

# Studying nucleon multiparton structure via nucleon fragmentation

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## Outline

- \* **Going beyond single parton characteristics**

**Strategy similar to one used now to probe short range correlations in nuclei:  
observe correlations through decay of the system after removal of a constituent**

**Physics of validity of leading twist dynamics in target fragmentation**

- \* ***What is known experimentally about nucleon fragmentation in hard processes***

- \* **Observables for EIC**

- \* ***Few examples of nuclear observables***

# ***Nucleon Fragmentation in DIS***

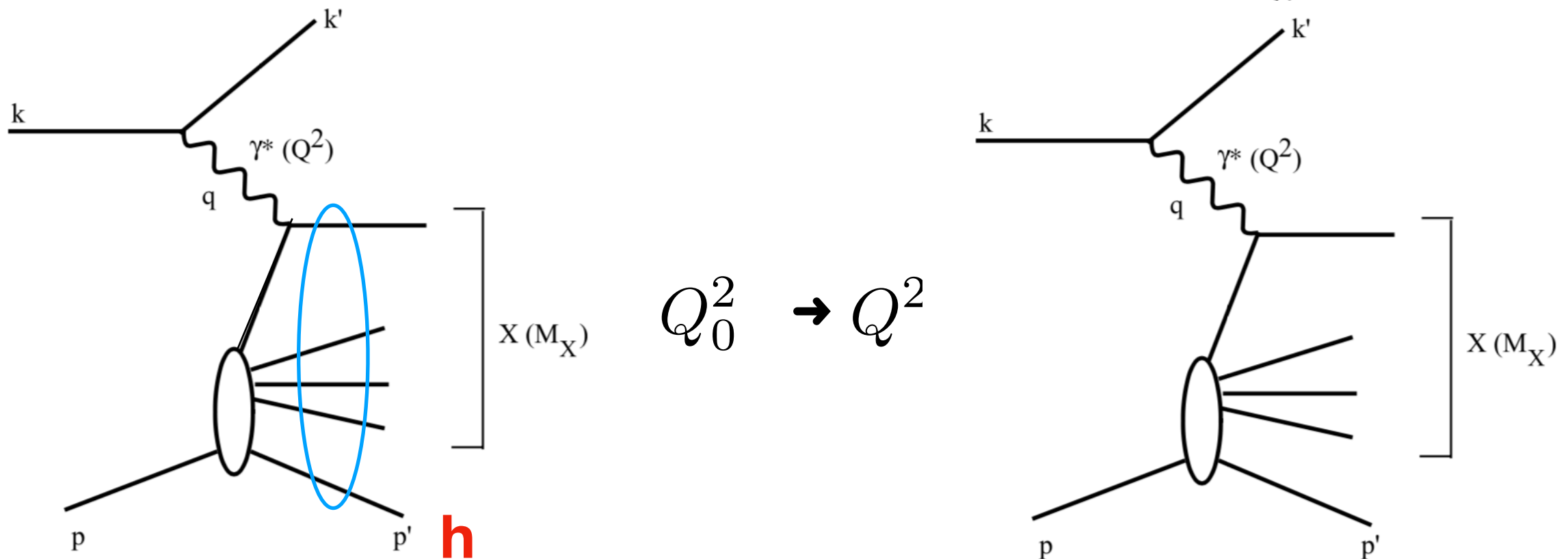
## **WHY'S**


- ❖ fragmentation is a **probe of multiparton nucleon structure** and QCD dynamics for which EIC can add a lot  
QCD  $Q^2$  evolution is simple - clear advantage as compared to TMD's
- ❖ information about fragmentation is necessary for optimizing studies of SRC - tagging, EMC effect,  $\Delta\Delta$  component of the deuteron
- ❖ Formation time of hadrons in eA scattering in nuclear fragmentation region

**Note that many requirements to detector for studying fragmentation are similar to the ones for studying of short range correlations**

Collins factorization theorem: consider hard processes like

$$\gamma^* + T \rightarrow X + T(T'), \quad \gamma^* + T \rightarrow jet_1 + jet_2 + X + T(T')$$



Interaction of partons which would form  $h$  with the rest of partons:  — does not change with  $Q^2$  since overall interaction does not resolve  $qg$  which are located at transverse distance  $\ll 1/Q_0$

