^-)| = 80\%$,
 - $|P(e^+)| = 30\% \dots 60\%$
- integral part of physics programme, for Higgs and beyond, c.f. <https://arxiv.org/abs/1801.02840>.
- construction under political consideration in Japan



Impact of Polarisation Uncertainty.

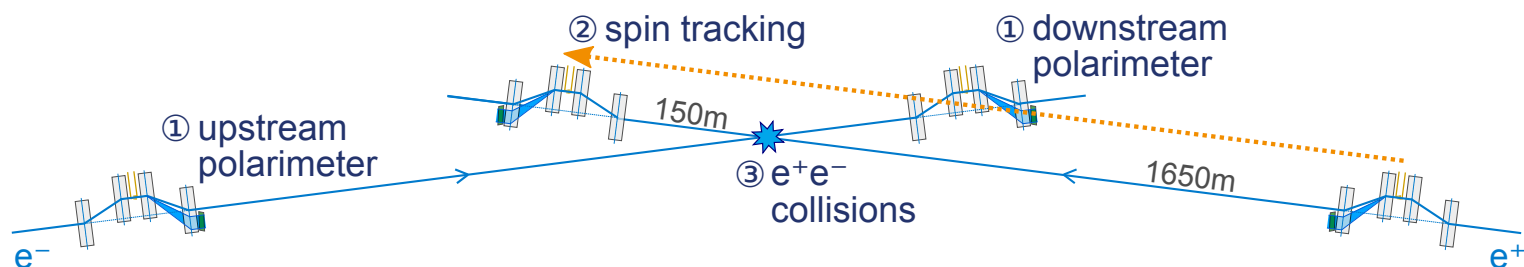
- SM precision measurements, eg. A_{LR} at Z pole will be limited by polarisation knowledge
→ simultaneous extraction of A_{LR} and $\langle P_{\text{eff}} \rangle_{IP}$
- BSM example: WIMP Dark Matter Search



- 500 fb^{-1} at 500 GeV,
 $|P(e^-, e^+)| = (0.8, 0.3)$
- ILD full simulation incl. systematics
- $dP/P = 0.25\%$
→ “evidence for”
- $dP/P = 0.1\%$
→ “discovery of”

Polarimetry concept for the ILC.

Goal for ILC polarimetry: per-mille level precision on luminosity weighted average polarisation at the IP, $\langle P_z \rangle_{IP} = \frac{\int P_z(t) \mathcal{L}(t) dt}{\int \mathcal{L}(t) dt}$



- ① Compton polarimeter measurements upstream and downstream of the e^+e^- interaction point
- ② Spin tracking to relate these measurements to the polarization at the e^+e^- interaction point
- ③ Long-term average determined from e^+e^- collision data as absolute scale calibration

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Polarisation Average from Collision Data.

Direct extraction from collision data

- any abundant, well-known, polarisation dependent process
- total cross-sections only (aka “(modified) Blondel scheme”):
 $P_+(e^-) = -P_-(e^-)$ and $P_+(e^+) = -P_-(e^+)$
- to lift condition: differential cross-sections *and* polarimeter constraints.

Methods studied so far

- total cross-sections: WW at 500 GeV and 1 TeV (ILD)
 single W etc at 3 TeV (CLIC)
- single-differential cross-sections: WW at 500 GeV & 1 TeV (ILD)
- double-differential cross-sections: WW at 1 TeV (SiD)
- **NEW:** global fit incl. differential cross-sections for all
 $e^+e^- \rightarrow f\bar{f}$ and $e^+e^- \rightarrow f\bar{f}'f''\bar{f}'''$ processes

Fast helicity reversal: .

... for both beams:

- collect data for all helicity configurations **simultaneously**
- ensures similar polarisation (absolute) values for all data sets
- enables cancellation of time dependent effects / systematic uncertainties for collider detector!

Counter example HERA:

- **slow** helicity reversal:
weeks between flips
- differences in $\langle P_e \rangle_{IP}$:
rely on polarimeters
- uncertainty $\sim 2\%$

Collisions	$P_e[\%]$	$\mathcal{L}[\text{pb}^{-1}]$
$e^+ p$	+32	98
$e^+ p$	-38	82
$e^- p$	+37	46
$e^- p$	-26	103

Phys. Lett. B704 (2011) 388 [arxiv:1107.3716] (H1 Leptoquarks)

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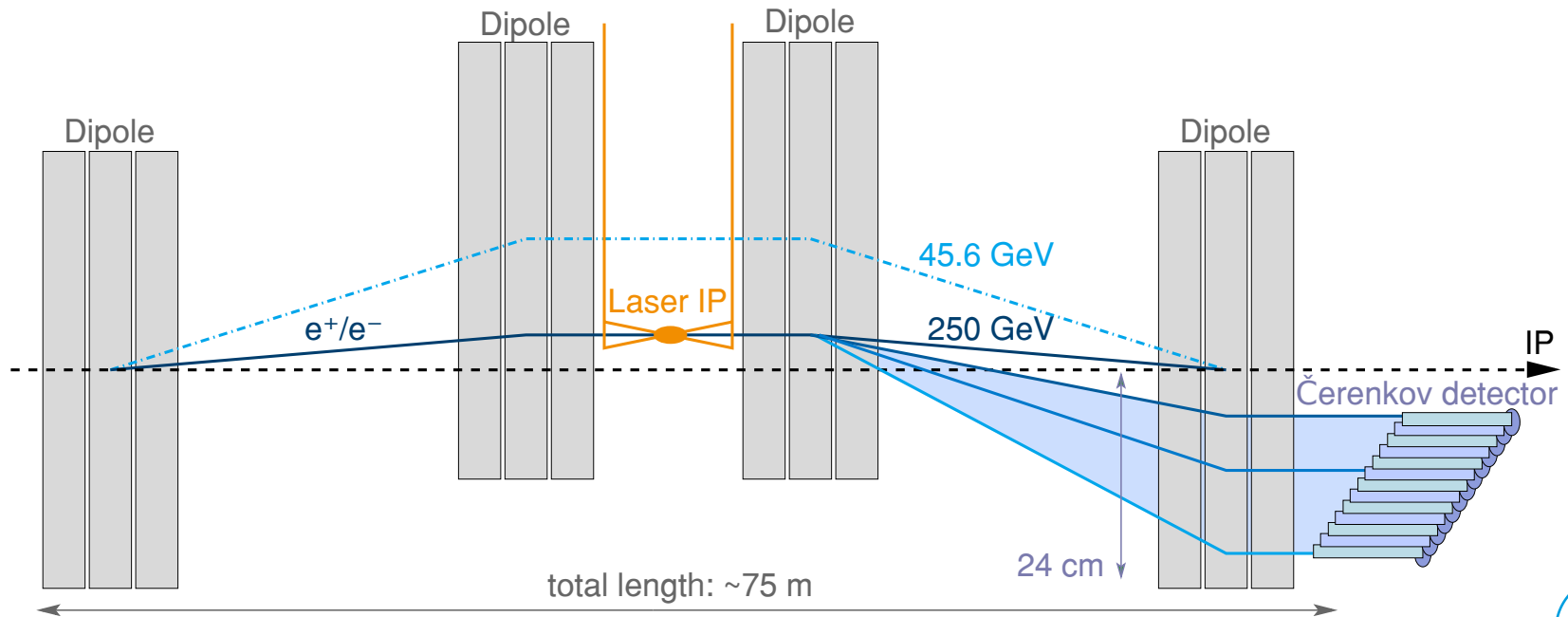
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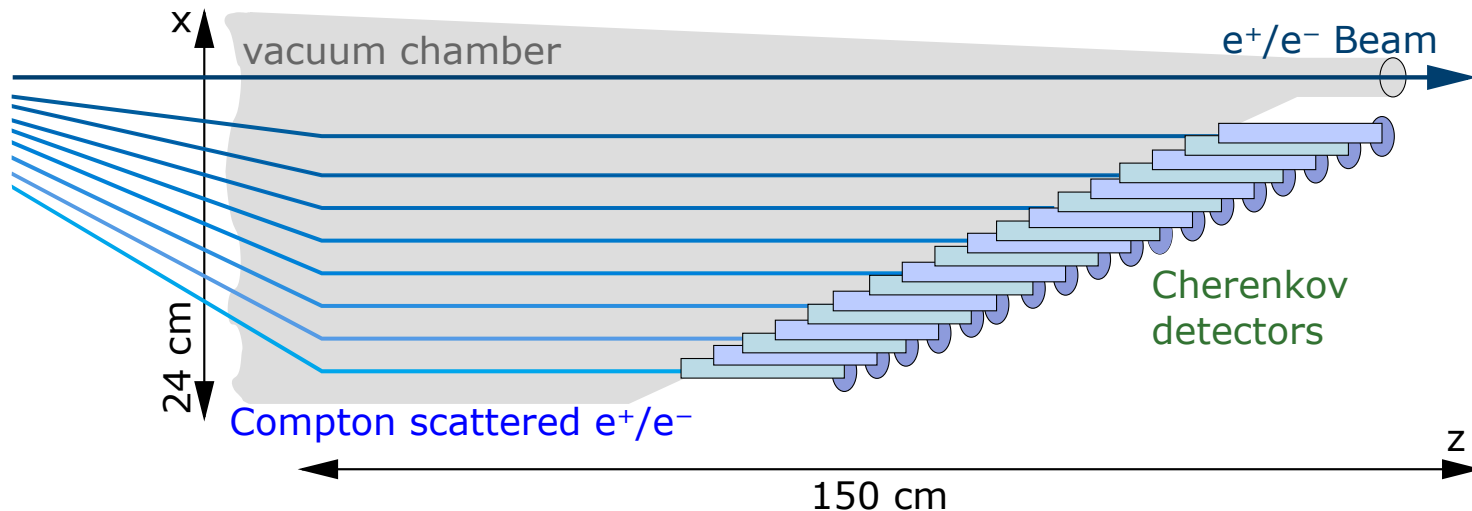
Compton polarimetry.

- fast measurement: $\mathcal{O}(10^3)$ Compton scatterings/bunch
- Energy spectrum of scattered e^+/e^- depends on product of lepton (\mathcal{P}) and laser (λ) polarisations
- Magnetic chicane: energy distribution \rightarrow position distribution
- Measure number of e^+/e^- per detector channel



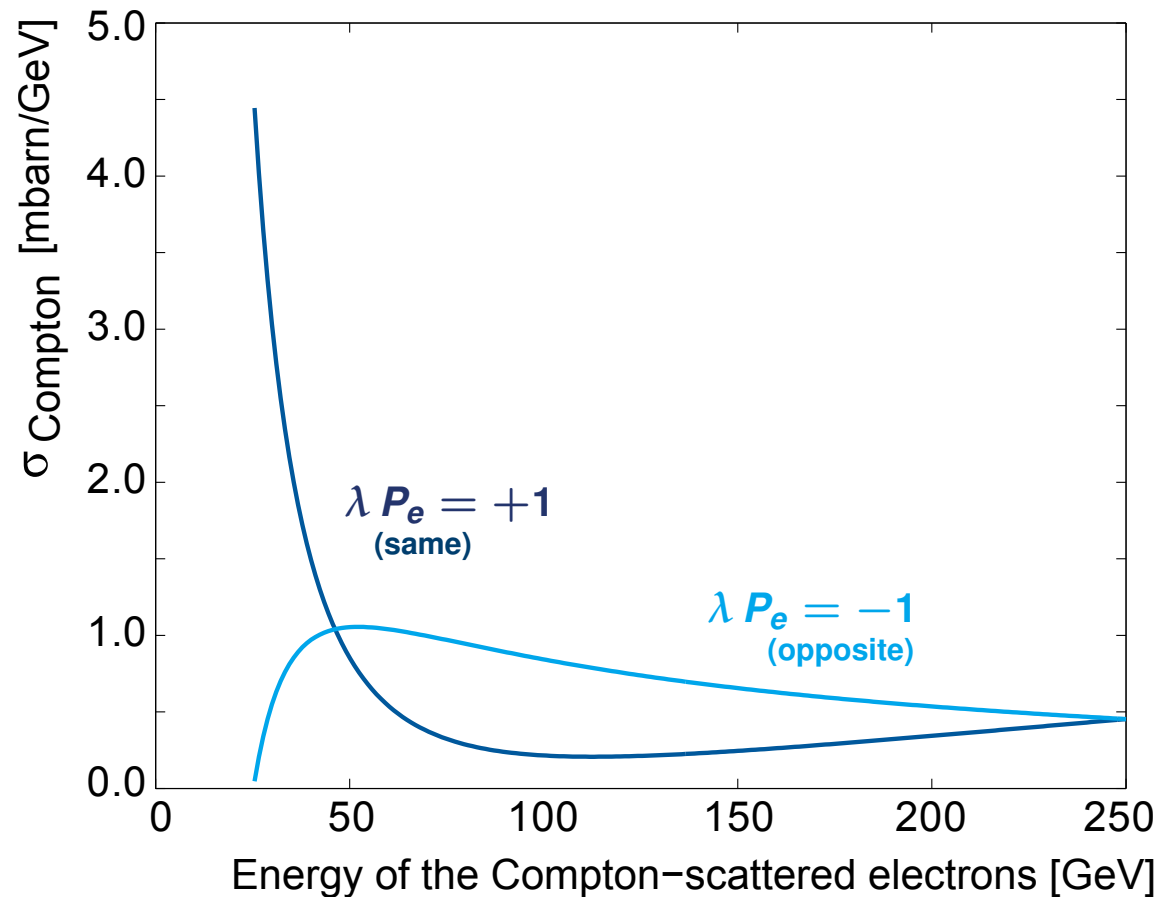
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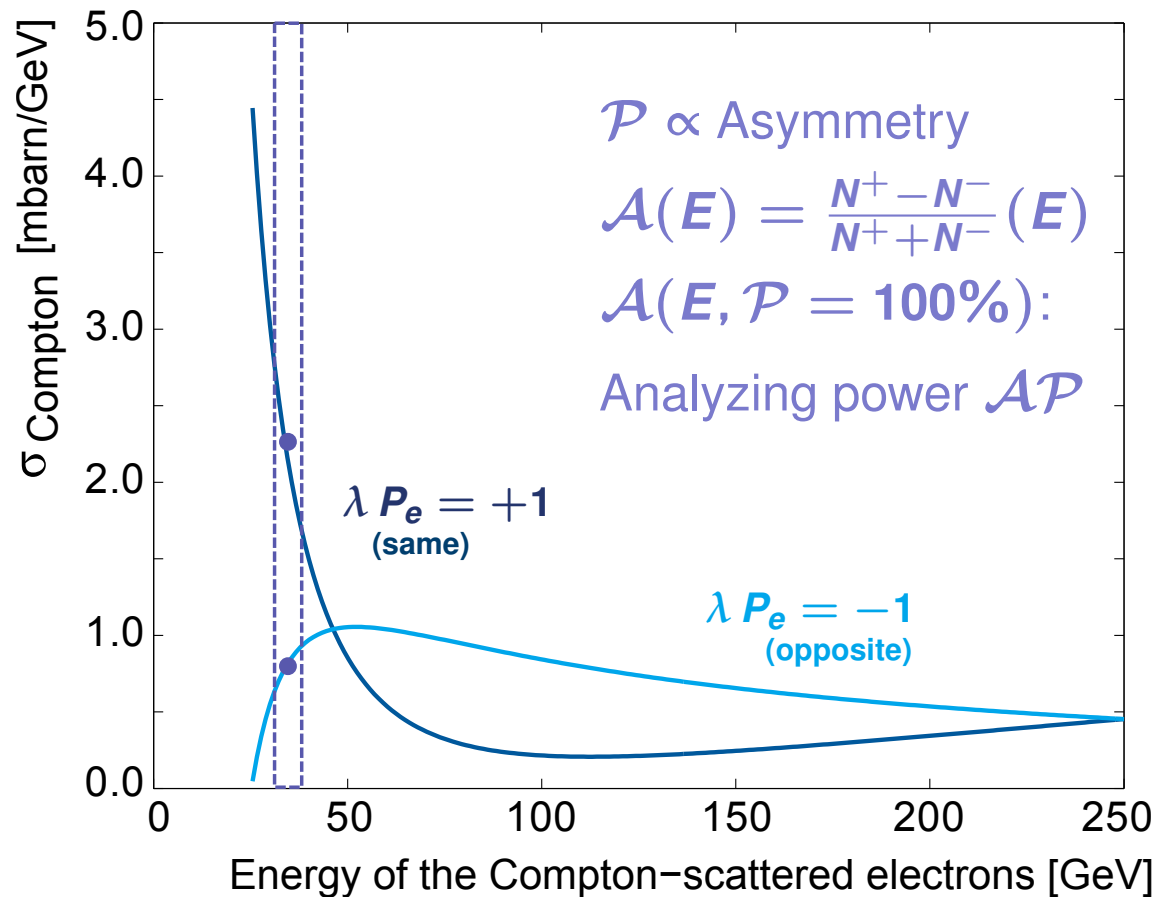
Measurement Principle.

