



# Constraints on heavy flavour PDFs from LHC data

#### Juan Rojo

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Opportunities with Heavy Flavour at the EIC Workshop Center for Frontiers in Nuclear Science, 04/11/2020

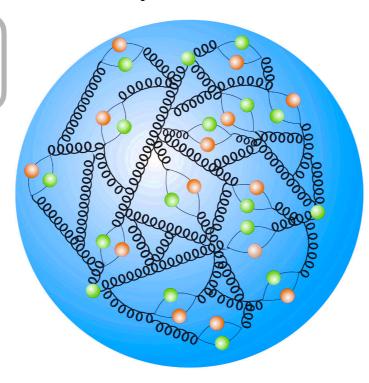


#### a QCD bound state of quarks and gluons

 $m_p \simeq 1 \text{ GeV}$ 

valence quarks: up & down

 $m_{u,d} \simeq \text{few MeV}$ 



gluons

$$m_g = 0$$

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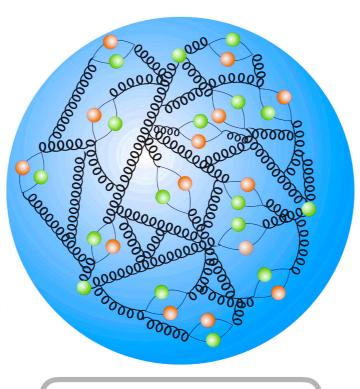
 $m_{u,d} \simeq 200 \text{ MeV}$ 

charm quarks?

 $m_{u,d} \simeq 1.5 \text{ GeV}$ 

a QCD bound state of quarks and gluons

valence quarks: up & down



gluons

strange quarks

bottom quarks?

charm quarks?

**leptons** 

photons

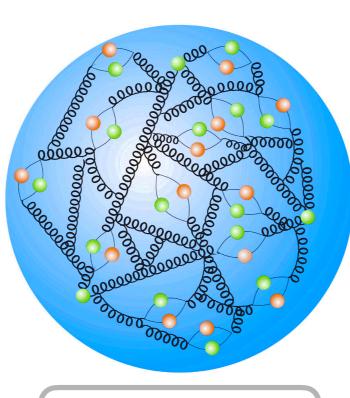
a QCD bound state of quarks and gluons

valence quarks: up & down

Higgs bosons?

strange quarks

*leptons* 



bottom quarks?

top quarks?

gluons

weak bosons?

charm quarks?

photons

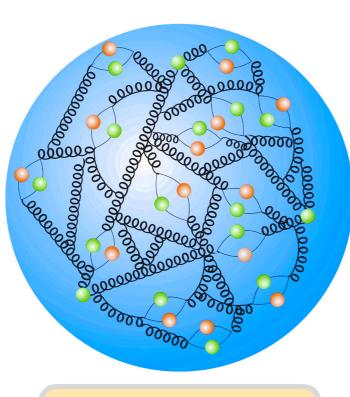
"Heavy" flavour PDFs:

valence quarks: up & down

Higgs bosons?

strange quarks

*leptons* 



gluons

weak bosons?

charm quarks?

bottom quarks?

top quarks?

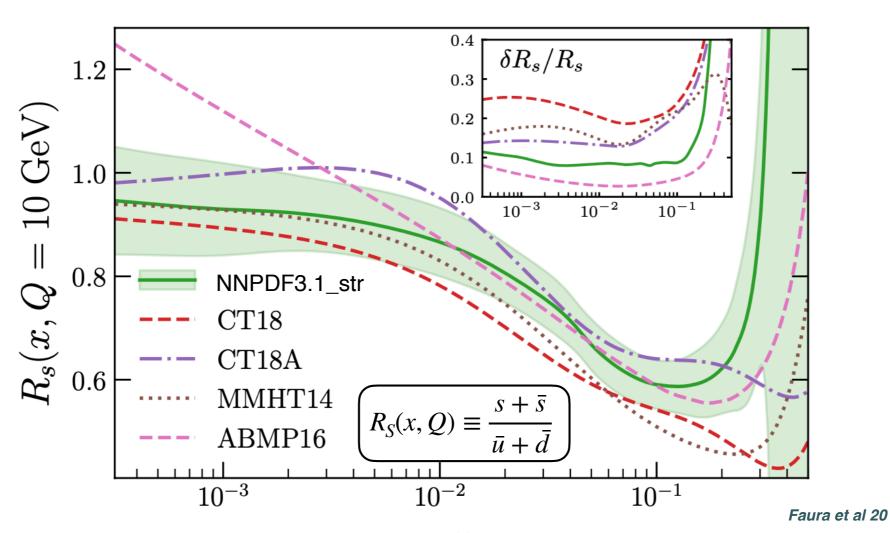
photons

# The strange PDF

See Faura, Iranipour, Nocera, JR, Ubiali 20 for all references

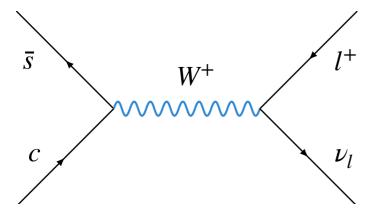
# How strange is the proton?

marked differences in the strange PDFs from **recent global analyses**, both for central values and for the size of its uncertainties



the strange PDF can be constrained by different processes, both collider and fixed-target