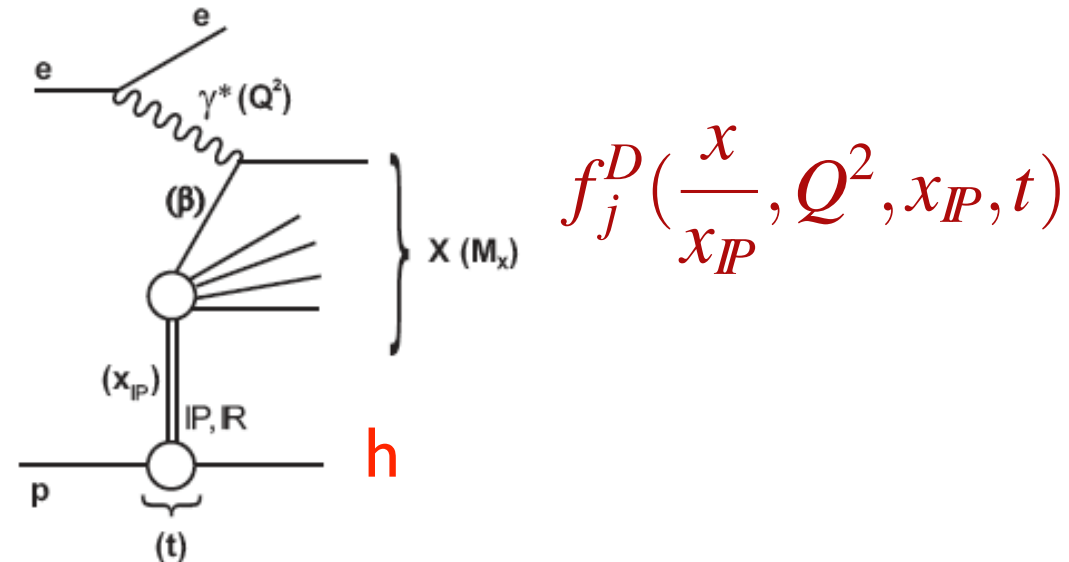
one can define fracture (Trentadue & Veneziano) parton distributions

at collider

$$x_L = p_h / p_{in} \quad x_L \leq 1 - x$$

$$x_{\mathbb{P}} = 1 - x_L$$

$$\beta \equiv x / x_{\mathbb{P}} = Q^2 / (Q^2 + M_X^2)$$



$$F_2(x, Q^2) = x \sum_{j=q, \bar{q}, g} \int_x^1 \frac{dy}{y} C_j \left( \frac{x}{y}, Q^2 \right) f_j(y, Q^2)$$

$$F_2^{D(4)}(x, Q^2, x_{\mathbb{P}}, t) = \beta \sum_{j=q, \bar{q}, g} \int_\beta^1 \frac{dy}{y} C_j \left( \frac{\beta}{y}, Q^2 \right) f_j^{D(4)}(y, Q^2, x_{\mathbb{P}}, t)$$

**Theorem:**

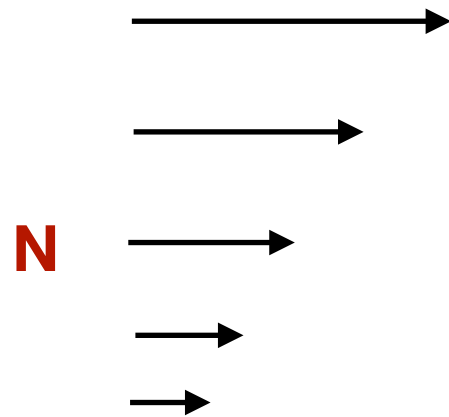
For fixed  $x_{\mathbb{P}}, t$  universal fracture pdf (same in different hard processes, but a priori depending on  $x_{\mathbb{P}}, t$ ) + the evolution is the same as for normal pdf's

*Comment:*  $x_{\mathbb{P}}$  is traditional notation due to focus on diffraction - notion of Pomeron is not necessary in the general factorization analysis

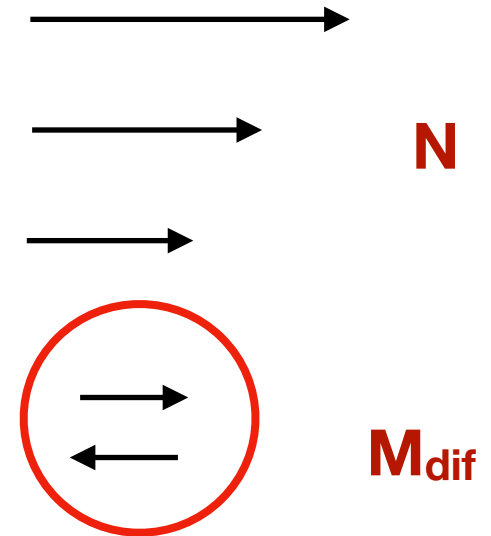
$x_{\mathbb{P}} < 0.01$  for diffraction but kinematically restricted to  $< 1$

