



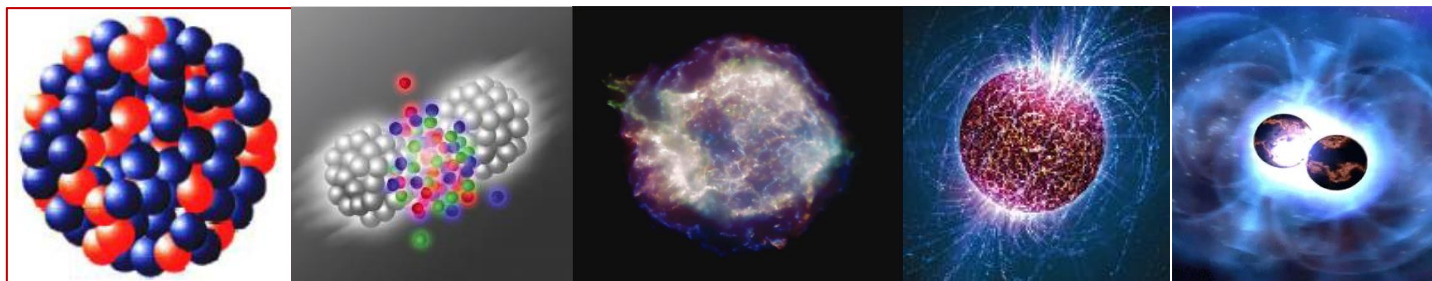
多信使时代核物质的状态方程

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□ Where does the dense nuclear matter exist?



Nucleus

HIC

SN

NStar

NS Merger

□ Can we build a unified Energy Density Functional to describe these dense systems?

复旦大学 核科学与技术系/现代物理研究所 (online)

2022年4月29日



Outline

- Nuclear matter EOS and the symmetry energy (E_{sym})
- Dense nuclear matter from *Nuclear experiments + Neutron star mass and radius + Tidal deformability from GW170817*
- Summary and outlook

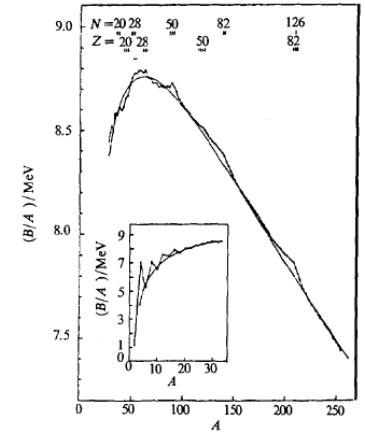


The Symmetry Energy of **Finite Nuclei**

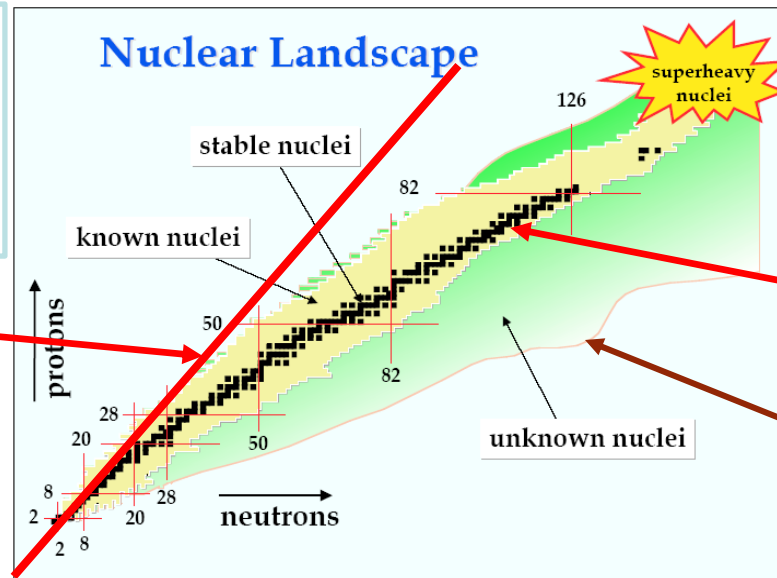
Liquid-drop model: **Bethe-Weizsäcker mass formula-1935**

$$a_v A - a_s A^{2/3} - a_4 \frac{(N - Z)^2}{A} - a_c \frac{Z(Z - 1)}{A^{1/3}} + a_p \frac{\Delta(N, Z)}{A^{1/2}}$$

Symmetry energy term
(对称能项)



对称能：
使中子数和质子数
趋于对称



库仑能：
使中子数和质子数
偏离对称

$N = Z$

$N > Z$

The $E_{sym}(2/3\rho_0)$ is very important for the global behaviors of N/P driplines!
R. Wang (王睿) /LWC, PRC92, 031303(R) (2015)



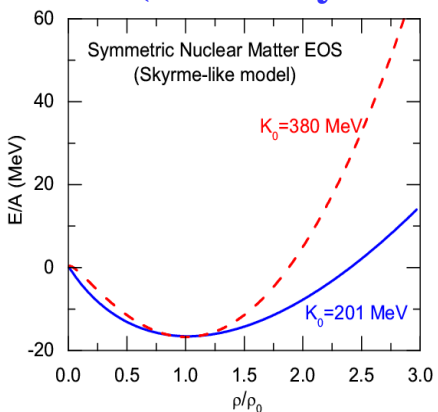
The Symmetry Energy of Nuclear Matter

EOS of Isospin Asymmetric Nuclear Matter (Parabolic law)

$$E(\rho, \delta) = E(\rho, 0) + E_{\text{sym}}(\rho)\delta^2 + O(\delta^4), \quad \delta = (\rho_n - \rho_p) / \rho$$

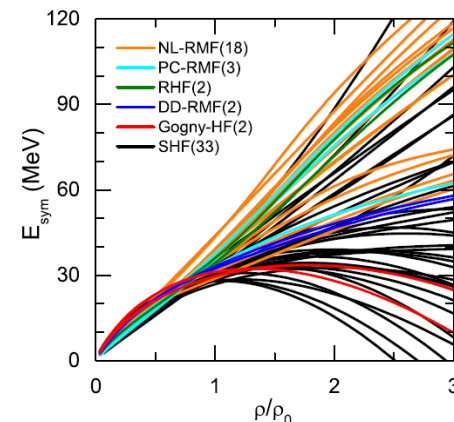
Symmetric Nuclear Matter
(relatively well-determined)

Isospin asymmetry
Symmetry energy term (largely uncertain)



Nuclear Matter Symmetry Energy

$$E_{\text{sym}}(\rho) \equiv \frac{1}{2} \frac{\partial^2 E(\rho, \delta)}{\partial \delta^2}$$



Saturation: $\rho_0 \approx 0.16 \text{ fm}^{-3} \approx 2.7 \times 10^{14} \text{ g/cm}^3$

LWC, EPJ Web of Conf. 88, 00017 (2015)

对称能是回答以下两个重大科学问题的关键量

- 中子星和致密核物质的性质是什么?
- 宇宙中的元素是怎样产生的?

The Frontiers of Nuclear Science
A LONG RANGE PLAN

《美国核科学前沿长期规划》

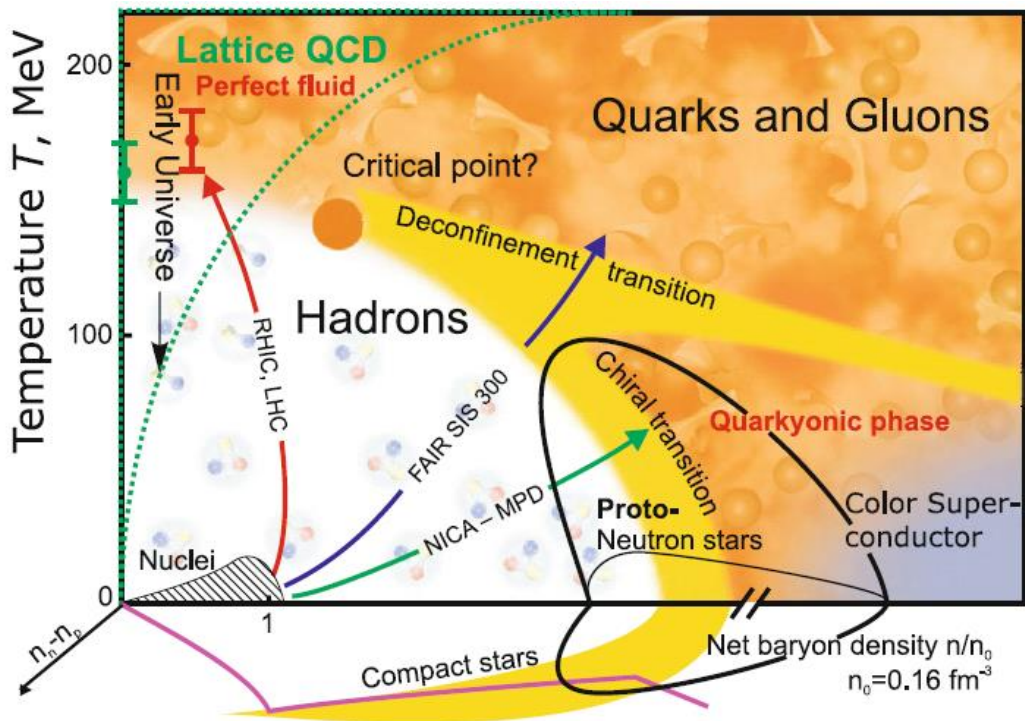
The Nuclear Science Advisory Committee

- ◆ 确定对称能已成为一些核物理大科学装置的重要物理目标, 比如: 中国兰州CSR/惠州HIAF、日本RIBF/RIKEN、美国FRIB/MSU和德国FAIR/GSI



QCD Phase Diagram

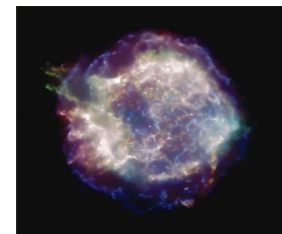
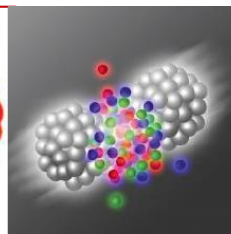
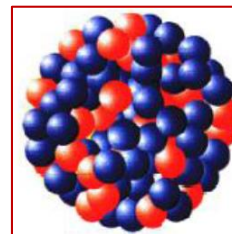
V.E. Fortov, Extreme States of Matter – on Earth and in the Cosmos, Springer-Verlag, 2011



Holy Grail
of
Nuclear
Physics



Probing QCD phase diagram in
Nuclei and Heavy Ion Collisions
in terrestrial labs
and in **NS Merger, SN, and Nstar**
in heaven?



Nuclei

HIC

SN

NStar

NS Merger

- **Small baryon chemical potential:**
Smooth Crossover Transition
- **Large baryon chemical potential:**
First-order Phase Transition
- **QCD Critical Endpoint (CEP):**
where the first-order phase transition ends

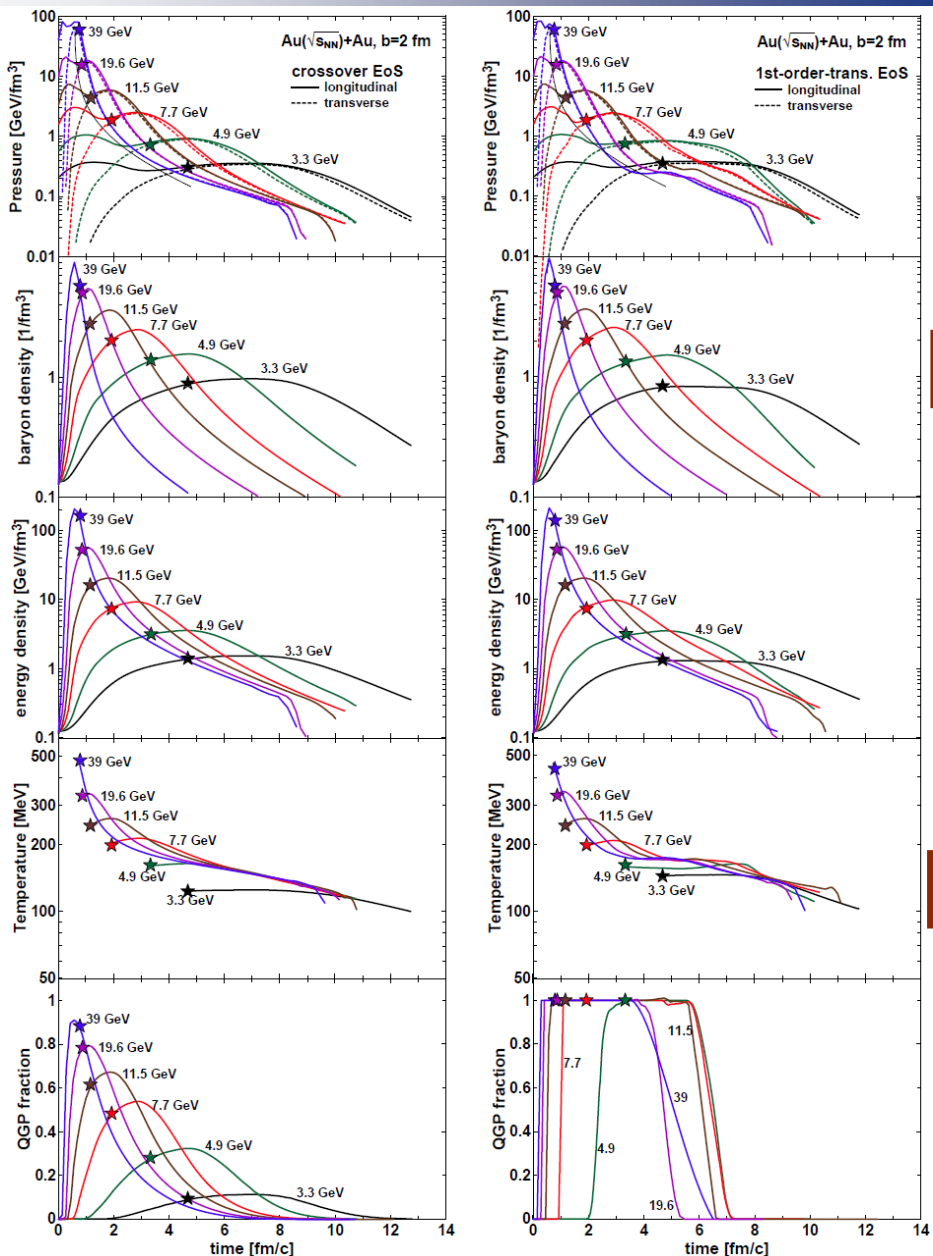
NS Mergers/NStar/CCSNe:

Ideal and unique site to probe QCD phase diagram at low T and high densities (and large isospin)!
Quark Matter Symmetry Energy?

M. Di Toro et al., NPA775 (2006); P.C. Chu (初鹏程)/LWC, ApJ780 (2014); LWC, 《原子核物理评论》 34, 20 (2017) [arXiv:1708.04433]

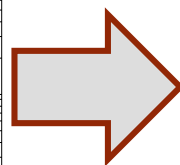


n_B and T in relativistic heavy-ion collisions

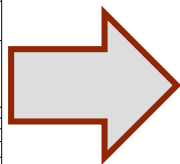


Head on Au+Au collisions (3-fluid dynamics)

Yu. B. Ivanov and A. A. Soldatov,
PRC 101, 024915 (2020) [arXiv:1911.10872]



The maximum baryon density could reach about $\sim 4\rho_0$ ($30\rho_0$) in head on AuAu collisions at $\sqrt{s_{NN}} = 3.3$ (39) GeV



The temperature could reach about ~ 110 ($500-110$) MeV in head on AuAu collisions at $\sqrt{s_{NN}} = 3.3$ (39) GeV

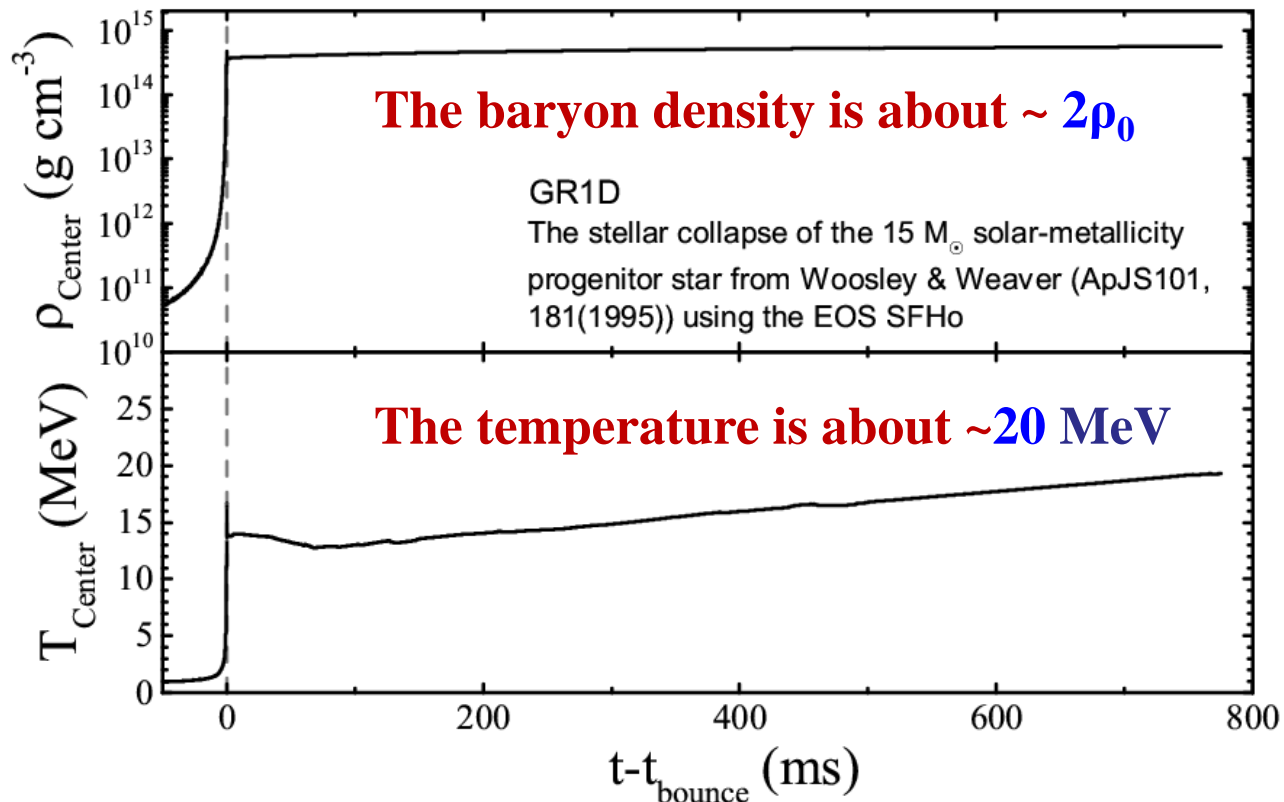
See also:

