

Study of multiplicity and transverse momentum fluctuations in the Monte-Carlo model of interacting quark-gluon strings

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Saint-Petersburg State University

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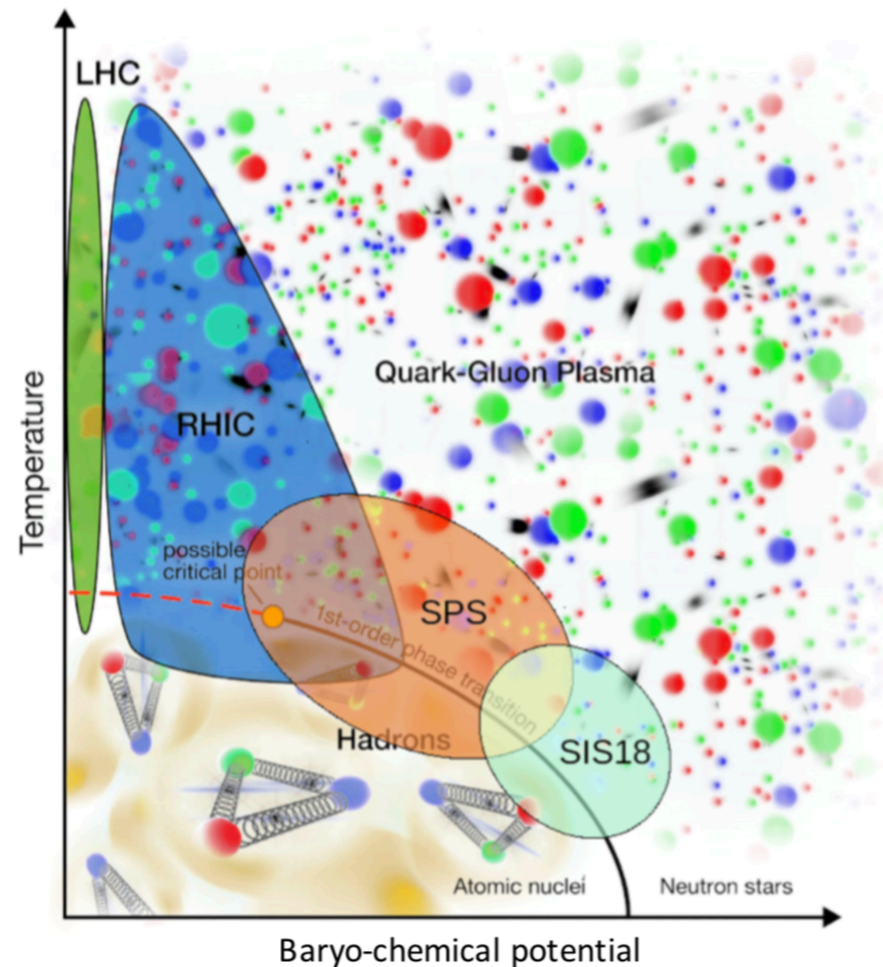
Outline

Study of multiplicity and transverse momentum fluctuations in the Monte-Carlo model of interacting quark-gluon strings

1. Study of fluctuations in the search for CP
2. Quark-gluon string models of particle production
3. String fusion in Monte-Carlo model
4. Results
5. Outlook

Study of fluctuations in the search for CP

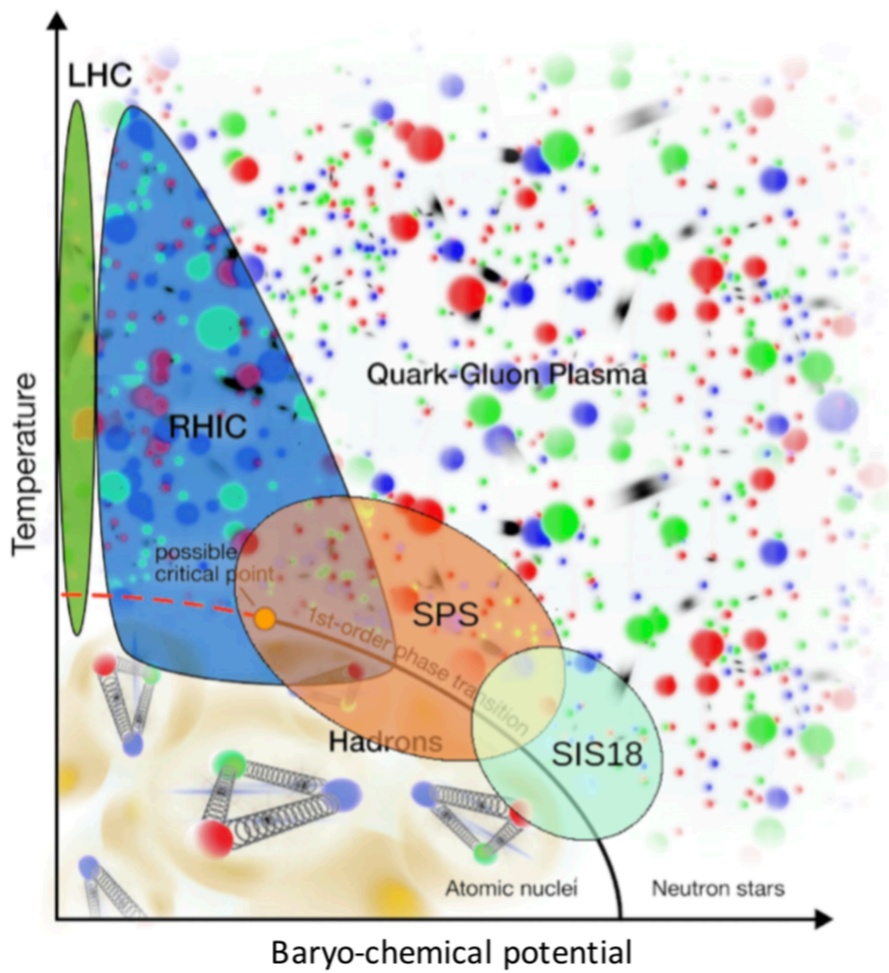
Expected phase diagram of strongly interacting matter



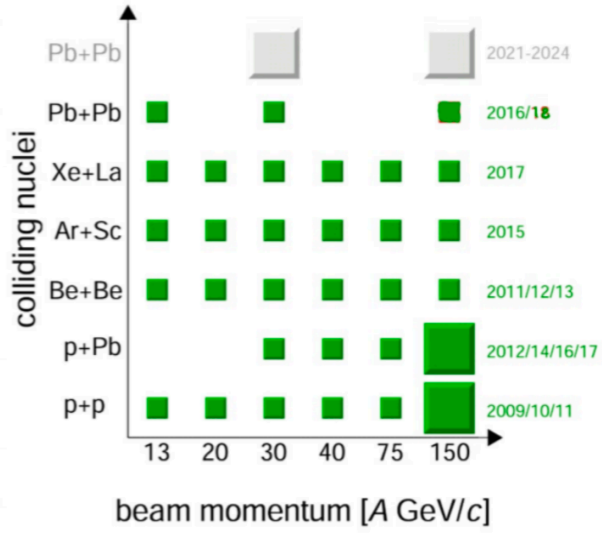
We know about hadron and QGP phases.
What is in between?

Study of fluctuations in the search for CP

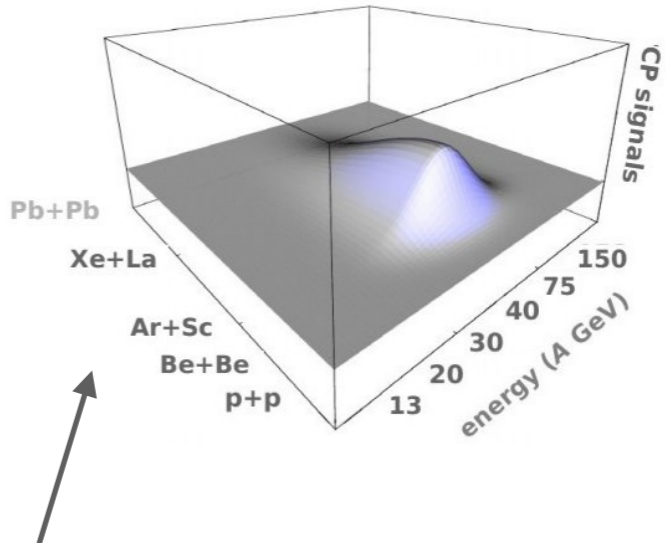
Expected phase diagram of strongly interacting matter



NA61/SHINE @ SPS CERN



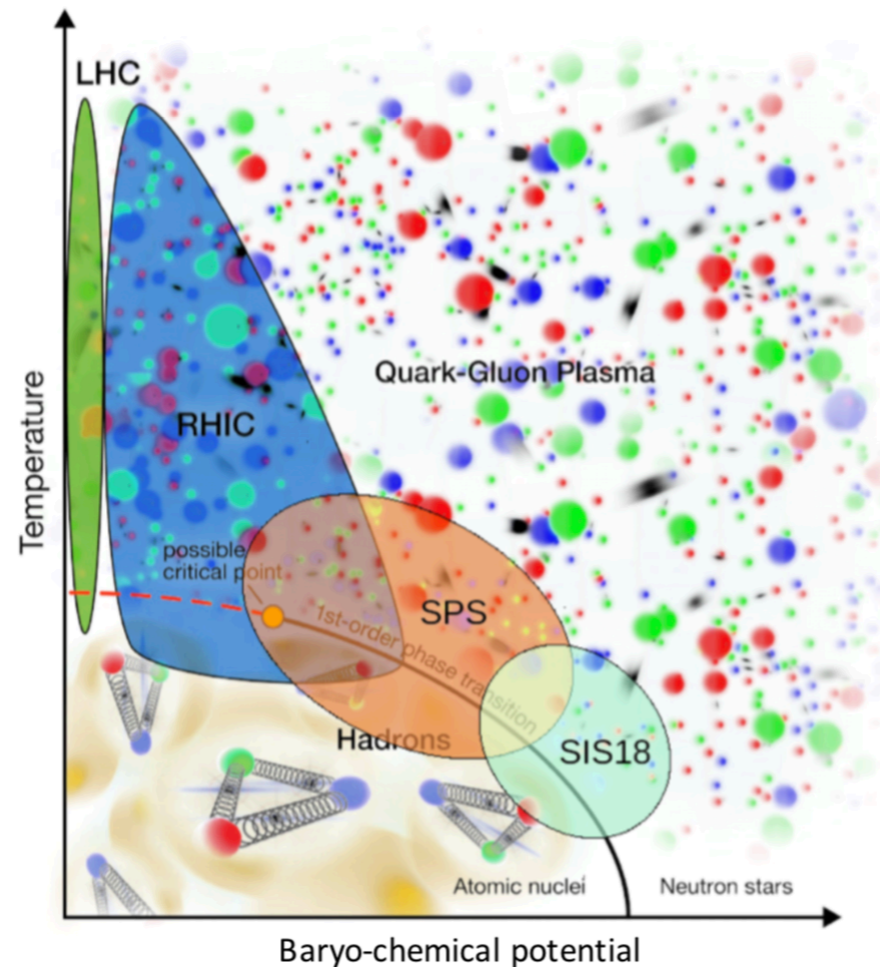
[Gazdzicki, M. and Seyboth, P. *Acta Phys. Pol. B* 47, 1201 (2016)]



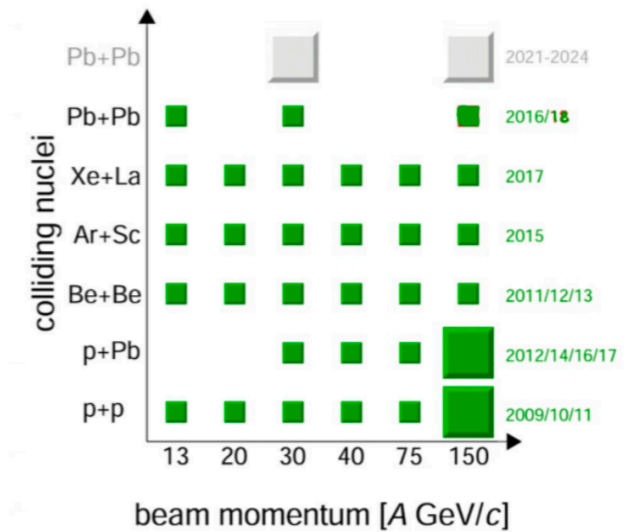
Experimental System size vs Energy scan to spot the critical point by the anticipated signal of enhanced fluctuations

Study of fluctuations in the search for CP

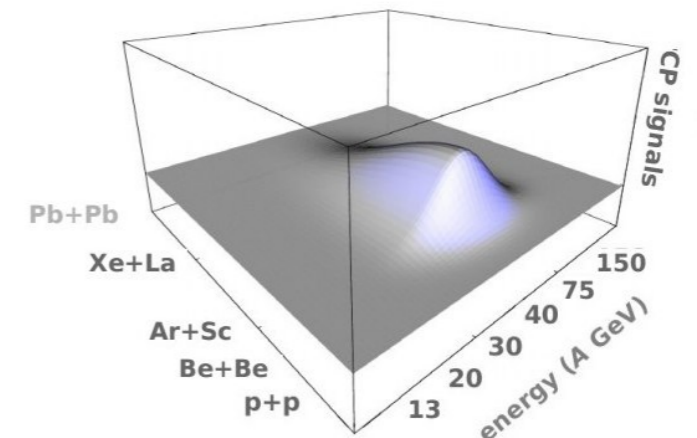
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NA61/SHINE @ SPS CERN



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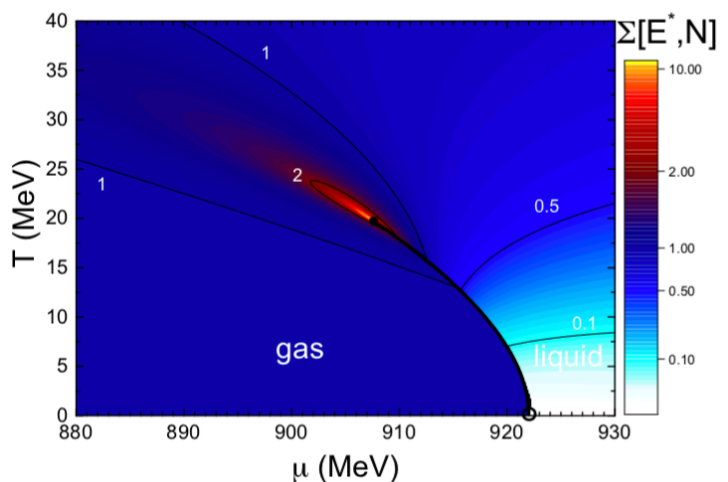
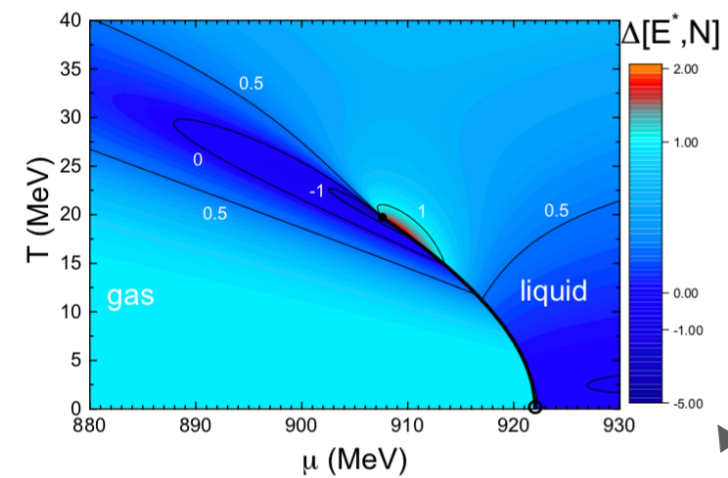
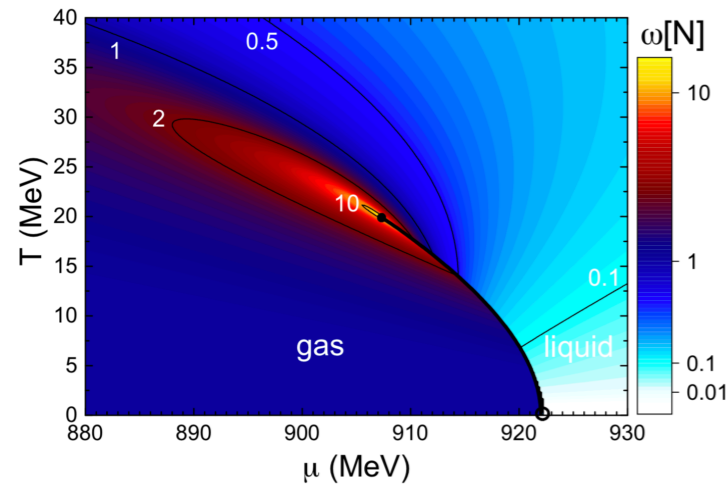
Study of the special strongly intensive quantities that are independent of the volume and fluctuations of the volume within GCA or independent sources model

$$\Delta[A, B] = \frac{1}{C_{\Delta}} [\langle B \rangle \omega[A] - \langle A \rangle \omega[B]]$$

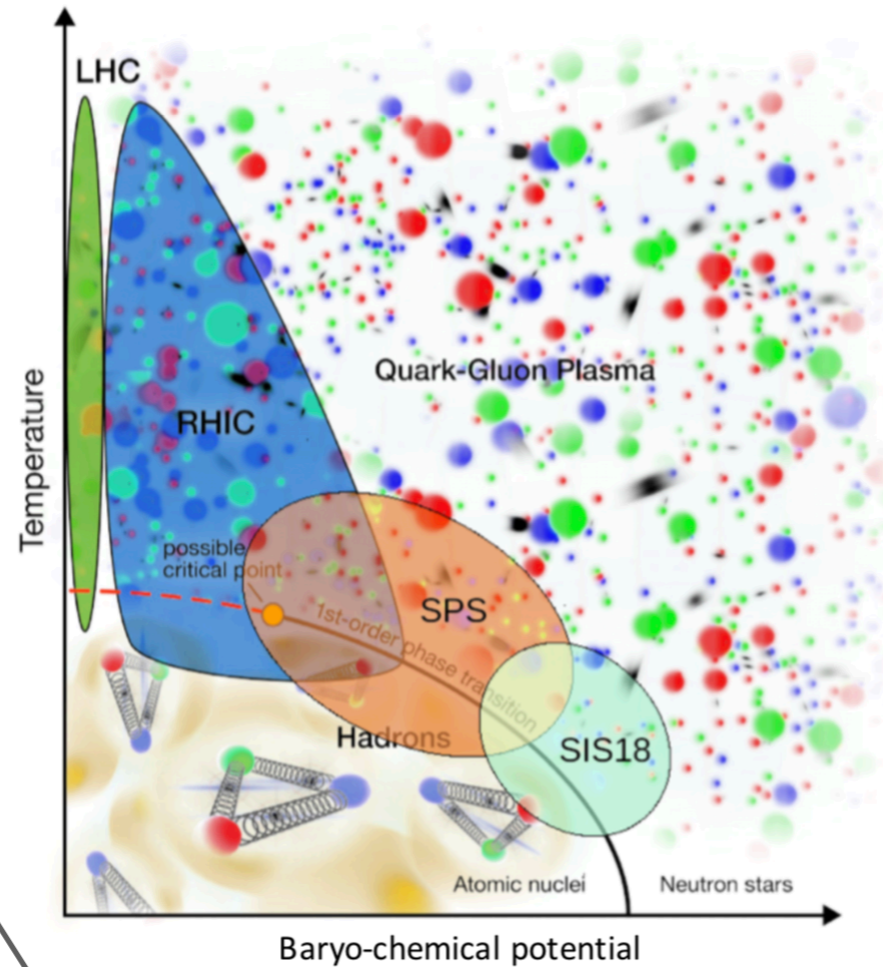
[M. I. Gorenstein and M. Gaździcki, *Physical Review C* 84, 014904 (2011)]

$$\Sigma[A, B] = \frac{1}{C_{\Sigma}} [\langle B \rangle \omega[A] + \langle A \rangle \omega[B] - 2(\langle AB \rangle - \langle A \rangle \langle B \rangle)]$$

Study of fluctuations in the search for CP



Expected phase diagram of strongly interacting matter

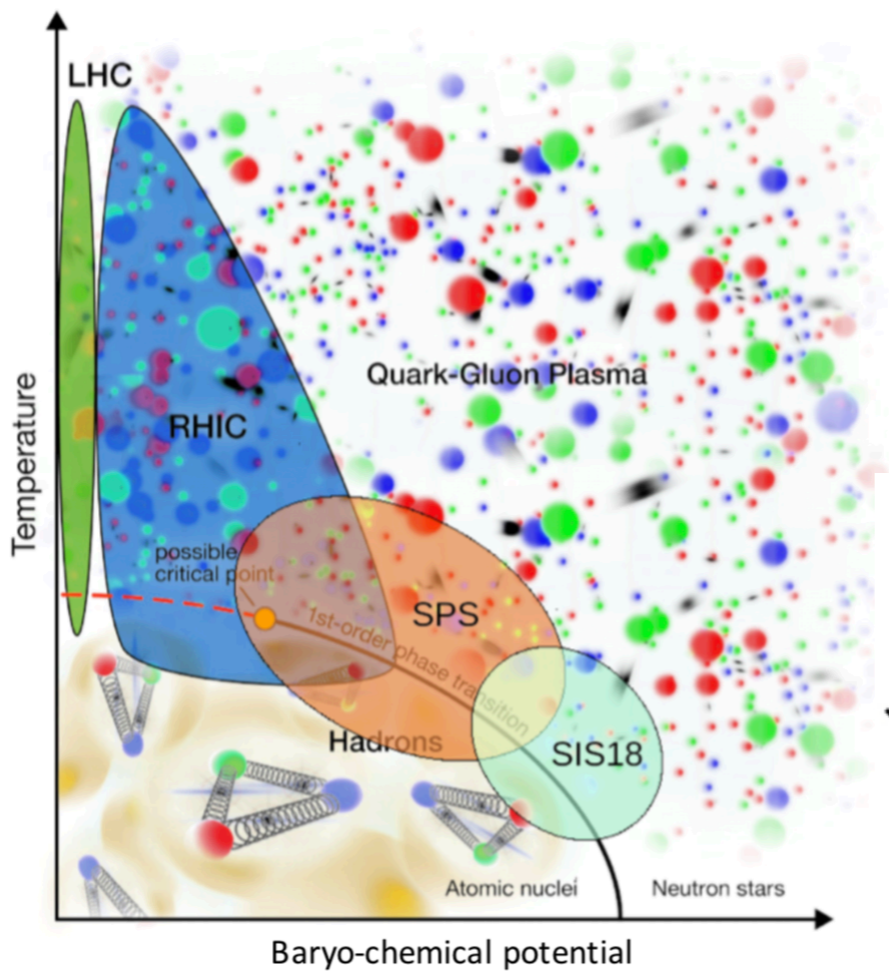


Possible sensitivity: nucleon system with van der Waals EOS in GCE formulation in the vicinity of the Critical Point

