# <span id="page-0-0"></span>Implementation of the Goloskokov-Kroll model pseudoscalar meson production in PARTONS

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### Chiral-odd GPDs

Four chiral-odd GPDs  $H_{T}, \tilde{H}_{T}, E_{T}, \tilde{E}_{T}$  at leading twist  $\frac{1}{2}\int\frac{dz^-}{2\pi}\,e^{ixP^+z^-}\langle p',\lambda'|\,\bar{\psi}(-\tfrac{1}{2}z)\,i\sigma^{+i}\,\psi(\tfrac{1}{2}z)\,\left|p,\lambda\right\rangle\Big|_{z^+=0,\,{\bf z}_T=0}$  $\hspace{.2cm} = \hspace{.2cm} \frac{1}{2P^+} \bar{u}(p', \lambda') \left[ H^{q}_T \, i \sigma^{+i} + \tilde{H}^{q}_T \, \frac{P^+ \Delta^i - \Delta^+ P^i}{m^2} \right]$  $+\,E^q_T\,\frac{\gamma^+\Delta^i-\Delta^+\gamma^i}{2m}+\tilde E^q_T\,\frac{\gamma^+P^i-P^+\gamma^i}{m}\Bigg]\,u(p,\lambda).$ 

where  $i = 1, 2$  is the transversity index  $\frac{1}{|D|}$  [Diehl '03]

• Accessible through exclusive meson production processes



## Chiral-odd GPDs

- Factorization for electroproduction of mesons, only for longitidunally polarized photons, has been proven [Collins-Frankfurt-Strikman '97]
- For transversely polarized photons, cross section is power suppressed  $\mathsf{bv}\;1/Q$  [Collins-Frankfurt-Strikman '97]
- But in some kinematics, it is apparent that transversely polarized photons contributes substantially [HERMES Collaboration, arXiv:0907.2596[hep-ex]]



Fig. 2 (Color online) The  $\sin \phi$ , moment for a transversely polarized target at  $Q^2 \simeq 2.45$  GeV<sup>2</sup> and  $W = 3.99$  GeV. The prediction from our handbag approach is shown as a solid line. The dashed line is obtained disregarding the twist-3 contribution. Data are taken from  $[10]$ 

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#### Figure: [Goloskokov-Kroll '10]

### <span id="page-3-0"></span>Goloskokov-Kroll Model

- Goloskokov-Kroll(GK) model for pseudoscalar meson production considers the region of small  $\xi$  and small  $-t$ , but large  $Q^2$  and  $W.$ [Goloskokov-Kroll '10]
- Contributions from transversely polarized photons can be computed as a twist-3 effect in the handbag mechanism
- In pseudoscalar meson production, the following amplitudes are relevant

$$
\mathcal{M}_{0+,0+} = \sqrt{1 - \xi^2} \frac{e}{Q} [\langle \tilde{H} \rangle - \frac{\xi^2}{1 - \xi^2} \langle \tilde{E} \rangle]
$$

$$
\mathcal{M}_{0-,0+} = \frac{e}{Q} \frac{-t'}{2m} [\xi \langle \tilde{E} \rangle]]
$$

$$
\mathcal{M}_{0-,++} = \sqrt{1 - \xi^2} e \langle H_T \rangle
$$

$$
\mathcal{M}_{0+, \mu+} = -\frac{e}{4m} \sqrt{-t'} \langle \bar{E}_T \rangle
$$

### <span id="page-4-0"></span>Goloskokov-Kroll Model

• Generically,  $\langle F \rangle$  represents a convolution of a GPD F with an appropriate subprocess amplitude

$$
\langle F \rangle = \sum_{\lambda} \int_{-1}^{1} dx \, \mathcal{H}_{\mu' \lambda, \mu \lambda}(x, \xi, Q^2) \, F(x, \xi, t)
$$

where  $\lambda$  denotes unobserved helicities of the partons.

- Subprocesses are calculated in the so-called modified perturbative approach: Transverse momenta of the quark and the anti-quark in the meson are kept and gluon radiations are taken into account through Sudakov factor
- **GPDs appear in the following combination**

$$
F^{0}(x,\xi,t) = \frac{1}{\sqrt{2}} \Big( e_{u} F^{u}(x,\xi,t) - e_{d} F^{d}(x,\xi,t) \Big)
$$

$$
F^+(x,\xi,t) = F^u(x,\xi,t) - F^d(x,\xi,t)
$$

<span id="page-5-0"></span>• In impact space

$$
\mathcal{H}_{\pi} = \int d\tau d^2 \vec{b} \,\hat{\Psi}_{\pi}(\tau, -\vec{b}) \hat{\mathcal{F}}_{\pi}^i(\bar{x}, \xi, \tau, Q^2, \vec{b}) \alpha_s(\mu_R) \exp(-S(\tau, \vec{b}, Q^2))
$$