I.C. Arsene et al., PRC75, 034902(2007)



A typical Core-Collapse Supernova Explosion

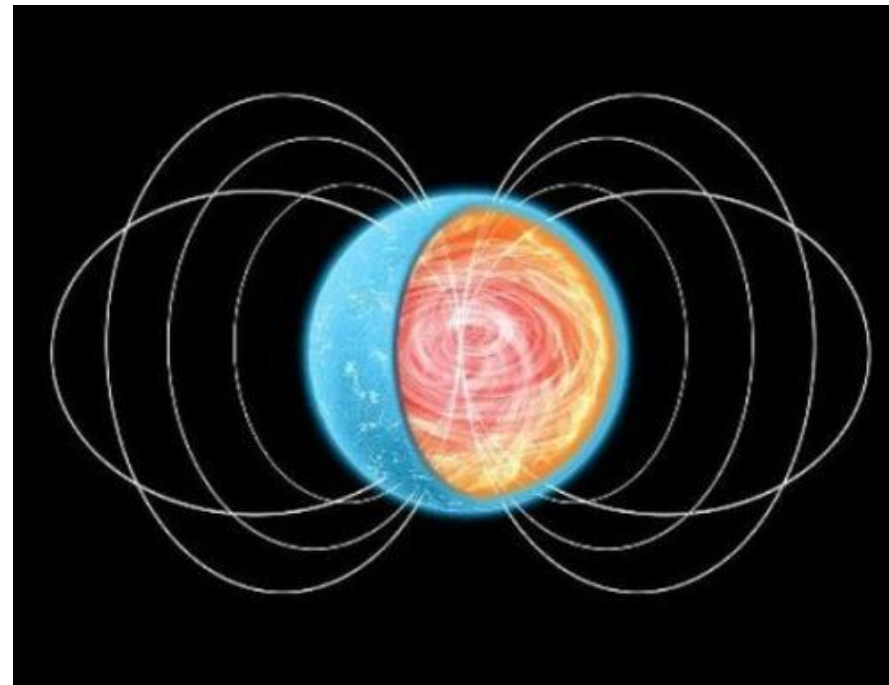
Xu-Run Huang (黄旭润), Shuai Zhao, and LWC, ApJL923, L26 (2021)



Density at the center: $\sim 2\rho_0$

Temperature at the center: ~ 20 MeV (~ 50 MeV around the shock wave)

Neutron stars

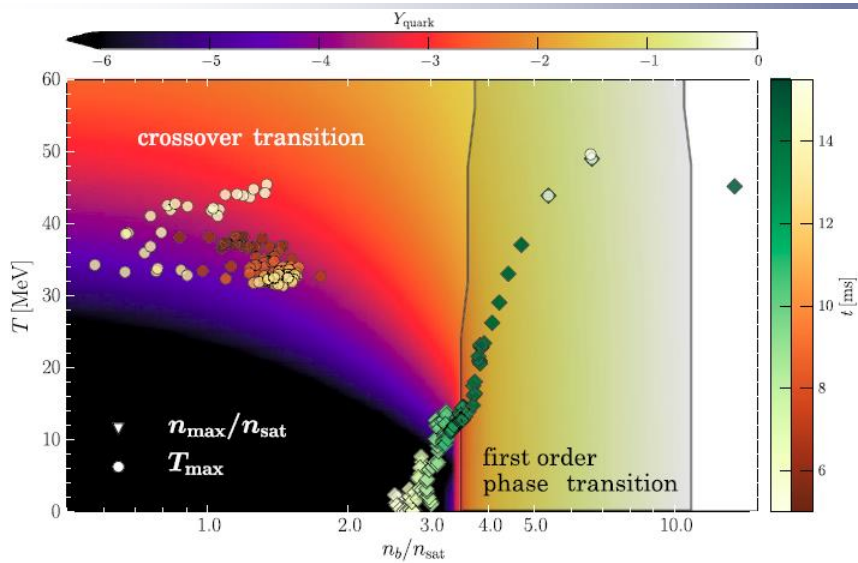


Center density: $\sim 6\rho_0$; Average density: $\sim 2.5\rho_0$

Temperature ~ 0



n_B and T in binary neutron star merger

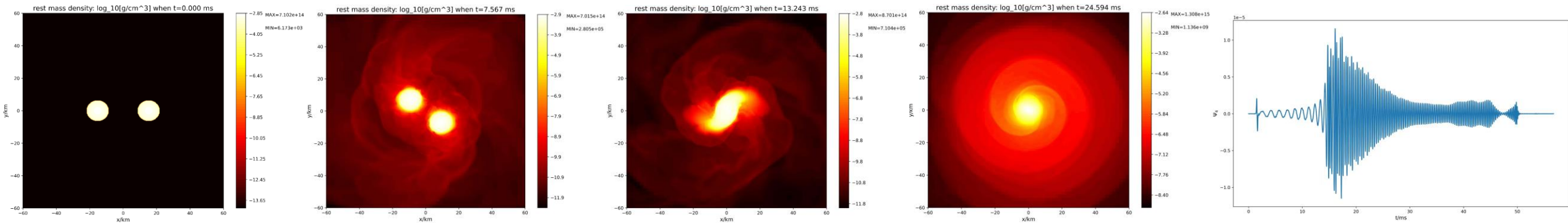


The overall evolution during and after the merger of the low-mass binary with total mass $2.8 M_{\odot}$

The baryon density in binary neutron star merger could reach more than $\sim 15\rho_0$ and the temperature could reach $\sim 0-50$ MeV but with large isospin density!

Elias R. Most et al., PRL122, 061101 (2019)

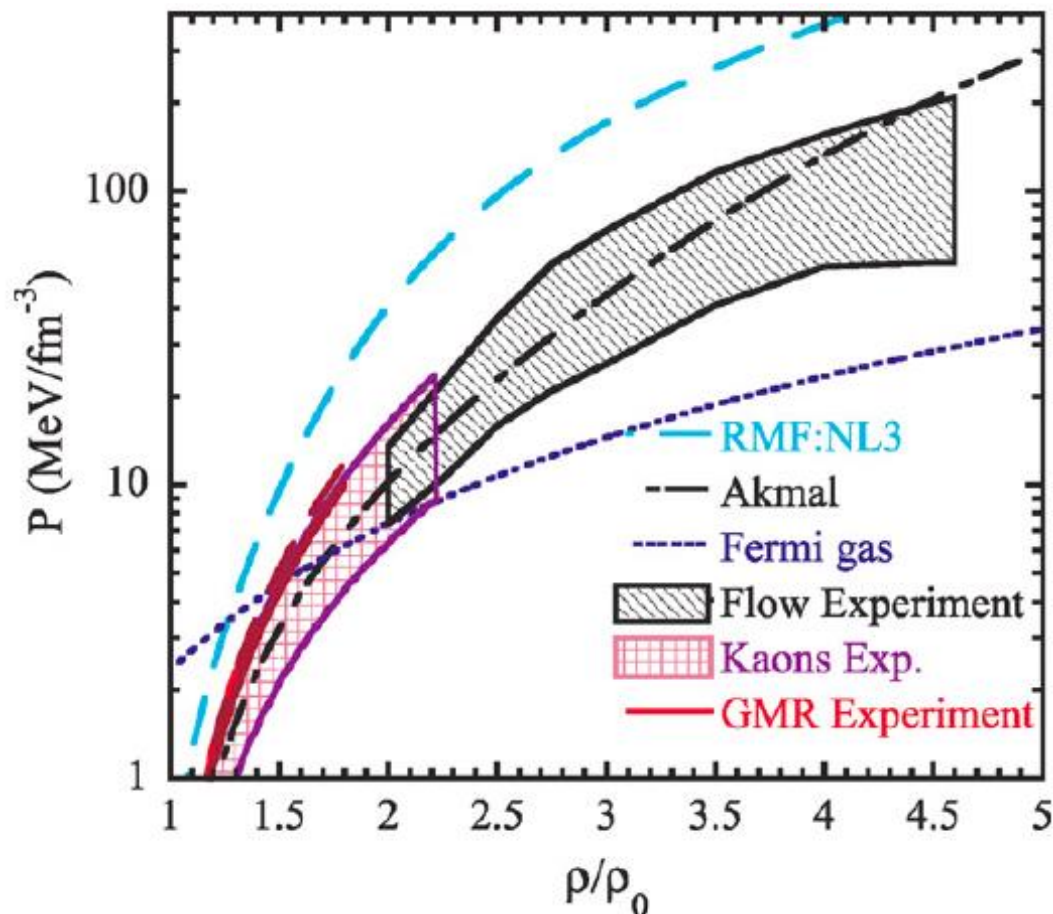
Matter density evolution of BNS mergers with equal mass $1.364M_{\odot}$ (LS220 EOS, Ji-Min Bai(白济民) et al., preliminary)



The high density Esym of nuclear/quark matter is particularly important for NS mergers/NStar/CCSNe!



W.G. Lynch et al., PPNP62, 427 (2009)



Curent status of the EOS of SNM:

- Around ρ_0 : $K_0=240 \pm 20$ MeV from **GMR** [U. Garg and G. Colo, PPNP101, 55 (2018)]
- $1\rho_0 < \rho < 3\rho_0$ from **K^+ production in HIC's**. The range of densities for the **kaon constraint** represents an **educated guess** (有根据的推测☹) based on available calculations [C. Fuchs, PPNP 56, 1 (2006)]
- $2\rho_0 < \rho < 5\rho_0$ using **flow data from BEVALAC, SIS/GSI and AGS** [P. Danielewicz, R. Lacey and W.G. Lynch, Science 298, 1592 (2002)]

The EOS of Symmetric NM has been relatively well constrained!
(~Soft EOS of SNM at high densities)



Esym: Many-Body Approaches

The nuclear matter EOS cannot be measured experimentally, its determination thus depends on theoretical approaches

- **Microscopic Many-Body Approaches**

- Non-relativistic Brueckner-Bethe-Goldstone (BBG) Theory
- Relativistic Dirac-Brueckner-Hartree-Fock (DBHF) approach
- Self-Consistent Green's Function (SCGF) Theory
- Variational Many-Body (VMB) approach
- Green's Function Monte Carlo Calculation
- $V_{\text{low}k}$ + Renormalization Group
- Nuclear Lattice Approach

- **Effective Field Theory**

- Density Functional Theory (DFT)
- Chiral Perturbation Theory (ChPT)
- QCD-based theory

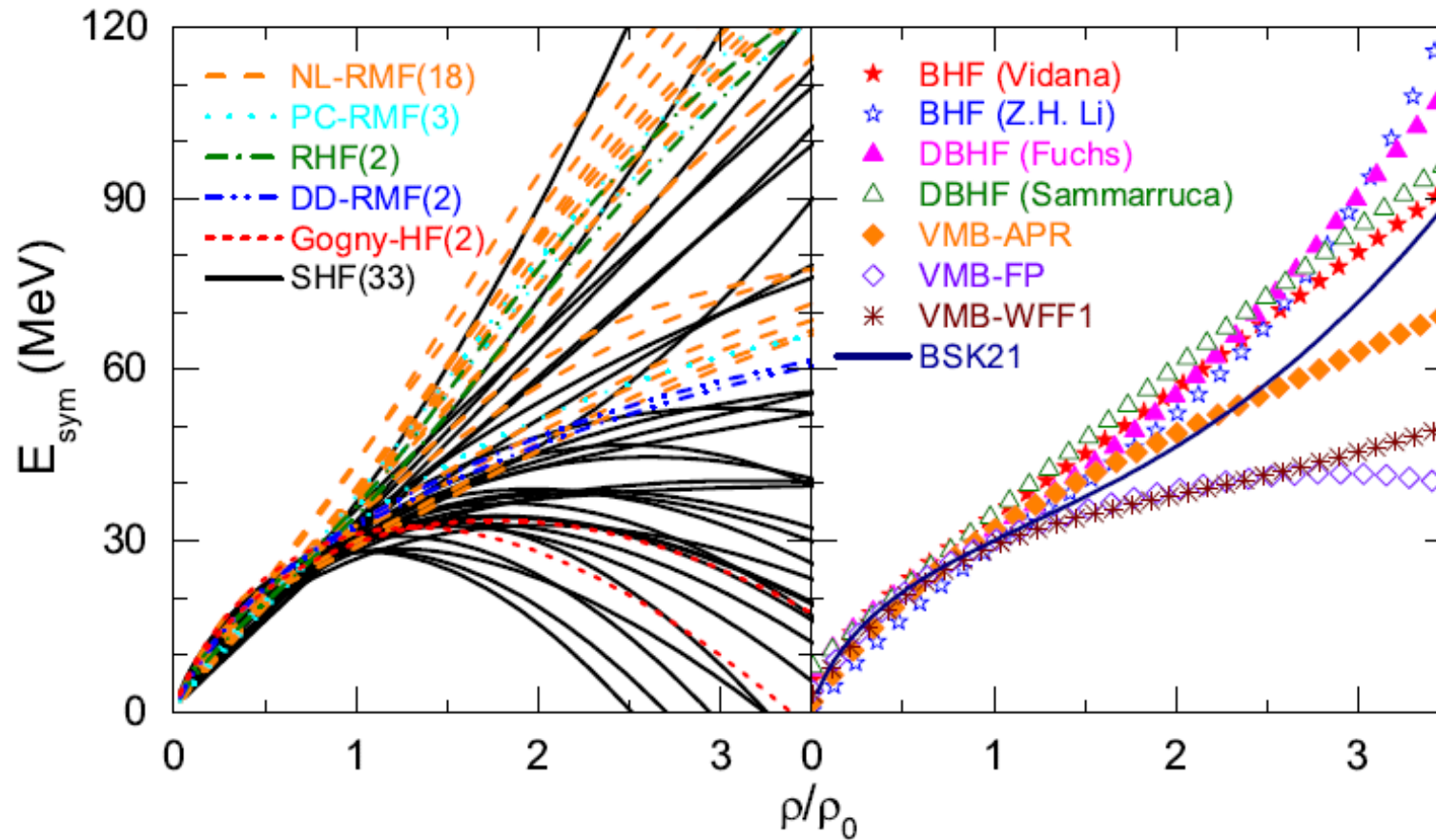
- **Phenomenological Approaches**

- Relativistic mean-field (RMF) theory
- Quark Meson Coupling (QMC) Model
- Relativistic Hartree-Fock (RHF)
- Non-relativistic Hartree-Fock (Skyrme-Hartree-Fock)
- Thomas-Fermi (TF) approximations



Esym: Many-Body Approaches

L.W. Chen, Nucl. Phys. Rev. (原子核物理评论) 34, 20 (2017)
[arXiv:1708.04433]





Esym: Experimental Probes

Promising Probes of the $E_{\text{sym}}(\rho)$

(an incomplete list !)

At sub-saturation densities (亚饱和密度行为)

- Sizes of n-skins of unstable nuclei from total reaction cross sections
- **Proton-nucleus elastic scattering in inverse kinematics**
- Parity violating electron scattering studies of the n-skin in ^{208}Pb
- **n/p ratio of FAST, pre-equilibrium nucleons**
- **Isospin fractionation and isoscaling in nuclear multifragmentation**
- **Isospin diffusion/transport**
- Neutron-proton differential flow
- **Neutron-proton correlation functions at low relative momenta**
- $t/{}^3\text{He}$ ratio
- **Hard photon production**
- Pigmy/Giant resonances
- Nucleon optical potential

Towards high densities reachable at CSR/Lanzhou, FAIR/GSI, RIKEN, GANIL and, FRIB/MSU (高密度行为)

- **π^-/π^+ ratio, K^+/K^0 ratio?**
- Neutron-proton differential transverse flow
- **n/p ratio at mid-rapidity**
- Nucleon elliptical flow at high transverse momenta
- **n/p ratio of squeeze-out emission**

B.A. Li, L.W. Chen, C.M. Ko
Phys. Rep. 464, 113(2008)
[arXiv:0804.3580] (Citations: 1082+)

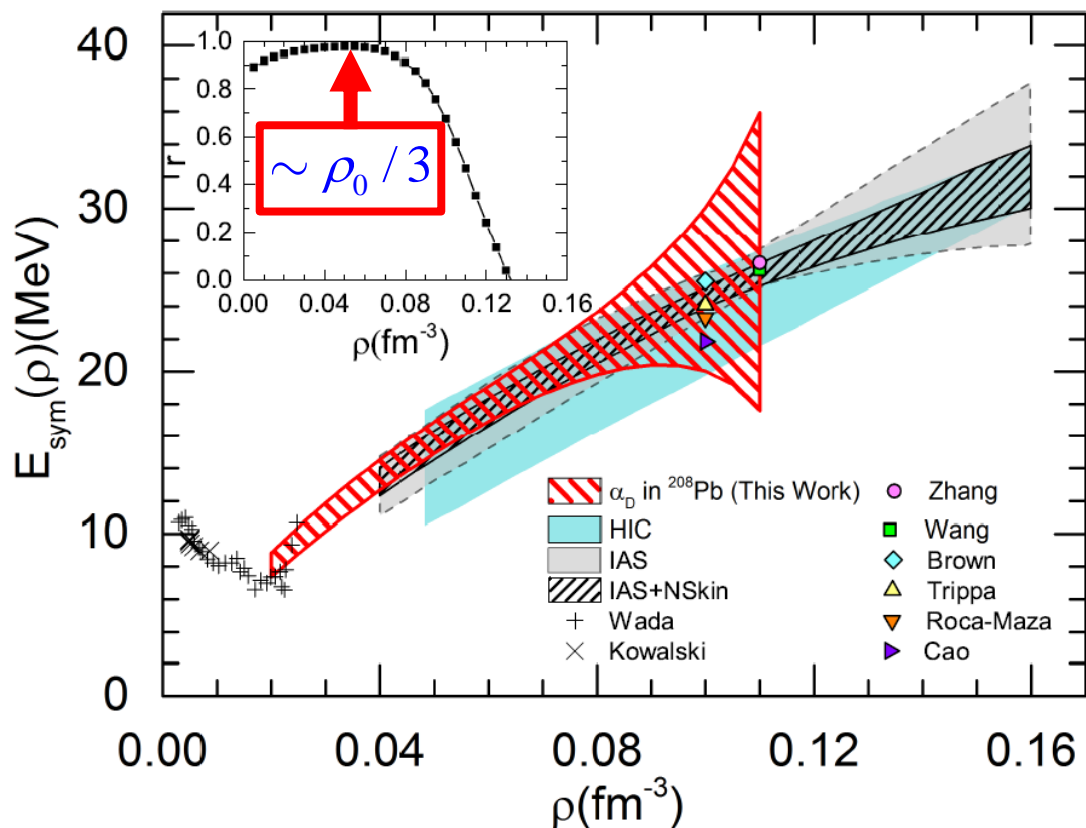


An incomplete list

- Cooling Storage Ring (**CSR**) Facility at HIRFL/Lanzhou in China (2008)
up to 500 MeV/A for ^{238}U (~1GeV/A for C)
<http://www.impcas.ac.cn/zhuye/en/htm/247.htm>
- Radioactive Ion Beam Factory (**RIBF**) at RIKEN in Japan (2007)
<http://www.riken.jp/engn/index.html>
- Texas A&M Facility for Rare Exotic Beams -T-REX (2013)
<http://cyclotron.tamu.edu>
- Facility for Antiproton and Ion Research (**FAIR**)/GSI in Germany (2022?)
up to 2 GeV/A for ^{132}Sn (**NUSTAR** - NUclear STructure, Astrophysics and Reactions)
http://www.gsi.de/fair/index_e.html
- **SPIRAL2**/GANIL in France (2013)
<http://pro.ganil-spiral2.eu/spiral2>
- Selective Production of Exotic Species (**SPES**)/INFN in Italy (2015)
<http://web.infn.it/spes>
- Facility for Rare Isotope Beams (**FRIB**)/MSU in USA (2022?)
up to 400(200) MeV/A for ^{132}Sn
<http://www.frib.msu.edu/>
- The Korean Rare Isotope Accelerator (KoRIA-**RAON**(**RISP Accelerator Complex**)) (2021?)
up to 250 MeV/A for ^{132}Sn , up to 109 pps
.....(NICA, J-PARC-HI)



