

Jet charge for spin asymmetries

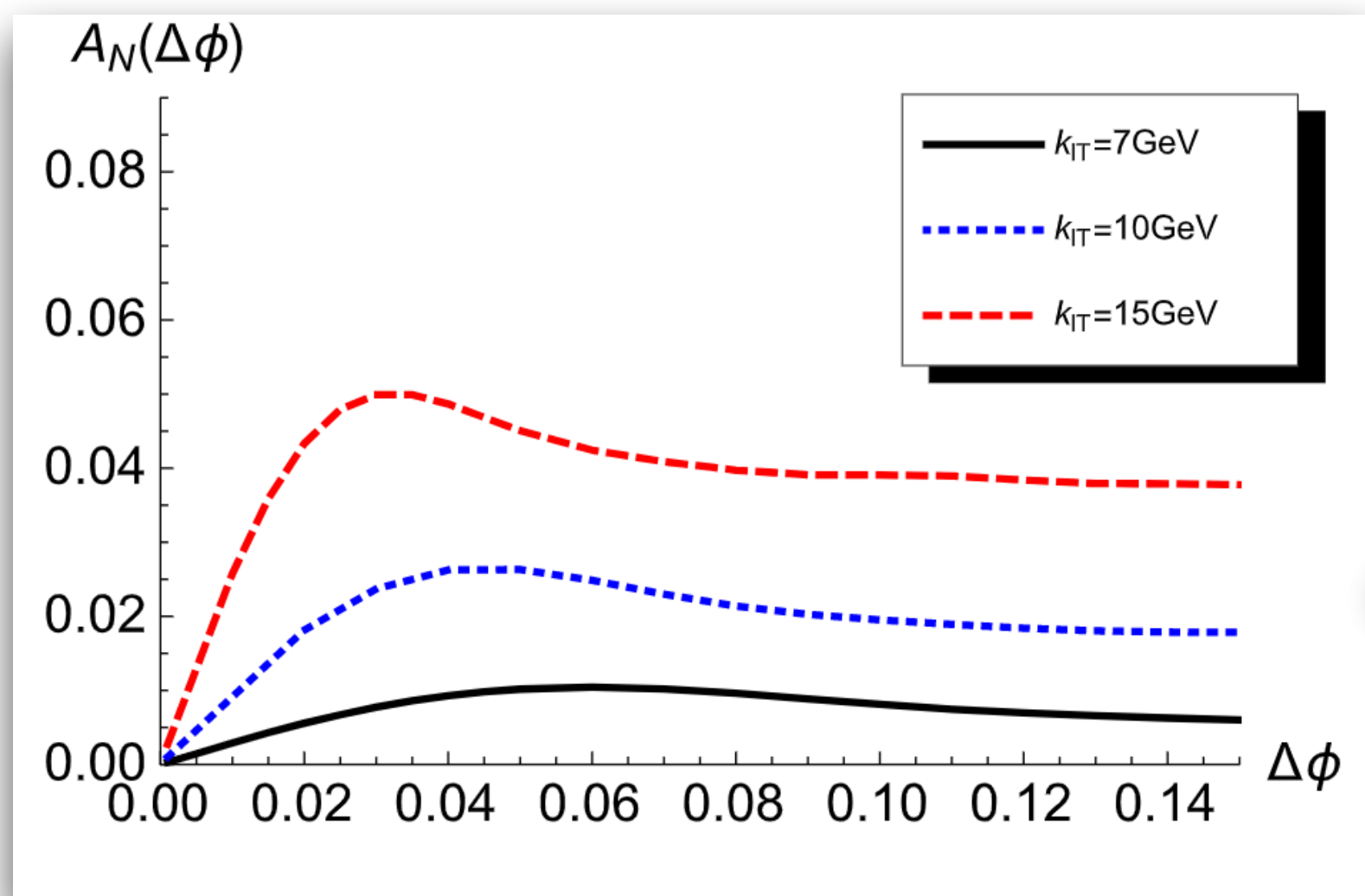
Xiaohui Liu

Jets for 3D imaging, 2020

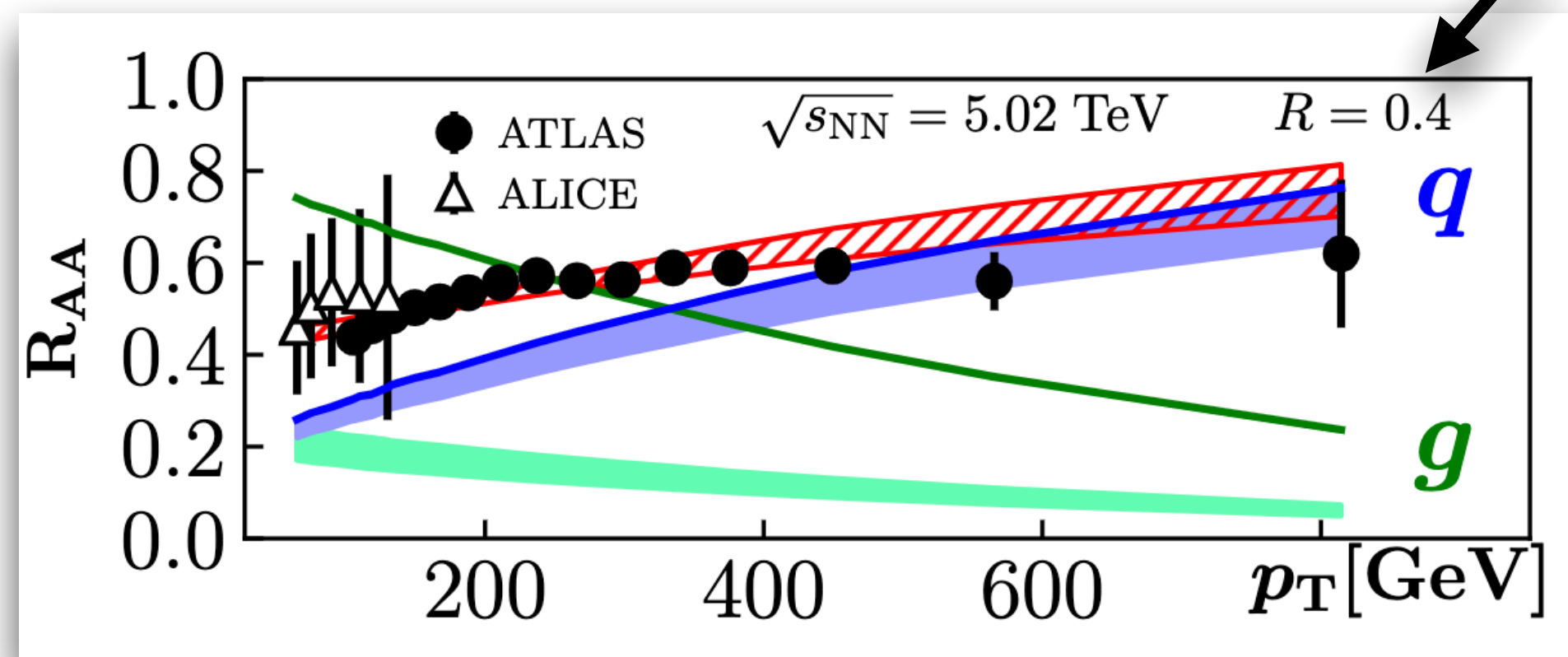


北京師範大學
BEIJING NORMAL UNIVERSITY

Jets @ EIC



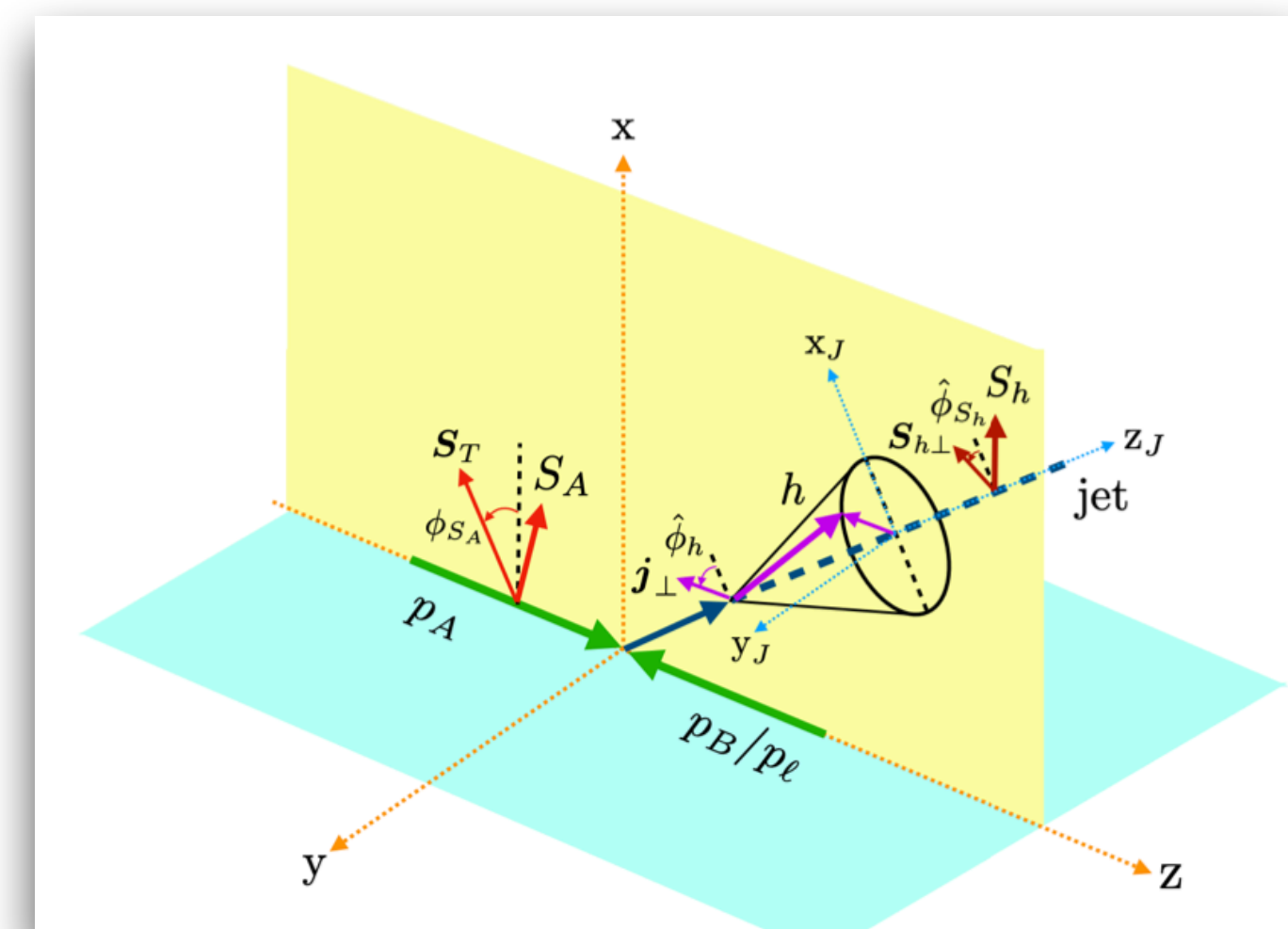
XL, Ringer, Vogelsang, Yuan, 19



Qiu, et al, 19

Proton polarization,
Sivers function

Proton polarization,
Collins function



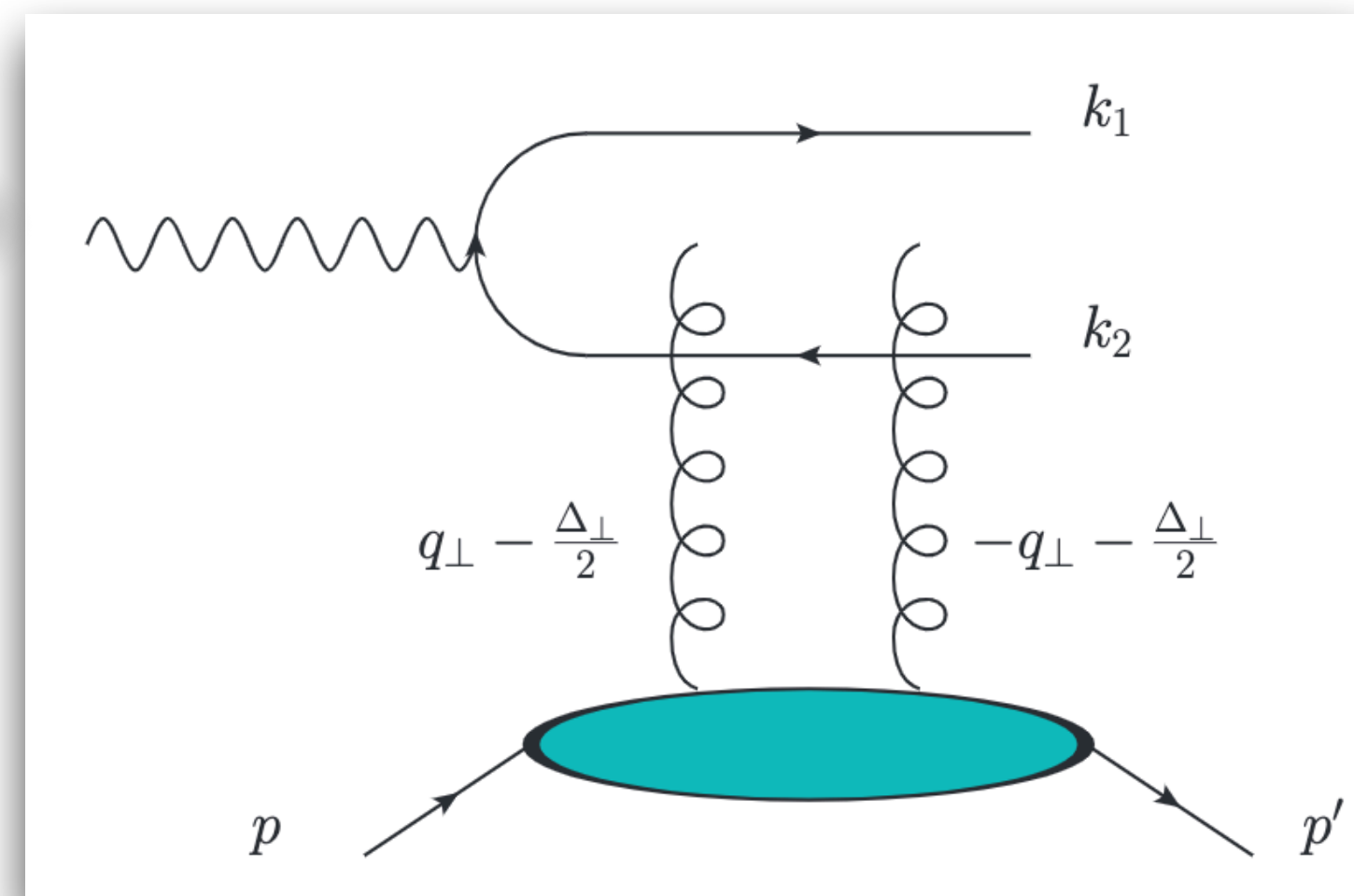
