

# Jet Substructure at EIC

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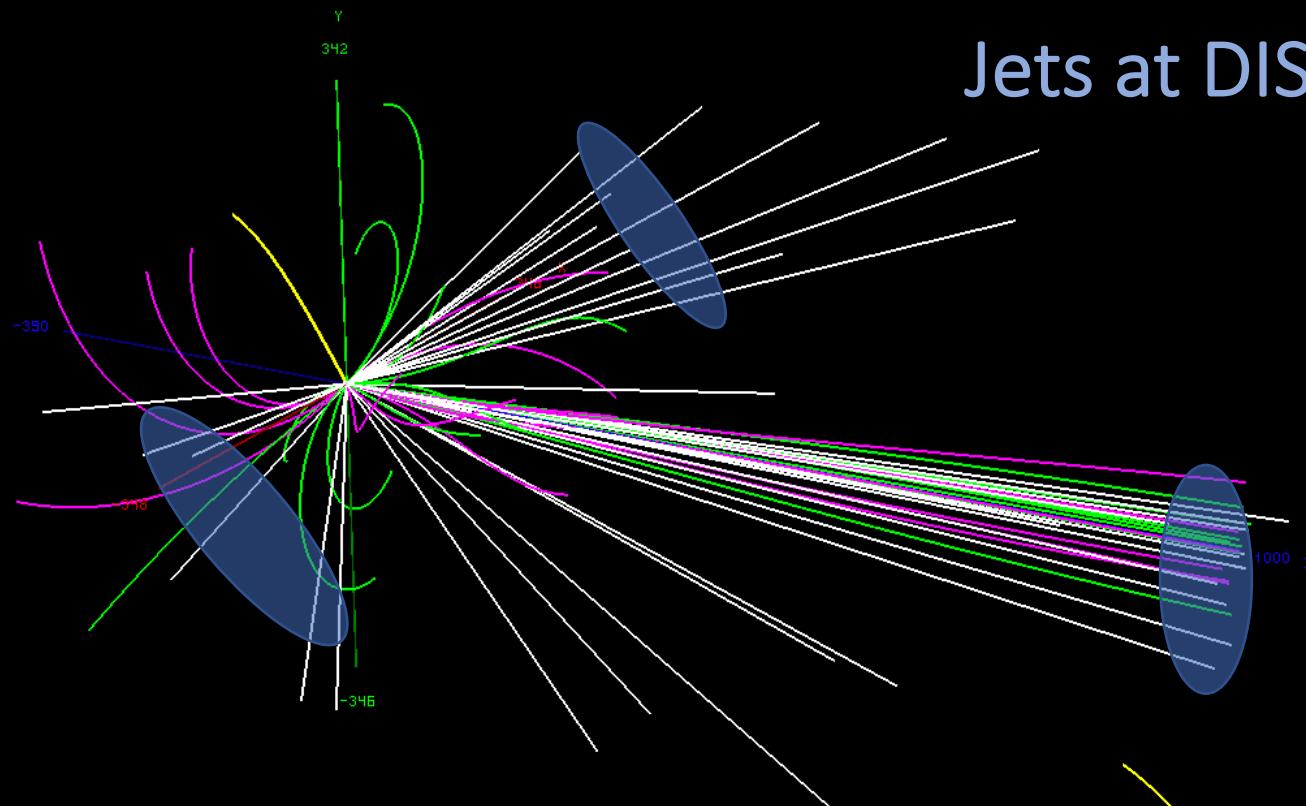
Wayne State University, Detroit

Experimentalist : Raghav Kunnawalkam Elayavalli

# Introduction

- Jets and their structure
  - Correlations in momentum, charge and flavor : leading and next to leading particles in a jet
  - access to the dynamics of fragmentation and color entanglement in QCD
- Observable
  - Charge asymmetry
  - Connection to dihedron fragmentation function
- Pythia event studies
  - Acceptance of Jet and constituent particles
  - Charge asymmetry and in combination of various particle species
  - At various center of mass energies at EIC
- Summary and Outlook