Z. Zhang (张振) and LWC, PRC92, 031301(R) (2015)



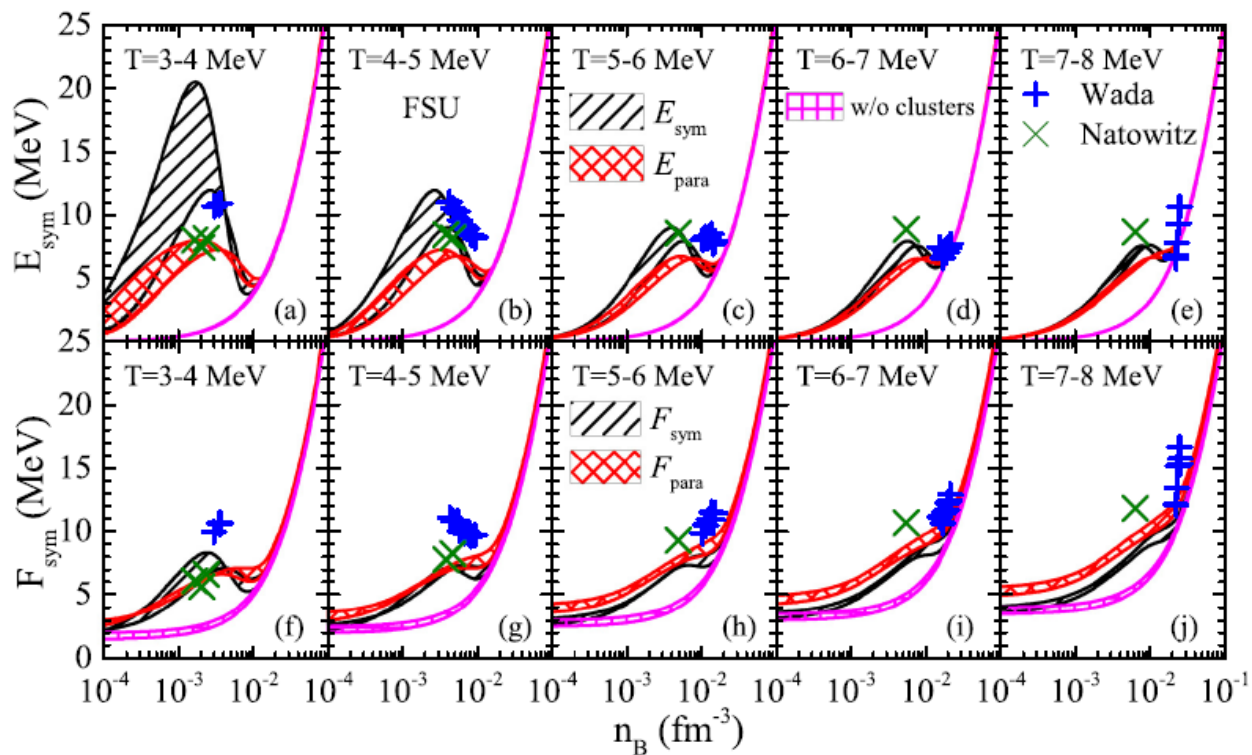
- $\Delta E(A \sim 208) \propto E_{\text{sym}}(\rho_{A=208}), \rho_{A=208} \approx 2/3\rho_0$
- $1/\alpha_D(A=208) \propto E_{\text{sym}}(\rho_{A=45}), \rho_{A=45} \approx 1/3\rho_0$

- **HIC: Sn+Sn**
M.B. Tsang *et al.*, Phys. Rev. Lett. **102**, 122701(2009)
- **IAS and IAS+NSkin**
P. Danielewicz and J. Lee, Nucl. Phys. **A922**, 1 (2014)
- **Zhang: Isotope binding energy difference**
Z. Zhang and L.W. Chen, Phys. Lett. **B726**, 234 (2013)
- **Wang: Fermi energy difference**
N. Wang *et al.*, Phys. Rev. C **87**, 034327 (2013)
- **Brown: Doubly magic nuclei**
B.A. Brown, Phys. Rev. Lett. **11**, 232502 (2013)
- **Trippa: Giant dipole resonance**
L. Trippa *et al.*, Phys. Rev. C **77**
- **Roca-Maza: Giant quadrupole resonance**
X. Roca-Maza *et al.*, Phys. Rev. C **87**, 034301 (2013)
- **Cao: Pygmy dipole resonance**
L.G. Cao and Z.Y. Ma, Chin. Phys. Lett. **25**, 1625 (2008)

Wada and Kowalski: experimental results of the symmetry energies at densities below $0.2\rho_0$ and temperatures in the range 3 ~11 MeV from the analysis of cluster formation in heavy ion collisions.

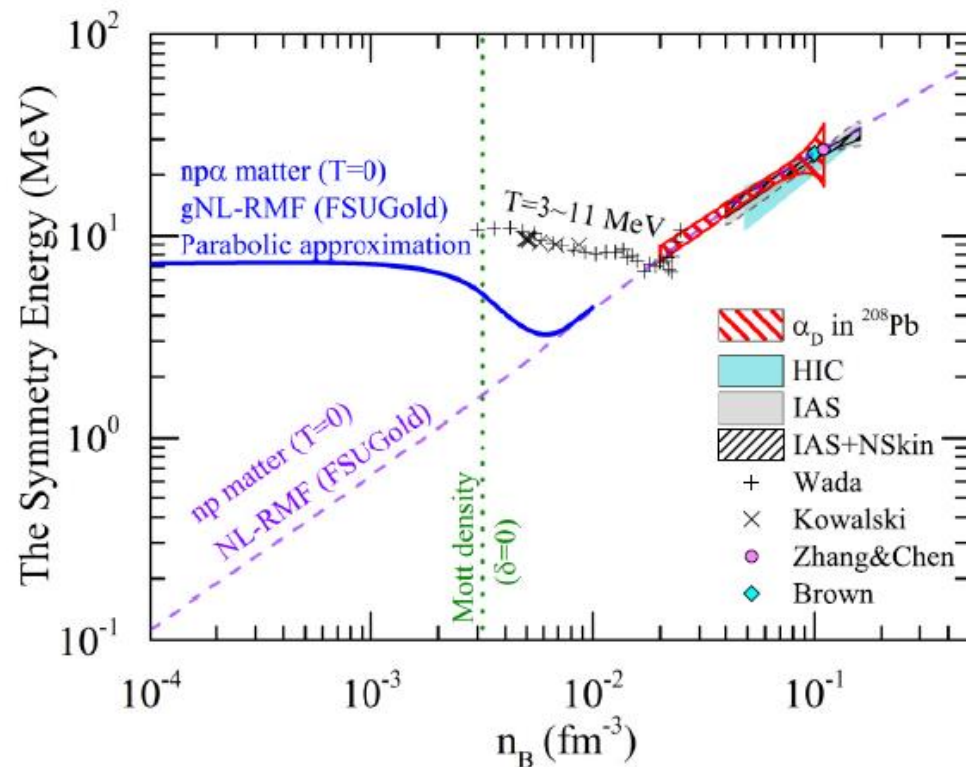
Wada *et al.*, Phys. Rev. C **85**, (2012) 064618; Kowalski *et al.*, Phys. Rev. C **75**, (2007) 014601. Natowitz *et al.*, Phys. Rev. Lett. **104**, (2010) 202501.

Clustering effects on E_{sym} within NL-RMF for n, p, t, h, α matter



Zhao-Wen Zhang (张肇文) and LWC, PRC95, 064330 (2017)

Alpha BEC effects on E_{sym} within NL-RMF for cold $np\alpha$ matter

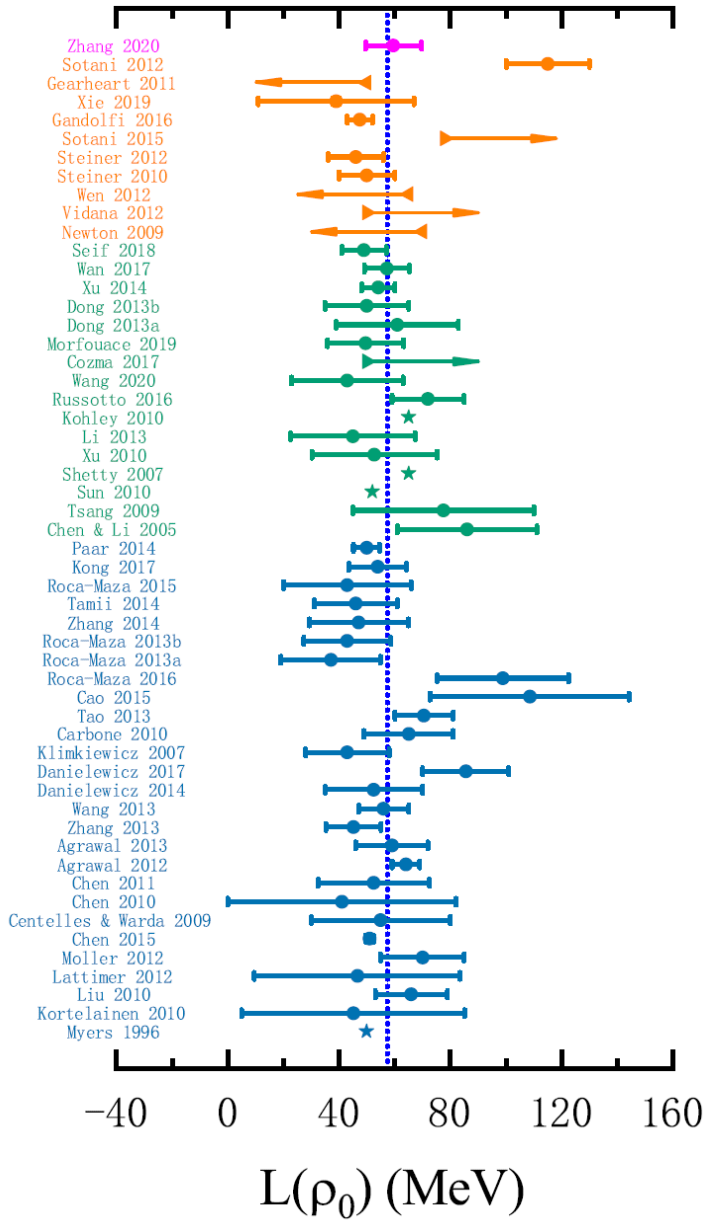
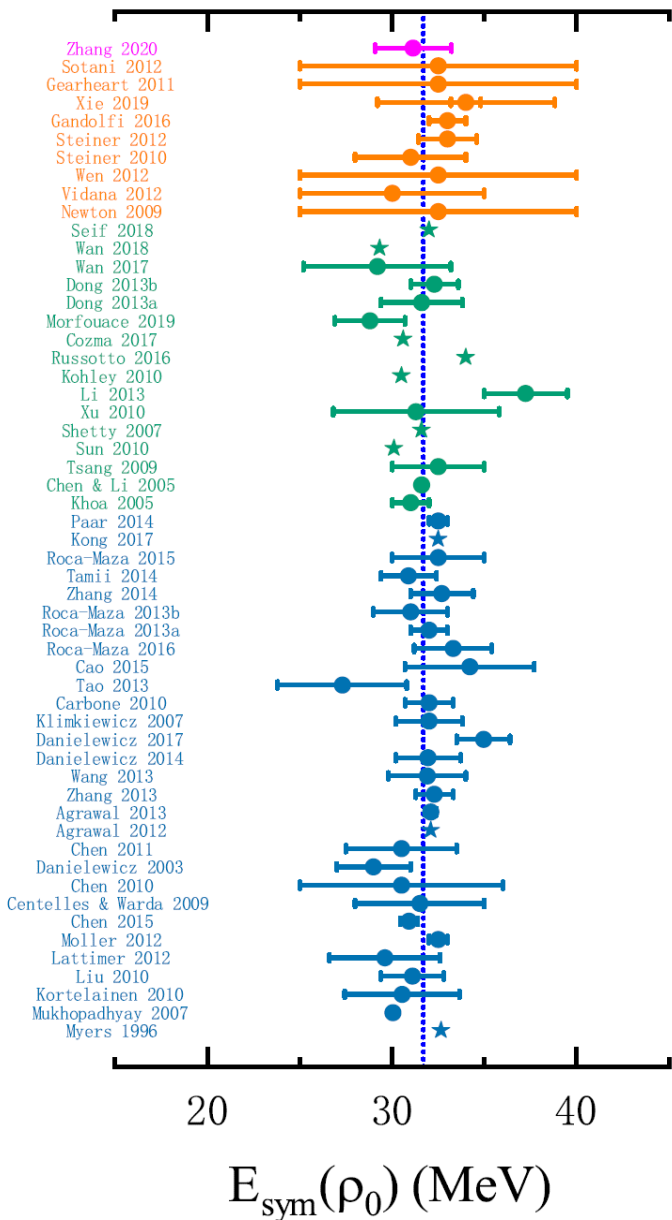


Zhao-Wen Zhang (张肇文) and LWC, PRC100, 054304 (2019)

See also: S. Typel, G. Röpke, T. Klähn, D. Blaschke, and H. H. Wolter, Phys. Rev. C 81, 015803 (2010).



E_{sym} : Around saturation density



LWC et al., Invited Review/PPNP
58 analyses of terrestrial nuclear experiments and astrophysical observations

$$E_{\text{sym}}(\rho_0) = 31.7 \pm 3.1$$

$$L = 57.5 \pm 24.5 \text{ MeV}$$

Similar conclusion has been obtained in:
B. A. Li and X. Han, Phys. Lett. B727, 276 (2013);
M. Oertel, M. Hempel, T. Klahn, and S. Typel, Rev. Mod. Phys. 89, 015007 (2017).



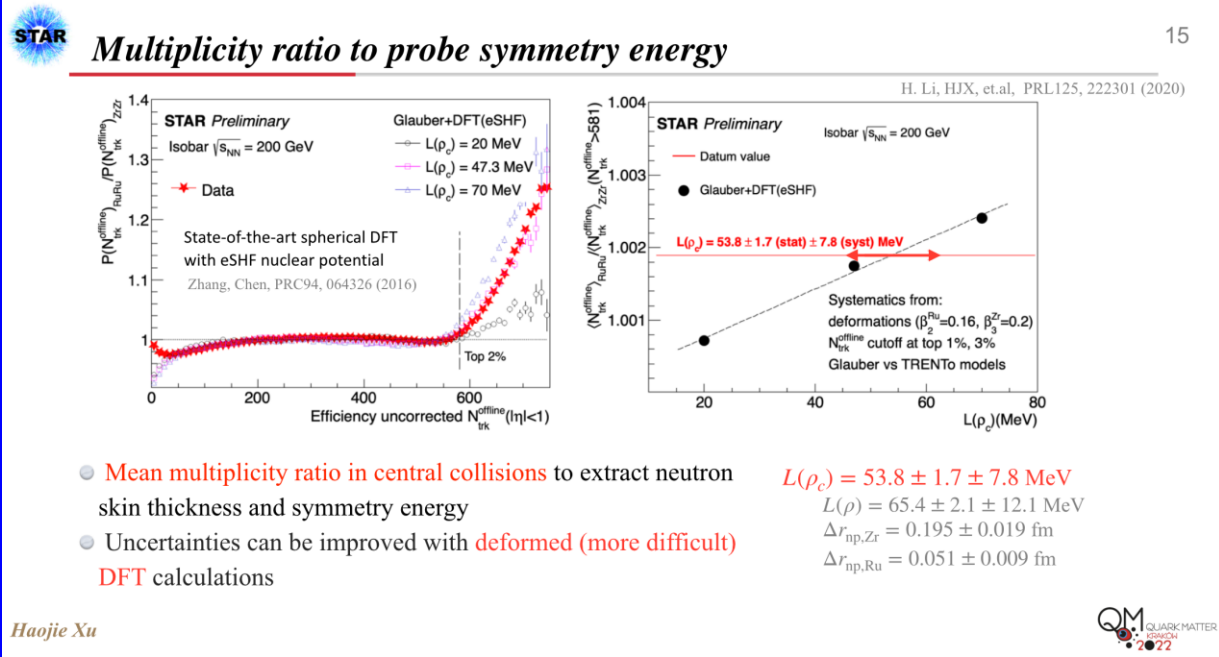
Assuming all the constraints are equally reliable !!!

Very recent PREX-II data suggest stiff E_{sym} around saturation density!

