# THE TRANSVERSE STRUCTURE OF THE PION IN MOMENTUM SPACE FROM ADS/QCD MODELS

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in collaboration with: Alessandro Bacchetta and Barbara Pasquini (UniPV) based on: Physics Letters B 771 (2017) 546–552

Workshop on Pion and Kaon Structure Functions at the EIC 2-5 June 2020

## HADRON STRUCTURE AND LFWFS



$$|P,\Lambda\rangle = \sum_{N,\beta} \int \left[\frac{dx}{\sqrt{x}}\right]_N \left[d^2 \mathbf{k}_T\right]$$

 $]_{N} \psi^{\Lambda}_{N,\beta}(r) | N; k_{1}, \cdots, k_{N}, \beta_{1}, \cdots, \beta_{N} \rangle$ 



## HADRON STRUCTURE AND LFWFS

$$|P,\Lambda\rangle = \sum_{N,\beta} \int \left[\frac{dx}{\sqrt{x}}\right]_N \left[d^2 \mathbf{k}_T\right]_N$$

LFWFs overlap representation formulae:

$$\sum_{\beta} \int d^{2} \boldsymbol{k}_{T} |\psi_{\beta}^{\Lambda}(x, \boldsymbol{k}_{T})|^{2} \qquad \sum_{\beta=\beta'} \int dx d^{2} \boldsymbol{k}_{T} |\psi_{\beta}^{\Lambda}(x, \boldsymbol{k}_{T})|^{2}$$

$$PDF \qquad Form$$



 $\Big]_N \psi^{\Lambda}_{N,\beta}(r) | N; k_1, \cdots, k_N, \beta_1, \cdots, \beta_N \rangle$ 

 $\boldsymbol{k}_{T}\psi_{\beta'}^{\Lambda'}(x,\boldsymbol{k}_{T}')\psi_{\beta}^{\Lambda}(x,\boldsymbol{k}_{T})$ 







# LIGHT-FRONT HOLOGRAPHIC (LFH) QCD

String theory in a 5d Anti-de Sitter space

Conformal field theory on the 4-dimensional boundary of the AdS

#### **Applicability to QCD?**

Soft-wall model (harmonic confining potential  $U(z) \sim \kappa^2 z^2 \rightarrow Confinement$ 

[Brodsky, de Téramond et al., 2004-present]

**Form Factors matching** 



[J.M. Maldacena, (1999)]





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[Brodsky, de Téramond et al., 2004-present]

Free current propagating in AdS space Confining current in a warped AdS space

Valence LFWF

$$\psi_{q\bar{q}/\pi}^{V}(x, \mathbf{k}_{T}) \sim \frac{1}{\kappa\sqrt{(1-x)x}} e^{-\frac{1}{2}\frac{\mathbf{k}_{T}^{2}}{\kappa^{2}x(1-x)}}$$



#### **Form Factors matching**

[J.M. Maldacena, (1999)]

**Effective LFWF** 

$$\psi_{q\overline{q}/\pi}^{E}\left(x, \ \boldsymbol{k}_{T}\right) \sim \frac{\sqrt{\log\left(\frac{1}{x}\right)}}{\kappa\left(1-x\right)} e^{-\frac{\log(1/x)}{(1-x)^{2}}\frac{\boldsymbol{k}_{T}^{2}}{2\kappa^{2}}}$$





# IN THE LITERATURE

#### Theory

- Dosch, et al. (2006-present)
- Soft-wall model for AdS/QCD: Karch, Katz, et al. (2006)
- Phenomenological studies
- et al (2017-2020), Kaur, Dahiya(2019), Chang, Raya, Wang (2020), etc...

Light-front holography (LFH) original approach and improvements: Brodsky, de Téramond, Deur,

Harmonic potential in LF corresponds to a linear (confinining) potential in IF: Trawinski, et al (2014)

PDFs, FFs, TMDs, GPDs, double PDFs Brodsky, Cao, deTéramond (2011), Forshaw, Sandapen (2012), Vega, Schmidt, Gutsche, Lyubovitskij, et al (2009-2020), Chakrabarti, Mondal, et al. (2013-2019), Bacchetta, Cotogno, Pasquini (2017), Rinaldi, Traini, Vento, et al. (2017-2020), Ahmady, Sandapen,



#### MODELS OF LFWFS

**Inclusion of the quark masses** → completion of the invariant mass operator

$$M^{2} = \sum_{i} \frac{m_{i}^{2} + k_{Ti}^{2}}{x_{i}} = \frac{m^{2} + k_{T}^{2}}{x(1-x)}$$

Valence LFWF (from bound state equation)

$$\psi_{q\bar{q}/\pi}^{V}(x, \ \boldsymbol{k}_{T}) = A \frac{4\pi}{\kappa \sqrt{(1-x)x}} e^{-\frac{1}{2\kappa^{2}} \left(\frac{m^{2}}{x(1-x)} + \frac{\boldsymbol{k}_{T}^{2}}{x(1-x)}\right)}$$