## Example: diffraction and nucleon structure

Consider Breit frame  $q=(0, -2xp_N)$



nucleon before  
absorption of virtual photon



nucleon and diffractive mass after  
absorption of virtual photon

*small  $x$  parton has a significant probability to have a nearby parton/partons with which it can form a white cluster*

Fracture pdfs are practically not explored except fragmentation in ep scattering in

$$e + p \rightarrow e + p + X, e + p \rightarrow e + n + X$$

Need high statistics as  $f_j$  are functions of  $(x, \beta, Q^2, t)$  not only  $\beta, Q^2, t$  like for quark fragmentation functions (Current fragmentation) .

Convenient quantity  $x_L = p_h/p_p$  — light cone fraction of proton carried by  $h$

$$z = x_L/(1-x) < 1 \quad \text{Maximal } x_L = (1-x)$$

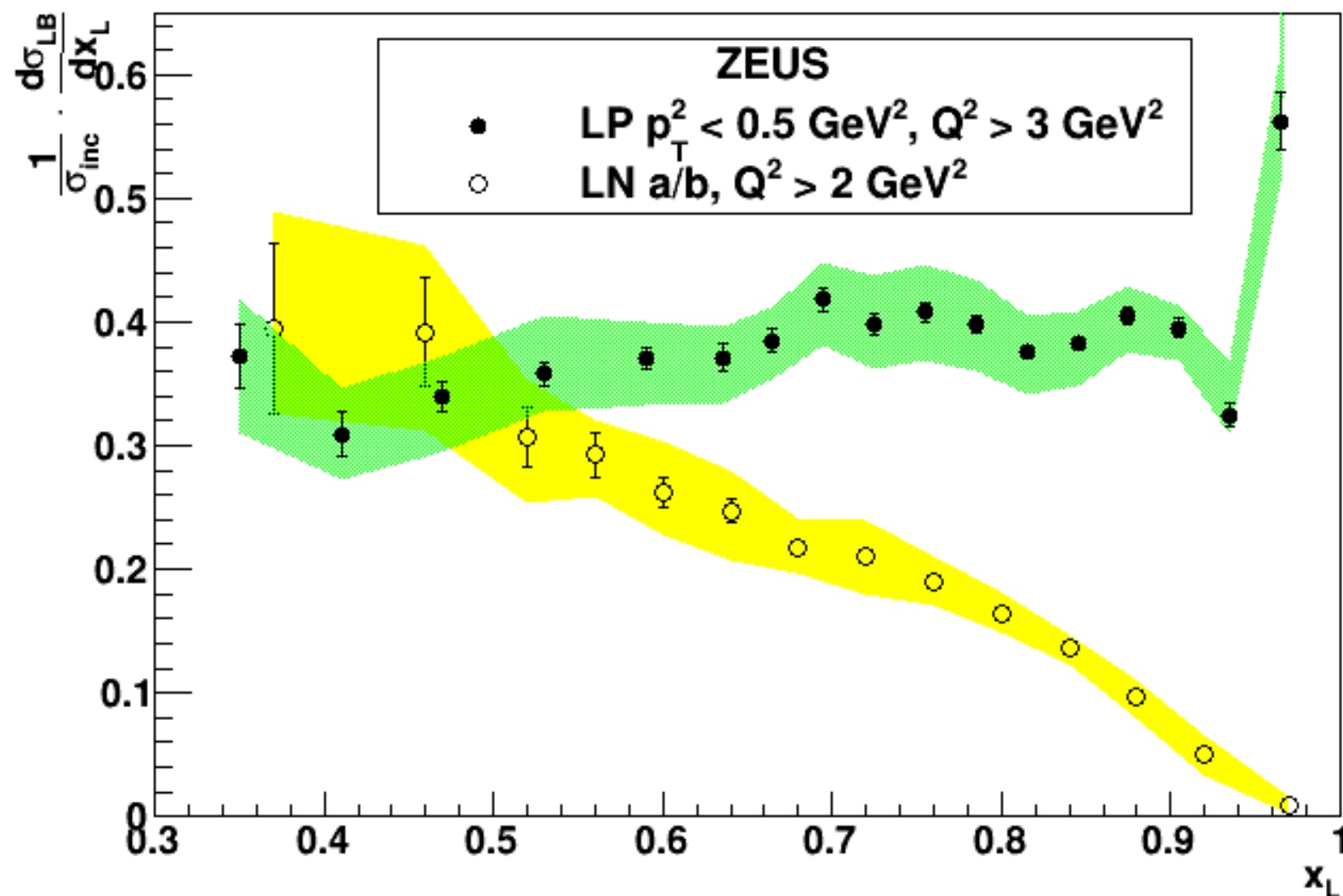
Currently except for diffraction all data are for  $x \ll 1 - z$   
integrals over  $x$  and  $\beta \ll 1$ .

