Vector-Meson Photoproduction @ Threshold

Igor Strakovsky* The George Washington University



J/ ψ photoproduction is sample of hard processes corresponding to relatively large scale $\mu_c \sim (0.5 - 1) * M_{J/\psi}$.

J/w is 'small size' object which can be used to study internal structure of proton (hadron), like in DIS case but now J/w feels not electric charge but gluon distribution.

*Supported by DE-SC0016583







PHYSICAL REVIEW LETTERS 123, 072001 (2019)

First Measurement of Near-Threshold J/ψ Exclusive Photoproduction off the Proton

First Measurement of Near-Threshold J/\(\psi\) Exclusive Photoproduction off the Proton

A. Ali, \(^{10}\) M. Amaryan, \(^{2}\) E. G. Anassontzis, \(^{2}\), A. Austregesilo, M. Boegin, M. Datton, Y. P. Eugeniz, M. Britton, M. Datton, Y. P. Eugeniz, M. Boegin, M. Boegin, M. Boegin, M. Boegin, M. Datton, Y. P. Eugeniz, M. Boegin, M. Boegin, M. Boegin, M. Boegin, M. Datton, J. P. Eugeniz, M. Georgi, M. Calo, J. R. Banka, J. M. Boegin, M. B

PHYSICAL REVIEW C 101, 045201 (2020)

Comparative analysis of ωp , ϕp , and $J/\psi p$ scattering lengths from A2, CLAS, and ClueX threshold measurements

Igor I. Strakovsky®, 1,* Lubomir Pentchev®, 2 and Alexander I. Titov® DC 20052, USA

**Department of Physics, Institute for Nuclear Studies, The George Washington University, 23606, USA

**Thomas Jefferson National Accelerator Facility, Newport News, Virginia 23606, USA Thomas Jefferson National Accelerator Facility, Newport News, Virginia 23606, U.

**Bogoliubov Laboratory of Theoretical Physics, JINR, Dubna 141980, Russia.

**Bogoliubov Laboratory of Theoretical Physics, JINR, Dubna 141980, Russia.



PHYSICAL REVIEW C 91, 045207 (2015)

Photoproduction of the ω meson on the proton near threshold I. I. Strakovsky, 1.* S. Prakhov, 1,2,3,† Ya. I. Azimov, 4 P. Aguar-Bartolomé, 2 J. R. M. Annand, 5 H. J. Arends, 2 K. Bantawa, 6 R. Beck, V. Bekrenev, H. Berghäuser, A. Braghieri, W. J. Briscoe, J. Brudvik, S. Cherepnya, R. F. B. Codling, S. R. F. B. Codling, S. Cherepnya, R. F. B. Codling, R K. Beck, V. Bekrenev, H. Bergnauser, A. Bragnien, W. J. Briscoe, J. Brudvik, S. Cherepnya, K. F. B. Codling, C. Collicott, 11.12 S. Costanza, B. T. Demissie, E. J. Downie, 1.2 P. Drexler, L. V. Fil'kov, 10 D. I. Glazier, 5.13 R. Gregor, R. Gr C. Collicott, "S. Costanza," B. I. Demissie, B. J. Downie, "P. Drexier," L. V. Fill kov, "D. I. Giazzier," R. Gregor, D. J. Hamilton, S. E. Heid, ^{1,2} D. Hornidge, ^{1,4} I. Jaegle, ^{1,5} O. Jahn, ² T. C. Jude, ¹³ V. L. Kashevarov, ^{2,10} I. Keshelashvili, ^{1,5} D. J. Hamilton, ² E. Heid, ^{1,2} D. Hornidge, ^{1,4} I. Jaegle, ^{1,5} O. Jahn, ² T. C. Jude, ^{1,5} V. L. Kashevarov, ^{2,10} I. Keshelashvili, ^{1,5} R. Kondratiev, ^{1,6} M. Korolija, ^{1,7} M. Kotulla, ⁸ A. Koulbardis, ⁴ S. Kruglov, ^{4,1} B. Krusche, ^{1,5} V. Lisin, ^{1,0} K. Livingston, ⁵ V. Metag, ⁸ D. G. Middleton, ^{2,14} A. Mushkarenkov, ⁹ B. M. K. Nefkens, ^{3,1} A. Nikolaev, ⁷ R. Novotny, ⁸ H. Ortega, ² D. Dadroni, ⁹ F. Dhason, ^{1,5} A. Dalanti, ^{1,6} I. Dabiason, ⁵ G. Danas, ⁵ T. Davison, ⁸ H. Ortega, ² M. Ostrick, ² C. Davison, ^{1,6} C. Davison, V. Metag., D. G. Middieton, S. A. Mushkarenkov, B. M. A. Neikens, S. A. Nikolaev, R. Novolny, H. Onega, M. Onica, P. B. Otte, B. Oussena, L. P. Pedroni, F. Pheron, J. A. Polonski, J. I. Robinson, G. Rosner, J. Rostomyan, J. S. Schumann, M. H. Gilbert, J. A. Clarostin, J. C. Charles, J. A. Clarostin, J. C. Charles, J. M. Tarbor, J. M. Ta J. Otte, B. Oussena, P. Pedroni, F. Pheron, A. Potonski, J. Robinson, G. Rosner, T. Rostomyan, S. Schuman, M. H. Sikora, J. A. Starostin, J. Supek, M. F. Taragin, C. M. Tarbert, M. Thiel, A. Thomas, M. Unverzagt, P. D. W. L. B. M. L. Rostomyan, M. Unverzagt, P. D. W. L. B. M. L. Rostomyan, M. Unverzagt, P. D. W. L. B. M. L. Rostomyan, M. Unverzagt, P. D. W. L. Rostomyan, M. L. Rostomyan

PHYSICAL REVIEW C 89, 055208 (2014)



Data analysis techniques, differential cross sections, and spin density matrix elements for the reaction $yp \rightarrow \phi p$

B. Dey, 1.* C. A. Meyer, 1 M. Bellis, 1.† M. Williams, 1.‡ K. P. Adhikari, 28 D. Adikaram, 28 M. Aghasyan, 17 M. J. Amaryan, 28 M. D. Anderson, ³⁶ S. Anefalos Pereira, ¹⁷ J. Ball, ⁷ N. A. Baltzell, ³ M. Battaglieri, ¹⁸ I. Bedlinskiy, ²¹ A. S. Biselli, ¹¹ J. Bono, ¹² S. Boiarinov, 34 W. J. Briscoe, 14 W. K. Brooks, 35,34 V. D. Burkert, 34 D. S. Carman, 34 A. Celentano, 18 S. Chandavar, 27 L. Colaneri, ¹⁹ P. L. Cole, ¹⁵ M. Contalbrigo, ¹⁶ O. Cortes, ¹⁵ V. Crede, ¹³ A. D'Angelo, ^{19,31} N. Dashyan, ³⁸ R. De Vita, ¹⁸ E. De Sanctis, ¹⁷ A. Deur, ³⁴ C. Djalali, ³³ D. Doughty, ^{8,34} M. Dugger, ⁴ R. Dupre, ³ A. El Alaoui, ³ L. El Fassi, ³ L. Elouadrhiri, ³⁴ G. Fedotov, 33 S. Fegan, 36 J. A. Fleming, 10 M. Garçon, 7 N. Gevorgyan, 38 Y. Ghandilyan, 38 G. P. Gilfoyle, 30 K. L. Giovanetti, 22 S. Pisano, 17 O. Pogorelko, 21 S. Pozdniakov, 21 J. W. Price, 5 S. Procureur, 7 D. Protopopescu, 36 A. J. R. Puckett, 9 D. Rimal, 12 M. Ripani, ¹⁸ B. G. Ritchie, ⁴ A. Rizzo, ¹⁹ P. Rossi, ^{24,17} P. Roy, ¹³ F. Sabatie, ⁷ M. S. Saini, ¹³ D. Schott, ¹² R. A. Schumacher, ¹ E. Seder, ⁹ I. Senderovich, ⁴ Y. G. Sharabian, ³⁴ A. Simonyan, ³⁸ E. S. Smith, ³⁴ D. I. Sober, ⁶ D. Sokhan, ²⁰ S. S. Stepanyan, ²³ P. Stoler, 29 I. I. Strakovsky, 14 S. Strauch, 33 V. Sytnik, 35 M. Taiuti, 27 W. Tang, 27 S. Tkachenko, 33 M. Ungaro, 9,29 B. Vernarsky, 1 A. V. Vlassov, 21 H. Voskanyan, 38 E. Voutier, 24 N. K. Walford, 6 D. P. Watts, 10 N. Zachariou, 33 L. Zana, 10 J. Zhang, 28 Z. W. Zhao,33 and I. Zonta19

PHYSICAL REVIEW C 101, 042201(R) (2020)

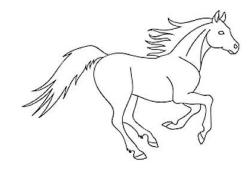
Rapid Communications

 $J/\psi p$ scattering length from GlueX threshold measurements

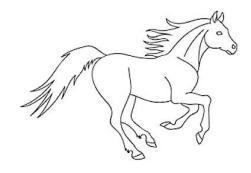
Igor I, Strakovsky O, 1,* Denis Epifanov O, 2,3 and Lubomir Pentchev O4 Department of Physics, Institute for Nuclear Studies, The George Washington University, Washington, D.C. 20052, USA ²Budker Institute of Nuclear Physics SB RAS, Novosibirsk 630090, Russia ³Novosibirsk State University, Novosibirsk 630090, Russia ⁴Thomas Jefferson National Accelerator Facility, Newport News, Virginia 23606, USA







- Vector-Meson domestic Zoo.
- Vector-Meson nucleon SL.
- Brief tour through experiments.
- σ_t fits.
- Brief tour through V-M SLs.
- Ongoing activities.
- Summary.





- Some *vector-mesons* can, compared to other mesons, be measured to very high precision.
- This stems from fact that *vector-mesons* have **same** quantum numbers as *photon*.

$$I^{G}(J^{PC}) = \mathbf{0}^{-}(\mathbf{1}^{--})$$



- Some *vector-mesons* can, compared to other mesons, be measured to very high precision.
- This stems from fact that *vector-mesons* have **same** quantum numbers as *photon*.