Constant A is set such that:  $\int_{0}^{1} dx \int_{-\infty}^{+\infty} \frac{d^{2} \boldsymbol{k}_{T}}{16\pi^{3}} |\psi_{q\bar{q}/\pi}^{V}(x, \boldsymbol{k}_{T})|^{2} = 1.$ 

[Brodsky, de Téramond et al., 2004-2015]

**Effective LFWF**  $\psi_{q\overline{q}/\pi}^{E}(x, \ \mathbf{k}_{T}) = 4\pi A \frac{\sqrt{\log\left(\frac{1}{x}\right)}}{\kappa(1-x)} e^{-\frac{\log(1/x)}{(1-x)^{2}}\frac{\mathbf{k}_{T}^{2}+m^{2}}{2\kappa^{2}}}$ 



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What we do: we fix the parameters of the model (m, K, Q<sub>0</sub>) using the experimental info on PDF and FF and calculate TMD

[Brodsky, de Téramond et al., 2004-2015]

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#### FF AND PDF OF THE PION - I

Physics Letters B 771 (2017) 546-552

$$f_1^V(x;Q_0) = A^2 e^{\left(-\frac{m^2}{\kappa^2 x} - \frac{m^2}{\kappa^2(1-x)}\right)}$$



#### $f_1^E(x;Q_0) = A^2 e^{-\frac{\log(1/x)}{(1-x)^2} \frac{m^2}{\kappa^2}}$



K. Wijesooriya, P. E. Reimer, and R. J. Holt Phys. Rev. C 72, 065203 6



#### FF AND PDF OF THE PION - II

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$$F_{\pi}^{V}(Q^{2}) = \int_{0}^{1} dx A^{2} e^{\left(-\frac{m^{2}}{\kappa^{2}x} - \frac{m^{2}}{\kappa^{2}(1-x)} - \frac{Q^{2}(1-x)}{4\kappa^{2}x}\right)}$$
 1.6

$$F_{\pi}^{E}(Q^{2}) = \int_{0}^{1} dx A^{2} e^{-\frac{\log(1/x)}{4\kappa^{2}} \left(Q^{2} + \frac{4m^{2}}{(1-x)^{2}}\right)} \qquad 0$$

0.4

S. R. Amendolia et al., Nucl. Phys. B277168 (1986) P. Braueln et al., Z. Phys. C 3,101 (1979) J. Volmer et al. Phys. Rev. Lett. 86 (2001), 1713 C. J. Bebek et al., Phys. Rev. D17, 1693 (1978)



LFWF	$m~({ m GeV})$	$\kappa ~({\rm GeV})$	$Q_0 ~({ m GeV})$	$\chi^2_{\rm d.o.f.} \left( \frac{\chi^2_{\rm FF} + \chi^2_{\rm PDF}}{N - N_{\rm par}} \right)$
$\psi^V_{q\overline{q}/\pi}$	0.005 (fixed) 0.200 (fixed) $0.0500 \pm 0.00004$	$\begin{array}{c} 0.397 \pm 0.003 \\ 0.351 \pm 0.003 \\ 0.371 \pm 0.002 \end{array}$	$0.500 \pm 0.003$ $0.491 \pm 0.003$ $0.498 \pm 0.002$	$3.15 \\ 11.76 \\ 2.25$
$\psi^E_{q\overline{q}/\pi}$	0.005 (fixed) 0.200 (fixed) 0. (fixed)	$\begin{array}{c} 0.261 \pm 0.002 \\ 0.322 \pm 0.002 \\ 0.262 \pm 0.002 \end{array}$	$0.498 \pm 0.003$ $0.630 \pm 0.008$ $0.498 \pm 0.003$	$5.44 \\ 12.96 \\ 5.38$

and the state

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Value of K systematically lower than the one obtained from fitting the Regge trajectories, but necessary to have a good description of the PDF at higher scales

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- The Effective model provides a worse description overall

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- necessary to have a good description of the PDF at higher scales
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- Initial scale Q<sub>0</sub> = 0.5 GeV is lower than the predictions of the LFH QCD (~I GeV)
- later in the "Universal" model)

Value of K systematically lower than the one obtained from fitting the Regge trajectories, but

Problem with the Valence and Effective model in reproducing the pole structure of FF (solved

### NOT IN THIS WORK: A UNIVERSAL LFWF FOR MESONS AND THE NUCLEONS

de Téramond, Liu, Sufian, Dosch, Brodsky, Deur (2018)





### NOT IN THIS WORK: A UNIVERSAL LFWF FOR MESONS AND THE NUCLEONS

de Téramond, Liu, Sufian, Dosch, Brodsky, Deur (2018)



- The pole structure of the FF is restored;
- The K parameter describes correctly the Regge trajectories, the poles in the FF pole expansion correspond to the physical ones, and the smallx behavior of the PDF is modified;
- The authors find an analytical expression for the PDF and GPD which depends on a universal function.