PHYSICAL REVIEW LETTERS 125, 222301 (2020)

Probing the Neutron Skin with Ultrarelativistic Isobaric Collisions

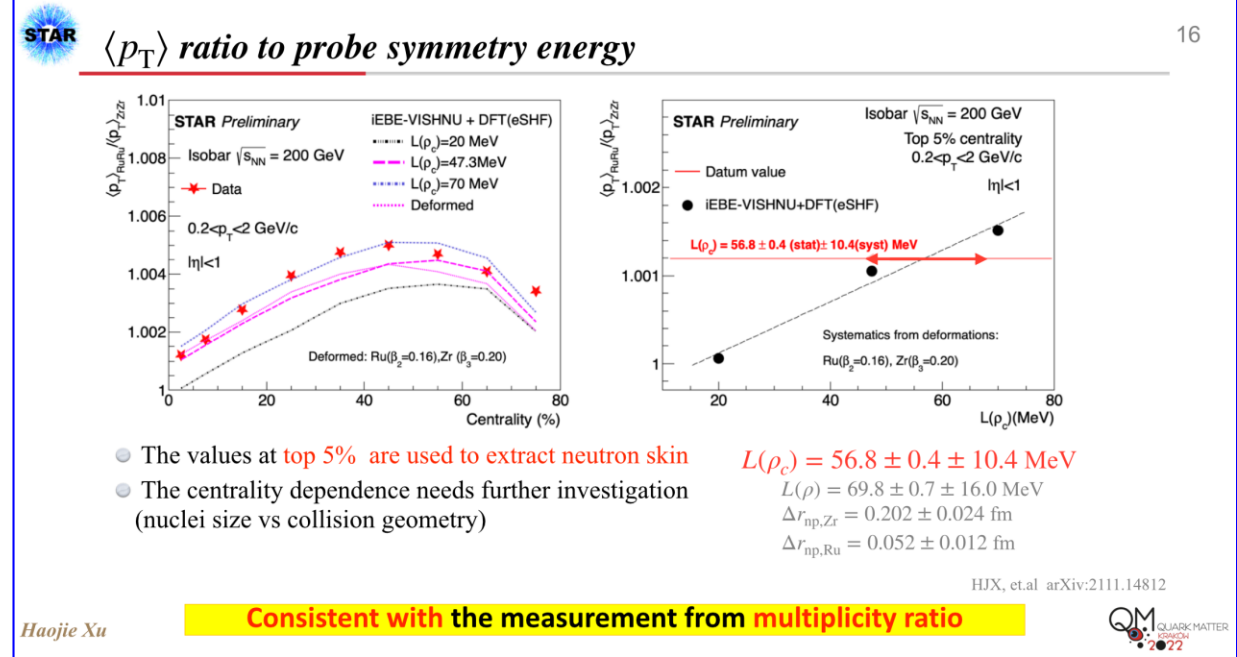
Hanlin Li,¹ Hao-jie Xu,^{2,*} Ying Zhou,³ Xiaobao Wang,² Jie Zhao,⁴ Lie-Wen Chen,^{3,†} and Fuqiang Wang^{2,4,‡}



H.-j. Xu et al., arXiv:2111.14812

Probing nuclear structure with mean transverse momentum in relativistic isobar collisions

Hao-jie Xu,¹ Wenbin Zhao,² Hanlin Li,³ Ying Zhou,⁴ Lie-Wen Chen,⁴ and Fuqiang Wang^{1,5}



**Multiplicity ratio: $L_c = 53.8 \pm 1.7 \pm 7.8 \text{ MeV}$
 $L = 65.4 \pm 2.1 \pm 12.1 \text{ MeV}$**

**$\langle p_T \rangle$ ratio: $L_c = 56.8 \pm 0.4 \pm 10.4 \text{ MeV}$
 $L = 69.8 \pm 0.7 \pm 16.0 \text{ MeV}$**

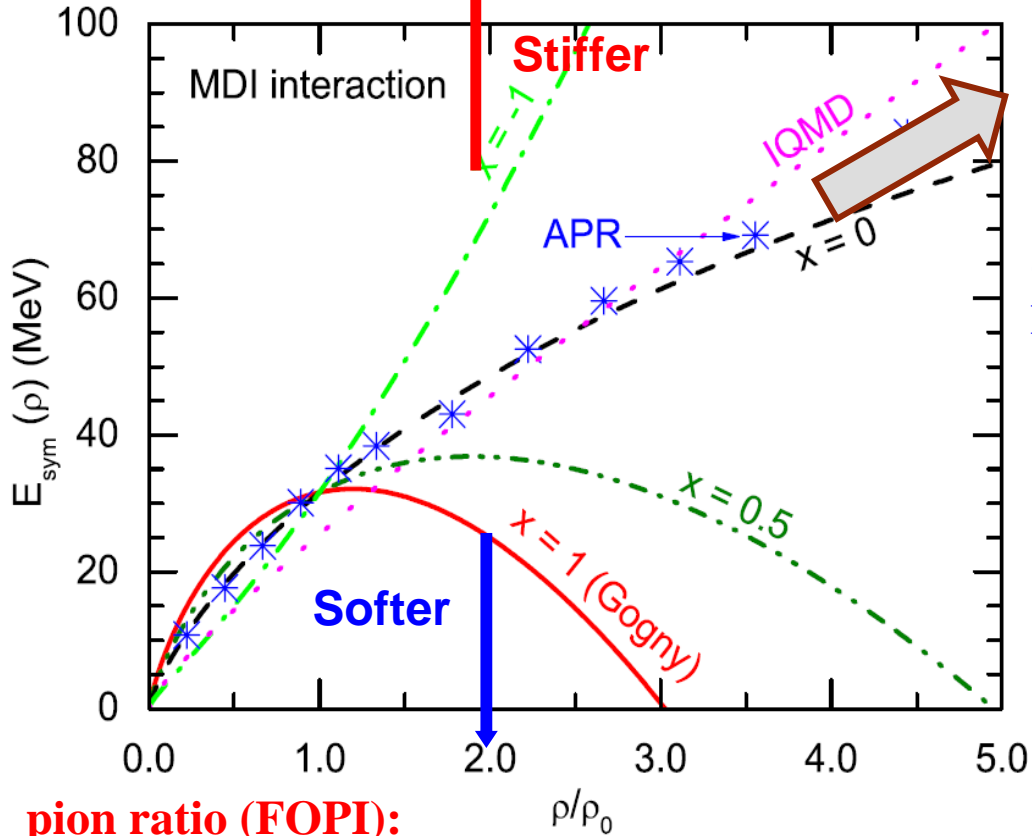
RHIC-STAR isobaric collisions data favor stiff E_{sym} around saturation density!



E_{sym} : Supra-saturation density

A Soft or Stiff E_{sym} at supra-saturation densities ???

pion ratio (FOPI): ImIQMD, Feng/Jin, PLB683, 140(2010)



pion ratio (FOPI):
IBUU04, Xiao/Li/Chen/Yong/Zhang, PRL102,062502(2009)
ImIBLÉ, Xie/Su/Zhu/Zhang, PLB718,1510(2013)

n/p v2 (FOPI): $(\rho/\rho_0)^\gamma$ with $\gamma = 0.9 \pm 0.4$

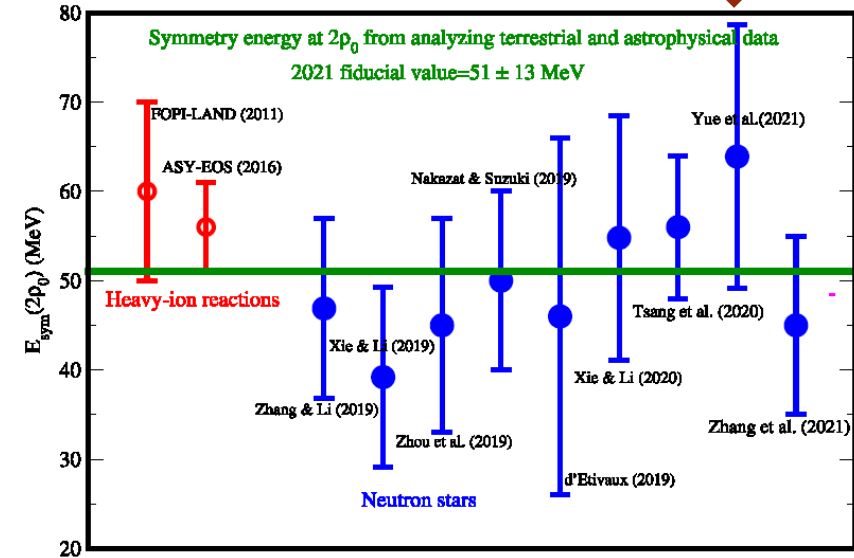
Rusotto/Trautmann/Li et al.,
PLB697, 471(2011) (UrQMD)

PRC94, 034608 (2016) $\gamma = 0.72 \pm 0.19$

Cozma/Trautmann/Li et al.,
PRC88, 044912 (2013) (Tubingen QMD - MDI)

Tong-Gang Yue (岳伺钢),
LWC, Z. Zhang (张振),
and Y. Zhou (周颖),
arXiv:2102.05267

PREX-II



B.A.Li, B.J. Cai, W.J. Xie, and N.B. Zhang, Universe 7, 187 (2021)



E_{sym} : Current Status

- There are MANY constraints on $E_{\text{sym}}(\rho_0)$ and L , and the world average values are:

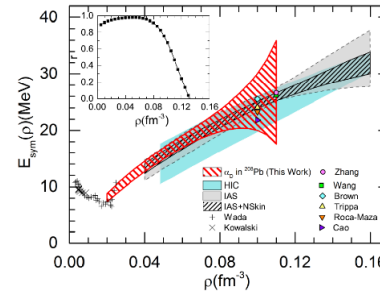
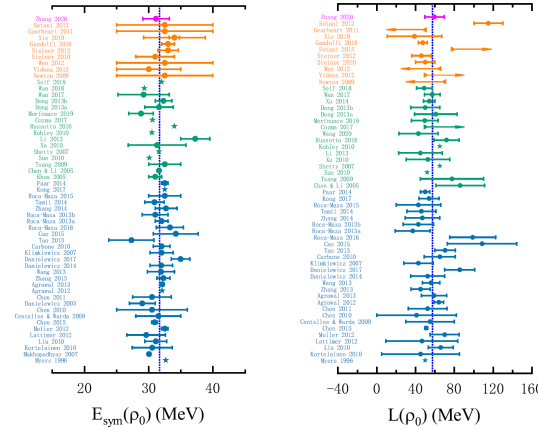
$$E_{\text{sym}}(\rho_0) = 31.7 \pm 3.1 \text{ MeV}$$

$$L = 57.5 \pm 24.5 \text{ MeV}$$

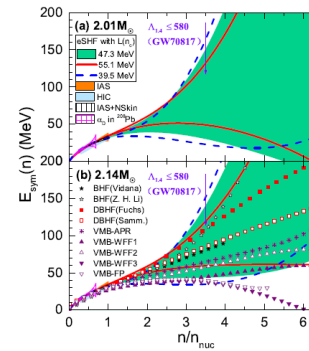
Very recent PREX-II makes the situation elusive !!! (This Talk)

- The symmetry energy at subsaturation densities have been relatively well-constrained

- Based on the GW multimessenger measurements, the high density E_{sym} cannot be too stiff or too soft but still with large uncertainty!!! (This Talk)



Z. Zhang(张振)/LWC,
PLB726, 234 (2013);
PRC92, 031301(R)(2015)



Y. Zhou (周颖), LWC,
ApJ886, 52(2019)
[arXiv:1907.12284]



Outline

- Nuclear matter EOS and the symmetry energy (E_{sym})
- Dense nuclear matter from *Nuclear experiments + Neutron star mass and radius + Tidal deformability from GW170817*
- Summary and outlook



Characteristic Parameters of NM EOS

PHYSICAL REVIEW C 80, 014322 (2009)

Higher-order effects on the incompressibility of isospin asymmetric nuclear matter

Lie-Wen Chen,^{1,2} Bao-Jun Cai,¹ Che Ming Ko,³ Bao-An Li,⁴ Chun Shen,¹ and Jun Xu³

¹Department of Physics, Shanghai Jiao Tong University, Shanghai 200240, People's Republic of China

²Center of Theoretical Nuclear Physics, National Laboratory of Heavy Ion Accelerator, Lanzhou 730000, People's Republic of China

³Cyclotron Institute and Physics Department, Texas A&M University, College Station, Texas 77843-3366, USA

⁴Department of Physics, Texas A&M University-Commerce, Commerce, Texas 75429-3011, USA

(Received 27 May 2009; published 30 July 2009)

$$E(\rho, \delta) = E_0(\rho) + E_{\text{sym}}(\rho)\delta^2 + E_{\text{sym},4}(\rho)\delta^4 + O(\delta^6)$$

$$E_0(\rho) = E_0(\rho_0) + \frac{K_0}{2!}\chi^2 + \frac{J_0}{3!}\chi^3 + \frac{I_0}{4!}\chi^4 + O(\chi^5)$$

$$\delta = (\rho_n - \rho_p)/\rho$$

$$E_{\text{sym}}(\rho) = E_{\text{sym}}(\rho_0) + L\chi + \frac{K_{\text{sym}}}{2!}\chi^2 + \frac{J_{\text{sym}}}{3!}\chi^3 + \frac{I_{\text{sym}}}{4!}\chi^4 + O(\chi^5)$$

$$\chi = \frac{\rho - \rho_0}{3\rho_0}$$

$$E_{\text{sym},4}(\rho) = E_{\text{sym},4}(\rho_0) + L_{\text{sym},4}\chi + \frac{K_{\text{sym},4}}{2}\chi^2 + \frac{J_{\text{sym},4}}{3!}\chi^3 + \frac{I_{\text{sym},4}}{4!}\chi^4 + O(\chi^5)$$

Order of the characteristic parameters according to the expansion with χ and δ :

Order-0: $E_0(\rho_0)$; **Order-2:** $K_0, E_{\text{sym}}(\rho_0)$;

Order-3: J_0, L ; **Order-4:** $I_0, K_{\text{sym}}(\rho_0), E_{\text{sym},4}(\rho_0)$


Jie Pu (普洁)
et al., to be
Submitted





Characteristic Parameters of NM EOS


Order of the characteristic parameters according to the expansion with χ and δ :

Order-0: $E_0(\rho_0)$; **Order-2:** $K_0, E_{\text{sym}}(\rho_0)$;
Order-3: J_0, L ; **Order-4:** $I_0, K_{\text{sym}}(\rho_0), E_{\text{sym},4}(\rho_0)$

Order-0  $E_0(\rho_0) = -16 \pm 1 \text{ MeV}$

Order-2  $K_0 = 240 \pm 20 \text{ MeV}, E_{\text{sym}}(\rho_0) = 32.5 \pm 2.5 \text{ MeV}???$

Order-3  $L = 55 \pm 25 \text{ MeV}???, J_0 = ???$

Order-4  $I_0 = ???, K_{\text{sym}}(\rho_0) = ???, E_{\text{sym},4}(\rho_0) = ???$

□ $J_0 \approx ?$ and $K_{\text{sym}} \approx ?$ High Density Behaviors?

Data of finite nuclei + Flow Data in HIC + Observed NStar Largest Mass +
NStar M-R from **NICER** + Tidal Deformability of Neutron Star (from **GW170817**)
analyzed simultaneously within a unified **Nuclear Energy Density Functional**



Why a unified Nuclear Energy Density Functional?

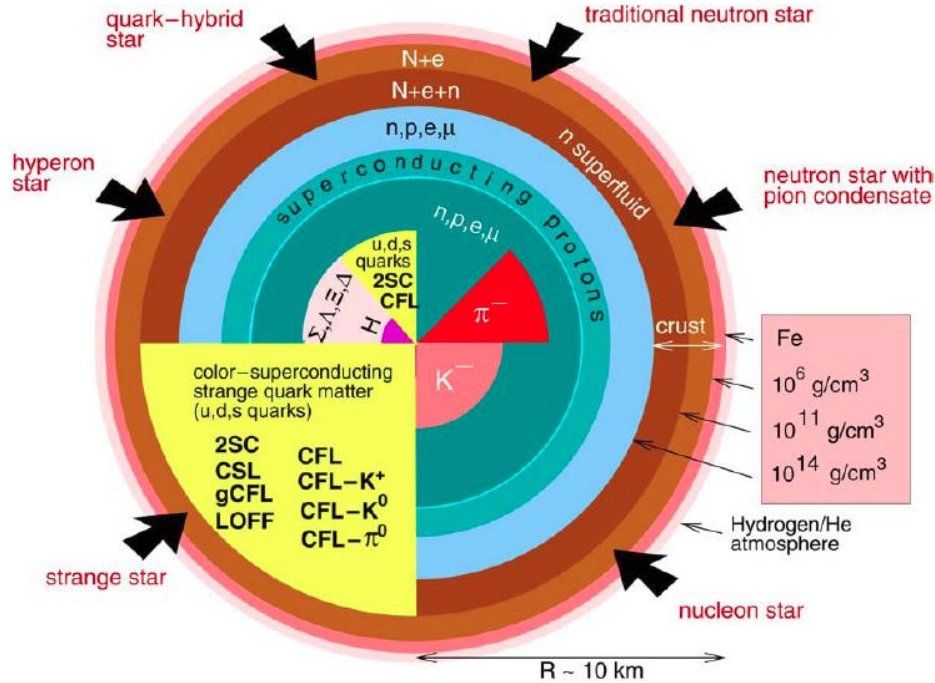
The nuclear energy density functional theory (e.g., SHF/RMF) is still the only realistic framework to simultaneously investigate the physics of heavy nuclei and neutron stars as well as heavy-ion collisions !!!



Composition of Neutron Star Matter

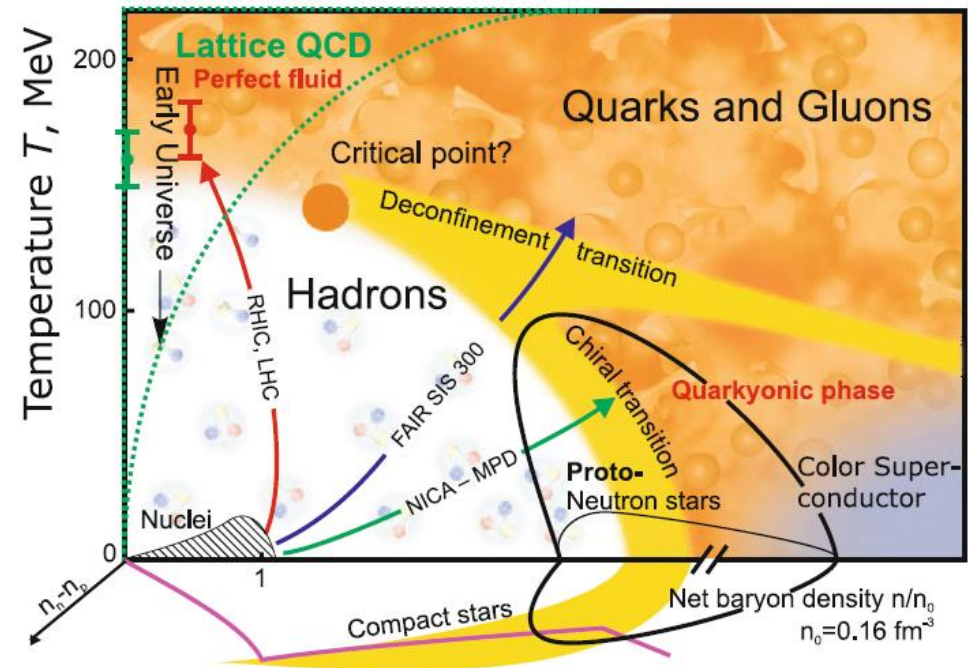
Pulsars: Neutron Stars? Quark Stars? Hybrid Stars? Others?