Soft factorization: for  $x \ll 1$ , we expect  $z, t$ -distribution to depend weakly on  $x$  at not very large  $Q^2$

Strong dependence of leading (large  $z$ ) baryon production on  $x$  (FS77):

$$f_j(x, z) \propto (1 - z)^{n(x)}$$

$n(x < 0.01) = -1$	diffraction + flat ( $n=0$ ) at smaller $x_L$
$n(x \sim 0.1) = 0? \ 1?$	onset of sea quark dominance
$n(x \sim 0.2) = 1$	valence quarks
$n(x \sim 0.5) = 2?$	fragmentation of two quarks with large relative momenta



plot prepared  
by W. Schmidke

$r_{LP} = 0.299 \pm 0.003 \text{ (stat.)} + 0.008 - 0.007 \text{ (syst.)}$  [not shown in the paper]

$r_{LN} = 0.159 \pm 0.008 \text{ (stat.)} + 0.019 - 0.006 \text{ (syst.)}$  [as shown in the paper]

**HERA studies missed a puzzle: where are baryons. Should be**

**#baryons - # anti baryons = 1** per event.

For small  $x$  and  $x_L > 0.3$  only **0.46** baryons are observed (70% p, 30% n) (strange baryons not measured but likely 30% correction of neutrons)

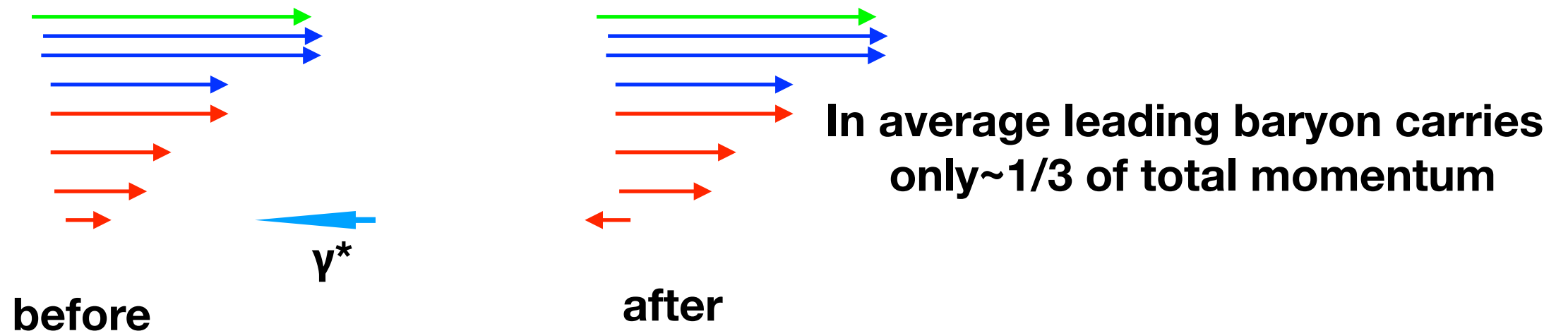
→ A lot (50%) of baryons are produced below  $x_L = 0.3$

In nucleus rest frame these baryons have large longitudinal momenta,  $p_L$

For example for  $x_L = 0.2$ ,  $p_L \sim 3 \text{ GeV}$ . Are they formed inside nucleus?



$x_{Bj}$  for these data is  $\sim 10^{-3}$ . It is highly nontrivial that a removal of a wee parton leads to a break up with large energy losses - nucleon seems to be pretty fragile



long range correlations in color?

high degree of coherence of small  $x$  partons with leading partons

*Emerging picture (small  $x$ ) from my analysis:*

leading protons  $x_L > 0.7$  — 3 valence quarks

protons  $0.7 > x_L > 0.3(?)$  — 2 valence quarks

protons  $0.3 > x_L > 0.1(?)$  — 1 valence quark

mostly protons & few neutrons

} comparable number  
of neutrons and protons

# EXPECTATIONS & OBSERVATIONS

if  $x \ll (1-x_L)$ , nucleon multiplicity for removal of (anti)quark or a gluon are the same.