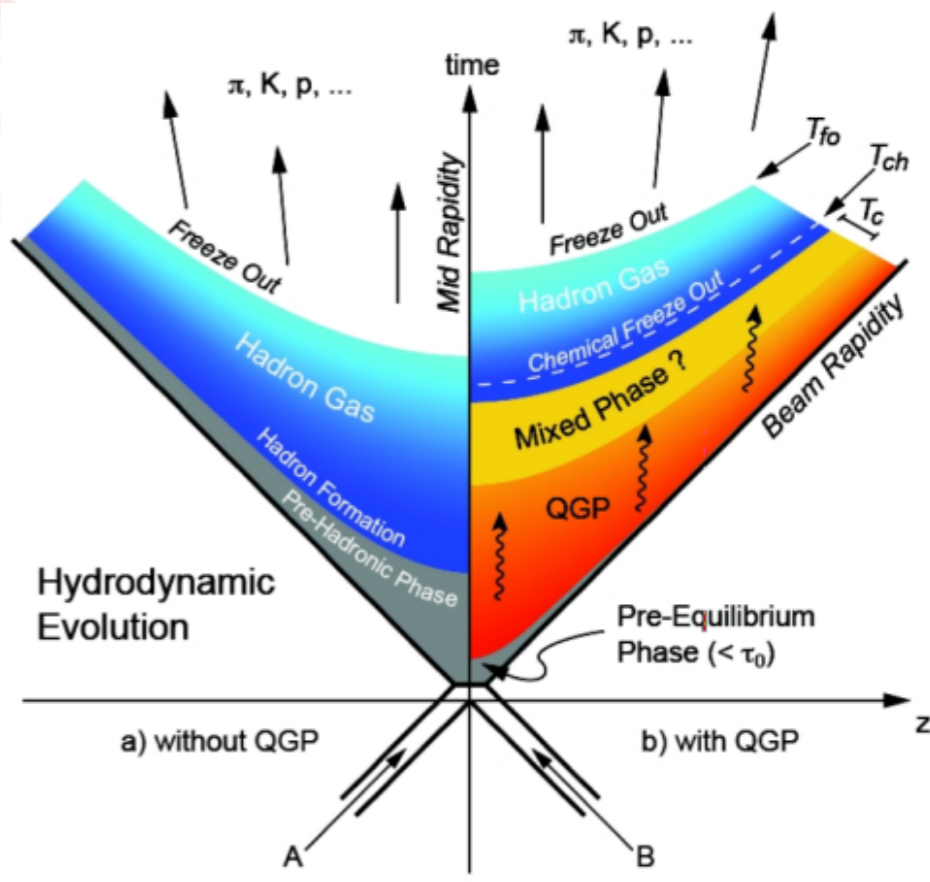
[Vovchenko, Gorenstein, Stoecker, PRL 118:182301]

Study of fluctuations in the search for CP

Expected phase diagram of strongly interacting matter



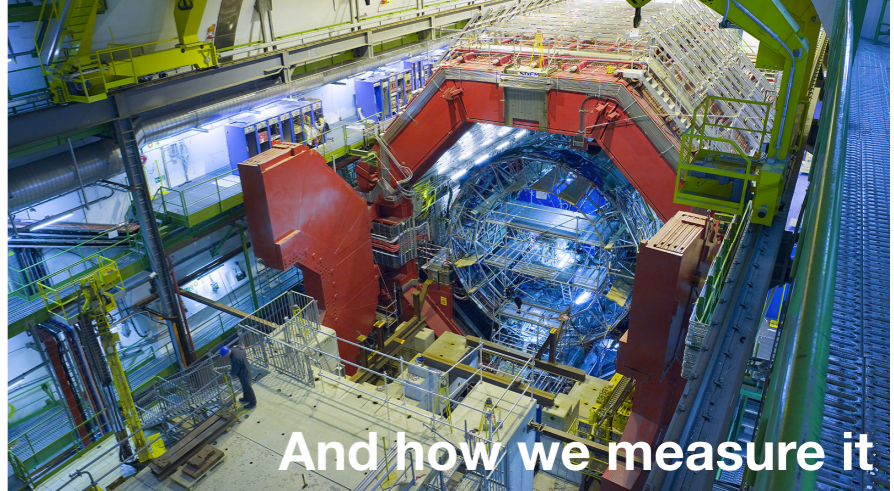
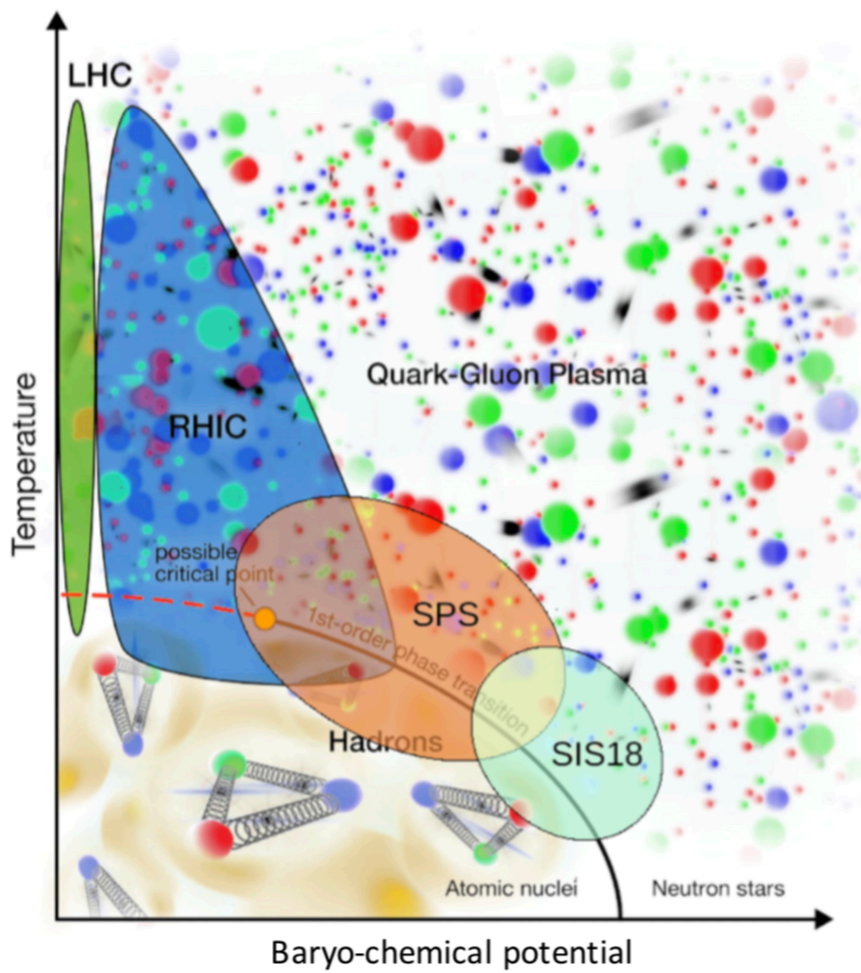
Picture of the collision as we think we understand it



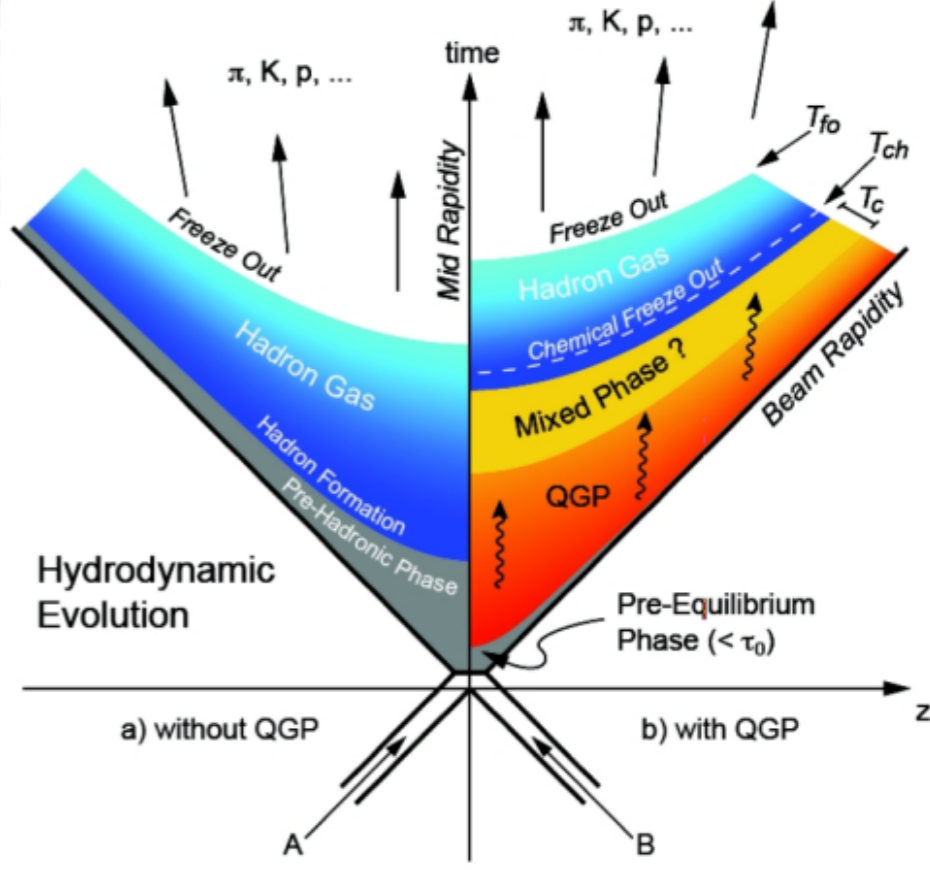
Study of fluctuations in the search for CP

Expected phase diagram of strongly interacting matter

Mona Schweizer, CERN



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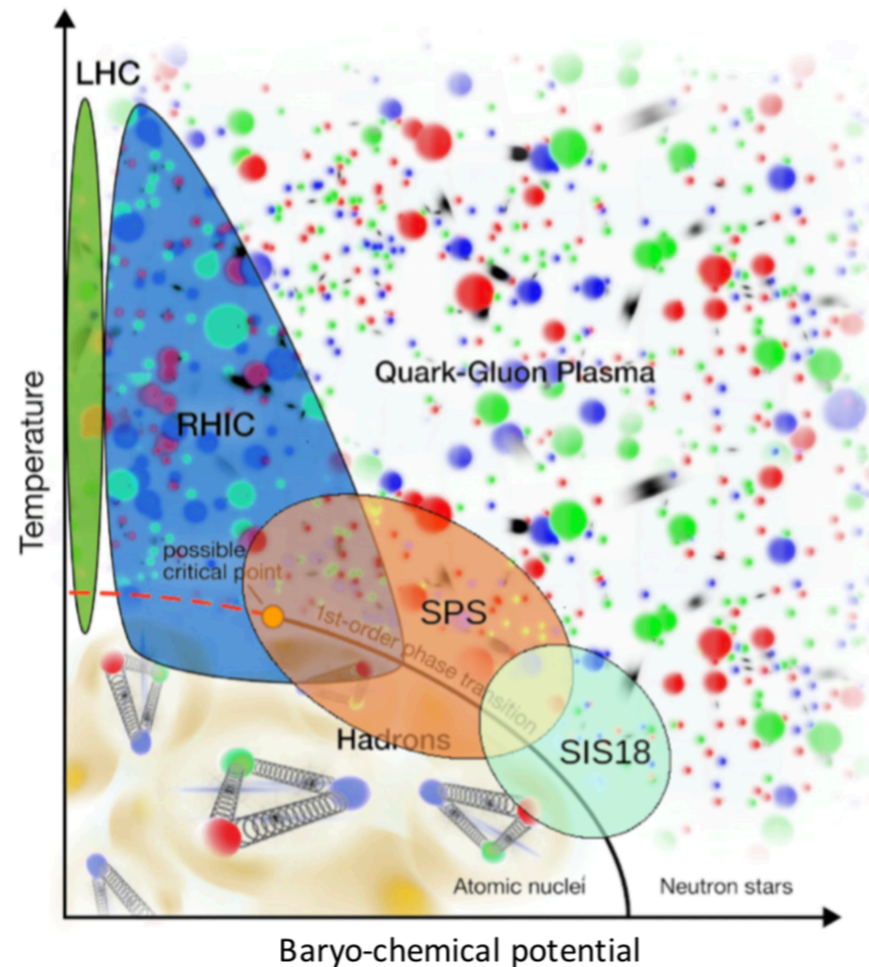
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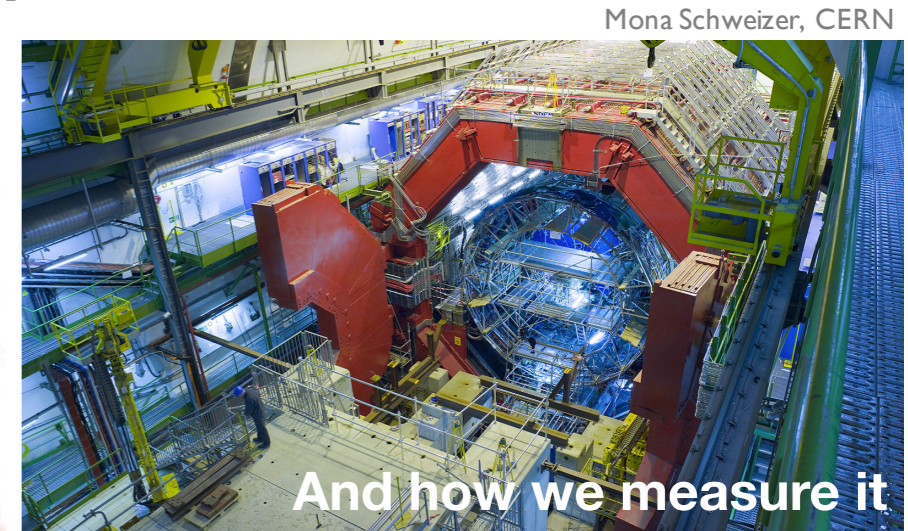
One can study:

- the fluctuations of multiplicity, transverse momentum
- moments of net electric charge or net baryon charge
- correlation coefficients to reveal the collective behavior

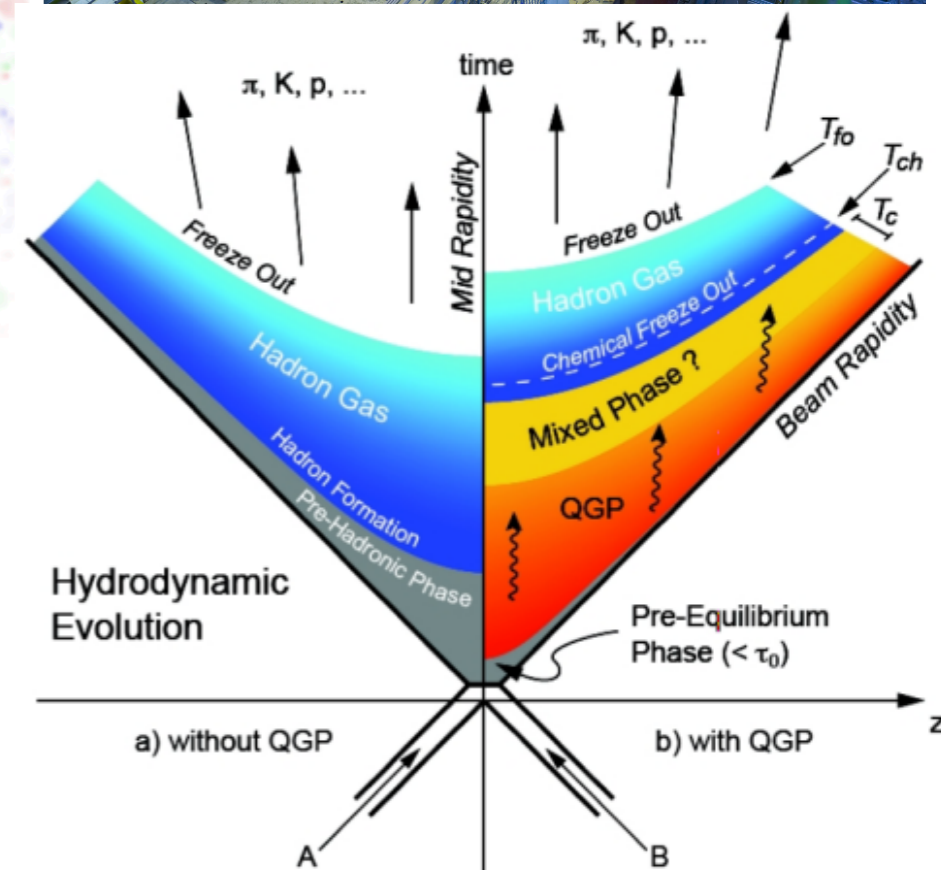
But the fluctuation background has to be subtracted to catch the CP!



Picture of the collision as we think we understand it



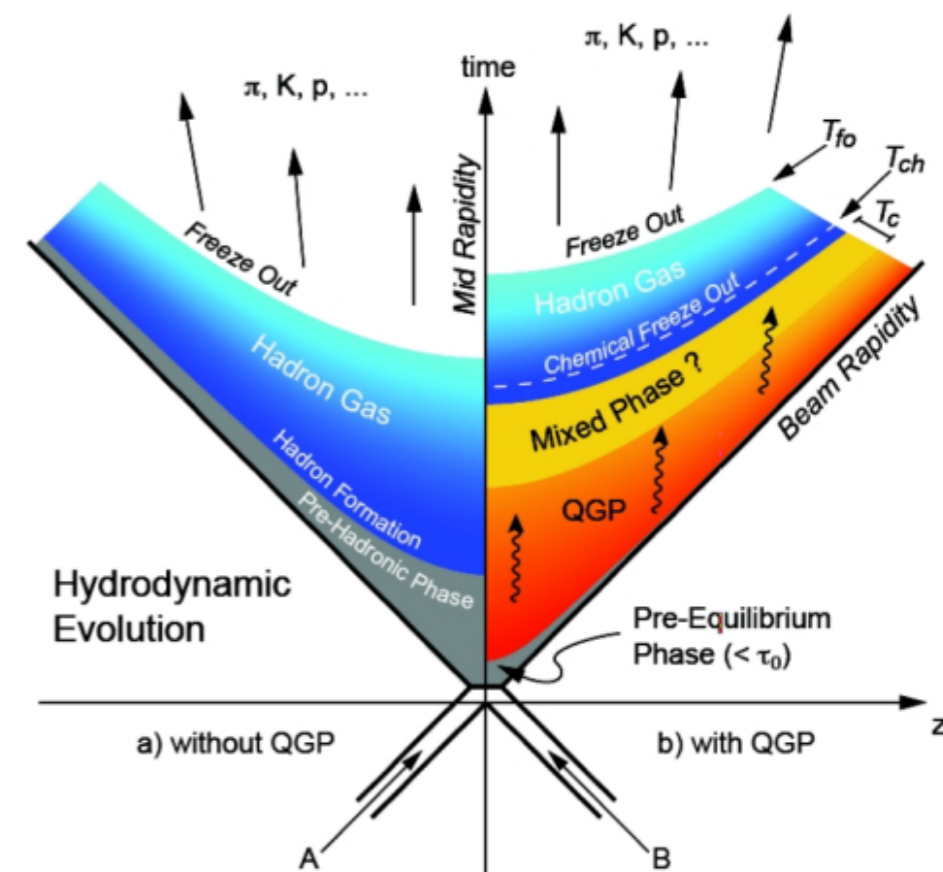
Mona Schweizer, CERN



Quark-gluon string models of particle production

What are they, string models:

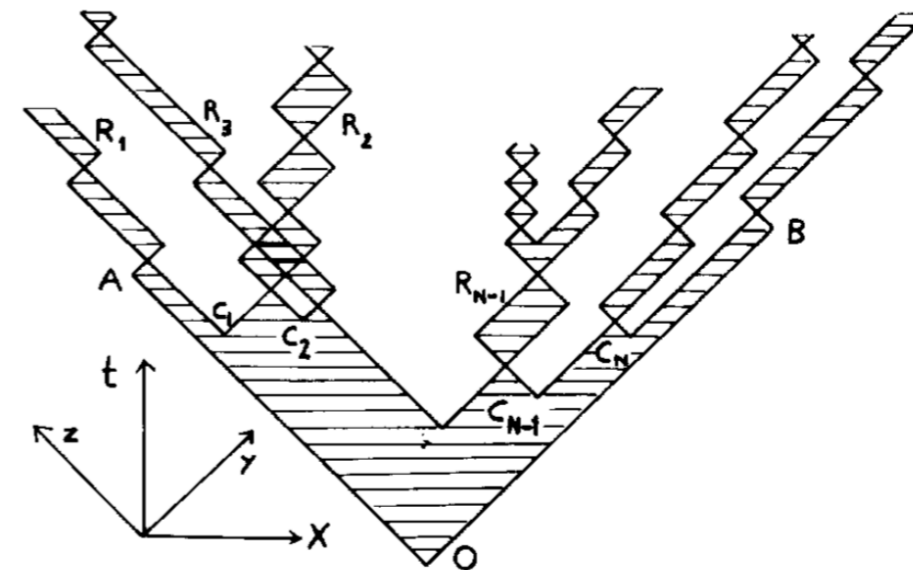
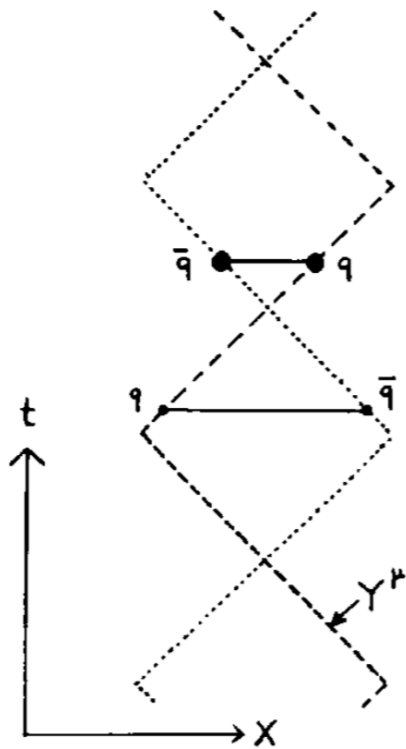
- Non-perturbative Regge approach to describe the soft particle spectra ($< 1 \text{ GeV}/c$)



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- Colorless hadron represented by the oscillating Jo-Jo solution

