

# Elastic Electron-Proton and Electron-Deuteron Scattering at the EIC

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(+ Douglas Higinbotham

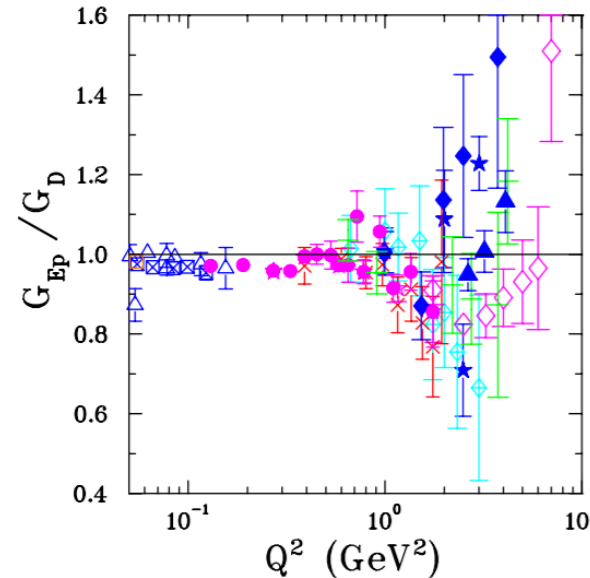
Andrew Puckett

Elena Long )

# Elastic Electron-Proton Scattering at the EIC

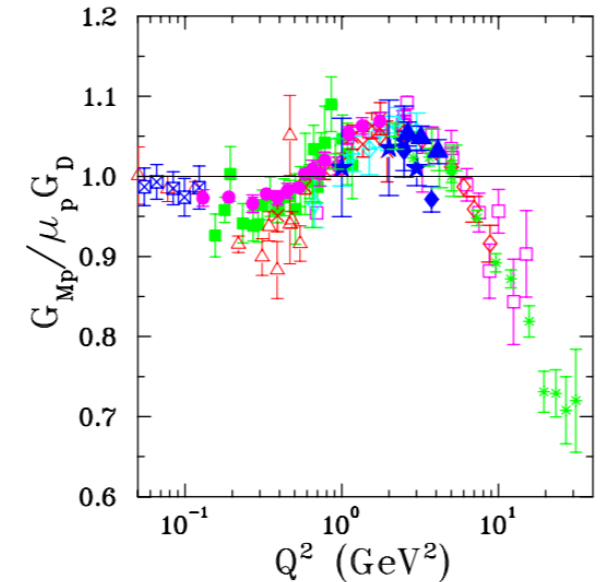
Elastic electron-proton scattering at high  $Q^2$  can be interesting in itself:

- Precision  $G_M$  required to study approach of QCD scaling in Dirac  $F_1$  Form Factor
- Constraints on GPDs at high- $x$  & high- $t$  via sum rules
- Possible increased sensitivity to hard two-photon exchange effects



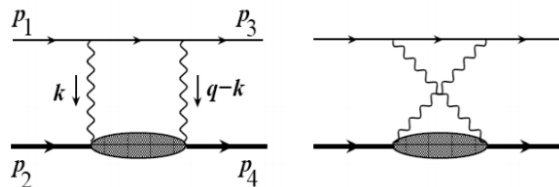
△ Han63  
◆ Lit70  
● Pri71  
× Ber71  
◇ Bar73  
☆ Han73

⊠ Bor75  
□ Sim80  
◇ And94  
★ Wal94  
+ Chr04  
▲ Qat05



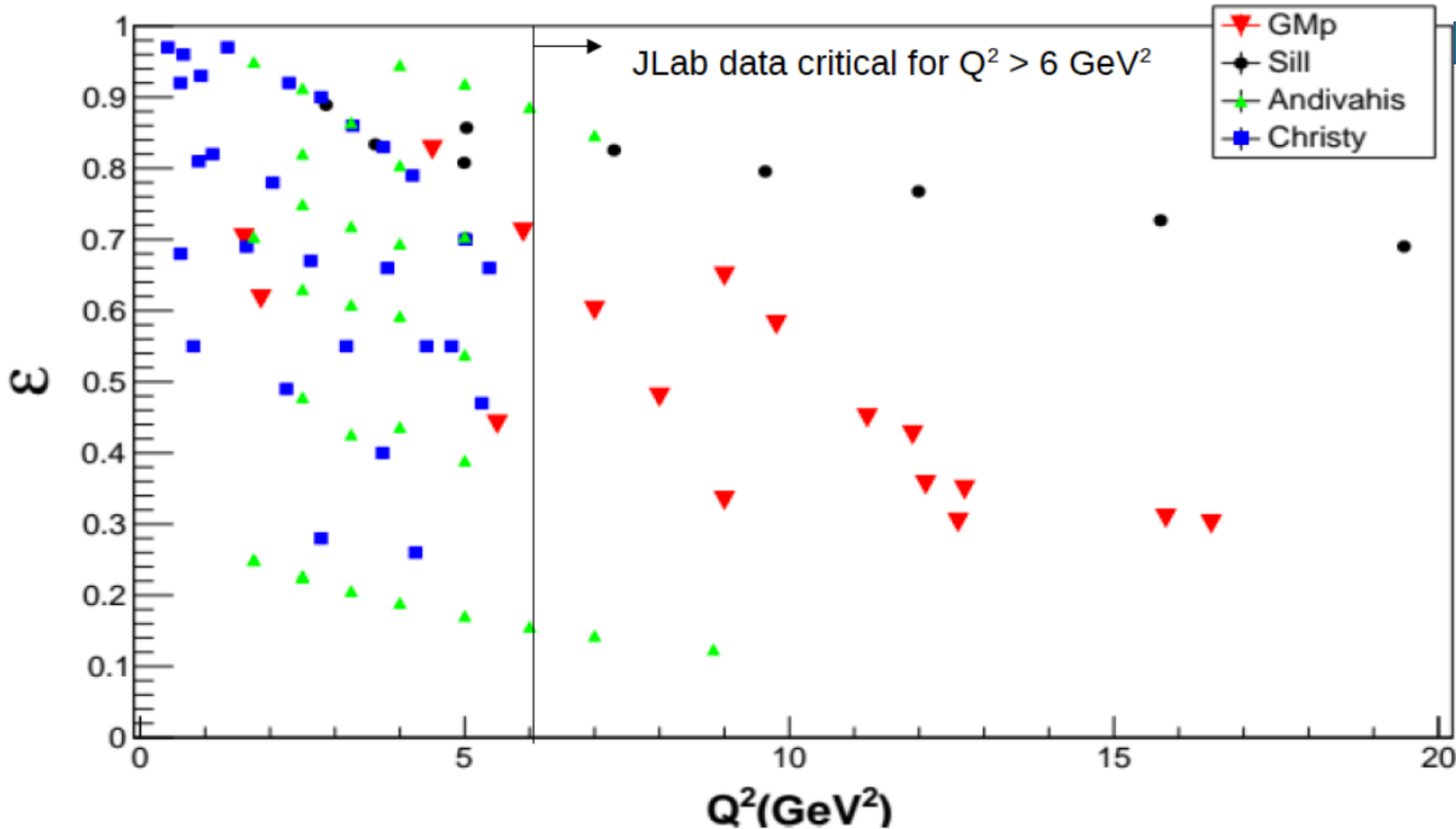
△ Han63  
■ Jan66  
□ Cow68  
◆ Lit70  
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◇ Bar73  
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\* Sil93  
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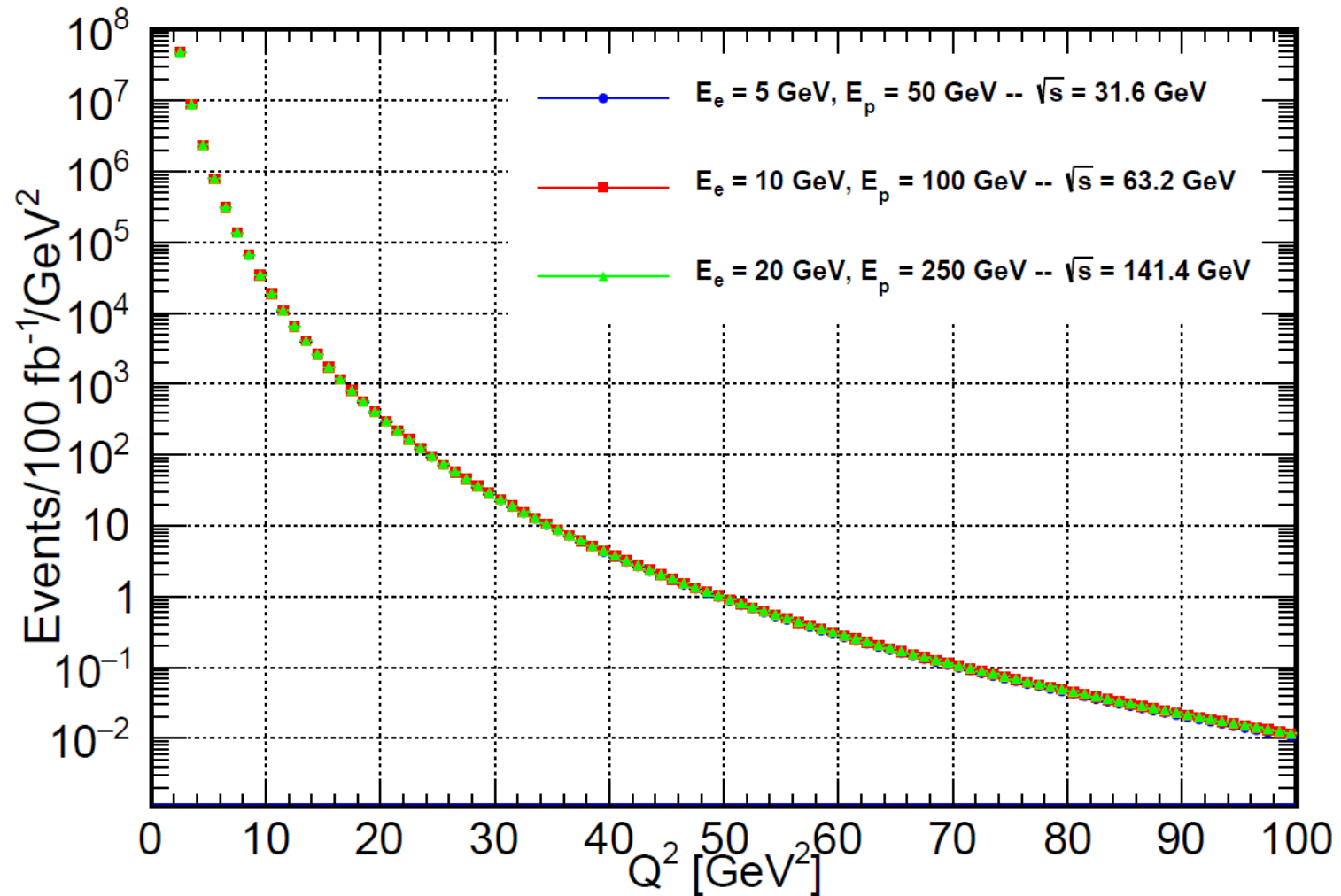
C.F Perdrisat, V. Punjabi, M. Vanderhaeghen, *Progress in Particle and Nuclear Physics* 59 (2007) 694–764