F. Weber, PPNP54, 193 (2005)



Mass: $\sim 1.4 M_{\odot}$, Radius: ~ 10 km
Extremely neutron-rich matter
Density at the center: $\sim 6\rho_0$
Average density: $\sim 2.5\rho_0$

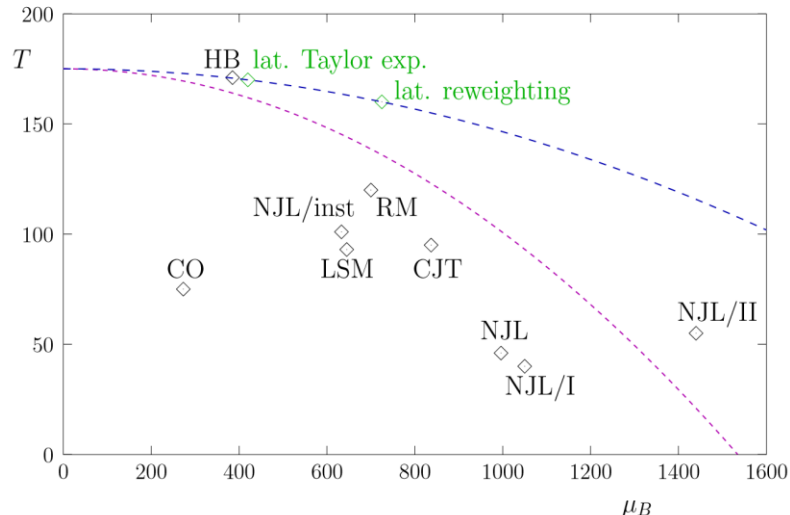
V.E. Fortov, Extreme States of Matter – on Earth and in the Cosmos, Springer-Verlag Berlin Heidelberg 2011



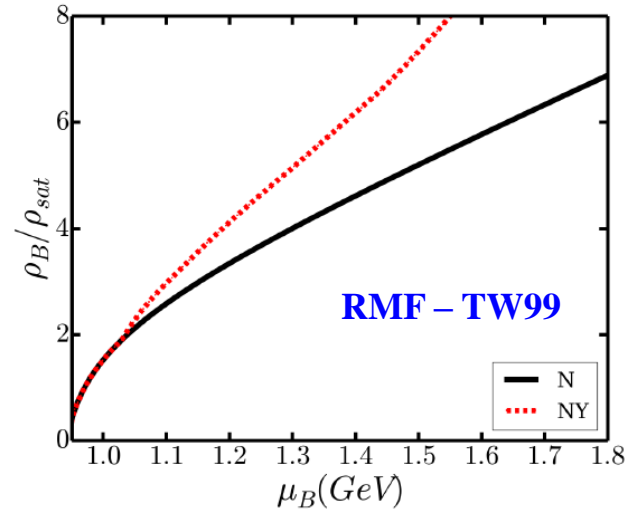
The mass of ud Quark Star can be $2.77M_{\odot}$, Z. Cao (曹政), LWC, P.C. Chu (初鹏程), and Y. Zhou (周颖), arXiv:2009.00942.



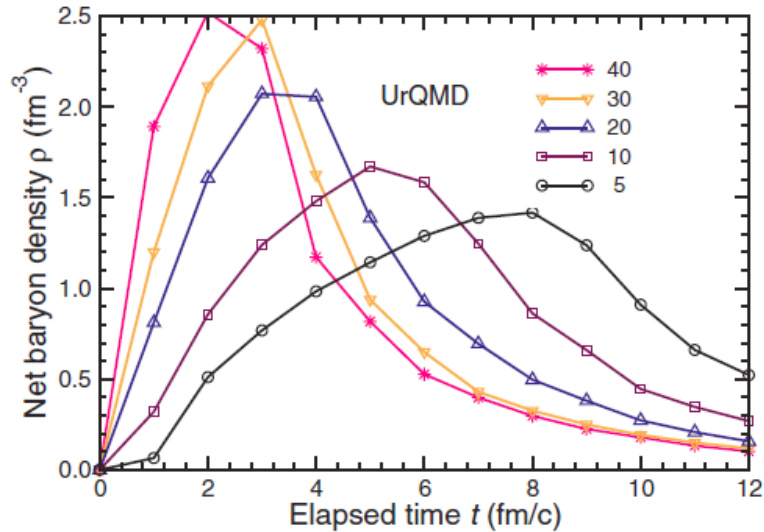
Composition of Neutron Star Matter



M.A. Stephanov, *Int. J. Mod. Phys. A*20, 4387 (2005).



Z. Bai/Y.X. Liu, arXiv:1904.01978



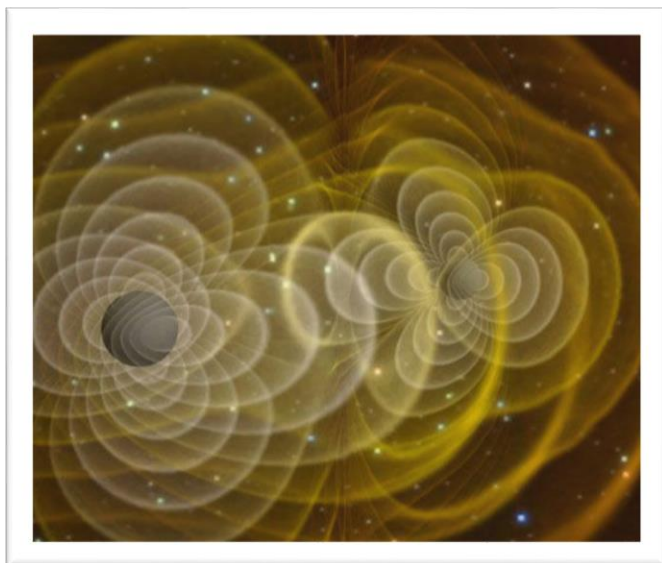
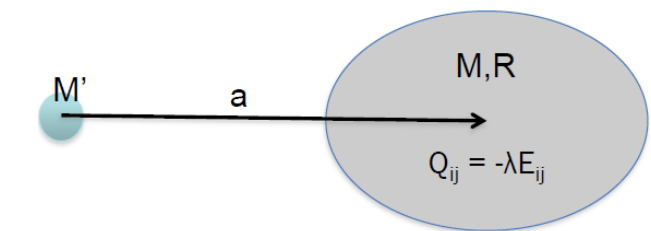
I.C. Arsene et al., *PRC*75, 034902(2007)

- The maximum baryon density could reach about $30\rho_0$ in heavy ion collisions at **but with high temperature** (>110 MeV)
- The critical density for deconfinement/chiral symmetry restoration at $T = 0$ could be $\gg 6\rho_0$ ($\mu_B \sim 1.6$ GeV)!
- Quark phase exists inside Nstars?? (E. Annala et al., *Nature Physics* (2020))



Tidal Deformability

Tidal Deformability (Polarizability) (oscillation response coefficient λ)



$$Q_{ij} = \lambda \varepsilon_{ij}$$

Q_{ij} : Quadrupole moment

ε_{ij} : Tidal field of companion

$$\lambda = \frac{2}{3} k_2 R^5$$

k_2 : Love number

R : Radius

M : Mass

Dimensionless Tidal Deformability

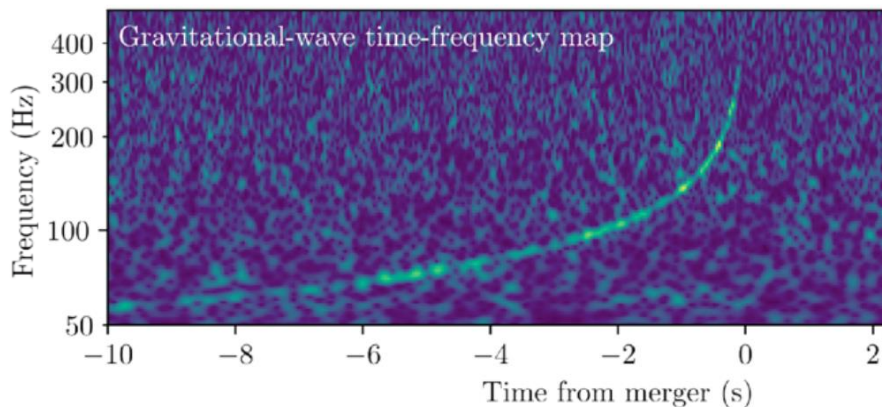
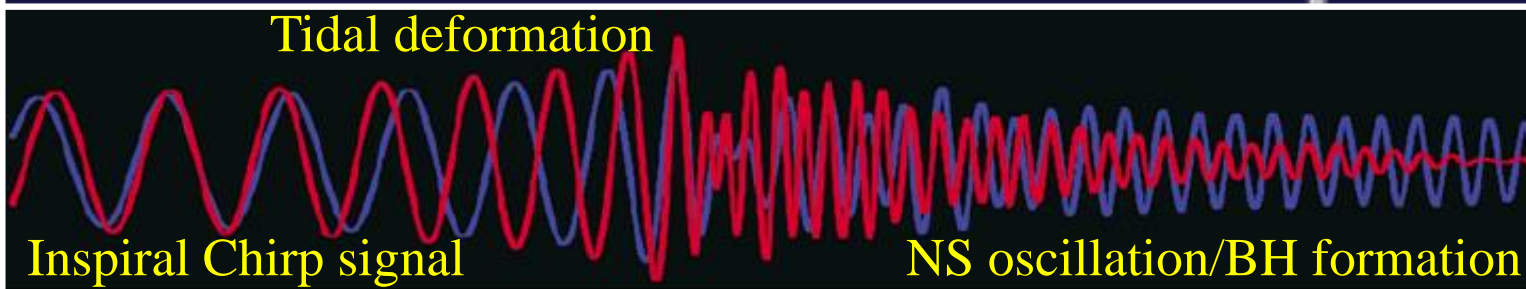
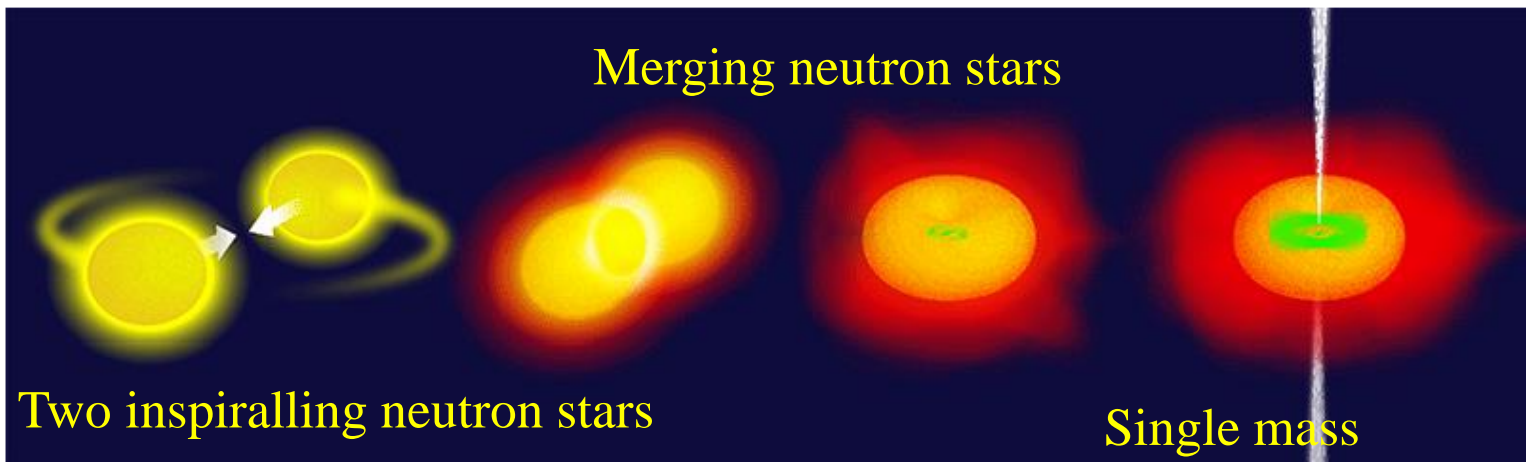
$$\Lambda = \frac{2}{3} k_2 (R / M)^5$$

Éanna É. Flanagan and Tanja Hinderer, Phys.Rev.D 77, 021502(R) (2008)

F.J. Fattoyev, J. Carvajal, W.G. Newton, and Bao-An Li, Phys. Rev. C 87, 015806 (2013)



Binary Neutron Star Merger



LIGO is sensitive to increase in orbital frequency as system loses energy to both gravitational waves and internal excitation of neutron stars. GW170817 data place limits on polarizability (deformability) Λ of NS and hence limits on NS radius.



Multimessenger era: GW170817

PRL 119, 161101 (2017)

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week ending
20 OCTOBER 2017



arXiv:1710.05832

GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral

B. P. Abbott *et al.**

(LIGO Scientific Collaboration and Virgo Collaboration)

Citations: 5669+



- ❑ On August 17, 2017, the merger of two neutron stars was observed with gravitational waves (GW) by the LIGO and Virgo detectors.
- ❑ The Fermi and Integral spacecrafts independently detected a short gamma ray burst.
- ❑ Extensive follow up observations detected this event at X-ray, ultra-violet, visible, infrared, and radio wavelengths.
- ❑ No ultra-high-energy gamma-rays and no neutrino candidates consistent with the source were found in follow-up searches. **These observations support the hypothesis that GW170817 was produced by the merger of two neutron stars in NGC 4993 followed by a short gamma-ray burst (GRB 170817A) and a kilonova/macronova powered by the radioactive decay of *r*-process nuclei synthesized in the ejecta**

Our assumptions

□ Core of the neutron stars consist of **infinite β -equilibrium $n_{pe\mu}$ matter with charge neutrality**. Its EOS is determined by **the Nuclear Energy Density Functionals**

□ The inner crust $2.46 \times 10^{-4} \text{ fm}^{-3} = n_{\text{out}} < n < n_t$

n_t is determined self-consistently by using dynamical method

(J Xu(徐骏)/LWC/Li/Ma, ApJ697,1549(2009))

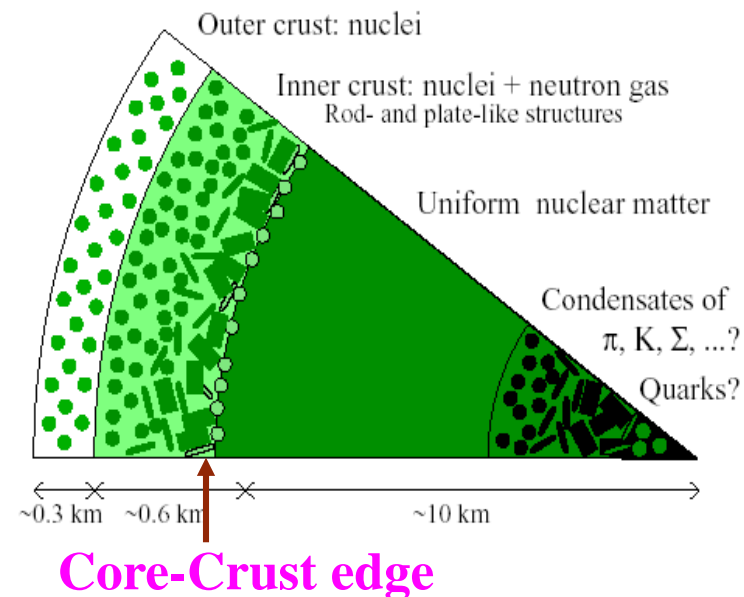
$$P = a + b\epsilon^{4/3}$$

$$a = \frac{P_{\text{out}}\epsilon_t^{4/3} - P_t\epsilon_{\text{out}}^{4/3}}{\epsilon_t^{4/3} - \epsilon_{\text{out}}^{4/3}} \quad b = \frac{P_t - P_{\text{out}}}{\epsilon_t^{4/3} - \epsilon_{\text{out}}^{4/3}}$$

□ The outer crust

$$6.93 \times 10^{-13} \text{ fm}^{-3} < n < n_{\text{out}} \quad (\text{EOS of BPS})$$

$$4.73 \times 10^{-15} \text{ fm}^{-3} < n < 6.93 \times 10^{-13} \text{ fm}^{-3} \quad (\text{EOS of Feynman-Metropolis-Teller})$$





The extended SHF Energy Density Functional

Extended Skyrme Interaction:

$$\begin{aligned}
 v_{i,j} = & t_0(1 + x_0 P_\sigma) \delta(r) \\
 & + \frac{1}{2} t_1(1 + x_1 P_\sigma) [\mathbf{K}'^2 \delta(r) + \delta(r) \mathbf{K}^2] \\
 & + t_2(1 + x_2 P_\sigma) \mathbf{K}' \cdot \delta(r) \mathbf{K} \\
 & + \frac{1}{6} t_3(1 + x_3 P_\sigma) n(\mathbf{R})^\alpha \delta(r) \\
 & + iW_0(\boldsymbol{\sigma}_i + \boldsymbol{\sigma}_j) \mathbf{K}' \cdot \delta(r) \mathbf{K} \\
 & + \frac{1}{2} t_4(1 + x_4 P_\sigma) [\mathbf{K}'^2 n(\mathbf{R})^\beta \delta(r) + \delta(r) n(\mathbf{R})^\beta \mathbf{K}^2] \\
 & + t_5(1 + x_5 P_\sigma) \mathbf{K}' \cdot n(\mathbf{R})^\gamma \delta(r) \mathbf{K}
 \end{aligned}$$

Momentum-dependence of many-body forces

N. Chamel, S. Goriely, and J.M. Pearson, PRC80, 065804 (2009)

Z. Zhang/LWC, PRC94, 064326 (2016)

LWC/Ko/Li/Xu, PRC82, 024321(2010)

13 Skyrme parameters: $\alpha, t_0 \sim t_5, x_0 \sim x_5$

13 macroscopic nuclear properties:

$$n_0, E_0, K_0, J_0, E_{\text{sym}}, L, K_{\text{sym}}, m_{s,0}^*, m_{v,0}^*, G_S, G_V, G_{SV}, G'_0$$

$$\begin{aligned}
 \mathcal{H} = & \mathcal{K} + \mathcal{H}_0 + \mathcal{H}_3 + \mathcal{H}_{\text{eff}} + \frac{G_S}{2} (\nabla \rho)^2 - \frac{G_V}{2} (\nabla \rho_1)^2 \\
 & - \frac{G_{SV}}{2} \delta \nabla \rho \nabla \rho_1 + \mathcal{H}_{\text{Coul}} + \mathcal{H}_{\text{SO}} + \mathcal{H}_{\text{sg}}, \quad (
 \end{aligned}$$



Why extended SHF EDF?