#### inclusive Drell-Yan

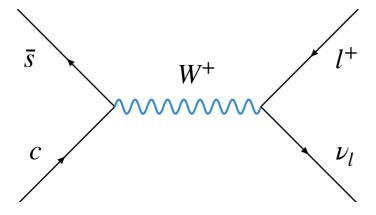


sensitivity via correlation of *W*+, *W*+, *Z* distributions

ATLAS 7 TeV data: preference for **R**<sub>S</sub>=1

the strange PDF can be constrained by different processes, both collider and fixed-target

inclusive Drell-Yan



sensitivity via correlation of *W*+, *W*+, *Z* distributions

ATLAS 7 TeV data: preference for **R**<sub>S</sub>=1

W+charm  $S \qquad W^ \bar{c}$ 

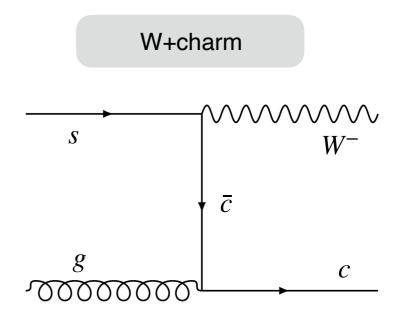
direct sensitivity, but larger uncertainties

NNLO QCD corrections recently calculated

Czakon et al 20

the strange PDF can be constrained by different processes, both collider and fixed-target

inclusive Drell-Yan <u>5</u>  $\boldsymbol{\mathcal{C}}$  $u_l$ Neutrino DIS  $u_{\mu}$ p(A)



Characteristic opposite-sign dimuon signature

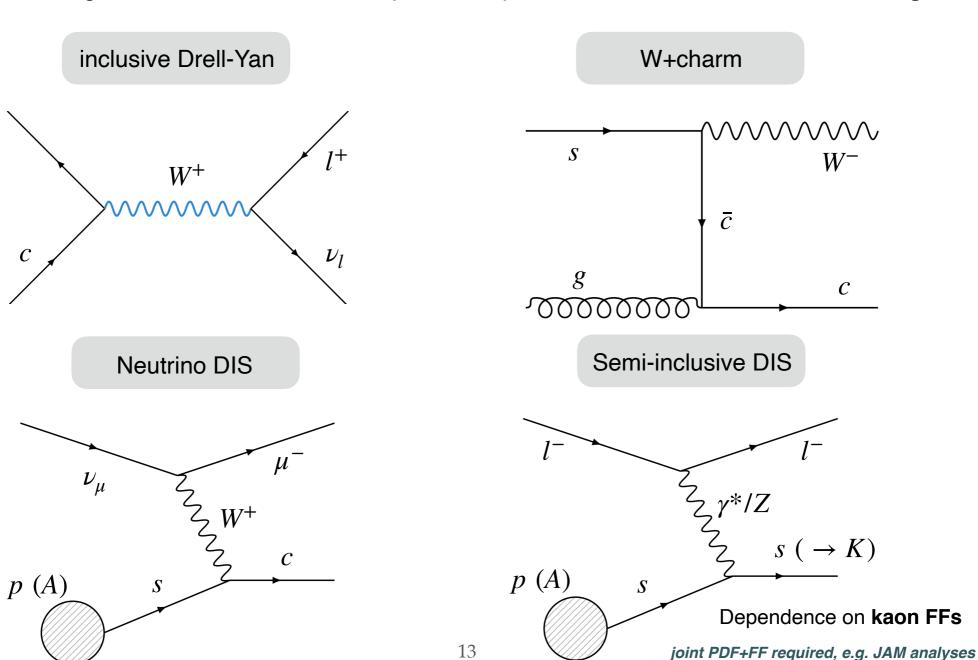
NNLO massive QCD corrections recently calculated

Gao et al 17

Traditionally associated with  $R_S \approx 0.5$ 

Theoretical issues w. nuclear PDFs and charm fragmentation

the strange PDF can be constrained by different processes, both collider and fixed-target



### How strange is the proton?

Why different groups obtain different strange PDF determinations?

- Different input datasets, e.g. including or not NuTeV/NOMAD, or ATLAS W,Z 2011
- Different theory settings, e.g. massive NNLO QCD corrections to neutrino DIS
- Methodological settings, e.g. treatment of charm PDF, uncertainty estimates

Do we really **understand** the proton strangeness?

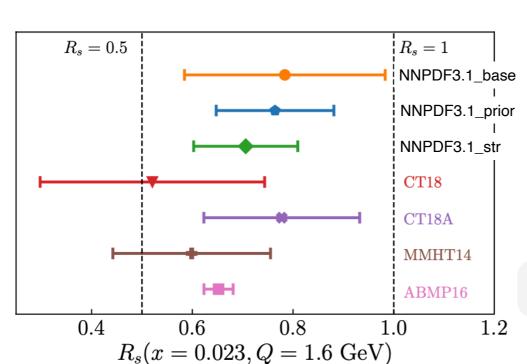
- Is there a tension between collider and neutrino DIS data?
- Is there a tension between inclusive *W,Z* and *W+charm* data?
- Are constraints from different neutrino data consistent among them?

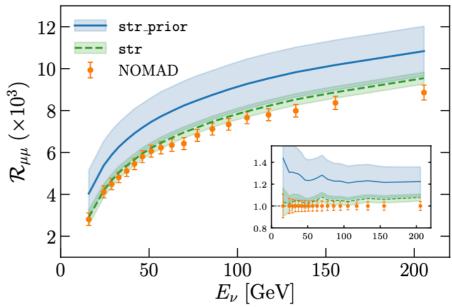
Is the proton strangeness suppressed or symmetric wrt the light quark sea?

