

# Pion-induced Drell-Yan and pion TMD distribution

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Workshop on Pion and Kaon Structure Functions at the EIC



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# Introduction

I present the analysis of  $q_T$ -spectrum of pion-induced DY  
within TMD factorization  
**JHEP 10 (2019) 090 [1907.10356]**

Main aim: Preparation for COMPASS  $\pi$ DY

## Plan of talk

- ▶ Reminder TMD factorization
- ▶  $\pi$ DY data TMD data
- ▶ Problems with normalization of fixed target DY
- ▶ TMDs for pion!



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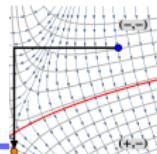
## TMD factorization formula for DY with TMD evolution

$$\frac{d\sigma}{dx dz dQ^2 d^2 \mathbf{q}_T} = \sum_{ff'} H_{ff'} \left( \frac{Q}{\mu} \right) \int d^2 b e^{i(\mathbf{b} \cdot \mathbf{q}_T)} F_{f \leftarrow p}(x_1, b, \mu, \zeta) F_{f' \leftarrow \pi}(x_2, b, \mu, \zeta)$$

TMD evolution (usual solution)

$$\frac{d\sigma}{dx dz dQ^2 d^2 \mathbf{q}_T} = \sum_{ff'} H_{ff'} \left( \frac{Q}{\mu} \right) \int \dots e^{\int_{\mu_0}^Q \gamma_V - \mathcal{D} \ln(\frac{\zeta}{\zeta_0})} F_{f \leftarrow p}(x_1, b, \mu_0, \zeta_0) F_{f' \leftarrow \pi}(x_2, b, \mu_0, \zeta_0)$$

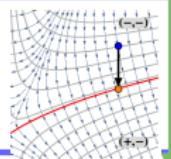
$$e^{\int_{\mu_0}^Q \gamma_V - \mathcal{D} \ln(\frac{\zeta}{\zeta_0})} = \exp \left( \int_{\mu_0}^\mu \frac{d\mu'}{\mu'} \left( \Gamma_{cusp}(\mu') \ln \left( \frac{\mu^2}{\sqrt{\zeta}} \right) - \gamma_V(\mu') \right) - D(b, \mu_0) \ln \left( \frac{\zeta}{\zeta_0} \right) \right)$$



## TMD factorization formula for DY with TMD evolution

$$\frac{d\sigma}{dxdzdQ^2d^2\mathbf{q}_T} = \sum_{ff'} H_{ff'} \left( \frac{Q}{\mu} \right) \int d^2b e^{i(\mathbf{b}\cdot\mathbf{q}_T)} F_{f \leftarrow p}(x_1, b, \mu, \zeta) F_{f' \leftarrow \pi}(x_2, b, \mu, \zeta)$$

TMD evolution (optimal solution)  
“ $\zeta$ -prescription”



$$\frac{d\sigma}{dxdzdQ^2d^2\mathbf{q}_T} = \sum_{ff'} H_{ff'} \left( \frac{Q}{\mu} \right) \int d^2b e^{i(\mathbf{b}\cdot\mathbf{q}_T)} \left( \frac{Q^2}{\zeta_\mu[\mathcal{D}]} \right)^{-2\mathcal{D}(b, \mu)} F_{f \leftarrow p}(x_1, b) F_{f' \leftarrow \pi}(x_2, b)$$

- ▶ Clean separation of TMDs from non-perturbative evolution  
(TMDs are defined at the point with  $\mathcal{D} = 0$ )
- ▶ Solution is made in terms of non-perturbative  $\mathcal{D}$
- ▶ Simple and fast expression for TMD evolution factor (just an algebraic function)
- ▶ Simpler expression for perturbative matching for TMDs



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## TMD factorization formula (in $\zeta$ -prescription)

Rapidity  
anomalous dimension  
 $\mathcal{D} \sim \langle 0 | F_{+b}[\text{staple link}] | 0 \rangle$   
 [AV, 2003.02288]

N	q	U	L	T
U	$f_1$			$h_1^\perp$
L			$g_1$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}$	$h_1$	$h_{1T}^\perp$

