

Run 13 direct photon ALL

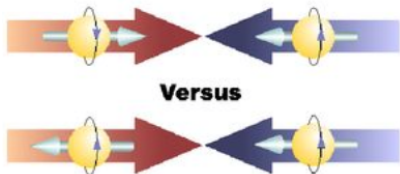
Zhongling Ji

Stony Brook University

September 10, 2020

Direct photon ALL

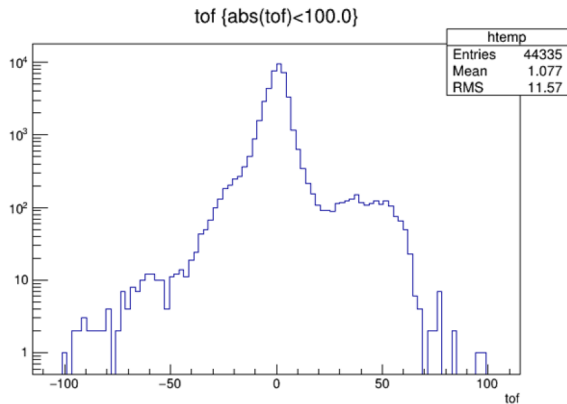
Yellow (Y) and Blue (B)



$$A_{LL} = \frac{\Delta\sigma}{\sigma} = \frac{\sigma_{++} + \sigma_{--} - \sigma_{+-} - \sigma_{-+}}{\sigma_{++} + \sigma_{--} + \sigma_{+-} + \sigma_{-+}}$$
$$= \frac{1}{P_B P_Y} \frac{N_{++} - RN_{+-}}{N_{++} + RN_{+-}}$$

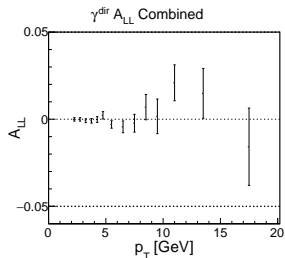
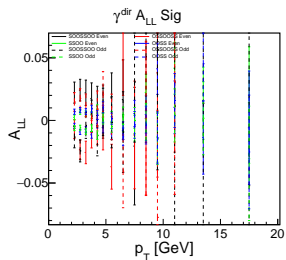
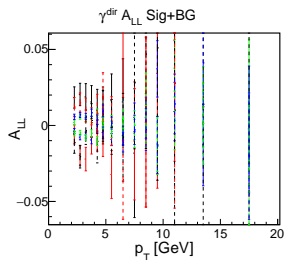
- ▶ Different runs have different polarization P_B and P_Y , so measured in a run-by-run basis.
- ▶ Even and odd crossings have different electric circuits, so measured separately.
- ▶ There are also four spin patterns, so total eight groups.
- ▶ For particles in isolation cone:
 - ▶ Loose cut: ToF < 50 ns, E > 0.15 GeV
 - ▶ Tight cut: ToF < 10 ns, E > 0.5 GeV

ToF distributions



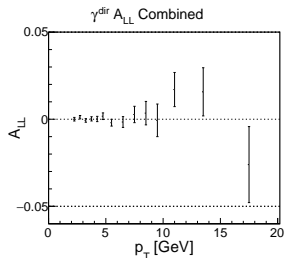
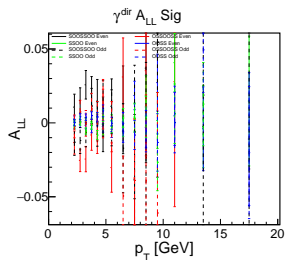
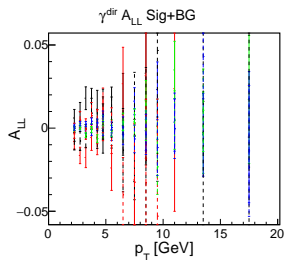
From John and Milap's jet ALL AN