For ep elastic scattering, the *EIC* will allow us to probe the highest-ever values of  $Q^2$

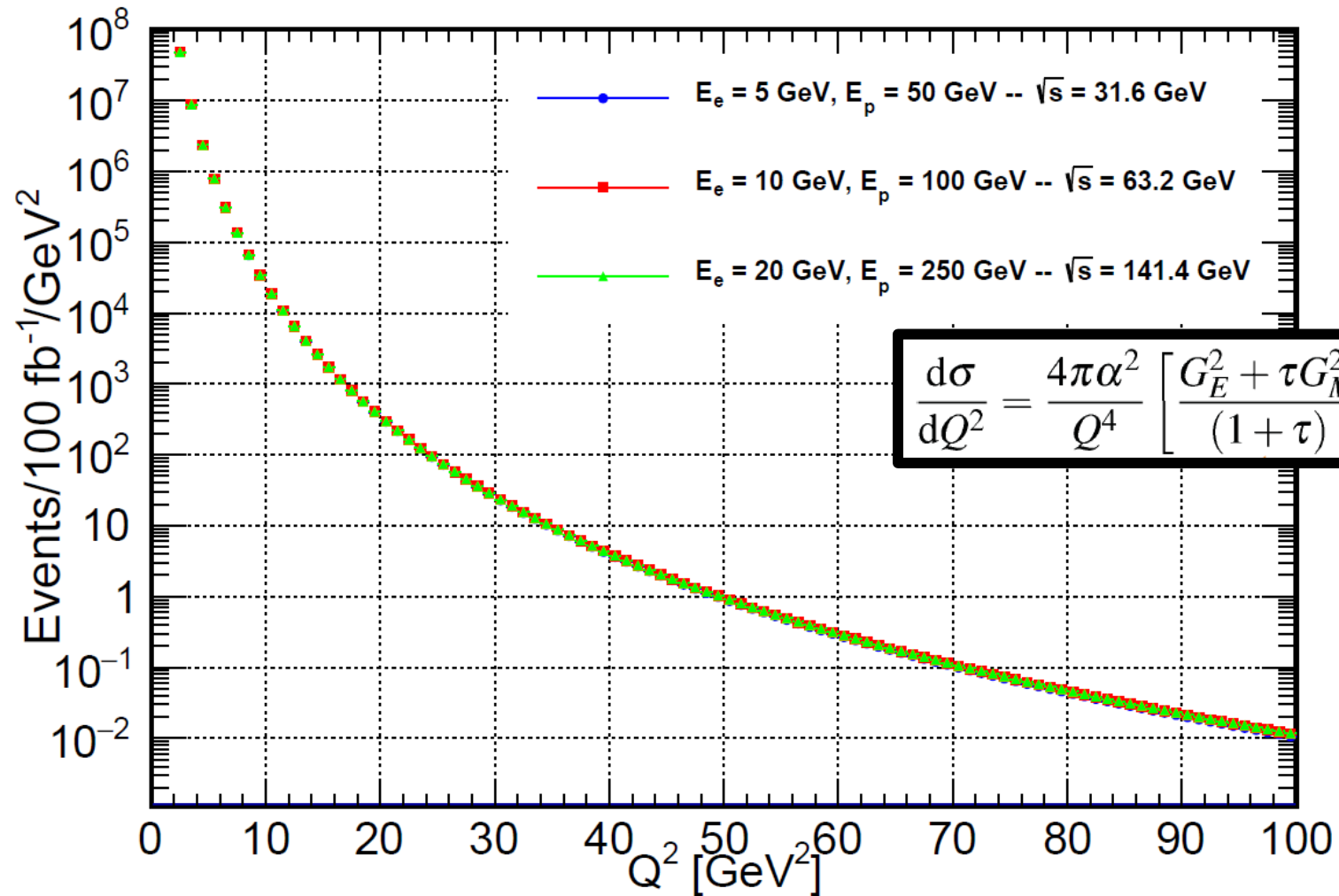


Up to  $Q^2 \sim 40 \text{ GeV}^2$   
or higher at the  
EIC – all at high  $\varepsilon$