$F \sim \langle P | \bar{q}[\text{staple link}] q | 0 \rangle$

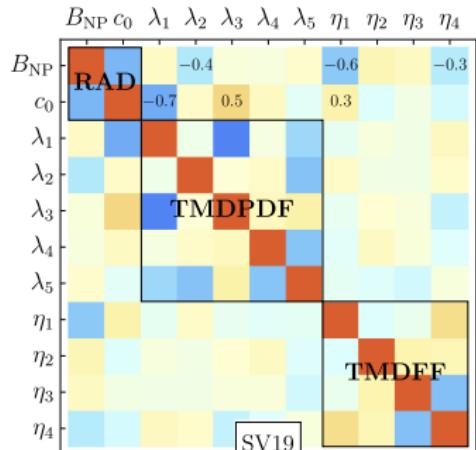
$$\frac{d\sigma}{dx dz dQ^2 d^2 \mathbf{q}_T} = \sum_{ff'} H_{ff'} \left( \frac{Q}{\mu} \right) \int d^2 b e^{i(\mathbf{b} \cdot \mathbf{q}_T)} \left( \frac{Q^2}{\zeta_\mu [\mathcal{D}]} \right)^{-2\mathcal{D}(b,\mu)} F_{f \leftarrow p}(x_1, b) F_{f' \leftarrow \pi}(x_2, b)$$

- ▶ Each data-point is a product (convolution) of **three independent non-perturbative** functions
- ▶ Functions do not “cross-talk” and could be modeled independently
- ▶ Each function is responsible for a separate kinematic variable
  - ▶ Rapidity AD:  $\mathcal{D} \rightarrow Q$  and  $b$
  - ▶ TMD N1:  $F_1 \rightarrow x_1$  and  $b$
  - ▶ TMD N2:  $F_2 \rightarrow x_2$  and  $b$

## Universality of TMDs

Global fit SV19  
unpolarized DY + SIDIS  
[1912.06532]

- ▶ DY: LHC, Tevatron, FermiLab, RHIC
- ▶ SIDIS: HERMES, COMPASS
- ▶ Large energy coverage ( $2 < Q < 150$  GeV) = decorrelation of RAD and TMDs
- ▶ NNLO matching +  $N^3LO$  evolution



TMD evolution and proton TMD PDF is known

To describe  $\pi$ DY one needs only  $\pi$ TMD PDF



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## Model for $\pi$ TMD PDF

$$F_{q \leftarrow \pi}(x, b) = \sum_f \int_x^1 \frac{dy}{y} C_{q \leftarrow f}(y, \mathbf{L}_\mu) f_{1, f \rightarrow \pi} \left( \frac{x}{y}, \mu \right) f_{\text{NP}}(x, b)$$

- ▶ NNLO matching to collinear distributions
- ▶ Collinear PDF is `JAM18pionPDFnlo`
- ▶ NP-part can be Gauss/Exponent, with three parameters  $a_{1,2,3}$

$$f_{\text{NP}}(x, b) = \exp \left( -\frac{(a_1 + (1-x)^2 a_2)b^2}{\sqrt{1+a_3 b^2}} \right)$$



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$\pi$ DY  $q_T$ -data  
not too much...