### How strange is the proton?

#### Reappraisal of the proton strangeness based combination of all relevant experimental inputs

Dataset	$n_{ m dat}$	$\chi^2_{\mathrm{base}}$	$\chi^2_{ m pr}$	$\chi^2_{ m str}$
	76/76/95/91/95	0.76	0.71	0.53
NuTeV [9]	76/76/76/76	0.76	0.71	0.53
NOMAD [10]	//19/15/19	[9.3]	[8.8]	0.55
	391/418/418/418/418	1.45	1.40	1.40
ATLAS [12]	34/61/61/61/61	1.96	1.65	1.67
	/37/37/37	[0.73]	0.68	0.60
CMS $[17, 18]$	/15/15/15/15	[1.04]	0.98	0.96
ATLAS [16]	$-\!\!-\!\!/22/22/22$	[0.52]	0.48	0.42
ATLAS [15]	/32/32/32	[1.58]	1.18	1.18
	3981/4077/4096/4092/4096	1.18	1.17	1.17
	NuTeV [9] NOMAD [10]  ATLAS [12]  CMS [17, 18] ATLAS [16]	76/76/95/91/95  NuTeV [9] 76/76/76/76  NOMAD [10] —/—/19/15/19  391/418/418/418/418  ATLAS [12] 34/61/61/61  —/37/37/37/37  CMS [17,18] —/15/15/15/15  ATLAS [16] —/22/22/22/22  ATLAS [15] —/32/32/32/32	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$





- Satisfactory simultaneous description of all datasets
- ✓ No evidence for tension between datasets or groups of processes
- Sizeable constraints from **NOMAD neutrino DIS**data, consistent with collider data

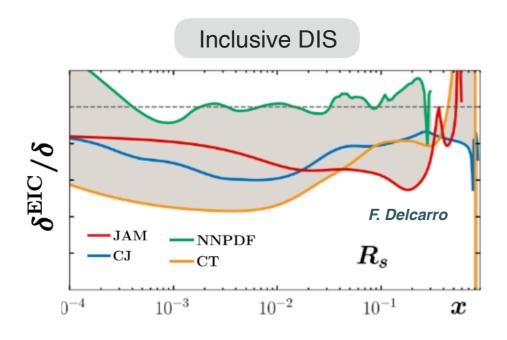
  see also ABM studies
- Strong preference for a moderately suppressed strangeness

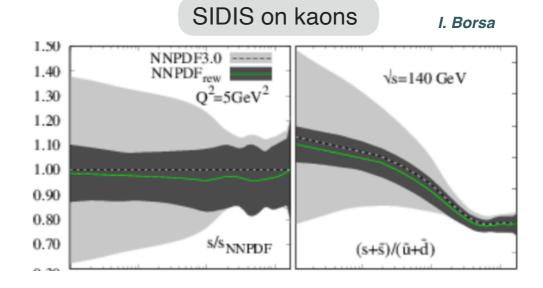
$$R_S(x = 0.023, Q = 1.6 \text{ GeV}) = 0.71 \pm 0.10$$

Faura et al 20

### Strangeness @ EIC

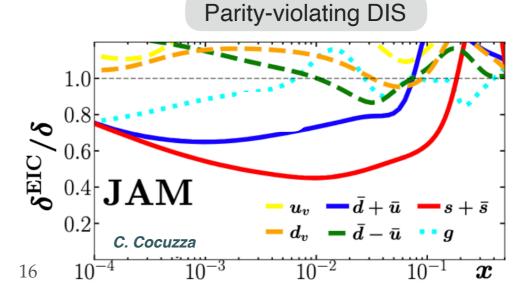
Can **future EIC measurements** provide information about the proton strangeness?





- ✓ Plenty of exciting opportunities to constrain strangeness at the EIC
- ☑ The combination of inclusive DIS, SIDIS, and PVDIS will make possible a stringent determination of the strange content of protons

from EIC Yellow Report, in preparation

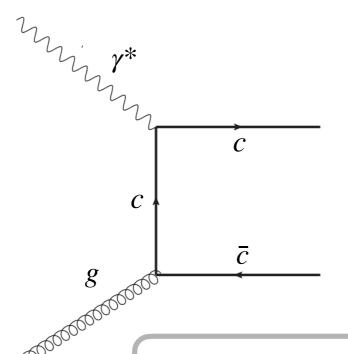


# The charm PDF

### A tale of three schemes

say you want to evaluate the **charm DIS structure function**. You have three options

Fixed-flavor scheme: no charm PDF, charm mass effects accounted for exactly



$$F_2^c(x,Q^2) \propto \sum_{i=g,u,d,s} C_i^{(n_f)}(\alpha_s,Q^2/m_c^2) \otimes f_i^{(n_f)}$$

exact in threshold region

not appropriate to describe  $Q^2 >> m_c^2$  region due to large unresummed logs

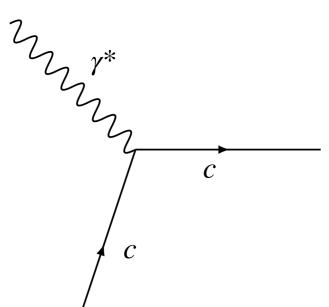
in the massive FFN scheme, nothing to say about the charm PDF

see also talks by Robert, Sven, Achim ...

### A tale of three schemes

say you want to evaluate the **charm DIS structure function**. You have three options

**Zero-mass scheme:** charm PDF treated on the same footing as all other quark flavours



$$F_2^c(x, Q^2) \propto \sum_{i=g,u,d,s,c} C_i^{(n_f+1)}(\alpha_s) \otimes f_i^{(n_f+1)}$$

exact in far from threshold region

not appropriate to describe  $Q^2 \approx m_c^2$  region due to lack of massive corrections

in this scheme, the charm PDF above threshold is constructed from the nf=3 PDFs via **matching**. If there is no charm PDF for  $n_f$ =3 (assumption) then

$$f_c^{(n_f+1)} \propto \alpha_s \ln \frac{Q^2}{m_c^2} \left( P_{qg} \otimes f_g^{(n_f+1)} \right) + \mathcal{O}\left(\alpha_s^2\right)$$

the charm PDF is deterministically generated from the gluon (and light quark) PDFs