Discrepancies in loose cuts



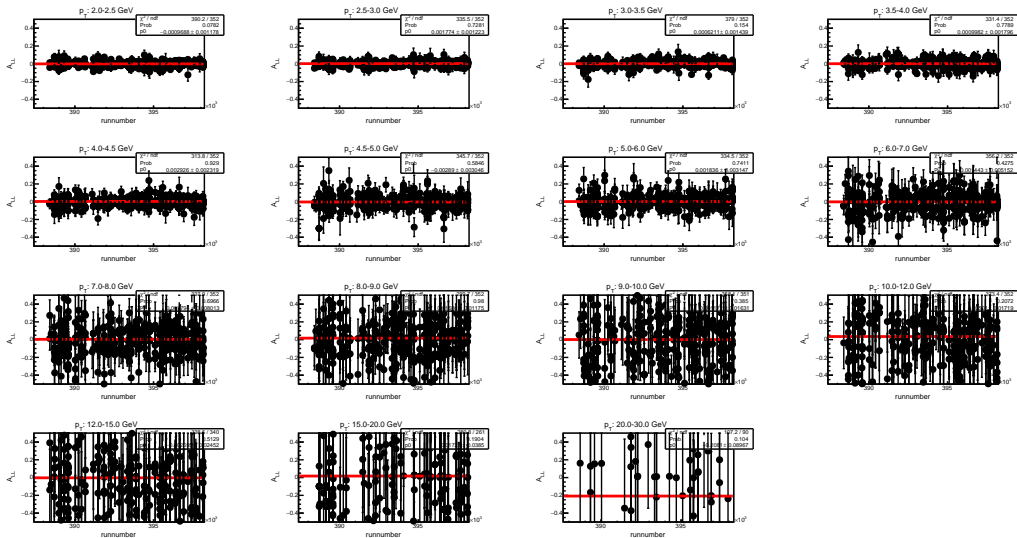
p_T [GeV]	F	p
2-2.5	5.901	8.709e-07
2.5-3	12.4	1.554e-15
3-3.5	8.116	9.932e-10
3.5-4	5.853	1.009e-06
4-4.5	3.334	0.00157
4.5-5	1.223	0.2864
5-6	1.704	0.104
6-7	1.256	0.2686
7-8	0.7887	0.5968
8-9	0.5234	0.8174
9-10	1.299	0.2471
10-12	0.697	0.6747
12-15	0.9198	0.4899
15-20	0.6332	0.7287

Better agreement in tight cuts

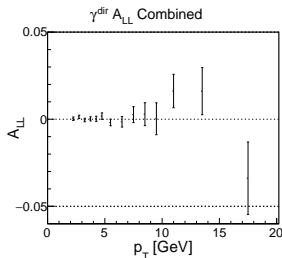
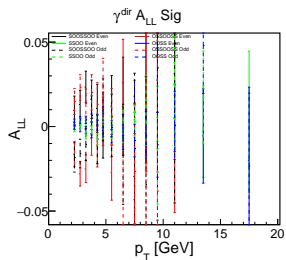
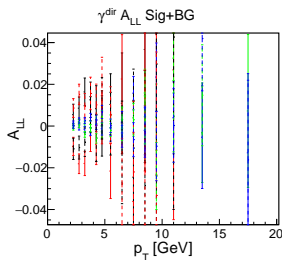


p_T [GeV]	F	p
2-2.5	1.254	0.2697
2.5-3	3.282	0.001815
3-3.5	3.933	0.0002903
3.5-4	1.777	0.08792
4-4.5	1.724	0.09921
4.5-5	1.125	0.3442
5-6	0.8223	0.5686
6-7	0.9596	0.4593
7-8	1.065	0.3837
8-9	0.4716	0.8556
9-10	0.9886	0.4377
10-12	0.9743	0.4482
12-15	0.6964	0.6752
15-20	0.8313	0.5613