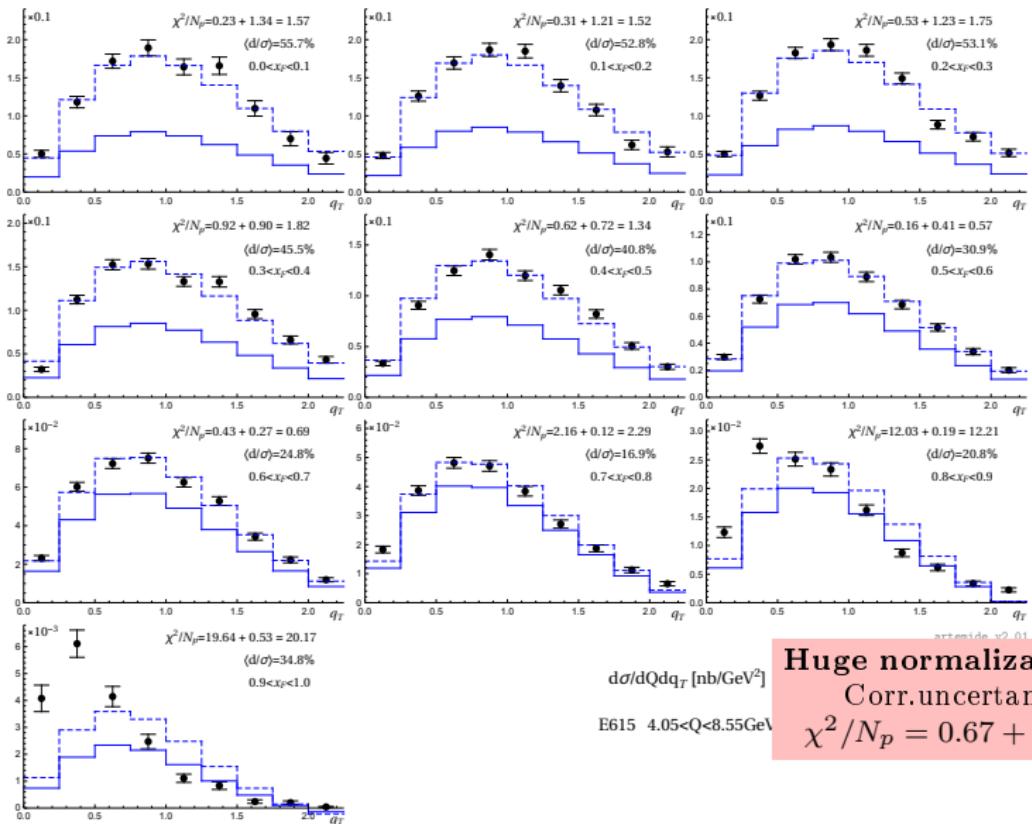
Experiment	$\sqrt{s}$ [GeV]	$Q$ [GeV]	$x_F$	$N_{pt}$	corr.err.	Typical stat.err.
E537 ( $Q$ -diff.)	15.3	$4.0 < Q < 9.0$ in 10 bins	$-0.1 < x_F < 1.0$	60/146	8%	$\sim 20\%$
E537 ( $x_F$ -diff.)	15.3	$4.0 < Q < 9.0$	$-0.1 < x_F < 1.0$ in 11 bins	110/165	8%	$\sim 20\%$
E615 ( $Q$ -diff.)	21.8	$4.05 < Q < 13.05$ in 10(8) bins	$0.0 < x_F < 1.0$	51/155	16%	$\sim 5\%$
E615 ( $x_F$ -diff.)	21.8	$4.05 < Q < 8.55$	$0.0 < x_F < 1.0$ in 10 bins	90/159	16%	$\sim 5\%$
NA3	16.8 , 19.4 22.9	$4.1 < Q < 8.5$ $4.1 < Q < 4.7$	$y > 0(?)$ $0 < y < 0.4$	—	15%	—

The usual TMD cut ( $q_T/Q < 0.25$ )

Data selected for fit

reason:

- ▶ E537 has too large uncertainties
- ▶ NA3 is available only as a plot
- ▶ It is better to use  $x$ -differential data, since  $Q$ -dependence is known.



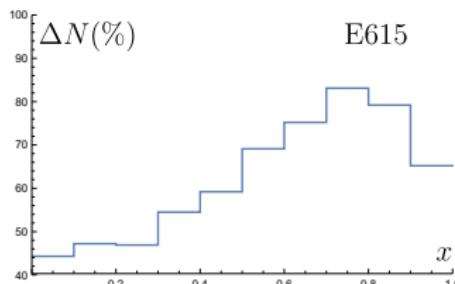
Huge normalization deficit  
Corr. uncertainty = 16%  
 $\chi^2/N_p = 0.67 + 0.77 = 1.44$



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## Possible sources of discrepancy