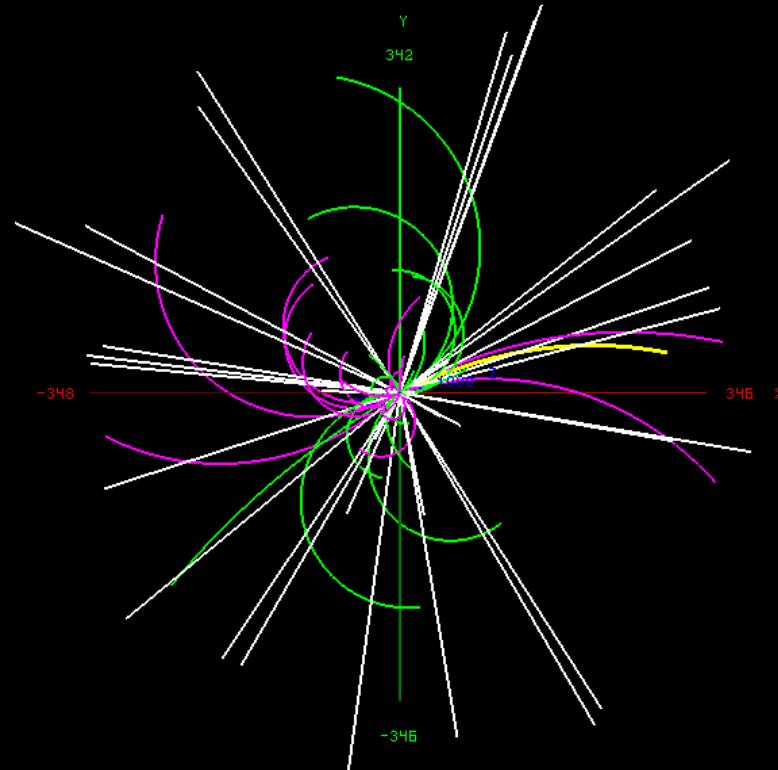
# Jets at DIS



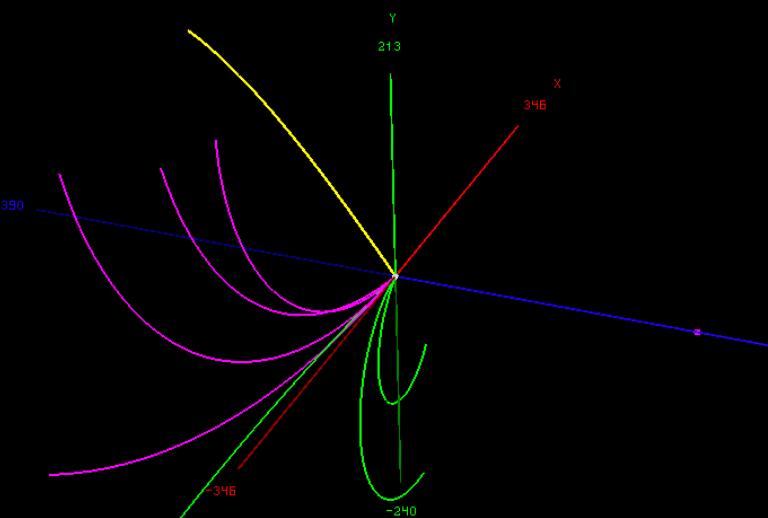
Leading momentum particles in the jet –  $K, \bar{p}$

Particles in Jet :

Px	Py	Pz	PID
-7.64	-4.41	-4.21	321
-2.06	-1.89	-1.40	-2212
-1.44	-0.87	-0.69	-211
-1.07	-0.08	0.04	2212
-0.58	-0.45	-0.53	-211
-0.70	-0.39	-0.39	-211
-0.65	-0.52	-0.04	211