	IG(JPC) = 0 ⁻ (1)		
	Name	Quark	Γ	
	PDG- particle data group	Content	(MeV)	
	ρ ⁺ (770)	ud	148	
	6 0(770)	$\frac{1}{\sqrt{2}}\left(\mathrm{u}\bar{\mathrm{u}}-\mathrm{d}\bar{\mathrm{d}}\right)$	149	
	∞ (782)	$\frac{1}{\sqrt{2}} \left(u\bar{u} + d\bar{d} \right)$	8.5	
	K*+(892)	us	51	
	K*0(892)	ds	47	
BROOKHAVEN NATIONAL LABORATORY	ø (1020)	ss	4.3	
	D*+(2010)	$c\overline{d}$	0.083	Open Ch
	D*0(2007)	cu	< 2.1	•
	J/ ψ(1S)(3097) cc	0.093	Charmor
	Y (1S)(9460)	b b	0.052	Quarkon

- Some *vector-mesons* can, compared to other mesons, be measured to very high precision.
- This stems from fact that *vector-mesons* have **same** quantum numbers as *photon*.

$I^{G}(J^{PC})=0^{-}(1^{})$

$ar{q}q$ I	Vonet D_s^{*+}
D	The second second
2-/-	K^{*0} $d\bar{s}$ ρ^0 ω $u\bar{s}$
ρ $d\bar{u}$	$s\bar{u} J/\psi \phi s\bar{d} u\bar{d} \rho^+$
	K^{*0} D^{*0}
	$D_{\underline{s}}^{*-}$

	Name	Quark	Γ	
	PDG- particle data group	Content	(MeV)	
	$\rho^{+}(770)$	ud	148	
	~ (770)	$rac{1}{\sqrt{2}}\left(uar{u}-dar{d} ight)$	149	
20	<u></u> <i>∞</i> (782)	$\frac{1}{\sqrt{2}} \left(u\bar{u} + d\bar{d} \right)$	8.5	
	K*+(892)	us	51	
	K* ⁰ (892)	ds	47	
BROOKHAVEN NATIONAL LABORATORY	ø (1020)	ss	4.3	
	D*+(2010)	cd	0.083	Open Char
	D*0(2007)	cu	< 2.1	
	J/ ψ(1S)(3097	CC	0.093	Charmonii
	Y (1S)(9460)	bb	0.052	Quarkoniu

• I will focus on 3 vector-mesons from $\overline{q}q$ *Nonet* which widths are **narrow** enough to study meson photoproduction @ threshold & where data are available.



Vector-Meson - Mucleon SI





Vector-Meson – Nucleon SL Determination

IS, L. Pentchev, & A.I. Titov, Phys Rev C 101 (2020)

 For evaluation of absolute value of VN SL, we apply VMD approach that links near-threshold photoproduction Xsections of $\gamma p \rightarrow Vp$ & elastic $Vp \rightarrow Vp$



$$\frac{d\sigma^{\gamma p \to Vp}}{d\Omega}|_{\rm thr} = \frac{q}{k} \frac{1}{64\pi} |T^{\gamma p \to Vp}|^2 = \frac{q}{k} \cdot \frac{\pi \alpha}{g_V^2} \frac{d\sigma^{Vp \to Vp}}{d\Omega}|_{\rm thr} = \frac{q}{k} \cdot \frac{\pi \alpha}{g_V^2} |\alpha_{Vp}|^2$$

k is **photon** CM momentum $k = (s - M^2) / 2 s^{1/2}$ q is vector-meson CM momentum

T^{yp→Vp} is the invariant amplitude of *vector-meson* photoproduction α is fine-structure constant



 g_v is VMD coupling constant, related to vector-meson EM decay width $\Gamma_{v\rightarrow e+e}$

$$g_V = \sqrt{\frac{\pi \alpha^2 m_V}{3\Gamma_{V \to e^+ e^-}}}$$

Vector-Meson — Nucleon SL Determination

IS, L. Pentchev, & A.I. Titov, Phys Rev C 101 (2020)

For evaluation of absolute value of VN SL,
 we apply VMD approach that
 links near-threshold photoproduction Xsections of γp→Vp & elastic Vp→Vp



$$\frac{d\sigma^{\gamma p \to Vp}}{d\Omega}|_{\rm thr} = \frac{q}{k} \frac{1}{64\pi} |T^{\gamma p \to Vp}|^2 = \frac{q}{k} \cdot \frac{\pi \alpha}{g_V^2} \frac{d\sigma^{Vp \to Vp}}{d\Omega} |_{\rm thr} = \boxed{\frac{q}{k}} \cdot \frac{\pi \alpha}{g_V^2} |\alpha_{Vp}|^2$$

k is **photon** CM momentum $k = (s - M^2) / 2 s^{1/2}$

q is vector-meson CM momentum

 $\mathsf{T}^{\gamma p \to \mathsf{V} p}$ is the invariant amplitude of *vector-meson* photoproduction α is fine-structure constant



 ${f g}_{{f v}}$ is VMD coupling constant, related to ${\it vector-meson}$ EM decay width $\Gamma_{{f v}
ightarrow {f e}+{f e}-{f e}}$

$$g_V = \sqrt{\frac{\pi \alpha^2 m_V}{3\Gamma_{V \to e^+ e^-}}}$$

 Finally, one can express absolute value of SL as product of pure EM VMD-motivated kinematic factor

$$R_V^2 = \alpha m_V k/12\pi \Gamma_{V \to e^+e^-}$$
 & $h_{Vp} = \sqrt{b_1}$,

where b_1 came from best fit $\sigma_t(q) = b_1 q + b_3 q^3 + b_5 q^5$,

that is determined by interplay of strong (hadronic) & EM dynamics as

$$|\alpha_{Vp}| = R_V h_{Vp}$$





Vector-Meson — Nucleon SL Determination

IS, L. Pentchev, & A.I. Titov, Phys Rev C 101 (2020)

For evaluation of absolute value of VN SL,
 we apply VMD approach that
 links near-threshold photoproduction Xsections of γp→Vp & elastic Vp→Vp



$$\frac{d\sigma^{\gamma p \to Vp}}{d\Omega}|_{\rm thr} = \frac{q}{k} \frac{1}{64\pi} |T^{\gamma p \to Vp}|^2 = \frac{q}{k} \cdot \frac{\pi \alpha}{g_V^2} \frac{d\sigma^{Vp \to Vp}}{d\Omega}|_{\rm thr} = \frac{q}{k} \cdot \frac{\pi \alpha}{g_V^2} |\alpha_{Vp}|^2$$

k is **photon** CM momentum $k = (s - M^2) / 2 s^{1/2}$

q is vector-meson CM momentum

 $\mathsf{T}^{\mathsf{vp} \to \mathsf{vp}}$ is the invariant amplitude of *vector-meson* photoproduction





 ${f g}_{{f v}}$ is VMD coupling constant, related to ${\it vector-meson}$ EM decay width ${f \Gamma}_{{f v}
ightarrow {f e}+{f e}-{f e}}$

$$g_V = \sqrt{\frac{\pi \alpha^2 m_V}{3\Gamma_{V \to e^+ e^-}}}$$

 Finally, one can express absolute value of SL as product of pure EM VMD-motivated kinematic factor

$$R_V^2 = \alpha m_V k/12\pi \Gamma_{V \to e^+e^-}$$
 & $h_{Vp} = \sqrt{b_1}$,

where b_1 came from best fit $\sigma_t(q) = b_1 q + b_3 q^3 + b_5 q^5$,

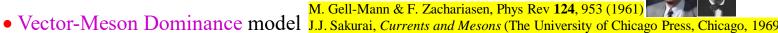
that is determined by interplay of strong (hadronic) & EM dynamics as

$$|\alpha_{Vp}| = R_V h_{Vp}$$

- To avoid theoretical uncertainties, we did not
 - determine sign of SL,
 - separate Re & Im parts of SL,
- extract Isospin 1/2 & 3/2 contributions.



VMD Approach

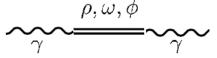


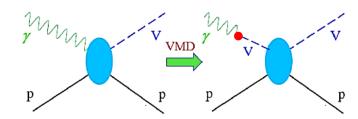


relying on transparent current-field identities N.M. Kroll, T.D. Lee, & B. Zumino, Phys. Rev. 157, 1376 (1967)



• In VMD, real photon can fluctuate into virtual vector-meson, which subsequently scatters off target proton.





• VMD does not contain *free parameters* & can be used for variety of qualitative estimates of observables in vector-meson photoproductions @ least as first step towards their more extended theoretical studies.

VMD Approach: EM Factor

VMD coupling constant $g_V = \sqrt{\frac{\pi \alpha}{3\Gamma_{V \to e^+e}}}$ $R_V^2 = \alpha m_V k / 12\pi \Gamma_{V \to e^+e^-}$ EM factor PDG-V $\Gamma_{
m v
ightarrow e+e-}$ $\mathbf{R}_{\mathbf{V}}$ $m_{\mathbf{V}}$ $\mathbf{g}_{\mathbf{V}}$ $(MeV^{1/2})$ (MeV) (keV) 390.5±6.4 782.65 0.60 ± 0.02 8.53 ± 0.14 ω 342.5±5.3 1019.461 1.27 ± 0.04 6.69±0.10 454.9 ± 4.1 J/ψ 3096.916 5.55±0.11 5.58 ± 0.07 2655.0±162.2 9460.30 1.340 ± 0.018 19.84 ± 0.14

• EM factor R_V for each vector-meson are close to each other.



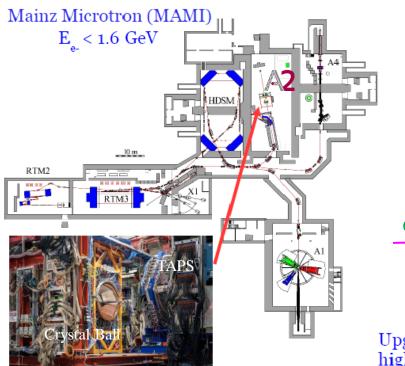
Brief Tour through Experiments



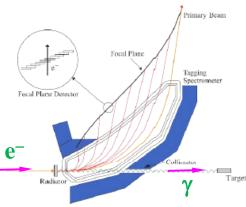








Tagger/End point tagger



$$\mathbf{E}_{_{\boldsymbol{\gamma}}} = \mathbf{E}_{_{\boldsymbol{e}^{\perp}}} - \mathbf{E}_{_{\boldsymbol{tagg}}}$$

Upgrade → experiments with ~4 times higher rates will be possible!