Kang, Lee, Zhao, 20

Jets/substructures

Heavy ion

Small-x, GPD

Growing interests
in jet probes of the
internal structures

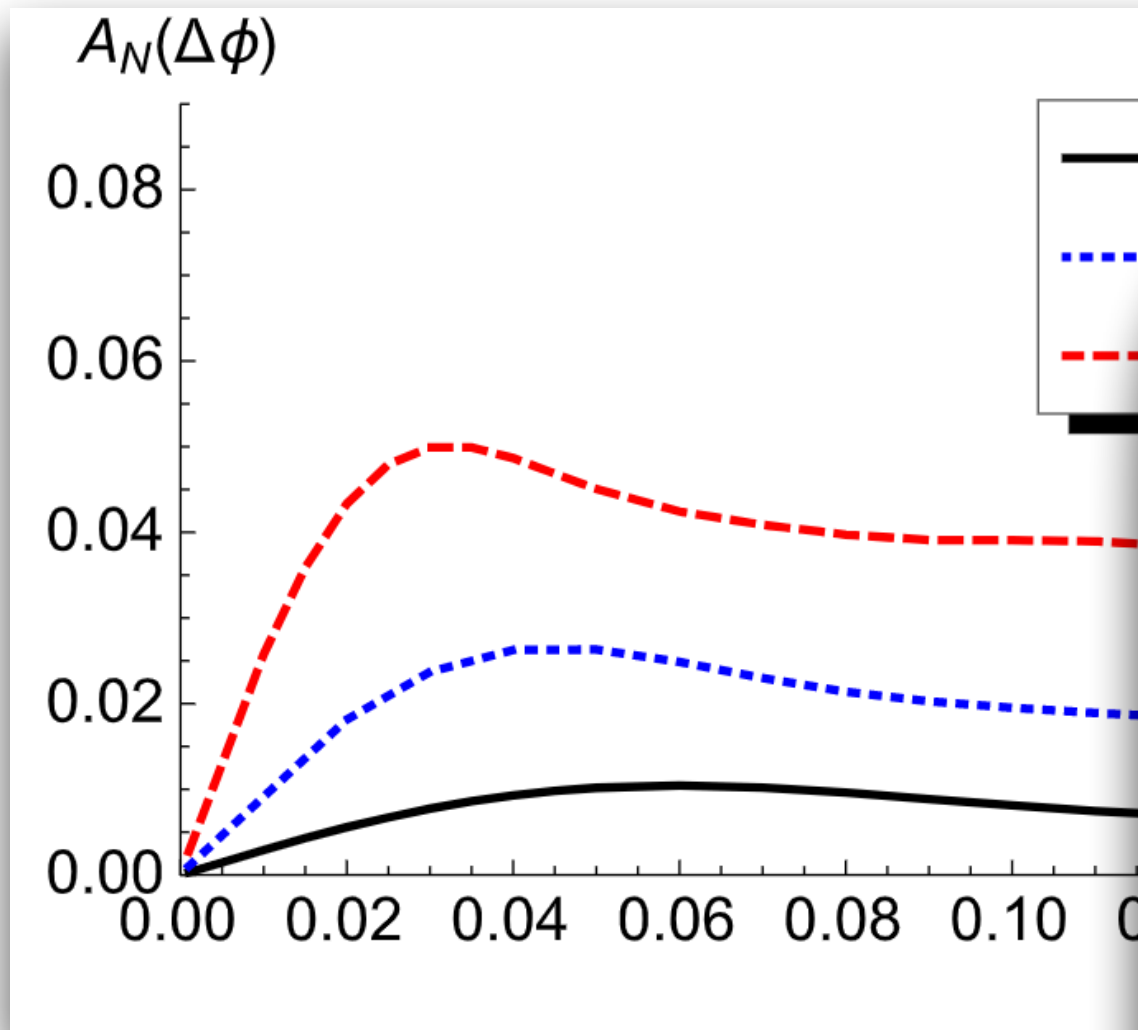


Hatta, Xiao, Yuan, 17

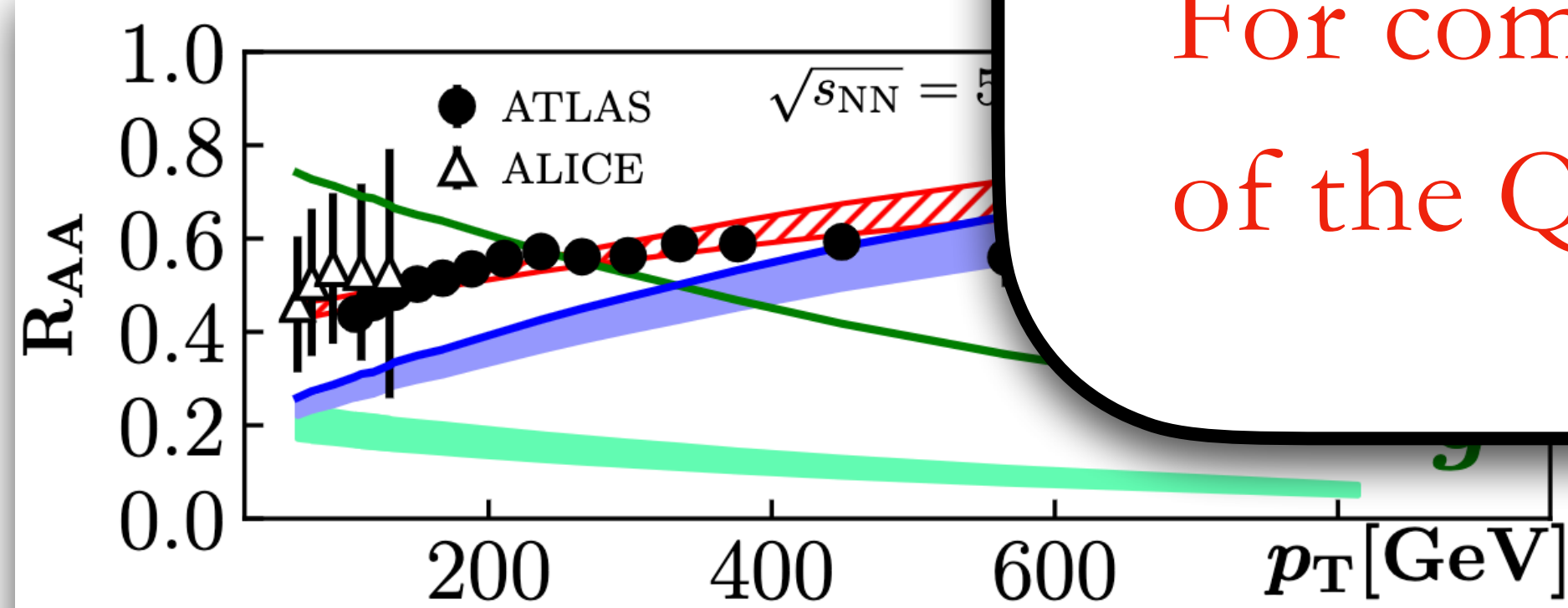
Jets @ EIC

- Better parton proxies, better theoretical control
 - Flexible, dis-entangle between different non-perturbative objects
 - Flavor separation ?
- For complete knowledge of the QCD dynamics