PHYSICAL REVIEW C 94, 064326 (2016)

Extended Skyrme interactions for nuclear matter, finite nuclei, and neutron stars

(张振, 中山大学) Zhen Zhang¹ and Lie-Wen Chen^{1,2,*}

¹Department of Physics and Astronomy and Shanghai Key Laboratory for Particle Physics and Cosmology,
Shanghai Jiao Tong University, Shanghai 200240, China

²Center of Theoretical Nuclear Physics, National Laboratory of Heavy Ion Accelerator, Lanzhou 730000, China

symmetry energy softer at subsaturation densities (favored by experimental constraints and theoretical predictions) but stiffer at higher densities (favored by the observation of $2M_{\odot}$ neutron stars) challenges the SHF model with the conventional Skyrme interactions. For example, the Skyrme interaction TOV-min [28], which is built by fitting properties of both finite nuclei and neutron stars, can successfully support $2M_{\odot}$ neutron stars but predicts a neutron matter EOS significantly deviating from the ChEFT calculations [14] as well as the constraint extracted from analyzing the electric-dipole polarizability in ^{208}Pb [49] at densities below about $0.5\rho_0$.

Furthermore, it is well known that a notorious shortcoming of the conventional standard Skyrme interactions is that they predict various instabilities of nuclear matter around saturation density or at supra-saturation densities, which in principle hinders the application of the Skyrme interactions in the study of dense nuclear matter as well as neutron stars. For instance, most of the conventional standard Skyrme interactions predict spin or spin-isospin polarization in the density region of about $(1 \sim 3.5)\rho_0$ [25,51], including the famous SLy4 interaction [19] which has been widely used in both nuclear physics and neutron star studies and leads to spin-isospin instability of symmetric nuclear matter at densities beyond about $2\rho_0$ [52]. On the other hand, the calculations

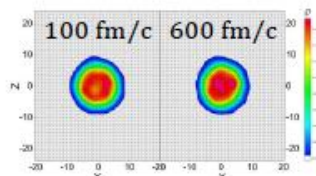
Momentum-dependence of many-body forces plays an important role for the high-density behaviors of EOS!

- The eSHF can describe simultaneously **nuclear matter, finite nuclei, and neutron stars!**
- The eSHF EDF is **very flexible to mimic various density behaviors for EOS (13 parameters)**

$$E_0(\rho) = a_2\rho^{2/3} + a_3\rho^{3/3} + a_4\rho^{(\alpha+3/3)} + a_5\rho^{5/3} + a_6\rho^{(\beta+5/3)} + a_7\rho^{(\gamma+5/3)}$$

Ground state LHV Calculations

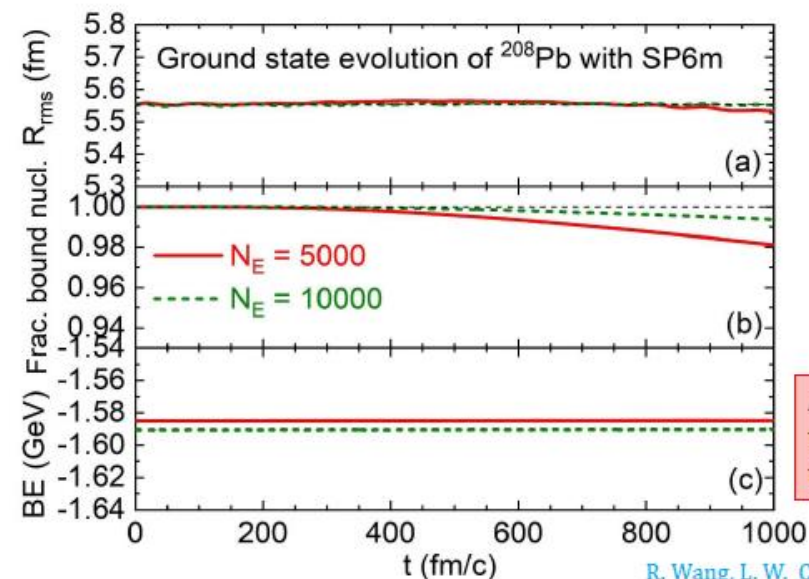
The rms radius, fraction of bound nucleons, binding energy of ground state evolution



Free test nucleons are those whose **form factor** do not overlap with that of other test nucleons

$$\langle r^2 \rangle = N_{\text{nucl}}^{-1} \int_{\text{bound}} d^3r r^2 \rho(r)$$

At 1000 fm/c, only about 2% of test nucleons evaporated from the nucleus for $N_E = 5000$ case



R. Wang, L. W. Chen, and Z. Zhang, PRC99, 044609 (2019)

Very stable ground state of the initialized nuclei!

R. Wang (王睿), LWC, and Z. Zhang (张振), PRC99, 044607 (2019)

The **GPU** parallel computing with large enough test particle number (up to **~100000** !)



Constraining the in-medium nucleon-nucleon cross section from the width of nuclear giant dipole resonance

Rui Wang^{a,b}, Zhen Zhang^c, Lie-Wen Chen^d, Che Ming Ko^e, Yu-Gang Ma^{a,b}

基于格点Hamiltonian方法和随机碰撞方法，发展了完整的格点Boltzmann-Uehling-Uhlenbeck(BUU)输运模型并成功用于描述原子核的集体运动

Invited Review



Nuclear Collective Dynamics in Transport Model With the Lattice Hamiltonian Method

Rui Wang^{1,2}, Zhen Zhang³, Lie-Wen Chen^{4*} and Yu-Gang Ma^{1,2}

王睿, 张振, LWC, and Y.G.Ma



Extended Skyrme forces with fixed J_0 & K_{sym}

$$n_0, E_0, K_0, J_0, E_{\text{sym}}, L, K_{\text{sym}}, m_{s,0}^*, m_{v,0}^*, G_S, G_V, G_{SV}, G'_0$$

TABLE I. Experimental data for 12 spherical even-even nuclei binding energies E_B [27], charge r.m.s. radii r_c [28–30], ISGMR energies E_{GMR} and its experimental error [31], and spin-orbit energy level splittings ϵ_{ls}^A [32].

$\frac{A}{Z}X$	$E_B(\text{MeV})$	$r_c(\text{fm})$	$E_{\text{GMR}}(\text{MeV})$	$\epsilon_{\text{ls}}^A(\text{MeV})$
^{16}O	-127.619	2.6991	...	6.30(1pν) 6.10(1pπ)
^{40}Ca	-342.052	3.4776
^{48}Ca	-416.001	3.4771
^{56}Ni	-483.995	3.7760
^{68}Ni	-590.408
^{88}Sr	-768.468	4.2240
^{90}Zr	-783.898	4.2694	17.81±0.35	...
^{100}Sn	-825.300
^{116}Sn	-988.681	4.6250	15.90±0.07	...
^{132}Sn	-1102.84
^{144}Sm	-1195.73	4.9524	15.25±0.11	...
^{208}Pb	-1636.43	5.5012	14.18±0.11	1.32(2dπ) 0.89(3pν) 1.77(2fν)

Our Strategy:

- Higher-order **J_0** and **K_{sym}** are fixed at various values
- **$E_{\text{sym}}(\rho_c)$** and **$L(\rho_c)$** at $\rho_c = 0.11 \text{ fm}^{-3}$ are fixed at **$E_{\text{sym}}(\rho_c) = 26.65 \text{ MeV}$** and **$L(\rho_c) = 47.3 \pm 7.8 \text{ MeV}$** using heavy isotope binding energy difference and α_D of ^{208}Pb (Z. Zhang(张振)/LWC, PLB726, 234(2013); PRC90, 064317(2014))
- Other **9 lower-order parameters** and **W_0** are calibrated to fit data of finite nuclei
- Causality

Lower-order parameters are basically fixed by nuclear properties, while **higher-order parameters** are not important for nuclei

Minimizing the Chi-square $\chi^2(p)$:

$$\chi^2(P) = \sum_{n=1}^N \left(\frac{\mathcal{O}_n^{(\text{th})}(P) - \mathcal{O}_n^{(\text{exp})}}{\Delta \mathcal{O}_n} \right)^2$$



Observed NStar max. mass and Tidal def.

28 OCTOBER 2010 | VOL 467 | NATURE | 1081

LETTER

doi:10.1038/nature09466

A two-solar-mass neutron star measured using Shapiro delay

P. B. Demorest¹, T. Pennucci², S. M. Ransom¹, M. S. E. Roberts³ & J. W. T. Hessels^{4,5}

Observed heaviest Nstar so far (before 2019):

A Massive Pulsar in a Compact Relativistic Binary

John Antoniadis *et al.*

Science **340**, (2013);

DOI: 10.1126/science.1233232



PSR J0348+0432

2.01 ± 0.04 solar mass (M_{\odot})

PRL 119, 161101 (2017)

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GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral

B. P. Abbott *et al.**

(LIGO Scientific Collaboration and Virgo Collaboration)

PRL121, 161101 (2018)

GW170817: Measurements of neutron star radii and equation of state

The LIGO Scientific Collaboration and The Virgo Collaboration
(compiled 30 May 2018)



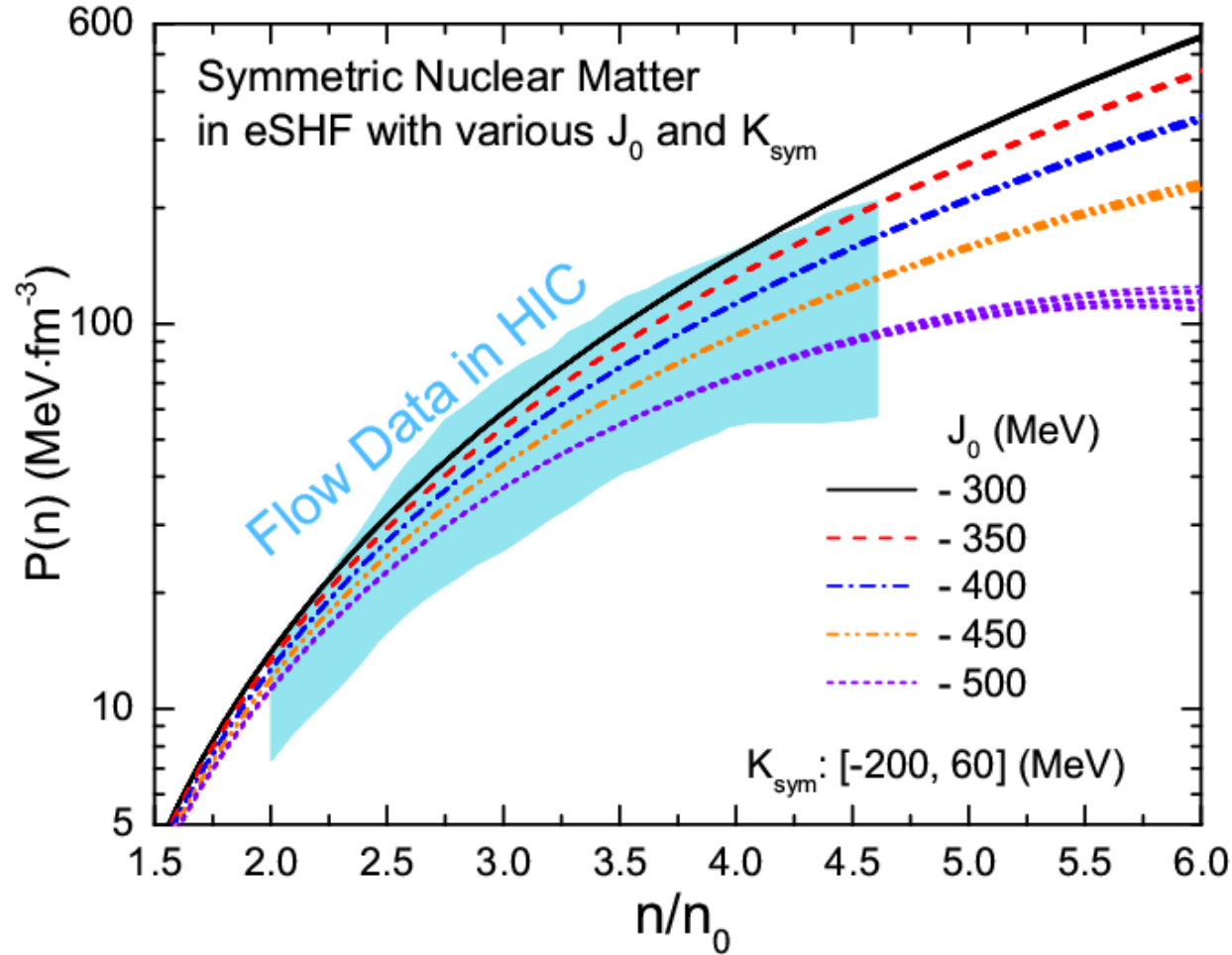
GW170817 (LIGO/Virgo):

70 < $\Lambda_{1.4}$ < 580 (Assuming NS)



J0: Flow data in HIC's

Y. Zhou(周颖)/LWC/Z. Zhang(张振), PRD99, 121301(R) (2019) [arXiv:1901.11364]



For various J_0 and
 K_{sym} : [-200, 60] MeV

$E_{\text{sym}}(\rho_c) = 26.65$ MeV
 $L(\rho_c) = 47.3$ MeV

Pressure of SNM is
very sensitive to J_0
but essentially
independent of K_{sym}

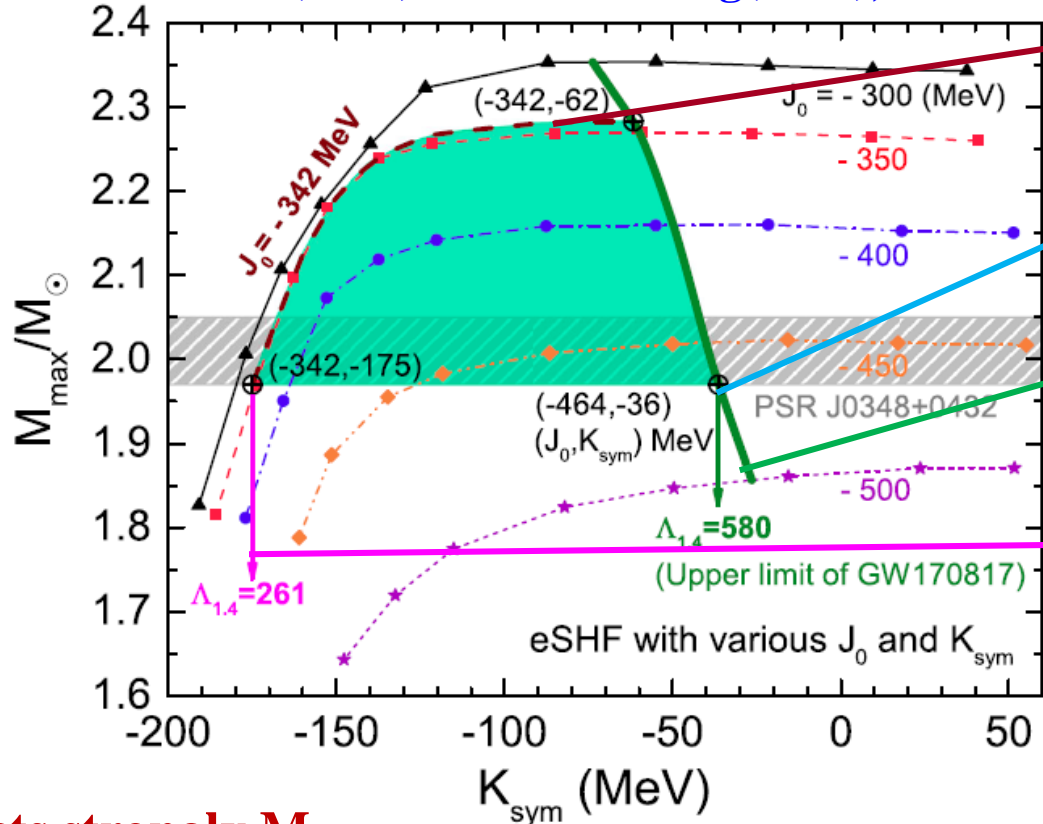
Lower-order parameters are
basically fixed by nuclear
properties

-550 MeV $\sim < J_0 < -342$ MeV: Flow Data in HIC's



J0 and Ksym: Flow data, NStar Mass, Λ

Y. Zhou(周颖)/LWC/Z. Zhang(张振), PRD99, 121301(R) (2019) [arXiv:1901.11364]



$J_0 = -342$ MeV (from flow data)

$J_0 = -464$ MeV

$K_{\text{sym}} < -36$ MeV

$\Lambda_{1.4} < 580$

GW170817 (LIGO/Virgo)

$\Lambda_{1.4} > 261$

Lower-order parameters are basically fixed by nuclear properties

K_{sym} affects strongly M_{\max}
for $K_{\text{sym}} < -100$ MeV

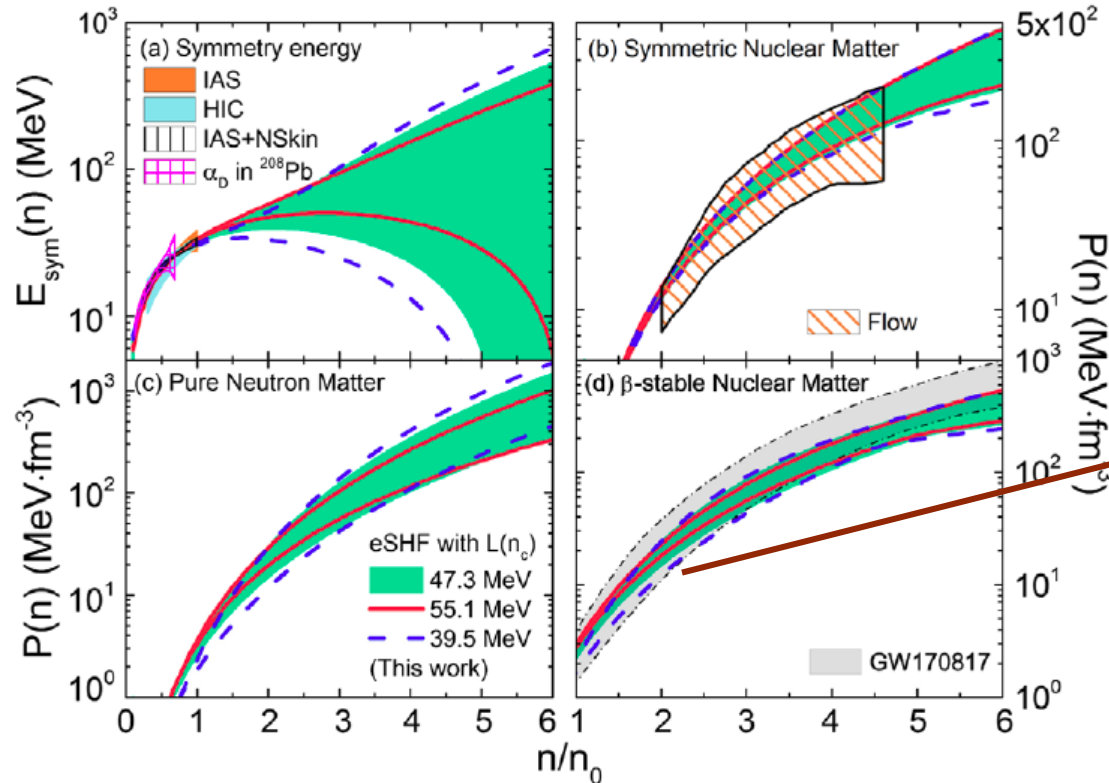
Flow Data in HIC + M_{\max} :
 $K_{\text{sym}} > -175$ MeV, $\Lambda_{1.4} > 261$

Flow Data in HIC + M_{\max} + $\Lambda_{1.4}$:
 -464 MeV $< J_0 < -342$ MeV; -175 MeV $< K_{\text{sym}} < -36$ MeV:



EOS: Flow data, NStar Mass, Λ

Y. Zhou(周颖)/LWC/Z. Zhang(张振), PRD99, 121301(R) (2019) [arXiv:1901.11364]



$L(\rho_c)=47.3\pm 7.8$ MeV using α_D of ^{208}Pb (Z. Zhang(张振), LWC, PRC90, 064317(2014))

$L(\rho_c)$ indeed affects the extraction of E_{sym} at high densities but does not change much the NStar matter EOS!