Compton rate asymmetry is proportional to the beam polarisation:



Measurement Principle.

Compton rate asymmetry is proportional to the beam polarisation:



Measurement Precision.

ILC Goal: total uncertainty $\delta P / P \approx 0.25\%$ for $|P| \simeq 80\%$

source of uncertainty	$\delta P / P$	
	SLC achieved	ILC goals
laser polarisation	0.1%	0.1%
analyzing power	0.4%	0.15% – 0.2%
detector linearity	0.2%	0.1%
electronic noise and beam jitter	0.2%	0.05%
Total	0.5%	0.25%

- **analysing power:** prediction of count rate asymmetry per detector channel
 ⇒ knowledge of beam parameters, design of chicane, beam-detector alignment, backgrounds
- **detector linearity, electronic noise:**
 ⇒ detector design & calibration
- **beam jitter:** much smaller at ILC due to luminosity requirements

Complementarity of Up- and Downstream.

Upstream Polarimeter

- 1.8 km upstream of IP
- rather clean environment
- begin beam cond.
- samples every bunch
- stat. error 1% after few μs
- reference for control of collision effects

Downstream Polarimeter

- 140 m downstream of IP
- high backgrounds
- disrupted beam
- samples one bunch / train
- stat. error 1% after $\simeq 1$ min
- access to depolarisation in collision

Combination

- without collisions: spin transport in Beam Delivery System and Extraction Line
- with collisions: depolarisation at IP
- **cross check each other!**

[c.f. "Spin Dance" Exp., Phys. Rev. ST Accel. Beams **7** 042802 (2004)]

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Detector Requirements.

Magnetic Chicane...

- transforms energy spectrum into spatial distribution
- behind chicane: ~ 20 cm wide
- detect Compton electrons over this area

Detector requirements:

- Total ionising dose up to 100 Mrad / year
- read out signals of 1000-2000 Compton electrons (25-250 GeV) **every** bunch crossing
- either very linear response or “counting” electrons
- alignment to ~ 100 μm and ~ 1 mrad
- suppression of background from low energetic particles

Detector Options.

Simple, robust, fast: Cherenkov detectors

- Cherenkov light emission proportional to number of electrons
- independent of electron energy (once relativistic)
- successfully used in best polarimeter so far at SLC
- gas or quartz option for Cherenkov medium

Detector Options.

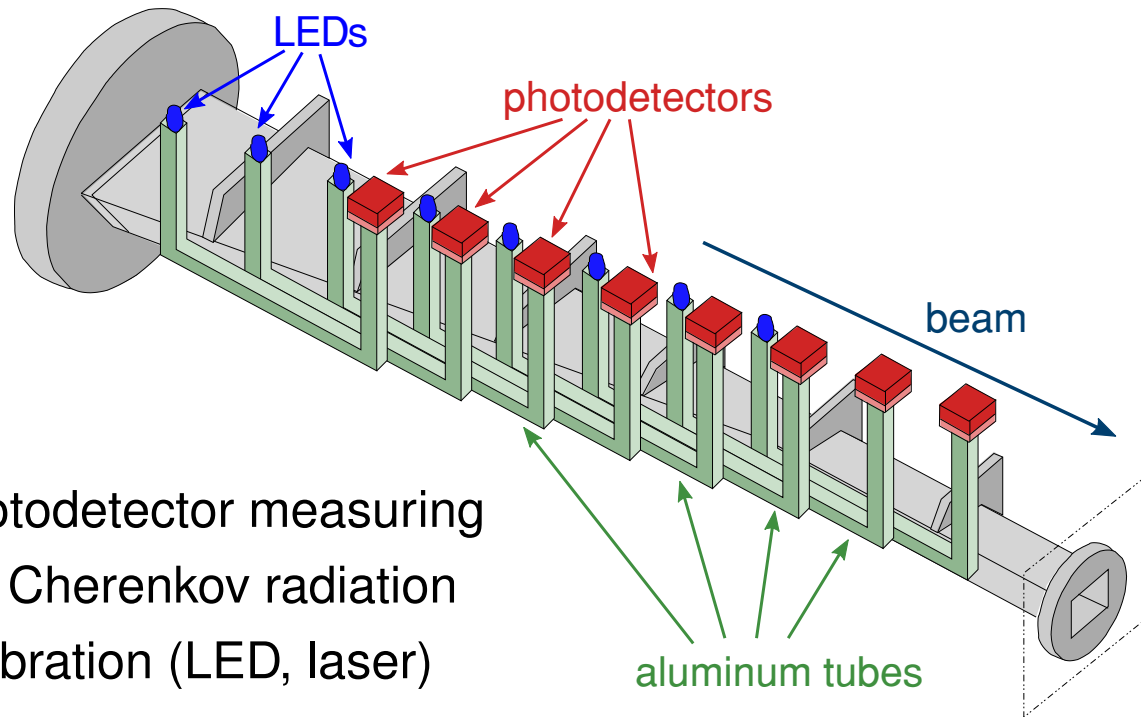
Simple, robust, fast: Cherenkov detectors

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Goal: total uncertainty $\Delta P / P \approx 0.25 \%$, of which

- laser: 0.1 %
- analysing power (i.e. asymmetry at $\mathcal{P} = 1$): 0.2 %
⇒ e.g. alignment
- detector linearity: 0.1 % ⇒ photodetector calibration

Gas Cherenkov detector.



- **hind U-leg:** photodetector measuring the Cherenkov radiation
- **front U-leg:** calibration (LED, laser)

Alignment: locate Compton edge in the spectrometer

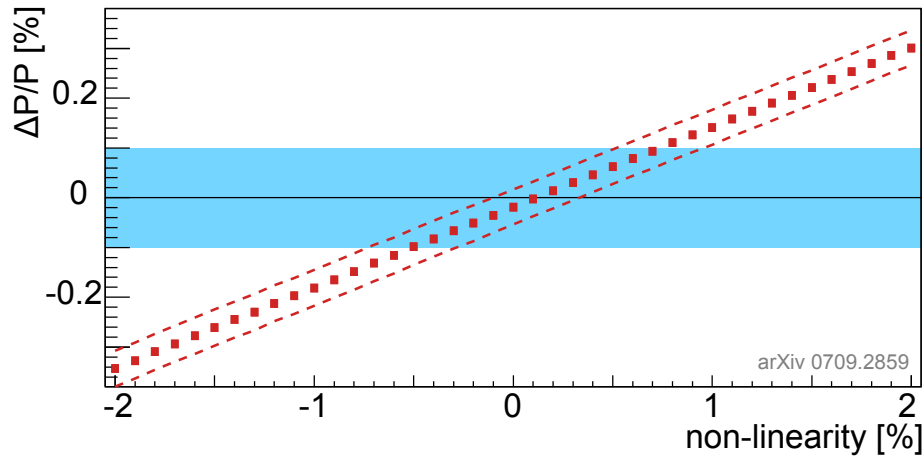
Segmented photodetectors: Tilt alignments via asymmetries

2-channel prototype tested at ELSA [JINST 7, P01019 (2012)]

⇒ tilt alignment of 0.1° , nearly fulfils alignment requirements

Calibration of detector non-linearity.

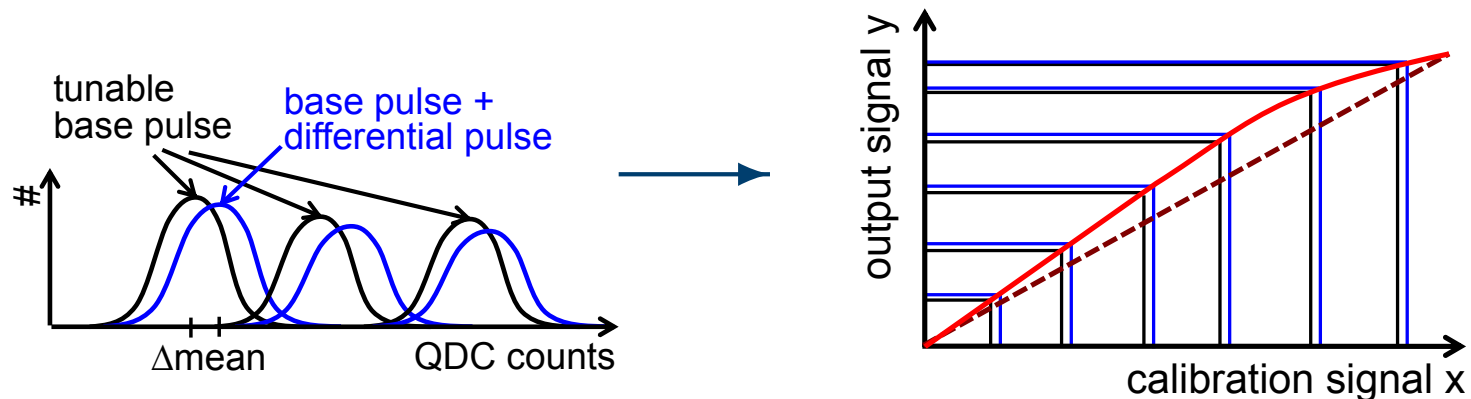
Goal: contribution to overall uncertainty $< 0.1\%$



PMTs have to be calibrated to non-linearity $< 0.5\%$.

$\mathcal{P} \propto \frac{N^+ - N^-}{N^+ + N^-}$: no absolute calibration needed.

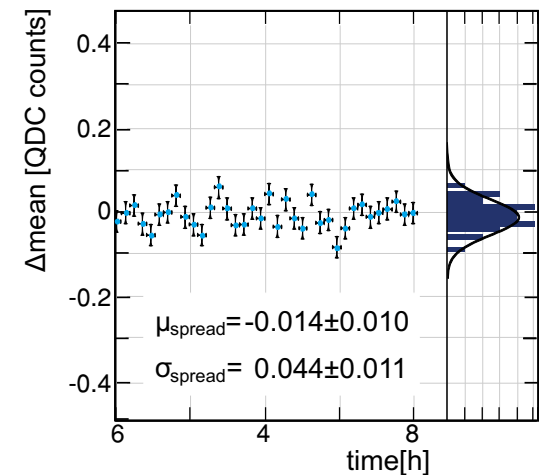
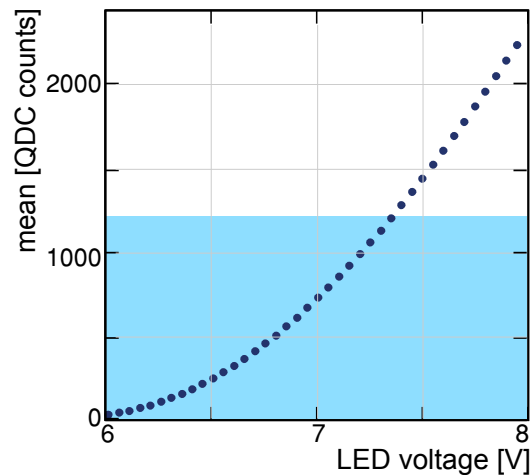
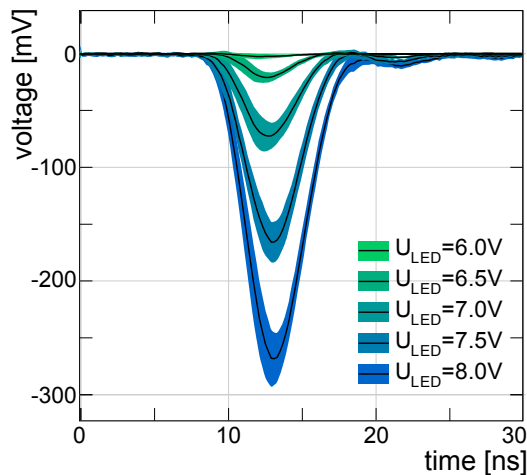
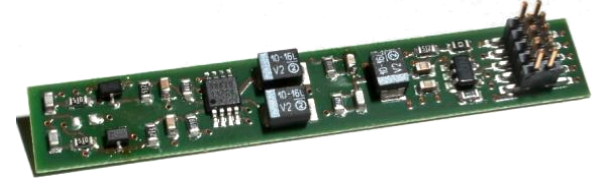
→ Differential calibration method using two LEDs:



Calibration source requirements.

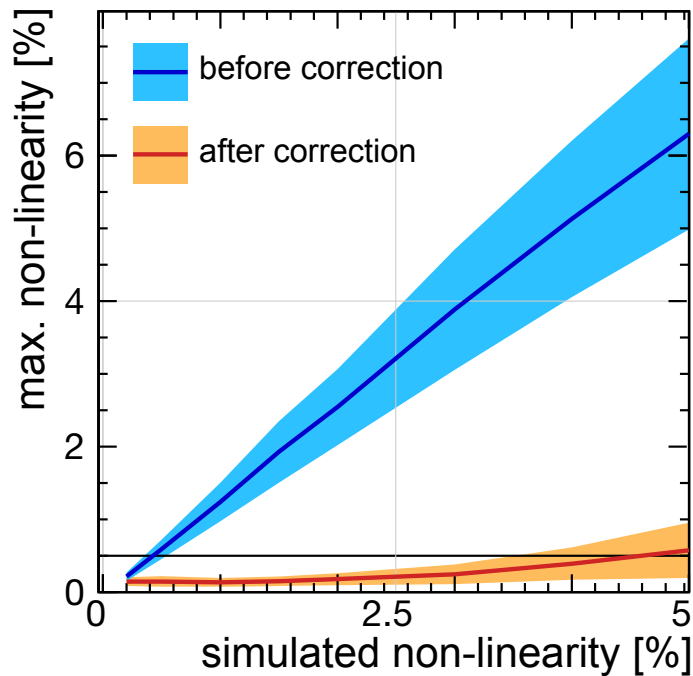
Requirements on the LED driver:

- wave length in UV range ($\lambda = 405 \text{ nm}$)
- applicable in detector design \rightarrow small
- short light pulses ($< 10 \text{ ns}$)
- coverage of the whole dynamic range of the expected signal
- reproducible and stable light pulses



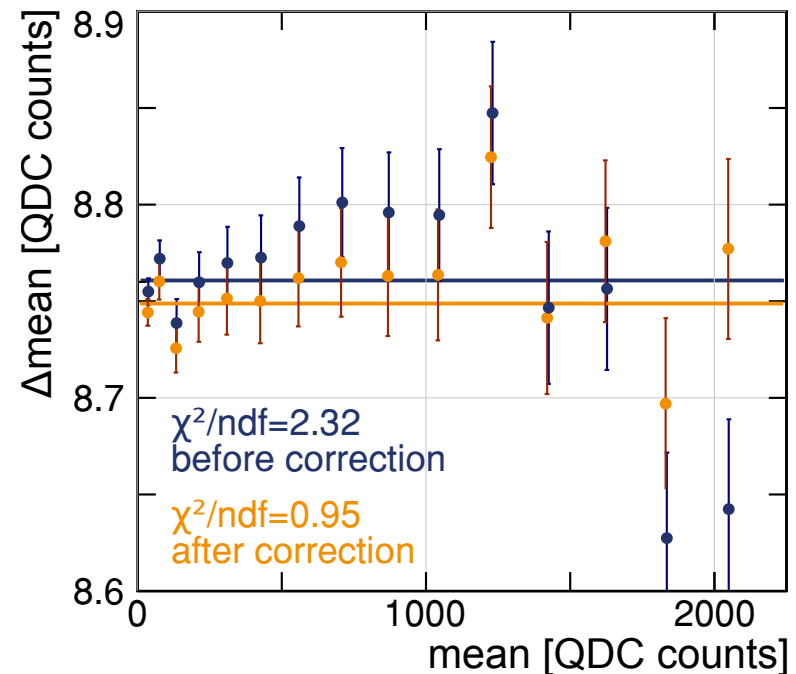
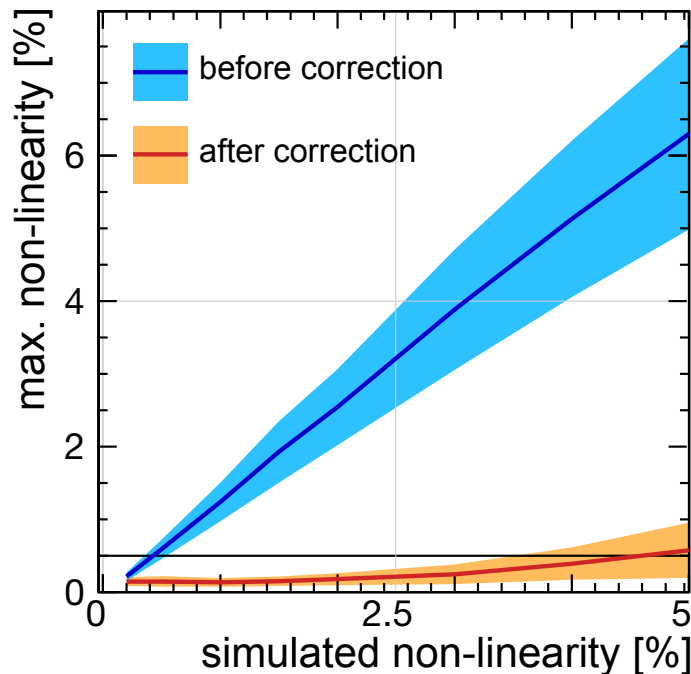
Test of non-linearity correction.