**Soft factorization**

Hence no dependence of the  $x_L$  distribution on  $W$ ,

**observed**

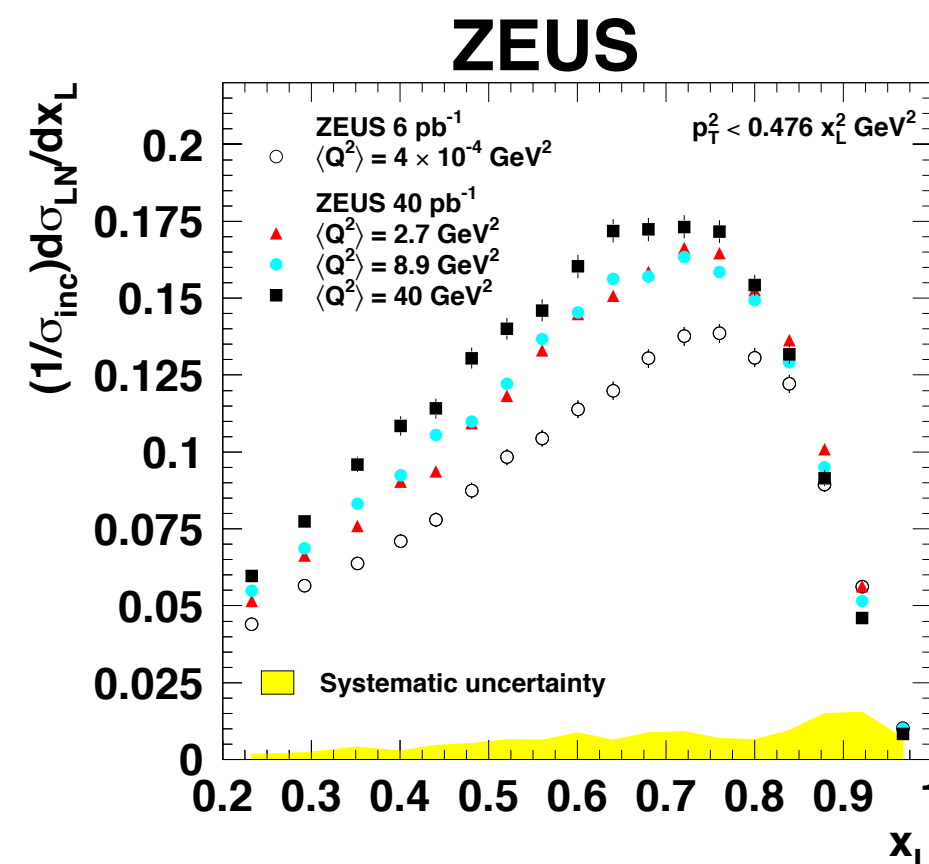
Transition from photoproduction to DIS:

disappearance of shadowing - reduction of nucleon yield at  $Q^2=0$ .

In Gribov - Regge theory presence of shadowing implies presence of a correlation between central rapidity multiplicity,  $n_h(y \sim 0)$  and nucleon yield:

larger  $n_h$  — smaller nucleon multiplicity at large  $x_L$

significant reduction for  $n_h \sim 2 \langle n_h \rangle$



There were numerous attempts to extract pion pdfs from reaction

$$ep \rightarrow en + X$$

this contribution requires approaching the pion pole which is very difficult:

$$t = -\frac{1}{x_L}(m_N^2(1-x_L)^2 + p_t^2) \quad x_L \geq 1 - m_\pi/m_N \sim 0.85, p_t \leq m_\pi$$

Space time interpretation - pion is well defined if its distance from the nucleon core is  $> 1/m_\pi$

Soft factorization leads to contributions to neutron production in all fragmentation processes

Simple example -

data:  $p/n > 2$  for  $x_L > 0.6$  (natural for fragmentation)

**Reggeon model:** pion exchange:  $n/p = 1/2$ ; need  $P'$ -exchange,  $\omega$ -exchange  $p/n \gg 1$

Warning: Regge exchange (triple reggeon limit) makes sense only for large rapidity intervals corresponds to  $1-x_L < 0.1$

For fragmentation observed :  $p/n > 2$  for  $x_L \sim 0.6$  is natural

If collective fragmentation of two valence quarks dominates at large  $z$

$$uu \rightarrow p \gg uu \rightarrow n, ud \rightarrow \frac{1}{2} p + \frac{1}{2} n$$

Probability that  $uu$ ,  $du$  diquark fragments is:  $\frac{1}{3} uu + \frac{2}{3} ud$