# Electron-Proton Elastic scattering expected yields



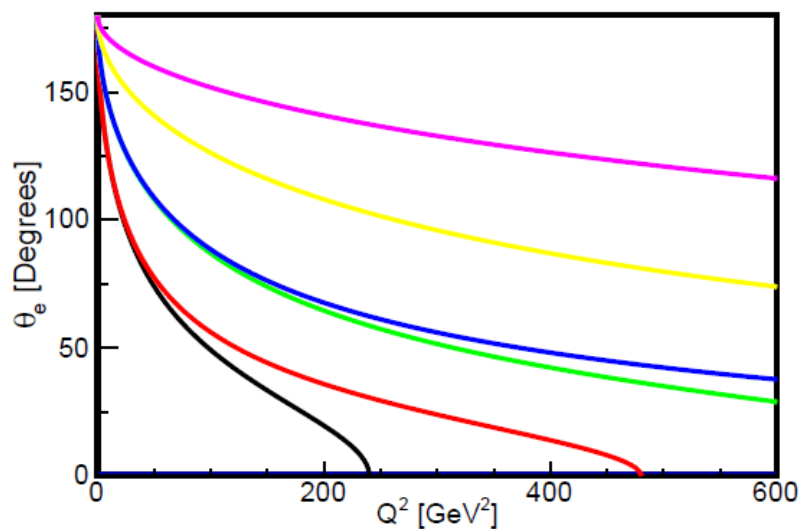
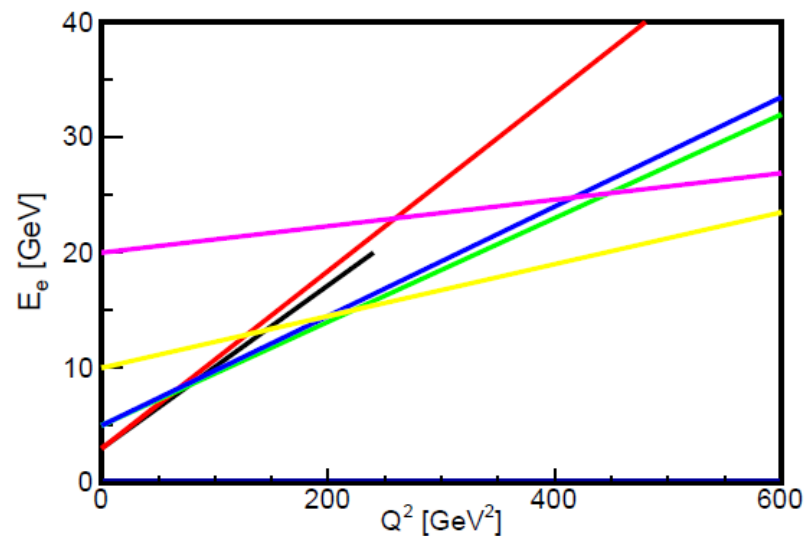
# Electron-Proton Elastic scattering expected yields



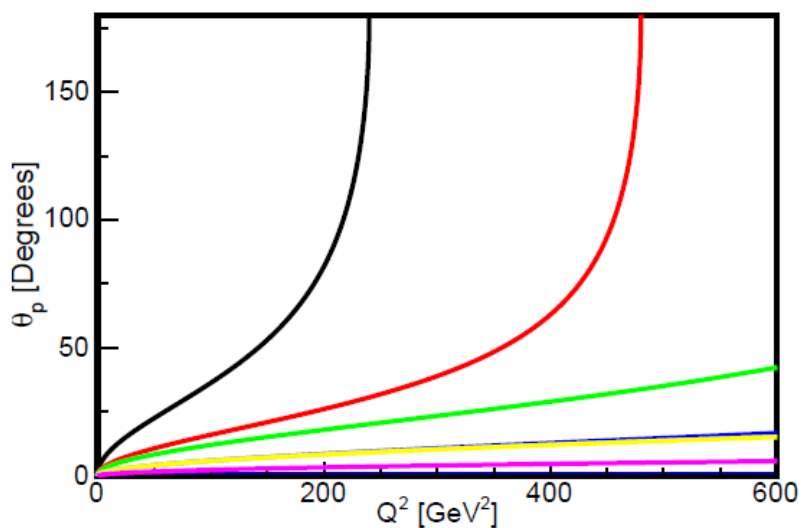
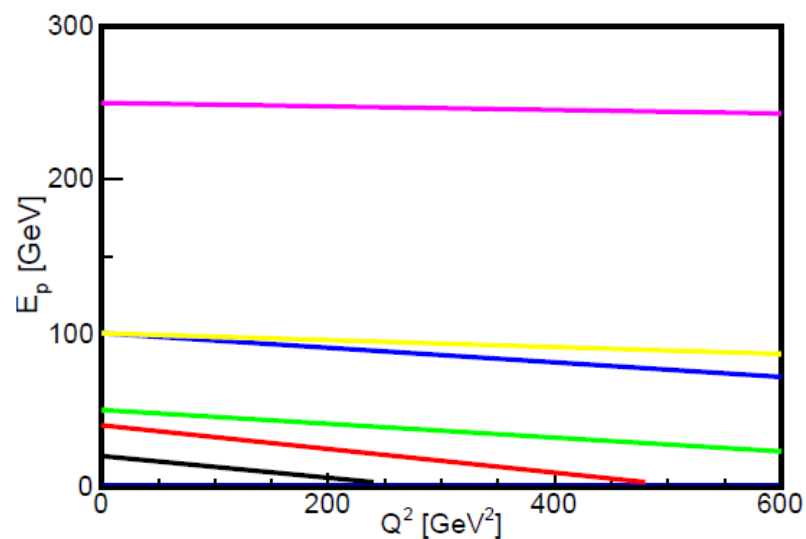
$$\frac{d\sigma}{dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[ \frac{G_E^2 + \tau G_M^2}{(1 + \tau)} \left( 1 - y - \frac{M^2 y^2}{Q^2} \right) + \frac{1}{2} y^2 G_M^2 \right]$$

**Lack of dependence on beam energy seen here is expected at the EIC since the data will all be at  $y \sim 0$  ( $\epsilon \sim 1$ ). See the Lorentz-invariant cross section formula above**