- ▶ Nuclear effects: multiplication by  $R(x)$  could give  $\sim 10\%$  effect (not enough +  $x$ -dependence)
- ▶ Effect of  $\pi$ PDF uncertainty: significant at a point (up to 20%) but only 2-3% in normalization
- ▶ Threshold logarithms: could be but the main disagreement is at  $x \sim 0.2$  and decreases to  $x \sim 0.7$
- ▶ Power corrections: unknown (but I would not expect more than 5%)
- ▶ Model bias: ?? (I don't think so)
- ▶ Resonance contamination: The bins go down to  $3 - 4\text{GeV}$ , they could be contaminated by  $J/\psi$  or  $\psi'$  resonances. However, typically, in this case theory overshoots the data.



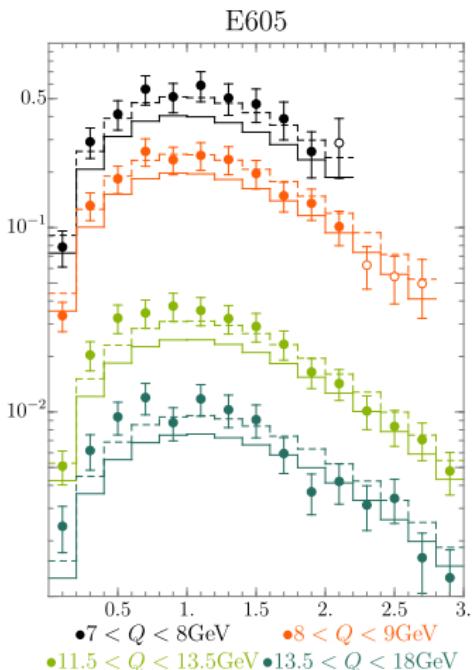
**Conclusion:**  
I do not see theory sources that would cover  
50% normalization gap.



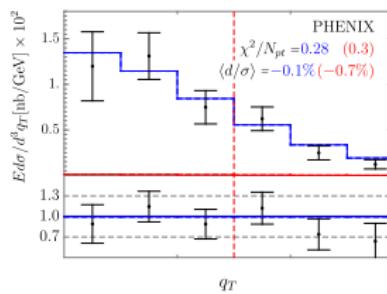
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The deficit in normalization is typical for TMD description of fixed-target experiments

However, it is usually 10-20% (not 50%)



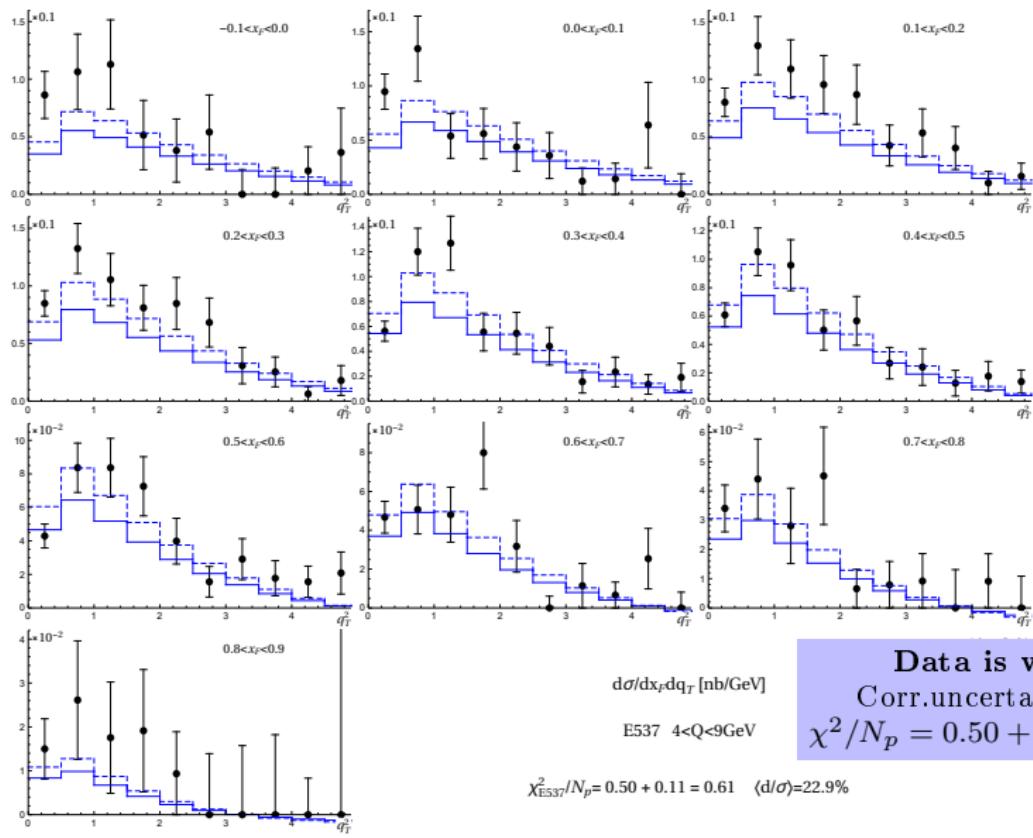
► Such problem does not exists for collider experiments