Hard scattering kernels has the following forms in momentum space

$$
\mathcal{F}_{\pi^0}^q = \frac{N_c^2 - 1}{2N_c} \sqrt{\frac{2}{N_c}} \frac{Q}{\xi} \Big[ \frac{1}{k_{\perp}^2 + \tau(\bar{x} + \xi) Q^2 / (2\xi) - i\epsilon} - \frac{1}{k_{\perp}^2 - \bar{\tau}(\bar{x} - \xi) Q^2 / (2\xi) - i\epsilon} \Big]
$$
  

$$
\mathcal{F}_{\pi^+}^q = \frac{N_c^2 - 1}{2N_c} \sqrt{\frac{2}{N_c}} \frac{Q}{\xi} \Big[ \frac{e_d}{k_{\perp}^2 + \tau(\bar{x} + \xi) Q^2 / (2\xi) - i\epsilon} - \frac{e_u}{k_{\perp}^2 - \bar{\tau}(\bar{x} - \xi) Q^2 / (2\xi) - i\epsilon} \Big]
$$

• A Gaussian meson wave function is used at twist-2

$$
\Psi_\pi(\tau, \vec{b}) \sim \tau(1-\tau) \text{exp}\Big[\frac{\tau(\tau-1)}{4} \frac{\vec{b}^2}{a_\pi^2}\Big]
$$

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• Sudakov factor has the form

$$
S(\tau, b, Q) = s(\tau, b, Q) + s(\bar{\tau}, b, Q) - \frac{4}{\beta_0} \ln \frac{\ln(\mu_R/\Lambda_{QCD})}{\hat{b}}
$$

where

$$
s(\tau, b, Q) = \frac{8}{3\beta_0} \left( \hat{q} \ln\left(\frac{\hat{q}}{\hat{b}}\right) - \hat{q} + \hat{b} \right) + NLL
$$

$$
\hat{b} = -\ln(b \Lambda_{QCD})
$$

$$
\hat{q} = \ln(\tau Q/(\sqrt{2}\Lambda_{QCD}))
$$

• Twist-3 meson wave function

$$
\Psi_{\pi}(\tau, \vec{b}) \sim \exp\Big[-\frac{\vec{b}^2}{8a_{\pi}^2}\Big] I_0(\frac{\vec{b}^2}{8a_{\pi}^2})
$$

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GPDs are constructed from double distribution ansatz

$$
F_i^a(\bar{x}, \xi, t) = \int_{-1}^1 d\rho \int_{-1+|\rho|}^{1-|\rho|} d\eta \, \delta(\rho + \xi \eta - \bar{x}) f_i^a(\rho, \eta, t)
$$

where for valence-quark GPDs;

$$
f_i(\rho, \eta, t) = \exp[(b_i - \alpha'_i \ln \rho)t] F_i^a(\rho, \xi = t = 0) \frac{3}{4} \frac{(1 - \rho)^2 - \eta^2}{(1 - \rho)^3} \Theta(\rho)
$$

### • Parameters of the forward limits

- $H$  : DSSV, Phys. Rev. D 80, 034030 (2009)
- $H_T$ : ABM, Phys. Rev. D 86, 054009 (2012) and DSSV, Phys. Rev. D 80, 034030 (2009)
- $\tilde{E}$  : LHPC Collaboration, Phys. Rev. D 77, 094502 (2008)
- $\bar{E}_T$ : QCDSF and UKQCD Collaborations, Phys. Rev. Lett. 98, 222001 (2007)

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- So altogether, to compute the longitudinal cross section, we need: Twist-2 meson wave function, kernel, running coupling, Sudakov factor and GPDs  $H$  and  $E$ .
- To compute the transverse cross section, we need: Twist-3 meson wave function, kernel, running coupling, Sudakov factor and GPDs  $H_{\mathcal{T}}$  and  $\bar{E}_{\mathcal{T}}$ .
- 3 dimensional integrals, over  $\bar{x}$ ,  $\tau$  and b, are performed in impact space
- In many different processes, meson wavefunction has the same structure.
- Sudakov factor has also the same structure
- $\pi^+$  electroproduction also receives a pion pole contribution, besides the handbag contribution

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Comparision between the PARTONS and Mathmematica codes (slightly different GPDs parametrization for  $\tilde{H}$ )



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<span id="page-10-0"></span>Comparision between the PARTONS and Mathmematica codes at amplitude level (the same GPDs parametrization).

$$
Q^2 = 5 GeV^2, \ W = 6 GeV
$$

 $M_{0p0p}|_{t=-0.2} = 0.173902 - 0.150831i \rightarrow 0.1737324 - 0.1509104i$  $M_{0m0p}|_{t=-0.2} = 0.00422273 - 0.0502173i$  → 0.0041224 - 0.0502987i  $M_{0p0p}|_{t=-0.4} = 0.136403 - 0.108212i \rightarrow 0.1362792 - 0.1082896i$  $M_{0m0p}|_{t=-0.4} = 0.00742887 - 0.0477006i \rightarrow 0.0073482 - 0.0477759i$ 

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