Simulations: Corrections of non-linearities up to 4 % possible.



Test of non-linearity correction.

Simulations: Corrections of non-linearities up to 4 % possible.
Applied method to on of the photodetectors used in testbeam:



⇒ Reached non-linearity $< 0.2\%$ in the expected dynamic range,
in single polarimeter channels even smaller.

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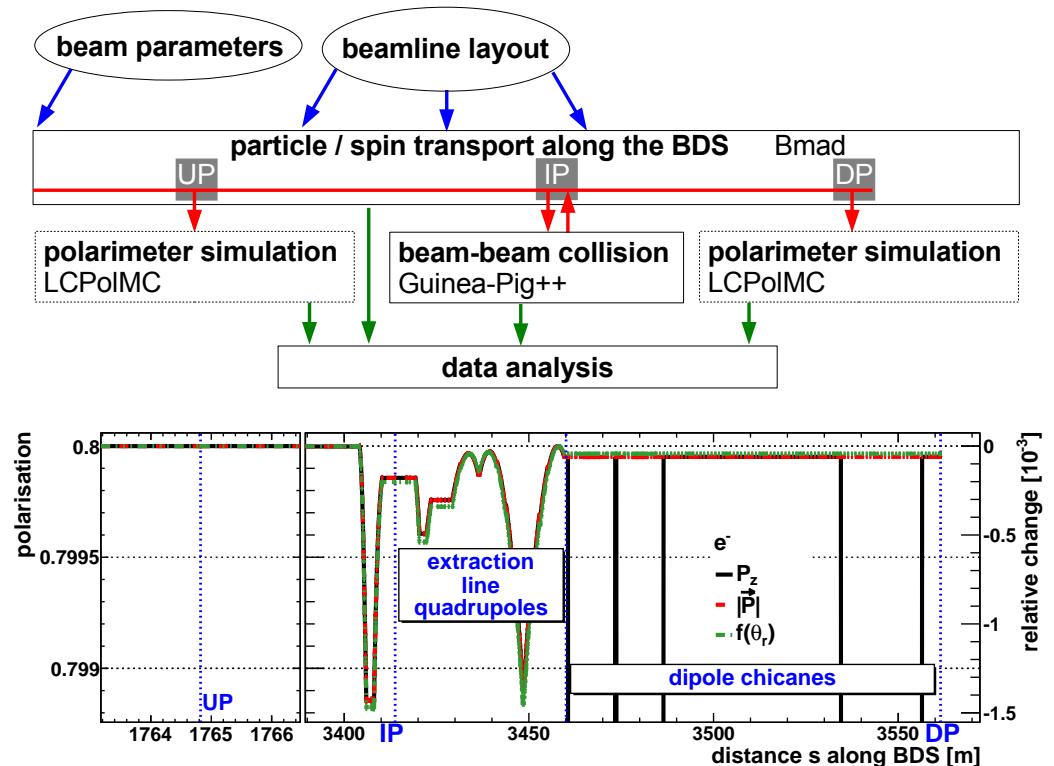
Conclusions

Spin Tracking in BDS and Extraction Line.

Based on SB2009-Nov10
lattice (PhD Thesis
M.Beckmann)

- developed simulation framework STaLC
- without collisions
⇒ cross-calibration of polarimeters
- with collisions
⇒ what does the downstream polarimeter measure?

Spin Transport At Linear Colliders



Cross-calibration of Polarimeters.

Without Collisions:

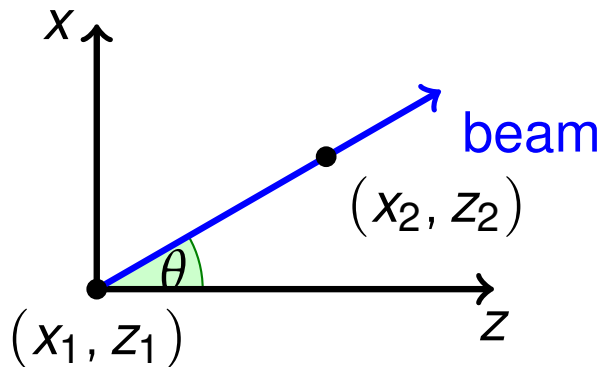
predict value at downstream location from upstream measurement

	effect on $P [10^{-3}]$
Beam and detector alignment at polarimeters ($\Delta\theta_{bunch} = 50 \mu\text{rad}$, $\Delta\theta_{pol} = 25 \mu\text{rad}$)	0.72
Variation in emittances	0.03
Crabbing	< 0.01
Detector magnets	0.01
Emission of synchrotron rad.	0.005
random misalignments (10 μm)	0.35
Total	0.80

Beam alignment.

Precision of Polarimeter Cross-Calibration

- is dominated by relative angle between beam directions at UP/DP $\Delta\theta_{bunch}$
- need *pairs* of BPMs at UP, IP and DP \Rightarrow “local” angle
- $\Delta x = 7 \mu\text{m}$, distance along $z =$ a few meters $\Rightarrow 1 - 2 \mu\text{rad}$
- challenge: absolute reference over $\simeq 2 \text{ km}$
 \Rightarrow is $50 \mu\text{rad}$ a realistic number?



$$\theta \approx \frac{x_2 - x_1}{L}; \quad L := z_2 - z_1; \quad L = 8\text{m}$$

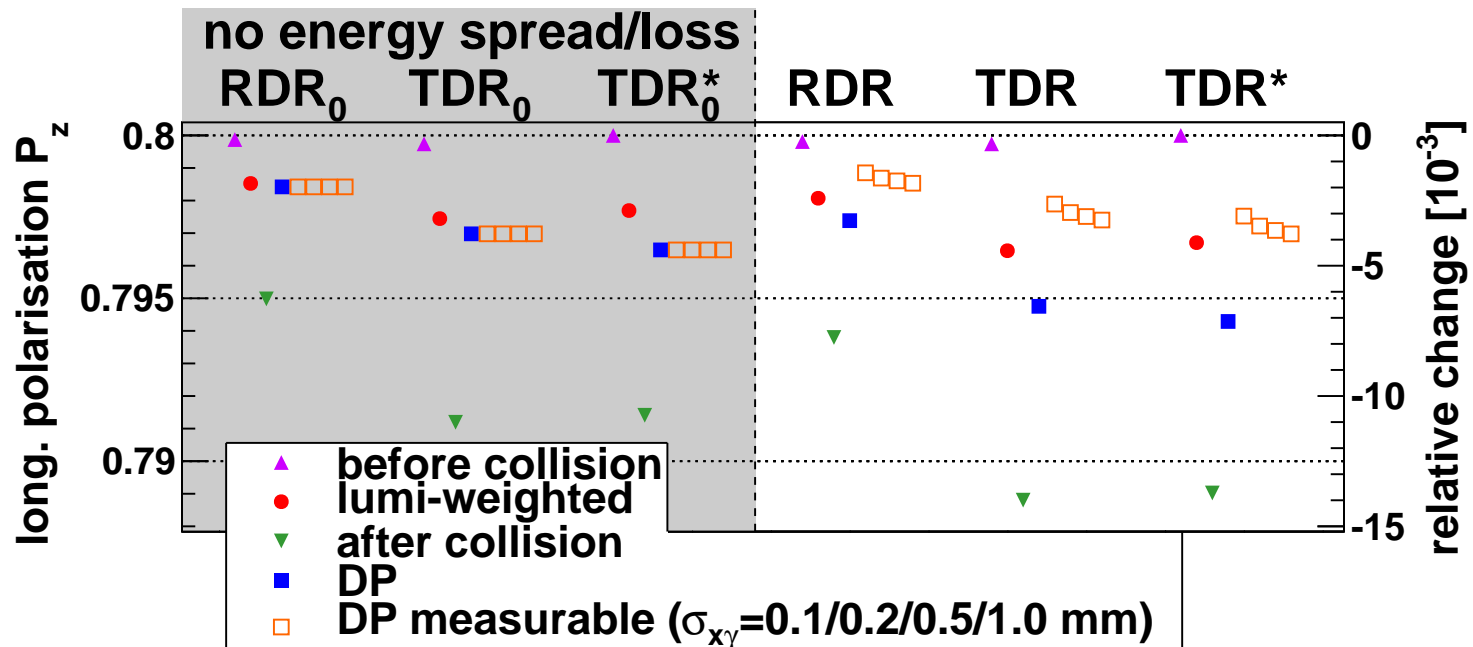
$$\Delta x_i = \Delta y_i = 7 \mu\text{m}; \quad \Delta L \simeq 0$$

$$\Rightarrow \Delta\theta^2 \leq \underbrace{2 \left(\frac{\Delta x_i}{L} \right)^2 + 2 \left(\frac{\Delta y_i}{L} \right)^2}_{:= (\Delta\theta_{BPM})^2} + \underbrace{\left(\frac{\theta \Delta L}{L} \right)^2}_{:= (\Delta\theta_L)^2}$$

$$\Rightarrow \Delta\theta \simeq 1.7 \mu\text{rad}$$

Collision Effects & Downstream Polarimeter.

- Extraction line optics designed to retrieve $\langle P \rangle_{IP}$ at downstream polarimeter
- Confirmed by STaLC w/o beamstrahlung (energy loss, grey background)
- **With beamstrahlung**: few permille difference to $\langle P \rangle_{IP}$
- **Measured polarisation depends on laser spot size** (here: perfect centering!)
- **Effect doubles from RDR \rightarrow TDR parameters**



Conclusions Spin Tracking.

In the presence of significant Beamstrahlung:

- Downstream polarimeter does *not* measure directly any more $\langle P \rangle_{IP}$
- difference DP measurement vs $\langle P \rangle_{IP}$ depends on
 - laser spot size & position
 - luminosity (\simeq energy loss in collision)
- \Rightarrow correcting for this requires
 - absolute reference from upstream polarimeter
 - luminosity & beam parameter monitoring
 - long-term scale of $\langle P \rangle_{IP}$ from collision data

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Conclusions.

Per mille-level precision on lumi-weighted average polarisation at IP required by physics, needs combination of

- long-term scale calibration from $e^+ e^-$ collision data (fast helicity reversal!)
- upstream (UP) and downstream (DP) polarimeters
 - **UP**: time resolution, absolute reference
 - **DP**: collision effects
 - **combined**: cross-calibration, lumi-weighted polarisation @ IP
- spin tracking and understanding of collision effects

Compton Polarimeters:

- beam-detector alignment & detector linearity crucial
- R&D well underway, requirements \simeq reached in prototypes
- cross-calibration without collisions: $\sim 0.1\%$ from alignment
 - esp. orbit and spin at UP and DP locations (2 km apart)

More Details.

➤ E&P workshop Zeuthen 2008

<http://indico.desy.de/conferenceDisplay.py?confId=585>

➤ its Executive Summary

arXiv:0903.2959 [physics.acc-ph]

➤ downstream polarimeter 6-magnet chicane

<http://www.slac.stanford.edu/cgi-wrap/getdoc/slac-pub-12425.pdf>

➤ publication on beam energy and polarisation measurements

JINST 4 (2009) P10015, arXiv:0904.0122 [physics.ins-det]

➤ publications on polarimeter detector R&D:

JINST 7 (2012) P01019, arXiv:1011.6314 [physics.ins-det]

arXiv:1502.06955 [physics.ins-det], DESY 15-028

➤ publication on BDS spin tracking

JINST 9 (2014) P07003, arXiv:1405.2156 [physics.acc-ph]

Backup Slides .

For historic information: Recommendations to GDE and Research Director (2008).

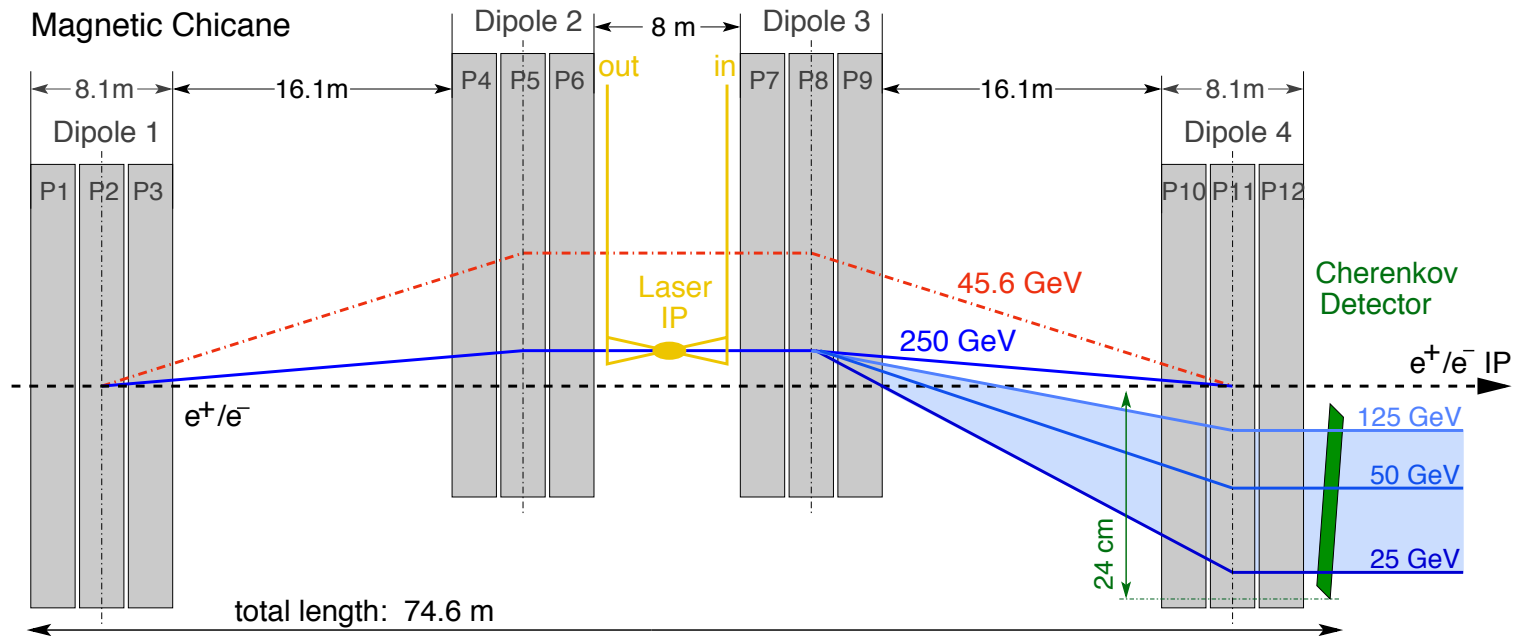
- Separate the functions of the upstream polarimeter chicane. Do not include an MPS energy collimator or laser-wire emittance diagnostics; use instead a separate setup for these two.
- Modify the extraction line polarimeter chicane from a 4-magnet chicane to a 6-magnet chicane to allow the Compton electrons to be deflected further from the disrupted beam line.
- Include precise polarisation and beam energy measurements for **Z**-pole calibration runs into the baseline configuration.
- Keep an initial positron polarisation of 30-45% for physics, don't reduce to 22% .
- Implement parallel spin rotator beamlines with a kicker system before the damping ring to provide rapid helicity flipping of the positron spin.
- Move the pre-DR positron spin rotator system from 5 GeV to 400 MeV. This eliminates expensive superconducting magnets and reduces costs.
- Move the pre-DR electron spin rotator system to the source area. This eliminates expensive superconducting magnets and reduces costs.



Design of the Upstream Polarimeter Chicane.

Why a 4-Dipole-Chicane?