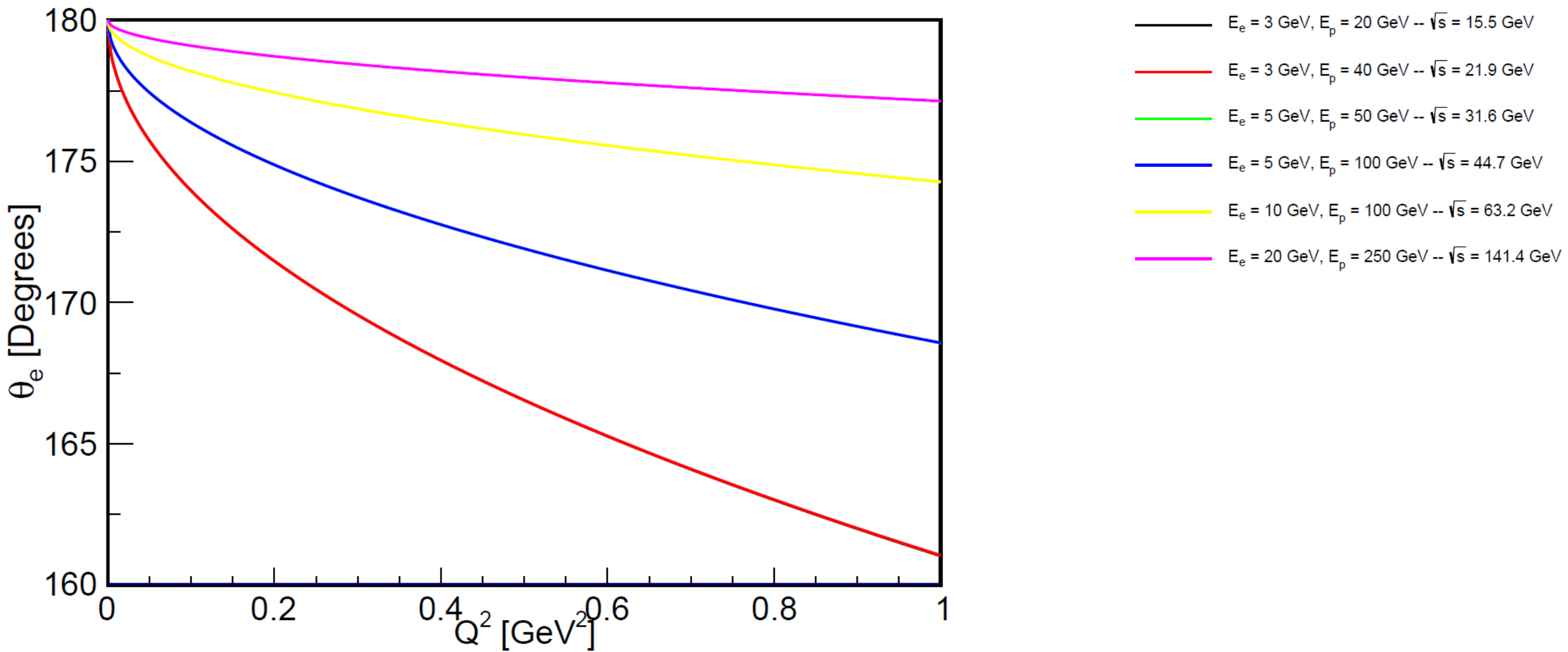
# Kinematics



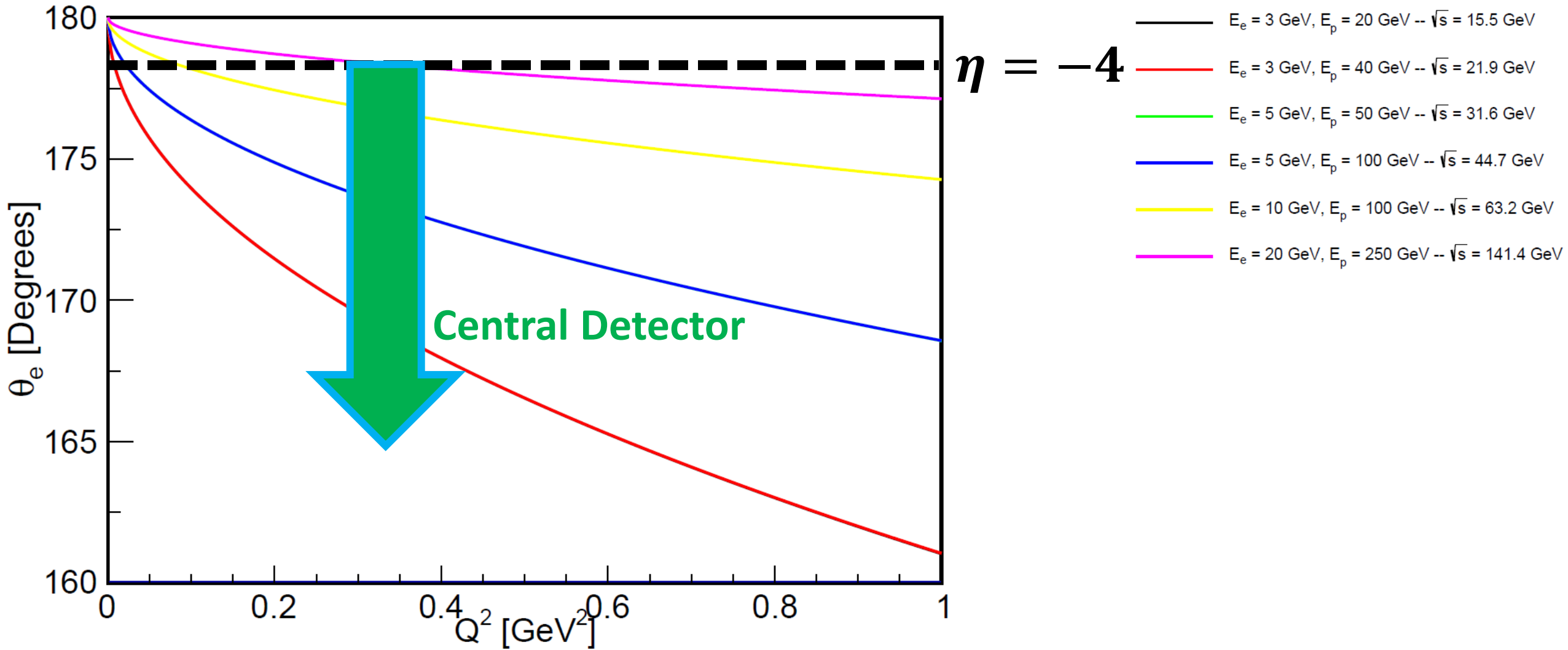
- $E_e = 3$  GeV,  $E_p = 20$  GeV --  $\sqrt{s} = 15.5$  GeV
- $E_e = 3$  GeV,  $E_p = 40$  GeV --  $\sqrt{s} = 21.9$  GeV
- $E_e = 5$  GeV,  $E_p = 50$  GeV --  $\sqrt{s} = 31.6$  GeV
- $E_e = 5$  GeV,  $E_p = 100$  GeV --  $\sqrt{s} = 44.7$  GeV
- $E_e = 10$  GeV,  $E_p = 100$  GeV --  $\sqrt{s} = 63.2$  GeV
- $E_e = 20$  GeV,  $E_p = 250$  GeV --  $\sqrt{s} = 141.4$  GeV