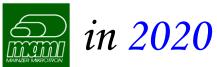
- High-Flux, Tagged, Bremsstrahlung Photon Beam: Unpolarized, Linear, and Circular
- Polarized and Unpolarized Targets
- Recoil polarimeter











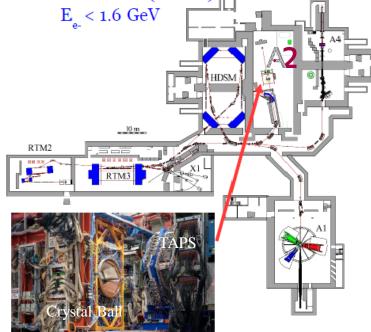
b quark physics @

for physics of hadrons &

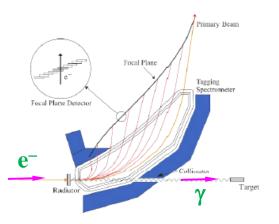
it is @ now.

Then worked @ BROOKHAVEN





Tagger/End point tagger



$$\mathbf{E}_{y} = \mathbf{E}_{e} - \mathbf{E}_{tagg}$$

Upgrade → experiments with ~4 times higher rates will be possible!

- High-Flux, Tagged, Bremsstrahlung Photon Beam: Unpolarized, Linear, and Circular
- Polarized and Unpolarized Targets
- Recoil polarimeter



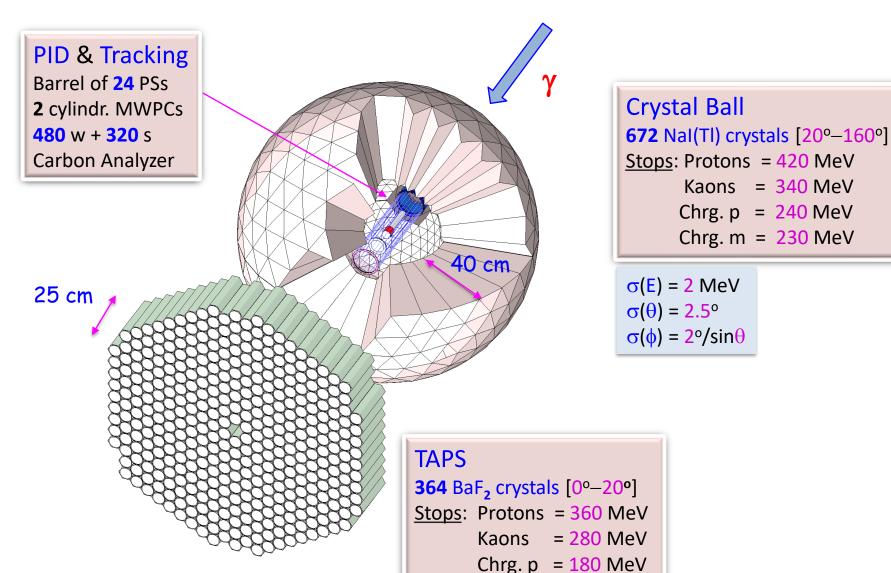








Crystal Ball, TAPS, & Tracking

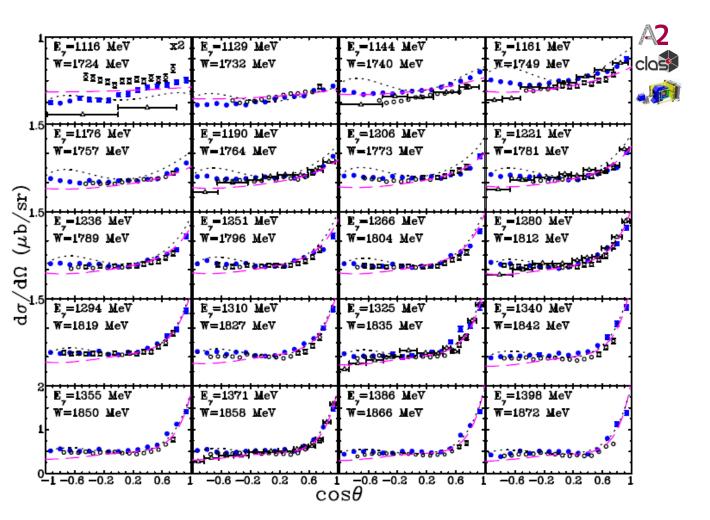






$\gamma p \longrightarrow \omega p \longrightarrow \pi^0 \gamma p \longrightarrow 3 \gamma p$ Measurements

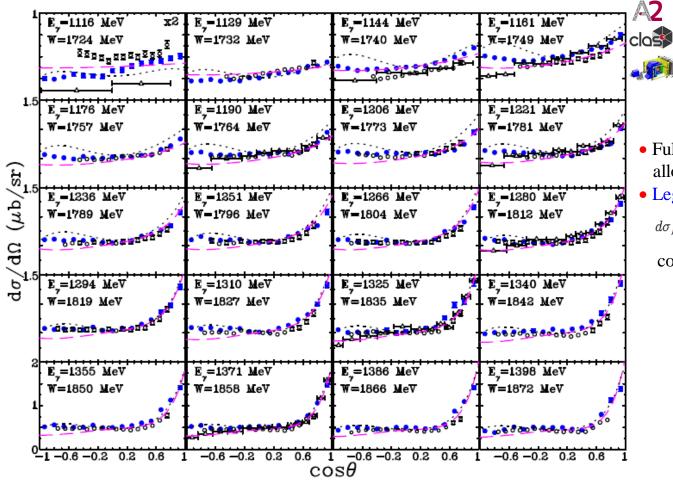
IS, S. Prakhov, Ya. Azimov *et al*, Phys Rev C **91**, 045207 (2015)

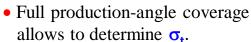


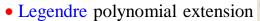


$\gamma p \longrightarrow \omega p \longrightarrow \pi^0 \gamma p \longrightarrow 3 \gamma p$ Measurements

IS, S. Prakhov, Ya. Azimov et al, Phys Rev C 91, 045207 (2015)









 $d\sigma/d\Omega(E_{\gamma},\cos\theta) = \sum A_{j}(E_{\gamma})P_{j}(\cos\theta)$ confirms of determination

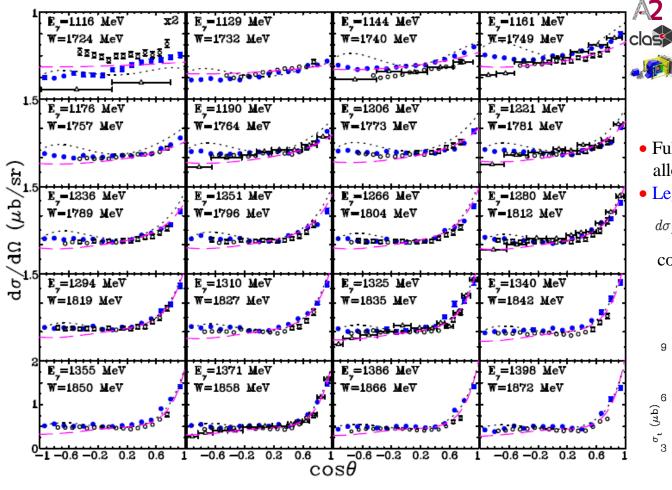
$$\sigma_t = 4\pi A_0(E_\gamma)$$





$\gamma p \longrightarrow \omega p \longrightarrow \pi^0 \gamma p \longrightarrow 3\gamma p$ Measurements

IS, S. Prakhov, Ya. Azimov et al, Phys Rev C 91, 045207 (2015)

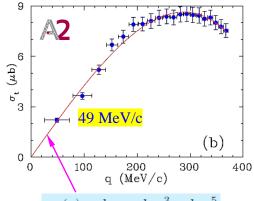


- Full production-angle coverage allows to determine σ_t.
- Legendre polynomial extension

$$d\sigma/d\Omega(E_{\gamma},\cos\theta) = \sum_{j=0} A_{j}(E_{\gamma})P_{j}(\cos\theta)$$

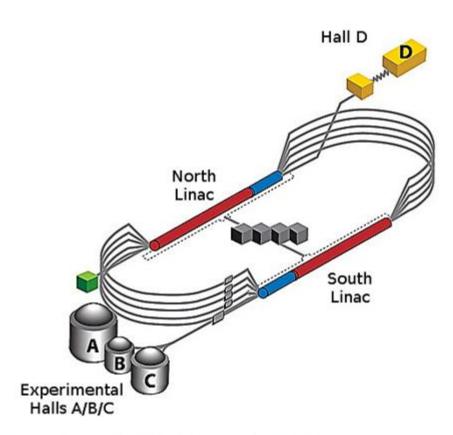
confirms σ_{\bullet} determination

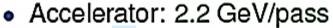
$$\sigma_t = 4\pi A_0(E_\gamma)$$



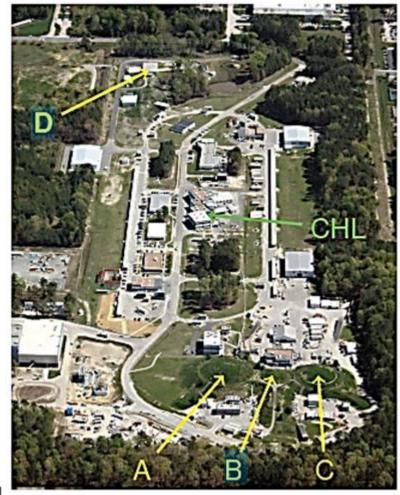
$$\sigma_t(q) = b_1 q + b_3 q^3 + b_5 q^5$$