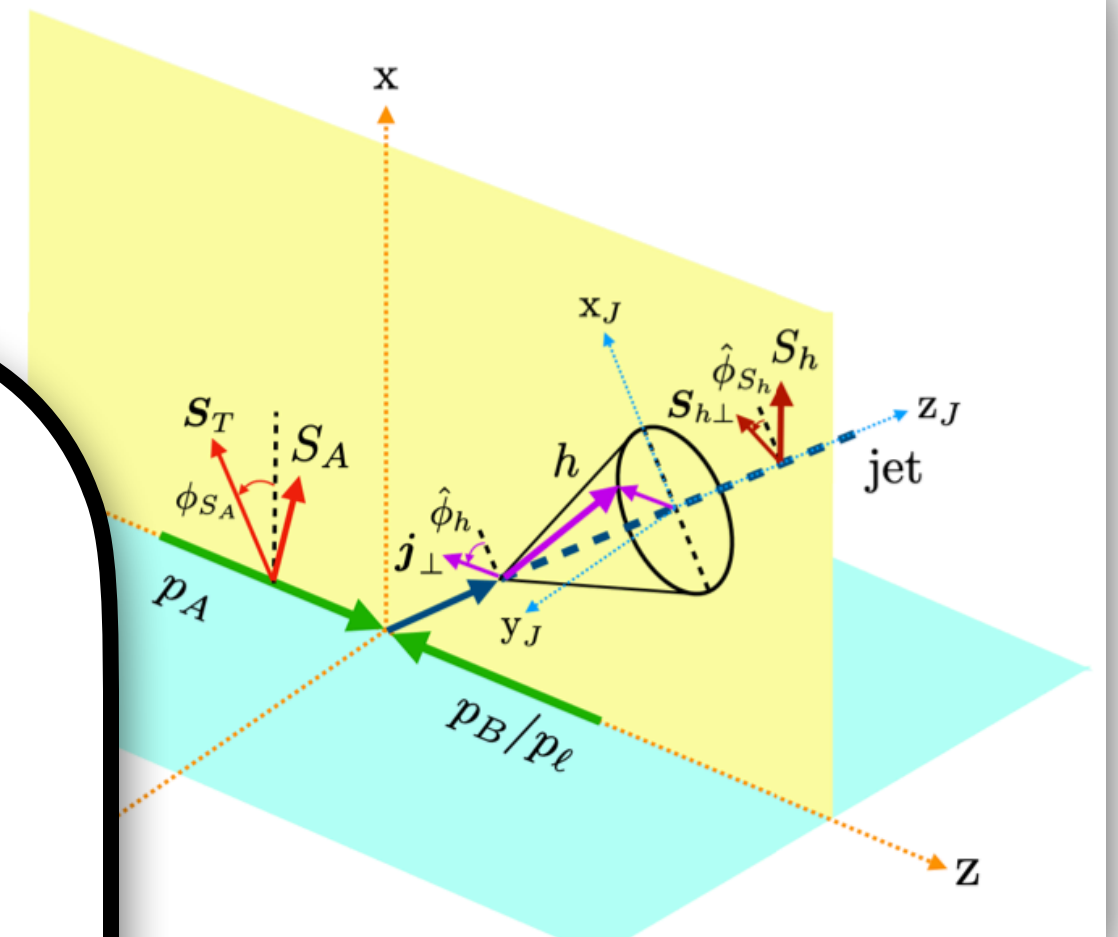
internal structures



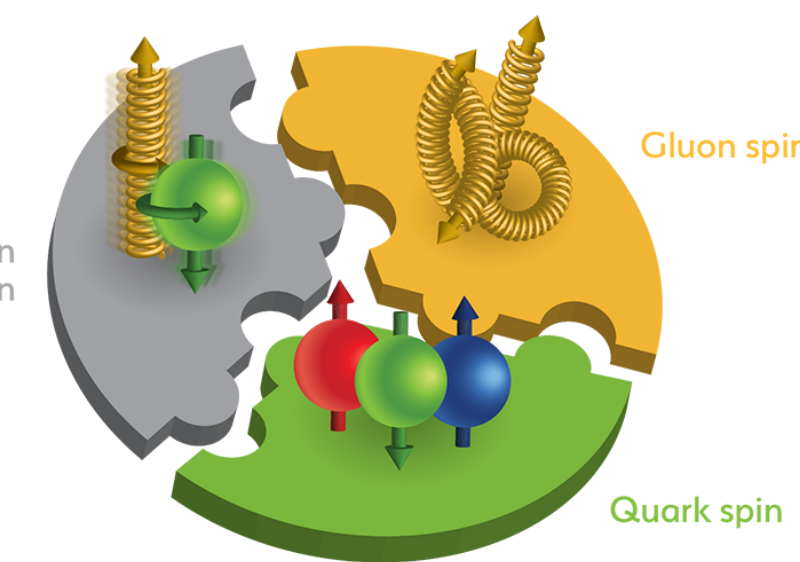
XL, Ringer, Vogelsang



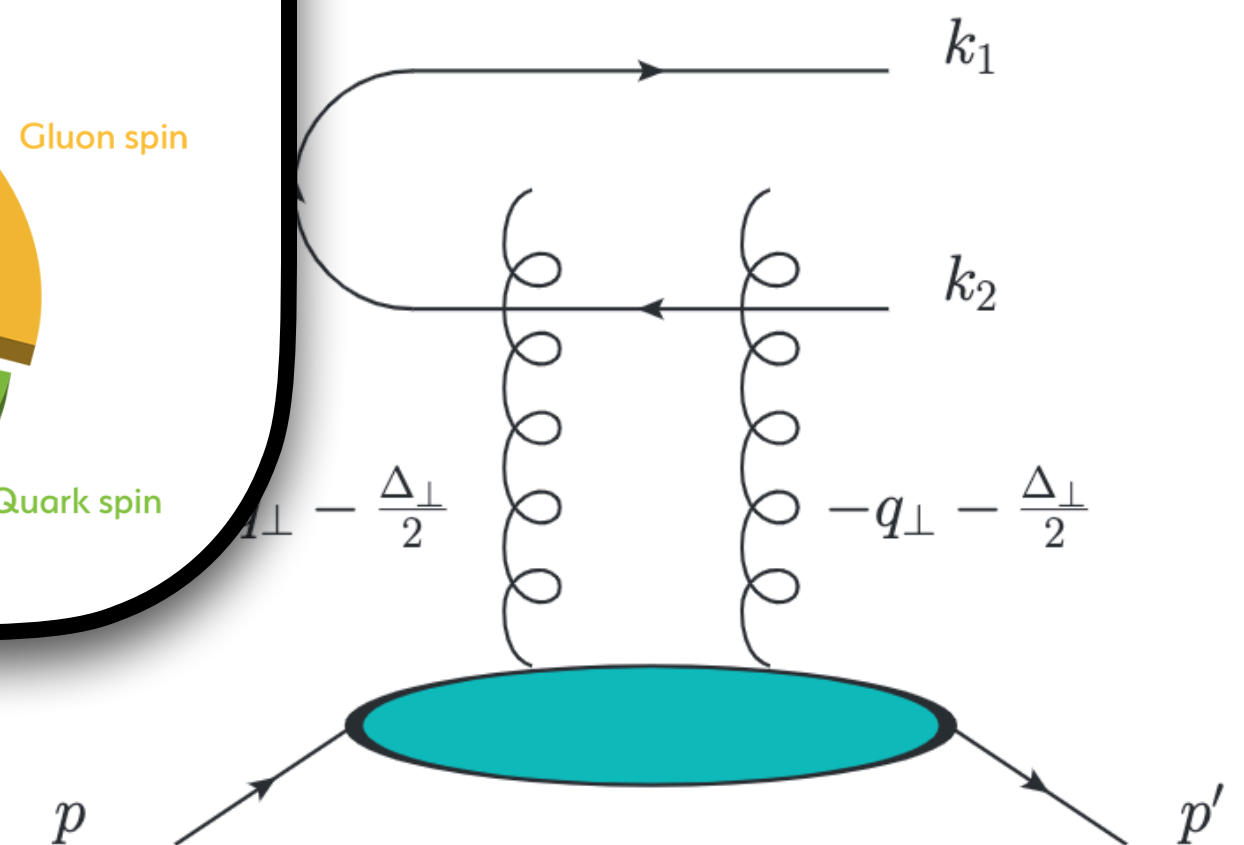
Qiu, et al, 19



Kang, Lee, Zhao, 20

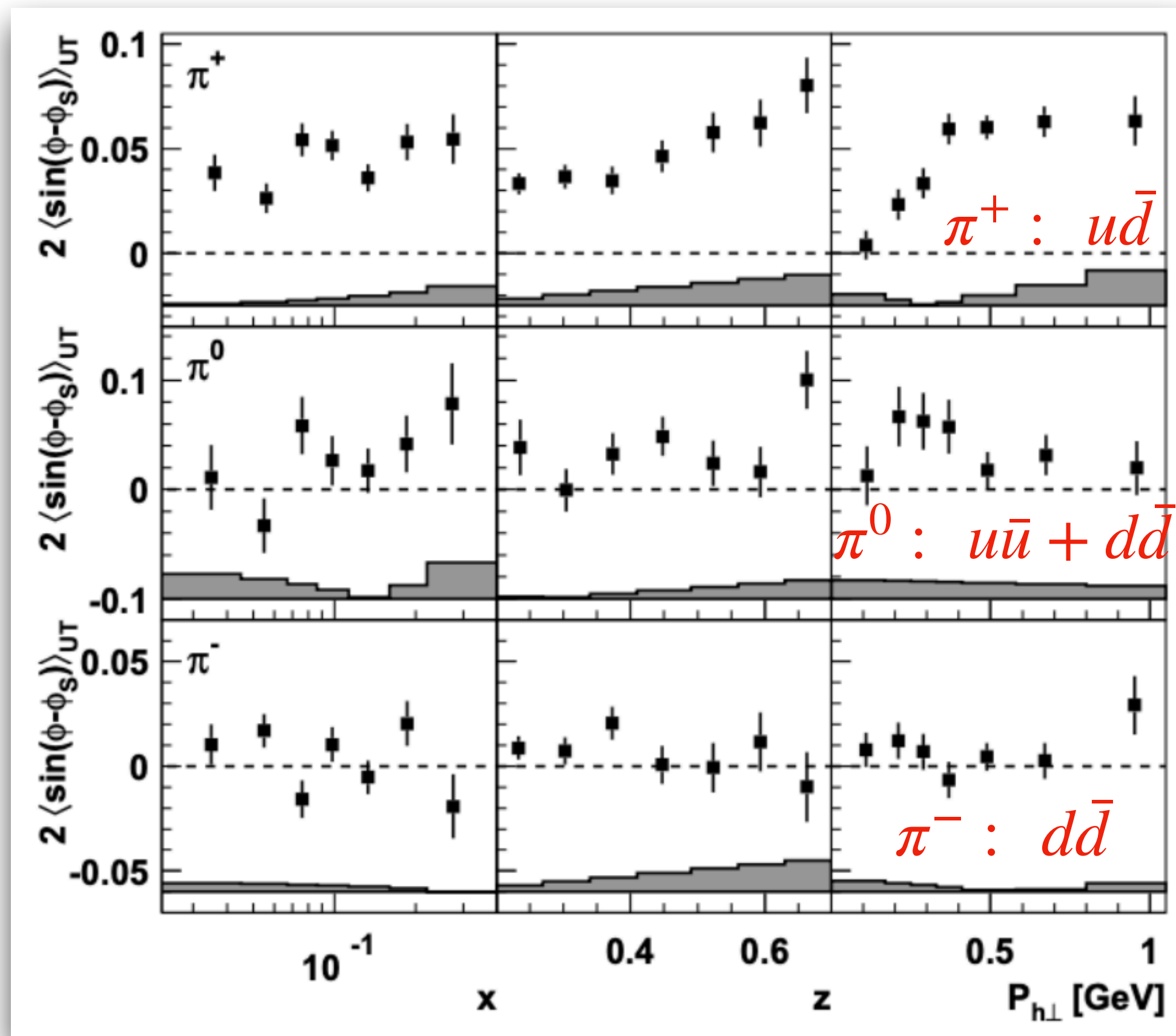


Quark and gluon internal motion



Hatta, Xiao, Yuan, 17

Jets @ EIC

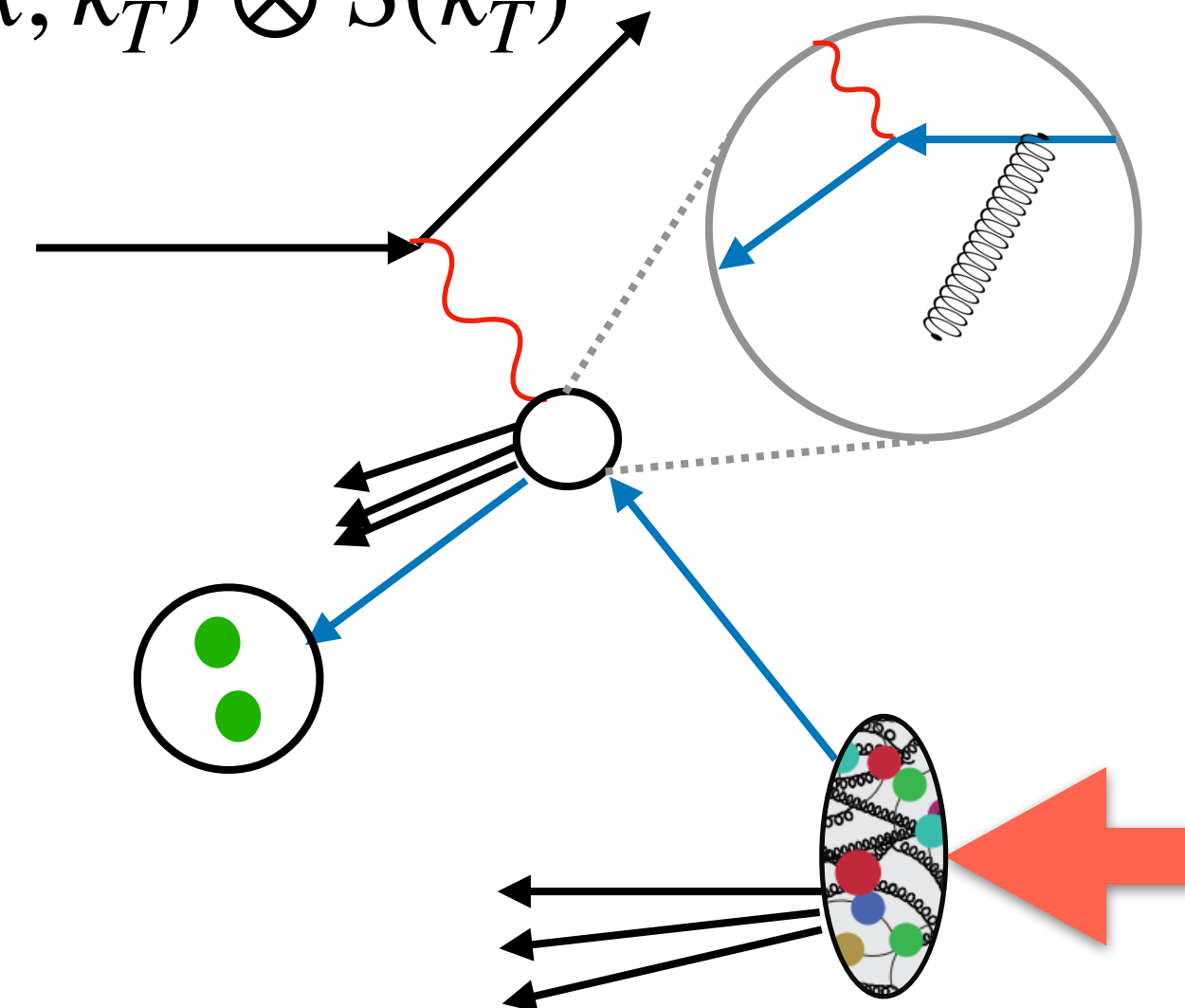


HERMES collaboration, 2009

Flavor separation in SIDIS

- Flavor correlated between initial and final
- Tagging final state hadrons for flavor discrimination
- Need TMD FFs / spin counterparts

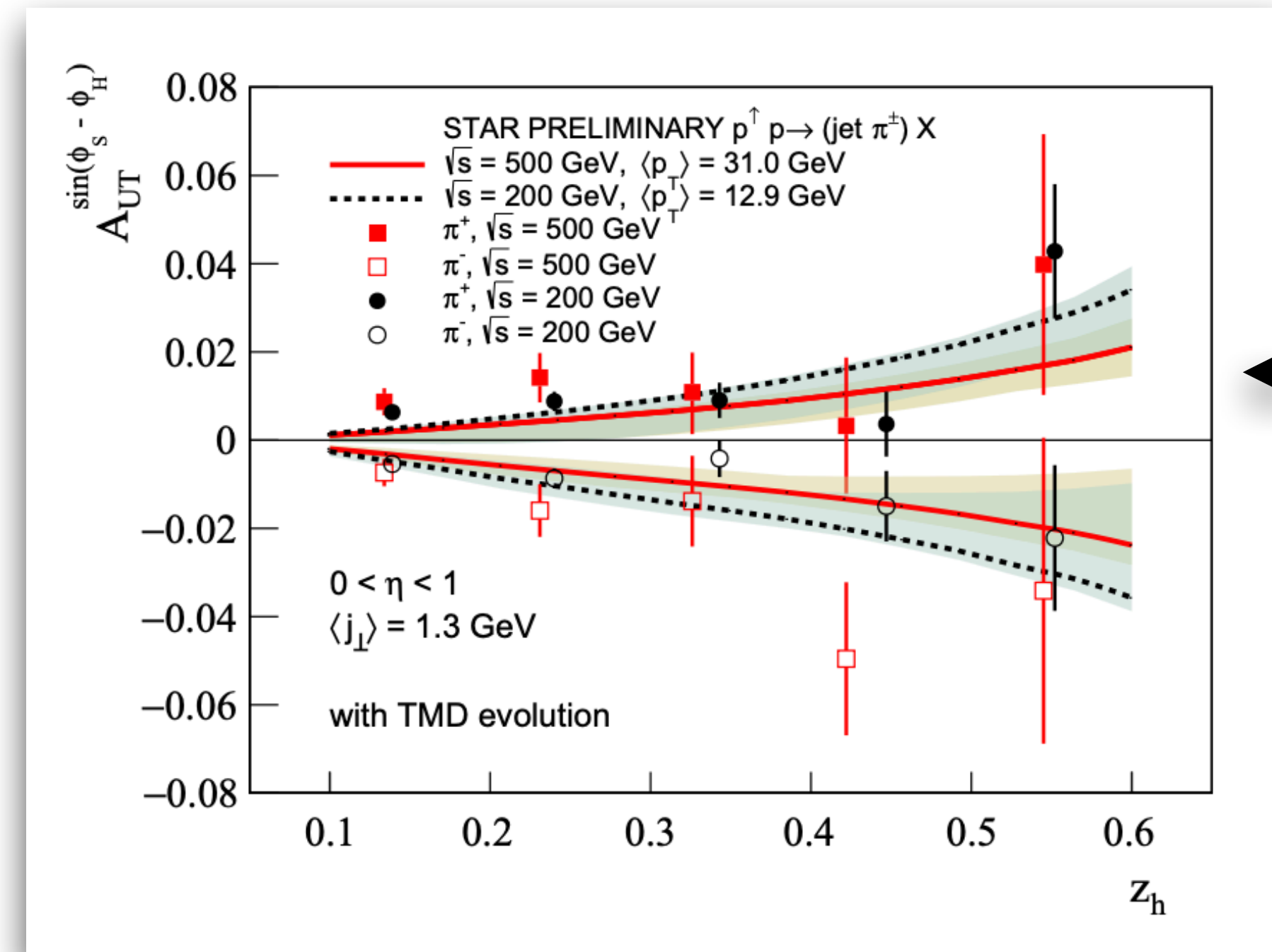
$$\frac{d\sigma}{dq_T} \sim \hat{\sigma}_{ii}(Q) D_{h/i}^{U/T}(z_h, k_T) \otimes f_{i/P}^{U/T}(x, k_T) \otimes S(k_T)$$



