# Implementation of the Goloskokov-Kroll model pseudoscalar meson production in PARTONS

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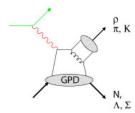
## Chiral-odd GPDs

ullet Four chiral-odd GPDs  $H_T, \tilde{H}_T, E_T, \tilde{E}_T$  at leading twist

$$\begin{split} &\frac{1}{2} \int \frac{dz^{-}}{2\pi} \, e^{ixP^{+}z^{-}} \langle p', \lambda' | \, \bar{\psi}(-\frac{1}{2}z) \, i\sigma^{+i} \, \psi(\frac{1}{2}z) \, | p, \lambda \rangle \Big|_{z^{+}=0, \, \mathbf{z}_{T}=0} \\ &= \, \frac{1}{2P^{+}} \bar{u}(p', \lambda') \left[ H_{T}^{q} \, i\sigma^{+i} + \tilde{H}_{T}^{q} \, \frac{P^{+}\Delta^{i} - \Delta^{+}P^{i}}{m^{2}} \right. \\ &\quad + E_{T}^{q} \, \frac{\gamma^{+}\Delta^{i} - \Delta^{+}\gamma^{i}}{2m} + \tilde{E}_{T}^{q} \, \frac{\gamma^{+}P^{i} - P^{+}\gamma^{i}}{m} \right] u(p, \lambda). \end{split}$$

where i = 1, 2 is the transversity index [Diehl '03]

Accessible through exclusive meson production processes



#### Chiral-odd GPDs

- Factorization for electroproduction of mesons, only for longitidunally polarized photons, has been proven [Collins-Frankfurt-Strikman '97]
- ullet For transversely polarized photons, cross section is power suppressed by 1/Q [Collins-Frankfurt-Strikman '97]
- But in some kinematics, it is apparent that transversely polarized photons contributes substantially [HERMES Collaboration, arXiv:0907.2596[hep-ex]]

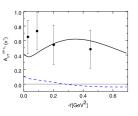


Fig. 2 (Color online) The  $\sin \phi_s$  moment for a transversely polarized target at  $Q^2 \simeq 2.45$  GeV<sup>2</sup> and W = 3.99 GeV. The prediction from our handbag approach is shown as a *solid line*. The *dashed line* is obtained disregarding the twist-3 contribution. Data are taken from [10]

Figure: [Goloskokov-Kroll '10]

- Goloskokov-Kroll(GK) model for pseudoscalar meson production considers the region of small  $\xi$  and small -t, but large  $Q^2$  and W. [Goloskokov-Kroll '10]
- Contributions from transversely polarized photons can be computed as a twist-3 effect in the handbag mechanism
- In pseudoscalar meson production, the following amplitudes are relevant

$$egin{aligned} \mathcal{M}_{0+,0+} &= \sqrt{1-\xi^2} rac{e}{Q} [\langle ilde{H} 
angle - rac{\xi^2}{1-\xi^2} \langle ilde{E} 
angle ] \ \mathcal{M}_{0-,0+} &= rac{e}{Q} rac{-t'}{2m} [\xi \langle ilde{E} 
angle ] ] \ \mathcal{M}_{0-,++} &= \sqrt{1-\xi^2} e \, \langle H_T 
angle \ \mathcal{M}_{0+,\mu+} &= -rac{e}{4m} \sqrt{-t'} \langle ar{E}_T 
angle \end{aligned}$$

• Generically,  $\langle F \rangle$  represents a convolution of a GPD F with an appropriate subprocess amplitude

$$\langle F 
angle = \sum_{\lambda} \int_{-1}^{1} dx \, \mathcal{H}_{\mu'\lambda,\mu\lambda}(x,\xi,Q^2) \, F(x,\xi,t)$$

where  $\lambda$  denotes unobserved helicities of the partons.

- Subprocesses are calculated in the so-called modified perturbative approach: Transverse momenta of the quark and the anti-quark in the meson are kept and gluon radiations are taken into account through Sudakov factor
- GPDs appear in the following combination

$$F^{0}(x,\xi,t) = \frac{1}{\sqrt{2}} \Big( e_{u} F^{u}(x,\xi,t) - e_{d} F^{d}(x,\xi,t) \Big)$$

$$F^{+}(x,\xi,t) = F^{u}(x,\xi,t) - F^{d}(x,\xi,t)$$

In impact space

$$\mathcal{H}_{\pi} = \int d\tau d^2\vec{b} \; \hat{\Psi}_{\pi}(\tau, -\vec{b}) \hat{\mathcal{F}}_{\pi}^i(\bar{x}, \xi, \tau, Q^2, \vec{b}) \alpha_s(\mu_R) \; exp\big(-S(\tau, \vec{b}, Q^2)\big)$$