Run-by-run fitting

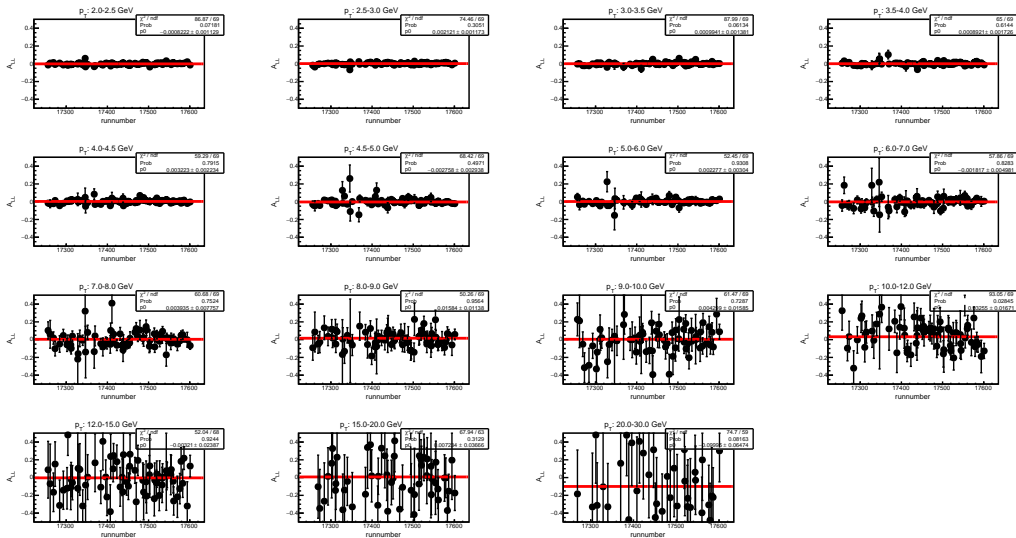


Cross check in fill-by-fill analysis



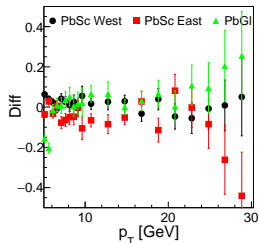
p_T [GeV]	F	p
2-2.5	1.824	0.08209
2.5-3	4.494	8.618e-05
3-3.5	3.533	0.001139
3.5-4	2.027	0.05136
4-4.5	1.883	0.07174
4.5-5	1.326	0.2373
5-6	0.8033	0.5851
6-7	0.6962	0.6753
7-8	1.029	0.4103
8-9	0.5332	0.8092
9-10	0.7743	0.6092
10-12	0.8879	0.5162
12-15	0.8807	0.5219
15-20	1.345	0.2288
20-30	0.6968	0.6748

Fill-by-fill fitting

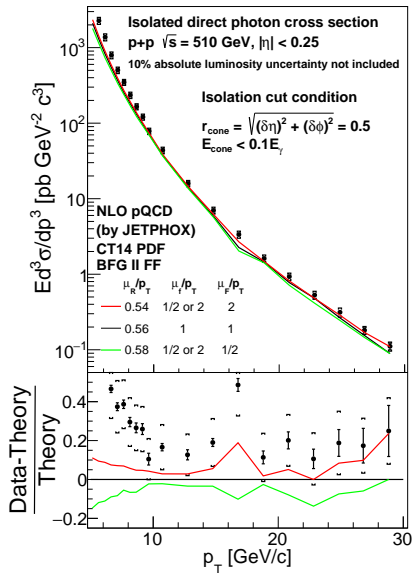
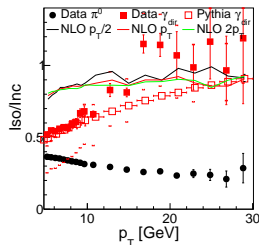


Cross section in tight cuts

Diff in parts



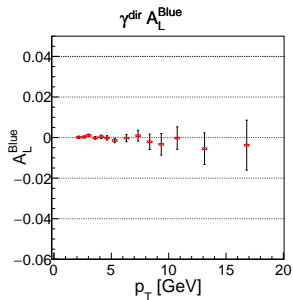
Isolated/Inclusive ratio



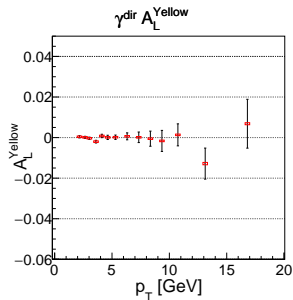
Systematic uncertainties

- ▶ False asymmetry in background due to ghost cluster: low p_T
- ▶ Uncertainty of relative luminosity: $3.853e-4$
- ▶ Global scaling uncertainty from polarization: 6.6%
- ▶ Uncertainty of background fraction estimation
- ▶ Uncertainty from eta background