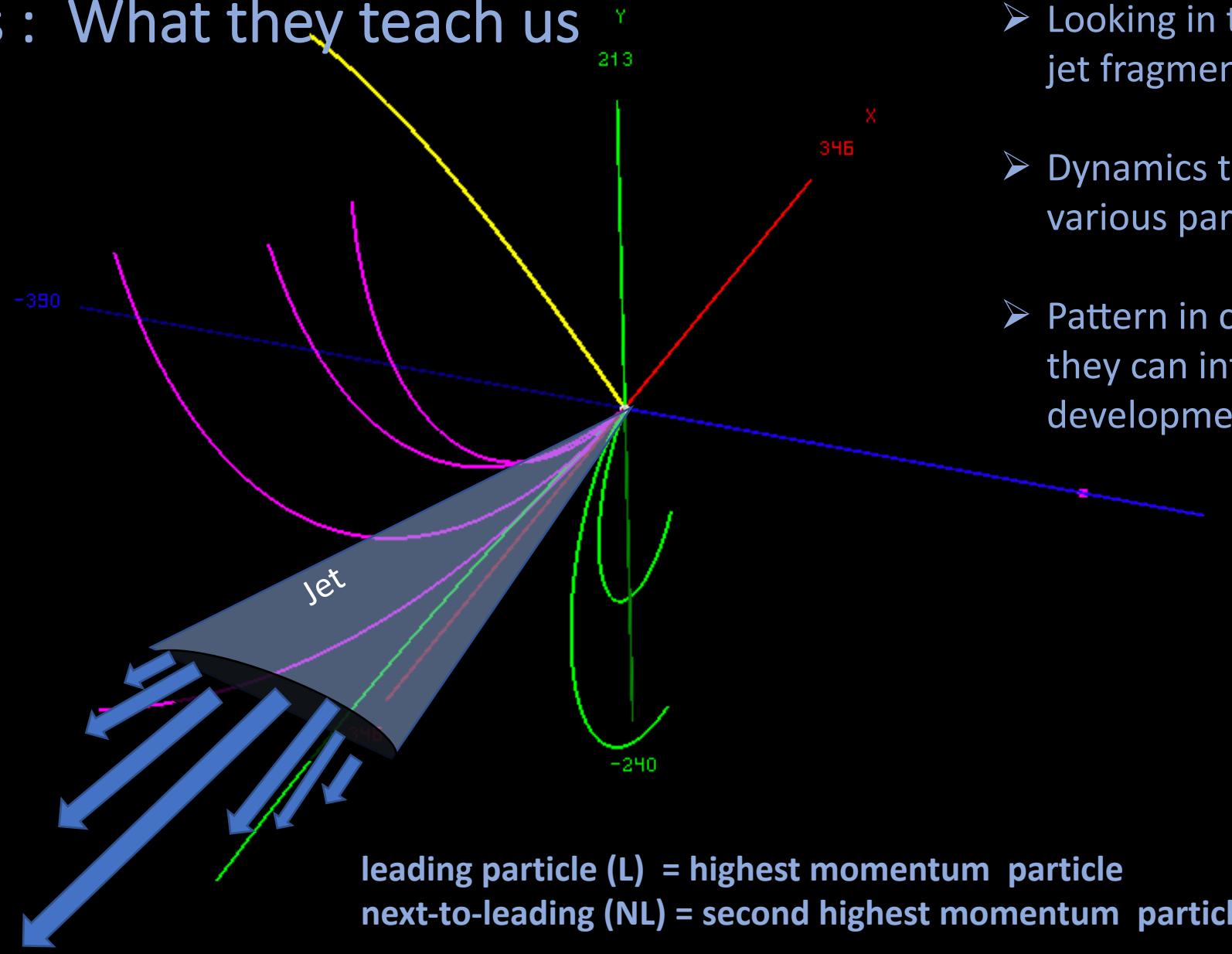


Scattered-electron  
+ve particles  
-ve particles  
Neutral particles



Charged Jet – anti- $k_T$  R = 0.7 ( $p_T$ -jet > 8GeV)

# Jets : What they teach us



- Looking in the nonperturbative aspect of jet fragmentation
- Dynamics that led to fragmentation of various particle species in certain ways
- Pattern in charge and flavor separation : they can inform in future theoretical development

# Momentum-charge correlations

- Leading particle (L) and next-to-leading (NL) are both pions generated in two pictures

i) “random” picture : L is fixed and NL is random and both L and NL pions are charged

$$N_{C\bar{C}}^{\text{random}} = N_{CC}^{\text{random}} = \frac{N^{\text{random}}}{2}$$

$C\bar{C}$  indicates opposite charges  
 $CC$  same charge

ii ) “alternating” picture : perturbative shower gives  $q_L$  followed by  $\bar{q}'_{NL}$ , which form pions by sharing a soft pair:

$$q_L + \bar{q}_{NL} \rightarrow q_L + (\bar{q}_s + q_s) + \bar{q}'_{NL} \rightarrow \pi(q_L, \bar{q}_s) + \pi(q_s, \bar{q}'_{NL})$$

$$N_{C\bar{C}}^{\text{alternating}} = N_{CC}^{\text{alternating}} \quad \text{and} \quad N_{CC}^{\text{alternating}} = 0$$

- The observable :

$$r_{\text{asy}} \equiv \frac{N_{CC} - N_{C\bar{C}}}{N_{CC} + N_{C\bar{C}}} = \frac{1-a}{2} - \left( \frac{1-a}{2} + a \right) = -a$$

- ✓ provided every event results from one of these two processes, with no interference
- ✓ percentage of “alternating” = a; and percentage of random events = 1-a

- $r_{\text{asy}}$  is a measurement of the fraction of hadronizations that are “string-like”, alternating between quark and antiquark (classical picture)

# Measurements of $r_{\text{asy}}$ and expressing in terms of di-hadron fragmentation functions

- Measurements of  $r_{\text{asy}}$ :
  - differentially in fractions  $z_L$  and  $z_{NL}$  in a jet,
  - “transverse” kinematic variables:
    - ✓ relative transverse momentum
    - ✓ pair invariant mass
    - ✓ pair formation time
    - ✓ including polarization where applicable

## □ $r_{\text{asy}}$ and its connection with generalized di-hadron fragmentation functions

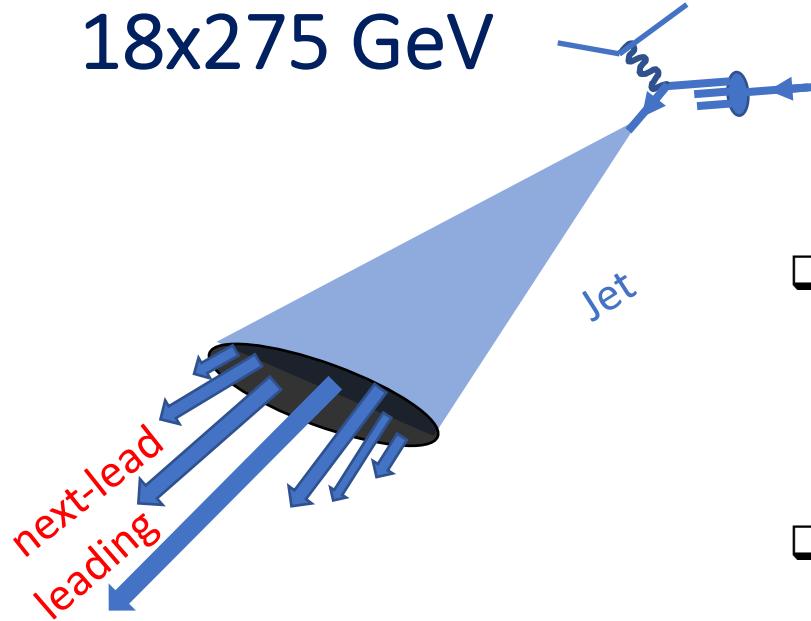
generalized di-hadron fragmentation functions for any hadrons  $h_1, h_2$ :  $D_{h_L, h_{NL}}^>(z_L z_{NL})$

$$N_{h_L, h_{NL}}^> = \int_0^1 dz_L \int^{\min(z_L, 1-z_L)} dz_{NL} D_{h_1, h_2}^>(z_L, z_{NL}, Q)$$

When  $z_L$  and  $z_{NL}$  are large enough, this is the usual di-hadron distribution  $D^>(x_1, x_2, Q) = D(x_1, x_2, Q)$  when  $x_2 > 1 - x_1 - x_2$

$$r_S = \frac{\sum_{h_1, h_2 \in S} Q_{h_1} Q_{h_2} N_{h_1, h_2}^>}{\sum_{h_1, h_2 \in S} |Q_{h_1} Q_{h_2}| N_{h_1, h_2}^>}$$