Hard scattering kernels has the following forms in momentum space

$$\begin{split} \mathcal{F}^{q}_{\pi^{0}} &= \frac{N_{c}^{2} - 1}{2N_{c}} \sqrt{\frac{2}{N_{c}}} \frac{Q}{\xi} \Big[ \frac{1}{k_{\perp}^{2} + \tau(\bar{x} + \xi)Q^{2}/(2\xi) - i\epsilon} - \frac{1}{k_{\perp}^{2} - \bar{\tau}(\bar{x} - \xi)Q^{2}/(2\xi) - i\epsilon} \Big] \\ \mathcal{F}^{q}_{\pi^{+}} &= \frac{N_{c}^{2} - 1}{2N_{c}} \sqrt{\frac{2}{N_{c}}} \frac{Q}{\xi} \Big[ \frac{e_{d}}{k_{\perp}^{2} + \tau(\bar{x} + \xi)Q^{2}/(2\xi) - i\epsilon} - \frac{e_{u}}{k_{\perp}^{2} - \bar{\tau}(\bar{x} - \xi)Q^{2}/(2\xi) - i\epsilon} \Big] \end{split}$$

• A Gaussian meson wave function is used at twist-2

$$\Psi_{\pi}( au,ec{b}) \sim au(1- au) exp \Big[rac{ au( au-1)}{4}rac{ec{b}^2}{ extstar{a}_{\pi}^2}\Big]$$

Sudakov factor has the form

$$S( au,b,Q) = s( au,b,Q) + s(ar{ au},b,Q) - rac{4}{eta_0} \, \, ext{In} \, rac{ ext{In}(\mu_R/\Lambda_{QCD})}{\hat{b}}$$

where

$$s(\tau, b, Q) = \frac{8}{3\beta_0} \left( \hat{q} \ln \left( \frac{\hat{q}}{\hat{b}} \right) - \hat{q} + \hat{b} \right) + NLL$$
$$\hat{b} = -\ln(b \Lambda_{QCD})$$
$$\hat{q} = \ln(\tau Q / (\sqrt{2}\Lambda_{QCD}))$$

Twist-3 meson wave function

$$\Psi_{\pi}( au,ec{b})\sim exp\Big[-rac{ec{b}^2}{8a_{\pi}^2}\Big]\,\mathit{I}_0(rac{ec{b}^2}{8a_{\pi}^2})$$

## Goloskokov-Kroll Model GPDs

GPDs are constructed from double distribution ansatz

$$F_i^{\mathsf{a}}(ar{x},\xi,t) = \int_{-1}^1 d
ho \int_{-1+|
ho|}^{1-|
ho|} d\eta \ \delta(
ho + \xi \eta - ar{x}) f_i^{\mathsf{a}}(
ho,\eta,t)$$

where for valence-quark GPDs;

$$f_i(\rho, \eta, t) = \exp[(b_i - \alpha_i' \ln \rho)t] F_i^a(\rho, \xi = t = 0) \frac{3}{4} \frac{(1 - \rho)^2 - \eta^2}{(1 - \rho)^3} \Theta(\rho)$$

Parameters of the forward limits

 $H: \mathsf{DSSV}$ , Phys. Rev. D **80**, 034030 (2009)

 $H_T$ : ABM, Phys. Rev. D 86, 054009 (2012) and DSSV, Phys. Rev. D 80, 034030 (2009)

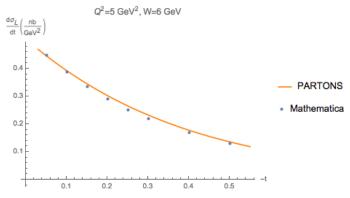
 $\tilde{E}$ : LHPC Collaboration, Phys. Rev. D 77, 094502 (2008)

 $\bar{E}_T$ : QCDSF and UKQCD Collaborations, Phys. Rev. Lett. 98, 222001 (2007)

- So altogether, to compute the longitudinal cross section, we need: Twist-2 meson wave function, kernel, running coupling, Sudakov factor and GPDs  $\tilde{H}$  and  $\tilde{E}$ .
- To compute the transverse cross section, we need: Twist-3 meson wave function, kernel, running coupling, Sudakov factor and GPDs  $H_T$  and  $\bar{E}_T$ .
- 3 dimensional integrals, over  $\bar{x}, \tau$  and b, are performed in impact space
- In many different processes, meson wavefunction has the same structure.
- Sudakov factor has also the same structure
- $\pi^+$  electroproduction also receives a pion pole contribution, besides the handbag contribution

# PARTONS vs. Mathmematica

ullet Comparision between the PARTONS and Mathmematica codes (slightly different GPDs parametrization for  $ilde{H}$ )



# PARTONS vs. Mathmematica

• Comparision between the PARTONS and Mathmematica codes at amplitude level (the same GPDs parametrization).

$$Q^2 = 5 GeV^2$$
,  $W = 6 GeV$ 

$$M_{0p0p}|_{t=-0.2} = 0.173902 - 0.150831i \rightarrow 0.1737324 - 0.1509104i$$
 $M_{0m0p}|_{t=-0.2} = 0.00422273 - 0.0502173i \rightarrow 0.0041224 - 0.0502987i$ 
 $M_{0p0p}|_{t=-0.4} = 0.136403 - 0.108212i \rightarrow 0.1362792 - 0.1082896i$ 
 $M_{0m0p}|_{t=-0.4} = 0.00742887 - 0.0477006i \rightarrow 0.0073482 - 0.0477759i$