# Kinematics: Low $Q^2$ Electron Angle

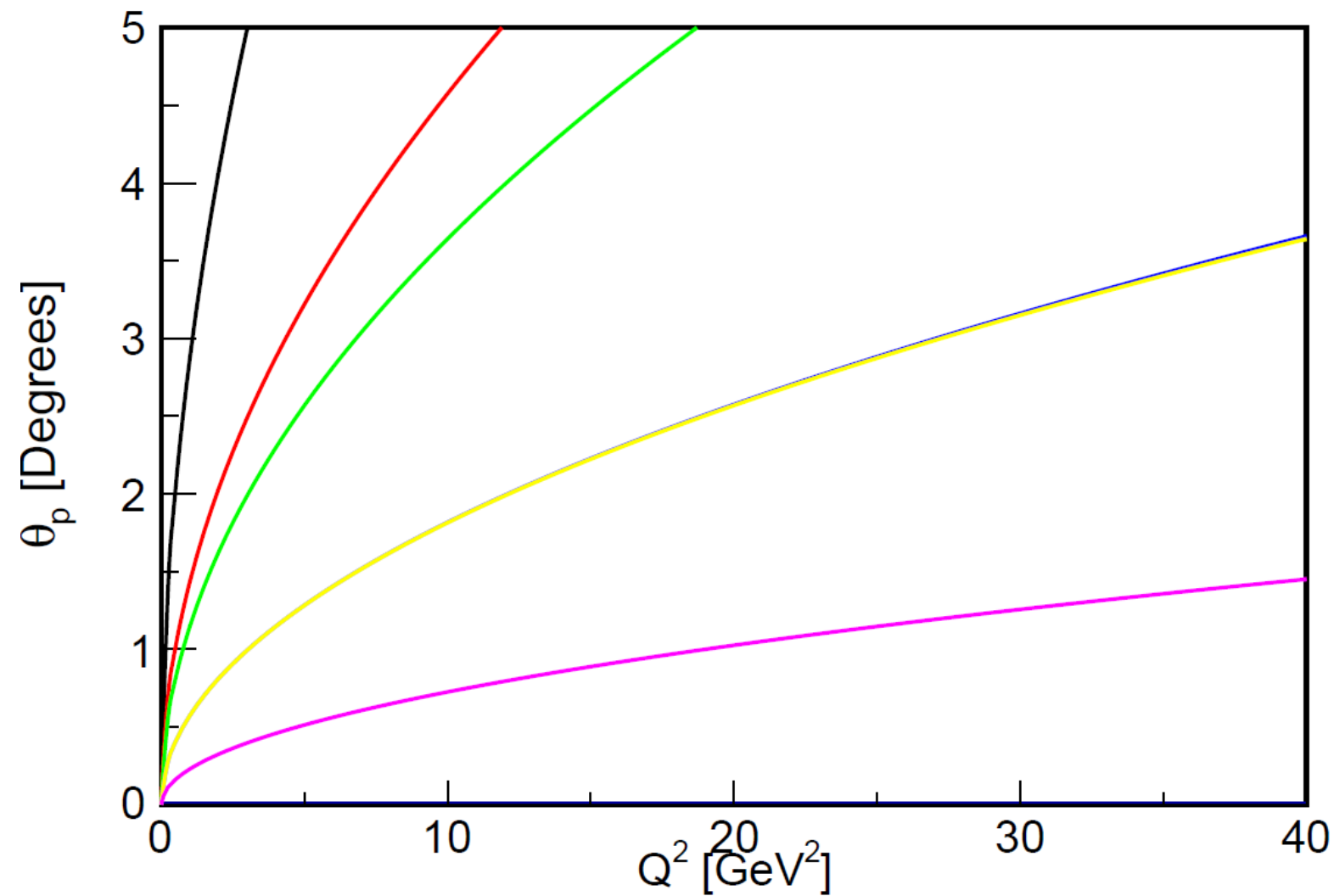


# Kinematics: Low $Q^2$ Electron Angle



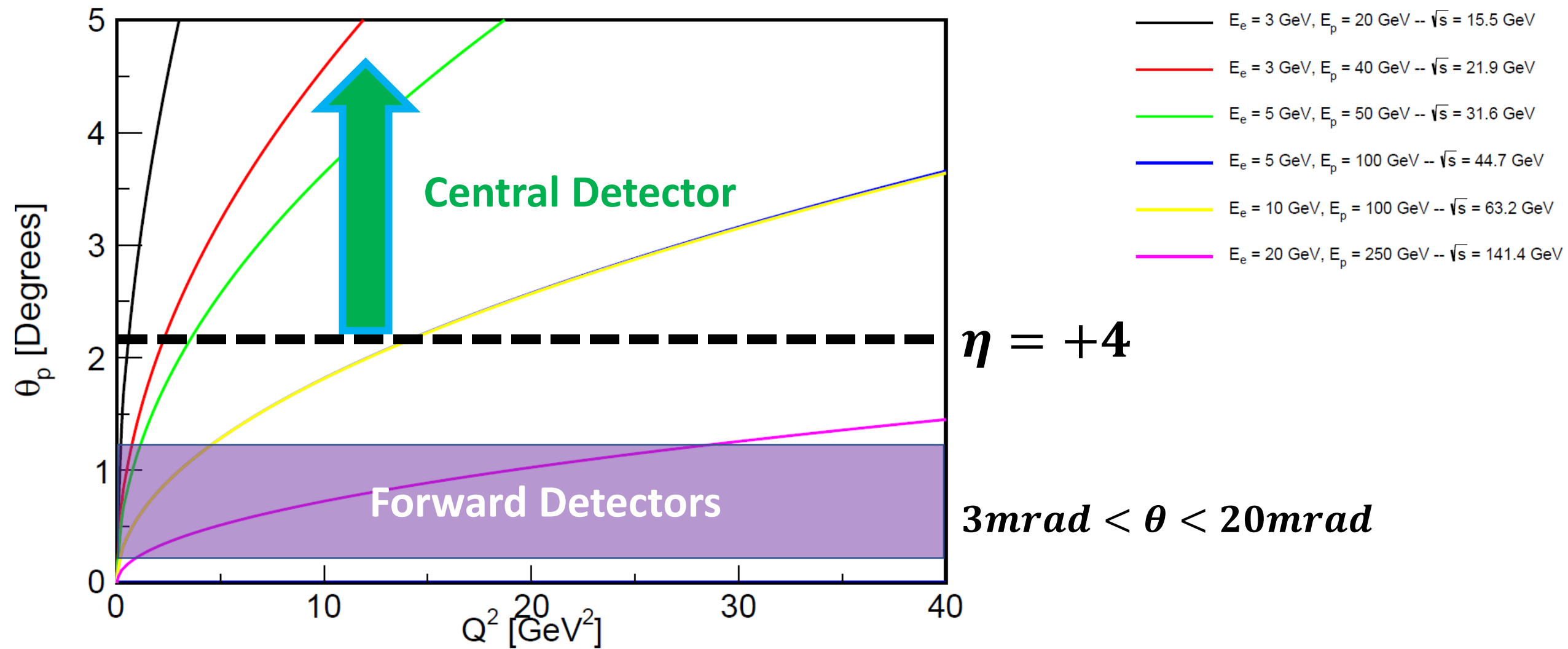


# Kinematics: Proton Angle



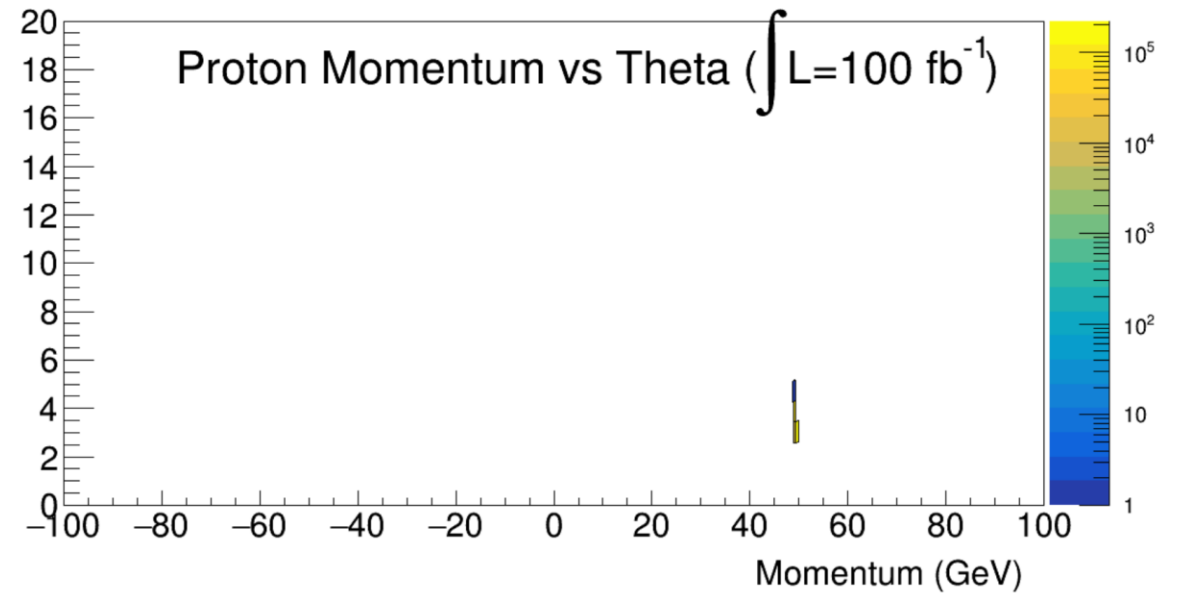
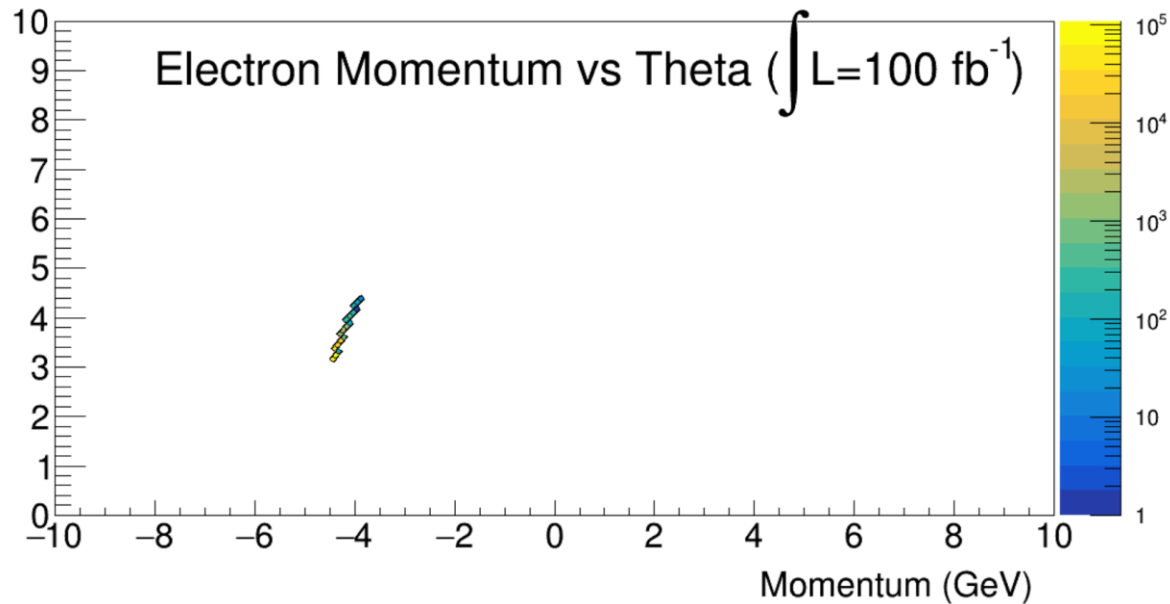
- $E_e = 3 \text{ GeV}, E_p = 20 \text{ GeV} \rightarrow \sqrt{s} = 15.5 \text{ GeV}$
- $E_e = 3 \text{ GeV}, E_p = 40 \text{ GeV} \rightarrow \sqrt{s} = 21.9 \text{ GeV}$
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# Kinematics: Proton Angle