$$4.8 < Q < 8.2$$



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$d\sigma/dx_F dq_T$  [nb/GeV]

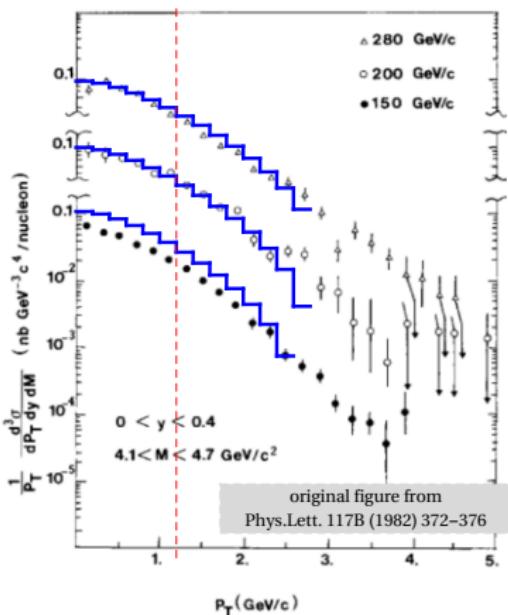
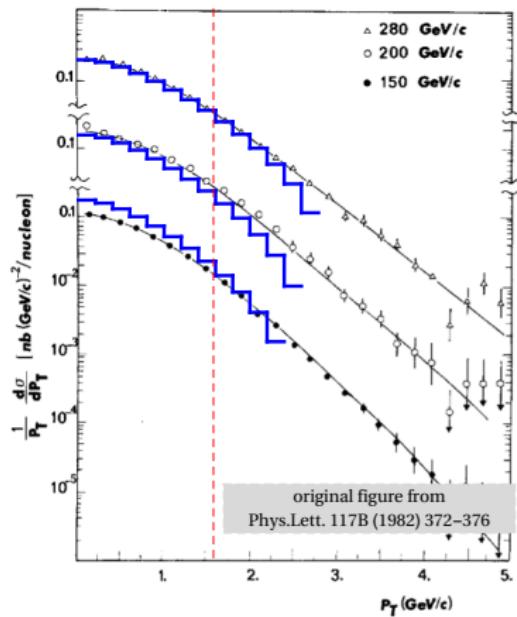
E537  $4 < Q < 9$  GeV

Data is worse  
Corr. uncertainty = 8%  
 $\chi^2/N_p = 0.50 + 0.11 = 0.61$

$$\chi^2_{E537}/N_p = 0.50 + 0.11 = 0.61 \quad \langle d\sigma \rangle = 22.9\%$$



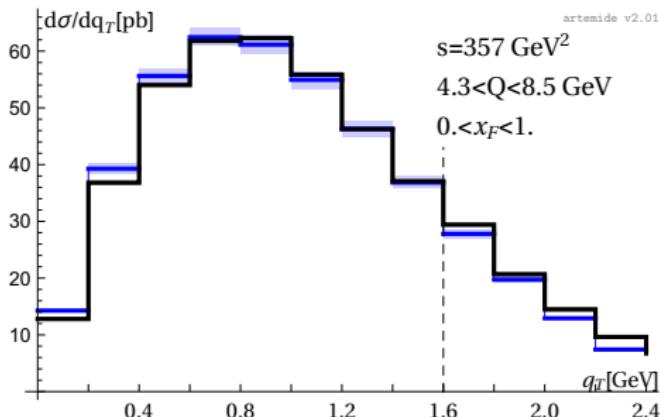
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$\pi$ DY-data from NA3 abd E537 does not show such anomalous behavior