Hence

$$\frac{p}{n} = \frac{\frac{1}{3} + \frac{2}{3} \cdot \frac{1}{2}}{\frac{2}{3} \cdot \frac{1}{2}} = 2$$

agrees with the HERA data for  $x \sim 0.5$

## x -dependence of fragmentation

For sea quark knock out up to  $x \sim 0.1$  -- approximate matching to HERA:

$$z = x_L / (1 - x) < 1 \quad r(z) = \frac{1}{\sigma_{inc}} \frac{d\sigma_{LN}}{dz} \propto (1 - z)^{n(x)}$$

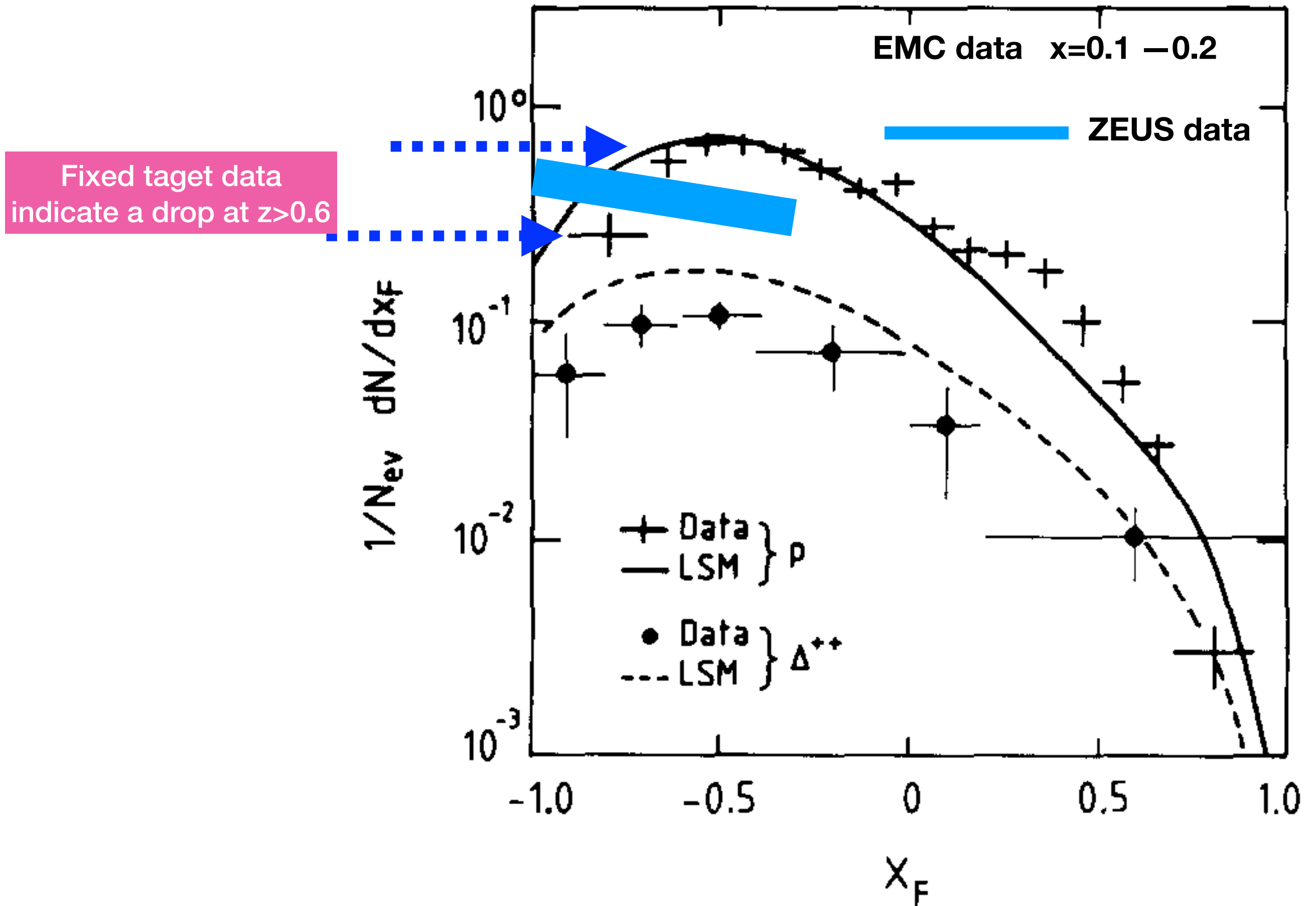
$$n_{p \rightarrow p}(x \sim 10^{-2}) \sim 0$$

$$n_{p \rightarrow n}(x < 10^{-2}) \sim -1$$

Based on our interpretation of  $p \rightarrow n$  as fragmentation of two valence quarks we expect

$$r_{p \rightarrow n}(z, x < 0.01) \propto r_{p \rightarrow p}(z, x = 0.2)$$

qualitatively consistent with the EMC data since in the EMC kinematics u-quark is knocked out and ud fragments



$W$  is not large enough to separate fragmentation and central regions for  $x_F > -0.3$  (?)