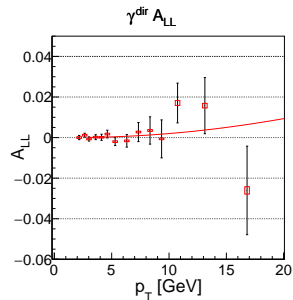
AL(L) with systematic uncertainties



Blue beam AL

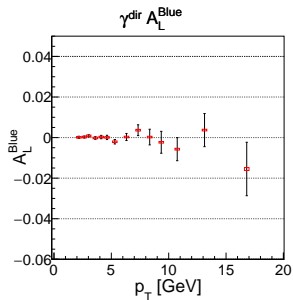


Yellow beam AL

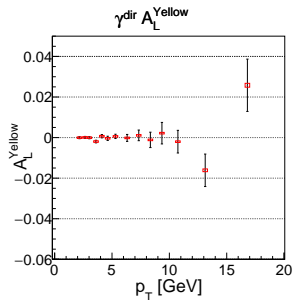


ALL

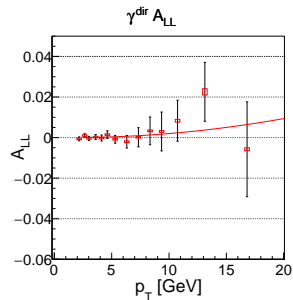
AL(L) with Inseok's calibration



Blue beam AL

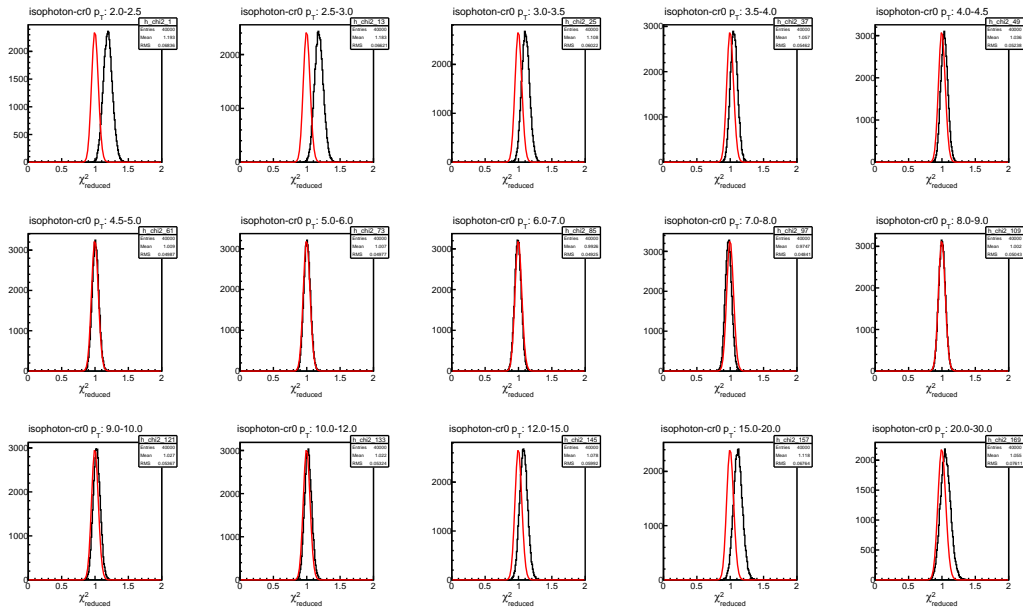


Yellow beam AL

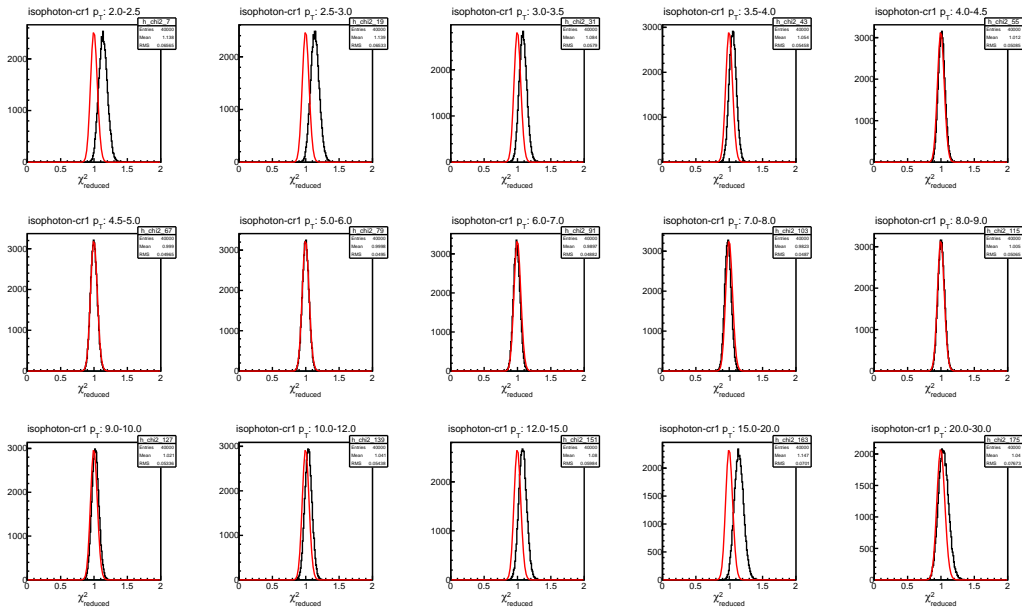


ALL

Bunch shuffling: even crossings



Bunch shuffling: odd crossings



Conclusions for ALL

- ▶ Discrepancies from spin patterns come from previous ghost events and can be removed by using tight cuts.
- ▶ Agreement with partonic NLO calculations in isolated cross section needs loose cut.
- ▶ Will use loose cut for cross section and tight cut for ALL.

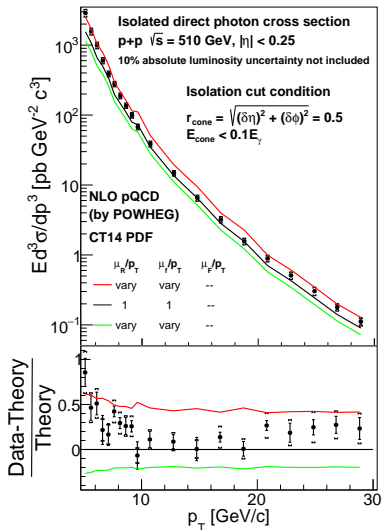
$$\frac{d\sigma_{\text{ME}}}{dx_1 dx_2} \sim \left| \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \end{array} \right|^2$$