Comparison with COMPASS preliminary [Phys.Rev.Lett. **119** (2017)]



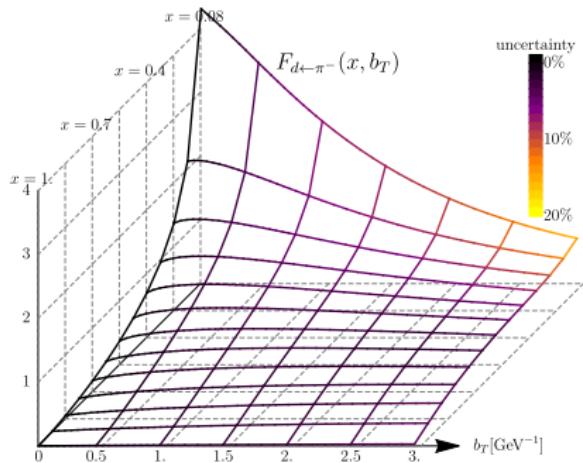
Normalization by theory



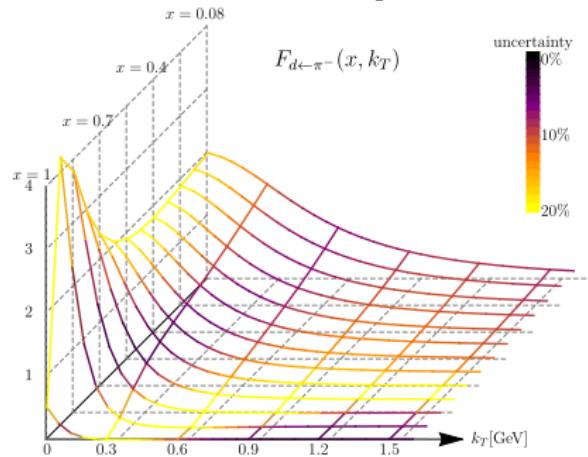
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$$a_1 = 0.17 \pm 0.11 \pm 0.03, \quad a_2 = 0.48 \pm 0.34 \pm 0.06, \quad a_3 = 2.15 \pm 3.25 \pm 0.32.$$