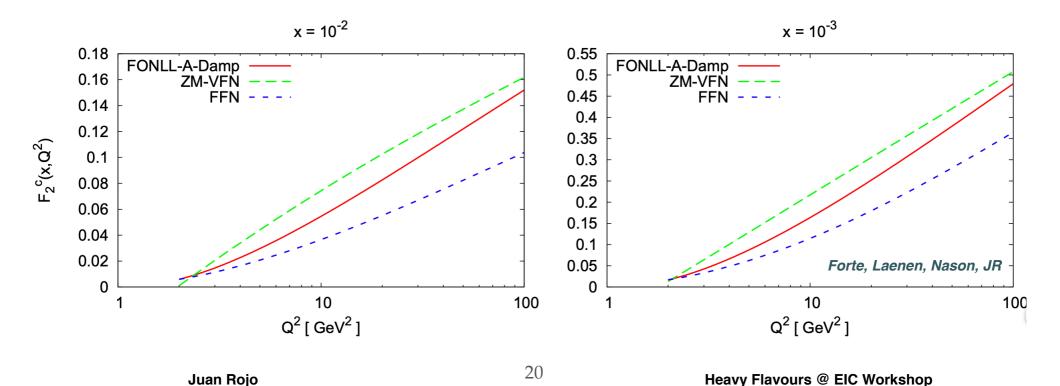
### A tale of three schemes

say you want to evaluate the **charm DIS structure function**. You have three options

General-mass VFN scheme: charm PDF treated on the same footing as all other quark flavours, massive effects included in coefficient functions

$$F_2^c(x, Q^2) \propto \sum_{i=g,u,d,s,c} C_i^{(GM)}(\alpha_s, Q^2/m_c^2) \otimes f_i^{(n_f+1)}$$

Systematically improvable, reliable for all values of  $Q^2$  from threshold to collider scales



### perturbative, intrinsic, and fitted charm

Let's work in the following in the GM-VFN scheme. The charm PDF above threshold is constructed from the  $n_f$ =3 PDFs via **matching** in three possible ways:

Perturbative charm: the charm PDF vanishes below threshold, then above threshold ( $\mu_c \approx m_c$ ) the charm PDF is deterministically generated from the gluon (and light quark) PDFs

$$f_c^{(n_f)} = 0 \quad \rightarrow \quad f_c^{(n_f+1)} \propto \alpha_s \ln \frac{Q^2}{m_c^2} \left( P_{qg} \otimes f_g^{(n_f+1)} \right) + \mathcal{O} \left( \alpha_s^2 \right) \quad \text{not much interesting to salphout the charm PDF here}$$

Intrinsic charm: a model for the charm PDF at the initial evolution scale (below or at threshold) is assumed. Then the charm PDF is this intrinsic component plus the perturbative component

$$f_c^{(n_f)}(x, Q_0) = Ax^2 \left[ 6x(1+x)\ln x + (1-x)(1+10x+x^2) \right]$$

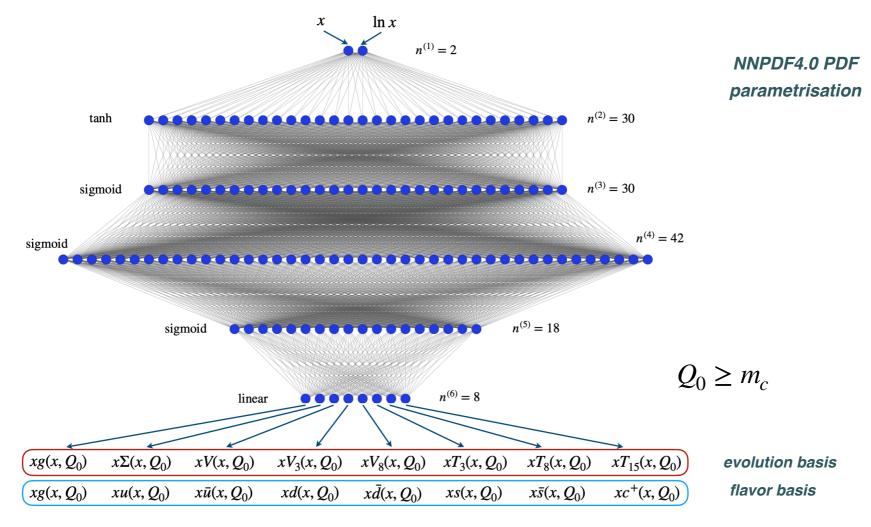
BHPS model (scale independent)

the model parameters (e.g. normalisation) are extracted from comparison with data

Fitted charm: no assumptions on possible intrinsic component are made. The charm is parametrised above threshold in exactly the same was as all other quark PDFs

$$f_c^{(n_f+1)}(x,Q_0)=x^{-\alpha_c}(1-x)^{\beta_c}\mathrm{NN}(x)$$
 NNPDF approach

### perturbative, intrinsic, and fitted charm

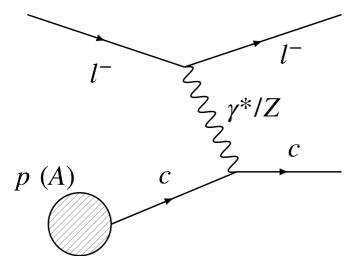


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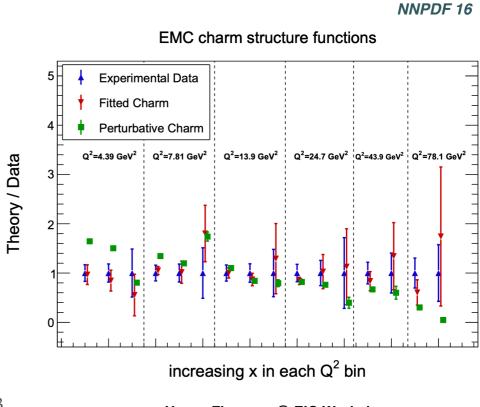
$$f_c^{(n_f+1)}(x,Q_0) = x^{-\alpha_c}(1-x)^{\beta_c} \mathrm{NN}(x)$$
 NNPDF approach

in the following, we assume that the charm PDF is either **fitted** or **intrinsic** (perturbative charm cannot be constrained, since it is generated by DGLAP evolution)