The diagram shows two Feynman diagrams for the Mott cross-section. In the first diagram, an incoming electron (blue line) from the left emits a photon (red wavy line) and then scatters. In the second diagram, an incoming electron (blue line) from the left scatters and then emits a photon (red wavy line). A green wavy line represents the incoming photon. The two diagrams are summed and then squared to give the Mott cross-section.

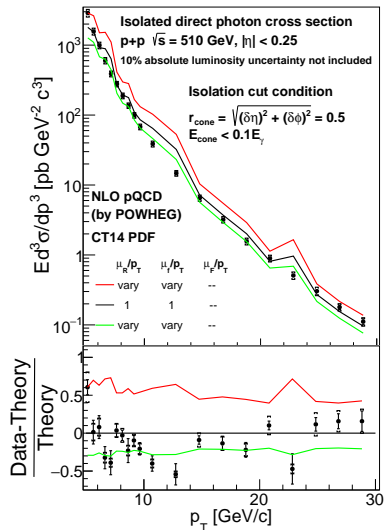
$$\frac{d\sigma_{\text{PS}}}{dx_1 dx_2} \sim \left| \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \end{array} \right|^2 + \left| \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \end{array} \right|^2$$

The diagram shows two Feynman diagrams for the plane-wave Born approximation (PS) cross-section. The first diagram is identical to the first diagram in the Mott cross-section. The second diagram is identical to the second diagram in the Mott cross-section. The two diagrams are squared individually and then summed to give the PS cross-section.