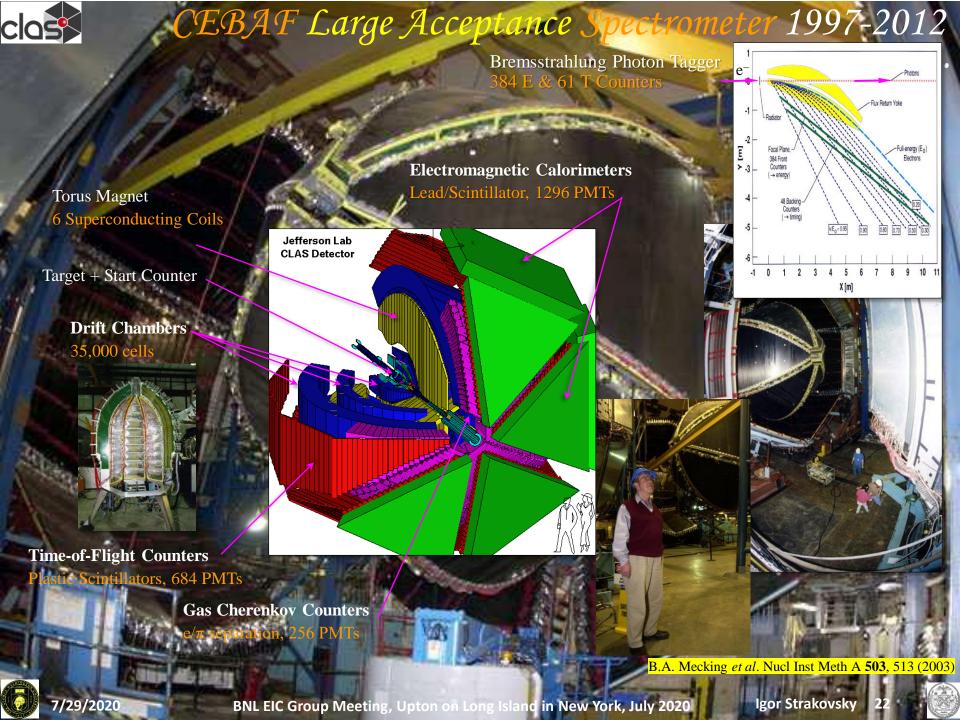
Jefferson Lab in 2020





- Halls A,B,C: e⁻ 1-5 passes ≤11 GeV
- Hall D: e^- 5.5 passes 12 GeV $\Rightarrow \gamma$ -beam
- Runs 2017-2018: 5.5 passes 11.7 GeV

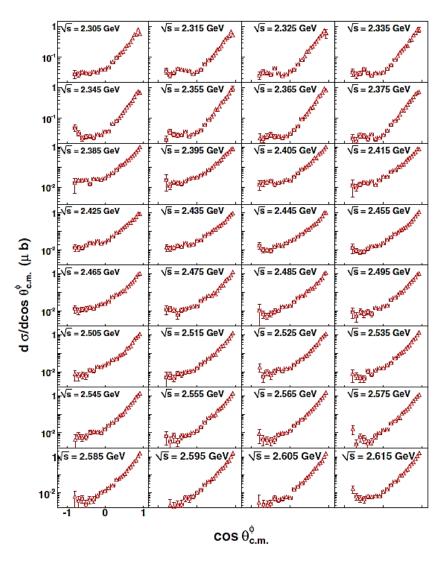






B. Dey et al, Phys Rev C 89, 055208 (2014)



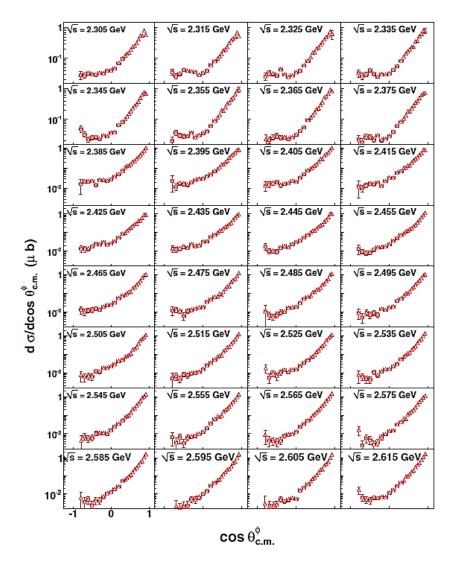




$\gamma p \rightarrow \phi p \rightarrow K^+K^-p$ Measurements

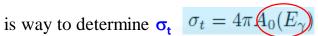
B. Dey et al, Phys Rev C 89, 055208 (2014)





- $\cos\theta$ of cosp spans from -0.80 to 0.93.
- Legendre polynomial extension

$$d\sigma/d\Omega(E_{\gamma},\cos\theta) = \sum_{j=0} A_{j}(E_{\gamma})P_{j}(\cos\theta)$$



IS, L. Pentchev, & A.I. Titov, Phys Rev C 101, 045201 (2020)

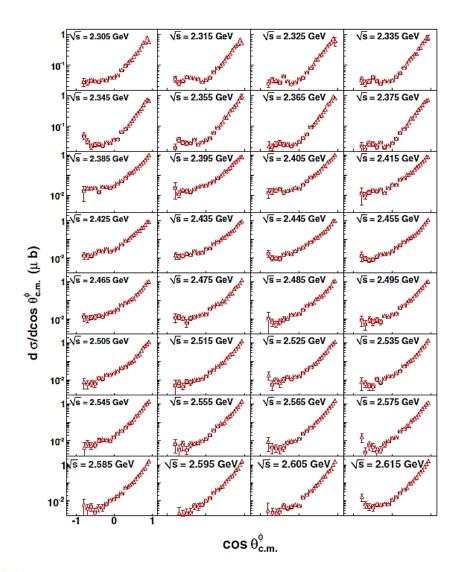




$\gamma p \rightarrow \phi p \rightarrow K^+K^-p$ Measurements

B. Dey et al, Phys Rev C 89, 055208 (2014)



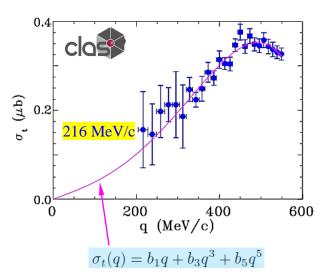


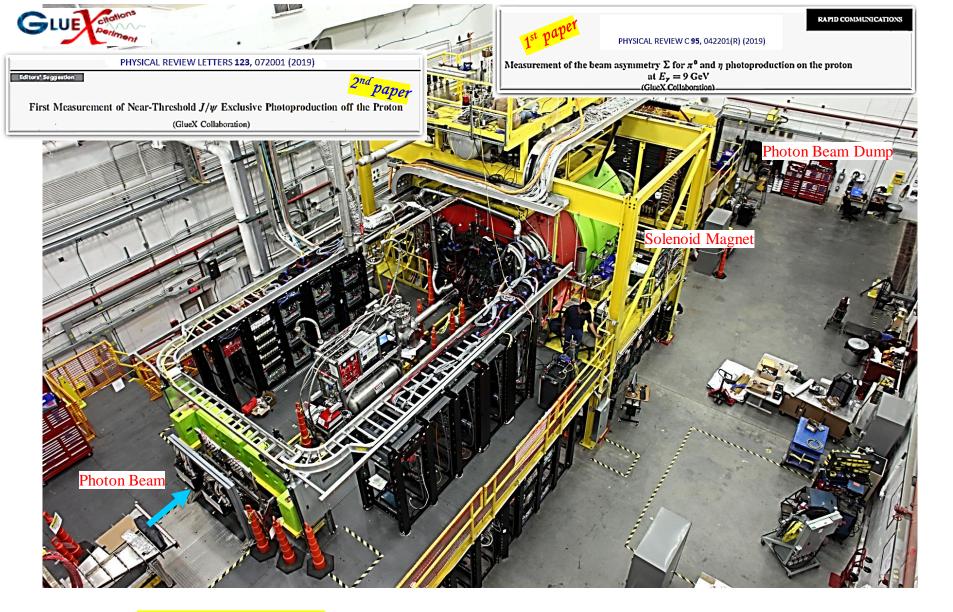
- $\cos\theta$ of closes spans from -0.80 to 0.93.
- Legendre polynomial extension

$$d\sigma/d\Omega(E_{\gamma},\cos\theta) = \sum_{j=0} A_j(E_{\gamma})P_j(\cos\theta)$$



is way to determine σ_t $\sigma_t = 4\pi A_0(E_\gamma)$







Statistics: Spring 2016: 10 pb⁻¹

25% of total statistics (2016-2018) up to date.



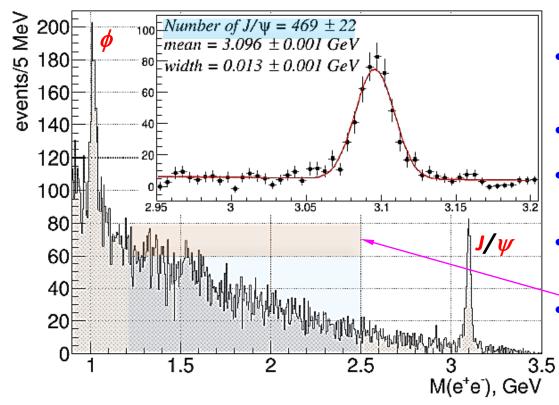




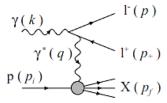
e⁺e⁻ Invariant Mass Spectrum

A. Ali *et al*, Phys Rev Lett **123**, 072001 (2019)





- Tagged photon beam (0.2% energy resolution) & exclusivity reaction: $\gamma p \rightarrow \mathcal{I}/\psi p \rightarrow e^+e^-p$.
- BR(J/ ψ \rightarrow e⁺e⁻) = (5.971±0.032)%.
- Pion contamination ~50% in continuum (using E/p fits to estimate it).
- Kinematic fit (constrained mostly recoil proton):
 13 MeV mass resolution; no rad tail.
- BH [M(e⁺e⁻) = 1.2 2.5 GeV] used for normalization.



• Overall **normalization** error is **27**%.





Total X section for $\gamma p \rightarrow J/\psi p \rightarrow e^+e^-p @ J/\psi Threshold$

A. Ali et al, Phys Rev Lett 123, 072001 (2019)

 $\sigma_{J/\psi}(E_{\gamma}) = \frac{N_{J/\psi}(E_{\gamma})}{N_{BH}(E_{\gamma})} \frac{\sigma_{BH}(E_{\gamma})}{\mathcal{B}_{J/\psi}} \frac{\varepsilon_{BH}(E_{\gamma})}{\varepsilon_{J/\psi}(E_{\gamma})}$

• SLAC experiment measured $d\sigma/dt$ at $t = t_{min}$ as function of E_{γ} .
• To determine σ_{i} from SLAC data, G_{min} used dipole t-dependence.