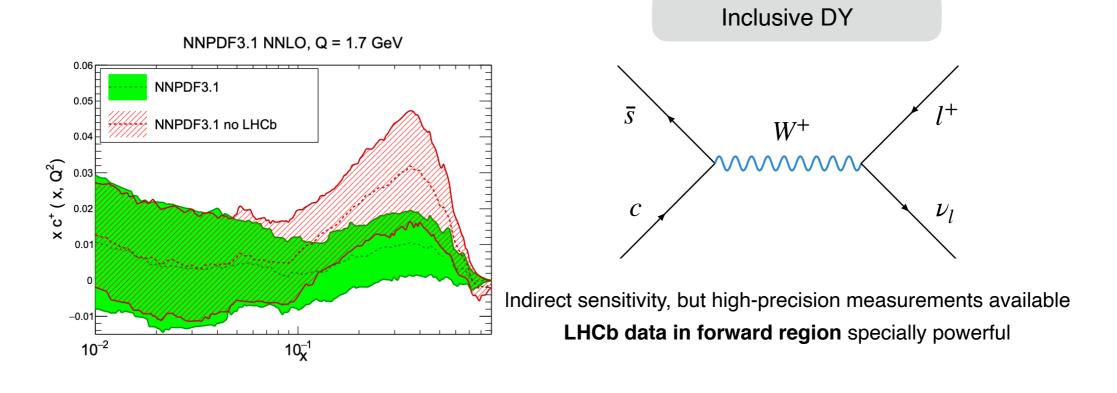
#### Charm production in DIS



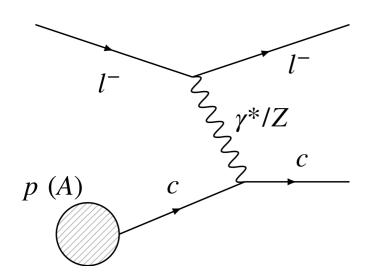
Only measurements of  $F_2^c$  at large-x are sensitive EMC data from the 80s available ...



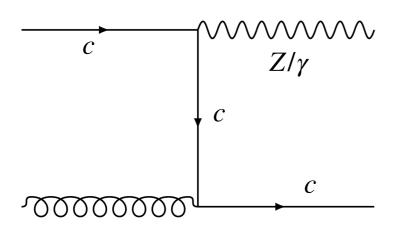
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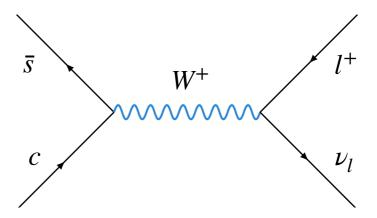
#### Charm production in DIS



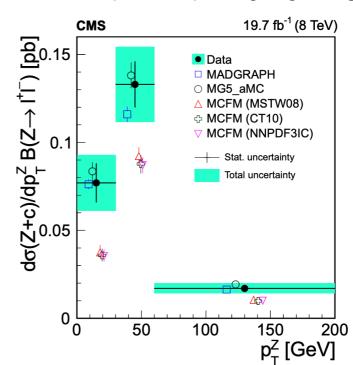
#### Z/γ+charm



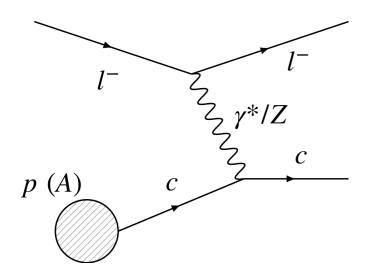
#### Inclusive DY



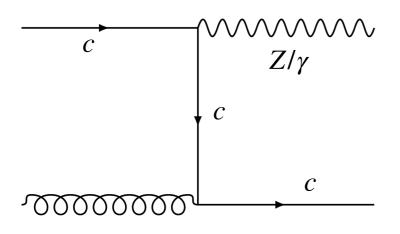
Direct sensitivity, but requires going to **high-p**<sub>T</sub>