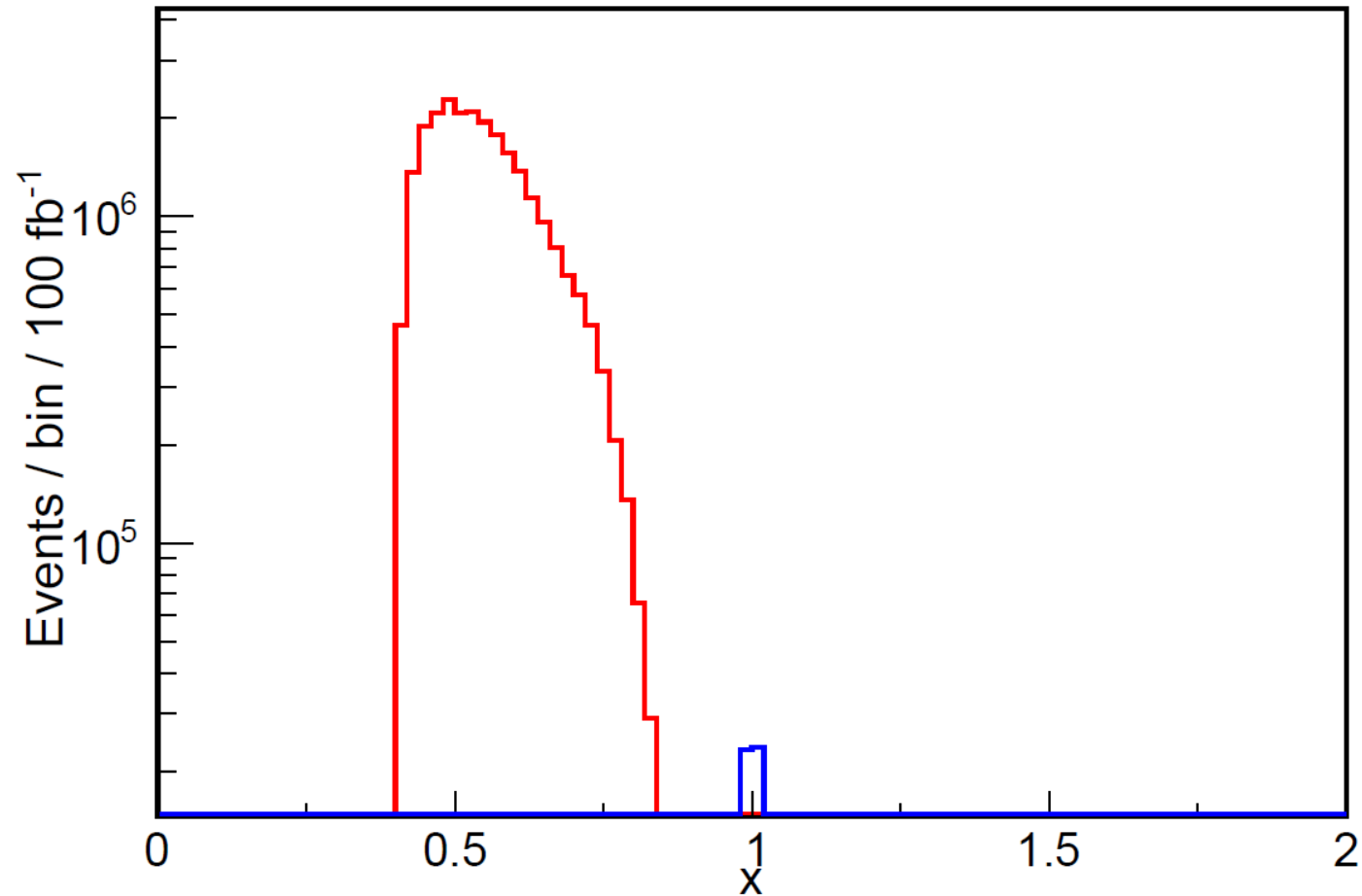
# Kinematics: Momentum vs. Angle

5 GeV electron on 50 GeV proton  
 $10 \text{ GeV}^2 < Q^2 < 20 \text{ GeV}^2$



# Mix With *Pythia6* Events – Generator Level

5 GeV  $e^{-1}$  on 50 GeV p,  $10 \text{ GeV}^2 < Q^2 < 20 \text{ GeV}^2$



**PYTHIA6 Events**  
**Elastic Events**

# Polarized Electron-Proton Elastic Scattering Asymmetry Measurements

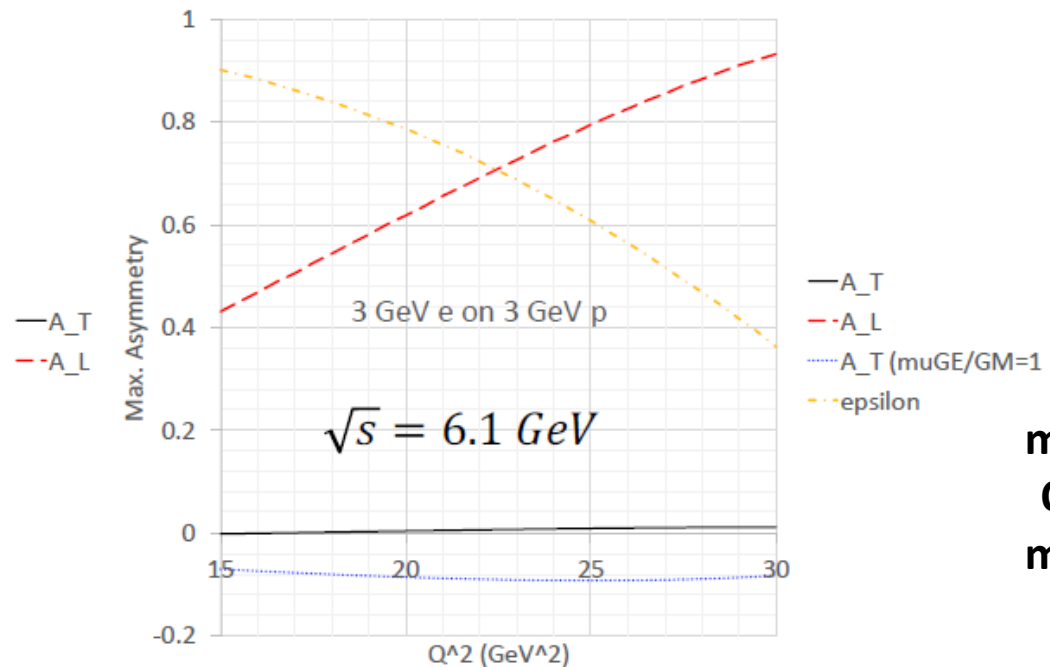
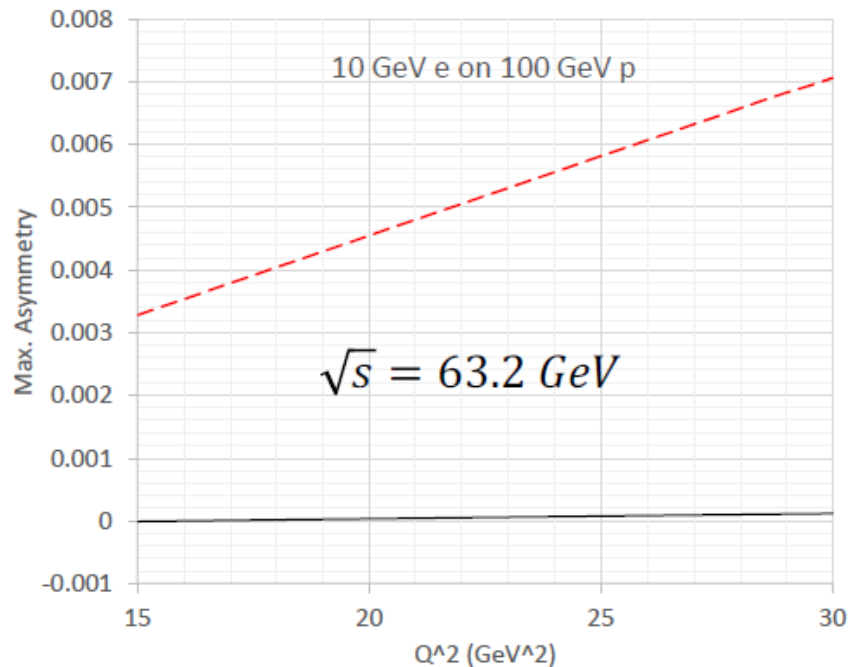
$$\begin{aligned}
 A_{eN} &= -\frac{P_{beam}P_{target}}{1 + \frac{\epsilon}{\tau}r^2} \left[ \left( \sqrt{\frac{2\epsilon(1-\epsilon)}{\tau}} \sin \theta^* \cos \phi^* \right) r + \sqrt{1-\epsilon^2} \cos \theta^* \right] \\
 &\equiv P_{target} [A_t \sin \theta^* \cos \phi^* + A_\ell \cos \theta^*]
 \end{aligned}
 \quad \Bigg| \quad r \equiv \frac{G_E}{G_M}$$

# Polarized Electron-Proton Elastic Scattering Asymmetry Measurements

$$A_{eN} = -\frac{P_{beam}P_{target}}{1 + \frac{\epsilon}{\tau}r^2} \left[ \left( \sqrt{\frac{2\epsilon(1-\epsilon)}{\tau}} \sin \theta^* \cos \phi^* \right) r + \sqrt{1-\epsilon^2} \cos \theta^* \right]$$

$$\equiv P_{target} [A_t \sin \theta^* \cos \phi^* + A_\ell \cos \theta^*]$$

$r \equiv \frac{G_E}{G_M}$



