Regions of DIS and Semi-Inclusive DIS

Jefferson Lab/Old Dominion University

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Variables and Observables

$$E'\frac{\mathrm{d}\sigma}{\mathrm{d}^3\mathbf{l}'} = \frac{2\,\alpha_{\mathrm{em}}^2}{(s-M^2)\,Q^4}\,L_{\mu\nu}W_{\mathrm{tot}}^{\mu\nu} \qquad \qquad \frac{\mathrm{d}\sigma}{\mathrm{d}x_{\mathrm{Bj}}\,\mathrm{d}y\,\mathrm{d}\psi} = \frac{\alpha_{\mathrm{em}}^2y}{Q^4}L_{\mu\nu}W_{\mathrm{tot}}^{\mu\nu}$$

$$W_{\rm tot}^{\mu\nu}(P,q) \equiv 4\pi^3 \sum_X \delta^{(4)}(P+q-P_X) \langle P, S|j^{\mu}(0)|X\rangle \langle X|j^{\nu}(0)|P,S\rangle$$

$$4P_{\rm B}{}^{0}E'\frac{{\rm d}\sigma}{{\rm d}^{3}{\bf l}'~{\rm d}^{3}{\bf P}_{\rm B}} = \frac{2\,\alpha_{\rm em}^{2}}{(s-M^{2})Q^{4}}\,L_{\mu\nu}W_{\rm SIDIS}^{\mu\nu} \qquad \frac{{\rm d}\sigma}{{\rm d}x_{\rm Bj}~{\rm d}y~{\rm d}\psi~{\rm d}z_{\rm N}~{\rm d}^{2}\boldsymbol{P}_{\rm B,T}} = \frac{\alpha_{\rm em}^{2}y}{4Q^{4}z_{\rm N}}\,L_{\mu\nu}W_{\rm SIDIS}^{\mu\nu}$$

$$W_{\text{SIDIS}}^{\mu\nu}(P,q,P_{\text{B}}) \equiv \sum_{X} \delta^{(4)}(P+q-P_{\text{B}}-P_{X}) \langle P,S|j^{\mu}(0)|P_{\text{B}},X\rangle \langle P_{\text{B}},X|j^{\nu}(0)|P,S\rangle$$

$$\sum_{B} \int \frac{\mathrm{d}^{2} \boldsymbol{P}_{\mathrm{B,T}} \, \mathrm{d} z_{\mathrm{N}}}{4 z_{\mathrm{N}}} W_{\mathrm{SIDIS}}^{\mu\nu} = \langle N \rangle W_{\mathrm{tot}}^{\mu\nu}$$

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Questions

• Factorization is based on mass/Q expansion.

Different factorization theorems for different regions of x,z,Q,P_T, etc.

• What are the boundaries to different regions in reality?





Subregions



Large and Small Transverse Momentum



SIDIS

• NLO study near the valence region. B. Wang, J. O. Gonzalez-Hernandez, TR, N. Sato <u>https://jeffersonlab.qithub.io/</u> <u>BigTMD/_build/html/index.ht</u> <u>ml</u>



Kinematics And Masses

• Final state momentum fraction

$$x_{\rm N} = -\frac{q^{+}}{P^{+}} = \frac{2x_{\rm Bj}}{1 + \sqrt{1 + \frac{4x_{\rm Bj}^{2}M^{2}}{Q^{2}}}}$$

$$z_{\rm N} = \frac{P_{\rm B}^{-}}{q^{-}}$$

$$z_{\rm N} = \frac{x_{\rm N}z_{\rm h}}{2x_{\rm Bj}} \left(1 + \sqrt{1 - \frac{4M^{2}M_{\rm B,T}^{2}x_{\rm Bj}^{2}}{Q^{4}z_{\rm h}^{2}}}\right)$$

$$= z_{\rm h} \left(1 - \frac{x_{\rm Bj}^{2}M^{2}}{Q^{2}} \left(1 + \frac{P_{\rm B,T}^{2}}{z_{\rm h}^{2}Q^{2}}\right) + \left(\frac{x_{\rm Bj}^{2}M^{2}}{Q^{2}}\right)^{2} \left(\frac{P_{\rm B,T}^{2}}{z_{\rm h}^{2}Q^{2}} - \frac{P_{\rm B,T}^{4}}{z_{\rm h}^{4}Q^{4}} + 2 - \frac{M_{\rm B}^{2}}{z_{\rm h}^{2}M^{2}x_{\rm Bj}^{2}}\right) + O\left(\left(\left(\frac{x_{\rm Bj}^{2}M^{2}}{Q^{2}}\right)^{3}\right)\right)$$

How low in Q?



How low in Q?



FIG. 10: (Color online) The ^{1,2}H(e,e' π^{\pm})X cross sections at fixed values of x = 0.40 and z = 0.55, as a function of Q^2 . The solid curves are the simple quark-parton model calculations following a high-energy factorized description. Solid symbols are data after events from diffractive ρ production are subtracted (see text). $P_t^2 < 0.2 ~ (\text{GeV}/c)^2$

$$\frac{\mathrm{d}\sigma}{\mathrm{d}x_{\mathrm{Bj}} \mathrm{d}y \mathrm{d}\psi \mathrm{d}z_{\mathrm{N}} \mathrm{d}^{2}\boldsymbol{P}_{\mathrm{B,T}}} = \frac{\alpha_{\mathrm{em}}^{2}y}{4Q^{4}z_{\mathrm{N}}} L_{\mu\nu}W_{\mathrm{SIDIS}}^{\mu\nu}$$

Conclusion

- The question of the where the between different factorizationbased and partonic pictures apply is an open (empirical) question.
- The question can nevertheless be systematized...