18x275 GeV



PYTHIA-6 : 1M events :  
Q2 > 65 GeV

Jet Reconstruction :  
anti-kt R = 0.7  
Jet pt > 8GeV  
pt-tracks > 0.2GeV  
track |eta| < 3.5  
Jet |eta| < 2.8

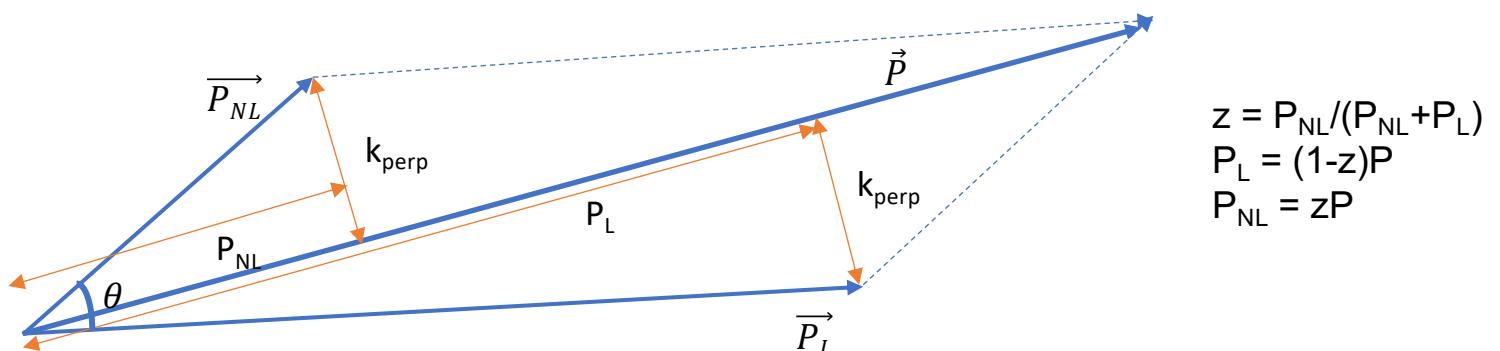
$$r_{\text{asy}} \equiv \frac{N_{CC} - N_{C\bar{C}}}{N_{CC} + N_{C\bar{C}}}$$

❑ Construct  $r_{\text{asy}}$  with particle compositions with various parameters

- ✓ Leading particle ( $\pi$ ) and next leading particles ( $\pi/K/p$ )
- ✓ Leading particle ( $K$ ) and next leading particles ( $K/\pi/p$ )
- ✓ Leading particle ( $p$ ) and next leading particles ( $p/\pi/K$ )