Total X section for $\gamma p \rightarrow J/\psi p \rightarrow e^+e^-p @ J/\psi Threshold$

A. Ali *et al*, Phys Rev Lett **123**, 072001 (2019)

$$\sigma_{J/\psi}(E_{\gamma}) = \frac{N_{J/\psi}(E_{\gamma})}{N_{BH}(E_{\gamma})} \frac{\sigma_{BH}(E_{\gamma})}{\mathcal{B}_{J/\psi}} \frac{\varepsilon_{BH}(E_{\gamma})}{\varepsilon_{J/\psi}(E_{\gamma})}$$

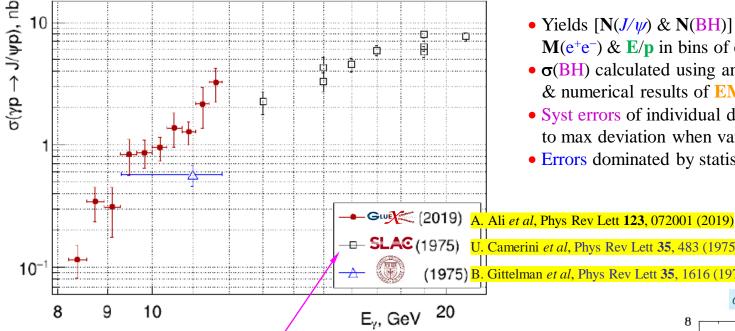
- Yields $[N(J/\psi) \& N(BH)]$ extracted from fits of $M(e^+e^-)$ & E/p in bins of energy.
- $\sigma(BH)$ calculated using analytical & numerical results of EM tree level diagrams.
- Syst errors of individual data points assigned to max deviation when varying fitting methods.

 $\gamma p \rightarrow J/\psi p$

(qu)

230 MeV

• Errors dominated by statistics.



- U. Camerini *et al*, Phys Rev Lett **35**, 483 (1975)

(1975) B. Gittelman et al, Phys Rev Lett 35, 1616 (1975)

IS, D. Epifanov, & L. Pentchev, Phys Rev C **101**, 042201 (2020)

SLAC experiment measured $d\sigma/dt$ at $t = t_{min}$ as function of E_v .

• To determine σ, from SLAC data, Guilland used dipole t-dependence.

a_i	GLUE CHONS	GLUE & SLAC
$a_1 [nb/(GeV/c)]$	0.46 ± 0.16	0.53 ± 0.12
$a_3 [nb/(GeV/c)^3]$	0.83 ± 0.91	0.78 ± 0.16
$a_5 [nb/(GeV/c)^5]$	0.28 ± 0.87	-0.06 ± 0.03
χ²/dof	0.67	0.98

• There is no discrepancy between GLUE & SLAC data.

7/29/2020





q (GeV/c)

 $\sigma_t(q) = b_1 q + b_3 q^3 + b_5 q^5$







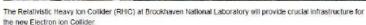
Electron Ion Collider Receives CD-0 Approval

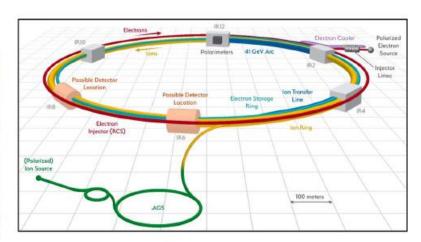
- EIC Panel evaluated proposals from JLAB and BNL (Aug Oct, 2019)
- CD-0 approved Dec. 19, 2019
- DOE announced selection of Brookhaven National Lab to host EIC Jan. 9, 2020

U.S. Department of Energy Selects Brookhaven National Laboratory to Host Major New Nuclear **Physics Facility**

January 9, 2020









7/29/2020

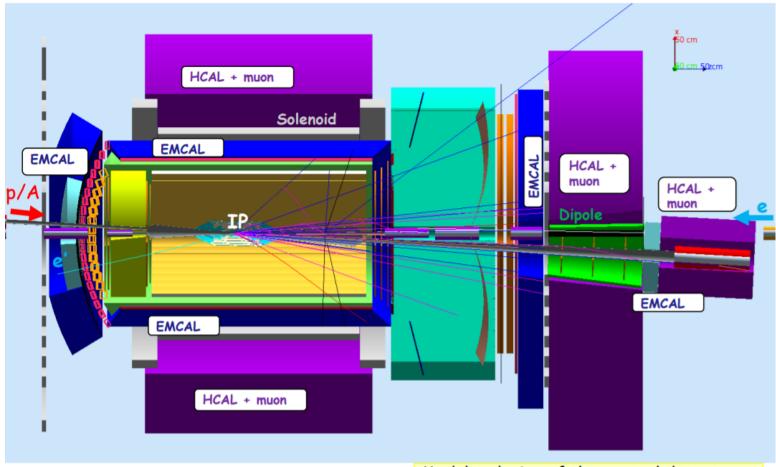






in 2020

EIC Central detector overview



Modular design of the central detector

Yulia Furletova

47

Courtesy of Yulia Furletova, June 2020



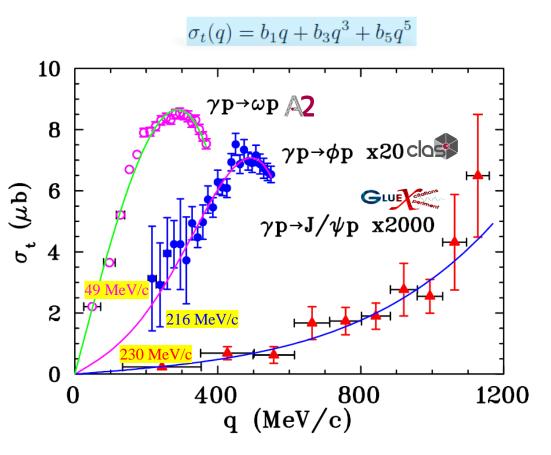
Total Cross Section Fils





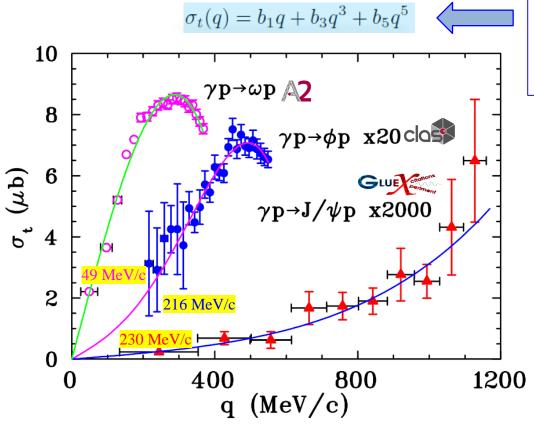
Total Cross Sections for Vector-Meson Photoproduction off Proton

• Traditionally, σ_t behavior of near-threshold binary inelastic reaction $m_a + m_b < m_c + m_d$ is described as series of odd powers in q(even powers in case of elastic).



Total Cross Sections for Vector-Meson Photoproduction off Proton

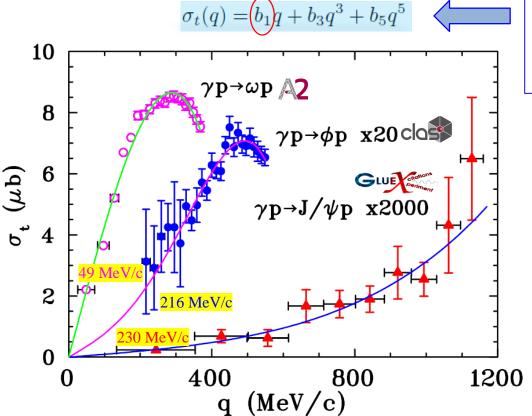
• Traditionally, σ_t behavior of near-threshold binary inelastic reaction $m_a + m_b < m_c + m_d$ is described as series of odd powers in q(even powers in case of elastic).



- Linear term is determined by two independent S-waves only with total spin 1/2 &/or 3/2.
- Contributions to cubic term come from both P-wave amplitudes & W dependence of S-wave amplitudes,
- Fifth-order term arises from *D*-waves &
 W dependencies of *S* & *P*-waves.

Total Cross Sections for Vector-Meson Photoproduction off Proton

• Traditionally, σ_t behavior of near-threshold binary inelastic reaction $m_a + m_b < m_c + m_d$ is described as series of odd powers in q (even powers in case of elastic).



- Linear term is determined by two independent S-waves only with total spin 1/2 &/or 3/2.
- Contributions to cubic term come from both P-wave amplitudes & W dependence of S-wave amplitudes,
- Fifth-order term arises from *D*-waves & W dependencies of *S* & *P*-waves.

$$b_1 = (4.42\pm0.14) \times 10^{-2} \, \mu b/(MeV/c)$$
IS, S. Prakhov, Ya. Azimov *et al*, Phys Rev C **91**, 045207 (2015)
$$b_1 = (3.40\pm1.15) \times 10^{-4} \, \mu b/(MeV/c)$$

 $b_1 = (3.40 \pm 1.15) \times 10^{-4} \,\mu\text{b/(MeV/c)}$ IS, L. Pentchev, & A.I. Titov, Phys Rev C **101**, 045201 (2020)

GLUE
$$b_1 = (0.46\pm0.16) \times 10^{-6} \mu b/(MeV/c)$$

IS, D. Epifanov, & L. Pentchev, Phys Rev C **101**, 042201 (2020)

• Dramatic differences in hadronic factors $\mathbf{h_{Vp}} = (\mathbf{b_1})^{1/2},$ as slopes (b_1) of σ_t @ threshold as function of q varies significantly from ω to ϕ to J/ψ .

• Therefore, such big difference in SL is determined mainly by hadronic factor h_{Vp} .



Brief Tour through S.C.



What is Known for ωNSL



 $|\alpha_{\omega p}| = (0.82 \pm 0.03)$ fm from phenomenology

IS, S. Prakhov, Ya. Azimov et al, Phys Rev C 91, 045207 (2015





 $\alpha_{\omega p} = (-0.97 \pm 0.16 \pm 0.03 + \mathbf{i}0.07 \pm 0.15 \pm 0.19)$ fm from phenomenology T. Ishikawa *et al*, Phys Rev C **101**, 052201(R) (2020)



 $\alpha_{\omega p} = (-0.026 + i0.28)$ fm from coupled-channel analysis of ω production in $\pi N \& \gamma N$ V. Shklyar et al, Phys Rev C 71, 055206 (2005)



 $\alpha_{\omega\rho} = (-0.41 \pm 0.05)$ fm from the QCD sum-rule

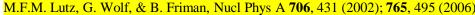
Y. Koike & A. Hayashigaki, Prog Theor Phys 98, 631 (1997)

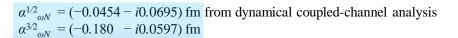


F. Klingl, T. Waas, & W. Weise, Nucl Phys A 650, 299 (1999)



 $\alpha_{up} = (-0.44 + i0.20)$ fm from coupled-channel unitary approach





M.W. Paris, Phys Rev C 79, 025208 (2009)



 $|\alpha_{\omega N}| = (0.81 \pm 0.41)$ fm from meson-nucleus optical potential

E. Friedman & A. Gal, Phys Rep 452, 89 (2007)

• That turned out to be in agreement, within uncertainties, with estimate made using 2 & o data.