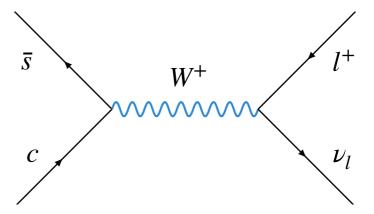
#### Charm production in DIS



#### Z/y+charm

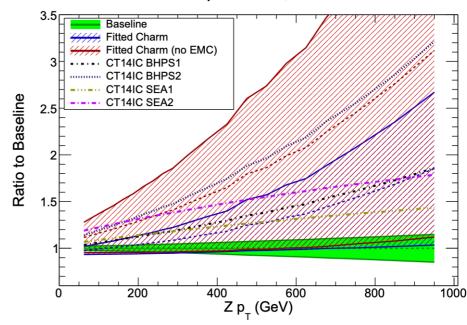


#### Inclusive DY

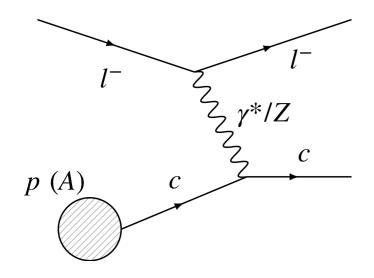


Direct sensitivity, but requires going to **high-** $p_T$ 

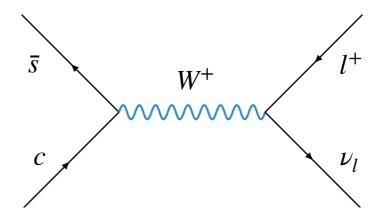
Z+Charm production, LHC 13 TeV



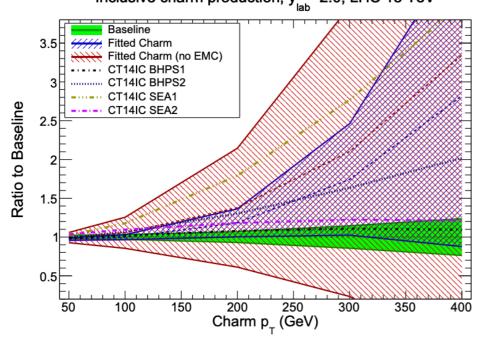
#### Charm production in DIS



#### Inclusive DY

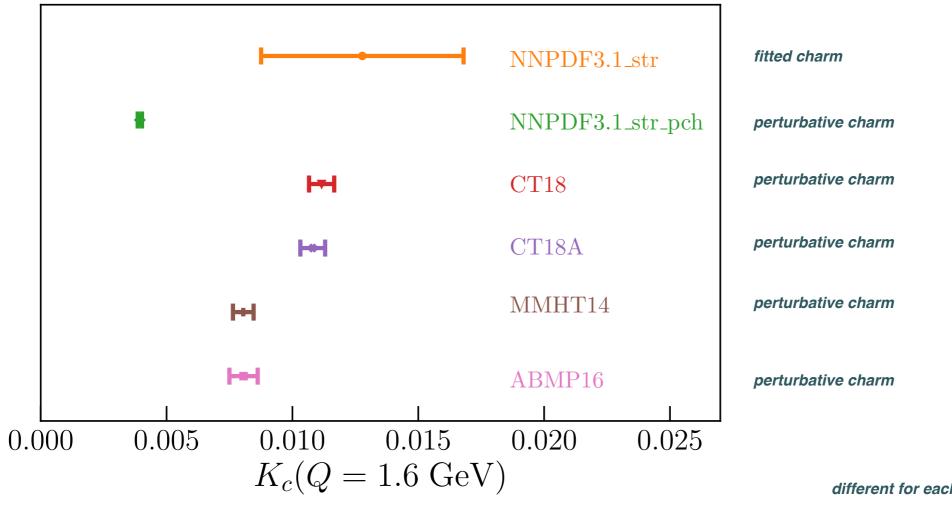


#### Inclusive charm production, $y_{lab}$ =2.0, LHC 13 TeV



#### Open charm production

also here going to **high-** $p_T$  and/or **large rapidities** key to optimise sensitivity



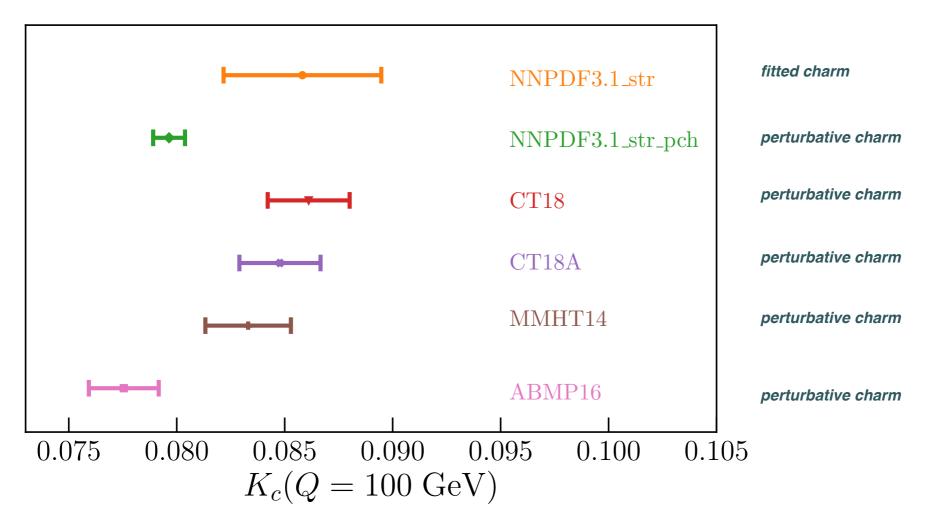
different for each group

$$K_c(Q) \equiv \frac{\int_0^1 dx \ (c + \bar{c})}{\int_0^1 dx \ (u^+ + d^+ + s^+)}$$