Jets @ EIC

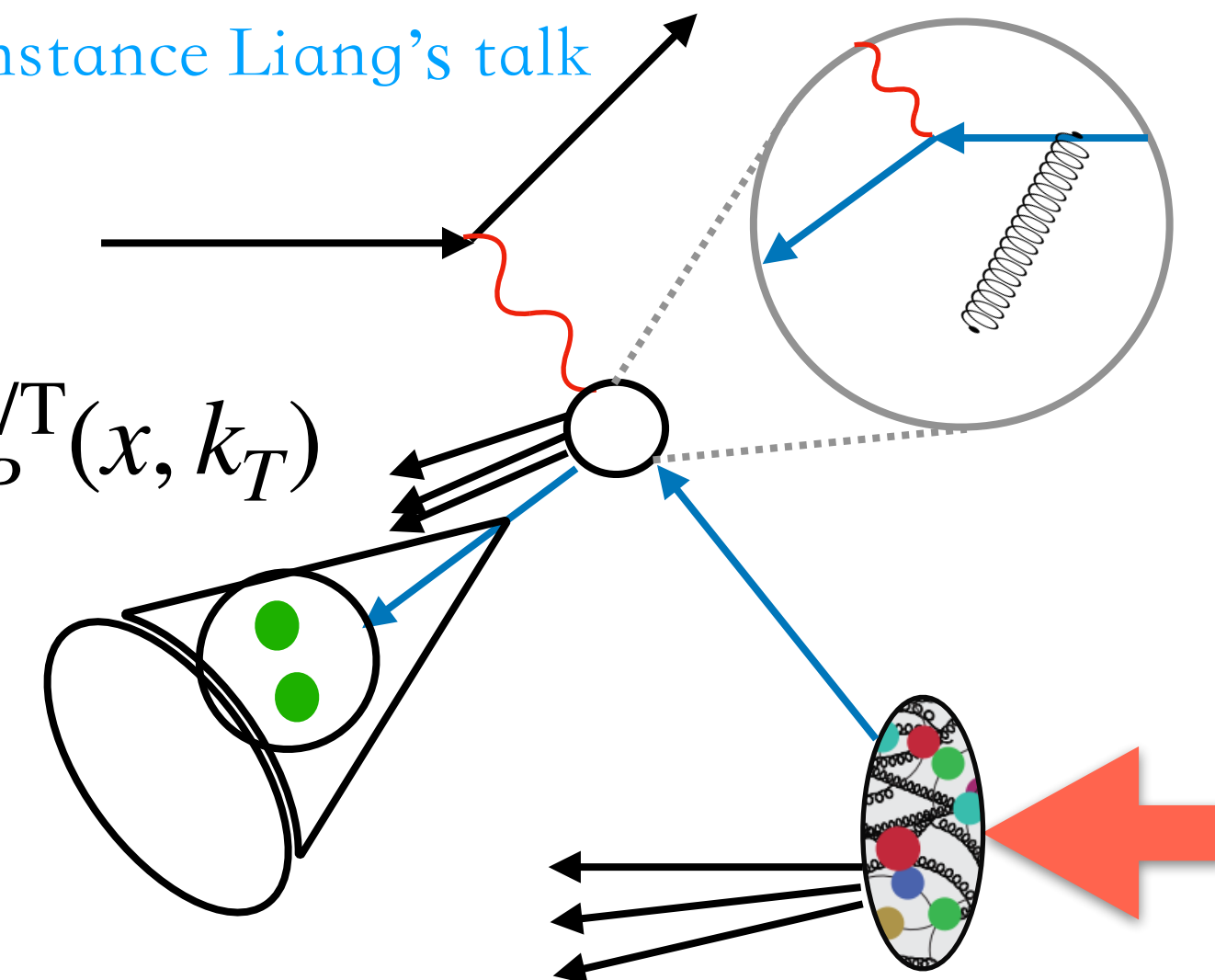
Flavor separation with hadrons in jets

- Tagging hadrons in jet for flavor discrimination
- Applied for the Collins function
- Also good for the flavor separation in Sivers function if small jet-lepton imbalance is measured. [For gluon Sivers, see for instance Liang's talk](#)
- Need FFs



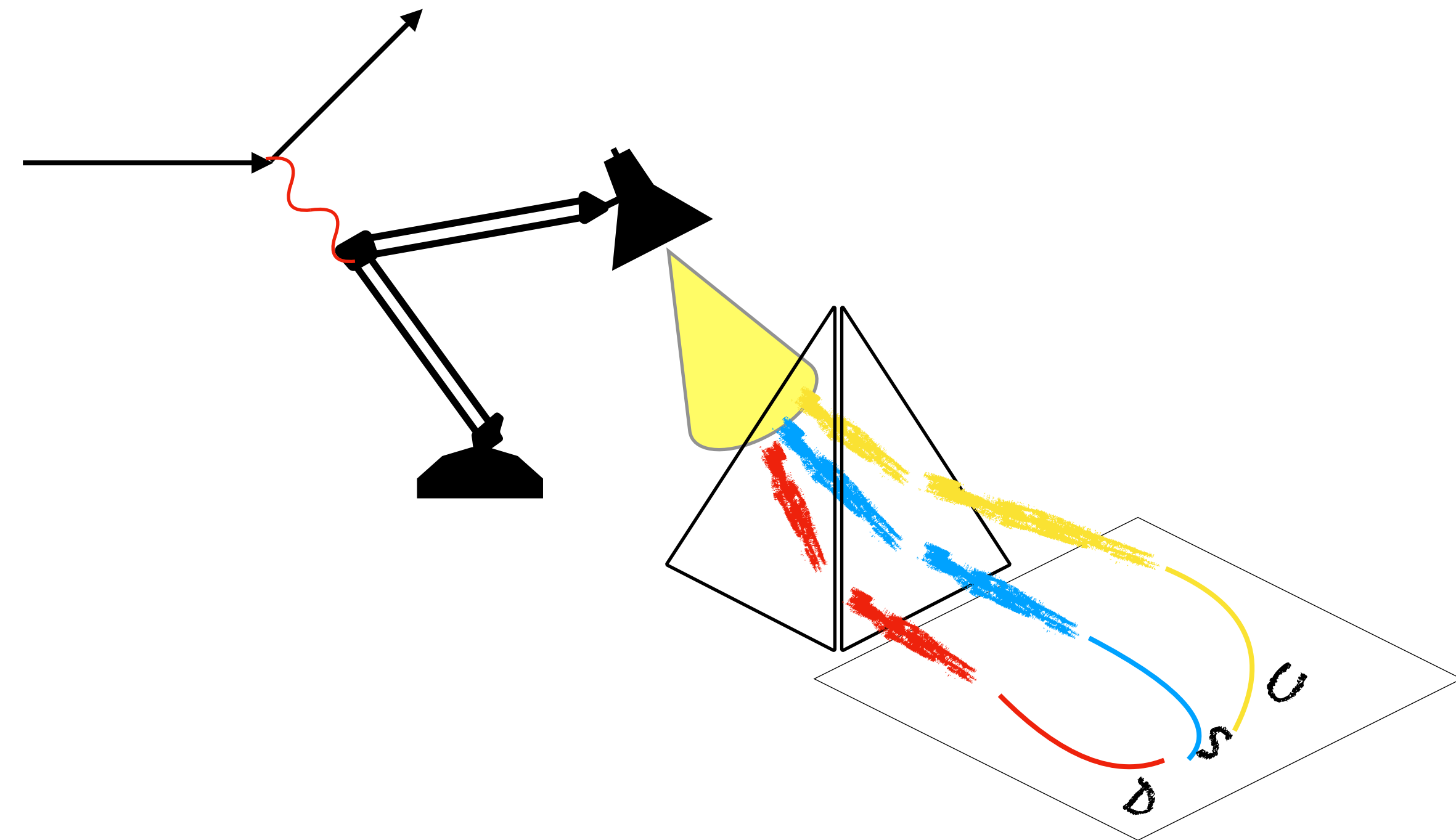
[Kang, Prokudin, Ringer, Yuan, 2017](#)

$$\frac{d\sigma}{dq_T} \sim \hat{\sigma}_{ii}(Q) D_{h/j}(z_h) J_{ji}(R) \otimes f_{i/P}^{U/T}(x, k_T)$$



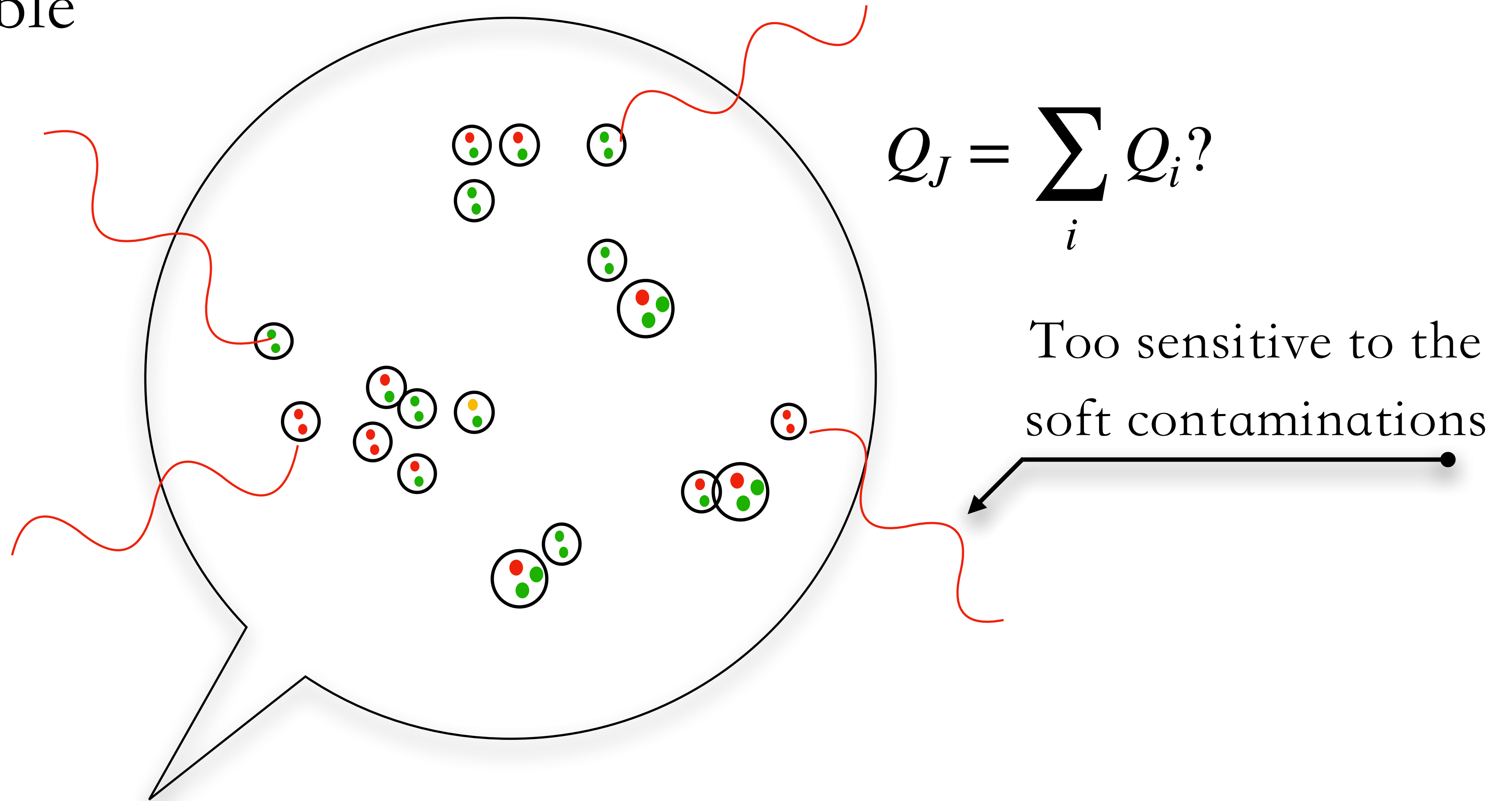
Other alternatives?

- Probe directly the light parton flavors
- No extra non-pert. distributions
- Information beyond QCD
- electro-charge? Robust against the hadronization?