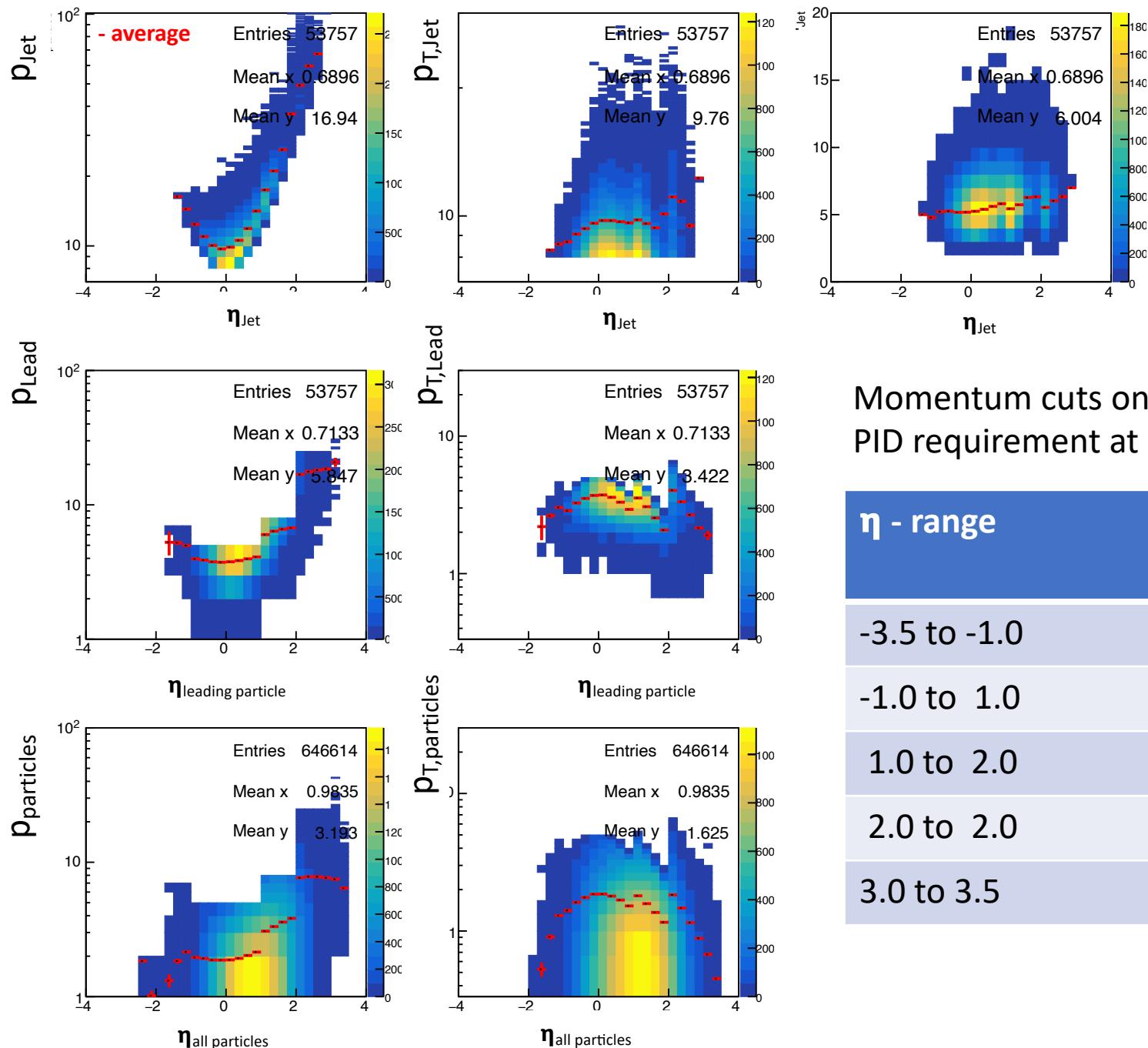
❑ Looking via

- ✓ Momentum-next lead particle/momentum of leading particle
- ✓ Fraction of jet momentum carried by leading particle
- ✓ Angle between the leading and next to leading particles ( $\Delta\theta$ )
- ✓ relative transverse momentum ( $k_{\text{perp}}$ )
- ✓ pair invariant mass
- ✓ Formation time :  $[2z(1-z) P] / k_{\text{perp}}^2$



# Acceptance

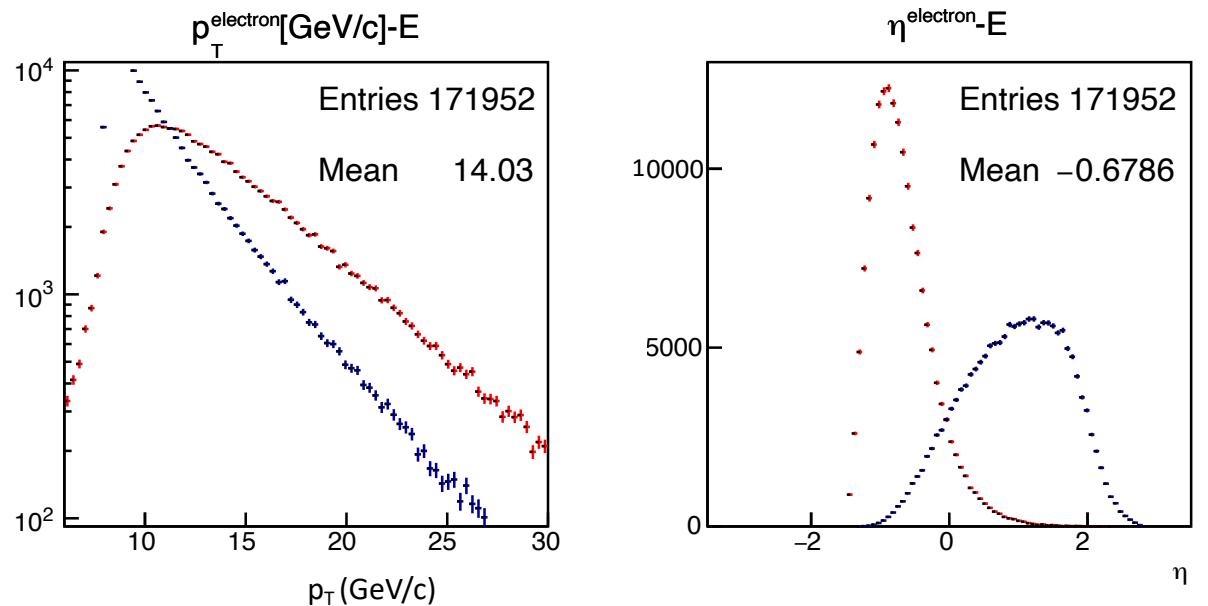
18x275 GeV



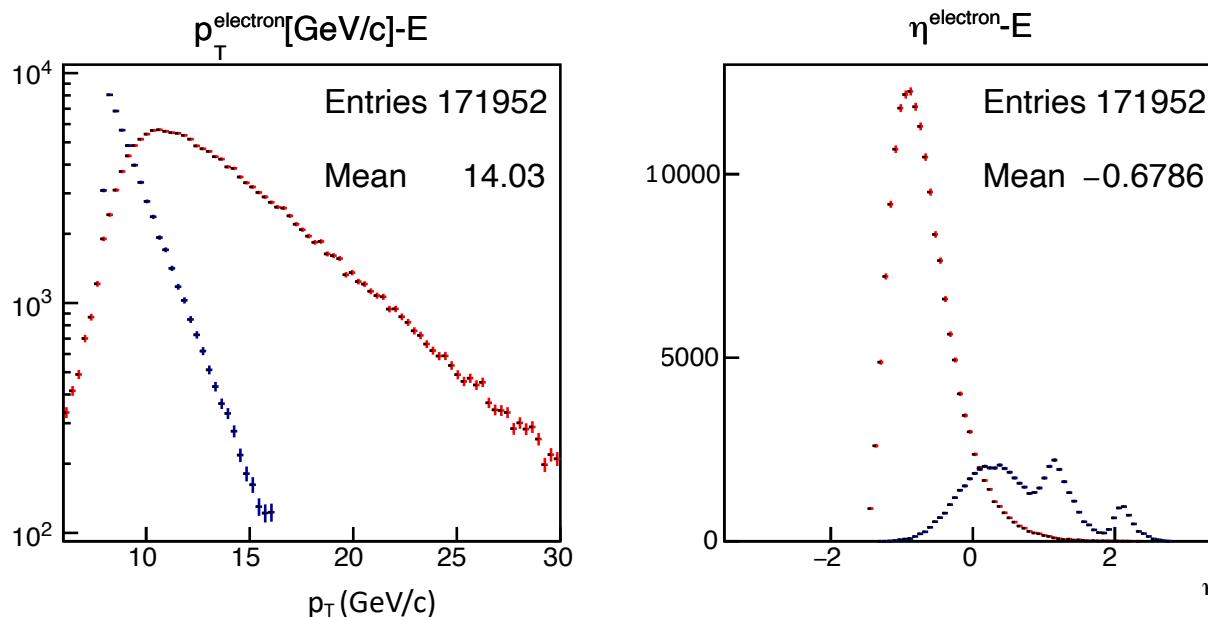
Momentum cuts on tracks for approximate PID requirement at EIC at different  $\eta$  regions