• Theoretical approaches lead to ImSL << ReSL.



What is Known for ϕNSL



 $|\alpha_{\phi p}| = (0.063 \pm 0.010)$ fm from phenomenology

IS, L. Pentchev, & A.I. Titov, Phys Rev C 101, 045201 (2020)



0.15 fm from forward coherent f-meson photoproduction from deuterons near threshold

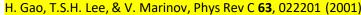
W.C. Chang et al, Phys Lett B 658, 209 (2008



 (-0.15 ± 0.02) fm from QCD sum rule analysis on spin-isospin averaged ρ , ω , & ϕ meson-N scattering

Y. Koike & A. Hayashigaki, Prog Theor Phys 98, 631 (1997)

~2.37 fm from QCD van der Waals attractive N potential for analysis of nucleus bound states





 This value is more than order of magnitude greater than results using exp data & provides problem for this particular potential model.





A.I. Titov et al, Phys Rev C 76, 048202 (2007)

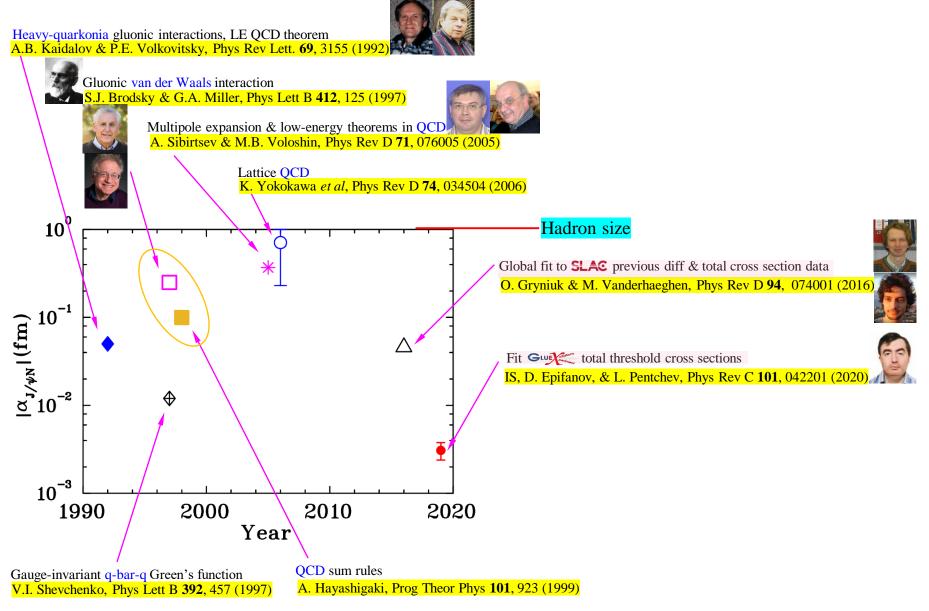








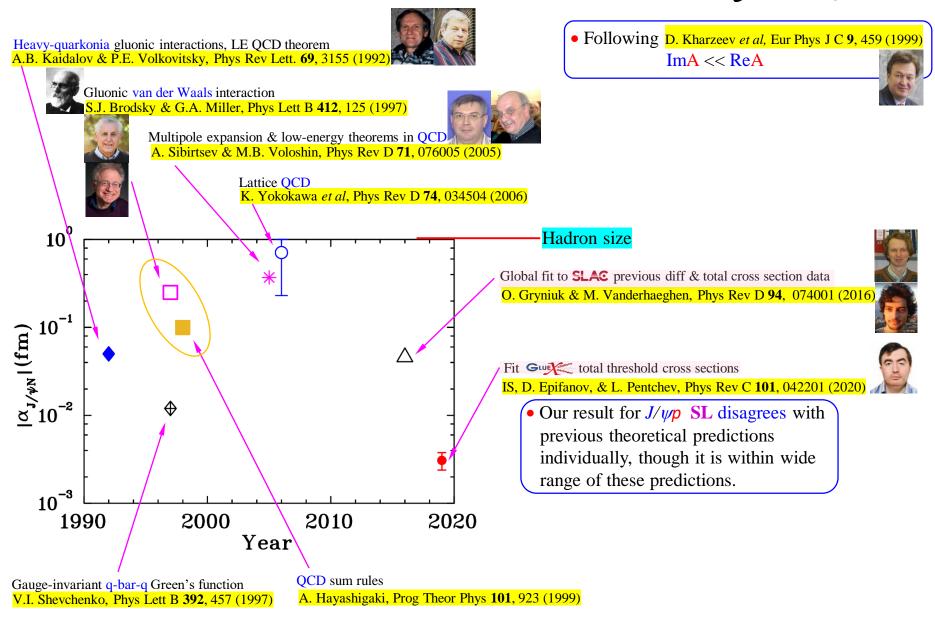
What is Known for J/\psi SL







What is Known for J/\psi SL







• There is no alternative VMD to get $J/\psi p$ SL from meson photoproduction.

Courtesy of A. Vainshtein & M. Ryskin, July 2020



• There is no alternative VMD to get $J/\psi p$ SL from meson photoproduction.

Courtesy of A. Vainshtein & M. Ryskin, July 2020

To estimate theoretical uncertainty related to VMD model, one refer to estimation of cross section of J/ψ photoproduction in *peripheral model* & found strong energy dependence close to threshold because non-diagonal γp→Vp & elastic Vp→Vp must have larger transfer momenta vs elastic scattering.
 This result in violation of VMD by factor of 5.

K.G. Boreskov & B.L. Ioffe, Sov J Nucl Phys 25, 331 (1977)





• There is no alternative VMD to get $J/\psi p$ SL from meson photoproduction.

Courtesy of A. Vainshtein & M. Ryskin, July 2020

To estimate theoretical uncertainty related to VMD model, one refer to estimation of cross section of J/ψ photoproduction in *peripheral model* & found strong energy dependence close to threshold because non-diagonal γp→Vp & elastic Vp→Vp must have larger transfer momenta vs elastic scattering.
 This result in violation of VMD by factor of 5.

K.G. Boreskov & B.L. Ioffe, Sov J Nucl Phys 25, 331 (1977)

• Color factor for *charmonium* is 1/9 while for *open charm* is 8/9.

B.Z. Kopeliovich, I. Schmidt, & M. Siddikov, Phys Rev C 95, 065203 (2017)



7/29/2020

• There is no alternative VMD to get $J/\psi p$ SL from meson photoproduction.

Courtesy of A. Vainshtein & M. Ryskin, July 2020

To estimate theoretical uncertainty related to VMD model, one refer to estimation of cross section of J/ψ photoproduction in *peripheral model* & found strong energy dependence close to threshold because non-diagonal γp→Vp & elastic Vp→Vp must have larger transfer momenta vs elastic scattering.
 This result in violation of VMD by factor of 5.

K.G. Boreskov & B.L. Ioffe, Sov J Nucl Phys 25, 331 (1977)

• Color factor for *charmonium* is 1/9 while for *open charm* is 8/9.

B.Z. Kopeliovich, I. Schmidt, & M. Siddikov, Phys Rev C 95, 065203 (2017)



• Additional suppression factor for $J/\psi N$ interaction @ threshold is OZI rule. OZI suppressed processes have larger number of independent fermion loops compared with non-suppressed processes.

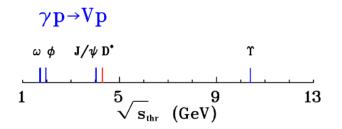


Courtesy of B. Kopeliovich & R. Jaffe, June 2020



Open Charm

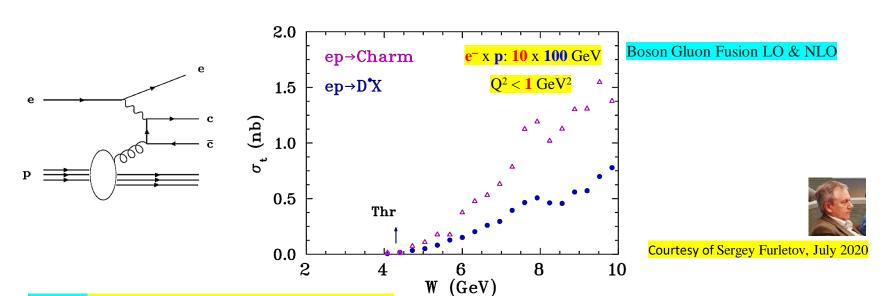
• It was shown that fluctuation of **photon** into open charm $\gamma p \rightarrow \overline{D}{}^0 \Lambda_c^+$ is preferable than into charmonium J/ψ . Suppression is just available phase space: $W_{thr}(\text{open charm}) = 4.30 \text{ GeV}$ while $W_{thr}(\text{charmonium}) = 4.03 \text{ GeV}$.





K. Boreskov *et al*, Phys Rev D **47**, 919 (1993)

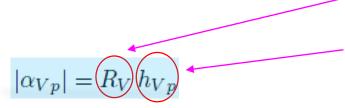
• There are no σ_t for open charm @ threshold.



HVQDIS: B.W. Harris & J. Smith, Phys Rev D 57, 2806 (1998)

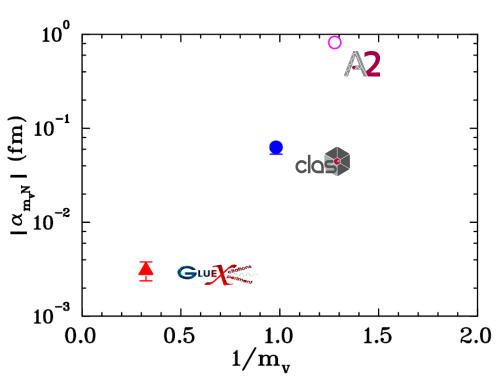


Vector-Meson - Nucleon SL

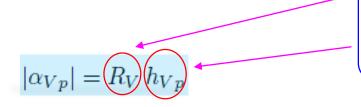


BNL EIC Group Meeting, Upton on Long Island in New York, July 2020

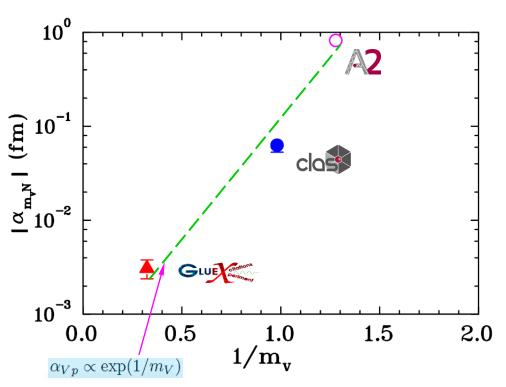
- EM factor R_v for each vector-meson are close to each other.
- Therefore, such big difference in SL is determined mainly by hadronic factor h_{Vp}.