Jet charge as a flavor prism @ EIC

- The observable



Jet charge as a flavor prism @ EIC

- The observable

$$Q_{\kappa} = \sum_i \left(\frac{p_{i,T}}{p_J} \right) Q_i$$

hadron transverse momentum

sum over hadrons within the jet. Flexibility to use a specific hadron species, “the hadron component” of the jet

jet transverse momentum

parameter to suppress contaminations $\kappa > 0$

hadron charge

sum over hadrons within the jet. Flexibility to use a specific hadron species, “the hadron component” of the jet

Field and R. Feynman, 1978

Krohn, Schwartz, Lin, and Waalewijn, 2013

Jet charge as a flavor prism @ EIC

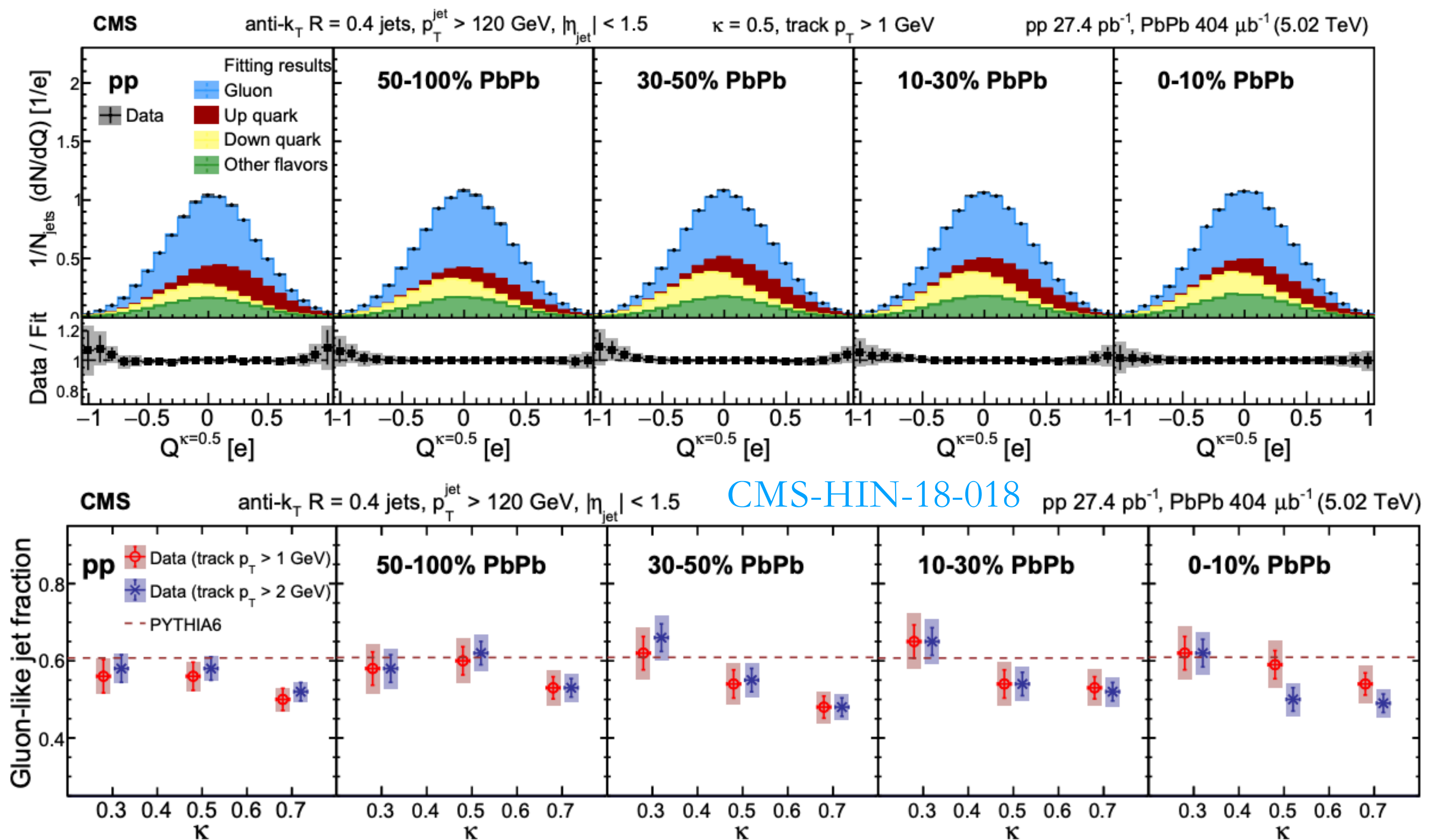
- The observable

hadron transverse
momentum

$$Q_K$$

sum over hadrons within
the jet

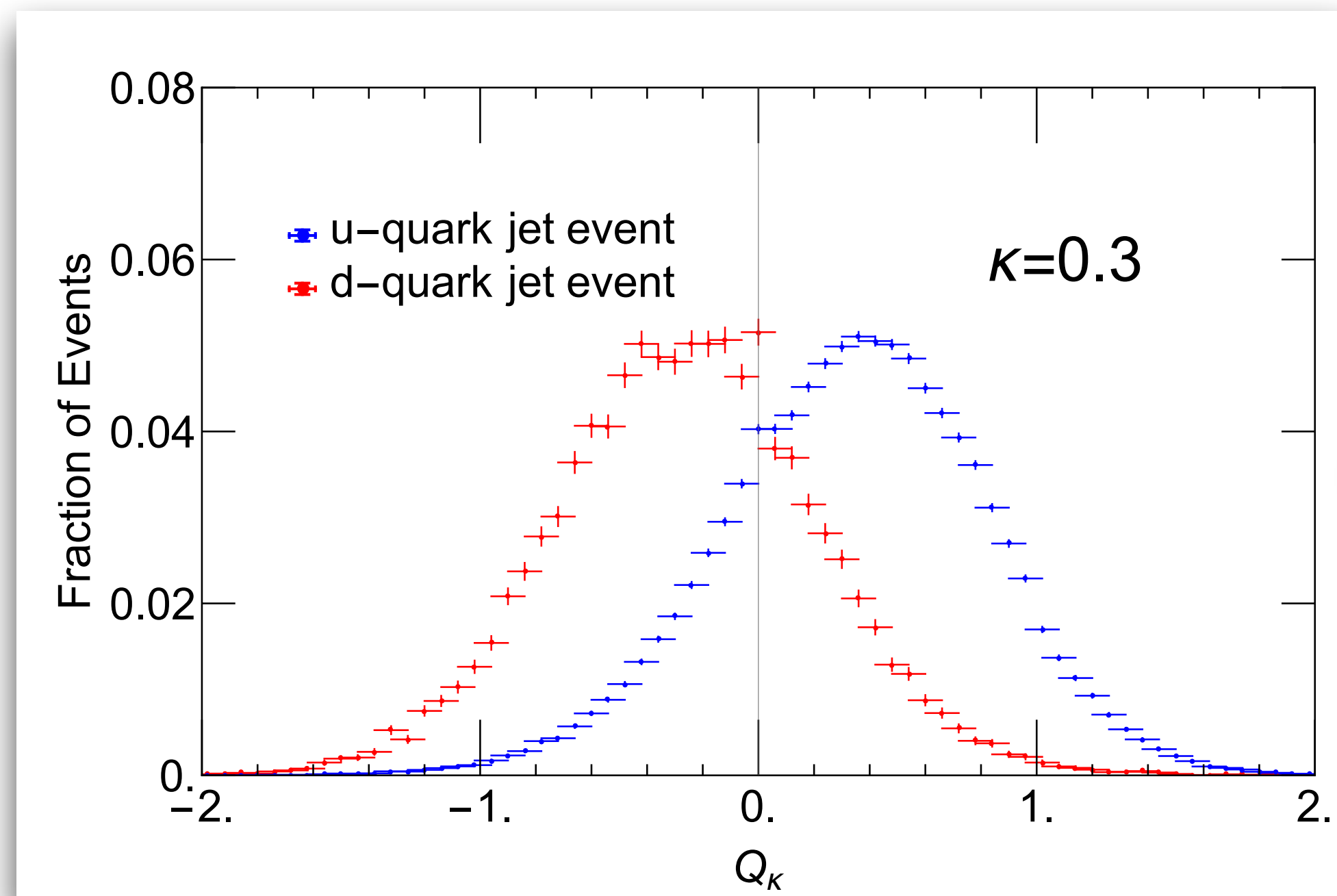
practical applications now



But mostly focus on the quark-gluon discrimination

Jet charge as a flavor prism @ EIC

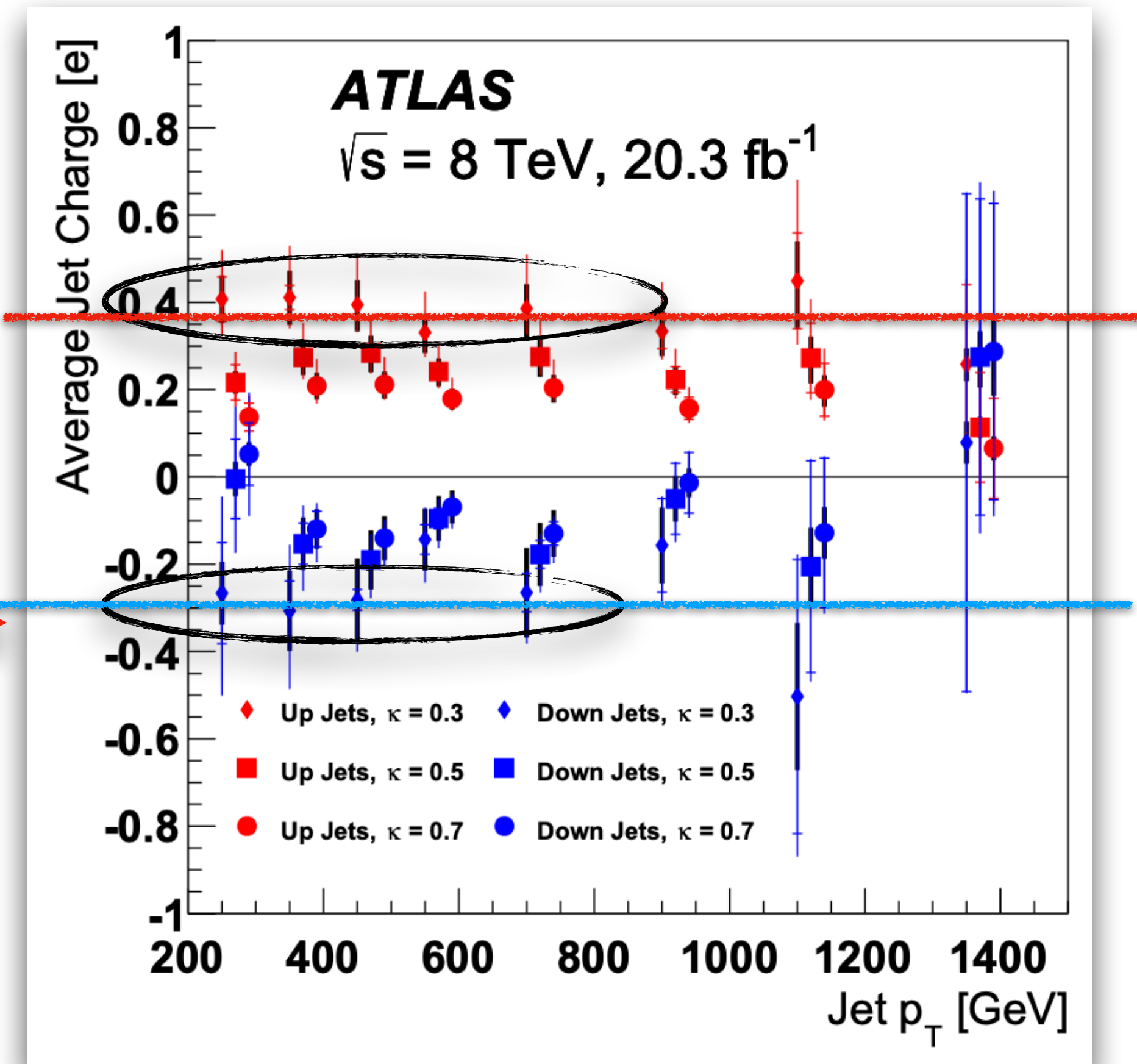
- The observable $Q_\kappa = \sum_i \left(\frac{p_{i,T}}{p_J} \right)^\kappa Q_i$



Pythia simulation

Consistent!

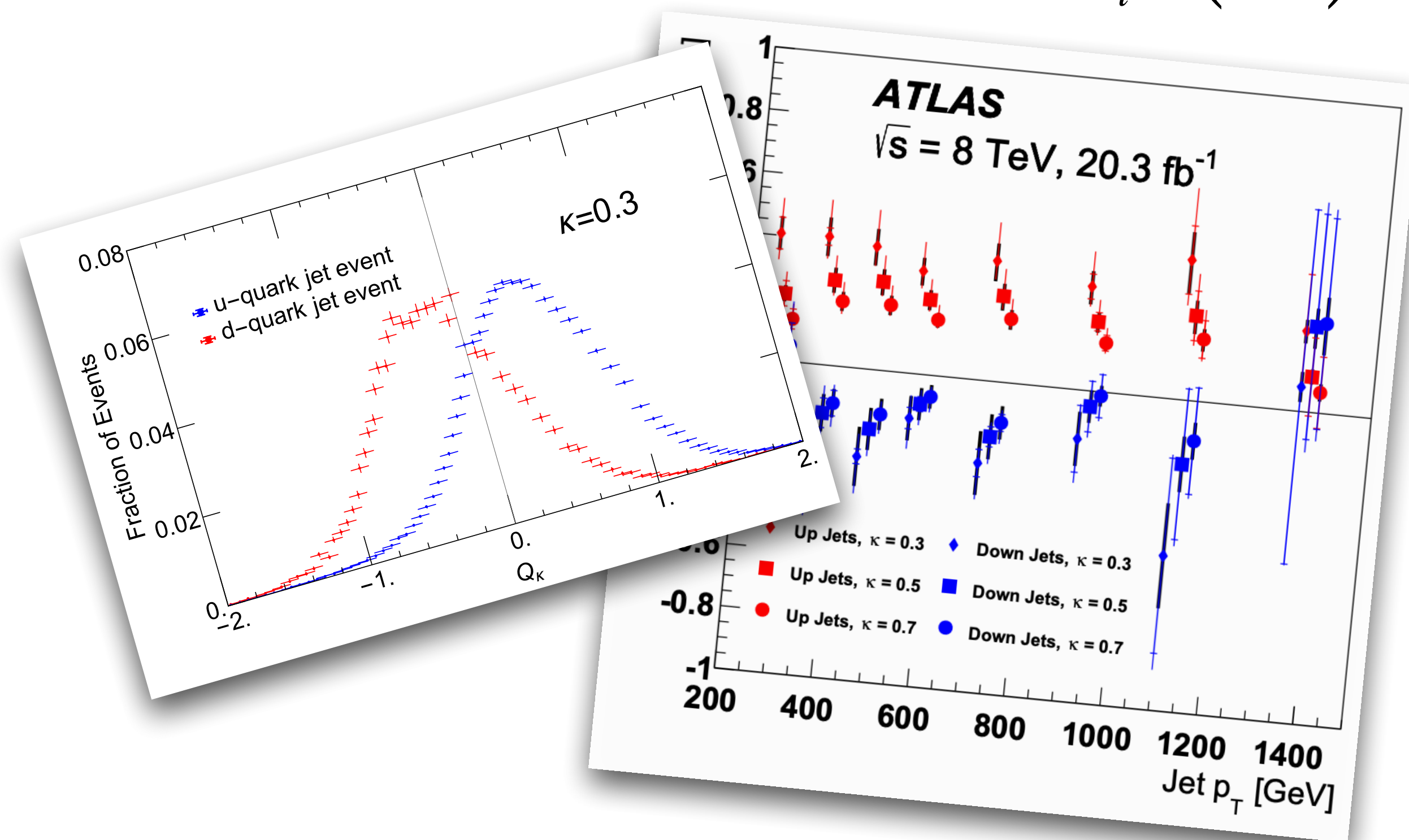
- $\kappa = 0.3$ seems to have the largest separation
- (Almost) independent of the jet p_T