Measured total  $\langle n_{\text{tot}} \rangle$ , forward  $\langle n_{\text{f}} \rangle$ , and backward  $\langle n_{\text{b}} \rangle$  average multiplicity of p (including those from  $\Delta$  decays),  $\Delta^{++}$ , and  $\Delta^0$

	$\langle n_{\text{tot}} \rangle$	$\langle n_{\text{F}} \rangle$	$\langle n_{\text{B}} \rangle$
p	$0.62 \pm 0.01$	$0.10 \pm 0.01$	$0.52 \pm 0.01$
$\Delta^{++}$	$0.10 \pm 0.02$	$0.02 \pm 0.01$	$0.08 \pm 0.02$
$\Delta^0$	$<0.18$		

Significant  $\langle n_{\text{F}} \rangle$ : at high energies central rapidity range with equal number of baryons and anti baryons. For moderate  $z$  and moderate  $W$  (fixed target energies) no sufficient separation of current and target fragmentation regions. Significant  $\Delta^{++}$  production in the proton fragmentation region for  $\langle x \rangle \sim 0.1 \div 0.2$ :

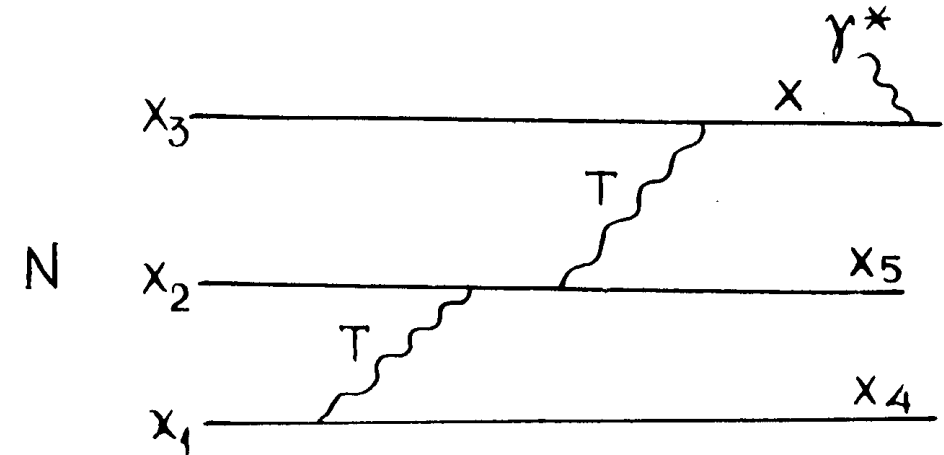
$$\Delta^{++}/p \text{ (direct)} = 1/3 \div 1/5$$

pretty large since mostly it is **ud fragmentation** since d-quark is hit in 10% of interactions

## Novel fragmentation pattern at $x > 0.5$

Expectation: **nucleon in large  $x$  configurations is smaller than in average.**

- \* In pQCD diagrams corresponding to  $x \rightarrow 1$  limit are dominated by configurations in which two spectator quarks carry very different light-cone fractions and rather large transverse momenta



- \* Analysis of centrality leading jets in pA at LHC: (Alvioli et al)  
area factor of 4 smaller than average
- ⇒ larger transverse momenta of spectators
- ⇒ enhanced probability of independent fragmentation of the spectator quarks

$$r(z) = \frac{1}{\sigma_{inc}} \frac{d\sigma_{LN}}{dz} \propto (1-z)^{n(x)} \quad \mathbf{n(x > 0.5) \geq 2}$$



*Tools available at EIC for obtaining qualitatively new information about dynamics of nucleon fragmentation, working of confinement and probing correlations in nucleons*

- ❖ **Double tagging:** detecting pions, kaon, charm, dijet in the current fragmentation region to separate processes where u, d, gluon, was involved in hard interaction.
- ❖ **Polarized ep scattering:** comparison of fragmentation for parallel and antiparallel helicities of quark and nucleon
- ❖ **Polarization/ spin alignment** of produced baryons
- ❖ **Comparison of proton and neutron fragmentation in polarized electron - polarized deuteron scattering**