POWHEG with different vetoes for isolated direct photon

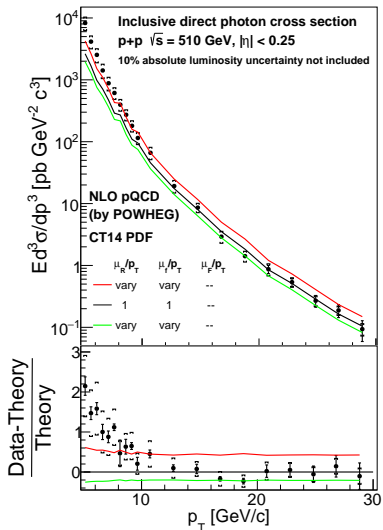


With MPI (default veto)

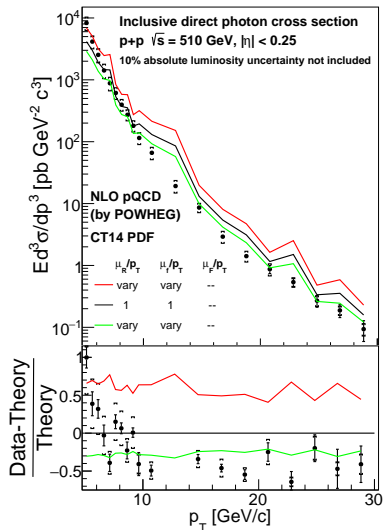


With MPI (QED-QCD veto)

POWHEG with different vetoes for inclusive direct photon

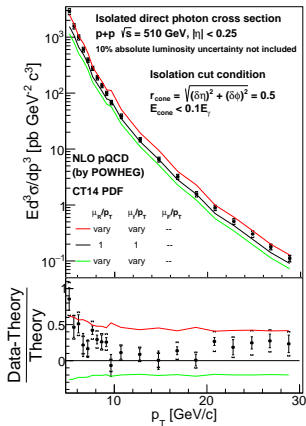


With MPI (default veto)

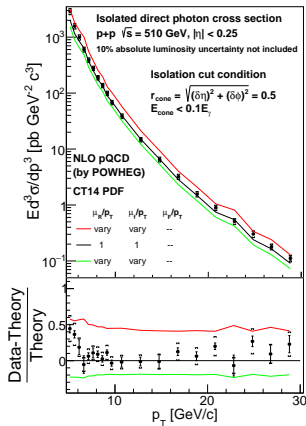


With MPI (QED-QCD veto)

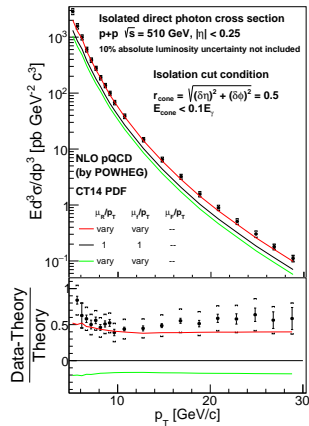
POWHEG w/o MPI for isolated direct photon



With MPI

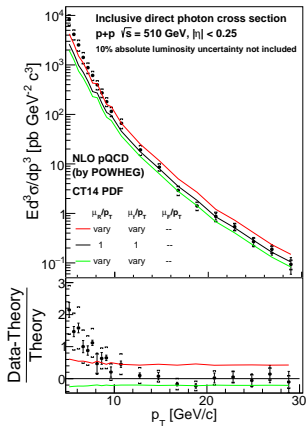


Without MPI

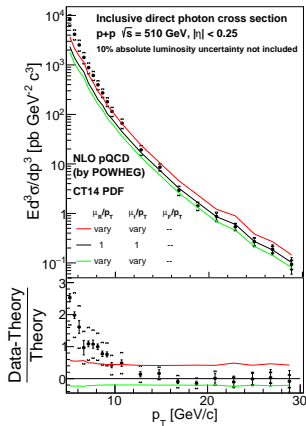


Pure hard processes

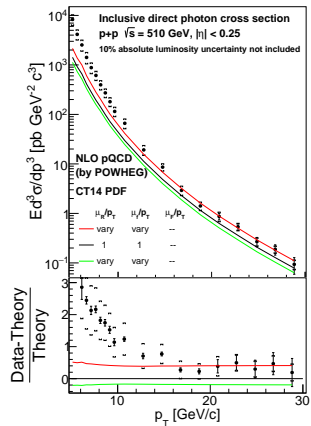
POWHEG w/o MPI for inclusive direct photon