The ATLAS Collaboration, 2015

Jet charge as a flavor prism @ EIC

- The observable $Q_\kappa = \sum_i \left(\frac{p_{i,T}}{p_J} \right)^\kappa Q_i$



- The sign of the jet charge is largely inherited from the initiating parton, robust against the hadronization,
- Very good u- and d-quark separation by dividing events into different charge bins

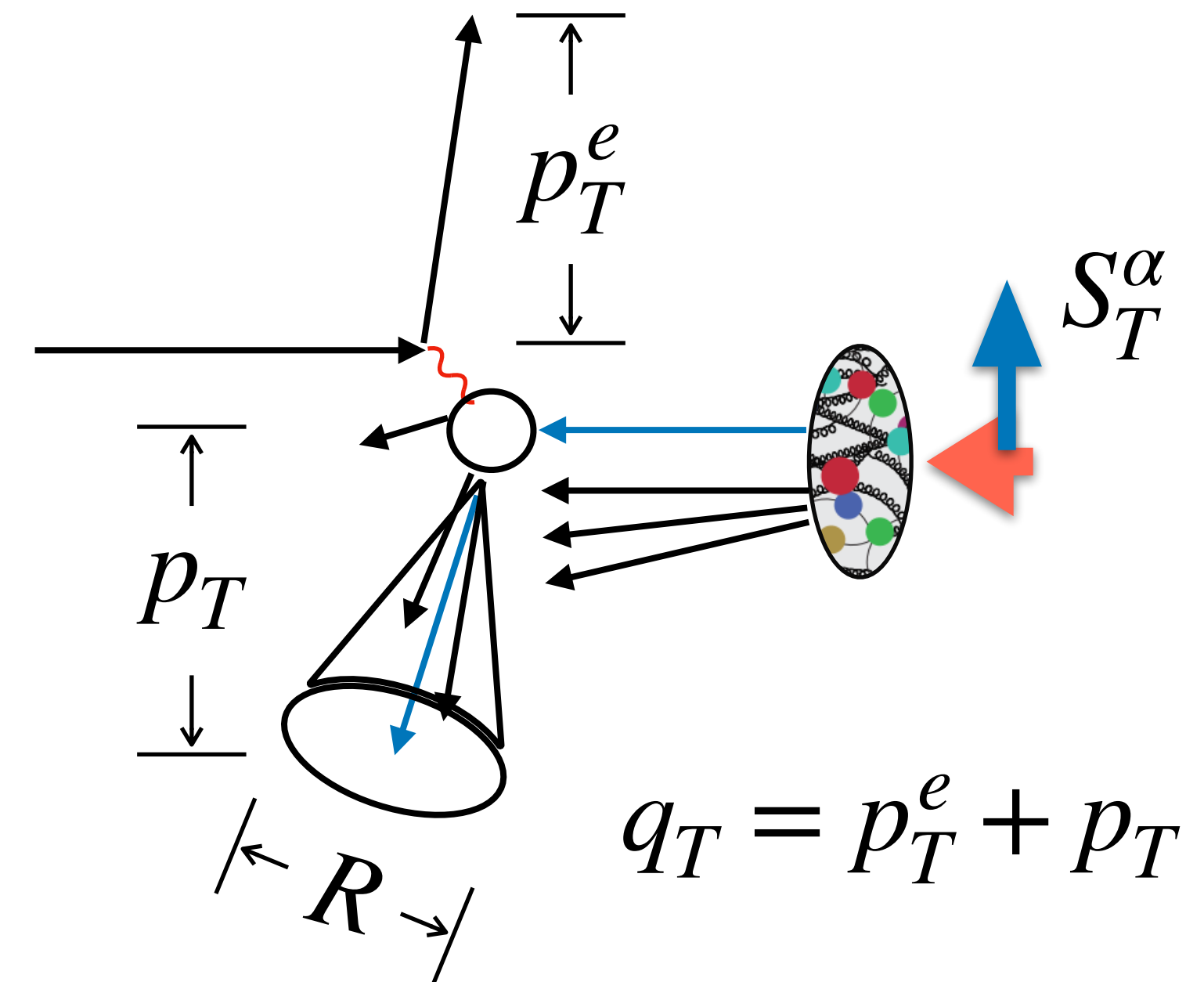
Jet charge as a flavor prism @ EIC

- The factorization [Kang, XL, Mantry, Shao, 2020](#)

$$\frac{d^5\sigma_{UT}^i(S_\perp)}{dy_e d^2p_T^e d^2q_T} = e_i^2 \sigma_0 \epsilon_{\alpha\beta} S_\perp^\alpha \int \frac{db_T^2}{(2\pi)^2} e^{iq_T \cdot b_T} \tilde{W}_{T,i}^\beta$$

$$\tilde{W}_{T,i}^\beta = H(Q) \tilde{f}_{1T,i}^{\perp,\beta}(x, b_T) S_J(b_T, R) \mathcal{J}_i(p_T R)$$

W/O the jet charge observation [XL, Ringer, Vogelsang, Yuan, 19](#)



Jet charge as a flavor prism @ EIC

- The factorization [Kang, XL, Mantry, Shao, 2020](#)

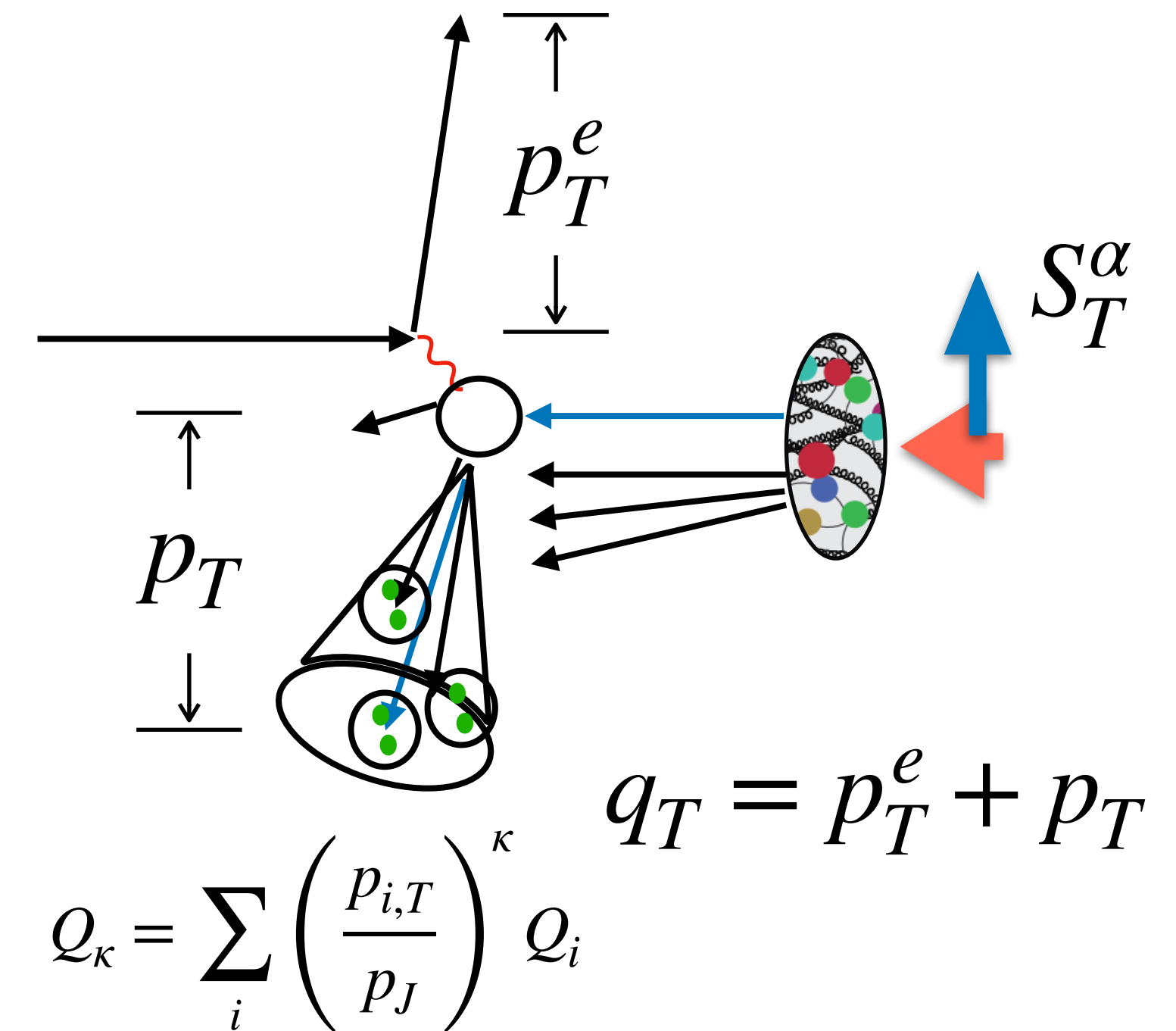
$$\frac{d^6\sigma_{UT}^i(S_\perp)}{dy_e d^2p_T^e d^2q_T dQ_\kappa} = e_i^2 \sigma_0 \epsilon_{\alpha\beta} S_\perp^\alpha \int \frac{db_T^2}{(2\pi)^2} e^{iq_T \cdot b_T} \tilde{V}_{T,i}^\beta$$

$$\tilde{V}_{T,i}^\beta = H(Q) \tilde{f}_{1T,i}^{\perp,\beta}(x, b_T) S_J(b_T, R) \mathcal{G}_i(Q_\kappa, p_T R)$$

[See also Krohn, et. al, 2013 for unpolarized pp](#)

$$\rightarrow \tilde{f}_{1T,i}^{\perp,\beta}(x, b_T) S_J(b_T, R) H(Q) \mathcal{J}_i(p_T R) \frac{\mathcal{G}_i(Q_\kappa, p_T R)}{\mathcal{J}_i(p_T R)}$$

$$\rightarrow W_i \frac{\mathcal{G}_i(Q_\kappa, p_T R)}{\mathcal{J}_i(p_T R)}$$



Jet charge as a flavor prism @ EIC

- The factorization [Kang, XL, Mantry, Shao, 2020](#)

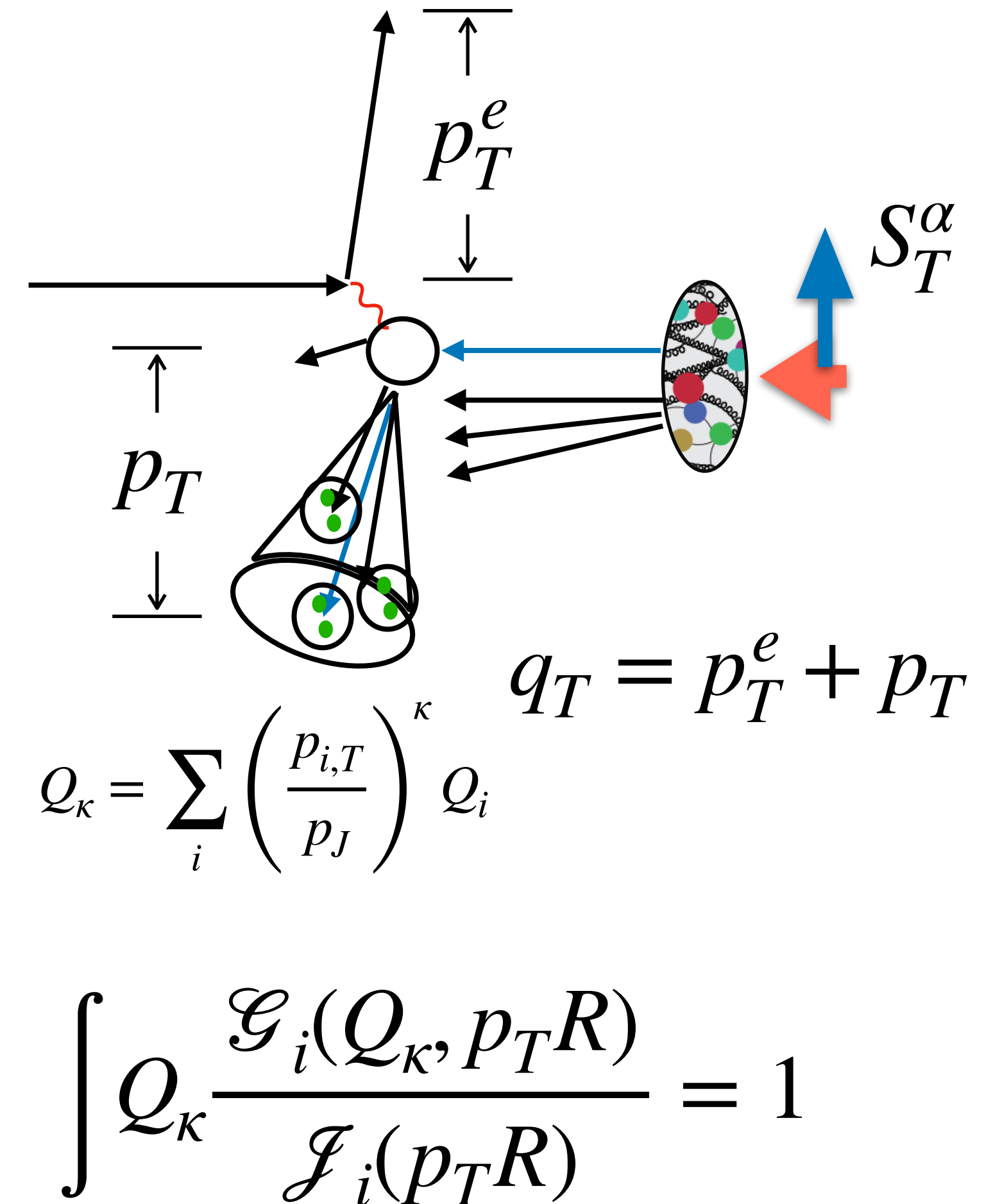
$$\frac{d^6\sigma_{UT}^i(S_\perp)}{dy_e d^2p_T^e d^2q_T dQ_\kappa} = \frac{d^5\sigma_{UT}^i(S_\perp)}{dy_e d^2p_T^e d^2q_T} \frac{\mathcal{G}_i}{\mathcal{J}_i}$$

$$\tilde{V}_{T,i}^\beta = H(Q) \tilde{f}_{1T,i}^{\perp,\beta}(x, b_T) S_J(b_T, R) \mathcal{G}_i(Q_\kappa, p_T R)$$