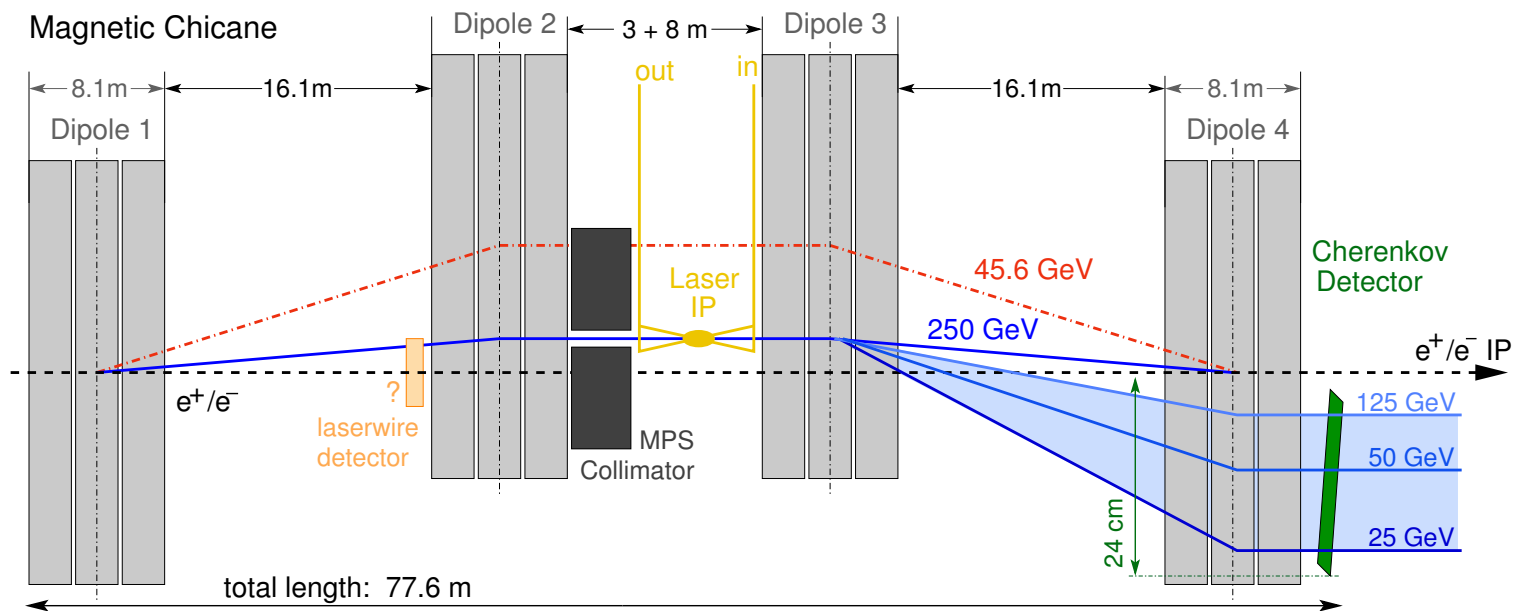
- Compton spectrum position at detector independent of E_{beam} if B -field constant
- price to pay: Compton IP moves laterally with E_{beam}



Design of the Upstream Polarimeter Chicane.

RDR Design:

- energy collimation and emittance diagnostics in same chicane
- \Rightarrow laterally fixed Compton IP \Rightarrow scaled field operation!
- collimator & laser-wire will create severe backgrounds

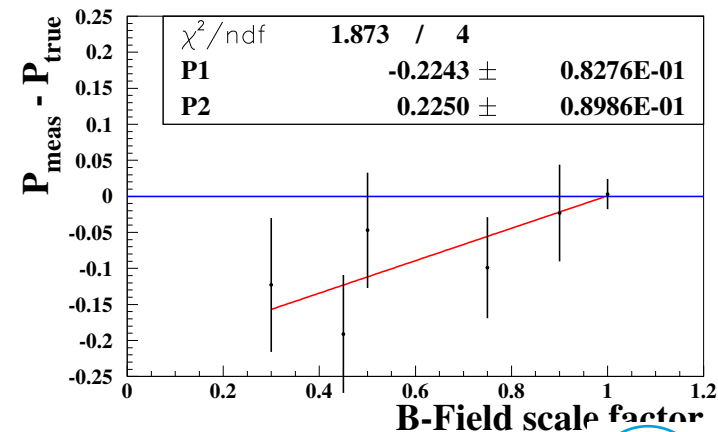
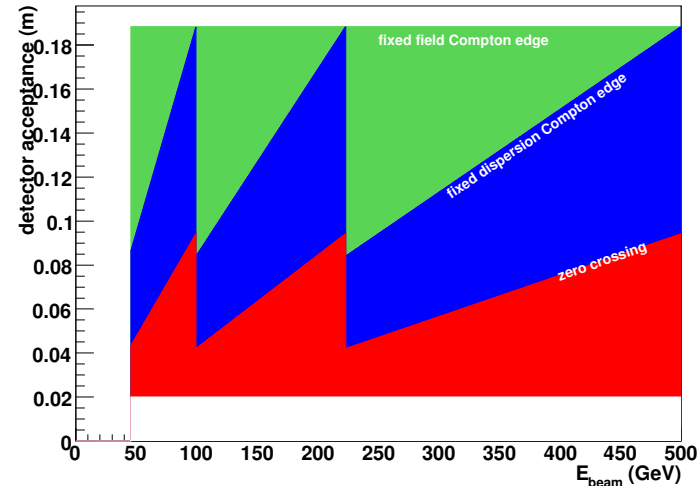


Scaled vs Fixed Field Operation.

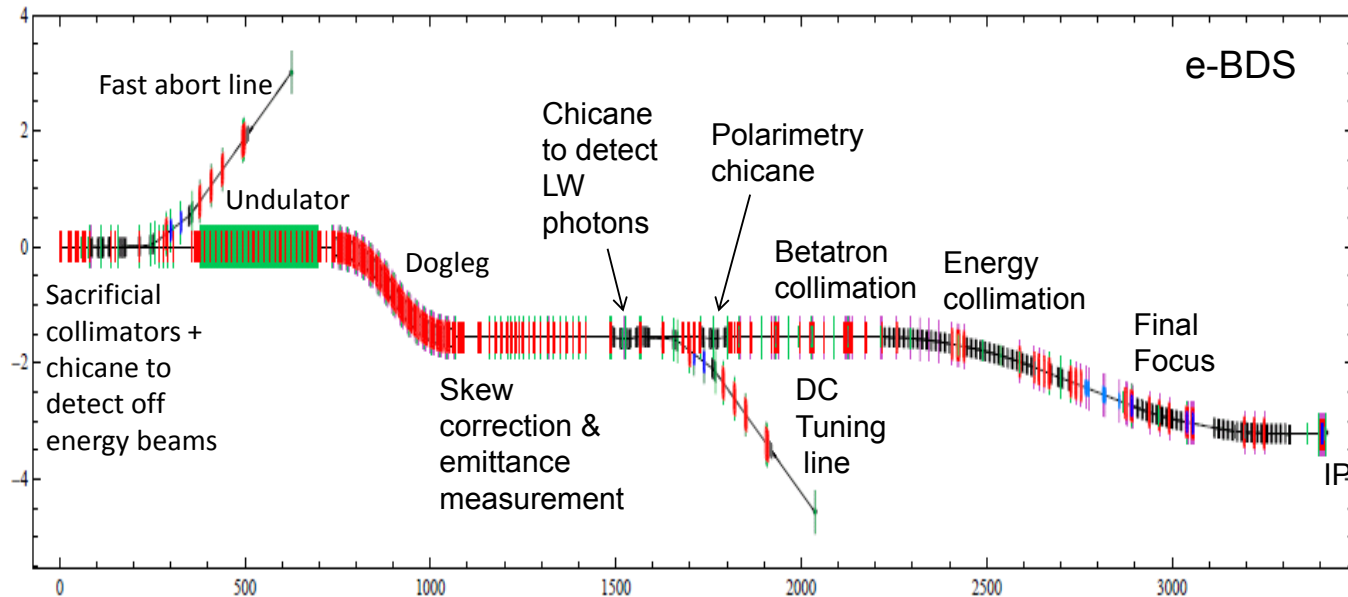
Effects of scaled field on measurement:

- acceptance: varies with E_{beam}
 ⇒ inhomogeneous quality of polarisation measurement
- alignment: via Compton edge position w.r.t. main beam
 ⇒ effect on $\delta\mathcal{P}/\mathcal{P}$ doubles
- systematic deviations for large scale factors

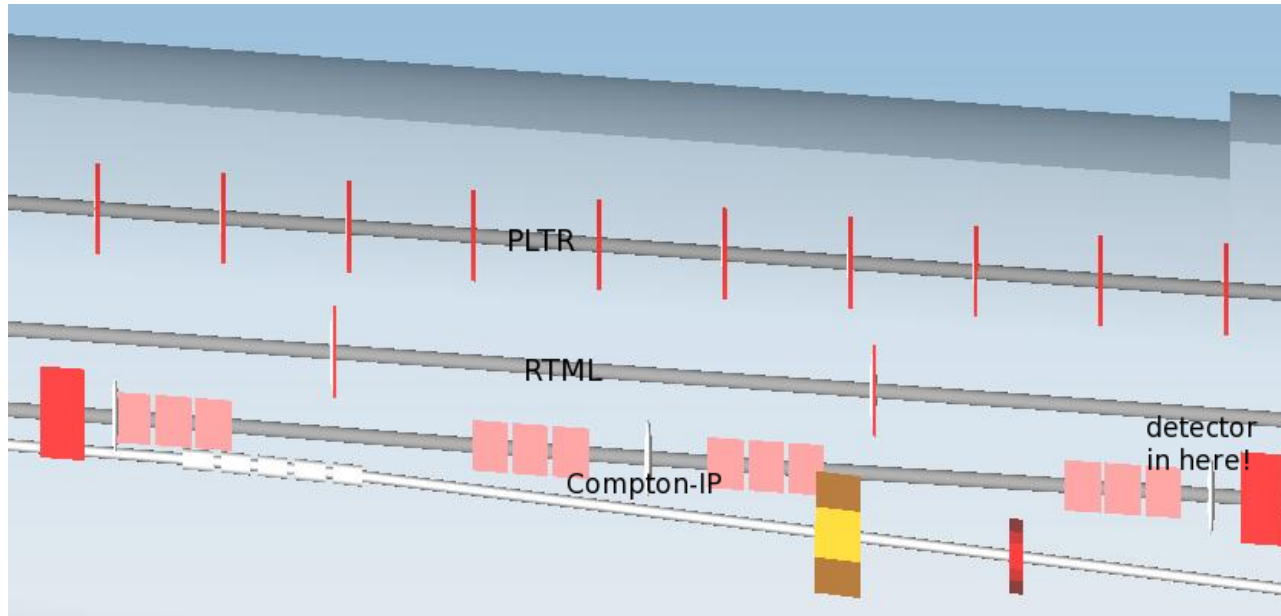
not compatible with extreme precision requirements c.f. ILC-NOTE-2008-047



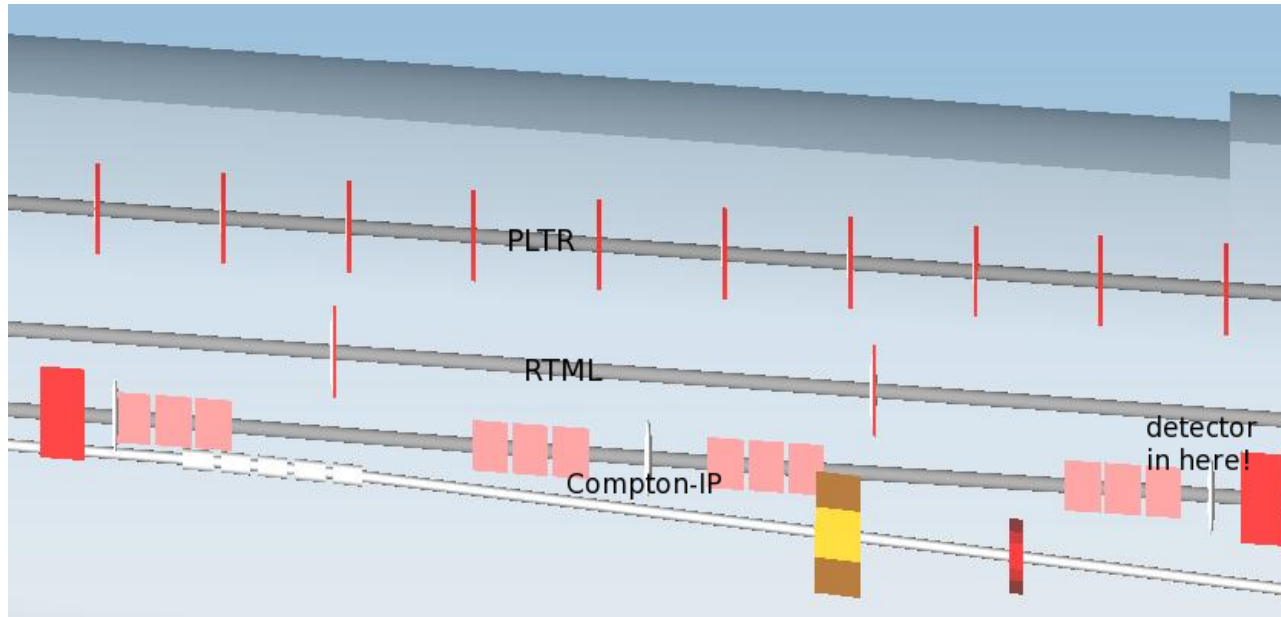
The Upstream Polarimeter in SB2009-Nov10 lattice.



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The Upstream Polarimeter in SB2009-Nov10 lattice.



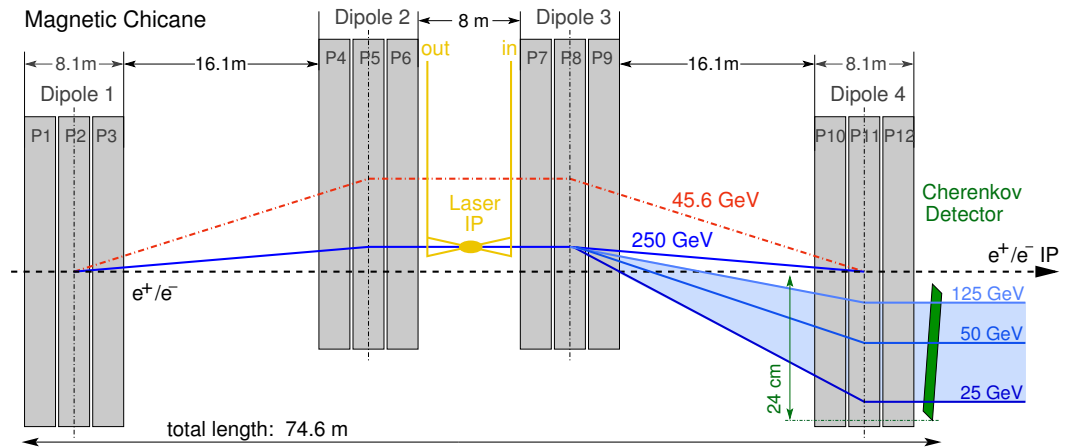
- distance Compton-IP to dump line ca 30 cm at 250 GeV
- down to ca 20 cm at lowest energies - enough?

[c.f. Baseline Technical Review Workshop 2011]

Vacuum Chamber in Chicane Region.

need special beam pipe through out whole chicane

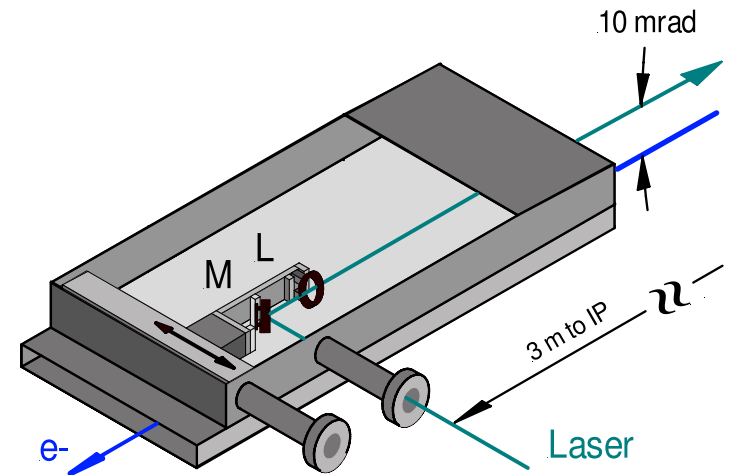
- to allow for varying bending angle
- to guide laser in and out
- to let fan of Compton scattered electrons pass
- to extract Compton fan to detector



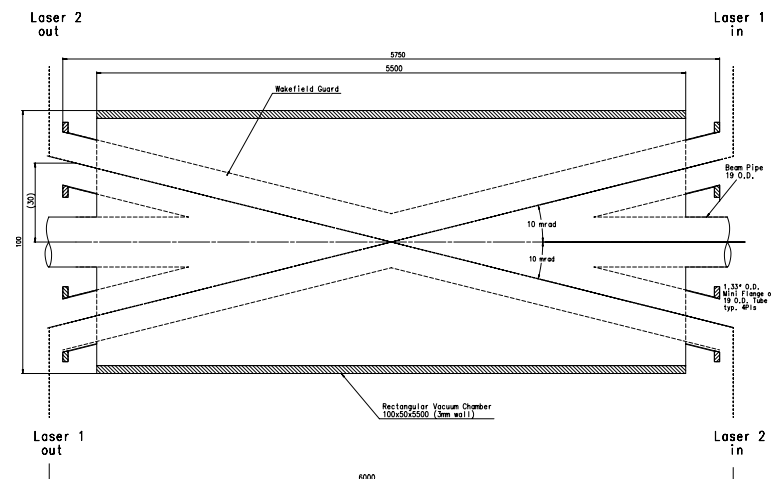
Attention: deflection of chicane the otherway round as on previous page!