Consistent with LIGO/Virgo constraints (see, e.g., D. Radice et al., ApJL852, L29(2018)) but more stringent due to nuclear data added

$L(\rho_c)=47.3$ MeV:
 $J_0: [-464, -342]$ MeV,
 $K_{\text{sym}}: [-175, -36]$ MeV
 $E_{\text{sym}}(2\rho_0): [39.4, 54.5]$ MeV

$L(\rho_c)=39.5$ MeV:
 $J_0: [-475, -342]$ MeV,
 $K_{\text{sym}}: [-203, -34]$ MeV
 $E_{\text{sym}}(2\rho_0): [33.0, 51.3]$ MeV

$L(\rho_c)=55.1$ MeV:
 $J_0: [-455, -342]$ MeV,
 $K_{\text{sym}}: [-138, -38]$ MeV
 $E_{\text{sym}}(2\rho_0): [46.9, 57.6]$ MeV



Relativistic Shapiro delay measurements of an extremely massive millisecond pulsar

H. T. Cromartie^{1*}, E. Fonseca², S. M. Ransom³, P. B. Demorest⁴, Z. Arzoumanian⁵, H. Blumer^{6,7}, P. R. Brook^{6,7}, M. E. DeCesar⁸, T. Dolch⁹, J. A. Ellis¹⁰, R. D. Ferdman¹¹, E. C. Ferrara^{12,13}, N. Garver-Daniels^{6,7}, P. A. Gentile^{6,7}, M. L. Jones^{6,7}, M. T. Lam^{6,7}, D. R. Lorimer^{6,7}, R. S. Lynch¹⁴, M. A. McLaughlin^{6,7}, C. Ng^{15,16}, D. J. Nice⁸, T. T. Pennucci¹⁷, R. Spiewak¹⁸, I. H. Stairs¹⁵, K. Stovall⁴, J. K. Swiggum¹⁹ and W. W. Zhu²⁰

edge-on) binary pulsar systems. By combining data from the North American Nanohertz Observatory for Gravitational Waves (NANOGrav) 12.5-yr data set with recent orbital-phase-specific observations using the Green Bank Telescope, we have measured the mass of the MSP J0740+6620 to be $2.14_{-0.09}^{+0.10} M_{\odot}$ (68.3% credibility interval; the 95.4% credibility interval is $2.14_{-0.18}^{+0.20} M_{\odot}$). It is highly likely to be the most massive neutron star yet observed, and serves as a strong constraint on the neutron star interior EoS.

THE ASTROPHYSICAL JOURNAL LETTERS, 915:L12 (15pp), 2021 July 1
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<https://doi.org/10.3847/2041-8213/ac03b8>



Refined Mass and Geometric Measurements of the High-mass PSR J0740+6620

E. Fonseca^{1,2,3,4}, H. T. Cromartie^{5,37}, T. T. Pennucci^{6,7}, P. S. Ray⁸, A. Yu. Kirichenko^{9,10}, S. M. Ransom⁶, P. B. Demorest¹¹, I. H. Stairs¹², Z. Arzoumanian¹³, L. Guillemot^{14,15}, A. Parthasarathy¹⁶, M. Kerr³, I. Cognard^{14,15}, P. T. Baker¹⁷, H. Blumer^{3,4}, P. R. Brook^{3,4}, M. DeCesar¹⁸, T. Dolch^{19,20}, F. A. Dong¹², E. C. Ferrara^{21,22,23}, W. Fiore^{3,4}, N. Garver-Daniels^{3,4}, D. C. Good¹², R. Jennings²⁴, M. L. Jones²⁵, V. M. Kaspi^{1,2}, M. T. Lam^{26,27}, D. R. Lorimer^{3,4}, J. Luo²⁸, A. McEwen²⁵, J. W. McKee²⁸, M. A. McLaughlin^{3,4}, N. McMann²⁹, B. W. Meyers¹², A. Naidu³⁰, C. Ng³¹, D. J. Nice³², N. Poj²⁹, H. A. Radovan³³, B. Shapiro-Albert^{3,4}, C. M. Tan^{1,2}, S. P. Tendulkar^{34,35}, J. K. Swiggum³², H. M. Wahl^{3,4}, and W. W. Zhu³⁶

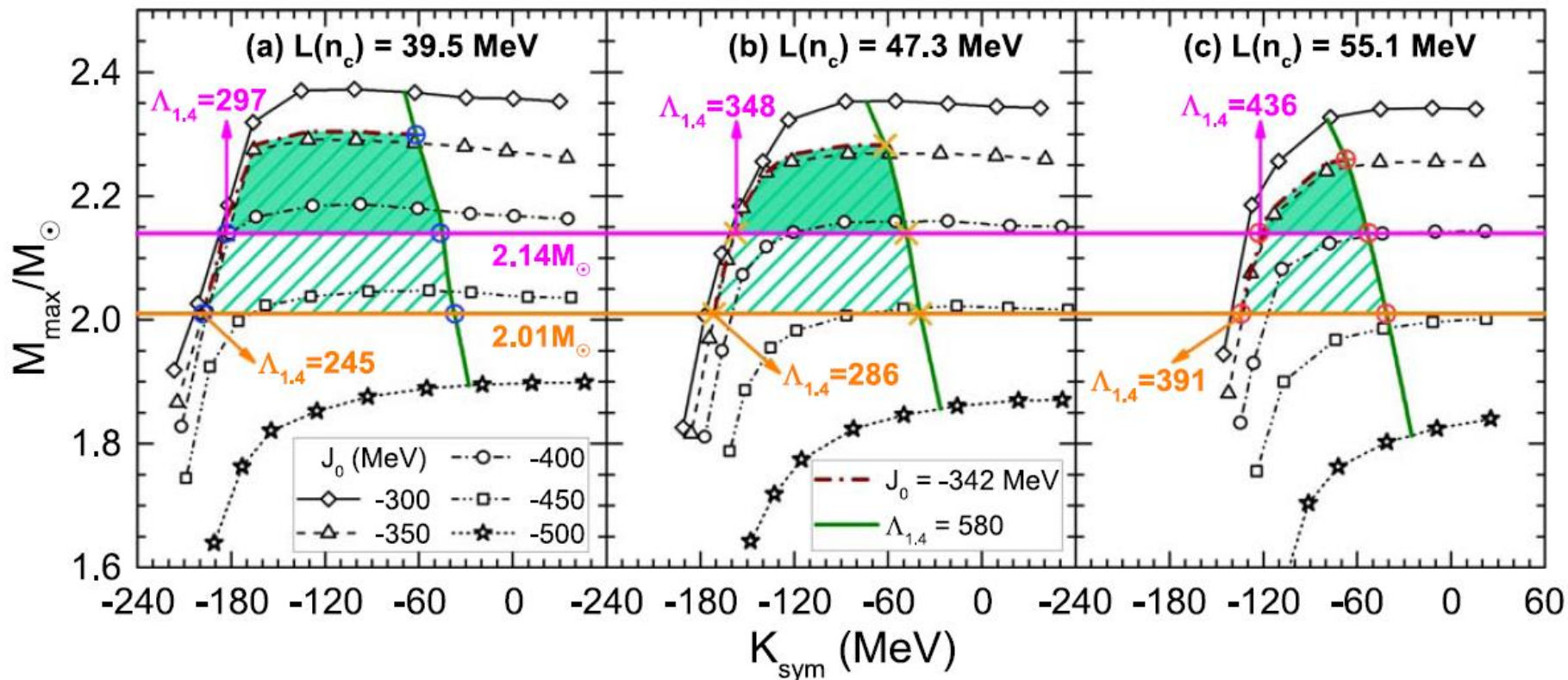
ABSTRACT

We report results from continued timing observations of PSR J0740+6620, a high-mass, 2.8-ms radio pulsar in orbit with a likely ultra-cool white dwarf companion. Our data set consists of combined pulse arrival-time measurements made with the 100-m Green Bank Telescope and the Canadian Hydrogen Intensity Mapping Experiment telescope. We explore the significance of timing-based phenomena arising from general-relativistic dynamics and variations in pulse dispersion. When using various statistical methods, we find that combining ~ 1.5 years of additional, high-cadence timing data with previous measurements confirms and improves upon previous estimates of relativistic effects within the PSR J0740+6620 system, with the pulsar mass $m_p = 2.08_{-0.07}^{+0.07} M_{\odot}$ (68.3% credibility) determined by the relativistic Shapiro time delay. For the first time, we measure secular variation in the orbital period and argue that this effect arises from apparent acceleration due to significant transverse motion. After incorporating contributions from Galactic differential rotation and off-plane acceleration in the Galactic potential, we obtain a model-dependent distance of $d = 1.14_{-0.15}^{+0.17}$ kpc (68.3% credibility). This improved distance confirms the ultra-cool nature of the white dwarf companion determined from recent optical observations. We discuss the prospects for future observations with next-generation facilities, which will likely improve the precision on m_p for J0740+6620 by an order of magnitude within the next few years.



New Observed Heaviest NStar

Y. Zhou(周颖), L.W. Chen*, ApJ886, 52(2019)
[arXiv:1907.12284]



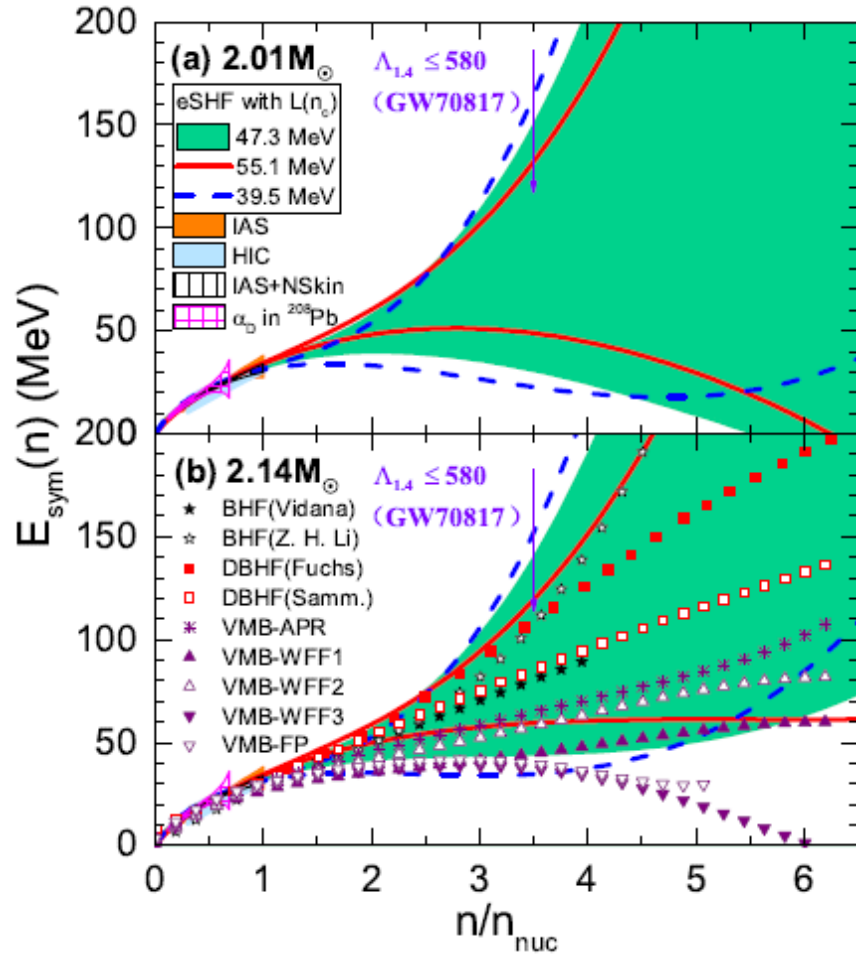
$\Lambda_{1.4} > 297$

The maximum mass of static NS is about $2.3M_{\odot}$!



Ruling out the “supersoft” E_{sym}

Y. Zhou (周颖), L.W. Chen*, ApJ886, 52(2019) [arXiv:1907.12284]



$$\frac{\hbar^2 k_F^2}{3m_s^*} (1 + F'_0) = \left. \frac{\partial^2 \mathcal{E}(\rho, \rho_1, s_0, s_1)}{\partial (\rho_1)^2} \right|_{\rho_1 = s_0 = s_1 = 0}$$

$$E_{\text{sym}}(\rho) = \frac{\hbar^2 k_F^2}{6m_s^*} (1 + F'_0) \quad \text{Z. Zhang(张振)/LWC, PRC94, 064326 (2016)}$$

The Landau stability conditions,

$$\mu_e = \mu_n - \mu_p = 4\delta E_{\text{sym}}(\rho)$$

$$F_l > -(2l + 1),$$

$$F'_l > -(2l + 1),$$

$$G_l > -(2l + 1),$$

$$G'_l > -(2l + 1),$$

Negative E_{sym} leads to
isospin instability:

Pure Neutron Matter will appear

Ruling out the “supersoft” E_{sym} means the pure
neutron matter cannot appear inside NStars



Nskin in ^{208}Pb : PREX-II

PHYSICAL REVIEW LETTERS 126, 172502 (2021)

Editors' Suggestion

Featured in Physics

Accurate Determination of the Neutron Skin Thickness of ^{208}Pb through Parity-Violation in Electron Scattering

PREX:

The Lead (Pb) Radius EXperiment

D. Adhikari,¹ H. Albataineh,² D. Androic,³ K. Aniol,⁴ D. S. Armstrong,⁵ T. Averett,⁵ C. Ayerbe Gayoso,⁵ S. Barcus,⁶

$$R_n - R_p = 0.283 \pm 0.071 \text{ fm.}$$

PREX-II: $R_{\text{skin}} = [0.212, 0.354] \text{ fm}$ for Pb208

Huge Nskin!

In PREX, the neutron density distribution in ^{208}Pb is determined by measuring the parity-violating electroweak asymmetry in the elastic scattering of polarized electrons off ^{208}Pb and thus is free from the strong interaction uncertainties

PHYSICAL REVIEW LETTERS 126, 172503 (2021)

Editors' Suggestion

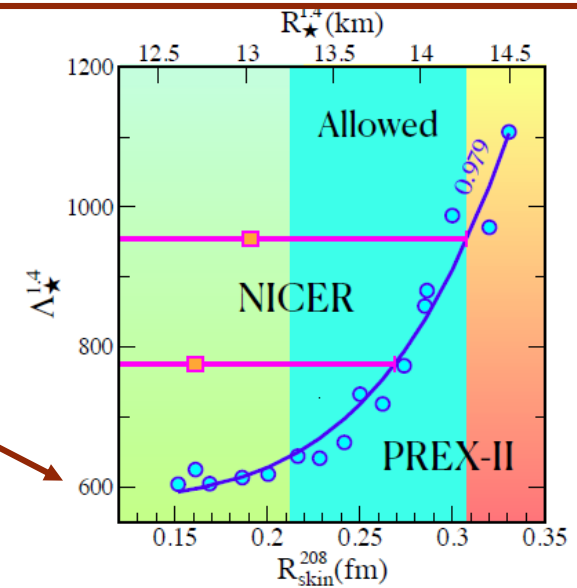
Implications of PREX-2 on the Equation of State of Neutron-Rich Matter

Brendan T. Reed^{1,2,*} F. J. Fattoyev^{3,†} C. J. Horowitz^{2,‡} and J. Piekarewicz^{4,§}

RMF: Inconsistent with GW170817: $\Lambda_{1.4} < 580$!!!

δ meson may be important in the RMF model !!!