**To make reasonable measurements in this higher Q<sup>2</sup> range, we would need a much lower energy than will be provided by the EIC**

# Summary

- For unpolarized electron-proton scattering, we see the possibility of making high  $Q^2$  measurements. We are working to determine how easily elastic events can be separated from high- $x$  inelastic events.
- We will most likely be unable to make useful beam-target double-spin asymmetry measurements at high  $Q^2$  using the *EIC*.
- We are beginning to conduct simulation studies on elastic electron-deuteron scattering for both form factor measurements and tensor polarization determination.

# BACKUP



# Description of rest-frame Elastic generator with anti-parallel beams

1. Boost from lab frame to proton's rest frame
2. Generate events according to Born cross section below
3. Form factor parameterization comes from *Kelly* (PHYSICAL REVIEW C **70**, 068202 2004, Phys. Rev. C **96**, 055203).
4. We choose to generate  $100 \text{ fb}^{-1}$  worth of simulation data

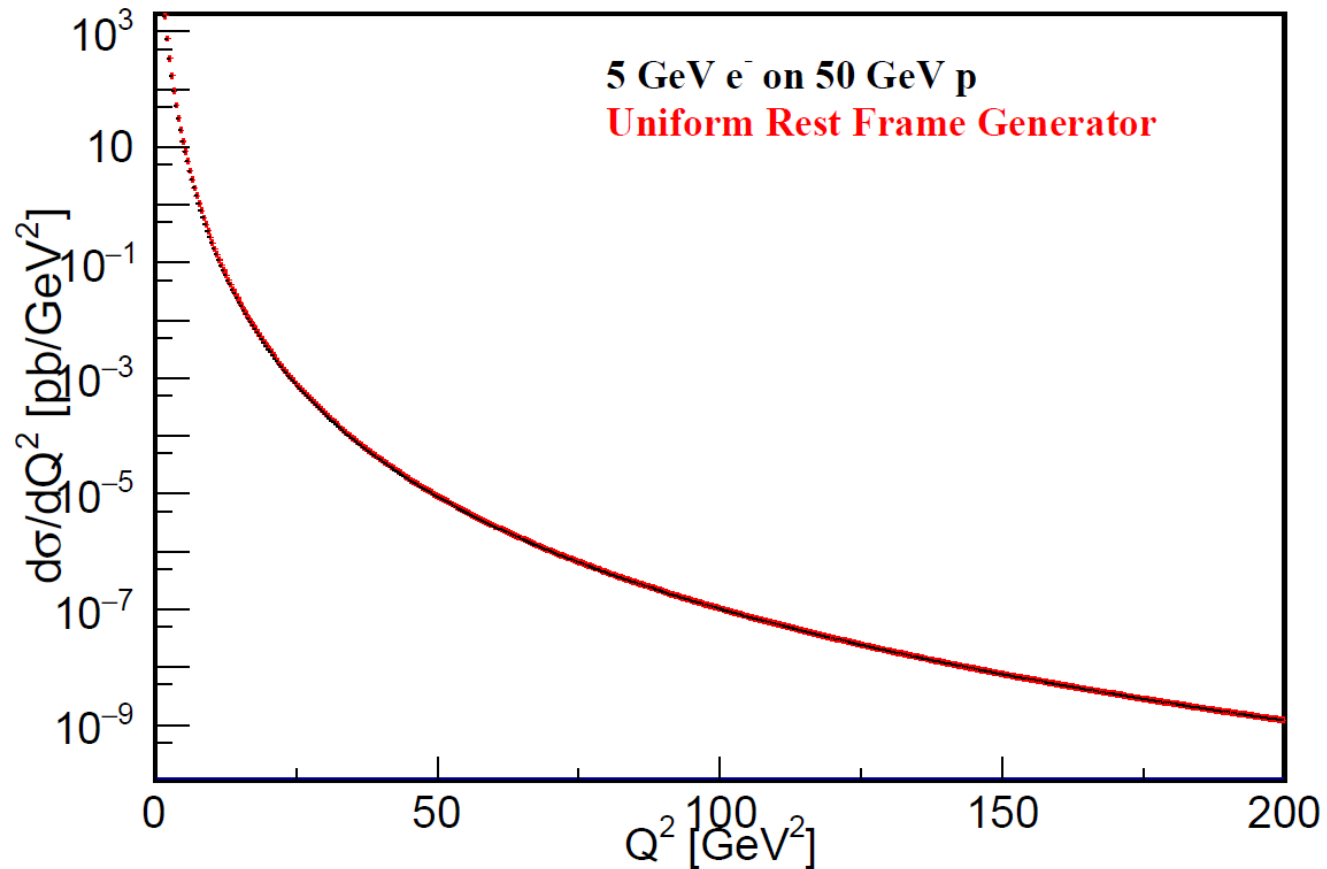
$$\begin{aligned} \frac{d\sigma}{d\Omega_e} &= \left( \frac{d\sigma}{d\Omega_e} \right)_{Mott} \frac{\epsilon G_E^2 + \tau G_M^2}{\epsilon(1 + \tau)} & \tau &\equiv \frac{Q^2}{4M_p^2} \\ \left( \frac{d\sigma}{d\Omega_e} \right)_{Mott} &= \frac{\alpha^2 \cos^2 \left( \frac{\theta_e}{2} \right) E'_e}{4E_e^2 \sin^4 \left( \frac{\theta_e}{2} \right) E_e} & \epsilon &\equiv \left[ 1 + 2(1 + \tau) \tan^2 \left( \frac{\theta_e}{2} \right) \right]^{-1} \\ \sigma_R &= \epsilon G_E^2 + \tau G_M^2 \end{aligned}$$