Vector-Meson - Nucleon SL



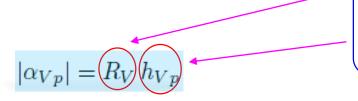
- EM factor R_V for each vector-meson are close to each other.
- Therefore, such big difference in SL is determined mainly by hadronic factor h_{Vp}



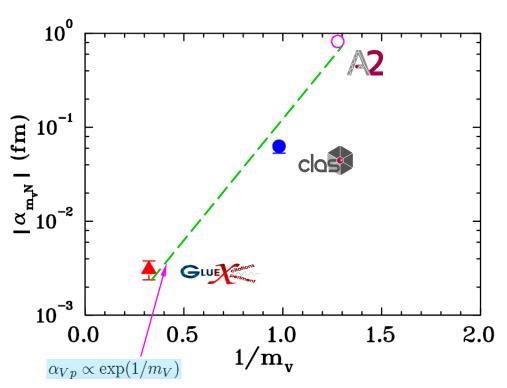
Such small value of φp SL compared to typical hadron size of 1 fm, indicates that proton is more transparent for φ-meson compared to ω-meson, & is much less transparent than for J/ψ-meson.

$$|\alpha_{J/\psi p}| \ll |\alpha_{\phi p}| \ll |\alpha_{\omega p}|$$

Vector-Meson - Nucleon SL

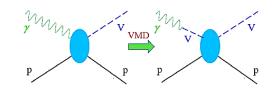


- EM factor R_V for each vector-meson are close to each other.
- Therefore, such big difference in SL is determined mainly by hadronic factor h_{Vp}



Such small value of φp SL compared to typical hadron size of 1 fm, indicates that proton is more transparent for φ-meson compared to ω-meson, & is much less transparent than for J/ψ-meson.

$$|\alpha_{J/\psi p}| \ll |\alpha_{\phi p}| \ll |\alpha_{\omega p}|$$



p→V coupling is proportional to α_s & separation of corresponding quarks.
 This separation (in zero approximation) is proportional to 1/m_V.





Ongoing Activities





• Present & future experiments @ AS12& 007^{IVIII} that are aimed to measure charmonium production on proton & nuclei will allow further studies of J/\(\psi N\) interaction & will give also access to variety of other interesting physics aspects that are present in near-threshold region.

• Present & future experiments @ AS12& 007^{J/*} that are aimed to measure charmonium production on proton & nuclei will allow further studies of J/\(\psi N\) interaction & will give also access to variety of other interesting physics aspects that are present in near-threshold region.

In particular, low-energy J/ψ photoproduction do/dt data can be used to extract fraction of nucleon's mass arising from gluons, & corresponding spatial distribution.

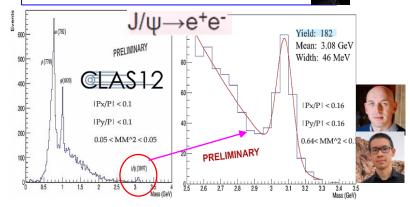
D. Kharzeev et al, Eur Phys J C 9, 459 (1999)

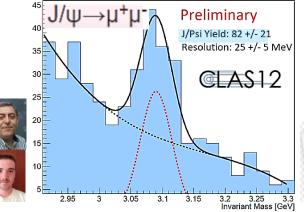


Fall 2018 dataset

Time-like Compton Scattering & J/ψ Photoproduction (E12-12-001A)

• Present & future experiments @ AS12& 007^{J/w} that are aimed to measure charmonium production on proton & nuclei will allow further studies of J/wN interaction & will give also access to variety of other *interesting* physics aspects that are present in near-threshold region.







Courtesy of Joseph Newton & Richard Tyson, CLAS Collab meeting, July 2020

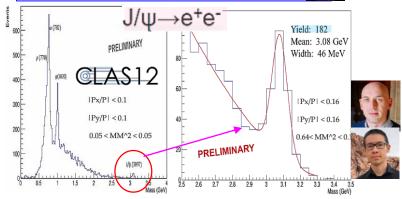


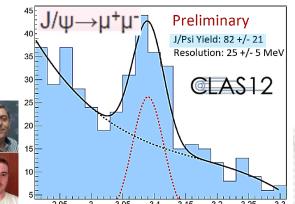


Fall 2018 dataset

*Time-like Compton Scattering &*J/\psi Photoproduction (E12-12-001A)

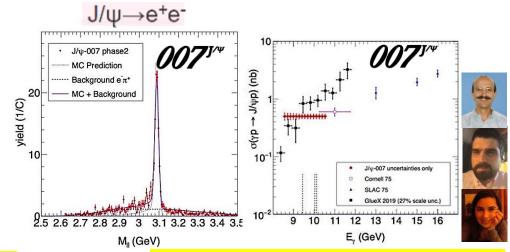
• Present & future experiments @ \bigcirc AS12& $007^{J/\psi}$ that are aimed to measure charmonium production on proton & nuclei will allow further studies of $J/\psi N$ interaction & will give also access to variety of other *interesting* physics aspects that are present in near-threshold region.







Search for Charmed "Pentaquark" using Photoproduction of J/ψ @ Threshold in Hall C @ Jefferson Lab (E12-16-007)



Courtesy of Burcu Duran & Sylvester Joosten, APS April 2020

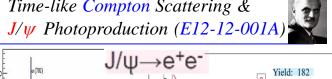
Courtesy of Joseph Newton & Richard Tyson, CLAS Collab meeting, July 2020

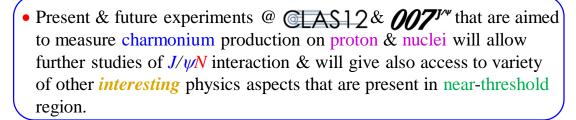




Fall 2018 dataset

Time-like Compton Scattering &

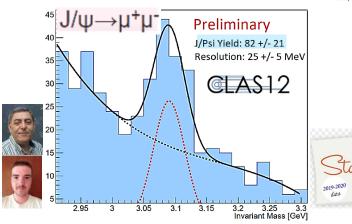


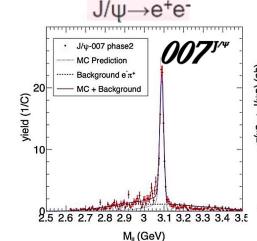


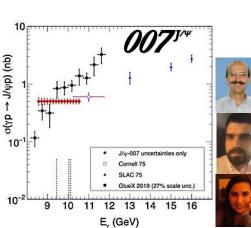
- σ_i is unmeasurable in both cases & requires extraction from $d\sigma/dt$.
- Critical factor is how close to threshold both experiments will go.

Mean: 3.08 GeV Width: 46 MeV ©I AS12 |Px/P| < 0.1 |Px/P| < 0.16 |Pv/P| < 0.1 |Pv/P| < 0.16 0.05 < MM^2 < 0.05 **PRELIMINARY**

Search for Charmed "Pentaquark" using Photoproduction of J/ψ @ Threshold in Hall C @ Jefferson Lab (E12-16-007)



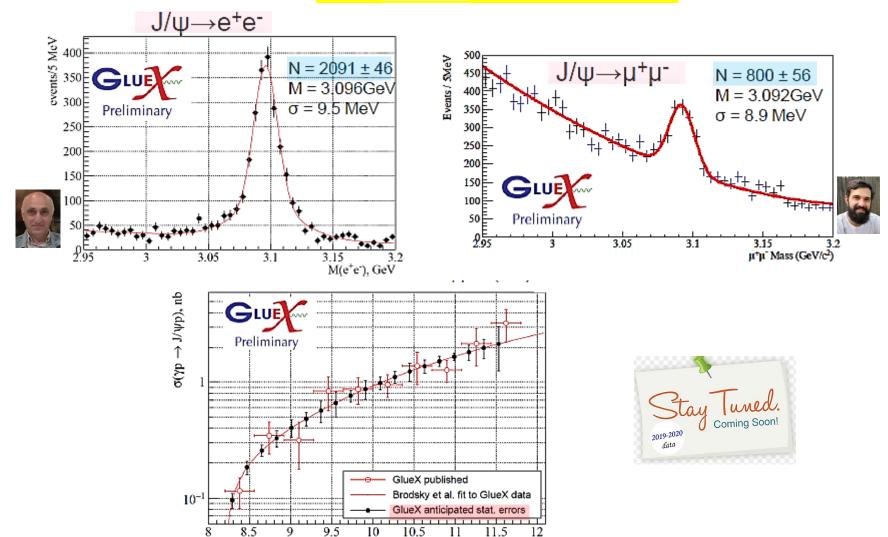




Courtesy of Joseph Newton & Richard Tyson, CLAS Collab meeting, July 2020

Courtesy of Burcu Duran & Sylvester Joosten, APS April 2020

• Total Statistics for 2016–2018



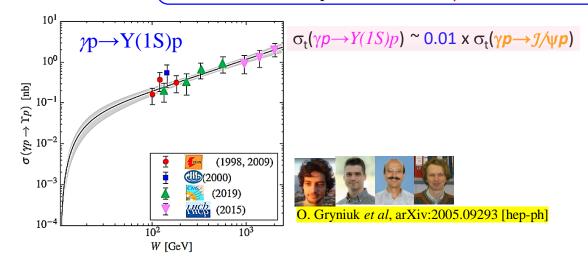




Expectation from



• Such small value of ϕp SL compared to typical hadron size of 1 fm, indicates that proton is more transparent for ϕ -meson compared to ω -meson, & is much less transparent than for J/ψ -meson.