[See also Krohn, et. al, 2013 for unpolarized pp](#)

$$\rightarrow \tilde{f}_{1T,i}^{\perp,\beta}(x, b_T) S_J(b_T, R) H(Q) \mathcal{J}_i(p_T R) \frac{\mathcal{G}_i(Q_\kappa, p_T R)}{\mathcal{J}_i(p_T R)}$$

$$\rightarrow W_i \frac{\mathcal{G}_i(Q_\kappa, p_T R)}{\mathcal{J}_i(p_T R)}$$



Jet charge as a flavor prism @ EIC

- The factorization [Kang, XL, Mantry, Shao, 2020](#)

$$\frac{d^5\sigma_{UT}(S_\perp)}{dy_e d^2p_T^e d^2q_T} = \sum_{i=u,d,\dots} \int_{Q_\kappa \in \text{bin}} dQ_\kappa \frac{\mathcal{G}_i(Q_\kappa, p_T R)}{\mathcal{J}_i(p_T R)} \frac{d^5\sigma_{UT}^i(S_\perp)}{dy_e d^2p_T^e d^2q_T}$$

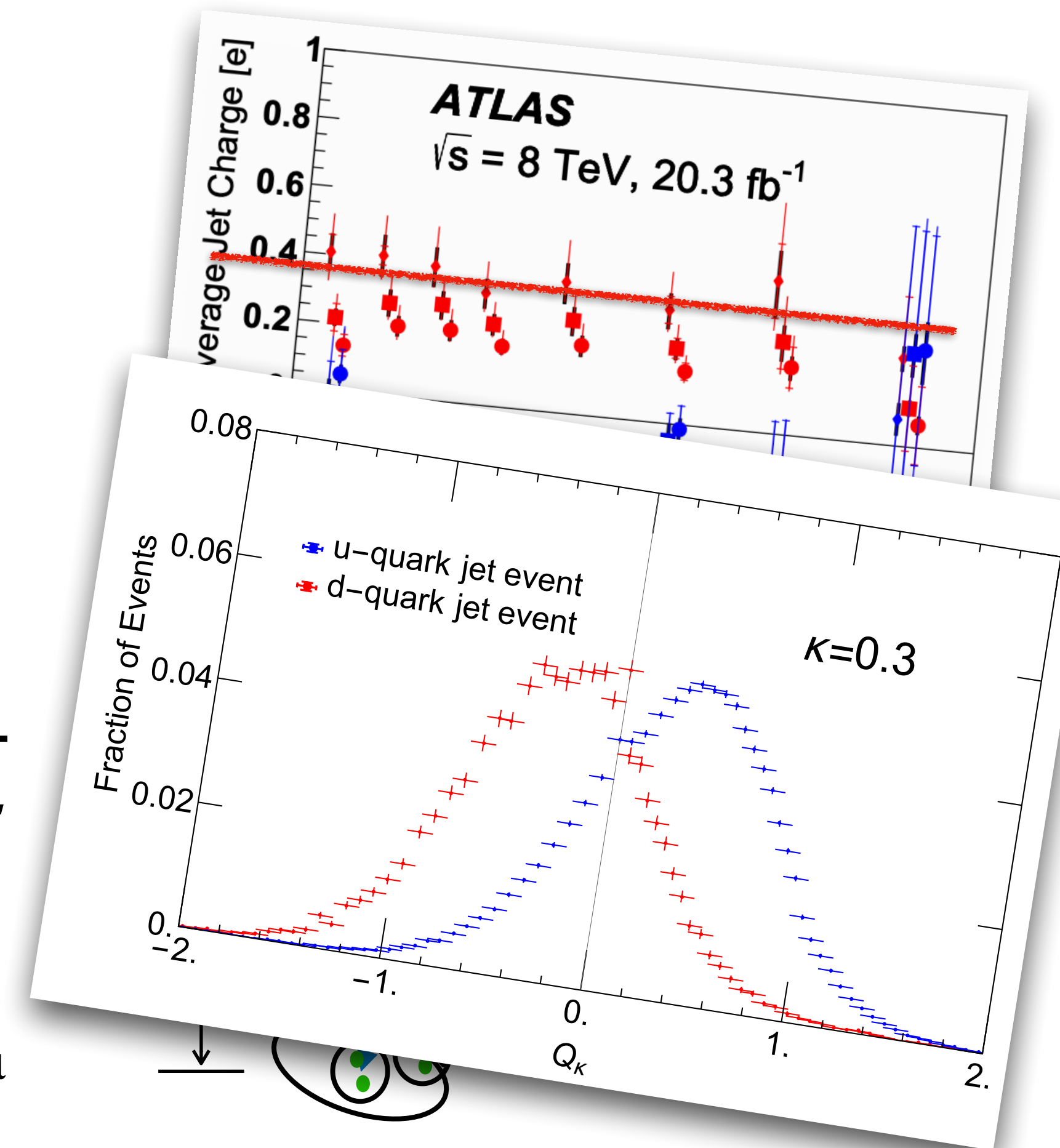
$$= \sum_{i=u,d,\dots} r_i \frac{d^5\sigma_{UT}^i(S_\perp)}{dy_e d^2p_T^e d^2q_T}$$

r_i the event fraction in a given charge bin

$$\int Q_\kappa \frac{\mathcal{G}_i(Q_\kappa, p_T R)}{\mathcal{J}_i(p_T R)} = 1 \quad \rightarrow \quad \sum_{\text{bin}} r_i = 1$$

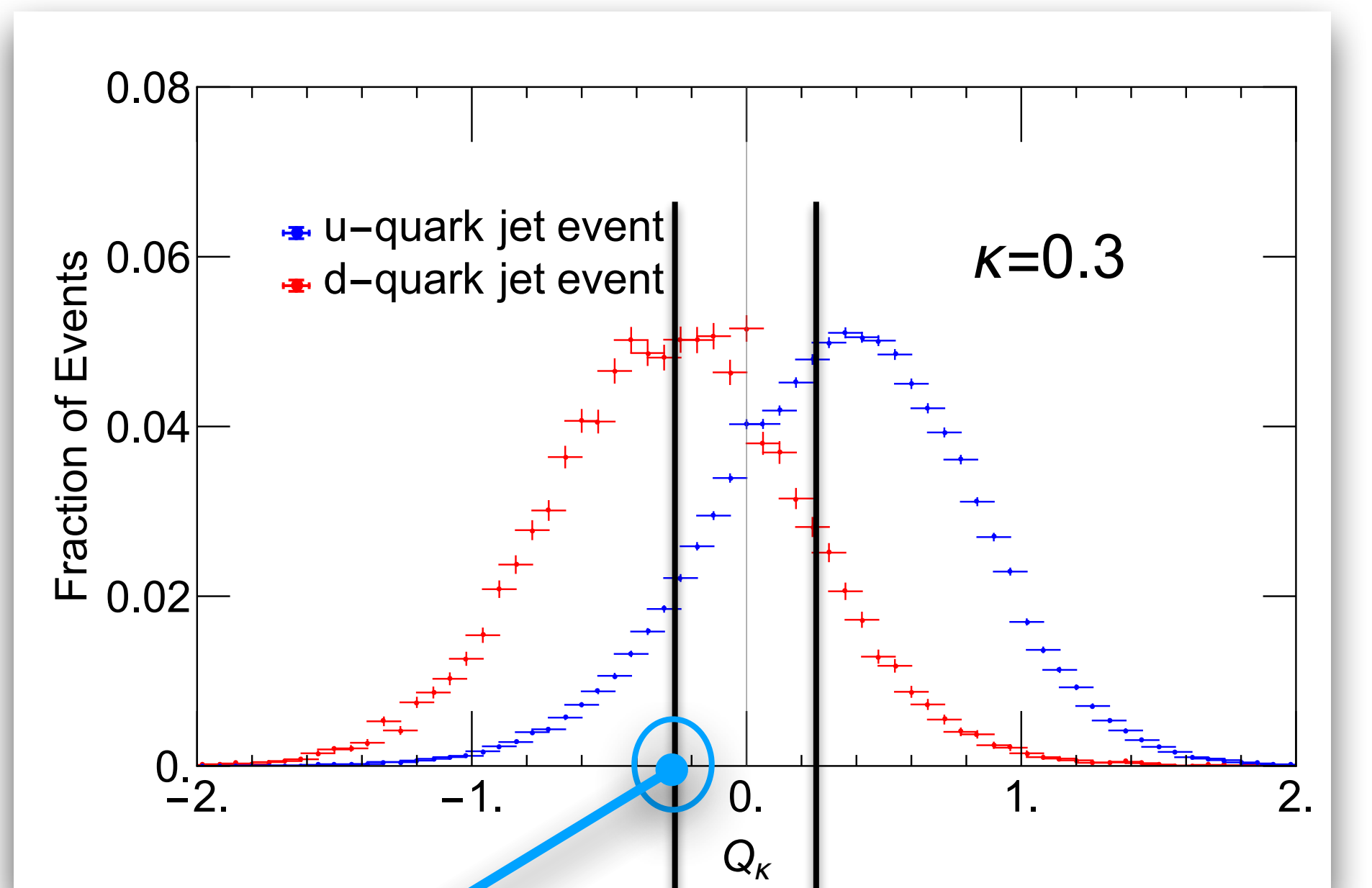
By choosing the charge bin, we can control the sensitivity to different quark flavors

- A non-perturbative number
- Scale independent
- mild $p_T R$ dependence via loops, small for small κ
- Universal, similar to the NRQCD quarkonium matrix element; with flavor symmetry



Jet charge as a flavor prism @ EIC

- As a flavor prism [Kang, XL, Mantry, Shao, 2020, see also STAR collaboration](#)



~ -0.25 ‘-’- bin ‘0’- bin ‘+’- bin

	u	\bar{u}	d	\bar{d}	s	\bar{s}
$r_{i,+}$	0.52	0.17	0.15	0.53	0.30	0.34
$r_{i,-}$	0.15	0.49	0.52	0.15	0.36	0.32
$r_{i,0}$	0.33	0.35	0.33	0.32	0.35	0.34