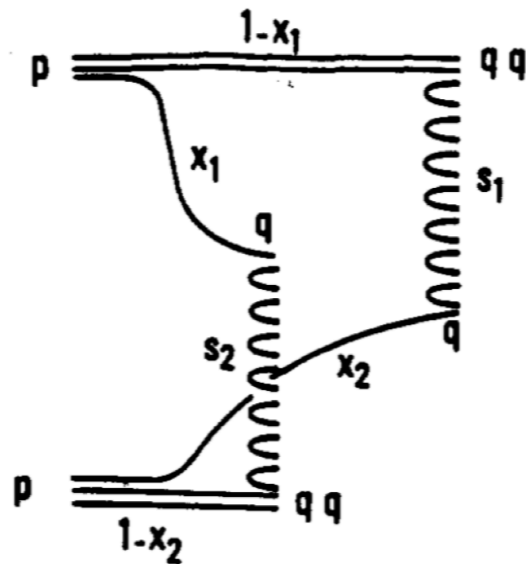


[X. Artru and G. Menessier Nuclear Physics B70 (1974) 93 - 115]

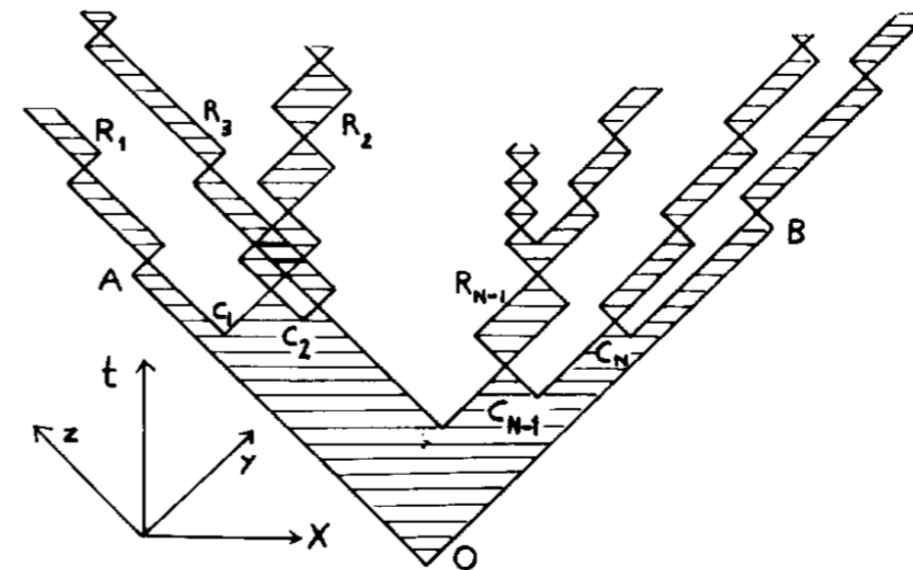
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- Unitarity cut of the cylindrical Pomeron results in two-chain diagram



[Capella A. et al Physics Reports 236, Nos. 4 & 5 (1994) 225 - 329]

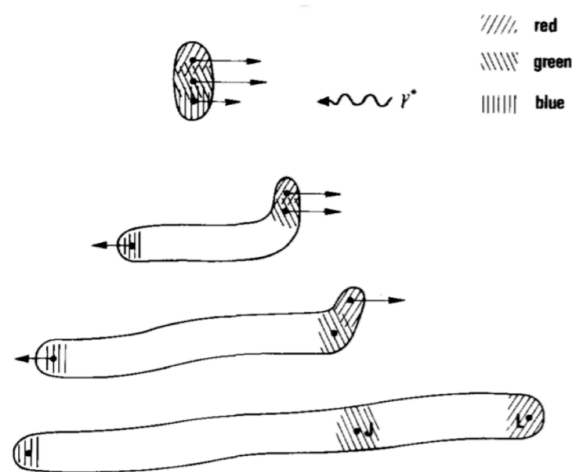


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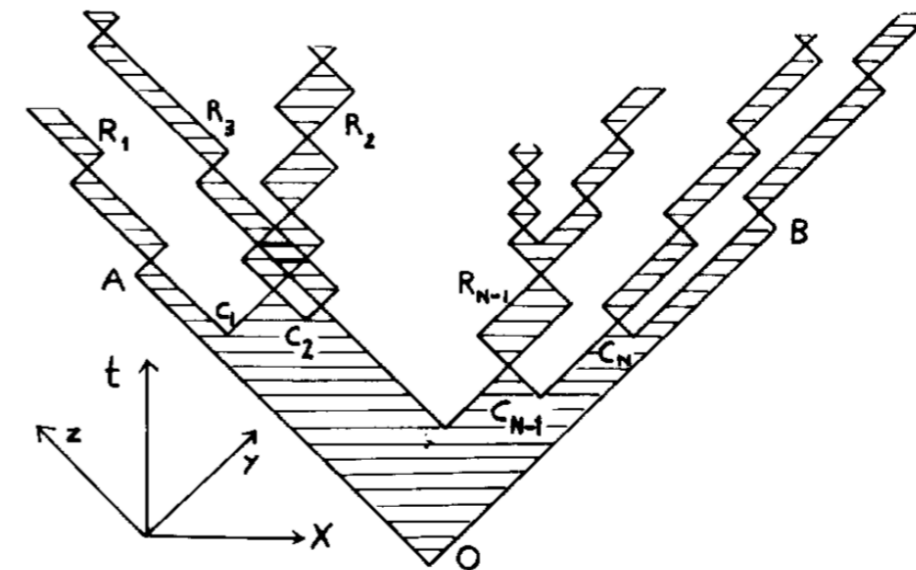
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- Unitarity cut of the cylindrical Pomeron results in two-chain diagram
- Lund model: longitudinally extended object stretched between the flying outwards wounded quarks and formed by the color field lines gathered together due to the gluon self-interaction



[Andersson B. et al *Physics Reports* 97, 31–145 (1983)]



[X. Artru and G. Menessier *Nuclear Physics B*70 (1974) 93 - 115]

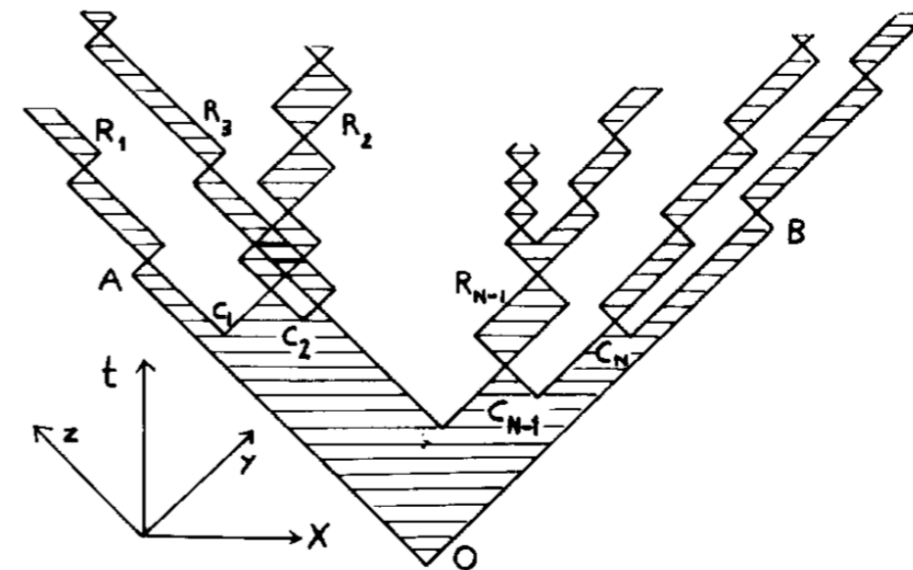
Quark-gluon string models of particle production

Motivation to study fluctuations in this framework:

[X. Artru, G. Mennessier, Nuclear Physics B **70**, 93 (1974), X. Artru, Physics Reports **97**, 147 (1983)]

- String as a particle emitting source impact to the **wide rapidity range** with the plateau at mid-rapidity → essential tool for studying long-range correlations and fluctuations

[N.S.Amelin, N.Armesto, M.A.Braun, E.G.Ferreiro, and C.Pajares, Phys Rev Let **73**, 2813 (1994)]



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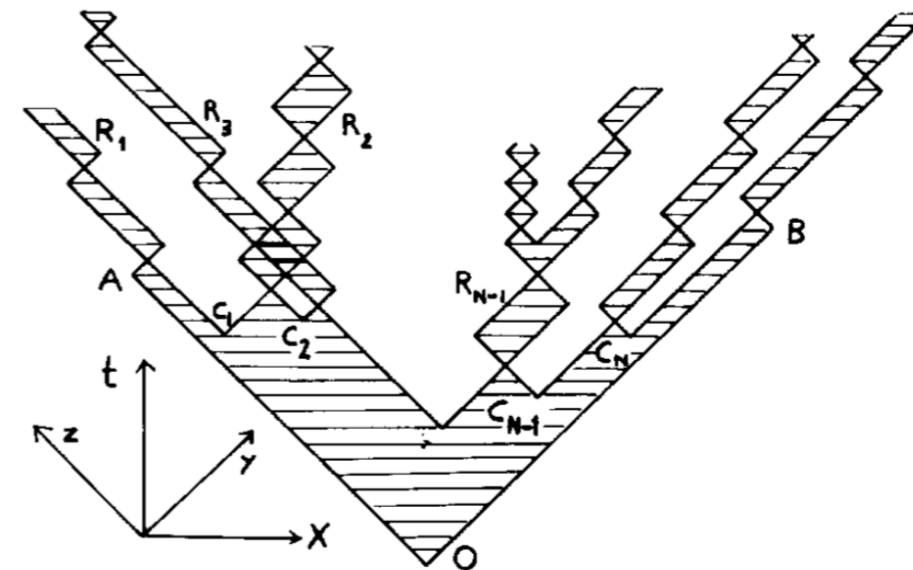
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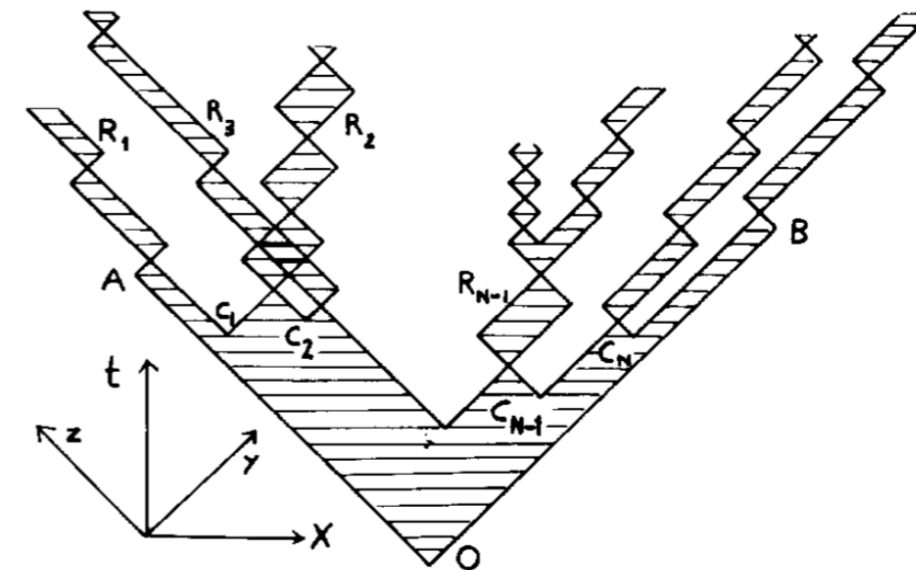
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- Variations of string ends positions in η \rightarrow estimate the influence of **initial conditions** and define the non-critical background of fluctuations [M. Rohrmoser, W. Broniowski, Phys Rev C **99**, 024904 (2019)]



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+ consider also string interactions

[V. Abramovskii and V. Kancheli, JETP letters **31**, 566 (1980)]

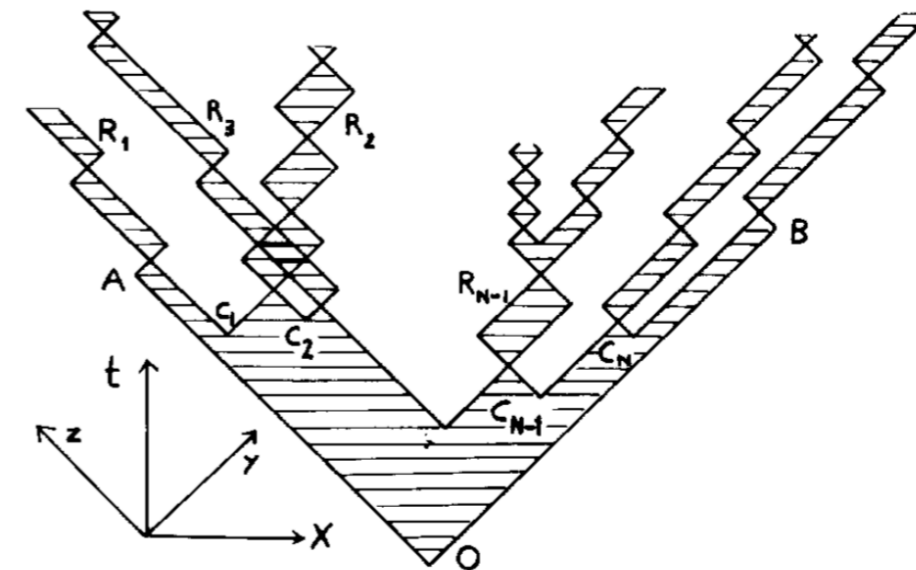
M. A. Braun and C. Pajares, Phys Let B **287**, 154 (1992)

M. Braun and C. Pajares, Nuclear Physics B **390**, 559 (1993)

N. S. Amelin, M. A. Braun, and C. Pajares, Phys Let B **306**, 312 (1993)

N.S.Amelin, M.A.Braun, and C.Pajares, Zeitschrift für Physik C Particles and Fields **63**, 507 (1994)

I. Altsybeev, AIP Conference Proceedings **1701**, 100002 (2016)]

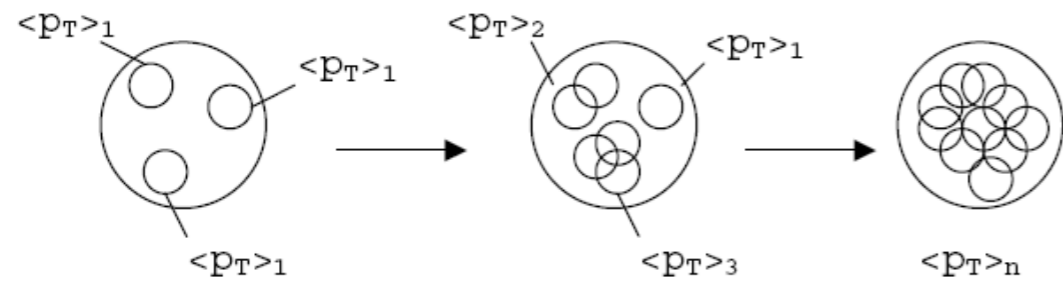


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String interactions

The string transverse position fluctuations changes the type of particle emitting sources

[Pajares, C. *Eur. Phys. J. C* 43, 9–14 (2005)]



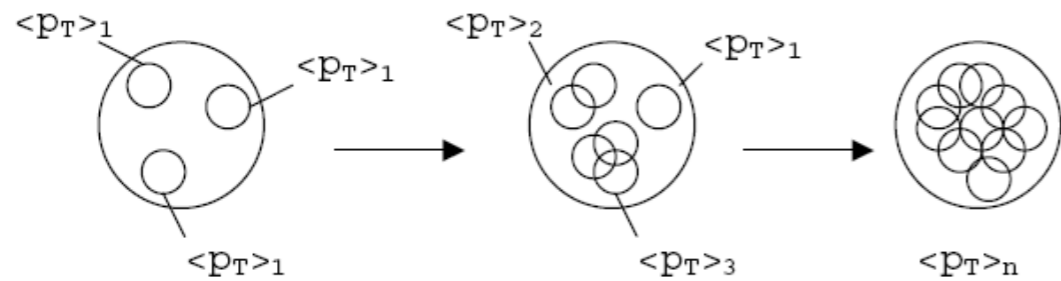
No fluctuations

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Simplification of the transverse picture

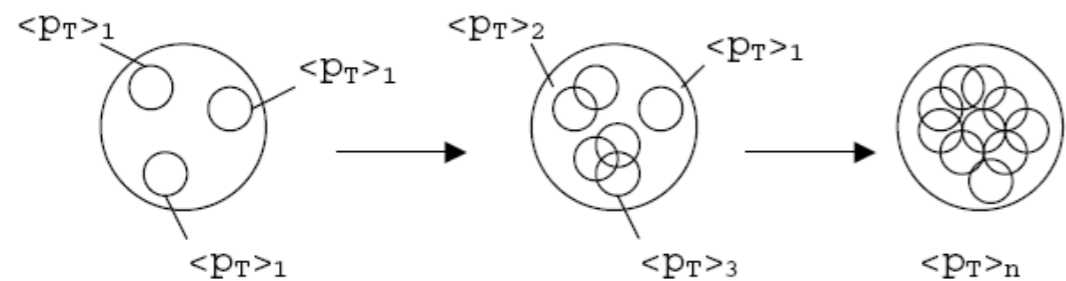
[V. Vechernin, I. Lakomov PoS(Baldin ISHEPP XXI)072 (2012)]

	"overlaps" (local fusion)	"clusters" (global fusion)
SFM	$C = \{S_1, S_2, \dots\}$ S_k – area covered k-times 	$C = \{S_1^{cl}, S_2^{cl}, \dots\}$ $k_i^{cl} = \frac{N_i^{str} \cdot \sigma_0}{S_i^{cl}}$
cellular analog of SFM	$C = \{N_{ij}^{str}\}$ $k_{ij} = N_{ij}^{str}$ – "occupation" numbers	$C = \{S_1^{cl}, S_2^{cl}, \dots\}$ $k_i^{cl} = \frac{N_i^{str} \cdot \sigma_0}{S_i^{cl}}$ $S_1^{cl} = 3\sigma_0; N_1^{str} = 5; k_1^{cl} = 5/3$ $S_2^{cl} = 2\sigma_0; N_2^{str} = 4; k_2^{cl} = 2$

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String fusion modifies the color field density [Braun, M. A., Kolevatov, R. S., Pajares, C. Vechernin, V. V. EPJ C, 32, 535–546 (2004)]

This affects the mean multiplicity by the string and the mean transverse momentum of produced particles

$$\langle \mu \rangle_k = \mu_0 \sqrt{k}$$

$$\langle p_T^2 \rangle_k = p_0^2 \sqrt{k}$$

Variations in the string length, location and interaction introduce additional fluctuations

Monte-Carlo model of interacting quark-gluon strings

Previous realization of the model:

[D. P., V.N. Kovalenko / (2020) / Study of Forward-Backward multiplicity fluctuations and correlations with pseudorapidity/ Journal Physics of Elementary Particles and Atomic Nuclei / Vol. 51 (3) / p. 323-326 / <https://doi.org/10.1134/S1063779620030247>]

Old one	New one (this report)
no quarks	quarks flavors and their directions of flight are sampled
string ends randomly distributed in rapidity space	string ends position in rapidity space calculated from quarks' PDFs
string impact calculated only in pseudorapidity intervals of interest	strings are discretized in rapidity with ϵ step
only multiplicities are calculated, «particles» do not exist	«particle» produced by a string has p_T and rapidity
mean number of strings in event is fixed, actual number is sampled from Poisson distribution for each event	
strings' transverse position is sampled, if strings are in the same cell - string fusion occurs	

Quantities of interest

Strongly intensive quantities $\Delta[PT,N]$ and $\Sigma[PT,N]$

$$\omega[N] = \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle}, \quad \omega[P_T] = \frac{\langle P_T^2 \rangle - \langle P_T \rangle^2}{\langle P_T \rangle}, \quad \omega(p_T) = \frac{\overline{p_T^2} - \overline{p_T}^2}{\overline{p_T}}$$

$$\Sigma[P_T, N] = \frac{1}{C_\Sigma} [\langle N \rangle \omega[P_T] + \langle P_T \rangle \omega[N] - 2 \cdot (\langle P_T \cdot N \rangle - \langle P_T \rangle \langle N \rangle)]$$

$$\Delta[P_T, N] = \frac{1}{C_\Delta} [\langle N \rangle \omega[P_T] - \langle P_T \rangle \omega[N]], \quad C_\Sigma = C_\Delta = \langle N \rangle \omega(p_T)$$

[M. I. Gorenstein and M. Gaździcki, Physical Review C **84**, 014904 (2011)]

Some NA61/SHINE results:

D.P. (2019) EPJ Web of Conferences **204** 07013

+ their dependence on the rapidity interval width

Normalization:

- $\Sigma[P_T, N] = \Delta[P_T, N] = 1$ for independent particle model
- $\Sigma[P_T, N] = \Delta[P_T, N] = 1$ for the IBG in GCE and CE
- $\Sigma[P_T, N] = \Delta[P_T, N] = 0$ in the absence of fluctuations

In the analysis of the experimental data this one window studies correspond to changing rapidity-averaged baryo-chemical potential at the freeze-out stage

[Becattini F, Manninen J and Gazdzicki M PRC 73 044905]

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Some NA61/SHINE results:

D.P. (2019) EPJ Web of Conferences **204** 07013

$$\omega[N] = \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle}, \quad \omega[P_T] = \frac{\langle P_T^2 \rangle - \langle P_T \rangle^2}{\langle P_T \rangle}, \quad \omega(p_T) = \frac{\overline{p_T^2} - \overline{p_T}^2}{\overline{p_T}}$$

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[Becattini F, Manninen J and Gazdzicki M PRC 73 044905]

Strongly intensive $\Sigma[N_F, N_B]$ in two kinematically separated regions of η :

$$\Sigma[N_F, N_B] = \frac{1}{C_\Sigma} [\langle N_B \rangle \omega[N_F] + \langle N_F \rangle \omega[N_B] - 2 \cdot (\langle N_F \cdot N_B \rangle - \langle N_F \rangle \langle N_B \rangle)]$$

[E. V. Andronov, Theoretical and Mathematical Physics **185**, 1383 (2015)]