Vacuum Chamber: Laser in / out.

- Laser enters chicane *horizontally* (far side from tune-up dump line!)
- final mirror / lens movable to adjust to e^- beam
- had been designed to some extent for TESLA (!) by N. Meyners, P. Schüler

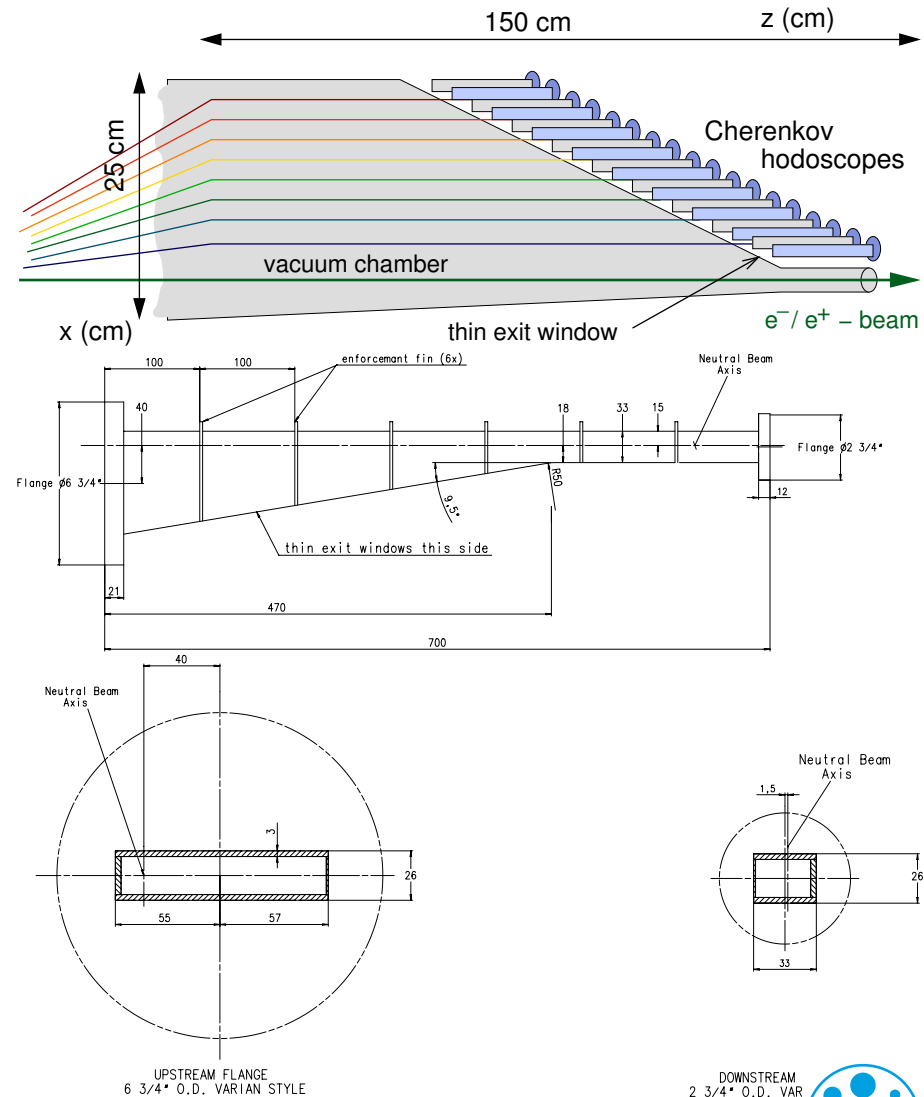


Movable Laser Beam

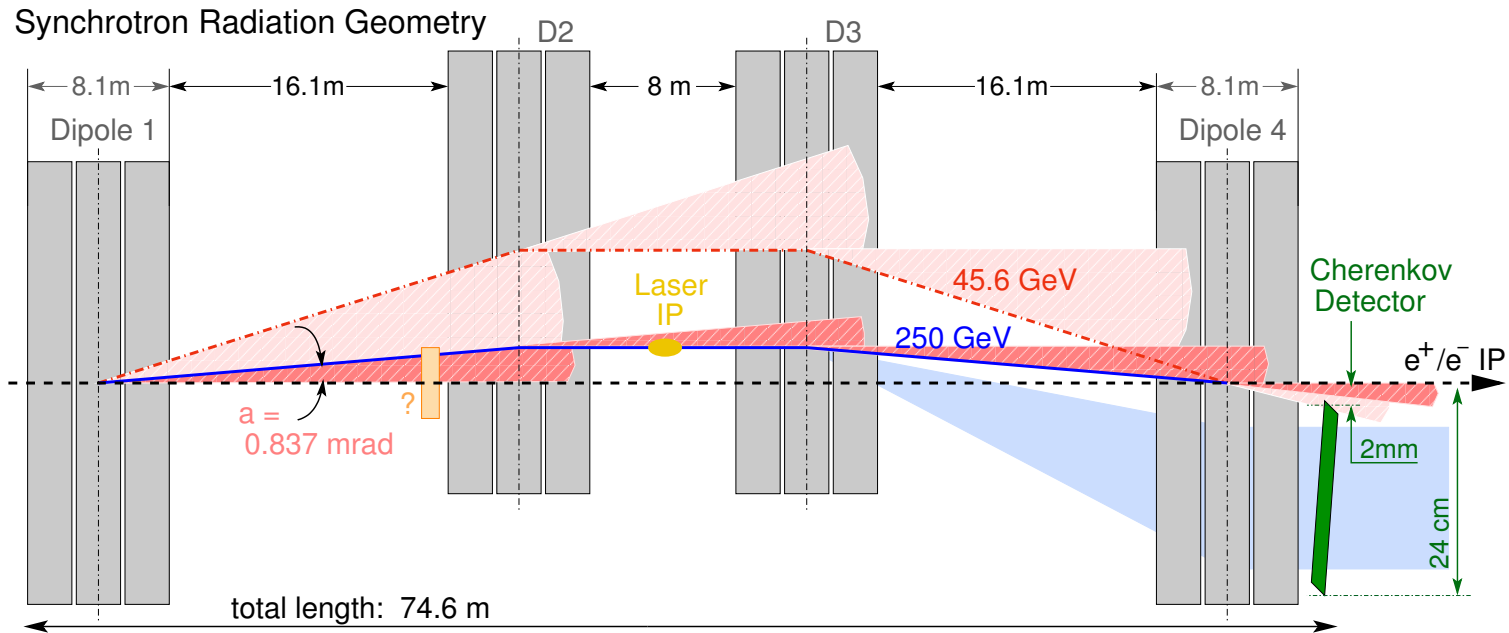


Vacuum Chamber: Compton fan exit.

- need tapered exit window to avoid wake fields
- again estimate from TESLA: $\simeq 10^\circ$ is fine (opinions?)
- need $\simeq 1.5$ m for detector array, make it 2 m for shielding, accessibility,...
- fine with SB2009-Nov10 lattice

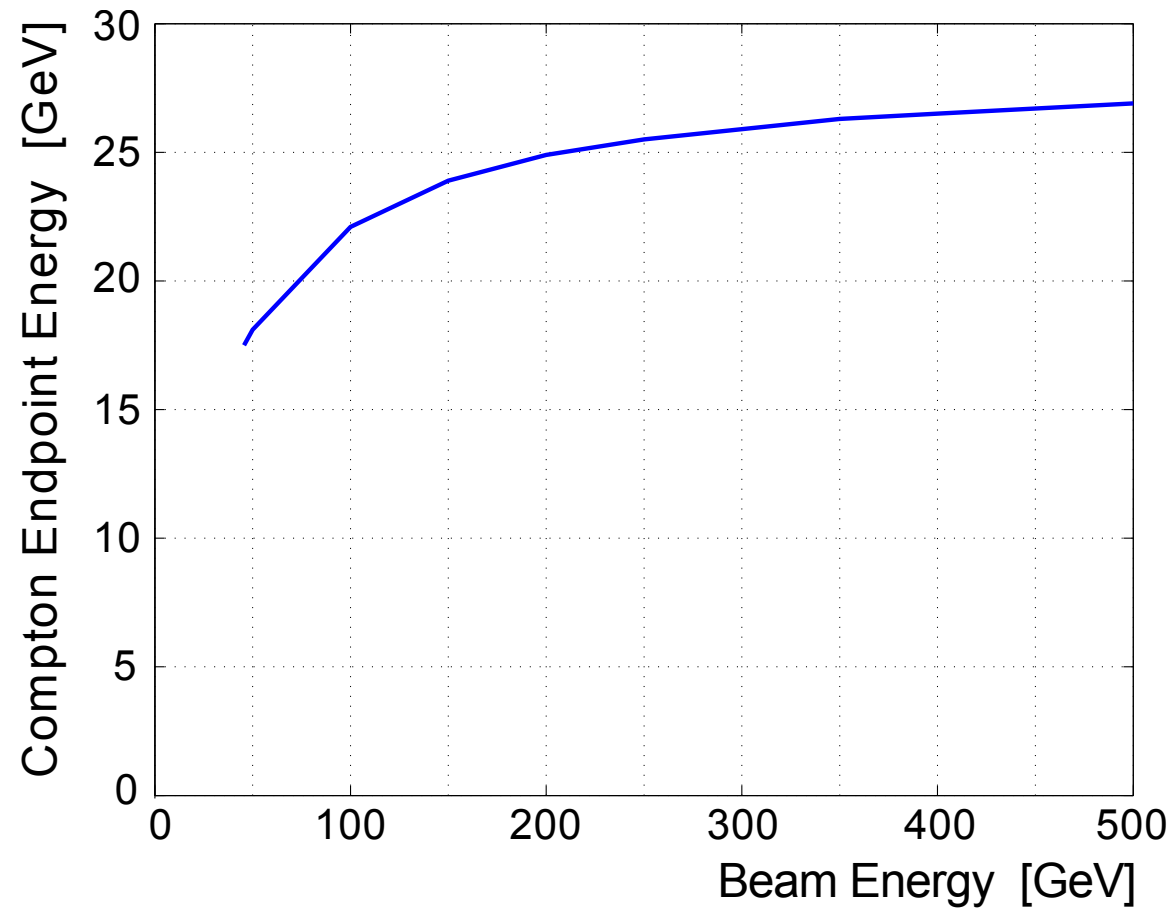


Synchrotron Radiation.



Compton edge.

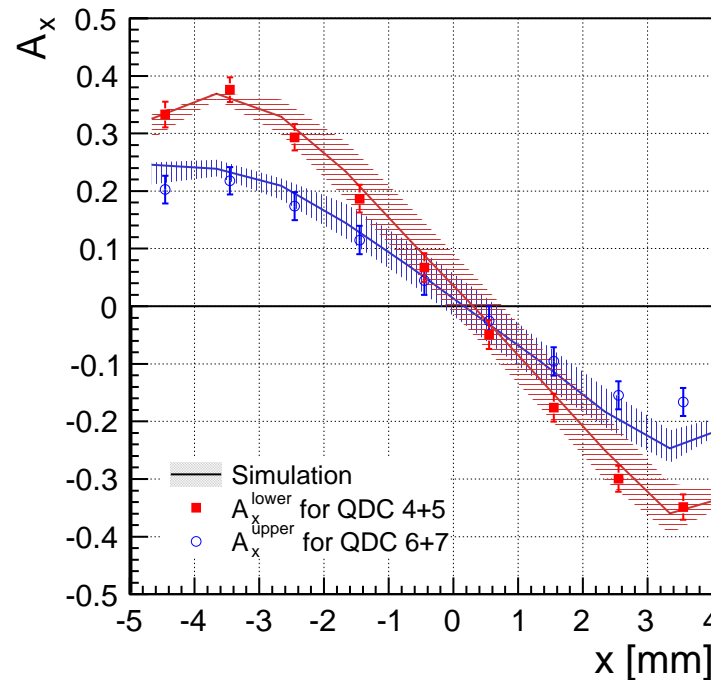
Compton edge position nearly independent of beam energy



Gas Cherenkov detector: Alignment.

If the detector is tilted

- beam path through the detector varies \Rightarrow different light path
- different light pattern on the photocathode
 \Rightarrow alignment via spatial assymetries possible:

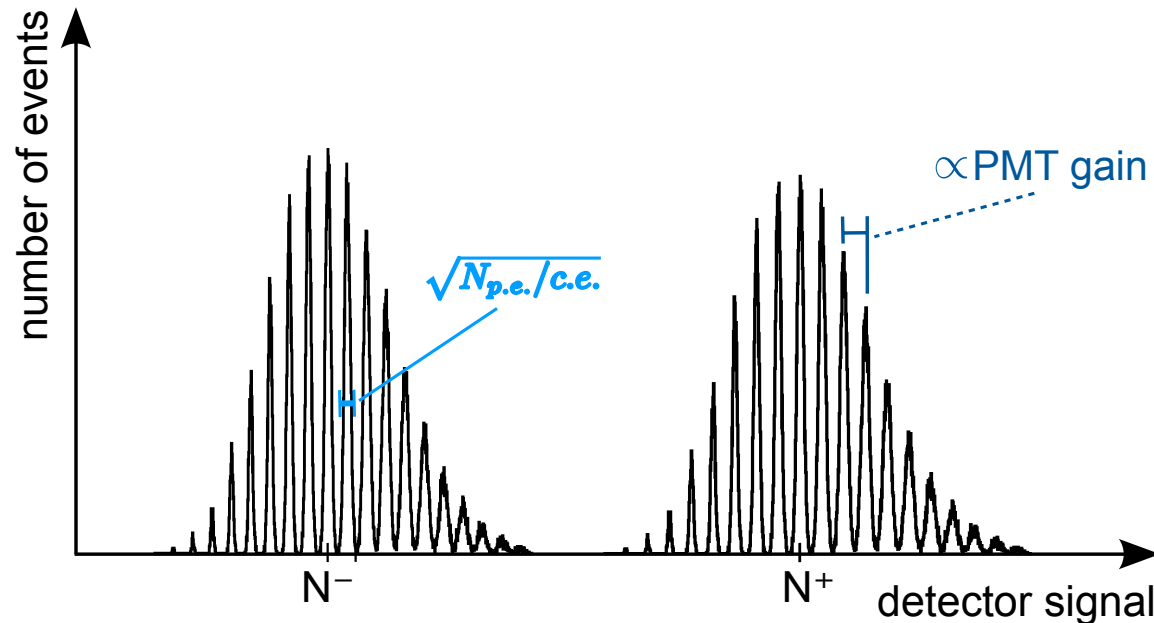


\Rightarrow Reached a tilt alignment of 0.1° . [JINST 7, P01019 (2012)]

Quartz Cherenkov detector.