$$Q_\kappa = \sum_i \left(\frac{p_{i,T}}{p_J} \right)^\kappa Q_i$$

choose $\kappa = 0.3$, large separation

choose π 's in this study
“ π ”-jet

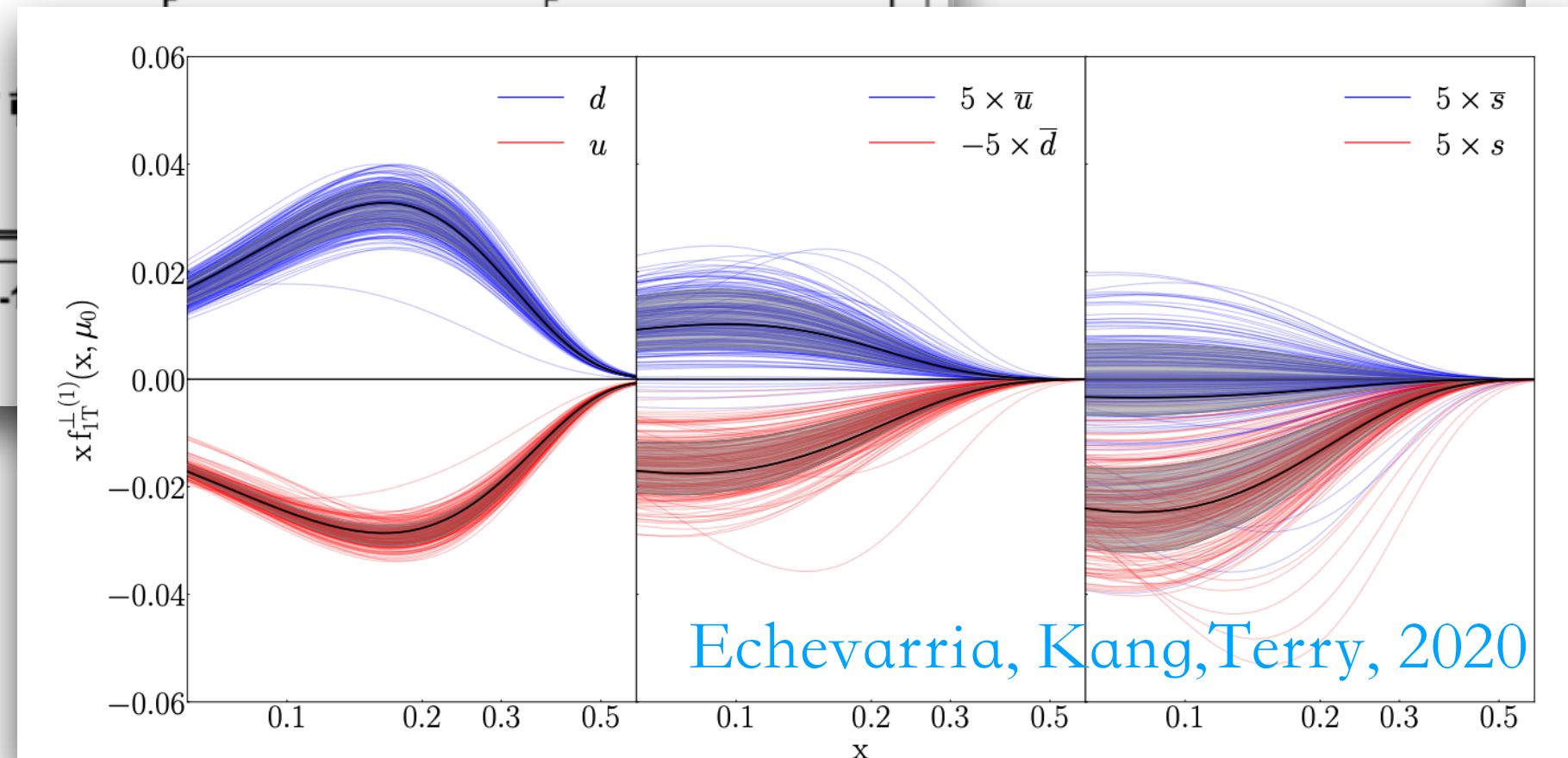
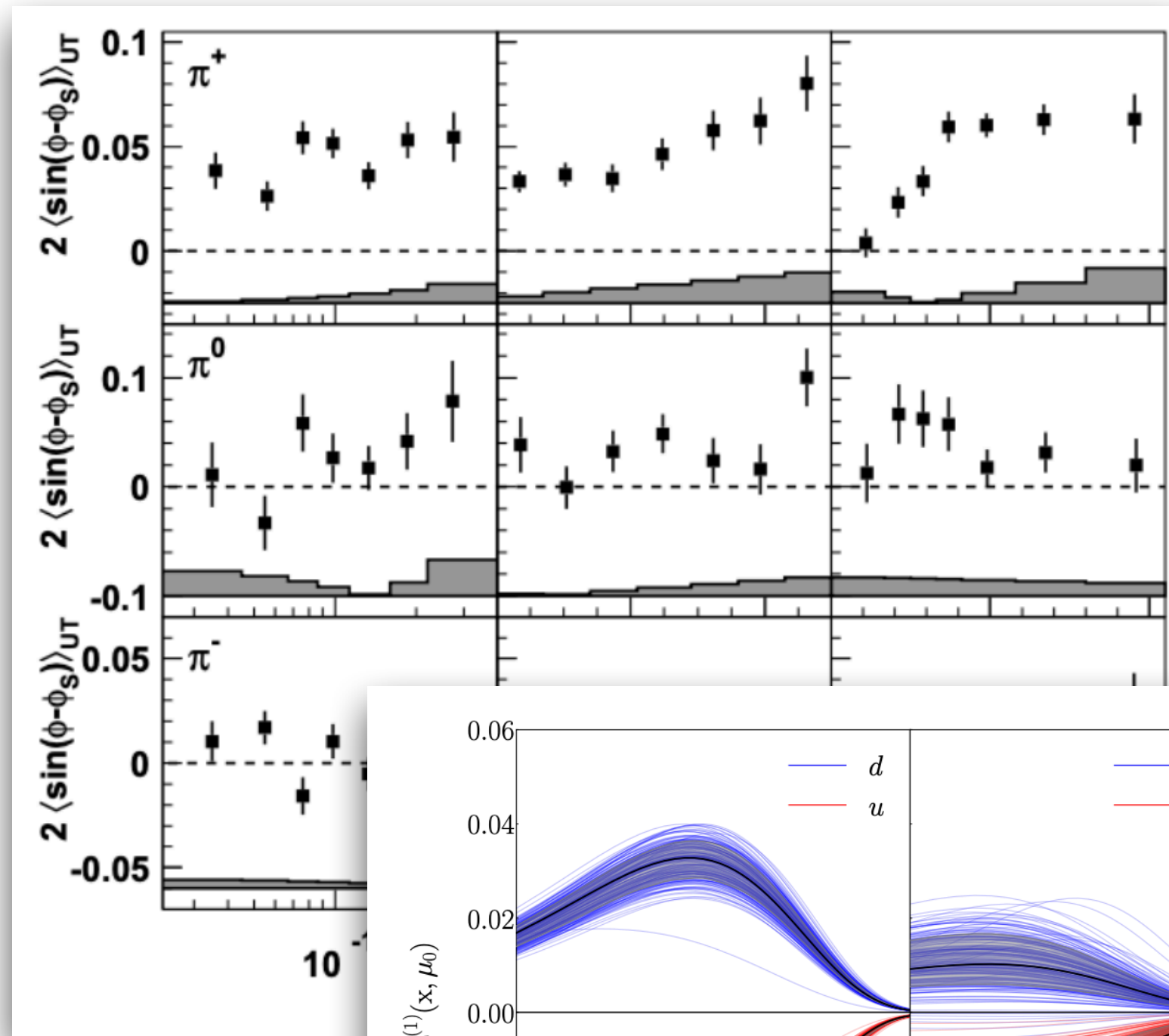
- Obvious flavor symmetries
- Expect enhancement of the d quark in the --bin
- Expect enhancement of the s quark if the kaon component used

Jet charge as a flavor prism @ EIC

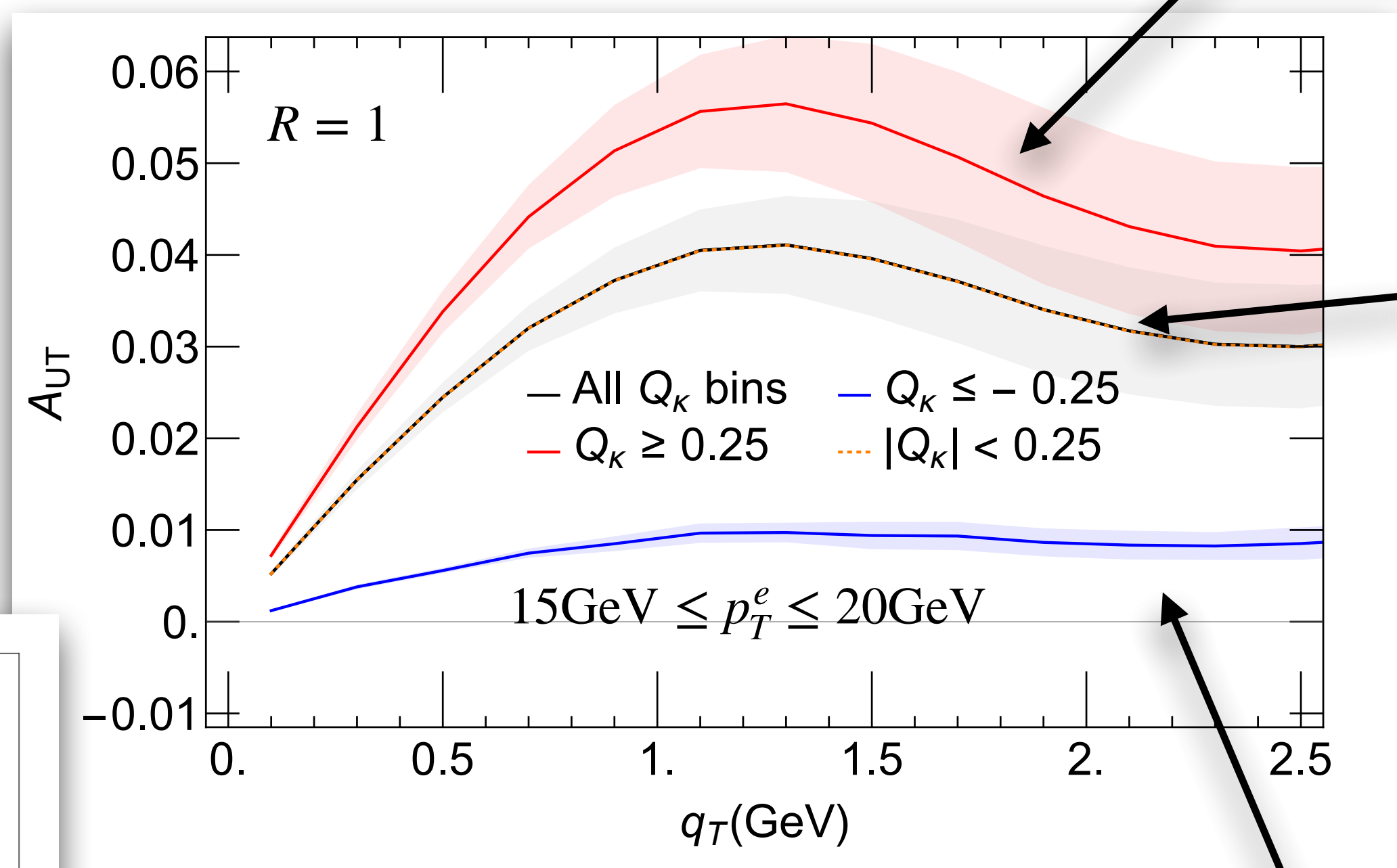
- As a flavor prism

Kang, XL, Mantry, Shao, 2020, see also STAR collaboration

$\sqrt{s} = 105\text{GeV}$,
 $Q^2 \geq 10\text{GeV}^2$,
 $0.1 \leq y \leq 0.9$



Echevarria, Kang, Terry, 2020



u -quark enhanced

u -quark
dominant
due to its
charge

d -quark
dominant

Cancellation
between u and d

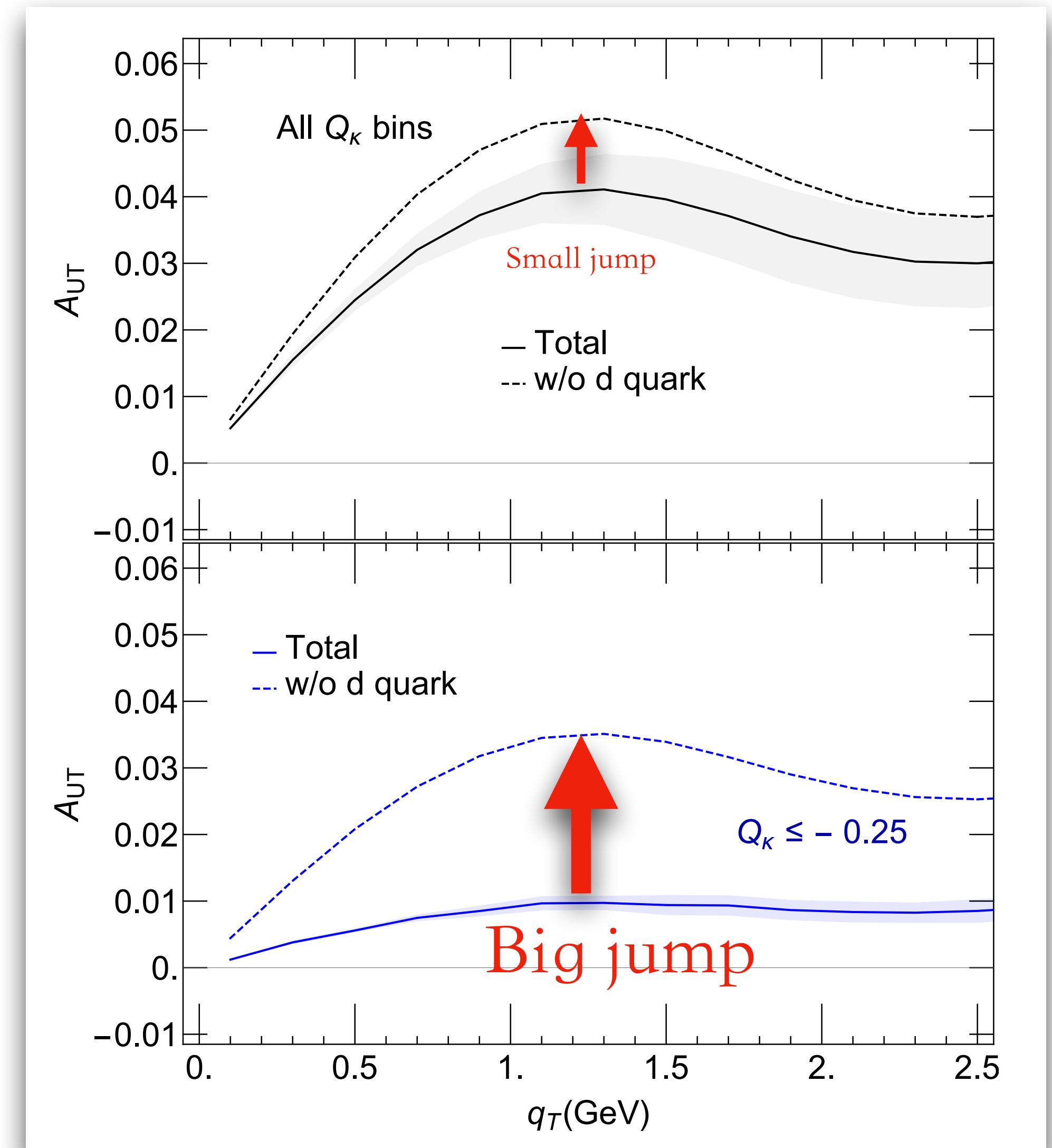
Very similar behavior as
SIDIS, map out the pion
component of the jet

Jet charge as a flavor prism @ EIC

- As a flavor prism

Kang, XL, Mantry, Shao, 2020, see also STAR collaboration

Dramatically improved
sensitivity to the d quark
in the negative charge bin



Conclusions

- The jet charge is sensitive to the parton species
- Maps out flavors; maps a specific hadron component of the jet
- Demonstrate its power in the flavor separation
- Expect to have applications in all hadron structure studies at EIC and other colliders

Happy HOLIDAYS