# Elastic Generator normalization for uniform generator

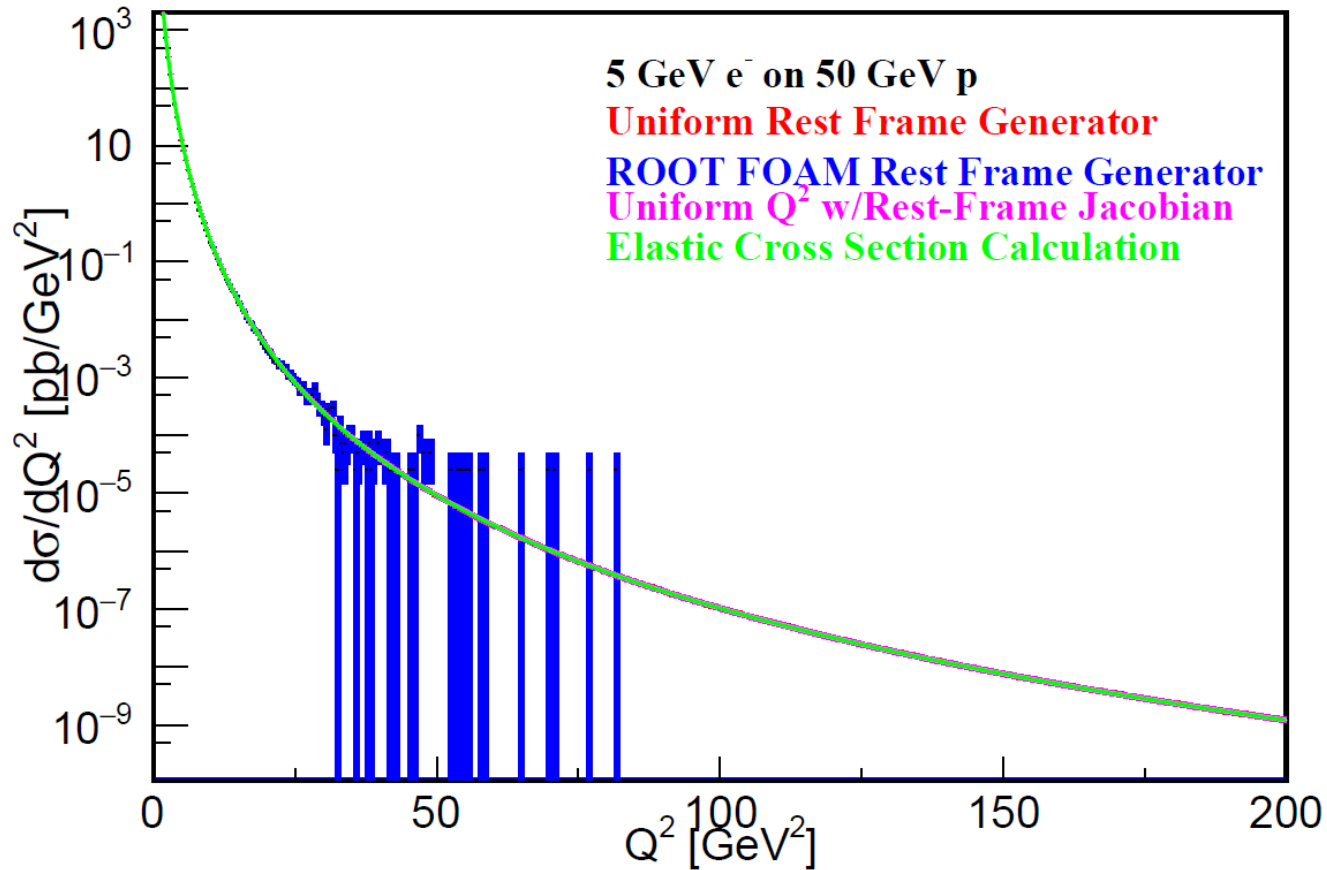
For Uniform Generator, to get expected yield (in bins of  $Q^2$ , for example), weight each event by:

$$\frac{\frac{d\sigma}{d\Omega_i}(\Omega_{tot})}{N_{tot} \Delta Q_i^2} \times \textit{luminosity}$$

# Generator agrees with cross section calculation



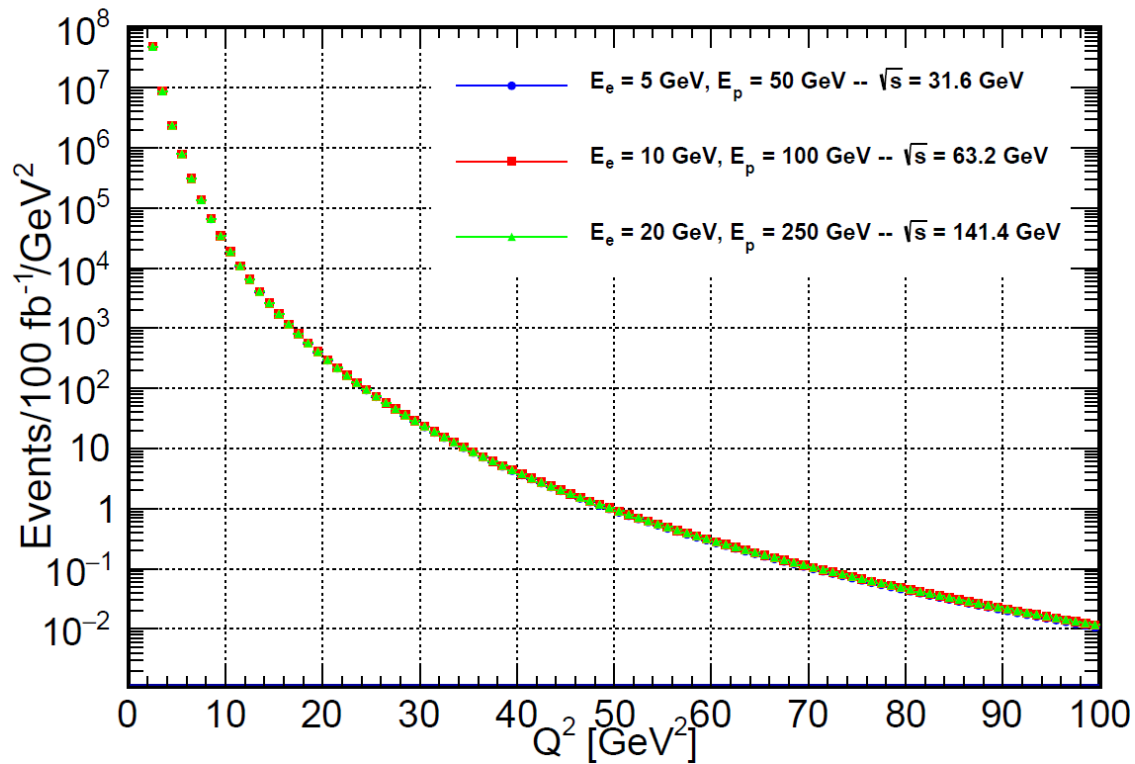
# Generator agrees with cross section calculation



$$\frac{d\sigma}{dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[ \frac{G_E^2 + \tau G_M^2}{(1+\tau)} \left( 1 - y - \frac{M^2 y^2}{Q^2} \right) + \frac{1}{2} y^2 G_M^2 \right]$$

# Generating in the proton rest frame and the lab frame also gives consistent results

## Rest Frame Generator



## Lab Frame Generator