momentum fraction carried by charm in units of that of the light quarks

perturbative charm results are sensitive to choice of value of charm mass

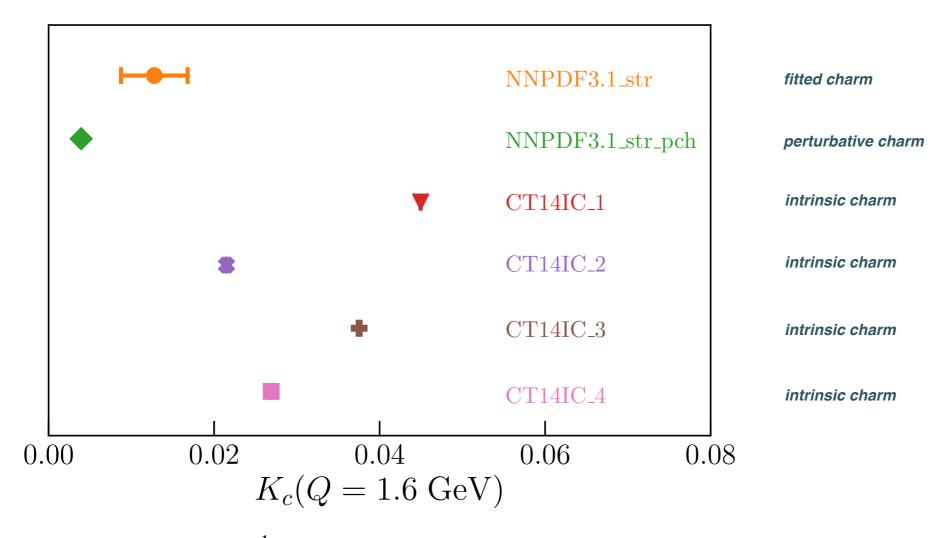
current data favour a non-zero charm component in the n<sub>f</sub>=3 scheme



$$K_c(Q) \equiv \frac{\int_0^1 dx \ (c + \bar{c})}{\int_0^1 dx \ (u^+ + d^+ + s^+)}$$

momentum fraction carried by charm in units of that of the light quarks

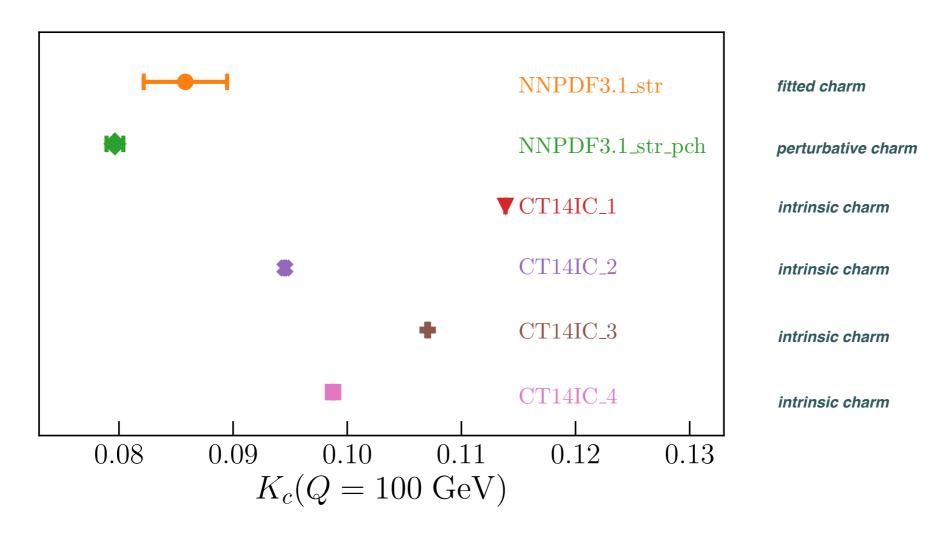
- perturbative charm results are sensitive to choice of value of charm mass
- component in the n<sub>f</sub>=3 scheme



$$K_c(Q) \equiv \frac{\int_0^1 dx \ (c + \bar{c})}{\int_0^1 dx \ (u^+ + d^+ + s^+)}$$

momentum fraction carried by charm in units of that of the light quarks

- ☑ The most extreme IC models from CT14 are in marked tension with the NNPDF3.1 fitted charm
- ✓ (Some of these models are extreme on purpose)



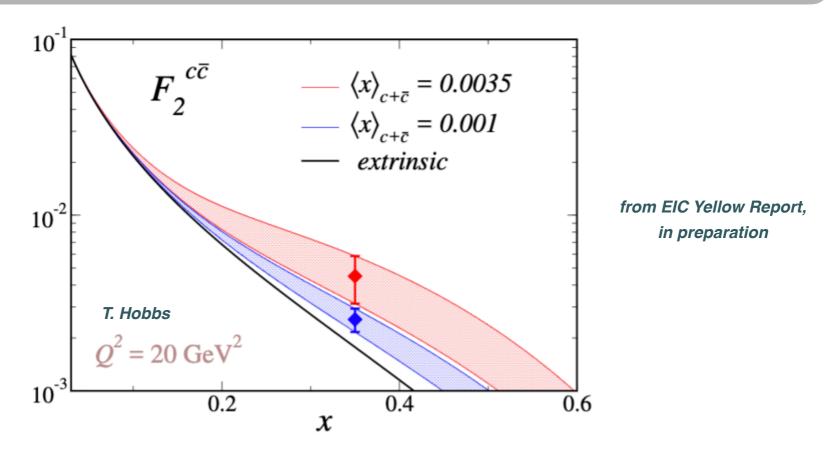
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momentum fraction carried by charm in units of that of the light quarks

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### Charm @ EIC

Can **future EIC measurements** provide information about the proton charmness?



- Plenty of exciting opportunities to constrain charmness at the EIC
- ☑ In particular, the measurement of the **charm structure function at large-***x* provides a very clean probe of the charm PDF and its various possible modifications

# The bottom PDF

### Intrinsic bottom?

Our considerations about the charm PDF (perturbative vs fitted vs intrinsic) apply equally well to the bottom PDF, though here on expects deviations from the perturbative picture to be quite suppressed

In **all PDF analysis to date**, the bottom PDF is always **generated dynamically via DGLAP** from light quarks and gluons

$$f_b^{(n_f+1)} \propto \alpha_s \ln \frac{Q^2}{m_b^2} \left( P_{qg} \otimes f_g^{(n_f+1)} \right) + \mathcal{O}\left(\alpha_s^2\right)$$

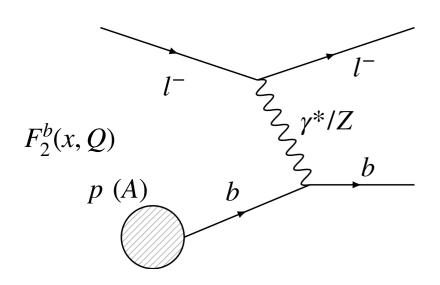
Assume that there is a "non-perturbative" component to the bottom PDF.