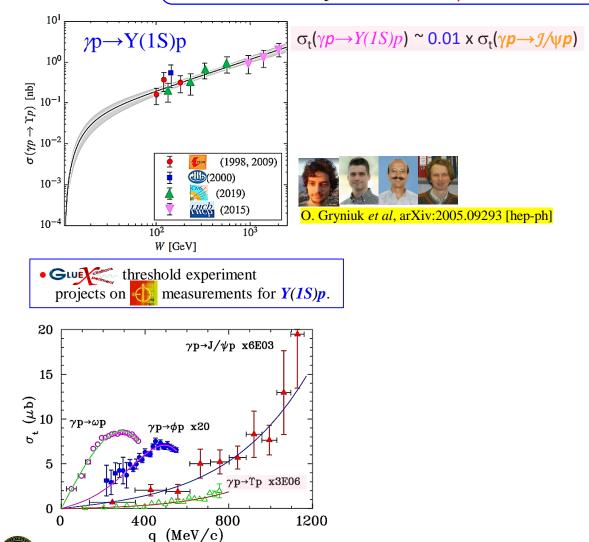




Expectation from



• Such small value of ϕp SL compared to typical hadron size of 1 fm, indicates that proton is more transparent for ϕ -meson compared to ω -meson, & is much less transparent than for J/ψ -meson.

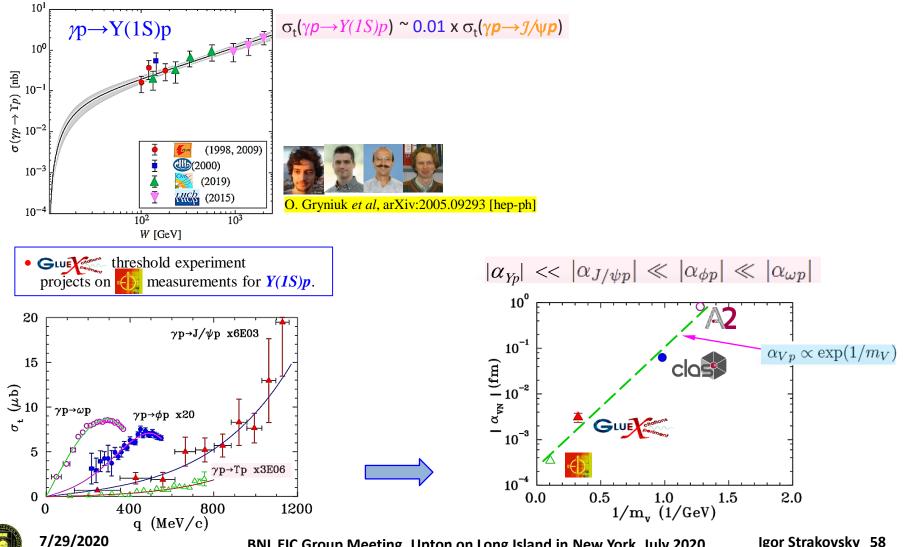


7/29/2020

Expectation from



• Such small value of ϕp SL compared to typical hadron size of 1 fm, indicates that proton is more transparent for ϕ -meson compared to ω -meson, & is much less transparent than for J/ψ -meson.





Summary







• High accurate measurements near-threshold by 2, close, & Guil allow to determine σ_t of reactions $\gamma p \rightarrow Vp \&$ to estimate absolute value of *Vp* SLs within VMD model.

$$|\alpha_{\rm op}|$$
 = (0.82 ± 0.03) fm
$$|\alpha_{\rm op}|$$
 = (0.063 ± 0.010) fm
$$|\alpha_{J/\psi p}|$$
 = (3.08 ± 0.55(stat.) ± 0.42(syst.)) mfm.



$$|\alpha_{\omega p}| = (0.82 \pm 0.03) \, \mathrm{fm}$$

$$|\alpha_{\phi p}| = (0.063 \pm 0.010) \, \mathrm{fm}$$

$$|\alpha_{J/\psi p}| = (3.08 \pm 0.55 (\mathrm{stat.}) \pm 0.42 (\mathrm{syst.})) \, \mathrm{mfm.}$$

• We found strong exponential increase of Vp SL with inverse mass of vector-meson $\alpha_{Vp} \propto \exp(1/m_V)$



• High accurate measurements near-threshold by 2, 6, 4 Gue, 4 allow to determine 5 of reactions 4 by 4 to estimate absolute value of 4 by 4 by 4 within 4 model.

$$|\alpha_{\omega p}| = (0.82 \pm 0.03) \, \mathrm{fm}$$

$$|\alpha_{\phi p}| = (0.063 \pm 0.010) \, \mathrm{fm}$$

$$|\alpha_{J/\psi p}| = (3.08 \pm 0.55 (\mathrm{stat.}) \pm 0.42 (\mathrm{syst.})) \, \mathrm{mfm.}$$

- We found strong exponential increase of Vp SL with inverse mass of vector-meson $\alpha_{Vp} \propto \exp(1/m_V)$
- It is remarkable that proton is quite so *transparent* to J/ψ , though general progression from ω to ϕ to J/ψ $|\alpha_{J/\psi p}| \ll |\alpha_{\phi p}| \ll |\alpha_{\omega p}|$ is perhaps qualitatively reasonable on account of OZI rule.



• High accurate measurements near-threshold by 2, 6, 4 Gue, 4 allow to determine 5 of reactions 4 by 4 to estimate absolute value of 4 by 4 by 4 within 4 model.

$$|\alpha_{\rm op}| = (0.82 \pm 0.03) \ {\rm fm}$$

$$|\alpha_{\rm op}| = (0.063 \pm 0.010) \ {\rm fm}$$

$$|\alpha_{J/\psi p}| = (3.08 \pm 0.55 ({\rm stat.}) \pm 0.42 ({\rm syst.})) \ {\rm mfm}.$$

- We found strong exponential increase of V_p SL with inverse mass of vector-meson $\alpha_{Vp} \propto \exp(1/m_V)$
- It is remarkable that proton is quite so *transparent* to J/ψ , though general progression from ω to ϕ to J/ψ $|\alpha_{J/\psi p}| \ll |\alpha_{\phi p}| \ll |\alpha_{\omega p}|$ is perhaps qualitatively reasonable on account of OZI rule.
- Due to small size of 'young' J/ψ vs 'old' J/ψ , measured SL is very small. J/ψ crated by photon @ threshold then most probably J/ψ is not formed completely & its radius is smaller than that for normal J/ψ . Therefore, one observe stronger $(r_{J/\psi}^2/R_{J/\psi}^2 < 100)$ suppression for J/ψ -p interaction.

E.L. Feinberg, Sov Phys Usp, **23**, 629 (1980) Courtesy of Michael Ryskin, July 2020



SUMMARY

High accurate measu

• High accurate measurements near-threshold by 2, 6, 4 Gue, 4 allow to determine 5 of reactions 4 by 4 to estimate absolute value of 4 SLs within 4 model.

$$|\alpha_{\omega p}| = (0.82 \pm 0.03) \, \mathrm{fm}$$

$$|\alpha_{\phi p}| = (0.063 \pm 0.010) \, \mathrm{fm}$$

$$|\alpha_{J/\psi p}| = (3.08 \pm 0.55 (\mathrm{stat.}) \pm 0.42 (\mathrm{syst.})) \, \mathrm{mfm.}$$

- We found strong exponential increase of Vp SL with inverse mass of vector-meson $\alpha_{Vp} \propto \exp(1/m_V)$
- It is remarkable that proton is quite so *transparent* to J/ψ , though general progression from ω to ϕ to J/ψ $|\alpha_{J/\psi p}| \ll |\alpha_{\phi p}| \ll |\alpha_{\omega p}|$ is perhaps qualitatively reasonable on account of OZI rule.
- Due to small size of 'young' J/ψ vs 'old' J/ψ , measured SL is very small. J/ψ crated by photon @ threshold then most probably J/ψ is not formed completely & its radius is smaller than that for normal J/ψ . Therefore, one observe stronger $(r_{J/\psi}^2/R_{J/\psi}^2 < 100)$ suppression for J/ψ -p interaction.

E.L. Feinberg, Sov Phys Usp, **23**, 629 (1980)

Courtesy of Michael Ryskin, July 2020

• Light vector-meson can be 'young' as well. This depends on particular kinematics. Another point is that for slow heavy quark, one need more time to reach equilibrium, i.e., to form final (long-living/static) vector-meson.





• High accurate measurements near-threshold by 2, 6, 4 Gue, 4 allow to determine 5 of reactions 4 by 4 to estimate absolute value of 4 by 4 by 4 within 4 model.

$$|\alpha_{\rm op}| = (0.82 \pm 0.03) \, {\rm fm}$$

$$|\alpha_{\rm op}| = (0.063 \pm 0.010) \, {\rm fm}$$

$$|\alpha_{J/\psi p}| = (3.08 \pm 0.55 ({\rm stat.}) \pm 0.42 ({\rm syst.})) \, {\rm mfm.}$$

- We found strong exponential increase of Vp SL with inverse mass of vector-meson $\alpha_{Vp} \propto \exp(1/m_V)$
- It is remarkable that proton is quite so *transparent* to J/ψ , though general progression from ω to ϕ to J/ψ $|\alpha_{J/\psi p}| \ll |\alpha_{\phi p}| \ll |\alpha_{\omega p}|$ is perhaps qualitatively reasonable on account of OZI rule.
- Due to small size of 'young' J/ψ vs 'old' J/ψ , measured SL is very small. J/ψ crated by photon @ threshold then most probably J/ψ is not formed completely & its radius is smaller than that for normal J/ψ . Therefore, one observe stronger $(r_{J/\psi}^2/R_{J/\psi}^2 < 100)$ suppression for J/ψ -p interaction.

E.L. Feinberg, Sov Phys Usp, **23**, 629 (1980) Courtesy of Michael Ryskin, July 2020

• Light vector-meson can be 'young' as well. This depends on particular kinematics. Another point is that for slow heavy quark, one need more time to reach equilibrium, i.e., to form final (long-living/static) vector-meson.



will open new window in solving **VN SL** puzzle. It will allow to make deal with 'young' **Y** as well.

Ocknowledgements

I thank

Yakov Azimov Konstantin Boreskov

Stanley Brodsky Daniel Carman

Eugene Chudakov Michael Eides

Denis Epifanov Sergey Furletov

Takatsugu Ishikawa Robert Jaffe

Dmitry Kharzeev Boris Kopeliovich

Jerry Miller Lubomir Pentchev

Michael Ryskin Alexander Titov

Arkady Vainstein Michael Voloshin

for useful remarks & continuous interest in project.