### Position-space



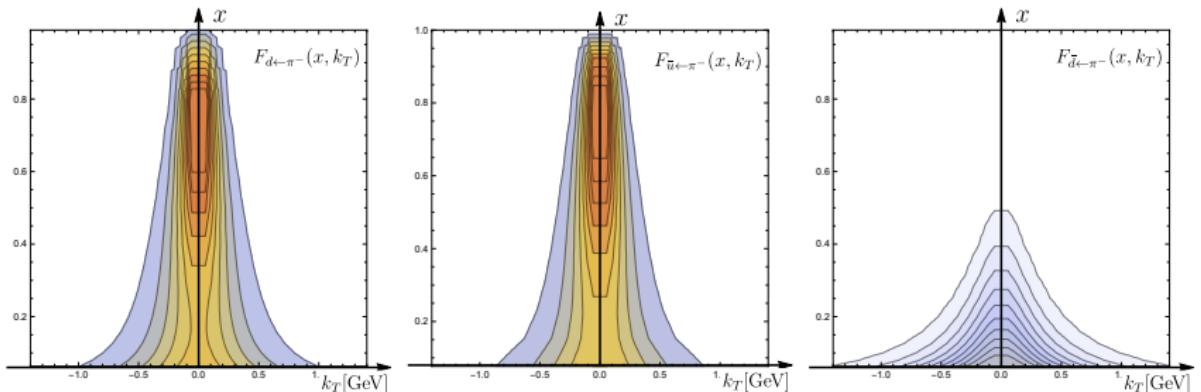
### Momentum-space



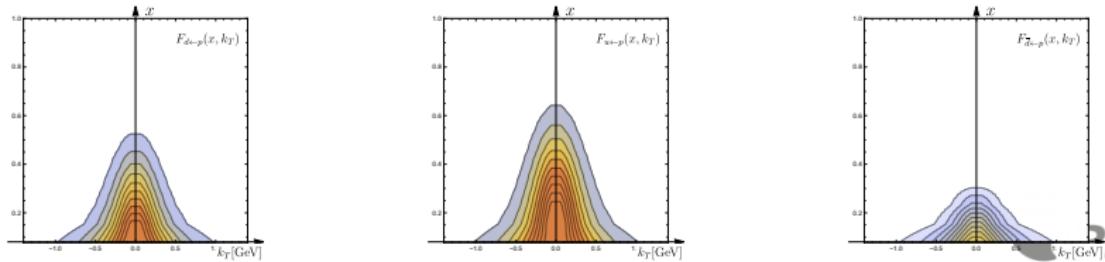
Only statistical uncertainties are shown



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Pion is “narrower” in the momentum space.



## Conclusion

- ▶ The analysis of  $\pi$ DY is made within the best available TMD-framework
  - ▶ NNLO matching,  $N^3LO$  evolution
  - ▶ Other components (proton TMD, NP TMD evolution) are from the global analysis SV19
  - ▶ Numerics is done by `artemide`
- ▶ Pion TMDs are extracted from the existing data
  - ▶ Normalization problem of  $q_T$ -dependent E615-data
  - ▶ Rest available data (lower quality) have no problem with normalization
- ▶ The extraction `Vpion19` is available as a part of default `artemide` distribution  
`github.com/VladimirovAlexey/artemide-public`

## Prospects

- ▶ The road to consistent global analysis ( $\pi$ DY angular modulations by COMPASS)
- ▶ Looking forward for COMPASS unpolarized  $\pi$ DY
- ▶ ...