$\eta$ - range	Momentum cut (GeV/c <sup>2</sup> )
-3.5 to -1.0	7
-1.0 to 1.0	5
1.0 to 2.0	8
2.0 to 2.0	25
3.0 to 3.5	45

# The effect of momentum cuts in traverse momentum and acceptance of jets

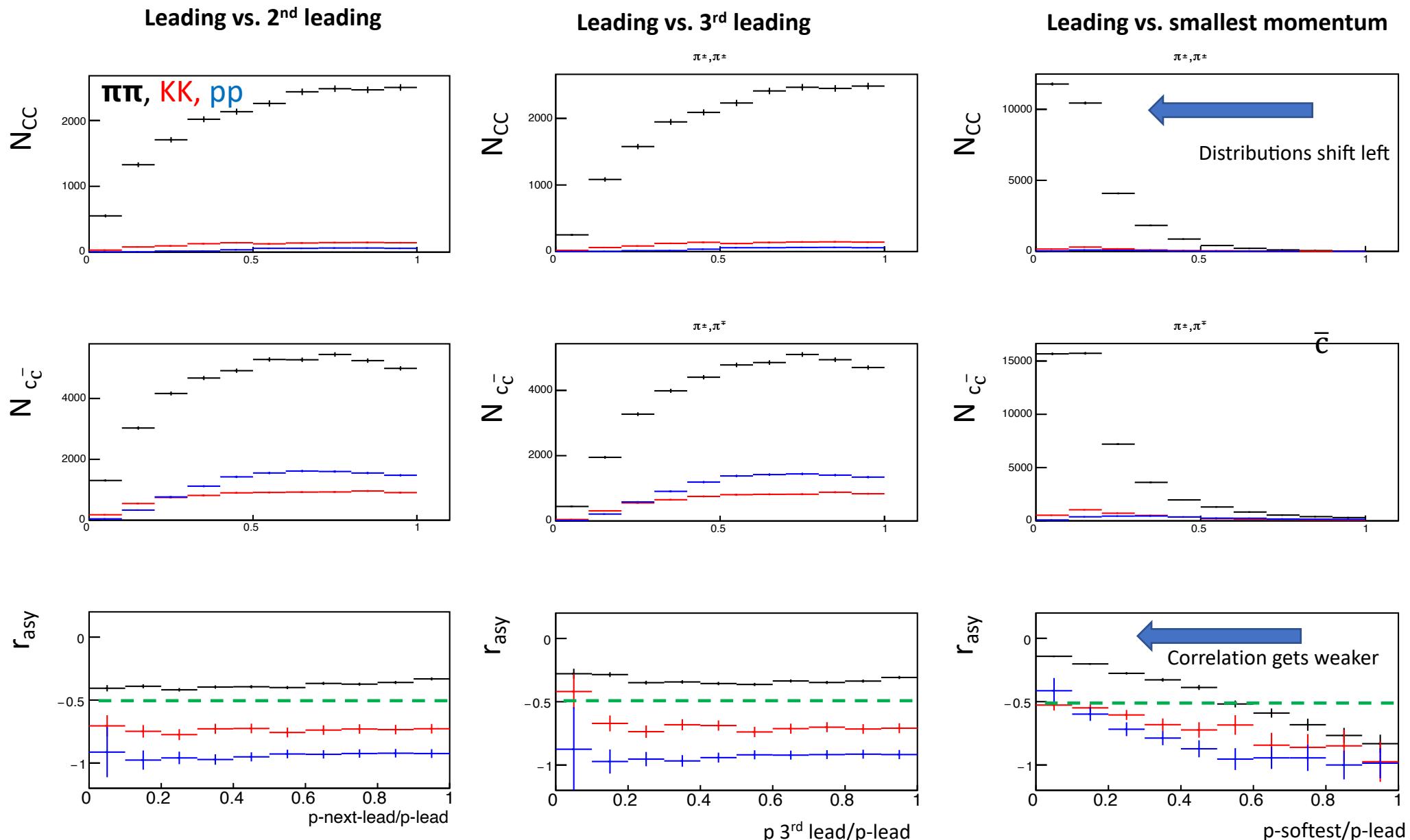


RED-electron  
BLUE-JET  
  
NO  
momentum  
cuts