Alternative detector concept: quartz detector

- Higher refractive index \rightarrow higher photon yield
- For enough photons per Compton e^- :
 \rightarrow calibrate gain directly from the data

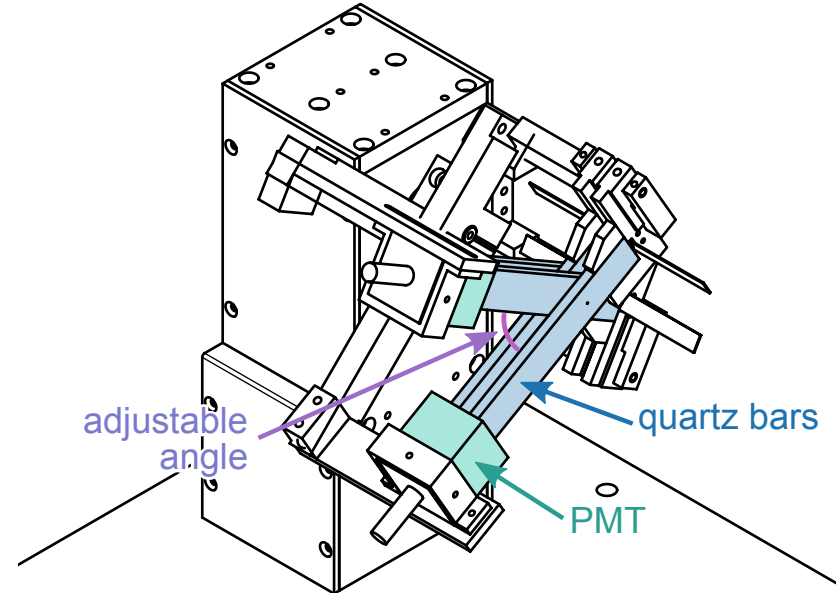


4-channel prototype operated at DESY II testbeam in 2014.

Quartz Cherenkov detector (2).

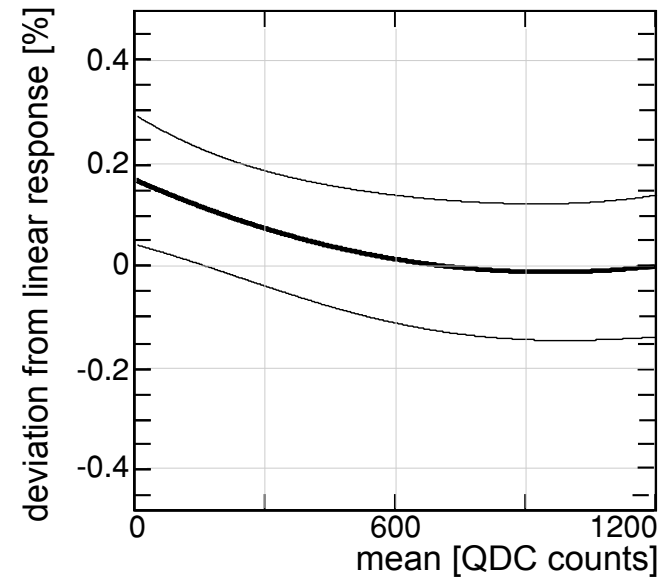
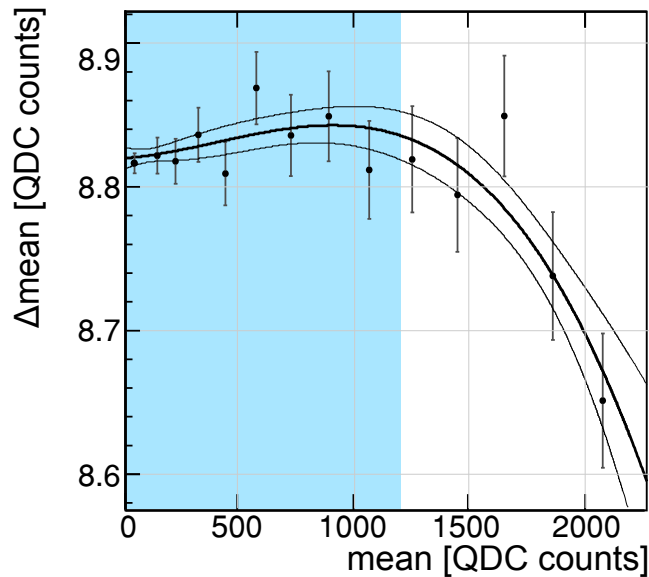
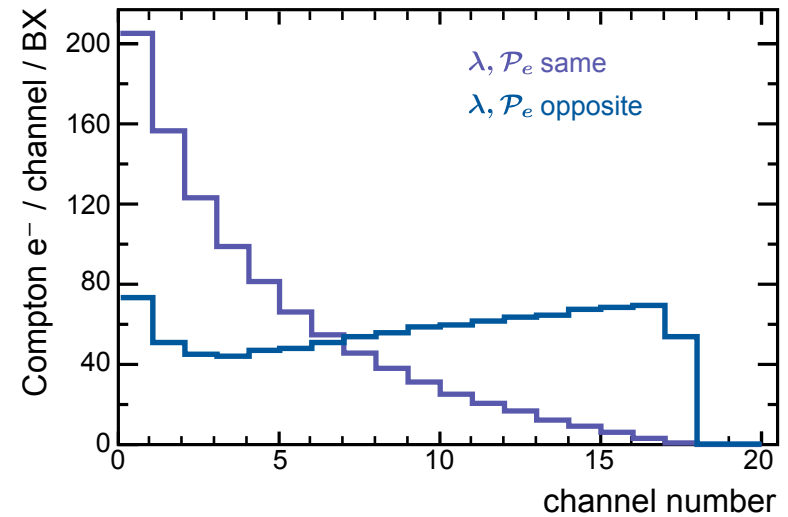
4-channel prototype operated at DESY II testbeam in 2014

- channels: quartz bars (5 mm x 18 mm x 100 mm)
- qualitative agreement with simulations (angular dependence, etc.)
- light yield smaller than predicted, studies ongoing



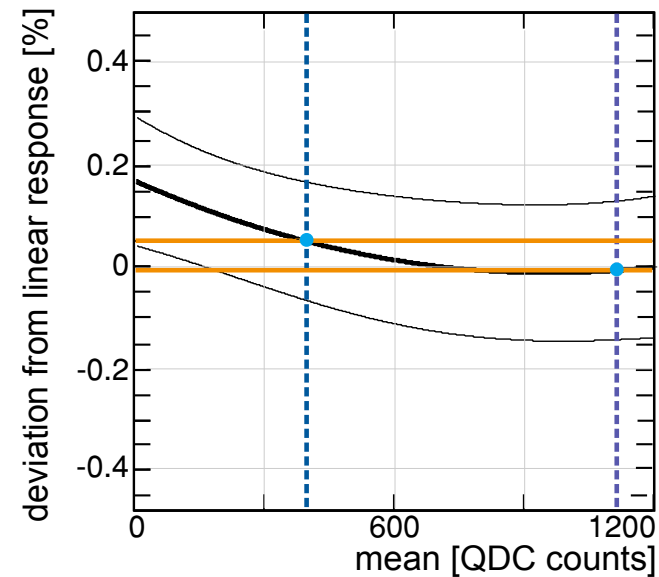
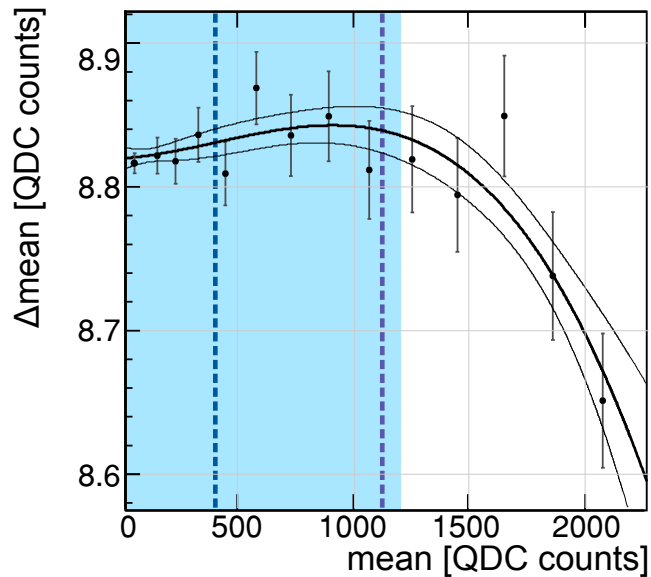
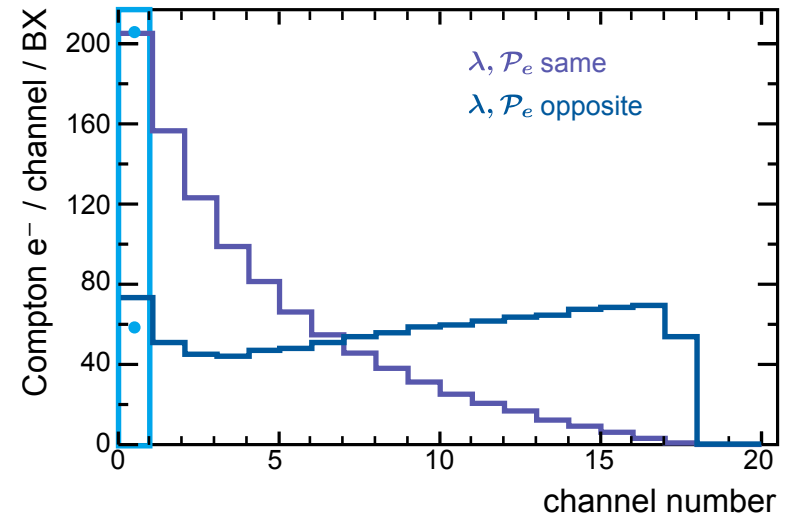
Non-linearity in extreme polarimeter channels.

- up to 210 Compton e^- (~ 1200 QDC counts)
- overall non-linearity already small in this range (max 0.2%)
- in single channels even smaller



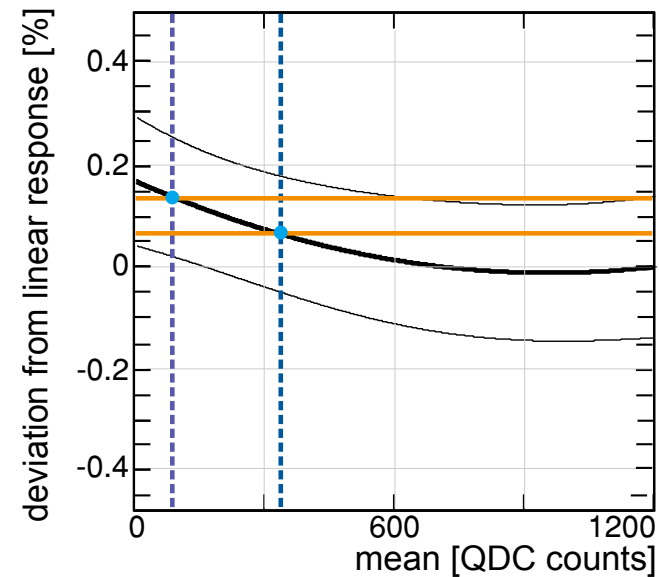
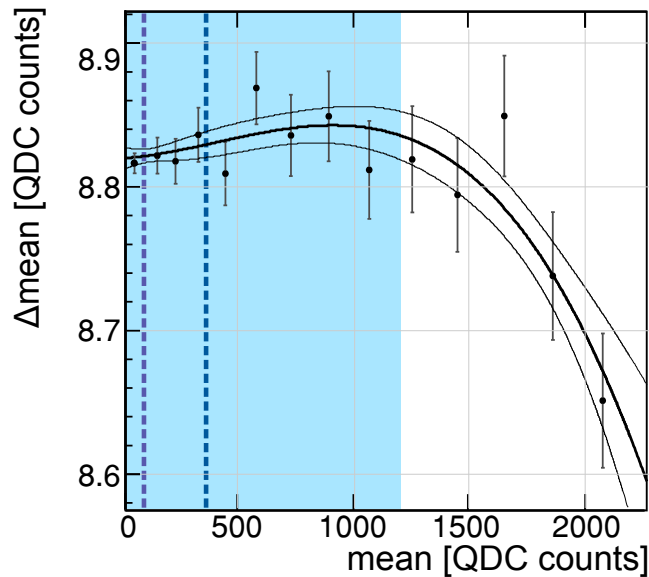
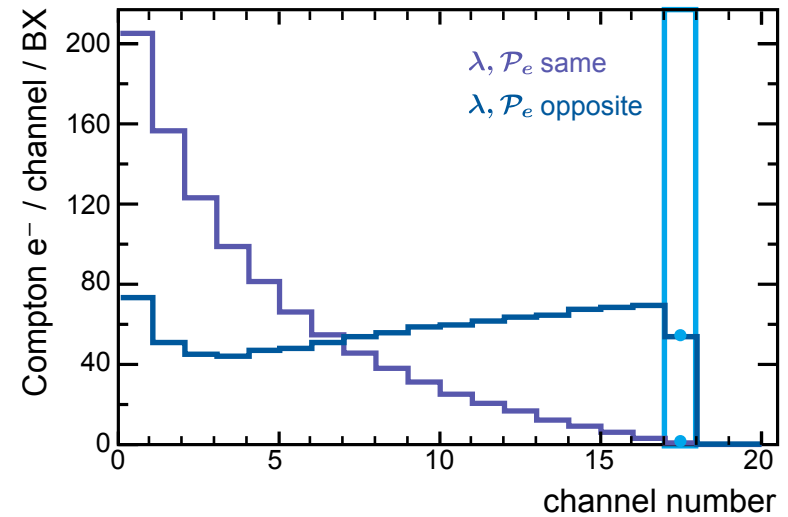
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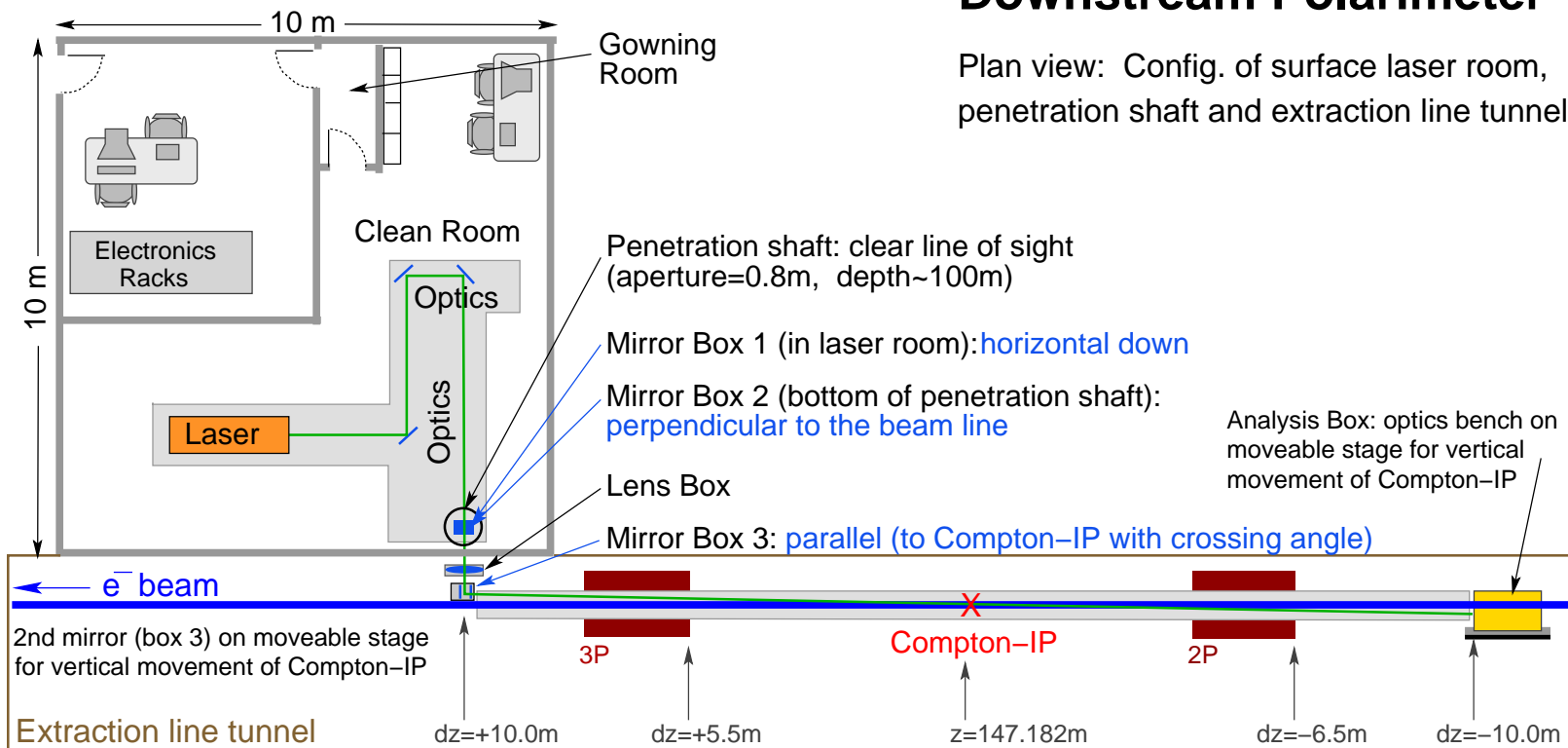
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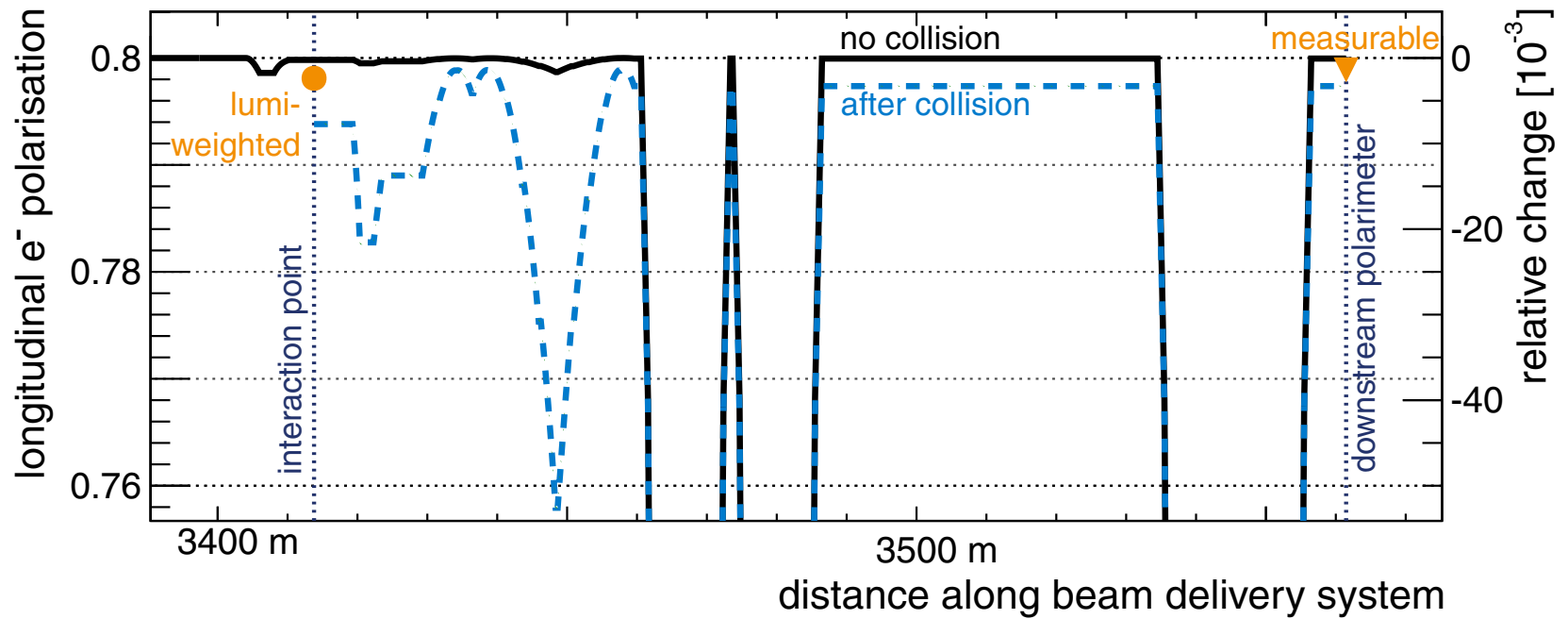