## A sample of interesting channels

- ▶ Removal of u (d) quark with helicity = /opposite to helicity proton can compare fragmentation of uu and ud with helicities 0 or 1.
  - ▶ **is**  $ud \rightarrow p = ud \rightarrow n$  **not guaranteed:**
    - ▶ different proportion of  $I=0$  and  $I=1$
    - ▶ Violation of SU(6) - large for large x
- $I=0, S=0$  diquarks?**
- ▶ how  $\Delta$  isobar production / spin alignment depends on helicity of diquark
  - ▶ longitudinal polarization of hyperons:
  - ▶ **octet baryons/ decuplet baryons - rate and x dependence**
  - ▶ z-dependence of the meson production
  - ▶ expect abundant baryon production for large x including rare/exotic baryons like 20-multiplet due to large angular momenta of spectator quarks (Feynman problem).
  - ▶ meson production at large z:  $(1-z)^n$ ,  $n -2 - 4?$
  - ▶ correlations of fragmentation and central multiplicity (easier at HERA)

Summary: from discussed studies we would get a precision knowledge of how a proton wave packet evolves when a parton with given  $x$  and flavor, helicity is removed from it.

Green pasture: Removal of color octet (dijet production) vs removal a triplet for large  $z$ , and  $x > 0.1$  -

Reference point for fragmentation in pp scattering with a hard (e.g.) dijet trigger. Screening, Multiparton interactions.

*Requirements to detector:*

$x_L$  range for protons down to 0.1,  $p_T$  range:  $0 < p_T < 0.7$  GeV/c

neutrons similar requirements or use D beams with tagged protons to select scattering off neutrons

$\Delta^{++}, \Delta^0$  pions with  $x_L$  range from 0.3 to 0.1

$\Lambda$  hyperon ?  $c\tau = 7.98$  cm

**Few applications for scattering off nuclei**

*Looking for non-nucleonic degrees of freedom ( a sample of processes)*

*Coherence in production of hadrons in the nucleus fragmentation region*

# Looking for $\Delta\Delta$ admixture in the deuteron in $eD$ scattering

$$\sigma(e^2H \rightarrow e + \Delta + X) = \sigma(x' = \frac{x}{(2-\alpha)}, Q^2) \frac{\Psi_{\Delta\Delta}^2(\alpha, k_t)}{(2-\alpha)}$$

spectator  
mechanism

$$\alpha_{\Delta} = \frac{p_{\Delta}}{p_D/2}$$

EIC frame

$$\alpha_{\Delta} = \frac{\sqrt{m_{\Delta}^2 + p^2} - p_3}{m_d/2}$$

Rest frame ,p is  $\Delta$  momentum

$\alpha=1, p_t=0$  corresponds to  $p_3 \sim 300$  MeV/c forward in lab

Competing mechanism -  $\Delta$ 's from nucleons=**direct mechanism**

$$\left. \frac{\sigma^{1D/\Delta}}{dx dy \frac{d\alpha}{\alpha} d^2k_t} \right|_{\text{direct}} = \int \frac{d\beta}{\beta} d^2p_t \rho_D^N(\beta, p_t) \times \quad (18)$$

$$\times \frac{d\sigma^{1N/\Delta}}{dx dy d\alpha/\alpha d^2k_t} \left[ \beta E_1, x/\beta, y, Q^2, \frac{\alpha}{\beta - x}, k_t - \frac{\alpha}{\beta} p_t \right]$$

For scattering of stationary nucleon

$$\alpha_{\Delta} < 1 - x$$

Also there is strong suppression for production of slow  $\Delta$ 's - larger  $x$  stronger suppression

$$x_F = \frac{\alpha_{\Delta}}{1 - x} \quad \sigma_{eN \rightarrow e + \Delta + X} \propto (1 - x_F)^n, n \geq 1$$

Numerical estimate for  $P_{\Delta\Delta} = 0.4\%$

$$\frac{\sigma^{ID/\Delta}}{dx \, dy \, \frac{d\alpha}{\alpha} \, d^2k_t} \Big|_{\text{direct}} \Big/ \frac{\sigma^{ID/\Delta}}{dx \, dy \, \frac{d\alpha}{\alpha} \, d^2k_t} \Big|_{\text{spect}} < 0.1$$

Tests possible to exclude rescattering mechanism:  $\pi N \rightarrow \Delta$  FS90

For the deuteron one can reach sensitivity better than 0.1 % for  $\Delta\Delta$  especially with quark tagging (FS 80-90)

for  $x > 0.1$  very strong suppression of two step mechanisms (FS80)

is confirmed by neutrino study of  $\Delta$ -isobar production off deuteron

Best limit on probability of  $\Delta^{++}\Delta^{-}$  component in the deuteron  
 $< 0.2\%$