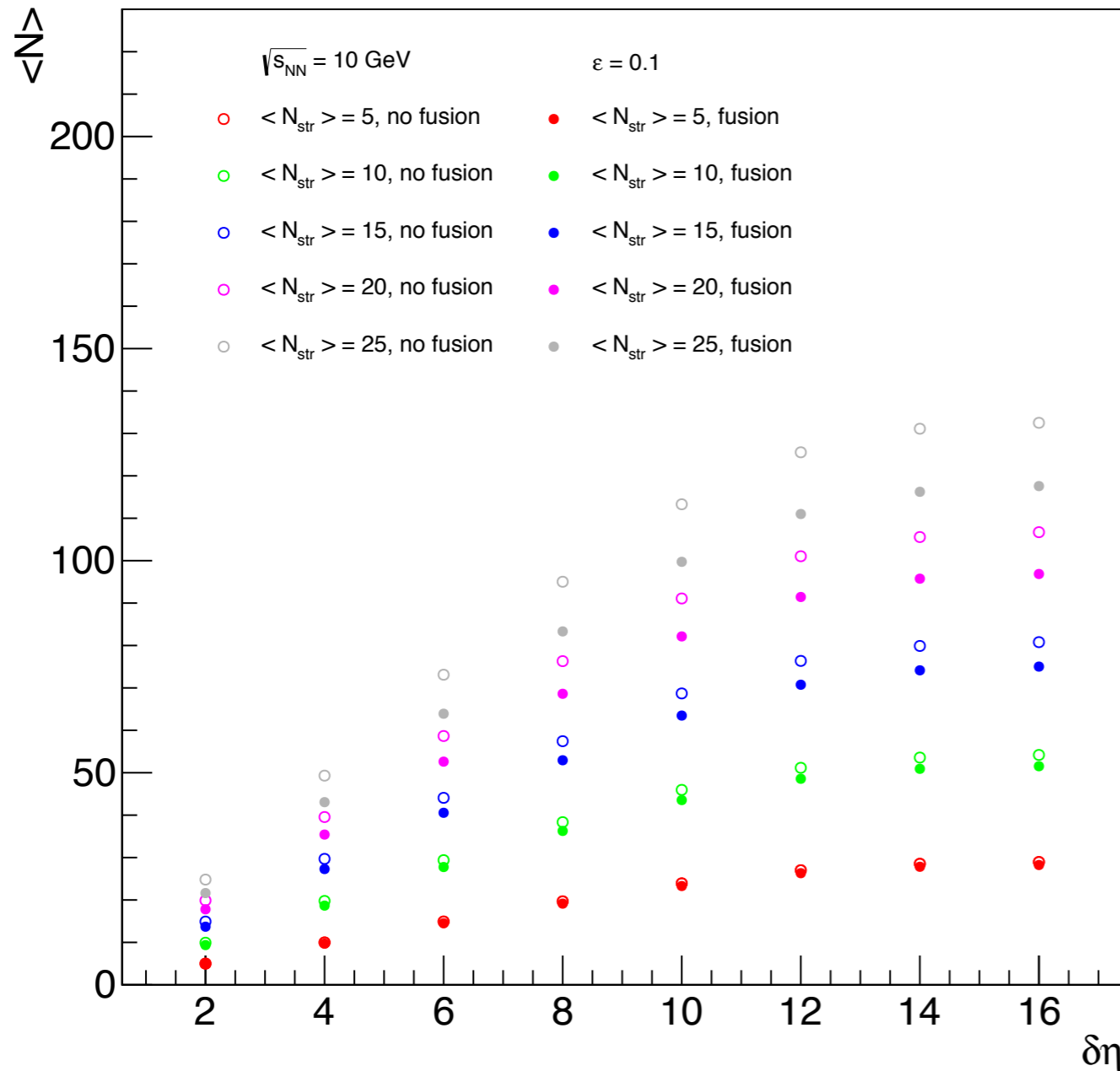
+ its dependence on the distance between Forward and Backward rapidity intervals

It is supposed to be sensitive to the initial conditions of particle production and short- and long-range multiplicity correlations

[E. Andronov, V. Vechernin, 1808.09770]

Results in MC model of interacting quark-gluon strings

$\langle N \rangle$ as a function of the rapidity window width $\delta\eta$



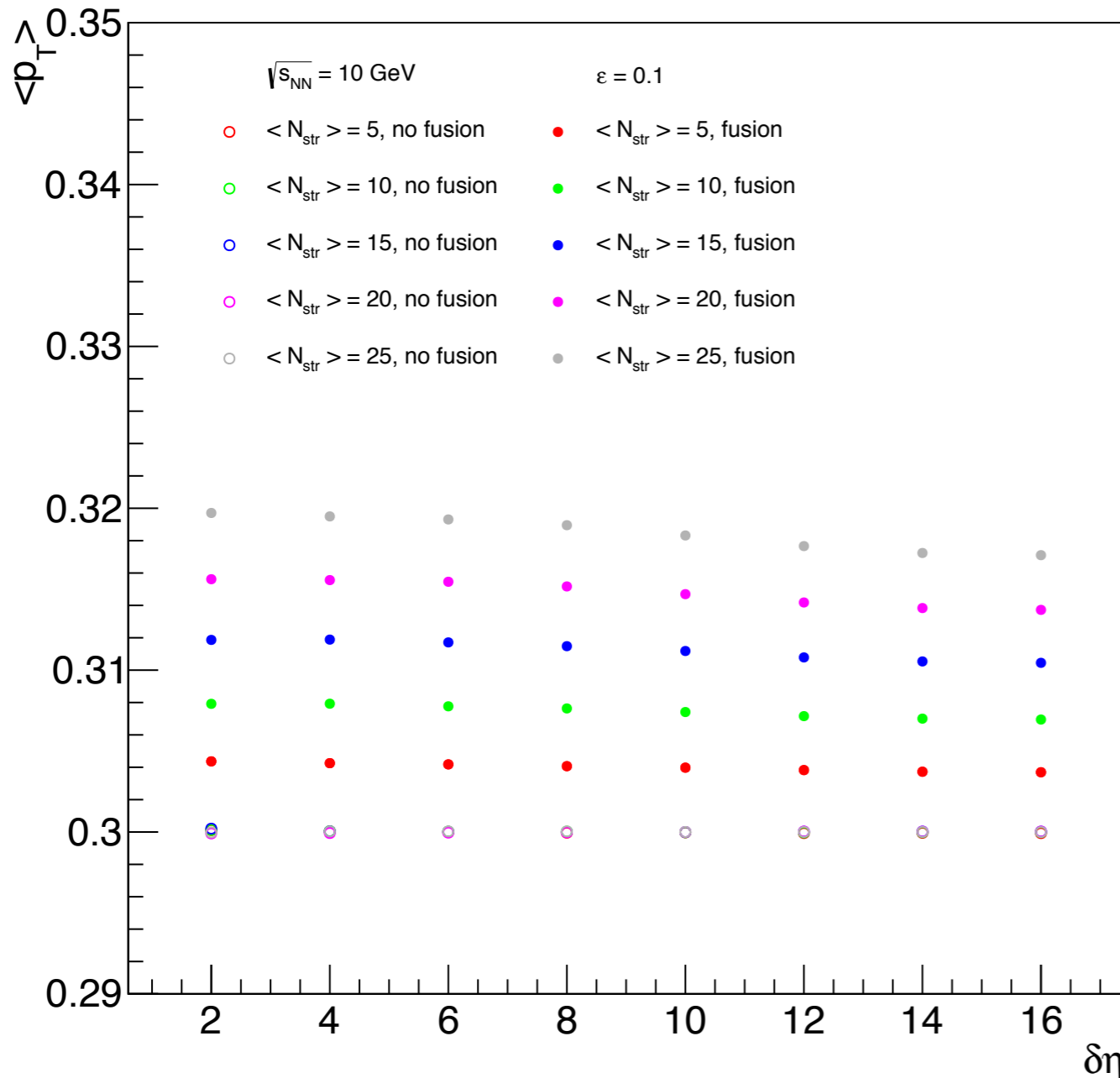
- $\langle N \rangle$ grows with the rapidity window width
- $\langle N \rangle$ grows with the number of strings in event
- string fusion causes the decrease of $\langle N \rangle$

$$\langle \mu \rangle_k = \mu_0 \sqrt{k-1}$$

- the larger the number of strings in event, the bigger the difference between fusion and non-fusion

Results in MC model of interacting quark-gluon strings

$\langle p_T \rangle$ as a function of the rapidity window width $\delta\eta$



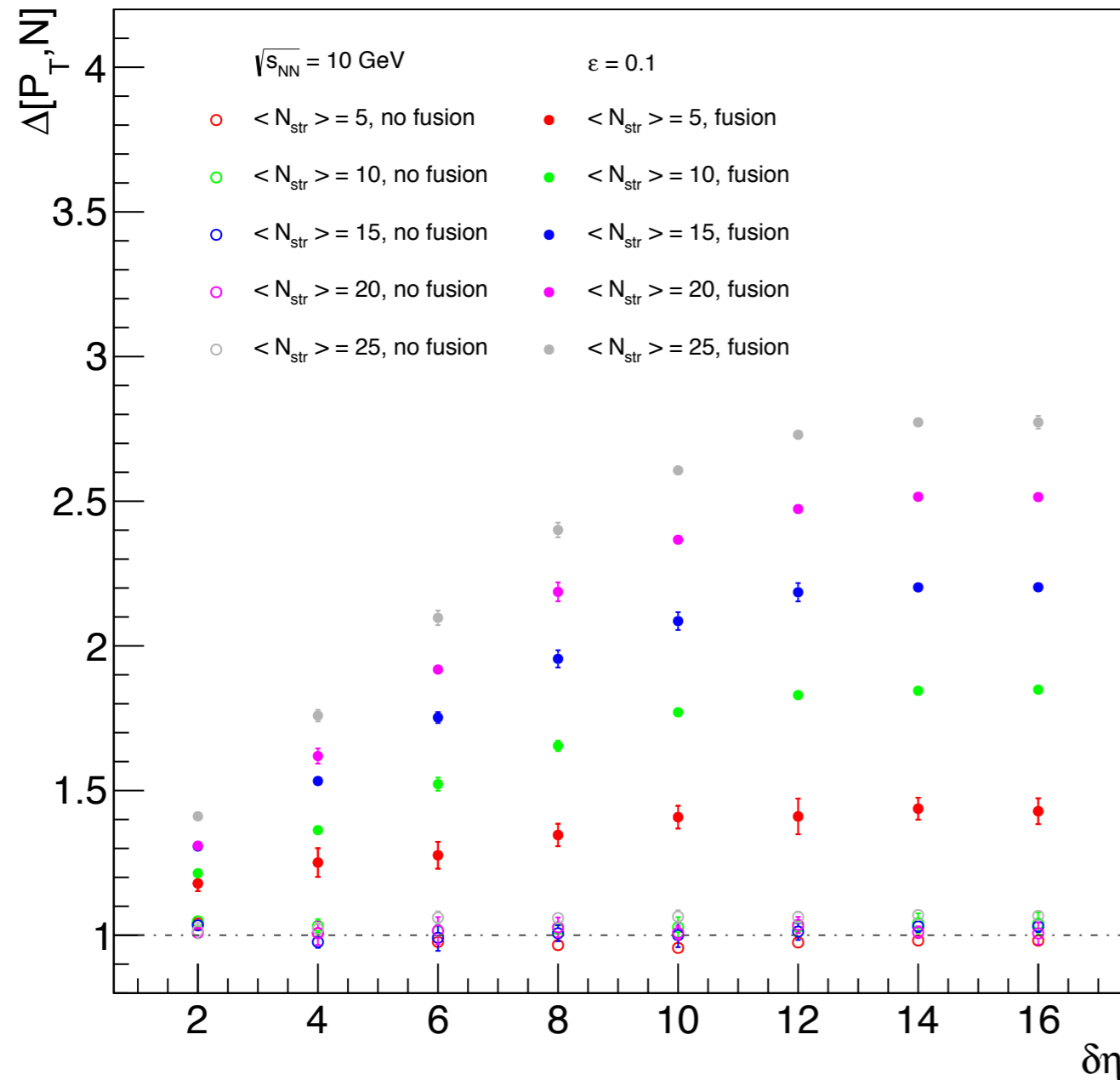
- $\langle p_T \rangle$ is constant with the rapidity window width
- $\langle p_T \rangle$ grows with the number of strings in event
- string fusion causes the increase of $\langle p_T \rangle$

$$\langle p_T^2 \rangle_k = p_0^2 \sqrt{k},$$

- the larger the number of strings in event, the bigger the difference between fusion and non-fusion

Results in MC model of interacting quark-gluon strings

$\Delta[PT,N]$ as a function of the rapidity window width $\delta\eta$

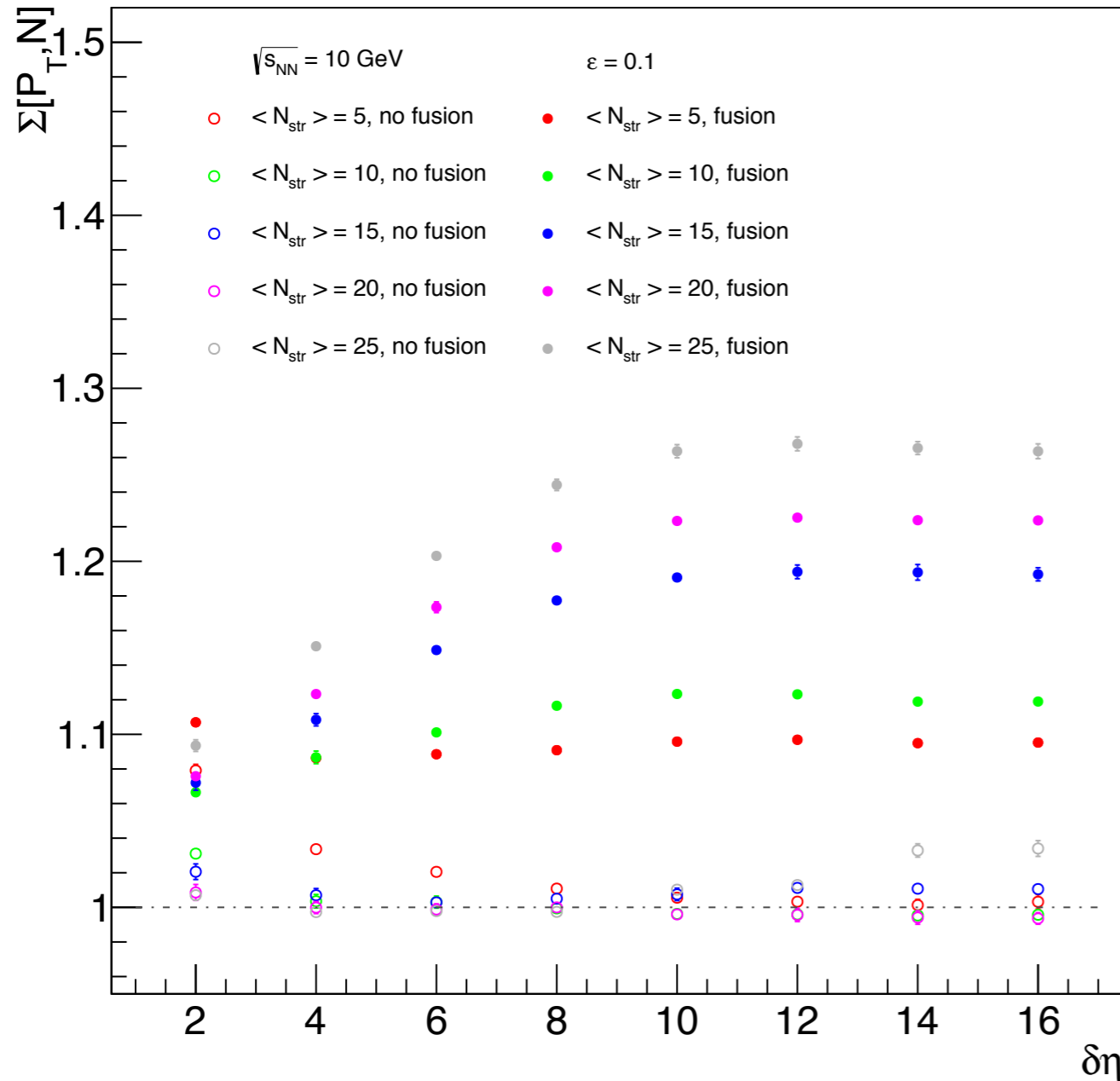


$$\Delta[PT, N] = \frac{1}{C_{\Delta}} [\langle N \rangle \omega[PT] - \langle PT \rangle \omega[N]]$$

- for no-fusion $\Delta[PT,N]$ fluctuates around 1
- string fusion causes the increase of $\Delta[PT,N]$ with the rapidity window width
- for fusion the larger number of strings in event, the larger the value of $\Delta[PT,N]$

Results in MC model of interacting quark-gluon strings

$\Sigma[PT,N]$ as a function of the rapidity window width $\delta\eta$

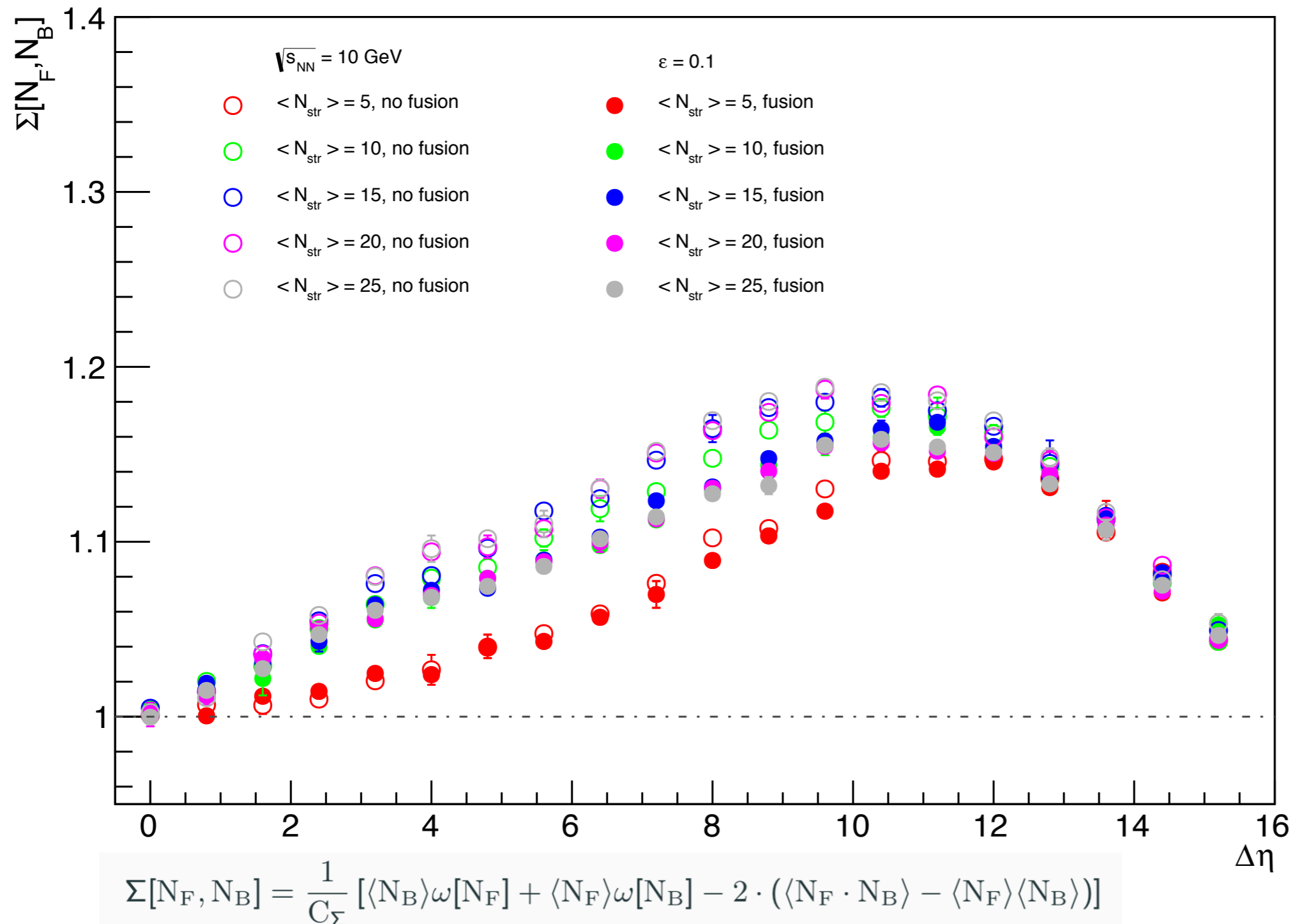


$$\Sigma[PT, N] = \frac{1}{C_\Sigma} [\langle N \rangle \omega[PT] + \langle PT \rangle \omega[N] - 2 \cdot (\langle PT \cdot N \rangle - \langle PT \rangle \langle N \rangle)]$$

- for no-fusion $\Sigma[PT,N]$ fluctuates around 1, but more then $\Delta[PT,N]$
- string fusion causes the increase of $\Sigma[PT,N]$ with the width of the rapidity window
- for fusion the larger the number of strings in event, the larger the value of $\Sigma[PT,N]$

Results in MC model of interacting quark-gluon strings

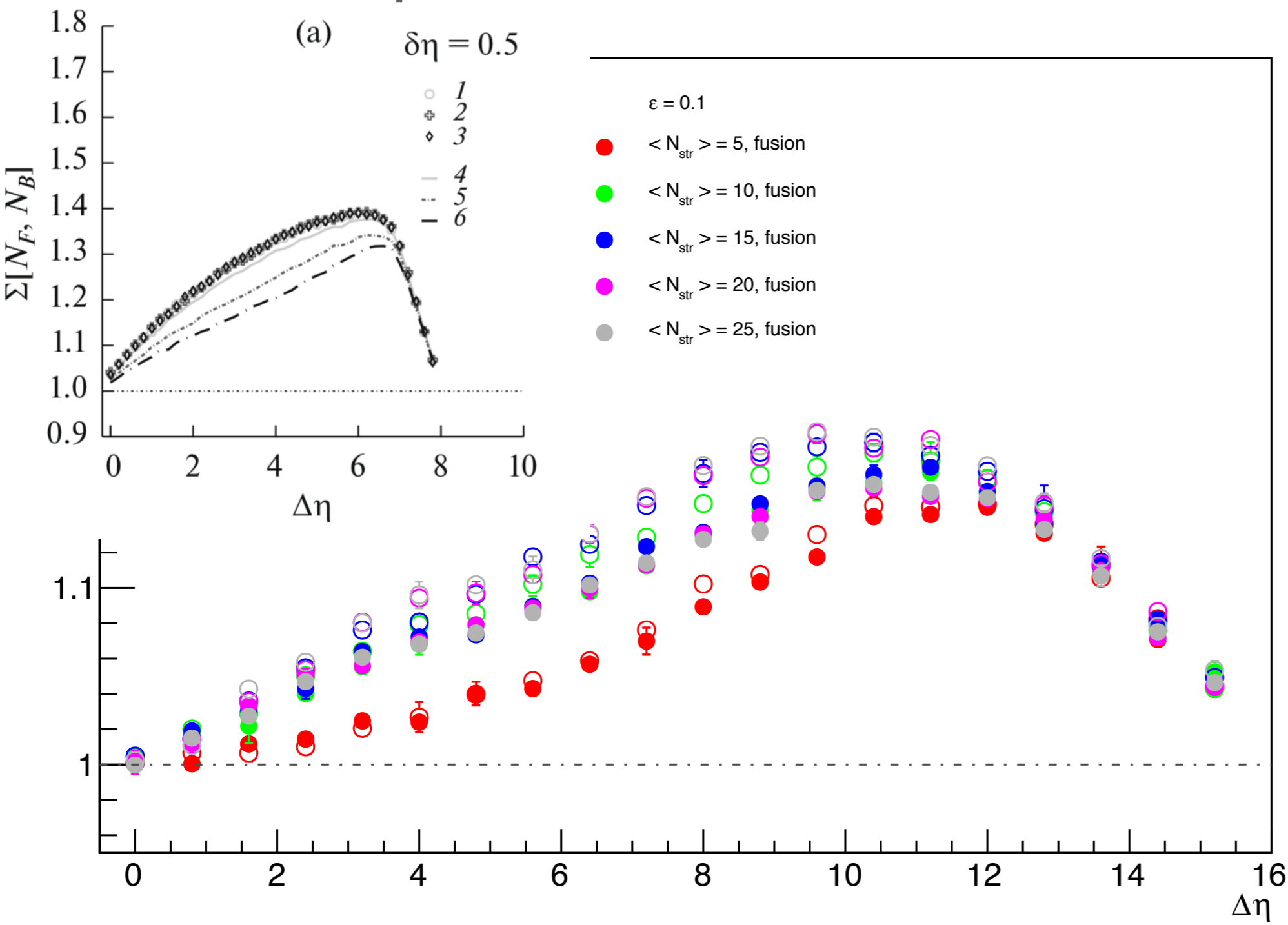
$\Sigma[N_F, N_B]$ as a function of the gap $\Delta\eta$ between two rapidity windows



- values of $\Sigma[N_F, N_B]$ for no-fusion are larger than the ones for fusion
- $\Sigma[N_F, N_B]$ grows with the distance between rapidity windows
- $\Sigma[N_F, N_B]$ decreases at high distances between rapidity windows due to the lack of statistics

Results in MC model of interacting quark-gluon strings

Results in the previous version of the model D. P., V.N. Kovalenko (2020) PEPAN 51 (3) 323-326

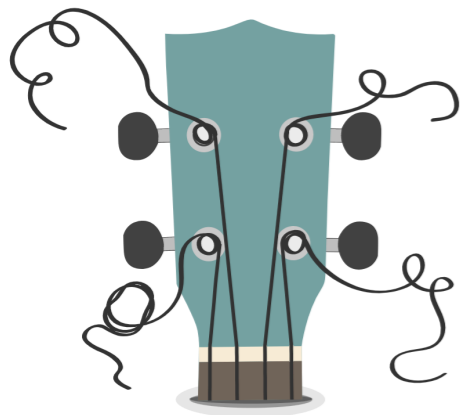


- values of $\Sigma[NF, NB]$ for no-fusion are larger than the ones for fusion
- $\Sigma[NF, NB]$ grows with the distance between rapidity windows
- $\Sigma[NF, NB]$ decreases at high distances between rapidity windows due to the lack of statistics

Outlook

Study of multiplicity and transverse momentum fluctuations in the Monte-Carlo model of interacting quark-gluon strings

1. Introduce charges of the produced particles
2. Introduce baryons and mesons
3. Implement the explicit energy dependence of the string number



Thank you for your attention!