#### TRANSVERSE STRUCTURE - MODEL SCALE

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No spin structure: only unpolarized TMD Spin effects constructions: Boer-Mulders TMD (Ahmady, Mondal, Sandapen, 2018-2019)









## TRANSVERSE STRUCTURE - EVOLUTION

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Note: Evolution'details and prescriptions as in Bacchetta, Delcarro, Pisano, Radici, Signori, (2017)





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Position of the max and k<sub>T</sub>-broadening after evolution



## TRANSVERSE STRUCTURE - EVOLUTION

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- Effective LFWF.
- at the initial scale (0.5 GeV).
- order of magnitude.
- The Gaussian shape is lost after Evolution.
- compared to the model scale 0.5 GeV.

At the scale of the model (0.5 GeV), the TMD has a Gaussian shape for both Valence and

The mean square transverse momentum is symmetric around x=0.5 for the valence LFWF

TMD Evolution of the pion TMD, from the initial scale of 0.5 GeV to a typical experimental scale of 5 GeV, increases the width of the distributions in momentum space of almost one

The x-dependence of the transverse momentum width at 5 GeV changes drastically



### THE QCD RUNNING COUPLING FROM ADS/QCD MODELS



 $\alpha_s(Q) = \begin{cases} \alpha_{\rm LFH}(Q) & Q \le Q_0\\ \alpha_{\overline{\rm MS}}(Q) & Q > Q_0, \end{cases}$ 



### THE QCD RUNNING COUPLING FROM ADS/QCD MODELS

![](_page_27_Figure_1.jpeg)

$$Q) = \begin{cases} \alpha_{\rm LFH}(Q) & Q \le Q_0 \\ \alpha_{\overline{\rm MS}}(Q) & Q > Q_0, \end{cases}$$

$$FH(Q^2) = \alpha_{\rm LFH}(0)e^{-Q^2/4\kappa^2}$$
  
ntinuity condition:  
$$\begin{cases} \alpha_{\rm LFH}(Q_0) = \alpha_{\overline{\rm MS}}(Q_0) \\ \beta_{\rm LFH}(Q_0) = \beta_{\overline{\rm MS}}(Q_0), \end{cases}$$

![](_page_27_Picture_5.jpeg)

### THE QCD RUNNING COUPLING FROM ADS/QCD MODELS

![](_page_28_Figure_1.jpeg)

$$Q) = \begin{cases} \alpha_{\rm LFH}(Q) & Q \leq Q_0 \\ \alpha_{\overline{\rm MS}}(Q) & Q > Q_0, \end{cases}$$

FH
$$(Q^2) = \alpha_{\text{LFH}}(0)e^{-Q^2/4\kappa^2}$$
  
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• The matching with the  $\overline{MS}$  scheme with  $\kappa = 0.51 \,\mathrm{GeV}$ gives  $Q_0^2 = 0.75 \,\mathrm{GeV}^2$  (black line)

![](_page_28_Picture_6.jpeg)

![](_page_29_Figure_1.jpeg)

$$\alpha_s(Q) = \begin{cases} \alpha_{\rm LFH}(Q) & Q \leq Q_0 \\ \alpha_{\overline{\rm MS}}(Q) & Q > Q_0, \end{cases}$$

#### Relaxed continuity condition:

 $\alpha_{\rm LFH}(Q_0) = \alpha_{\overline{\rm MS}}(Q_0)$ 

![](_page_29_Picture_5.jpeg)

![](_page_30_Figure_1.jpeg)

$$\alpha_s(Q) = \begin{cases} \alpha_{\rm LFH}(Q) & Q \leq Q_0 \\ \alpha_{\overline{\rm MS}}(Q) & Q > Q_0, \end{cases}$$

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Imposing the continuity of both  $\alpha$  and the derivative  $\beta$  at the transition point  $Q_0$  is too rigid for our approach.

![](_page_30_Picture_6.jpeg)

![](_page_31_Figure_1.jpeg)

$$\alpha_s(Q) = \begin{cases} \alpha_{\rm LFH}(Q) & Q \le Q_0 \\ \alpha_{\overline{\rm MS}}(Q) & Q > Q_0, \end{cases}$$

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Fixing by hand the value of  $Q_0 \sim I$  in our fit greatly deteriorates the results on the PDF

![](_page_31_Picture_7.jpeg)

![](_page_32_Figure_1.jpeg)

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K=0.37 GeV gives a range which is compatible with previous works

![](_page_32_Picture_8.jpeg)

## SUMMARY AND CONCLUSIONS

- We study two functional forms of pion LFWFs from LF Holographic models, with minimal modifications
- We fix the free parameters of the LFWFs using the experimental information of PDF and FF (possible to update this part with the new data)
- We obtain a predictions on the pion TMD (dependence on the non perturbative parameters used);
- We test the matching between the perturbative and non-perturbative physics deriving from this approach
- Possible update: use the new available and forthcoming data on PDF to improve the study.

![](_page_33_Picture_6.jpeg)