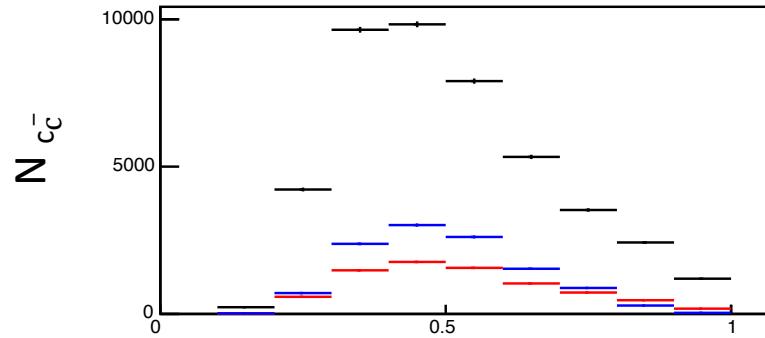
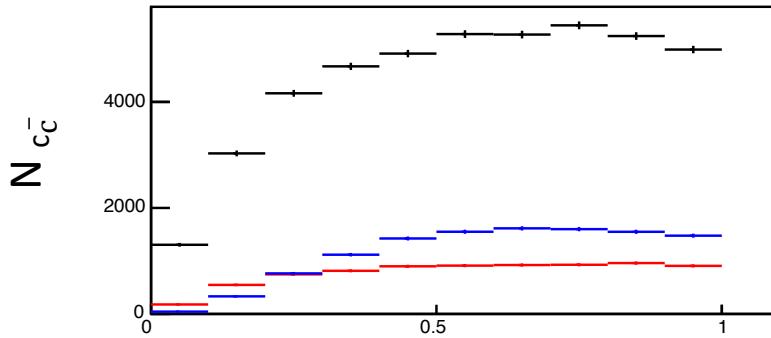
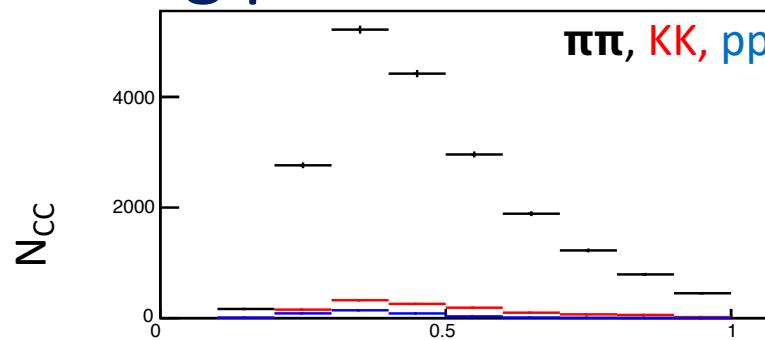
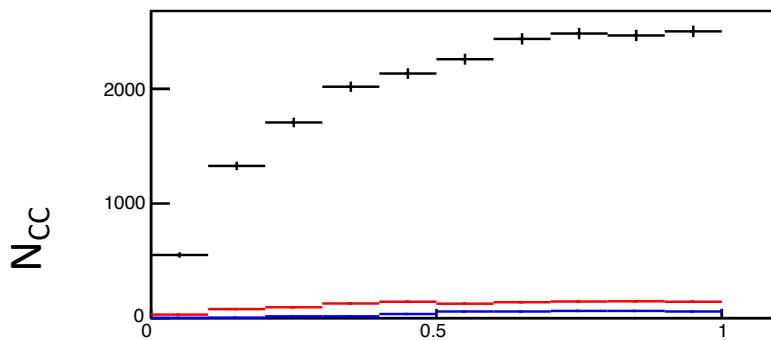
After  
momentum  
cuts

# Correlation with different particles

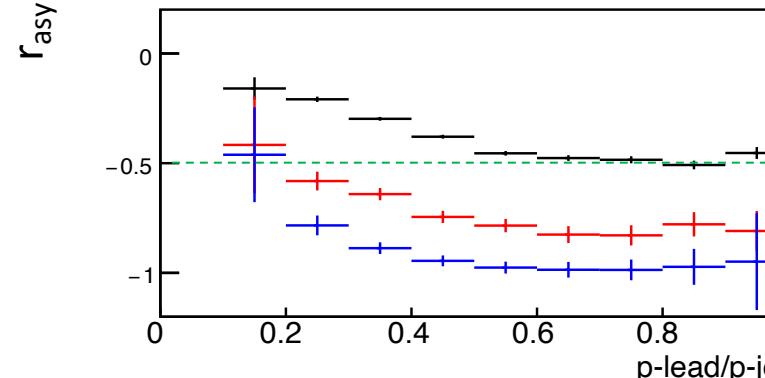
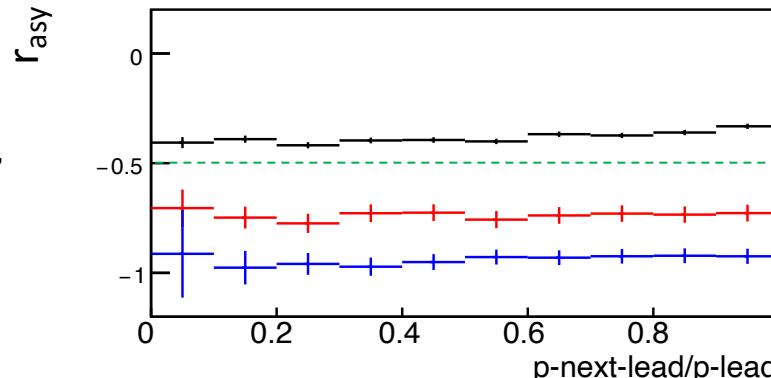