This work is supported by the RFBR research project no. 18-02-40097

Backup

Monte-Carlo model of interacting quark-gluon strings

Procedure: prerequisites

- 1) Take quarks x_f probability distribution functions from LHAPDF
- 2) Number of strings in event is sampled from the poisson distribution with the given mean number
- 3) Mean number of particles produced by the rapidity unit of a string: $\mu=1$
- 4) Mean p_T of particles produced by the string: $p_{T\text{mean}}=0.3$
- 5) $\sqrt{s_{NN}}=10\text{GeV}$
- 6) Mean active piece of the string: 0.1, 0.2, 0.4
- 7) Size of a transverse plane of cells: 25

Monte-Carlo model of interacting quark-gluon strings

Procedure: string formation

- 1) For each event sample a number of strings N_{str}
- 2) Define number of quarks as $2 \cdot N_{str}$
- 3) Generate the quark content by the random number and fulfilling the ratio $u:\bar{u}:d:\bar{d}:s:\bar{s}:c:\bar{c} = 6:2:4:2:2:2:1:1$
- 4) Generate the direction of the quark moving: +1 or -1
- 5) Define quark rapidities (+y for positive direction and -y for negative) according to:

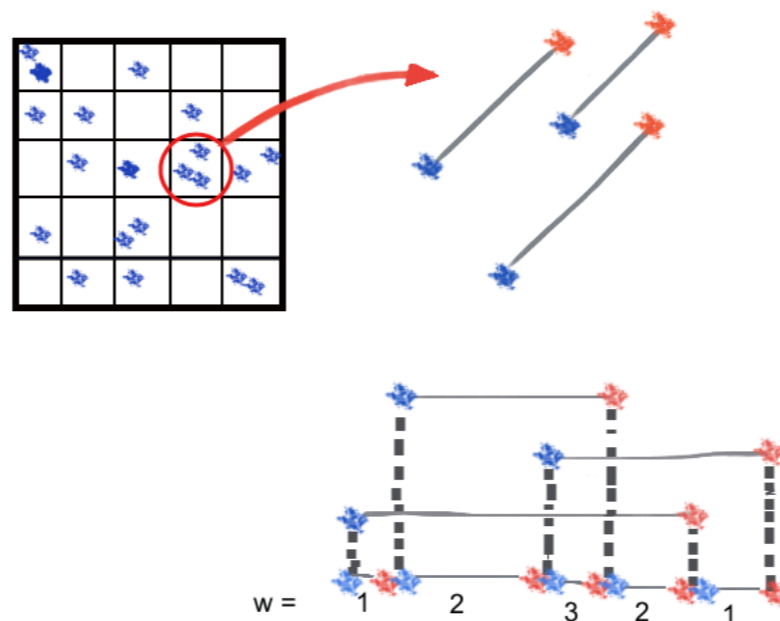
$$y_q = \operatorname{arcsinh} \left(x_q \sqrt{\frac{s}{4m_q^2} - 1} \right)$$

- 6) Define the largest one as a forward y_q and vice versa
- 7) String is formed if $\text{forward_end} - \text{backward_end} > 0.1$, otherwise the loop is repeated
- 8) Created primary strings in event, sample their position in transverse plane

Monte-Carlo model of interacting quark-gluon strings

Procedure: string fusion

- 1) For no fusion option strings for particle production are our prepared strings
- 2) For no fusion ON option we check if strings overlap in transverse plane and if so form a new list of strings (shorter ones, but with a higher weights)



Monte-Carlo model of interacting quark-gluon strings

Procedure: particles

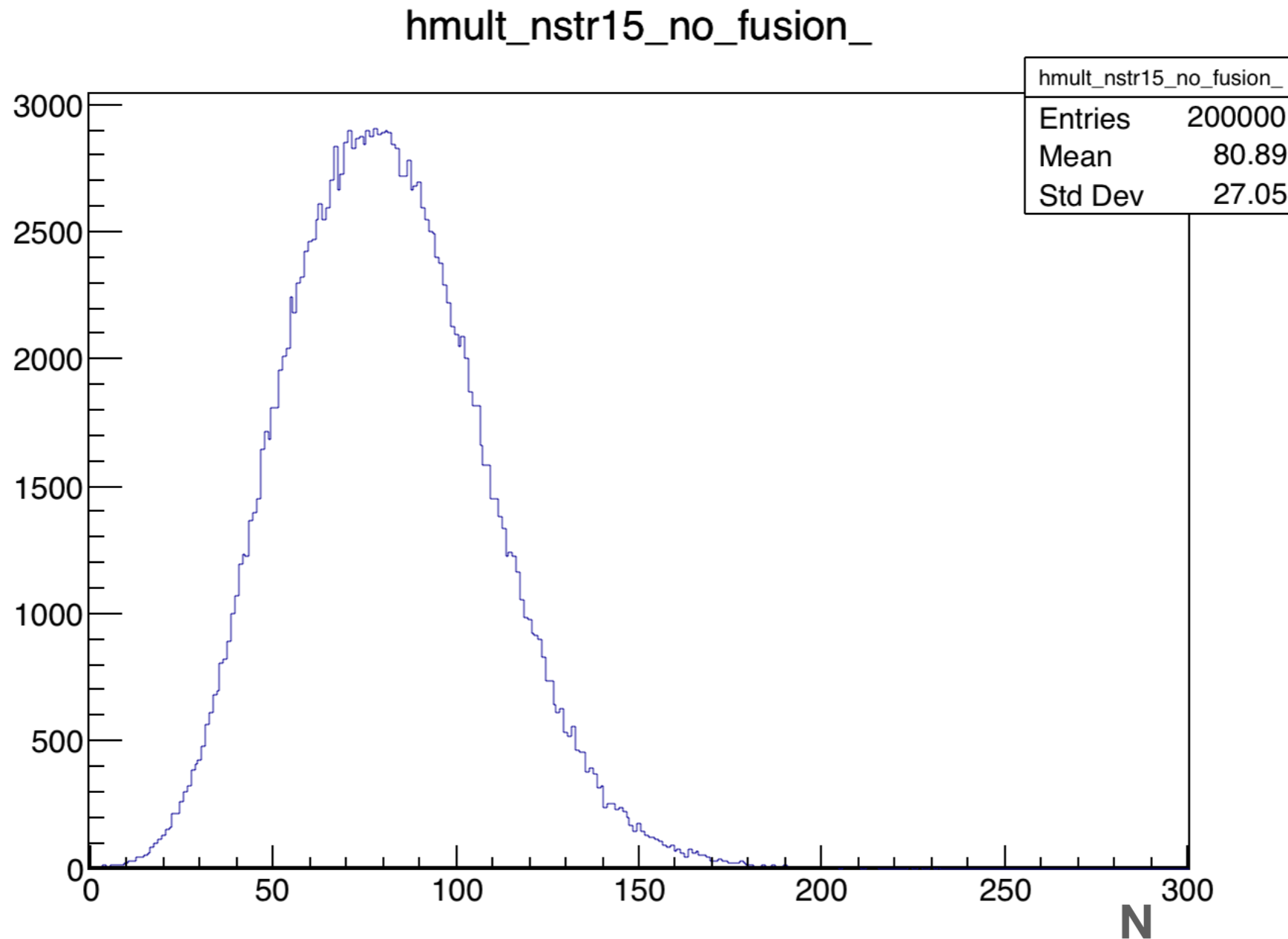
- 1) For the prepared set of strings find all active pieces of a predefined length and calculate their mean eta and mean multiplicity proportional to this length and weight defined by the number of the primary strings fused at this eta range due to the overlap in the transverse plane
- 2) Then find multiplicity by this piece from the Poisson distribution with the found mean - this is the number of particles, for each the pT is sampled from with the mean 0.3:

$$f(p_t) = \frac{\pi p_t}{2 \langle p_t \rangle_k^2} \exp\left(-\frac{\pi p_t^2}{4 \langle p_t \rangle_k^2}\right)$$

- 3) Fill histograms of final quantities

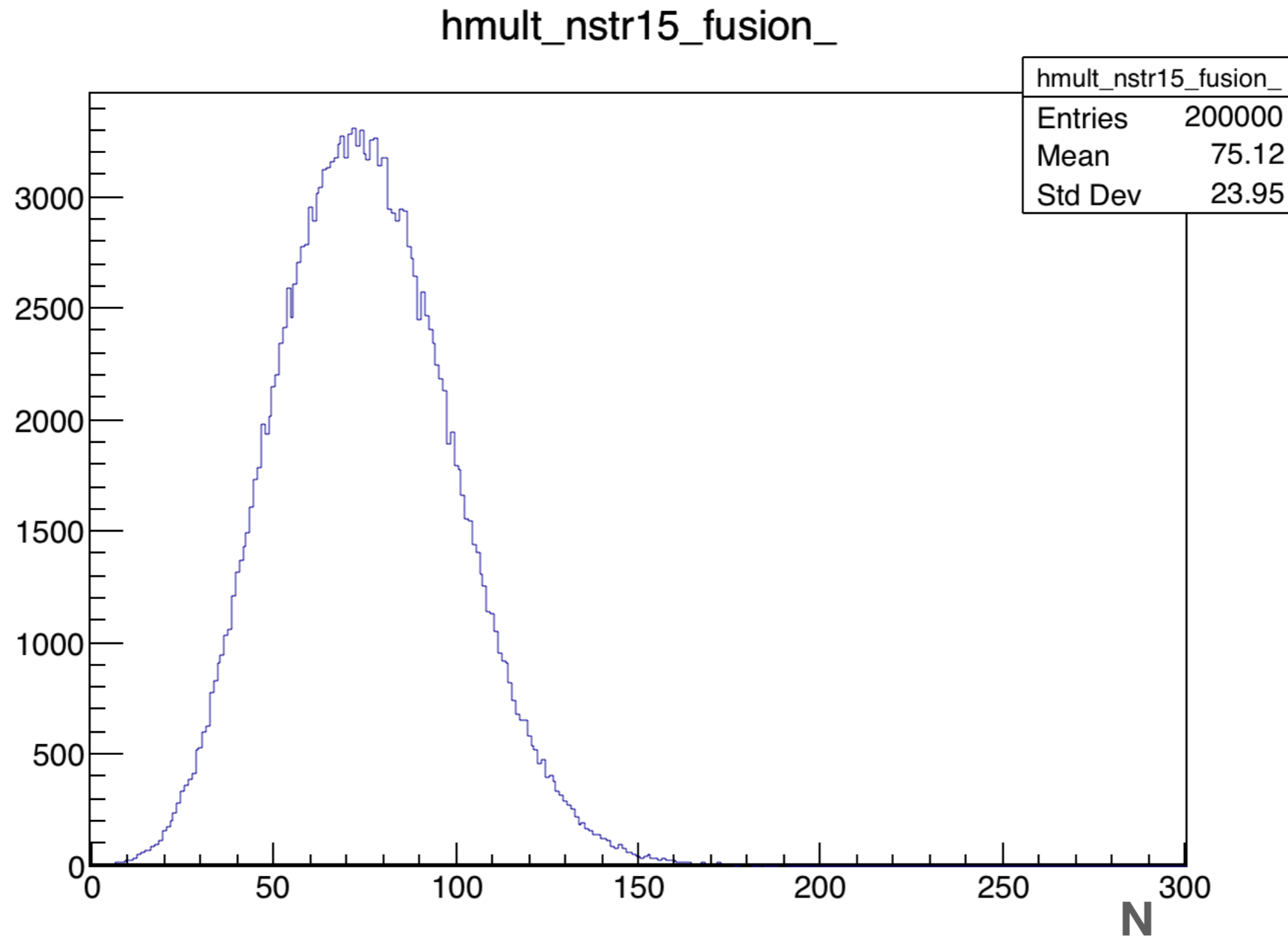
Monte-Carlo model of interacting quark-gluon strings

Basic distributions



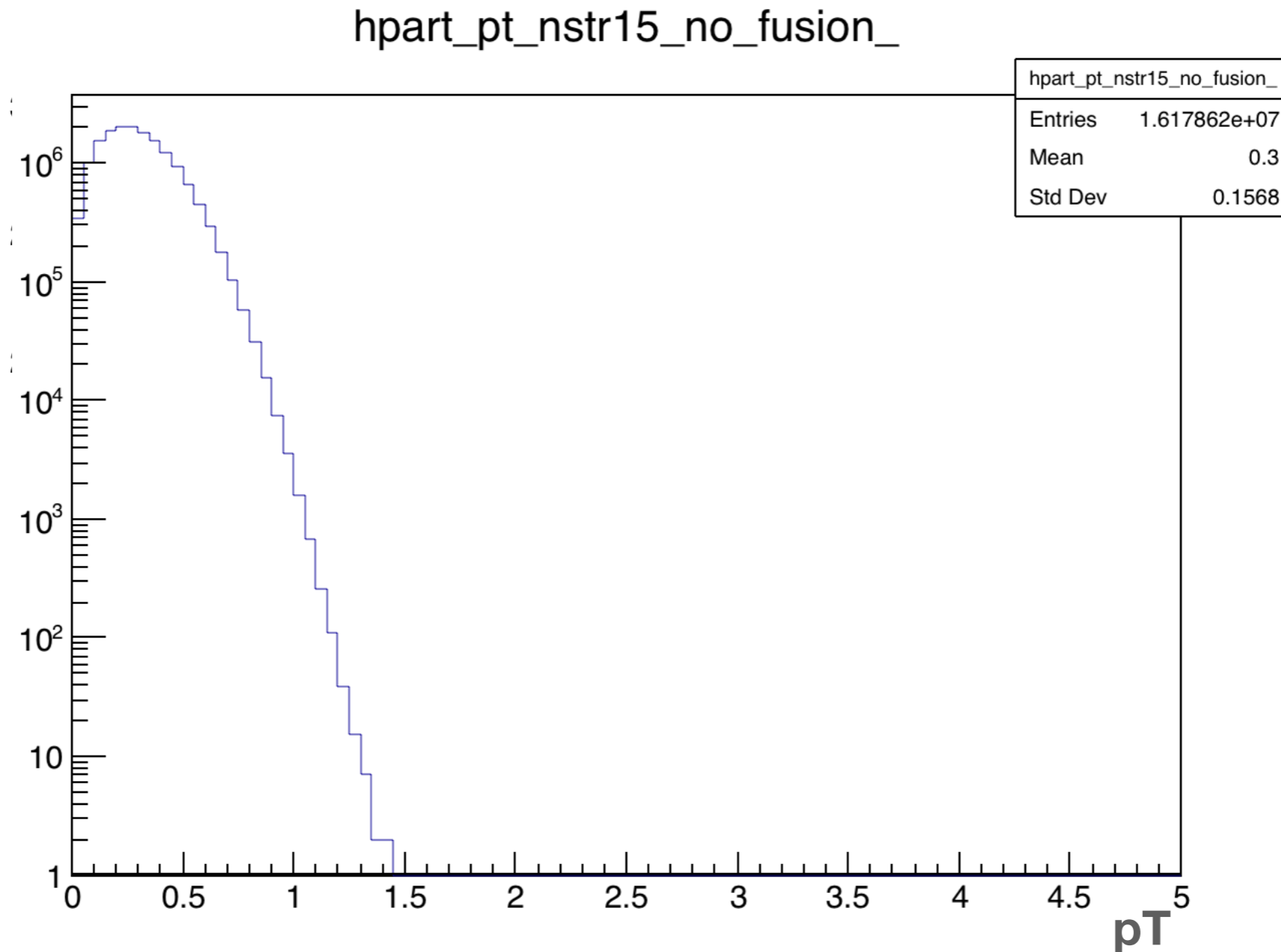
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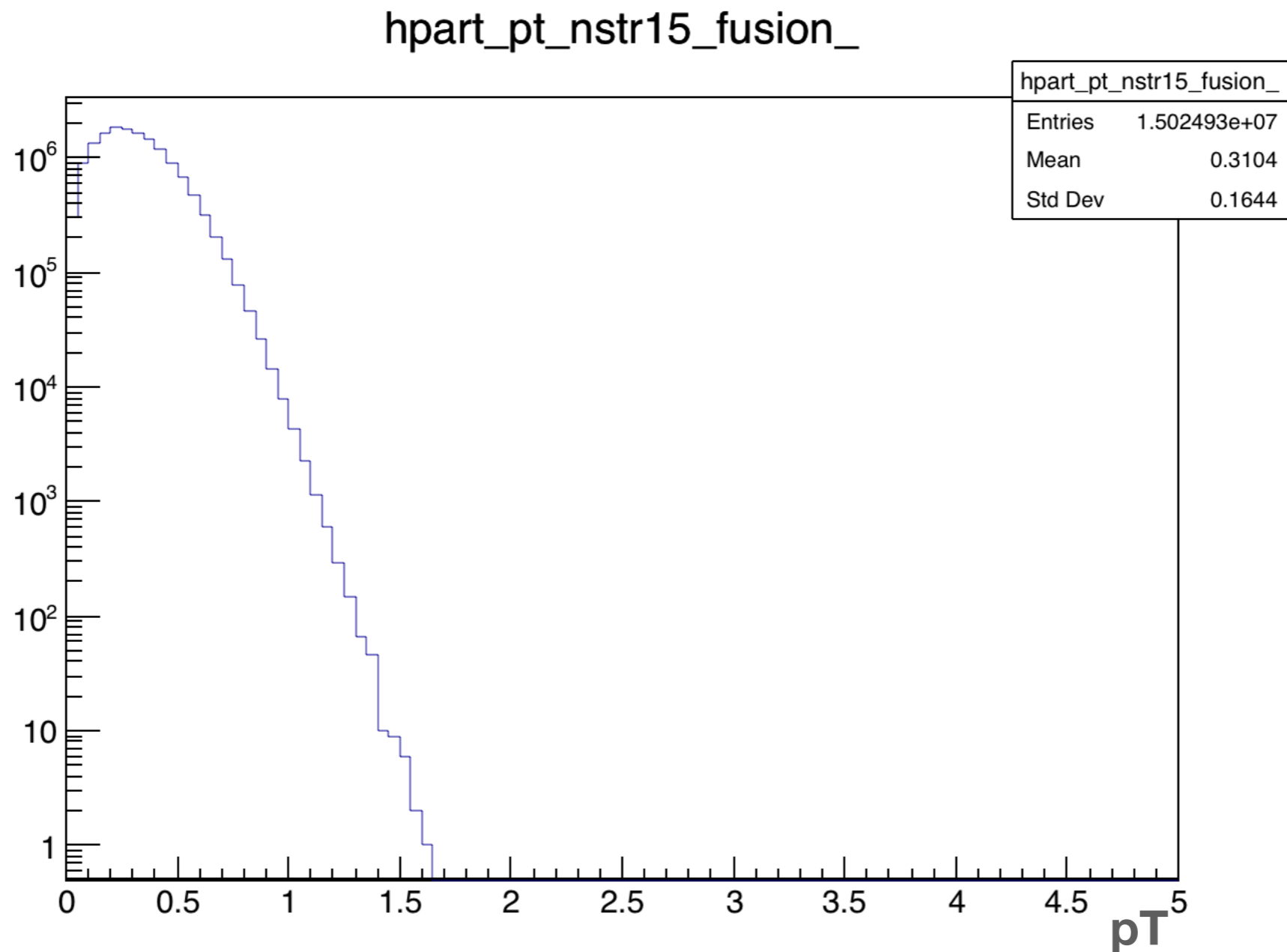
Monte-Carlo model of interacting quark-gluon strings