How we could constrain it?

same as for strange and charm: look for processes directly or indirectly sensitive to **bottom quarks in** the initial state of the reaction

## Constraining bottom

#### Bottom production in DIS

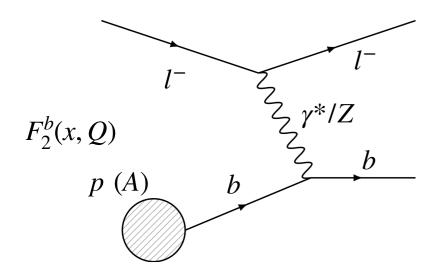


much smaller rates than for charm.....

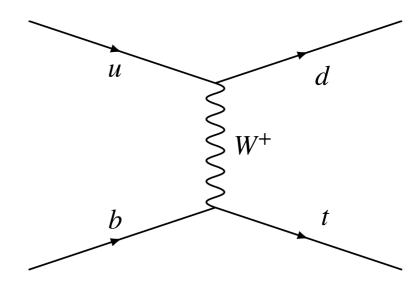
e.g. at HERA charm can be up to 25% while bottom is 1%

### Constraining bottom

#### Bottom production in DIS



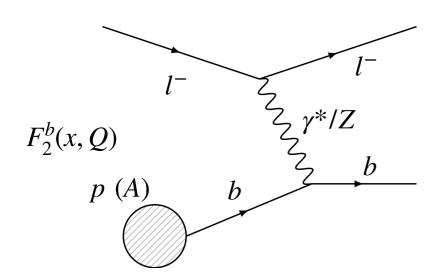
#### Single top production



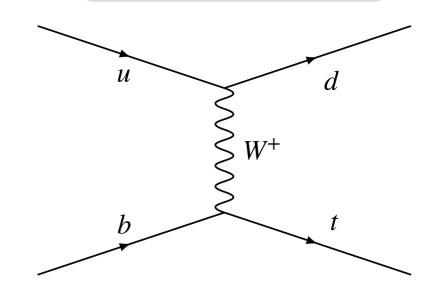
requires **matched scheme** to account for bottom quark mass corrections

# Constraining bottom

#### Bottom production in DIS



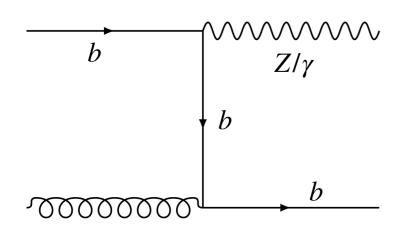
Single top production

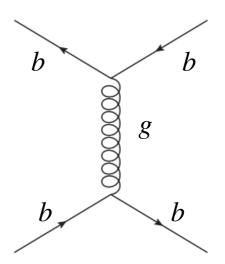


Z/y+bottom

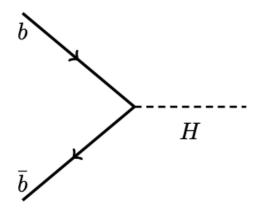
so far no one has interpreted these processes assuming an "intrinsic" bottom PDF ....

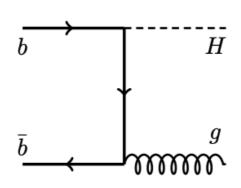
Open charm production

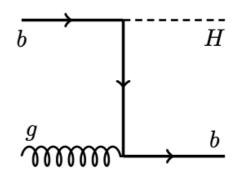




### Impact of parametrised bottom



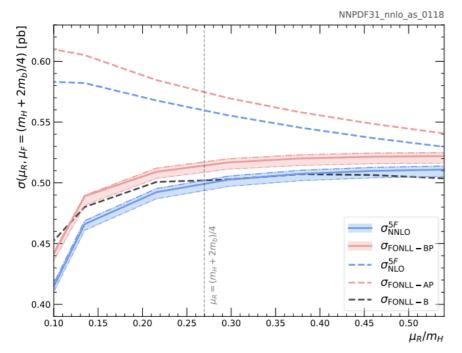




It has been recently shown, for the case of **Higgs production in bottom fusion**, how a matched scheme for processes involving initial-state bottom quarks simplifies the calculation provided a **fitted** 

#### bottom PDF is introduced

Forte, Giani, Napoletano 19



Irrespective of whether or not the proton contains a non-perturbative bottom component, fitting the bottom PDF might be advantageous for precision phenomenology

### Summary and outlook

The accurate determination of **heavy flavour PDFs** is an essential ingredient for **LHC phenomenology** and will play a key role for the **EIC physics program** 

- ☑ A global interpretation of strangeness-sensitive processes reveals a nice overall consistency
  and the lack of any significant tension
- ☑ In particular, the NOMAD neutrino DIS data imposes stringent constraints, without worsening description of collider data
- ☑ Going further requires accounting for SIDIS data: simultaneous PDF+FF determination
- Much progress in theoretical and phenomenological understanding of the **charm PDF**, with several measurements providing direct information (but still open issues)
- Some preference for a non-perturbative charm component, but most extreme IC models disfavoured by LHC data
- Fitting the bottom PDFs has never been attempted: could be beneficial for precision LHC phenomenology!

# Summary and outlook

The accurate determination of **heavy flavour PDFs** is an essential ingredient for **LHC phenomenology** and will play a key role for the **EIC physics program** 

The Electron Ion Collider will open a new frontier for heavy flavour PDF studies!