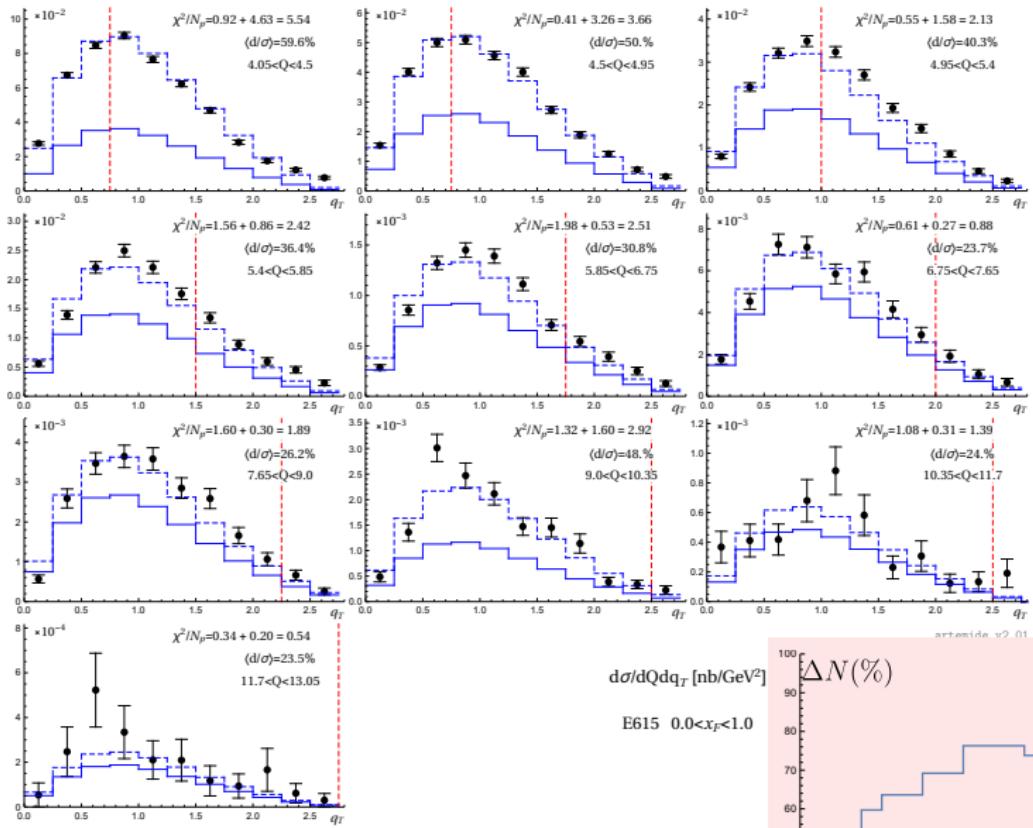


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# Backup slides



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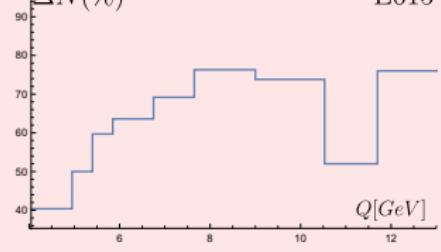


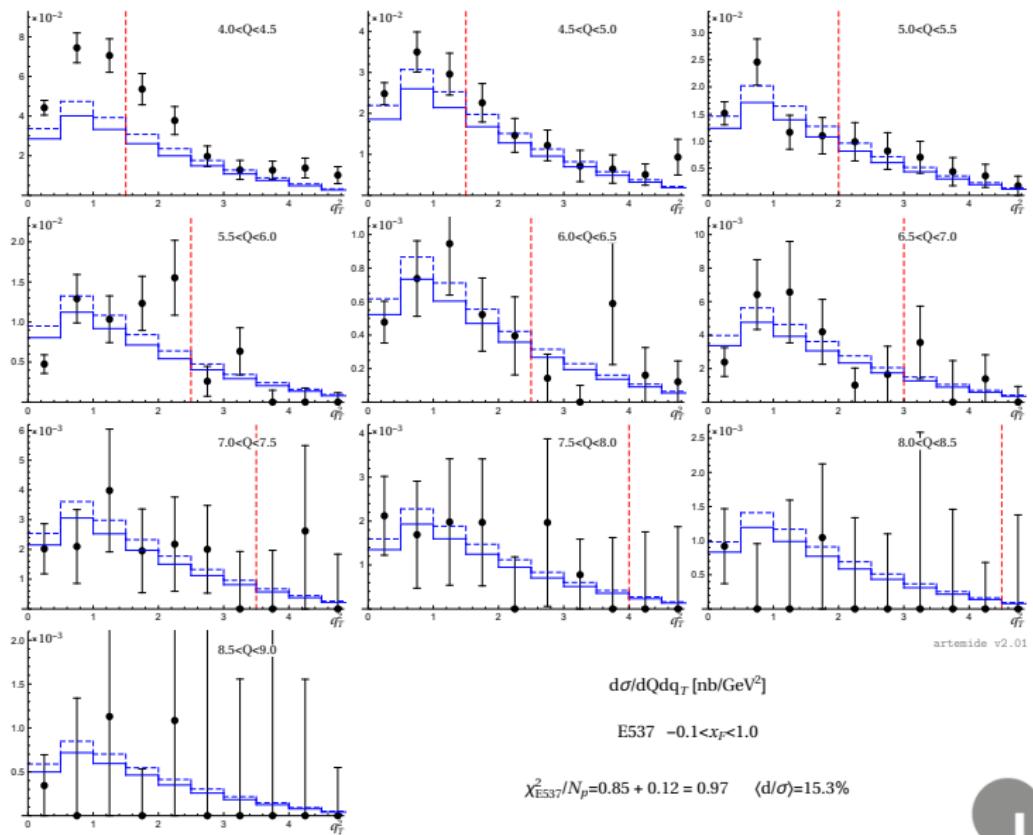
$d\sigma/dQdqt [nb/GeV^2]$

E615 0.0 <  $x_F$  < 1.0

$\Delta N(\%)$

E615





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