Basic distributions



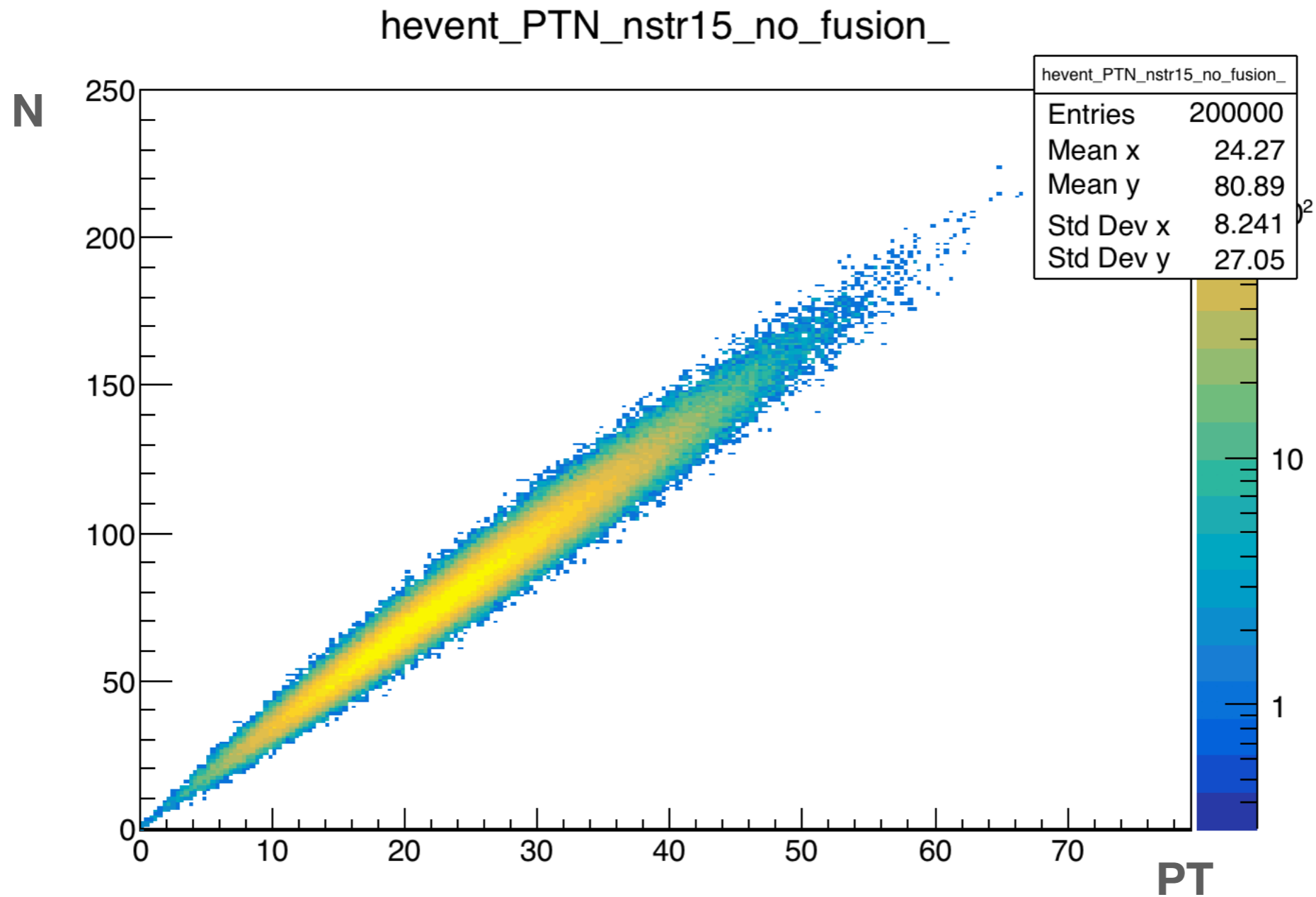
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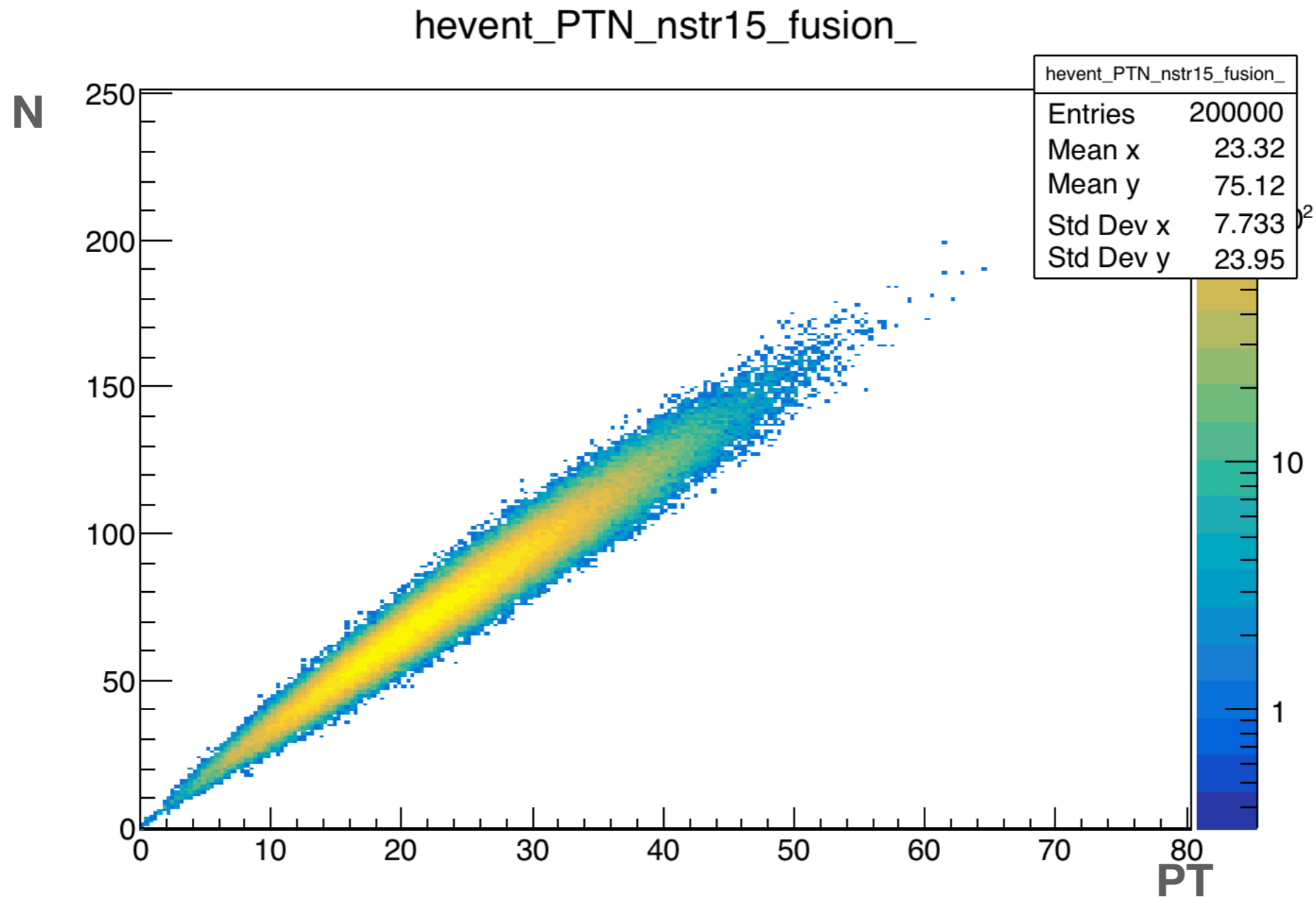
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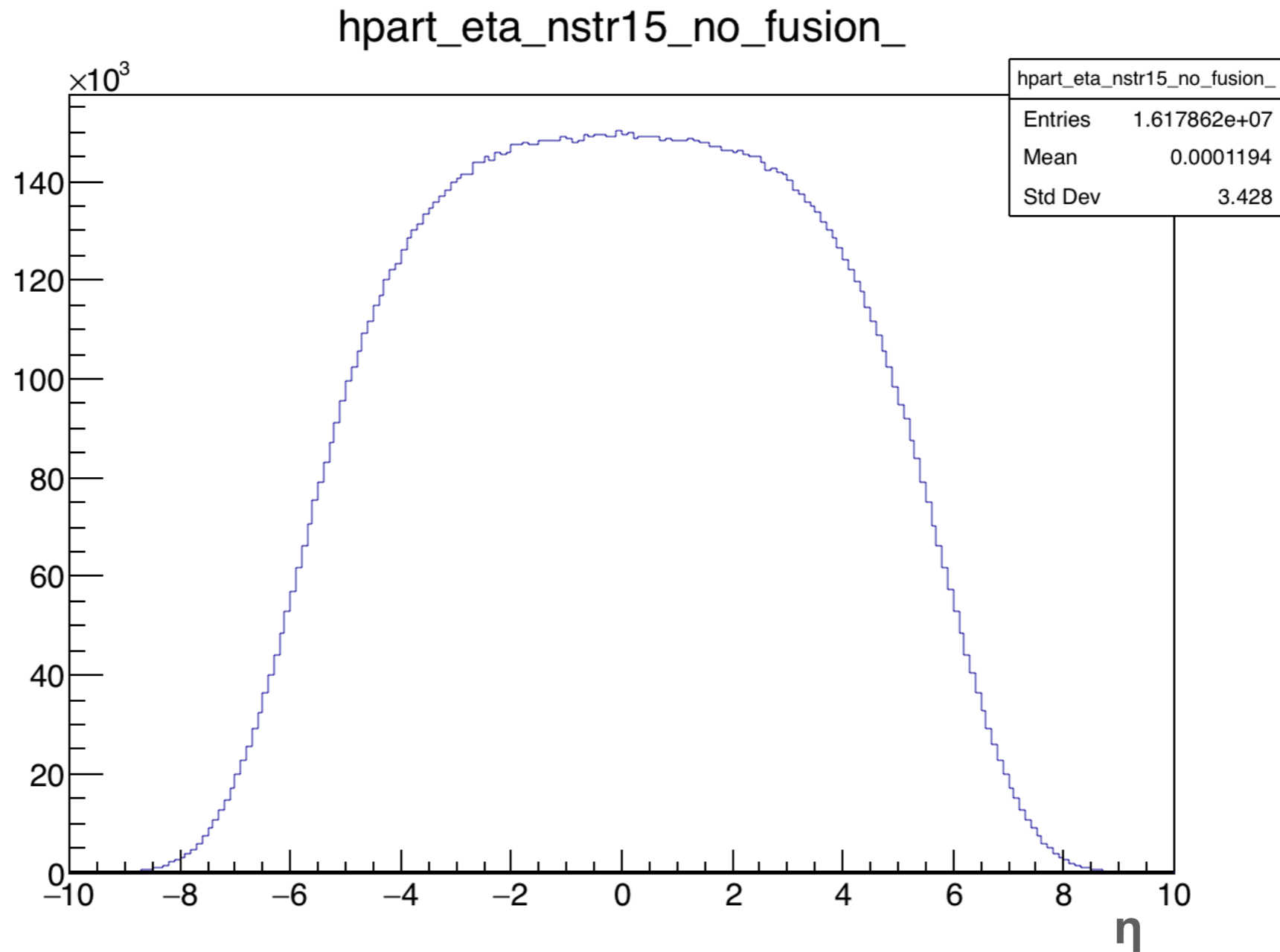
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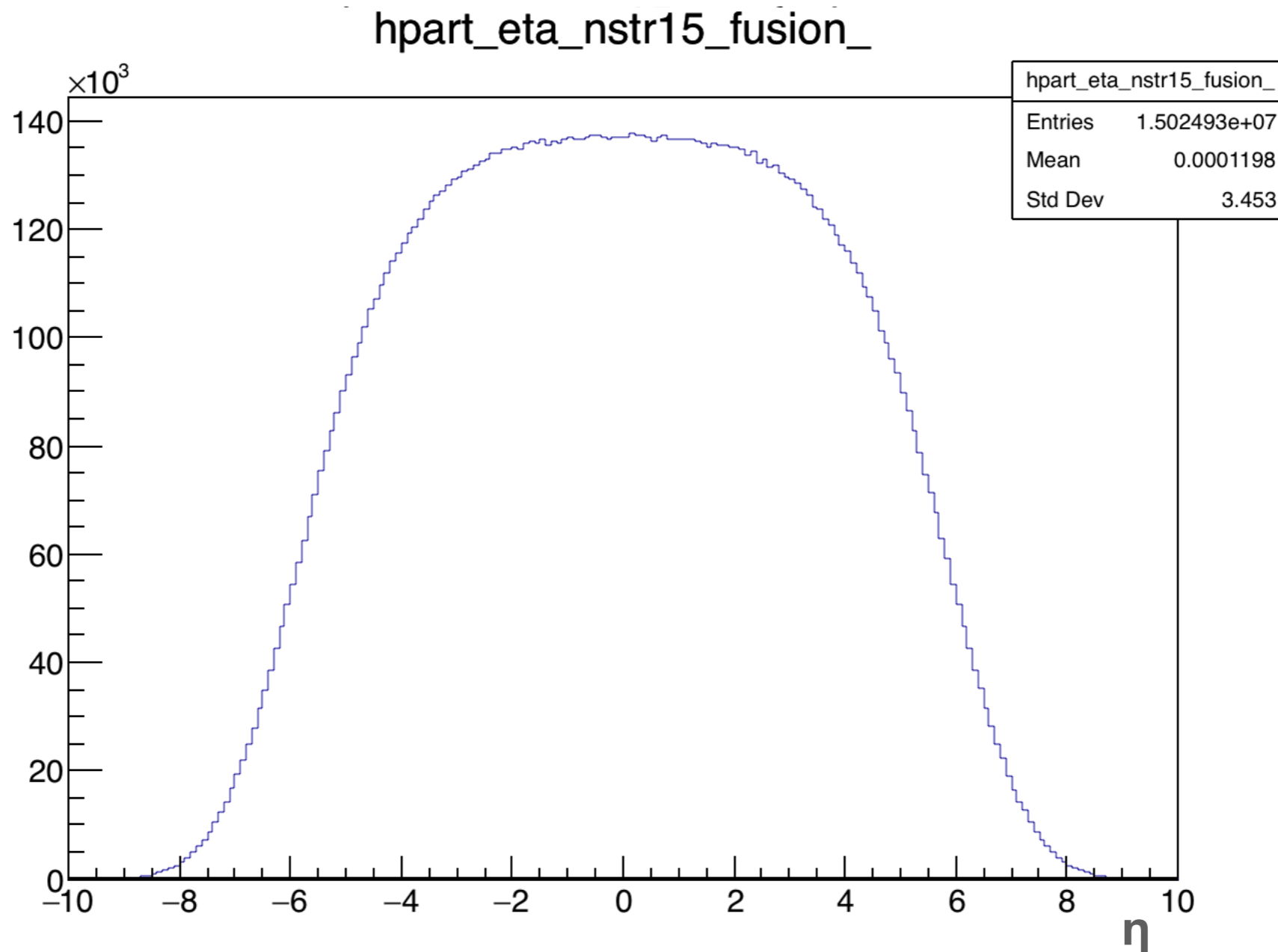
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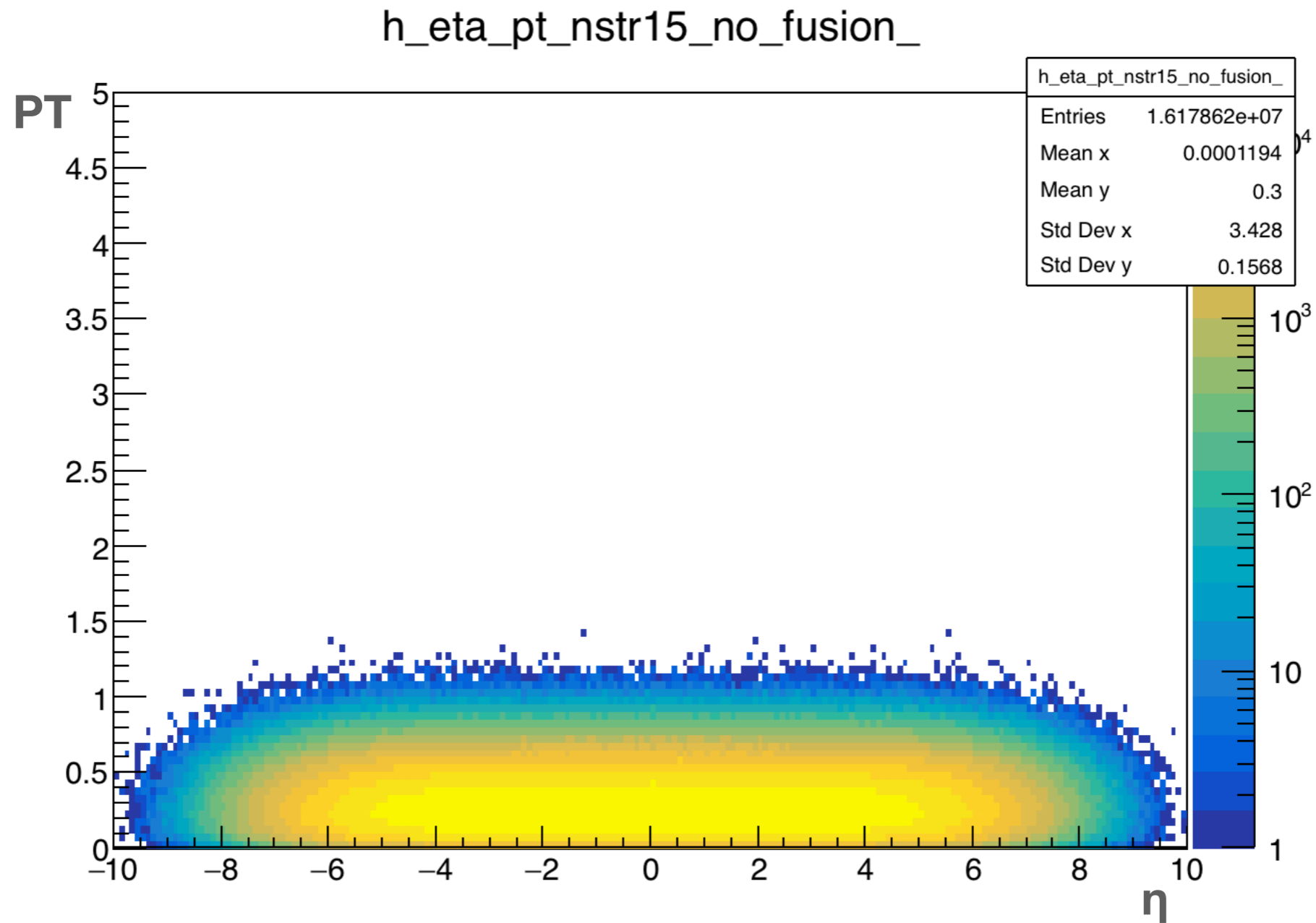
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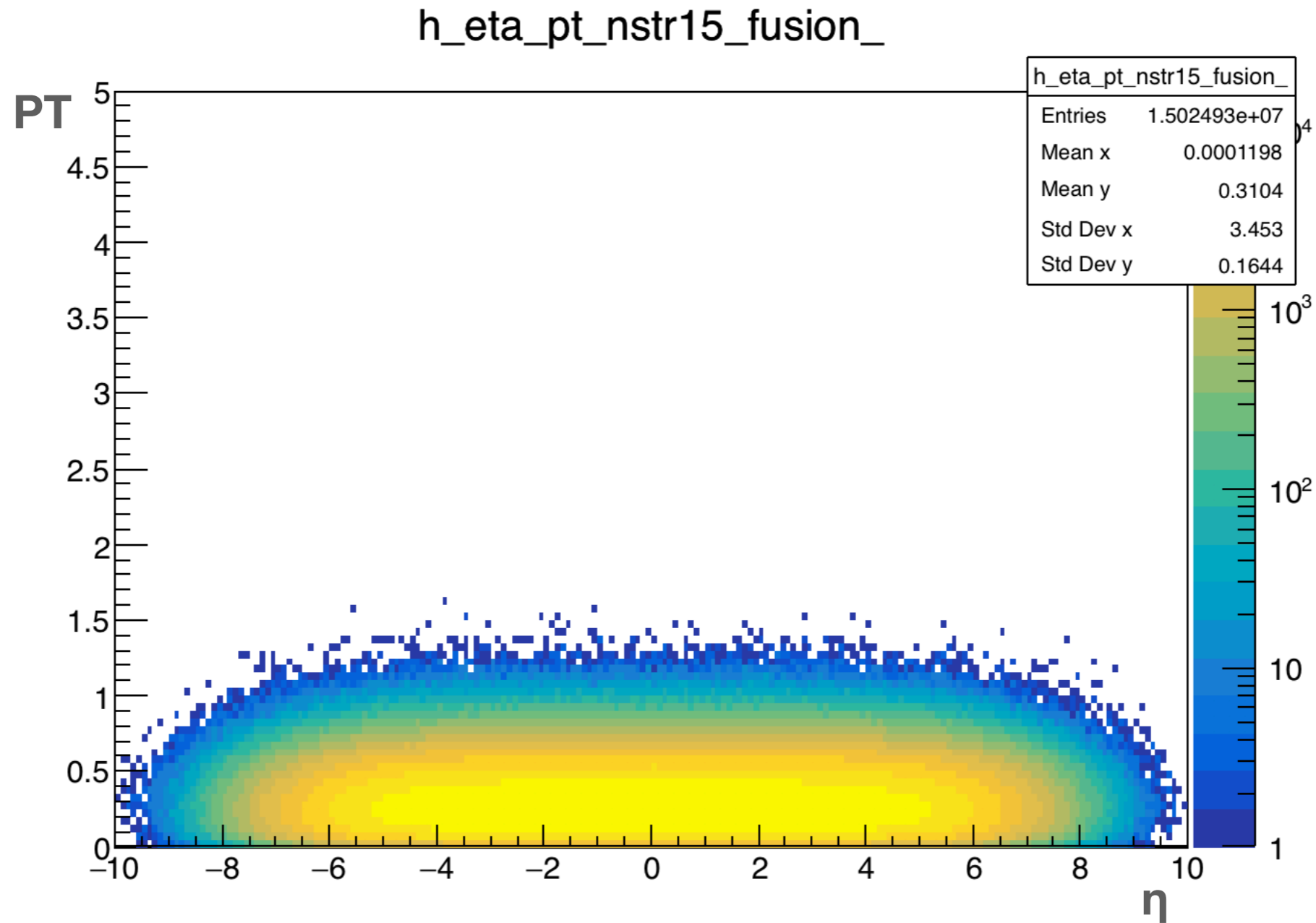
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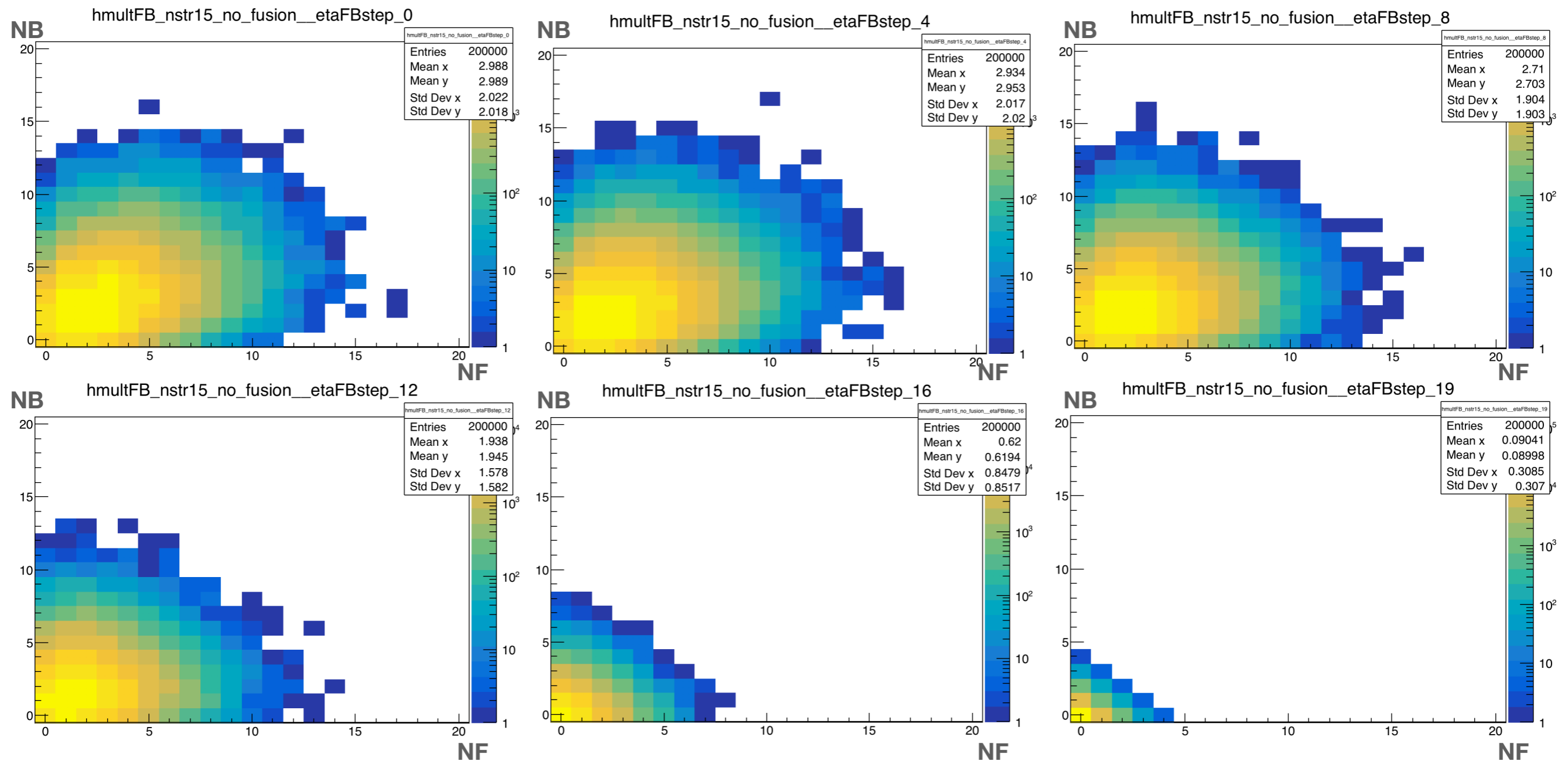
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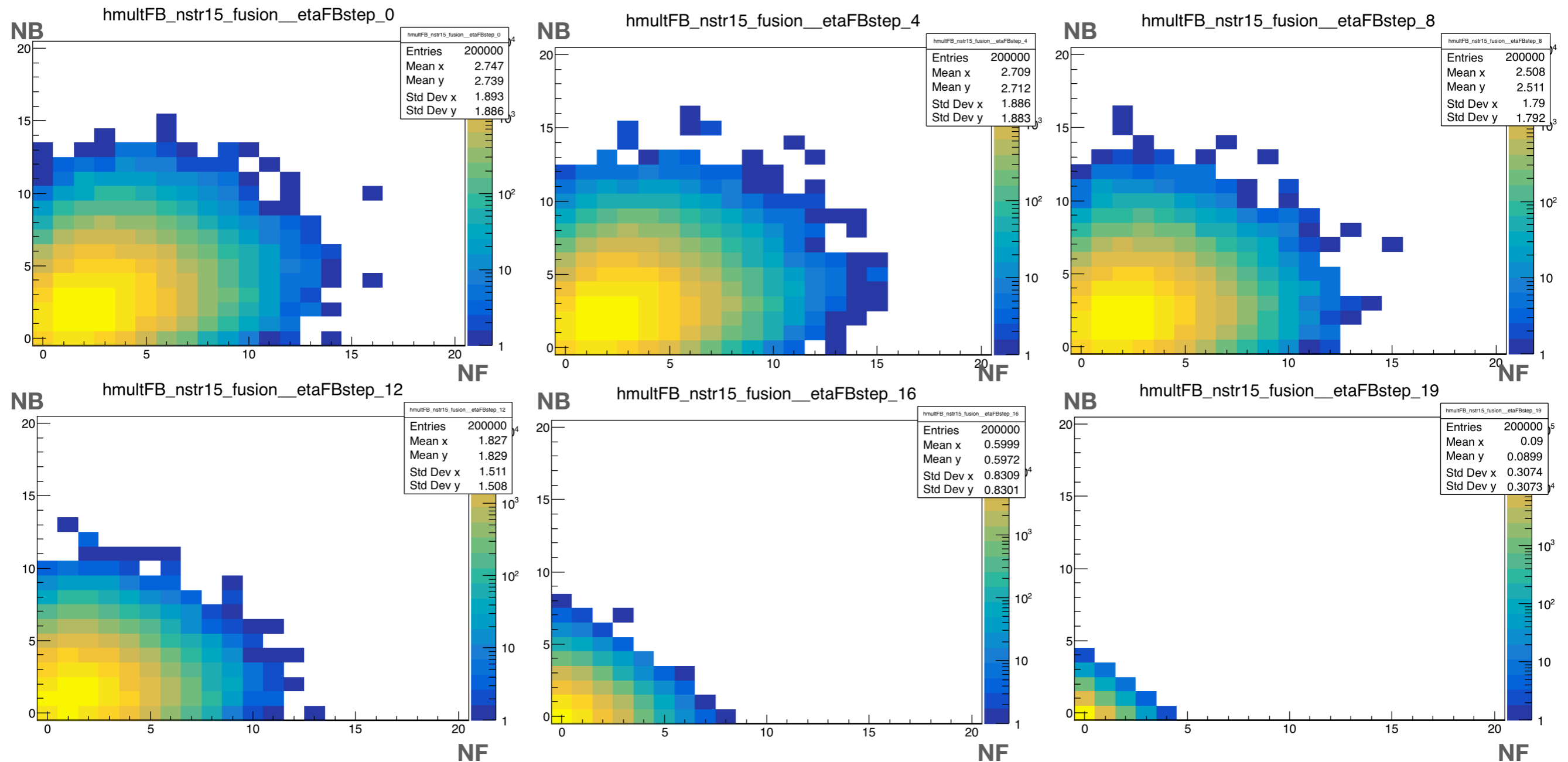
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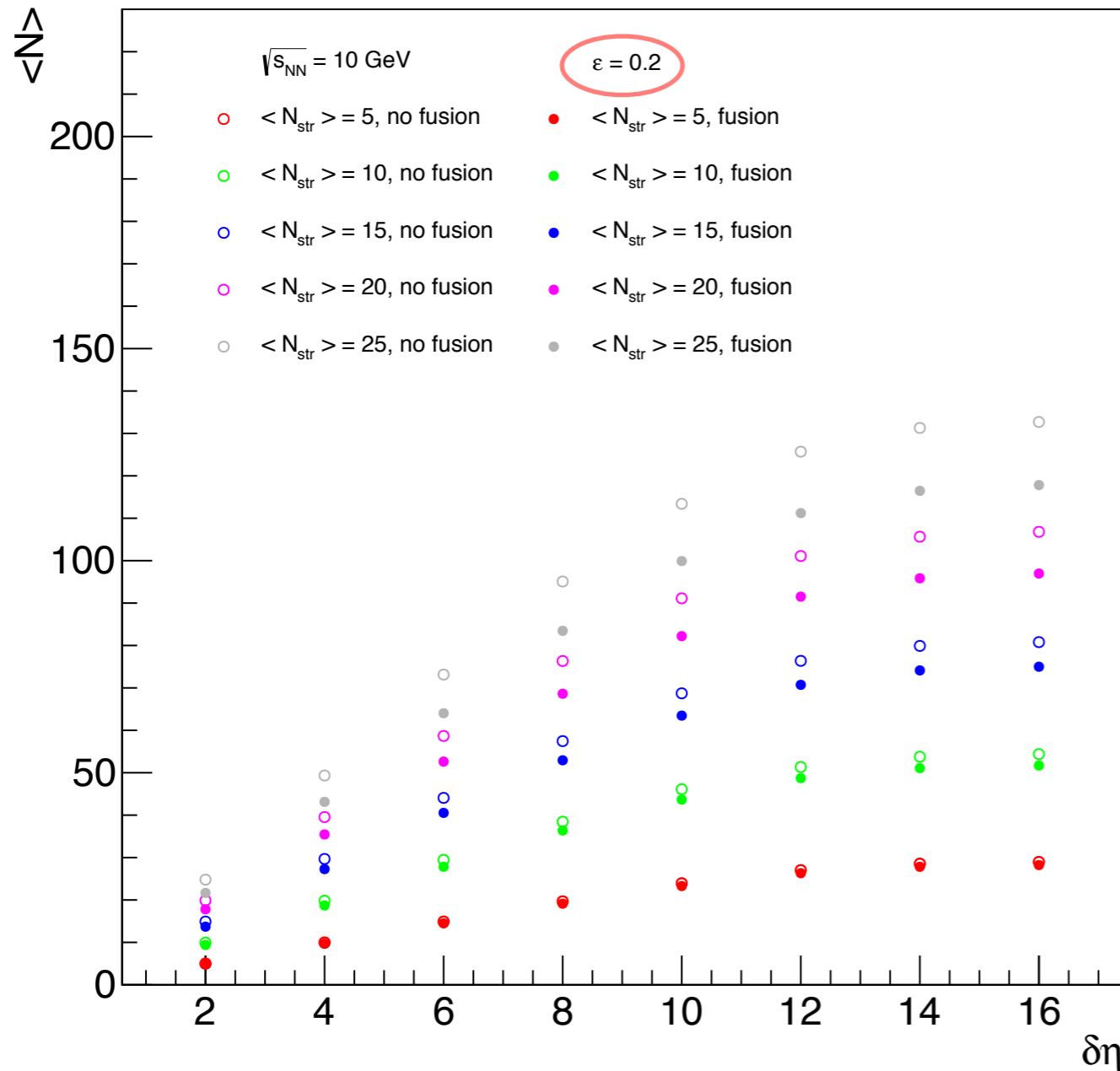
Monte-Carlo model of interacting quark-gluon strings