Laser Room.

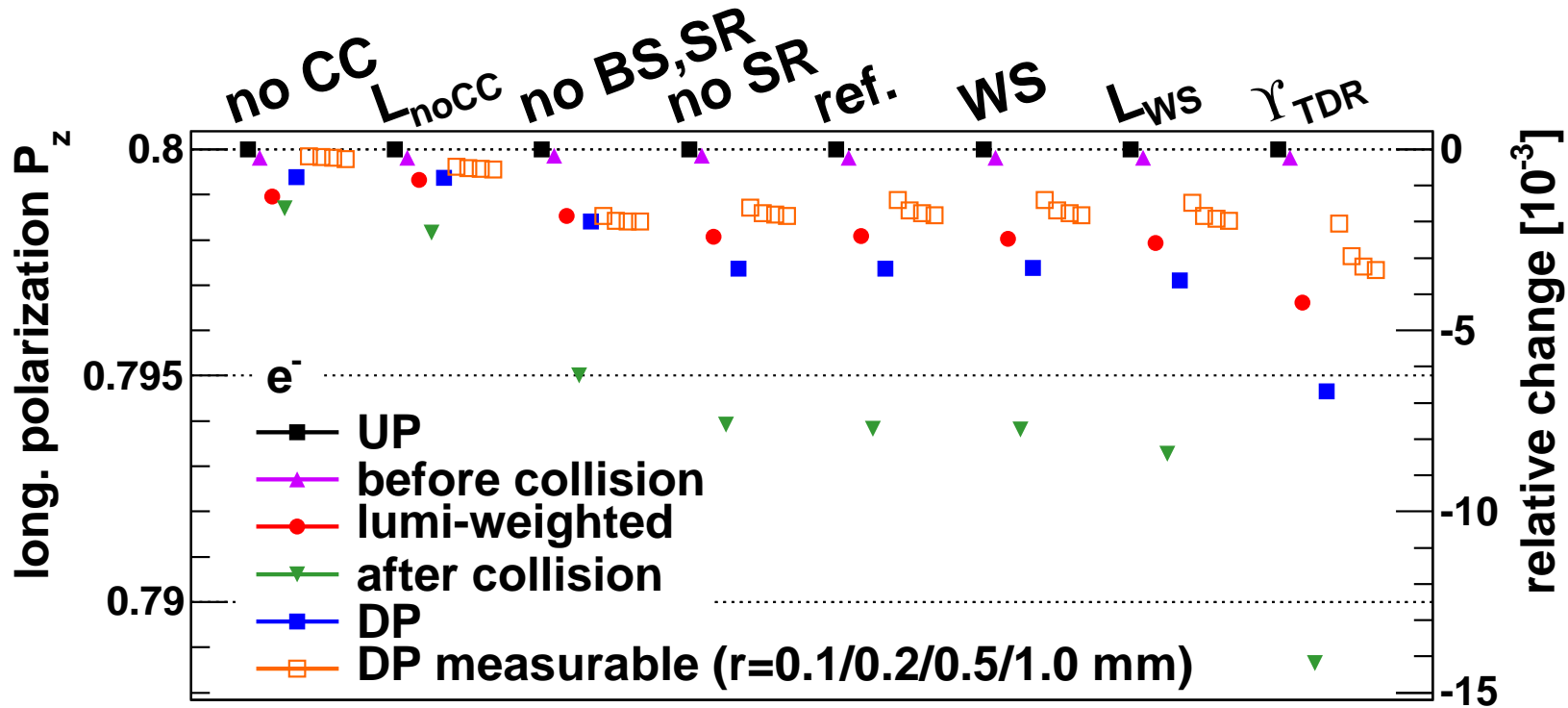
Laser Room on surface (10m x 10m x 3m)



Spin transport.



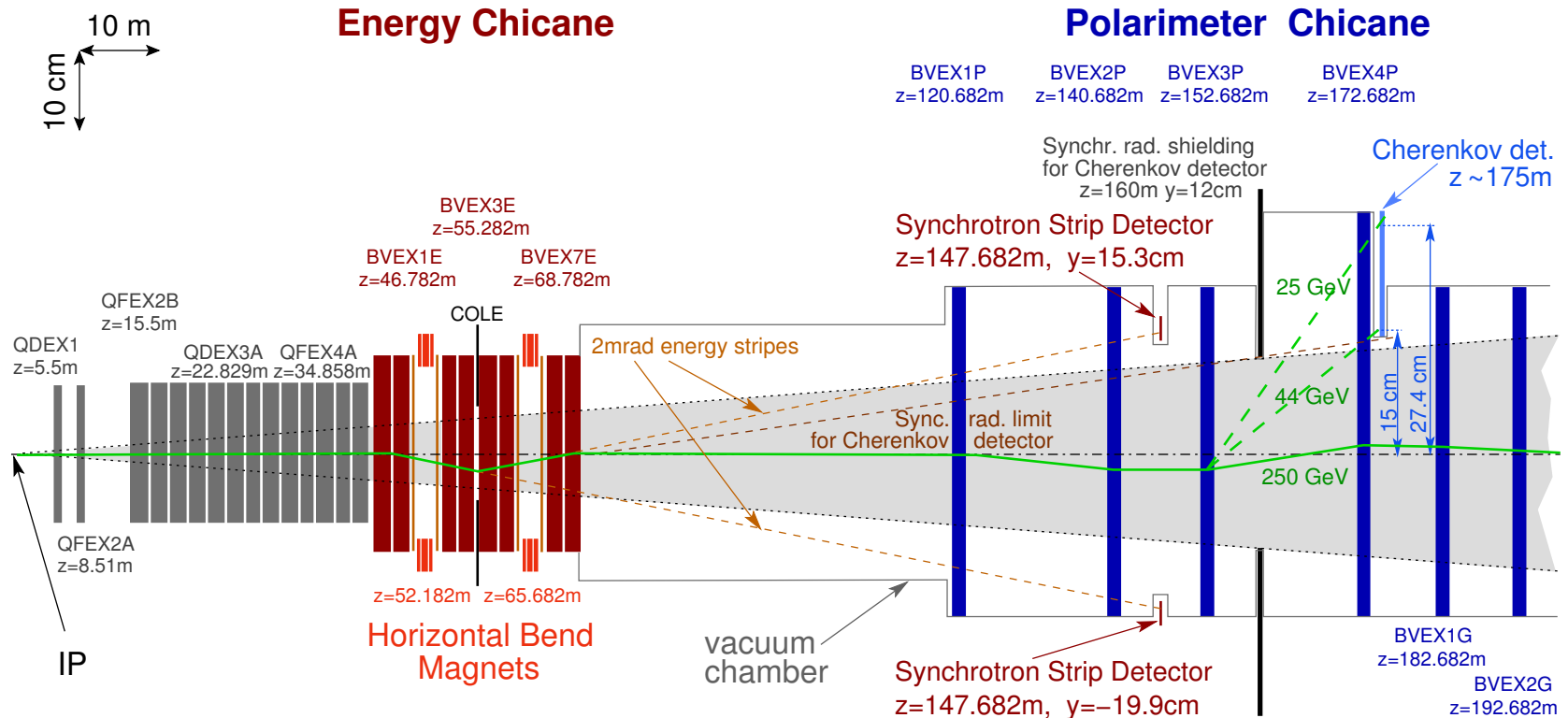
Spin tracking (more).



Downstream Polarimeter.

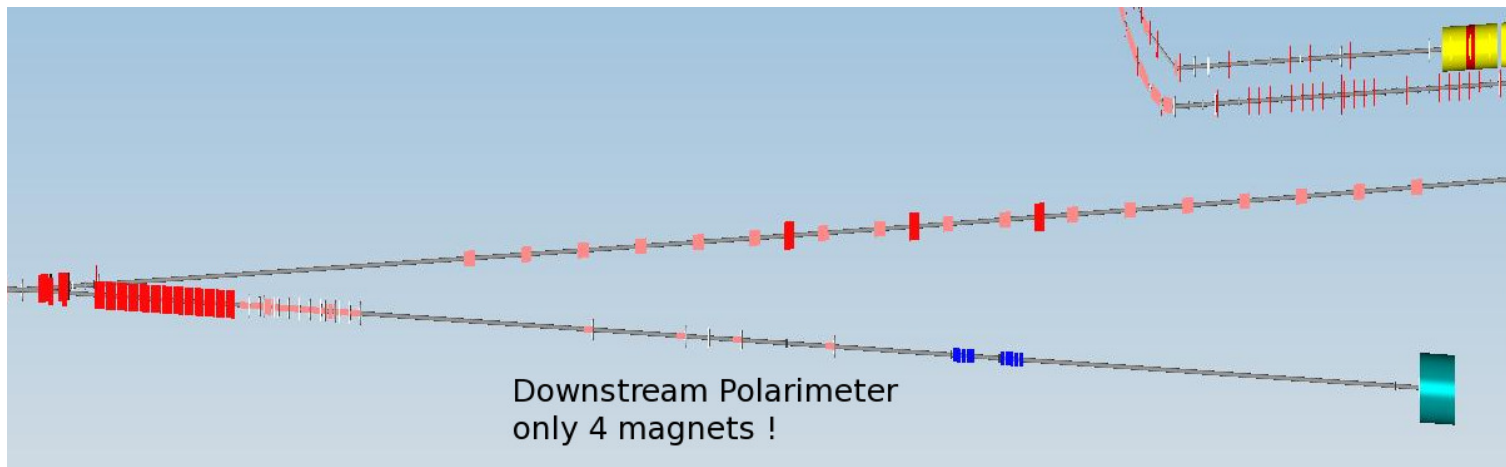
6-magnet chicane suggested in 2007 by Ken Moffeit et al:

- kick Compton e^- further out of the synchrotron radiation fan



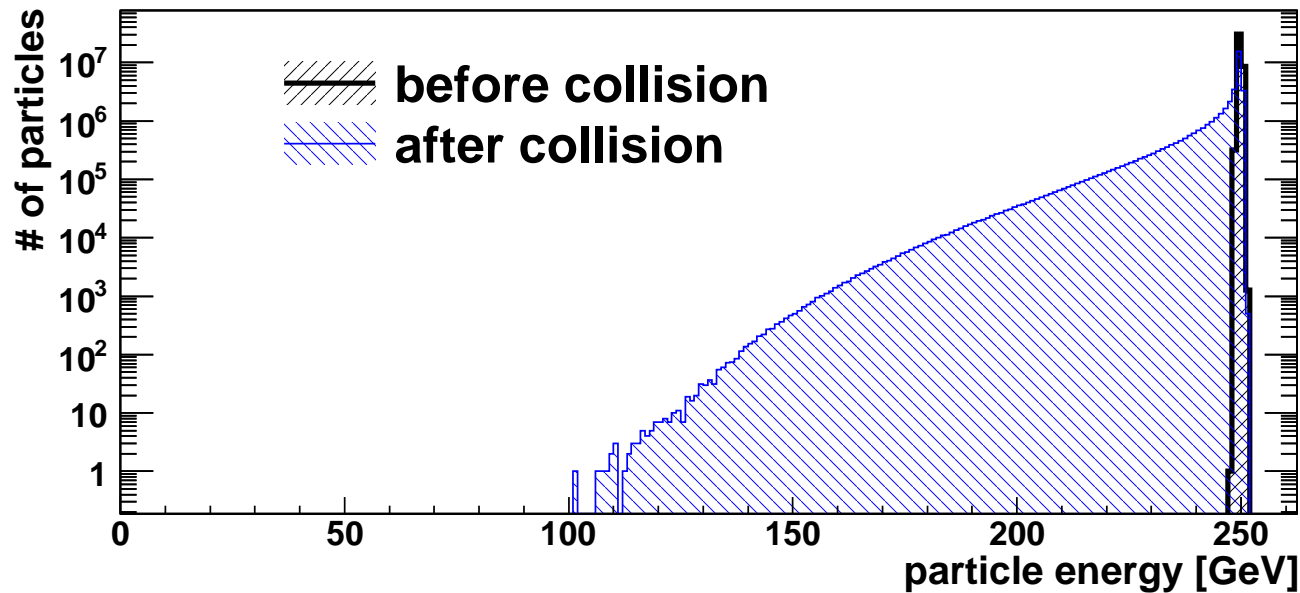
Downstream Polarimeter in SB2009-Nov10 lattice.

- still 4-magnet chicane - should be upgraded to 6-magnet design as proposed in SLAC-PUB-12425
- necessary due to push-pull related changes to the extraction SC quadrupoles
- at the same time gives better shielding of magnets due to additional collimators
- even more impact due to worse spent beam in low power configuration....



Beam Energy Spectrum with Collisions.