Fan Li (李帆)/Cai/Zhou/Jiang/LWC, ApJ 929, 183 (2022)

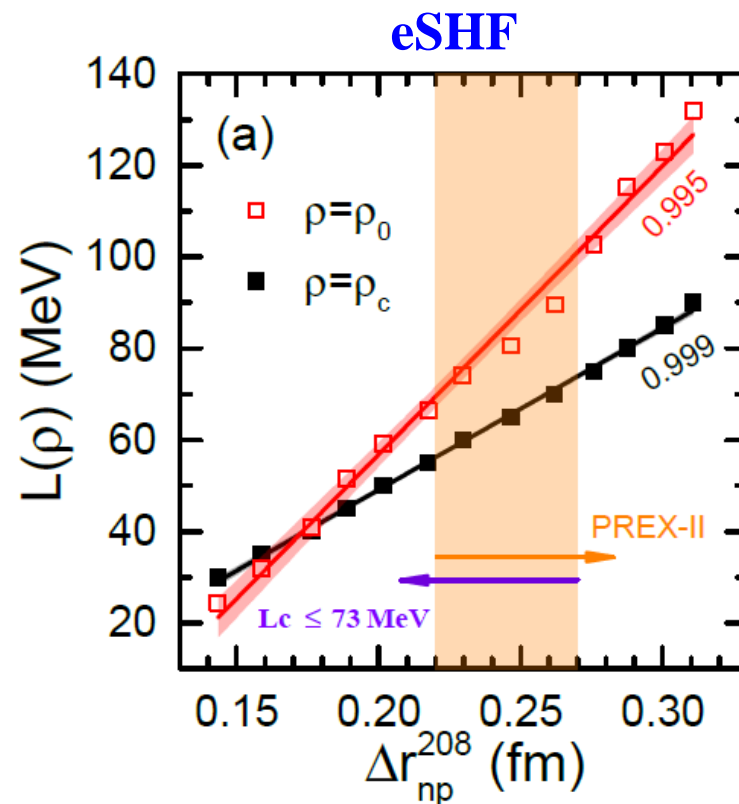
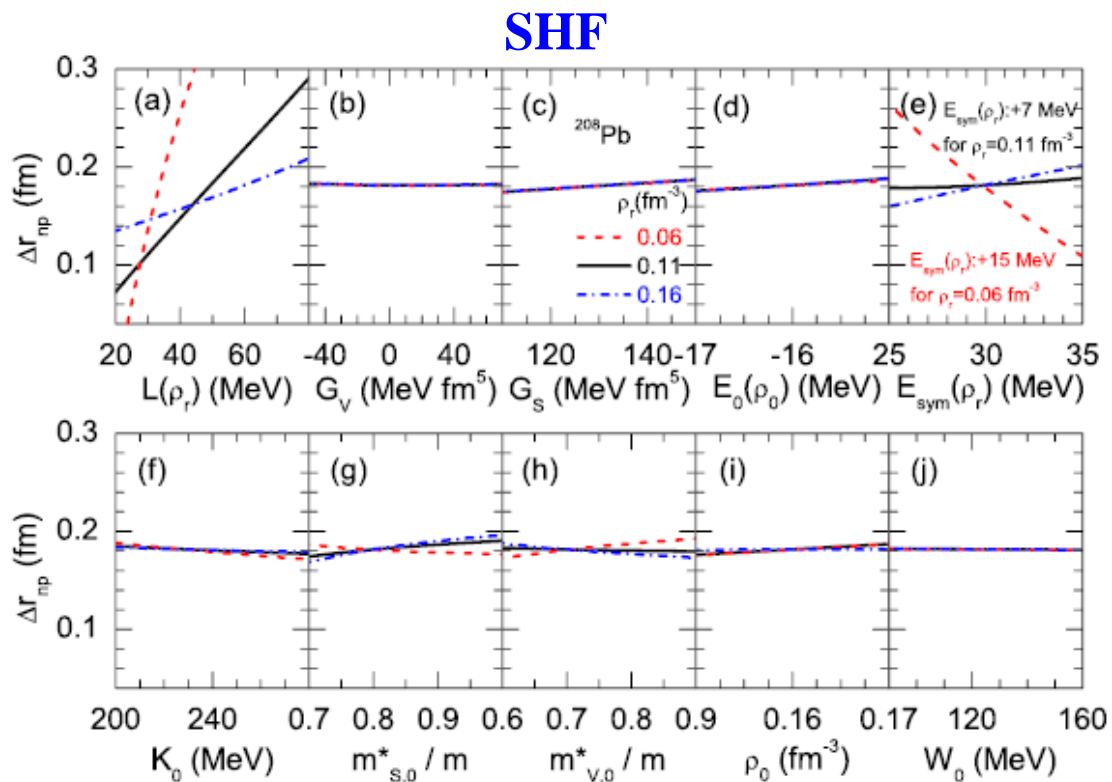




Nskin in ^{208}Pb : PREX-II

What determines the Nskin thickness? ---The density slope parameter L , especially L_c at $\rho_c \sim 2\rho_0/3$ (average density of nuclei)!!!

B. A. Brown, PRL85, 5296 (2000); R. J. Furnstahl, NPA706, 85 (2002); L.W. Chen et al., PRC72, 064309 (2005); M. Centelles et al., PRL102, 122502 (2009); L. W. Chen et al., PRC 82, 024321 (2010); X. Roca-Maza et al.,PRL106, 252501 (2011),



Nskin is determined by L_c !!!

Z. Zhang (张振) and LWC*, PLB726, 234 (2013) (Citations: 146+)

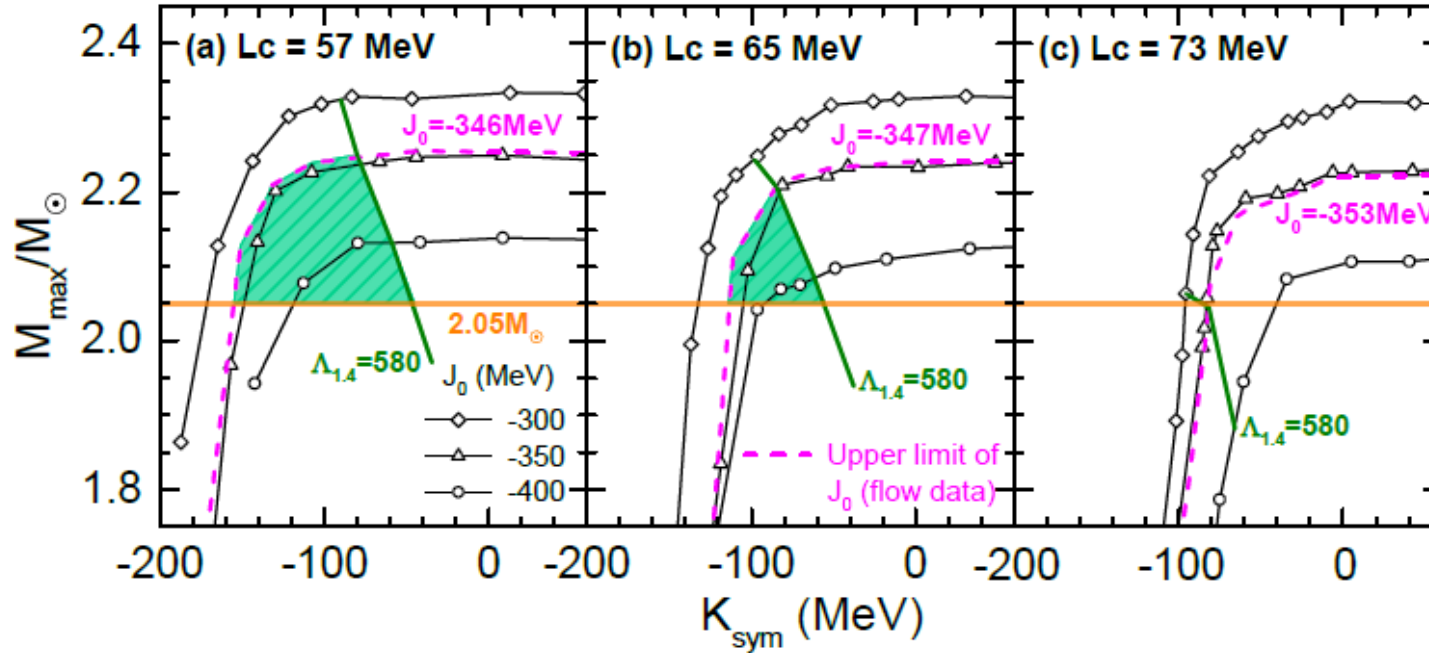
T.G. Yue (岳伺钢), LWC*, Z. Zhang (张振), and Y. Zhou (周颖), arXiv:2102.05267



Nskin in ^{208}Pb : PREX-II

Implications of Rskin from PREX-II

T.G. Yue (岳伺钢), LWC*, Z. Zhang (张振), and Y. Zhou (周颖), arXiv:2102.05267



L_c cannot be too big!!! $L_c < 73$ MeV and then set an upper limit on Rskin: < 0.27 fm for Pb208

eSHF provides a single unified framework to simultaneously describe the finite nuclei (Eb, Rc, GMR, Nskin-PREX-II) + Flow data in HIC+NStar (e.g., NICER)+GW170817

$$E_{\text{sym}}(\rho_0) = 34.5 \pm 1.5 \text{ MeV and } L = 85.5 \pm 22.2 \text{ MeV}$$



Nskin in ^{208}Pb : PREX-II

THE ASTROPHYSICAL JOURNAL, 929:183 (7pp), 2022 April 20

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<https://doi.org/10.3847/1538-4357/ac5e2a>



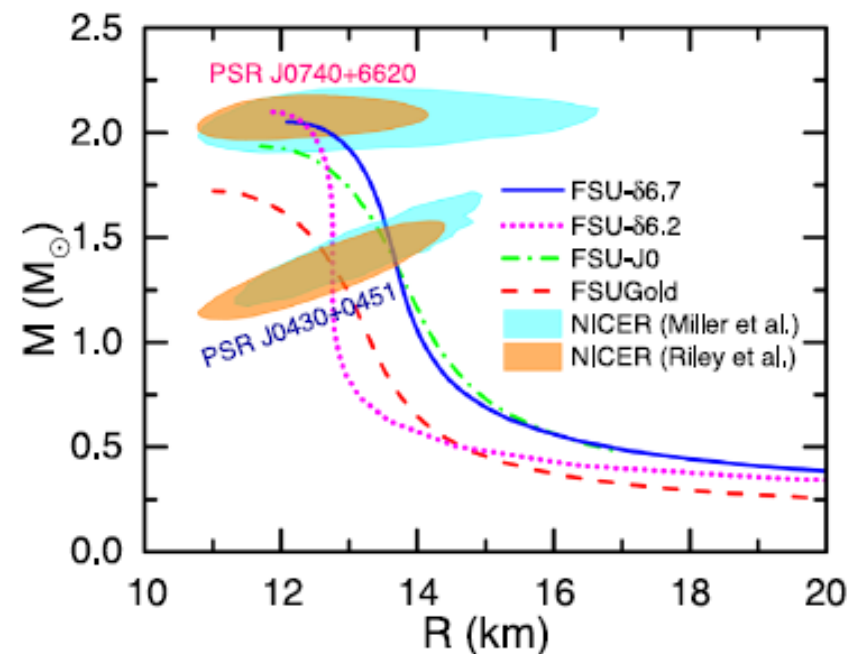
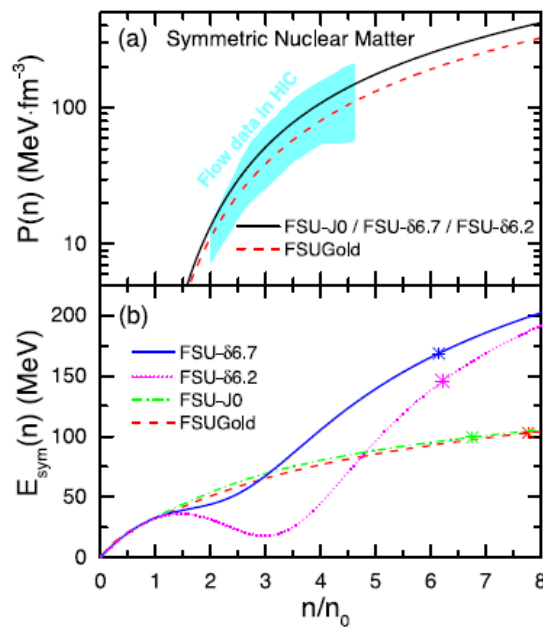
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Effects of Isoscalar- and Iovector-scalar Meson Mixing on Neutron Star Structure

Fan Li¹, Bao-Jun Cai², Ying Zhou², Wei-Zhou Jiang³, and Lie-Wen Chen^{1,4}

Fan Li (李帆), B.J. Cai, Y. Zhou, W.Z. Jiang, and LWC, ApJ 929, 183 (2022)

$$\begin{aligned} \mathcal{L} = & \bar{\psi}(i\partial_\mu\gamma^\mu - m)\psi \\ & + g_\sigma\sigma\bar{\psi}\psi - g_\omega\omega_\mu\bar{\psi}\gamma^\mu\psi - g_\rho\rho_\mu\bar{\psi}\gamma^\mu\tau\psi + g_\delta\delta\bar{\psi}\tau\psi \\ & + \frac{1}{2}(\partial_\mu\sigma\partial^\mu\sigma - m_\sigma^2\sigma^2) - \frac{1}{3}b_\sigma m(g_\sigma\sigma)^3 - \frac{1}{4}c_\sigma(g_\sigma\sigma)^4 \\ & - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} + \frac{1}{2}m_\omega^2\omega_\mu\omega^\mu + \frac{1}{4}c_\omega(g_\omega^2\omega_\mu\omega^\mu)^2 \\ & - \frac{1}{4}\rho_{\mu\nu}\rho^{\mu\nu} + \frac{1}{2}m_\rho^2\rho_\mu\rho^\mu + \frac{1}{2}\Lambda_V(g_\rho^2\rho_\mu\rho^\mu)(g_\omega^2\omega_\mu\omega^\mu) \\ & + \frac{1}{2}(\partial_\mu\delta\partial^\mu\delta - m_\delta^2\delta^2) + \frac{1}{2}C_{\delta\sigma}(g_\sigma^2\sigma^2)(g_\delta^2\delta^2), \end{aligned}$$



δ meson and its coupling to σ meson may play an important role in the RMF model !!!

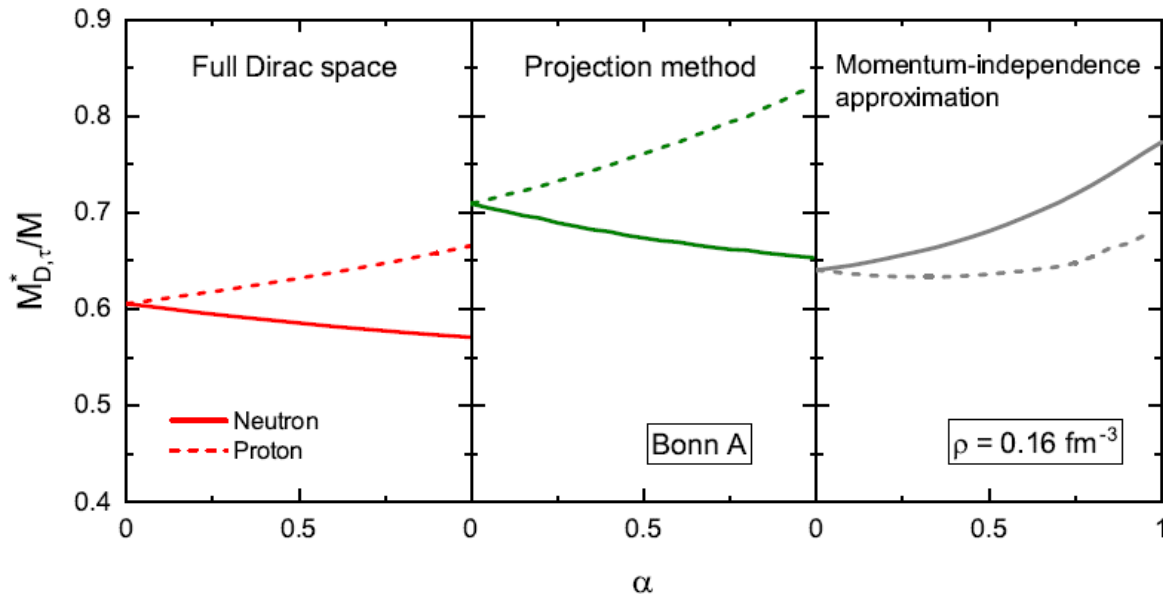
FSUGold + δ - σ mixing provides a single unified framework to simultaneously describe the finite nuclei (Eb, Rc, GMR, Nskin-PREX-II) + Flow data in HIC+NStar (e.g., NICER)+GW170817



Nskin in ^{208}Pb : PREX-II

Asymmetric Nuclear Matter and Neutron Star Properties in Relativistic *ab initio* Theory in the Full Dirac Space

Sibo Wang,¹ Hui Tong,^{2,3} Qiang Zhao,⁴ Chencan Wang,⁵ Peter Ring,⁶ and Jie Meng^{7,8,*}



S. Wang, H. Tong, Q. Zhao, C. Wang, P. Ring, and J. Meng, arXiv:2203.05397

Table 2
Nuclear Matter Bulk Parameters, NS Properties, and Ground-state Properties of ^{208}Pb and ^{48}Ca in FSUGold, FSU-J0, FSU- $\delta 6.7$, and FSU- $\delta 6.2$

Quantity	FSUGold	FSU-J0	FSU- $\delta 6.7$	FSU- $\delta 6.2$
$n_0(\text{fm}^{-3})$	0.148	0.148	0.148	0.148
$e_0(\text{MeV})$	-16.31	-16.31	-16.31	-16.31
$K_0(\text{MeV})$	229.2	229.2	229.2	229.2
$J_0(\text{MeV})$	-521.6	-322.0	-322.0	-322.0
m_{Dirac}^*	0.610	0.610	0.610	0.610
$m_{\text{Dirac},n,0.5}^*$	0.610	0.610	0.576	0.578
$m_{\text{Dirac},p,0.5}^*$	0.610	0.610	0.643	0.640

Fan Li (李帆), B.J. Cai, Y. Zhou, W.Z. Jiang, and LWC, ApJ 929, 183 (2022)

The isospin splitting of the nucleon Dirac mass in FSUGold + δ - σ mixing framework is in surprisingly good agreement with the microscopic RBHF calculations in the Full Dirac Space !!!



Outline

- Nuclear matter EOS and the symmetry energy (E_{sym})
- Dense nuclear matter from *Nuclear experiments + Neutron star mass and radius + Tidal deformability from GW170817*
- Summary and outlook



Summary and Outlook

- 通过重离子碰撞，原子核结构性质(mass, neutron skin, GR/PG...), 以及核子光学势的研究，我们已经对亚饱和密度以及饱和密度附近核物质对称能有比较好的认识：

- 亚饱和密度区的核物质对称能 – 比较精确
- 饱和密度附近的核物质对称能：

World Average: $E_{\text{sym}}(\rho_0) = 31.7 \pm 3.1 \text{ MeV}$ and $L = 57.5 \pm 24.5 \text{ MeV}$

But REX-II data make the situation quite elusive!!!

更精确的约束需要更精确的实验数据和更可靠的理论方法!

- GW170817给出了对称能高密行为的上限，2.14倍太阳质量的存在排除了“超软”的高密对称能，意味着中子星内部不存在纯中子物质
- 决定核物质状态方程的高密行为依然是一个巨大的，丰中子核引起的重离子碰撞的实验数据以及中子星/引力波观测数据将变得非常重要!
- 单个(相对论和非相对论)能量密度泛函能成功描述核实验数据和天文学观测数据(高密核物质中超子自由度、介子凝聚、夸克自由度???)



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Fan Li, Xu-Run Huang, Si-Pei Wang, Zhen Cao, Tong-Gang Yue, Ji-Min Bai (SJTU, Shanghai)

Funding:

National Natural Science Foundation of China

Major State Basic Research Development Program (973 Program) in China

Shanghai Rising-Stars Program, Shanghai “Shu-Guang” Project, Shanghai “Eastern Scholar”

Science and Technology Commission of Shanghai Municipality



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谢谢!
Thanks!