Basic distributions



Results in MC model of interacting quark-gluon strings

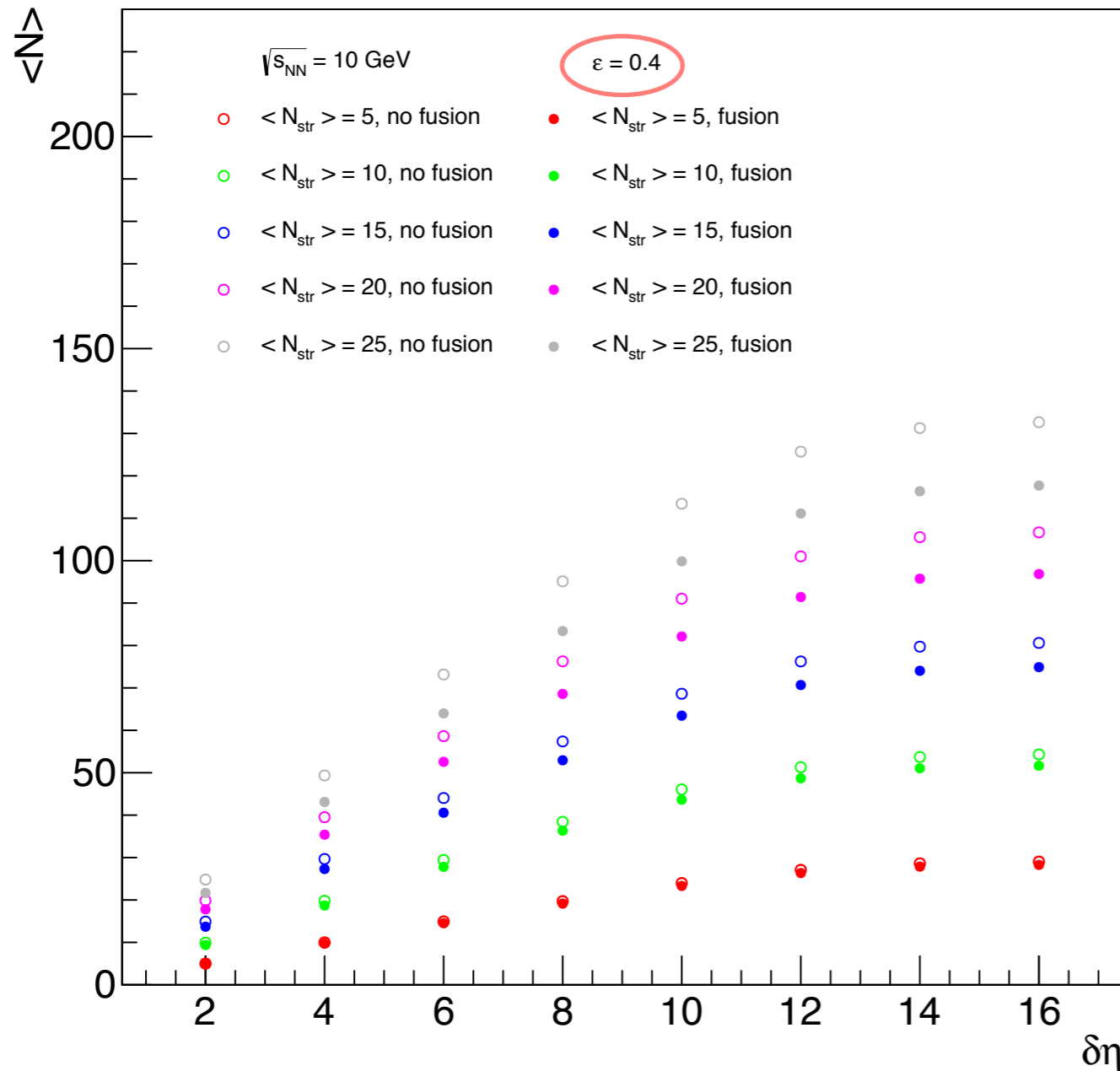
$\langle N \rangle$ as a function of the rapidity window width $\delta\eta$



- $\langle N \rangle$ grows with the rapidity window width
- $\langle N \rangle$ grows with the number of strings in event
- string fusion causes the decrease of $\langle N \rangle$
- the larger the number of strings in event, the bigger the difference between fusion and non-fusion

Results in MC model of interacting quark-gluon strings

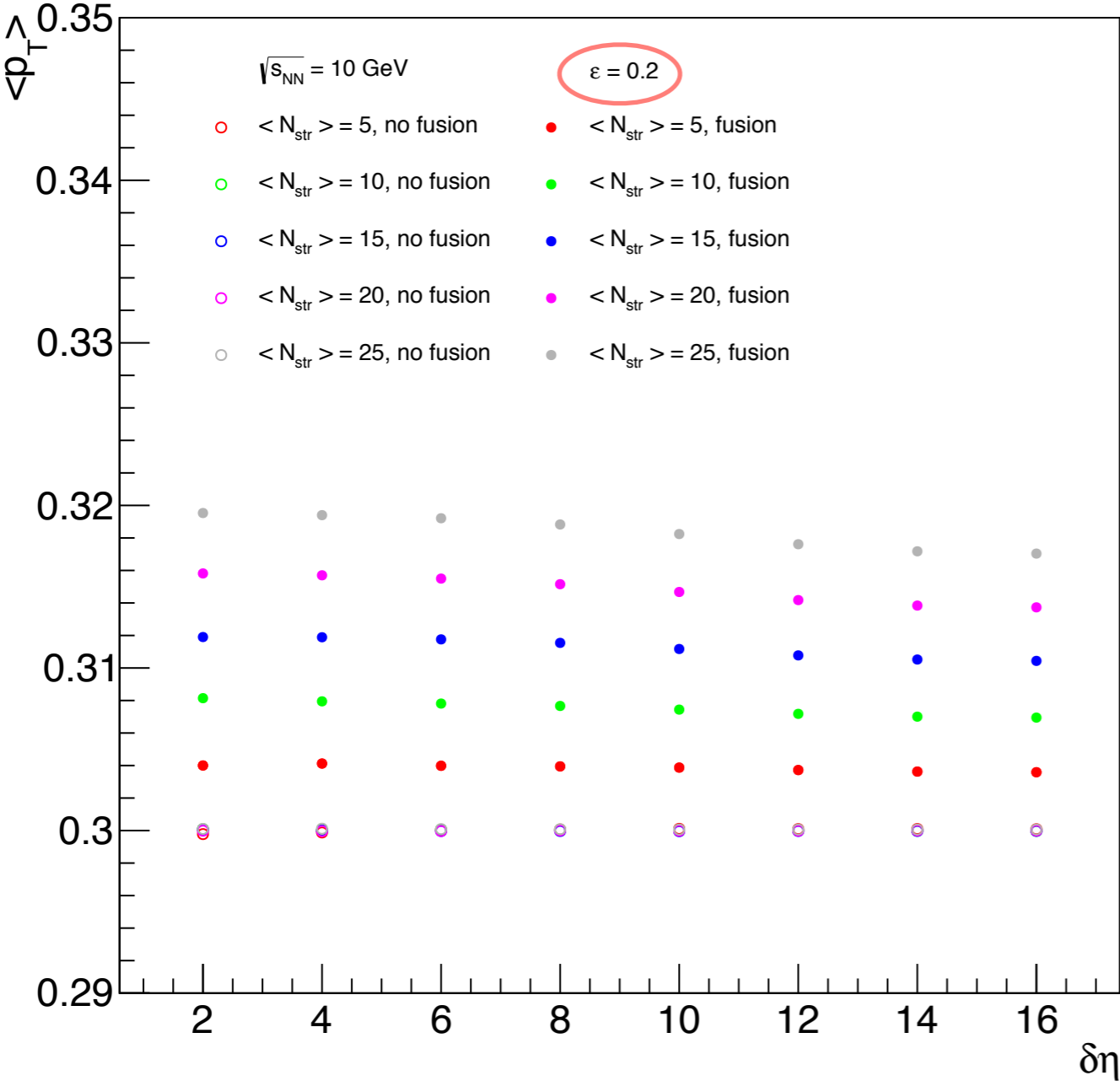
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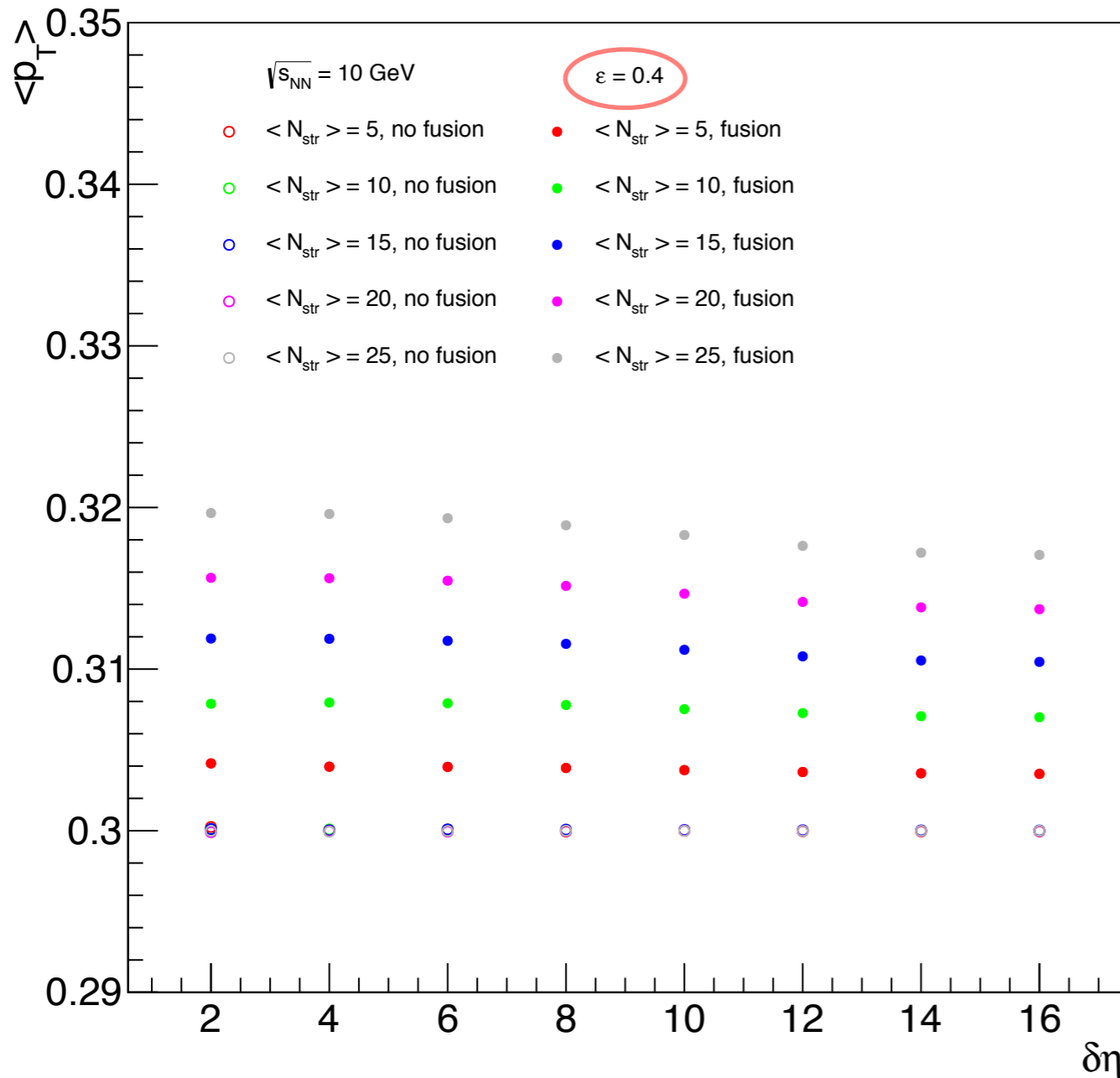
$\langle p_T \rangle$ as a function of the rapidity window width $\delta\eta$