With MPI

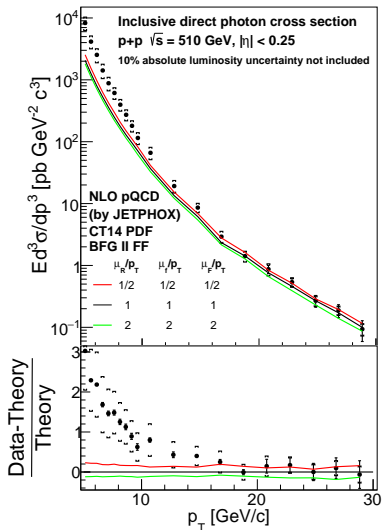


Without MPI

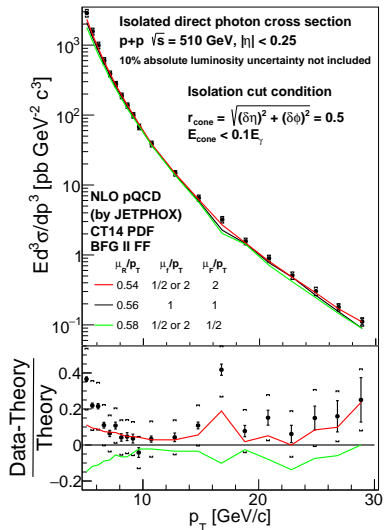


Pure hard processes

JETPHOX for direct photon

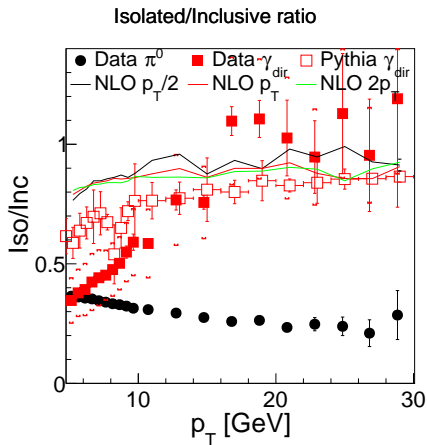


Inclusive

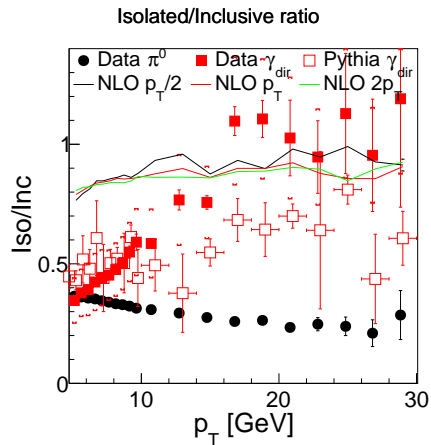


Isolated

POWHEG with MPI (different vetoes) for isolated over inclusive ratio

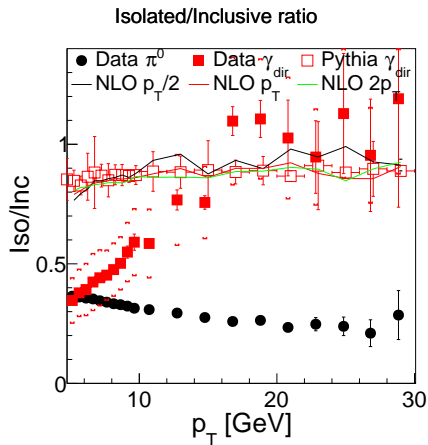


With MPI (default veto)

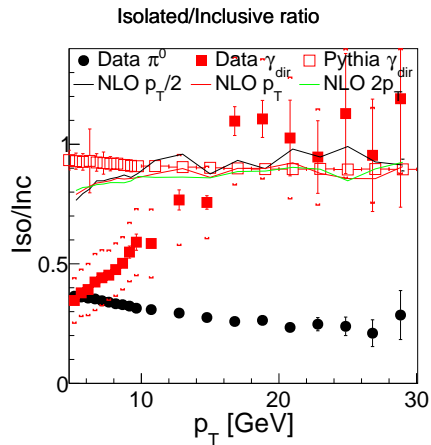


With MPI (QED-QCD veto)

POWHEG without MPI for isolated over inclusive ratio



Without MPI



Pure hard processes

Conclusions for POWHEG study

- ▶ POWHEG with MPI and default veto agrees data at high- p_T , while with QED-QCD veto agrees data at low- p_T .
- ▶ MPI is important in the isolated to inclusive ratio.
- ▶ POWHEG with pure hard processes is just like JETPHOX without fragmentation processes, as expected.