GuineaPig++, RDR nominal, $\sqrt{s} = 500$ GeV

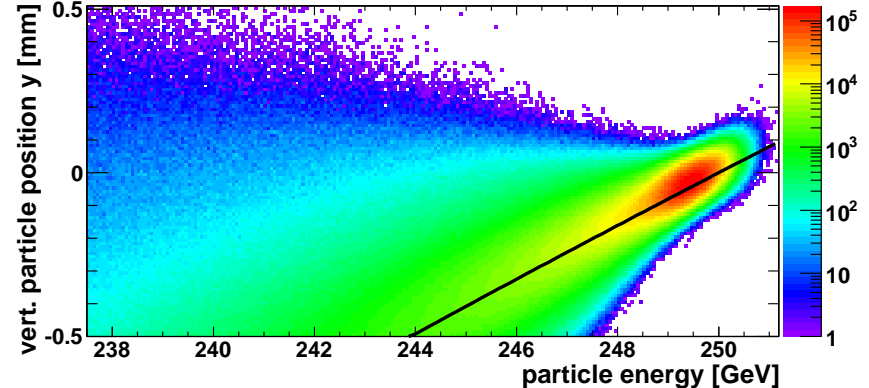
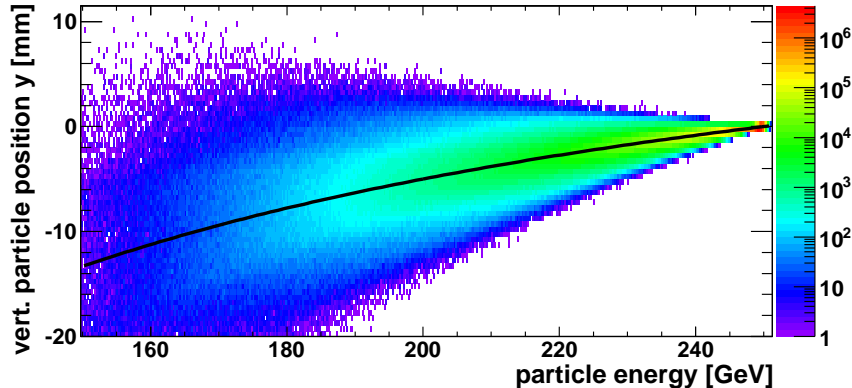


How will this influence the measurement at the downstream polarimeter?

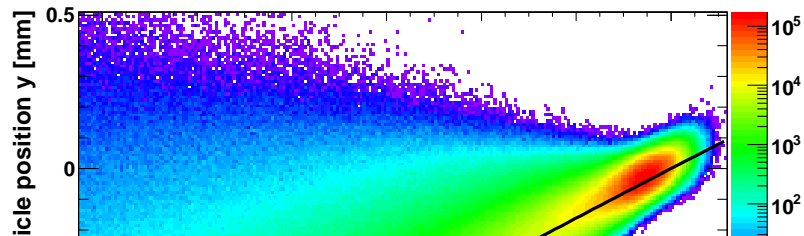
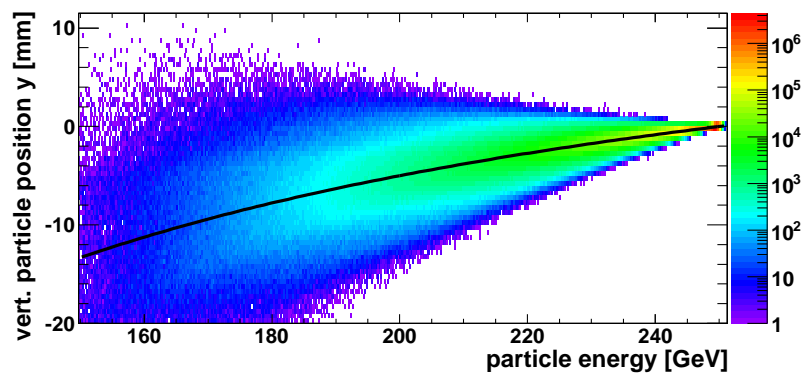
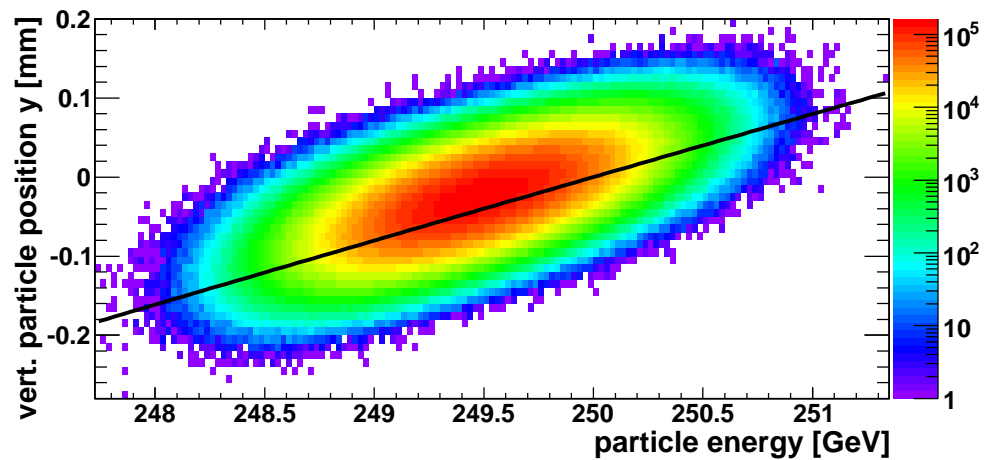
y vs E at DP IP.

Particle Tracking through SB2009-Nov10 lattice (M.Beckmann)

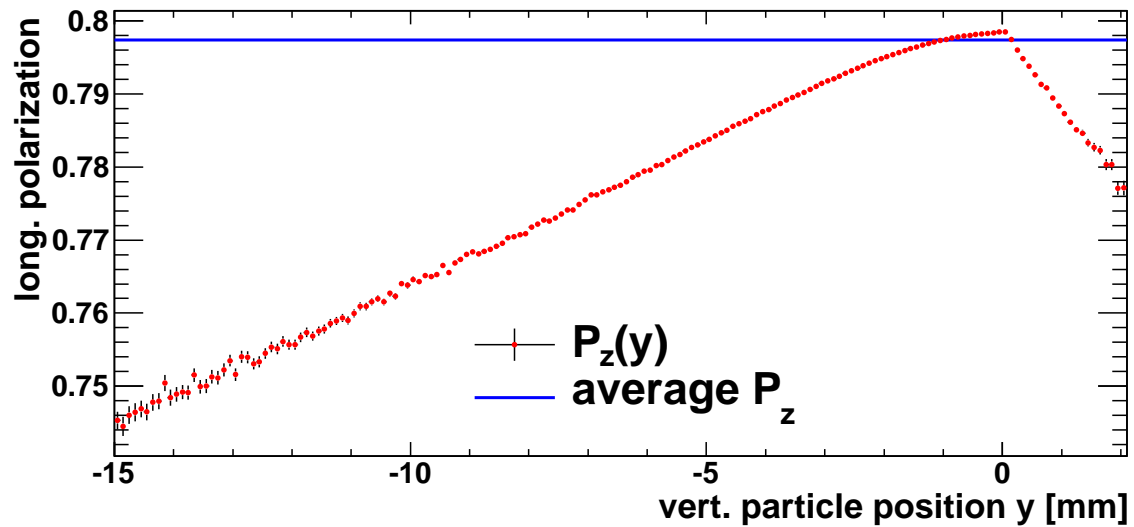
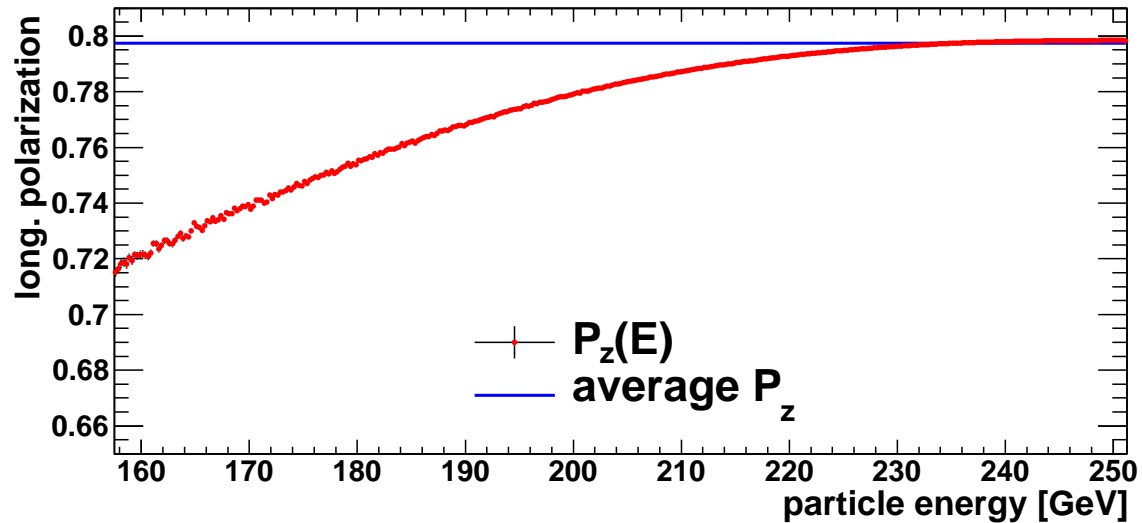
- vertical extension of spent beam at DP IP $\mathcal{O}(\text{cm})$
- “core” size still $\simeq 0.5 \text{ mm}$
- sizable correlation of energy and position
- which part will the laser sample?
- expect dependence of measured polarisation on laser spot size and laser-beam alignment



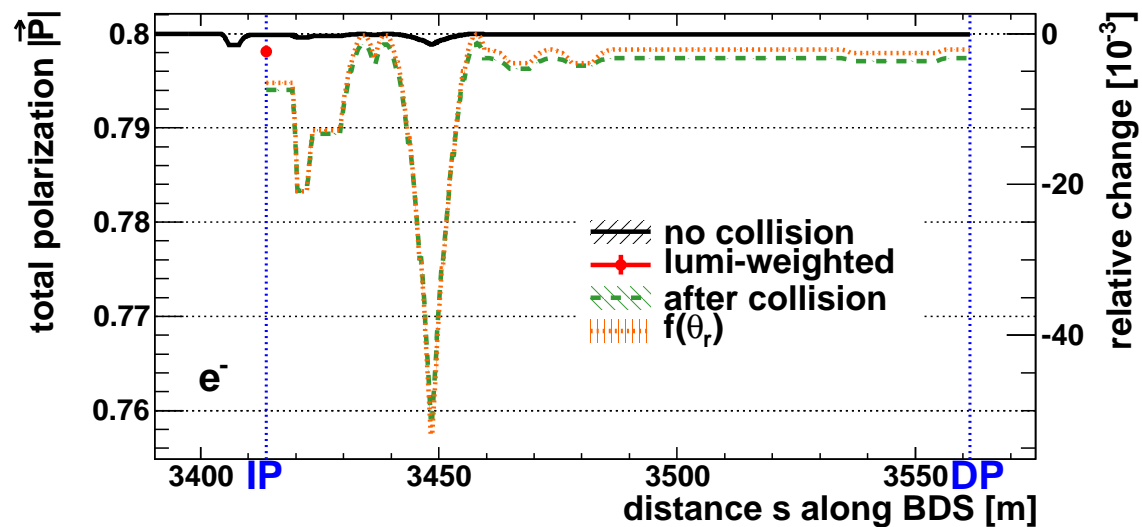
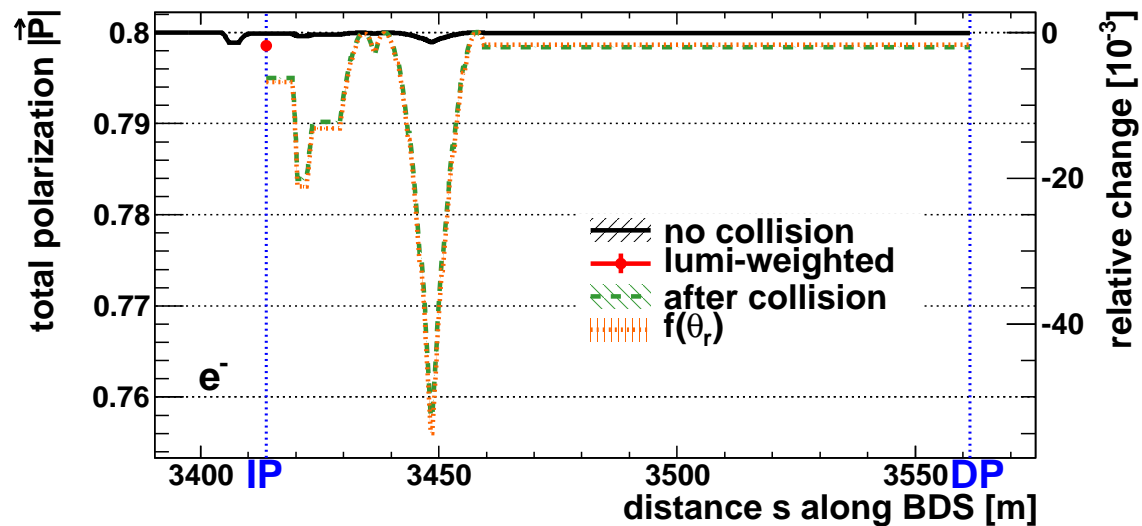
Downstream Polarimeter: y vs E .



Downstream Polarimeter: P_z vs E , P_z vs y .



Total Polarisation IP \rightarrow DP



Polarisation for Physics.

Longitudinal polarisation $P_z = \frac{N_R - N_L}{N_R + N_L}$

with $N_{R,L}$: number of right-/left-handed particles in bunch

- SM & BSM: left- and righthanded particles couple differently
 - polarised cross-sections are important observables carrying **qualitatively** new information!
 - beam polarisation can suppress background / enhance signal

- wanted for physics: **luminosity weighted average polarisation**

at the IP, $\langle P_z \rangle_{IP} = \frac{\int P_z(t) \mathcal{L}(t) dt}{\int \mathcal{L}(t) dt}$

- Note: most physics studies sofar assume this average is known exactly and independently for e^- and e^+ beam.

$P \equiv P_z$ in the following.

Polarised Cross-sections.

$$\sigma_{P_{e^-} P_{e^+}} = \frac{1}{4} \left\{ \begin{aligned} &(1 + P_{e^-})(1 + P_{e^+})\sigma_{RR} + (1 - P_{e^-})(1 - P_{e^+})\sigma_{LL} \\ &+ (1 + P_{e^-})(1 - P_{e^+})\sigma_{RL} + (1 - P_{e^-})(1 + P_{e^+})\sigma_{LR} \end{aligned} \right\}$$

processes with s-channel Z/γ exchange only:

- $\sigma_{RR} = \sigma_{LL} = 0$
- $4\sigma_{P_{e^-} P_{e^+}} = (1 - P_{e^-} P_{e^+})(\sigma_{LR} + \sigma_{RL})[1 - P_{\text{eff}}^- A_{LR}]$
- with $P_{\text{eff}}^- = 1 - \frac{P_{e^-} - P_{e^+}}{1 - P_{e^-} P_{e^+}}$ and $A_{LR} = \frac{\sigma_{LR} - \sigma_{RL}}{\sigma_{LR} + \sigma_{RL}}$

general case:

- $\sigma_{RR} \neq \sigma_{LL} \neq 0$
- $4\sigma_{P_{e^-} P_{e^+}} = (1 + P_{e^-} P_{e^+})(\sigma_{LL} + \sigma_{RR})[1 + P_{\text{eff}}^+ A_{LLRR}] + \text{above}$
- with $P_{\text{eff}}^+ = 1 + \frac{P_{e^-} + P_{e^+}}{1 + P_{e^-} P_{e^+}}$ and $A_{LLRR} = \frac{\sigma_{LL} - \sigma_{RR}}{\sigma_{LL} + \sigma_{RR}}$

Polarisation Averages.

Absolute cross-section measurements require:

- $\langle P_{e^\pm} \rangle_{IP} = \frac{\int P_{e^\pm}(t) \mathcal{L}(t) dt}{\int \mathcal{L}(t) dt}$
- $\langle P_{e^-} P_{e^+} \rangle_{IP} = \frac{\int P_{e^-}(t) P_{e^+}(t) \mathcal{L}(t) dt}{\int \mathcal{L}(t) dt}$
- correlations between lumi and polarisation?!

Direct extraction from collision data

- any abundant, well-known, polarisation dependent process:
- $\langle | P_{e^\pm} | \rangle_{IP} = \sqrt{\frac{(\sigma_{-+} + \sigma_{+-} - \sigma_{--} - \sigma_{++})(\pm\sigma_{-+} \mp \sigma_{+-} + \sigma_{--} - \sigma_{++})}{(\sigma_{-+} + \sigma_{+-} + \sigma_{--} + \sigma_{++})(\pm\sigma_{-+} \mp \sigma_{+-} - \sigma_{--} + \sigma_{++})}}$
- σ_{+-} is total cross-section for $P(e^-, e^+) = (+x\%, -y\%)$, etc.
- assumes $P_+(e^-) = -P_-(e^-)$ and $P_+(e^+) = -P_-(e^+)$

Correction to modified Blondel scheme.

$$P_+(e^\pm) = P^\pm + \epsilon^\pm \text{ and } P_-(e^\pm) = P^\pm - \epsilon^\pm$$