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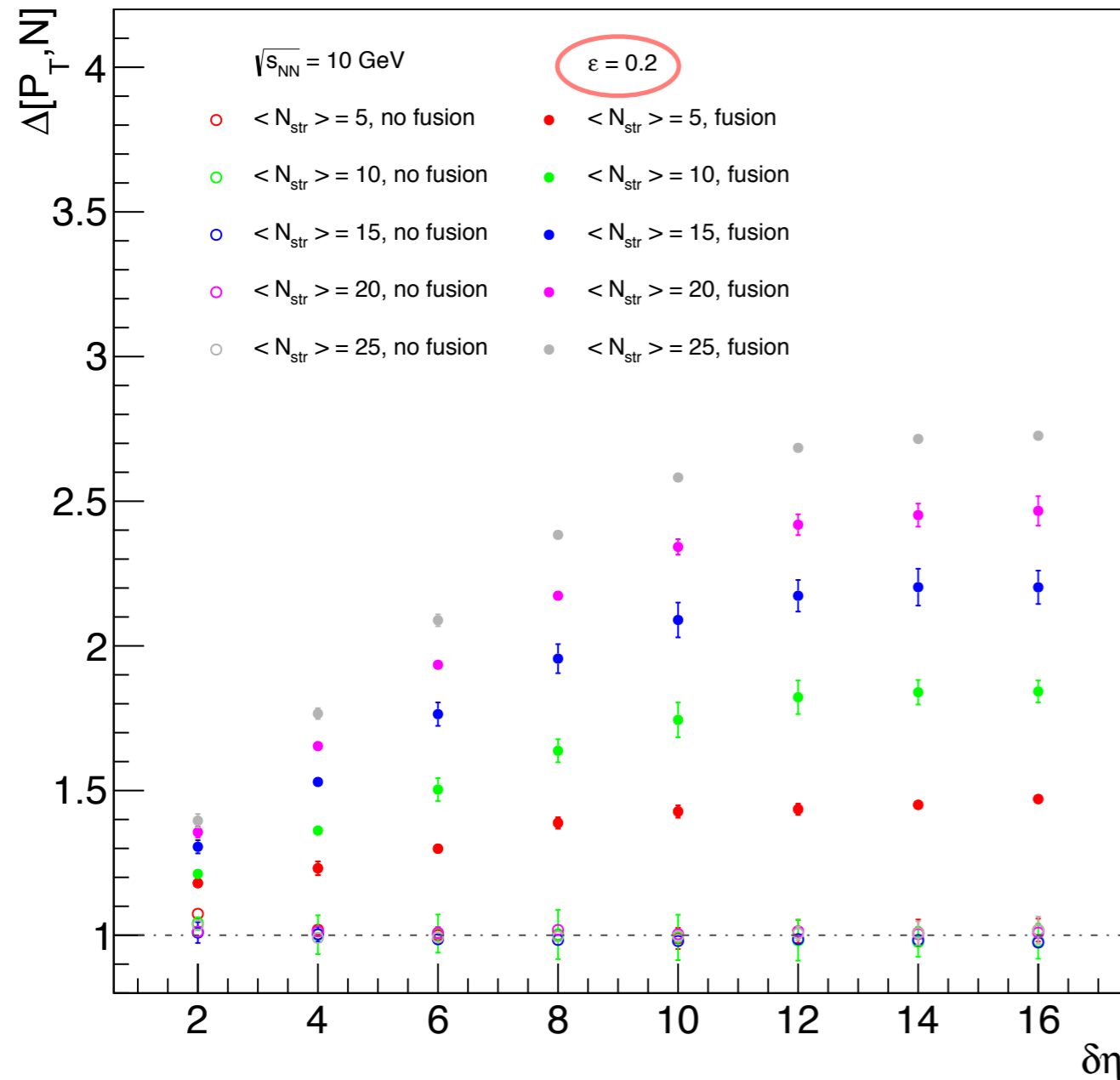
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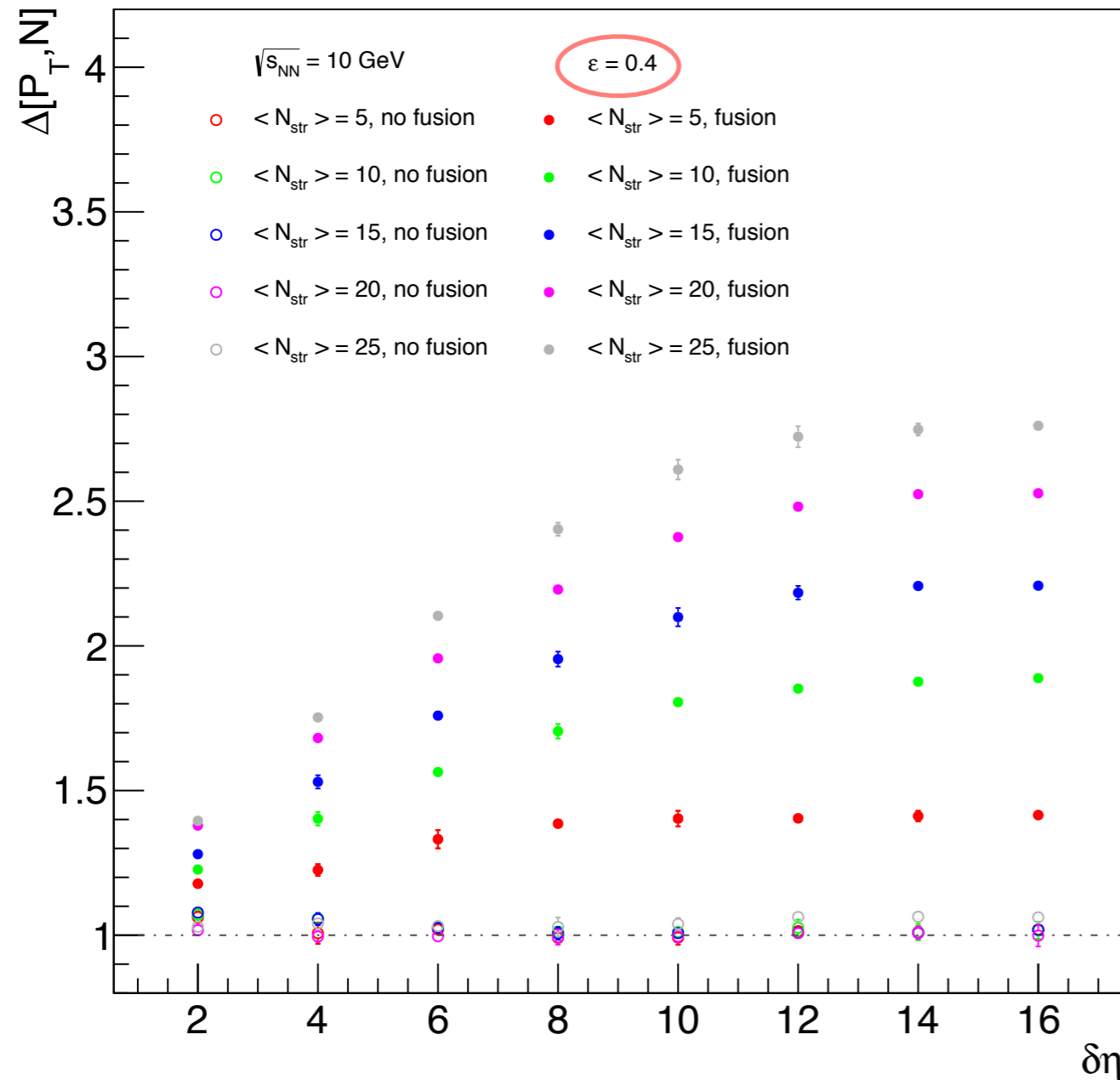
$\Delta[PT,N]$ as a function of the rapidity window width $\delta\eta$



- for no-fusion $\Delta[PT,N]$ fluctuates around 1
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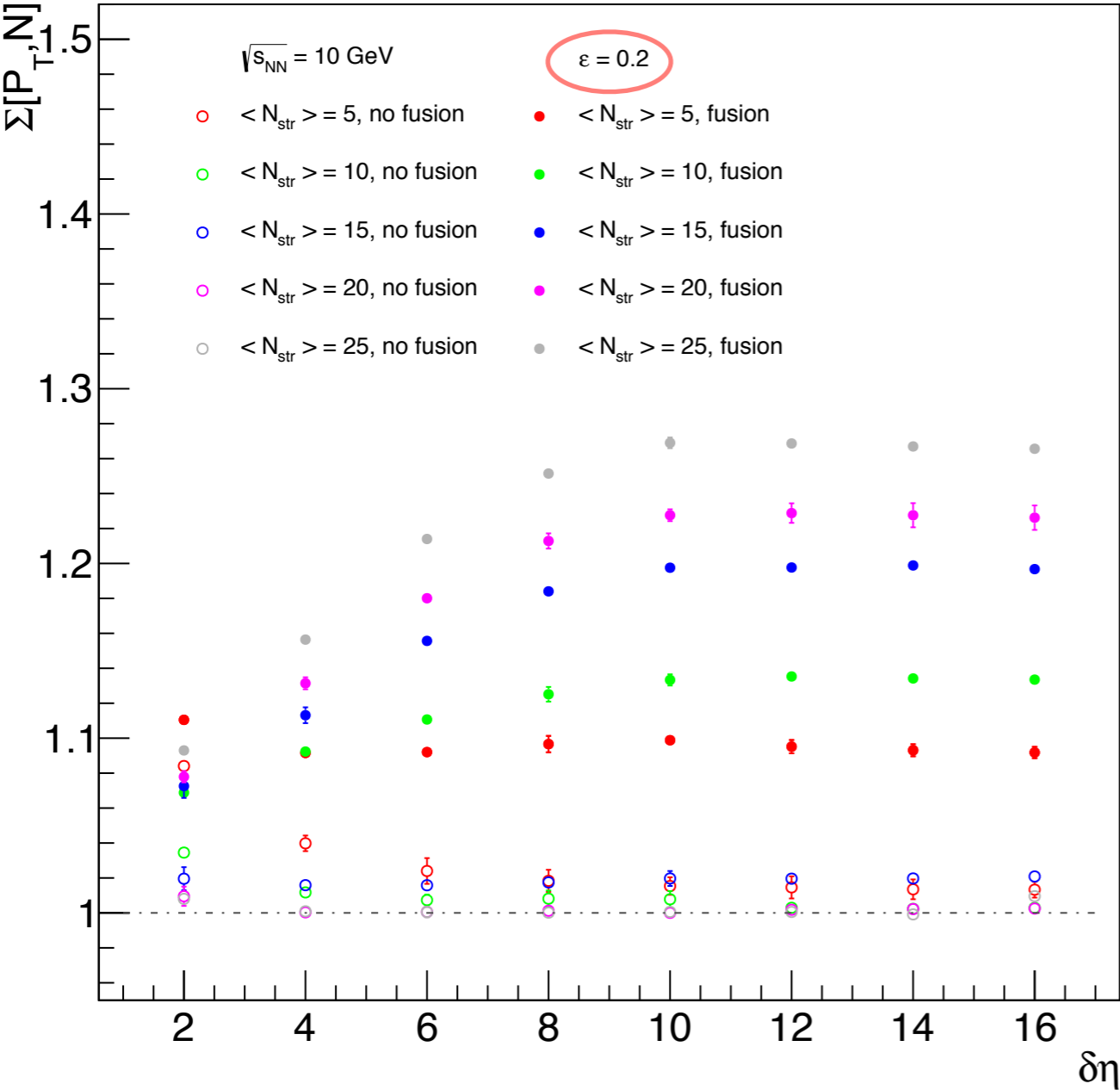
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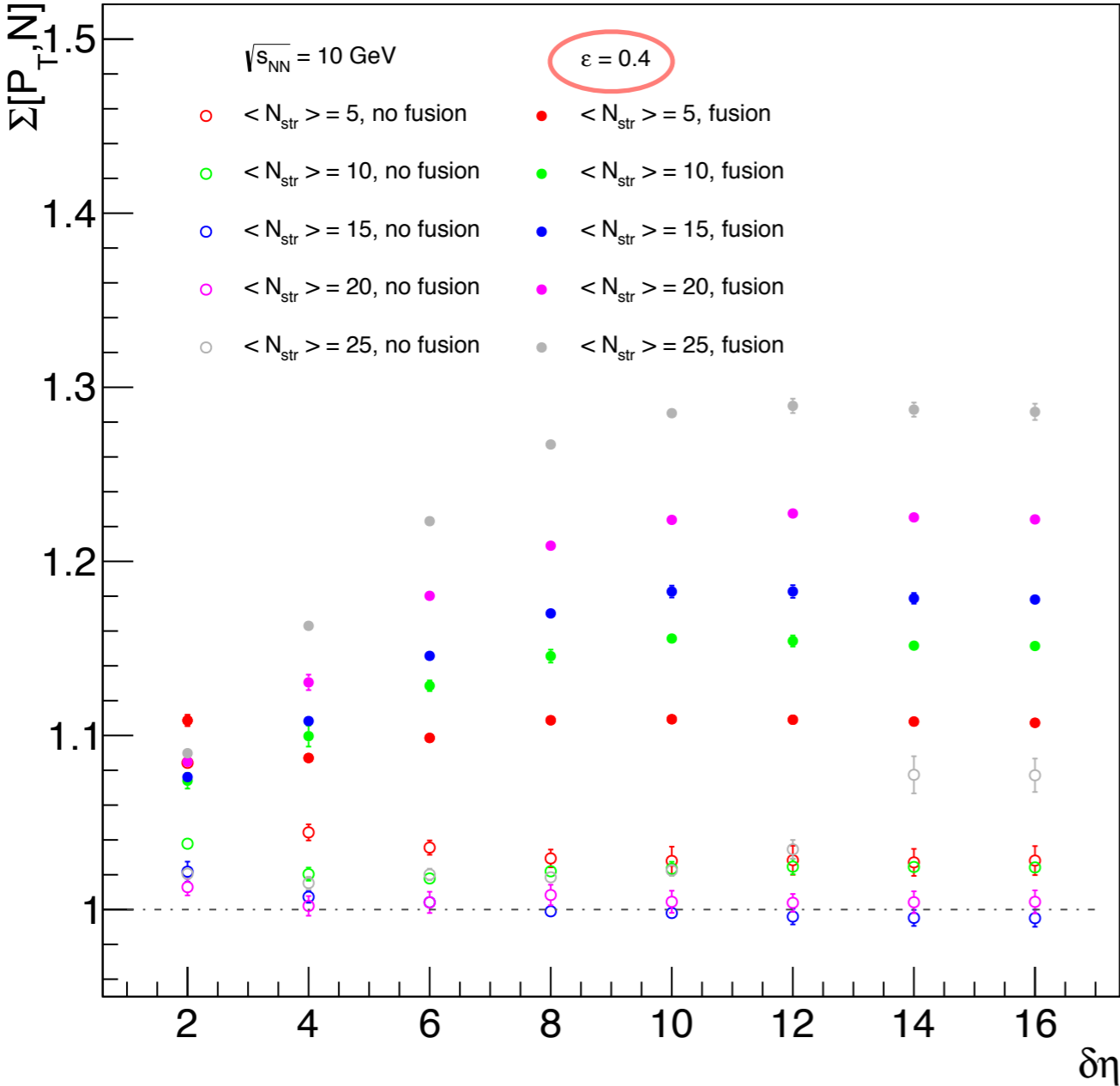
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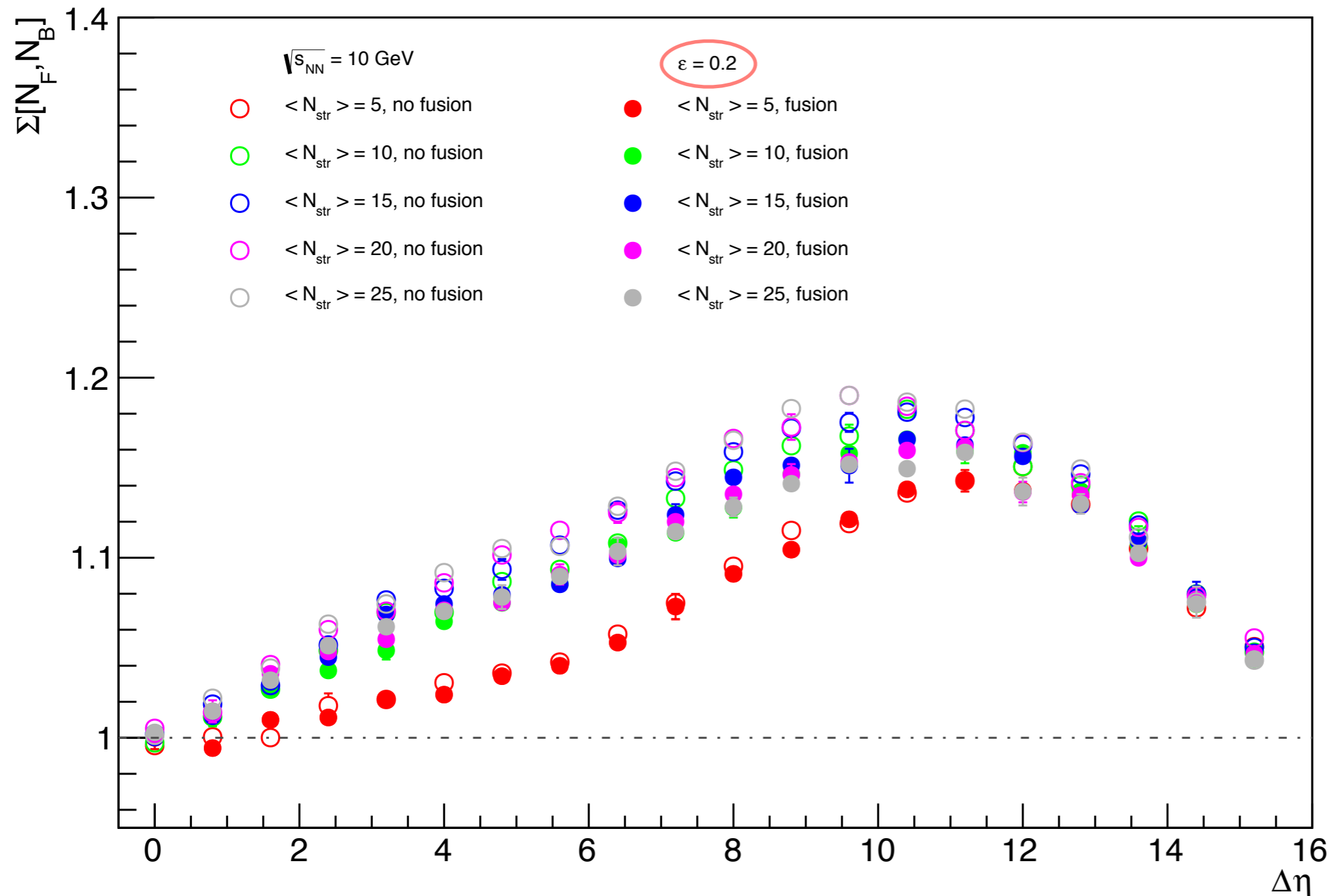
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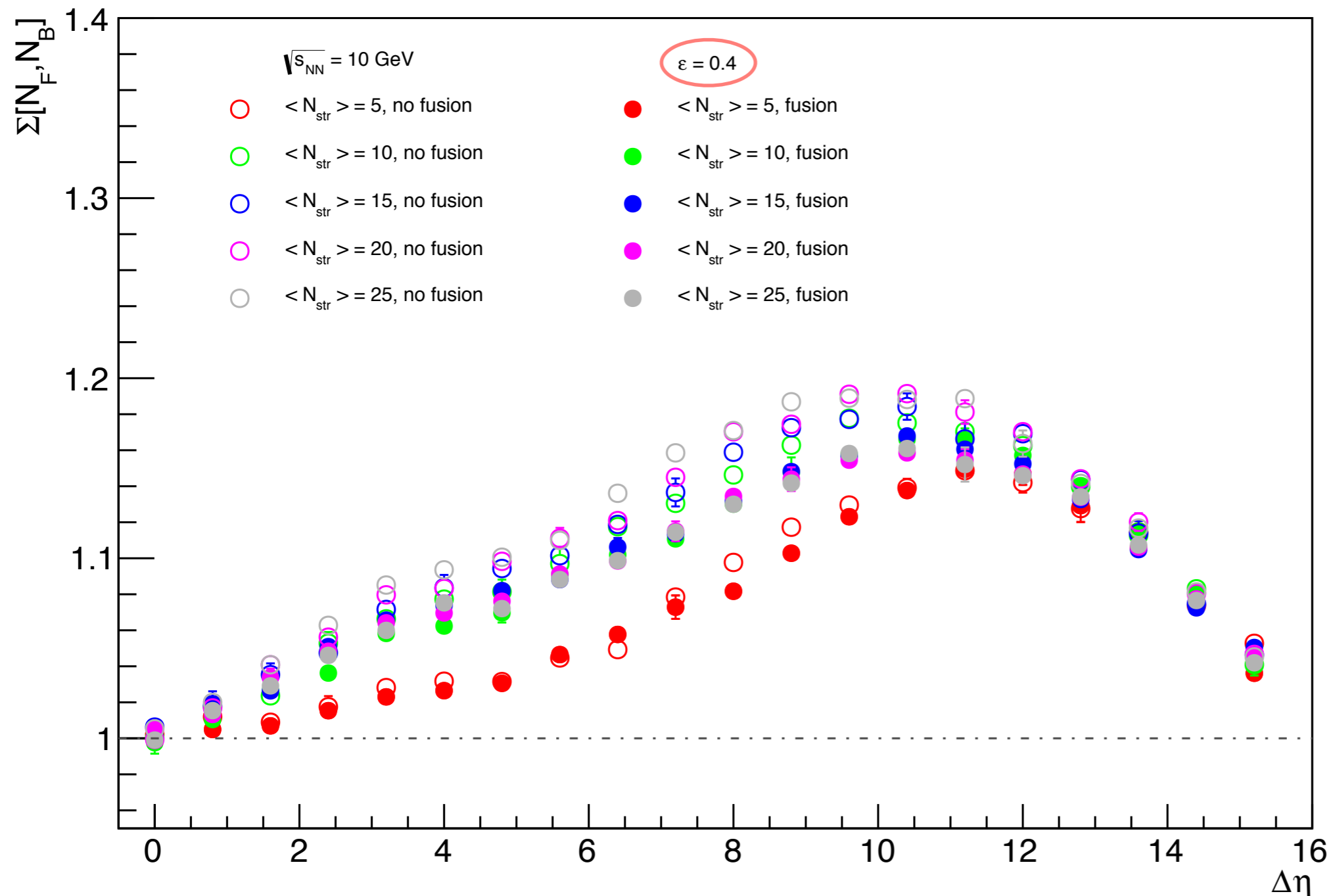
$\Sigma[N_F, N_B]$ as a function of the gap $\Delta\eta$ between two rapidity windows



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