# Correlation with two leading particles

NO  
momentum  
cuts



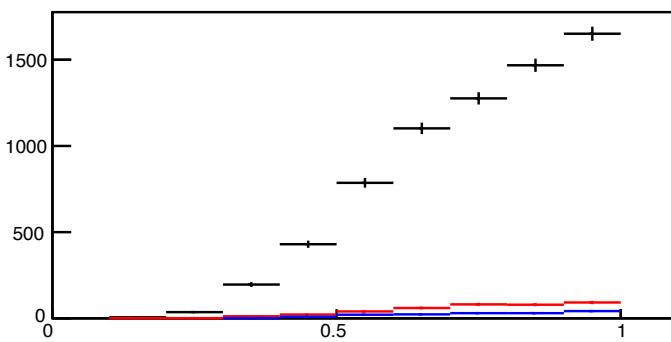
- The strength of correlations are different for pions, kaons and proton



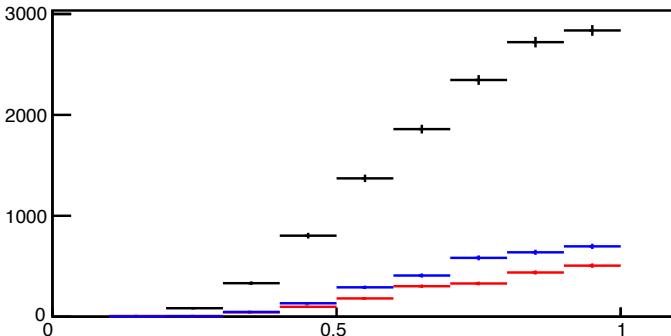
**After  
momentum  
cuts**

- The strength of correlations are different decrease as an acceptance effect
- Acceptance changes in certain regions due to momentum cuts

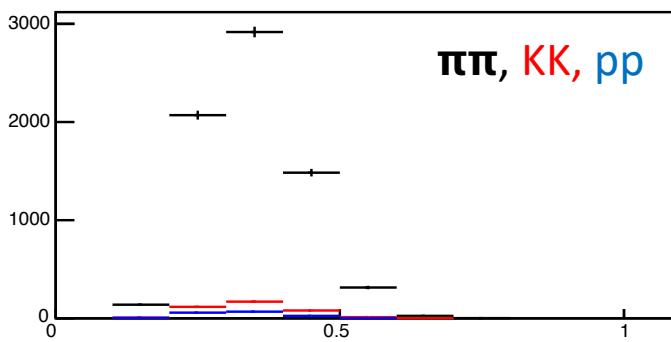
$N_{CC}$



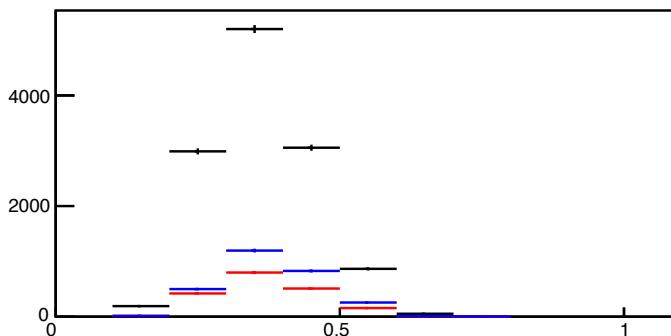
$N_{c\bar{c}}$



$N_{CC}$

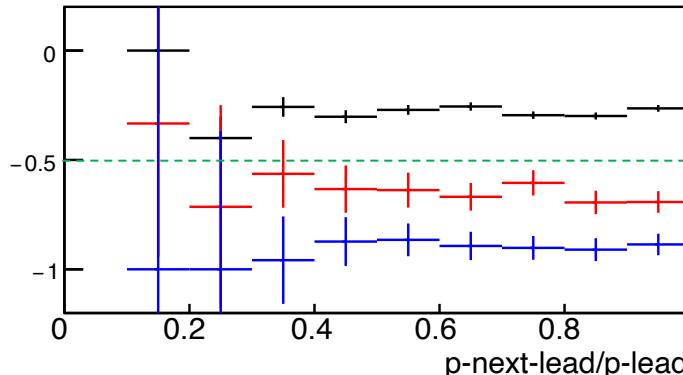


$N_{c\bar{c}}$

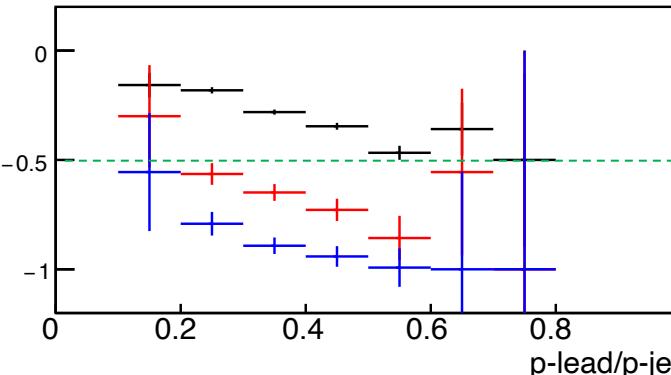


$r_{asy}$

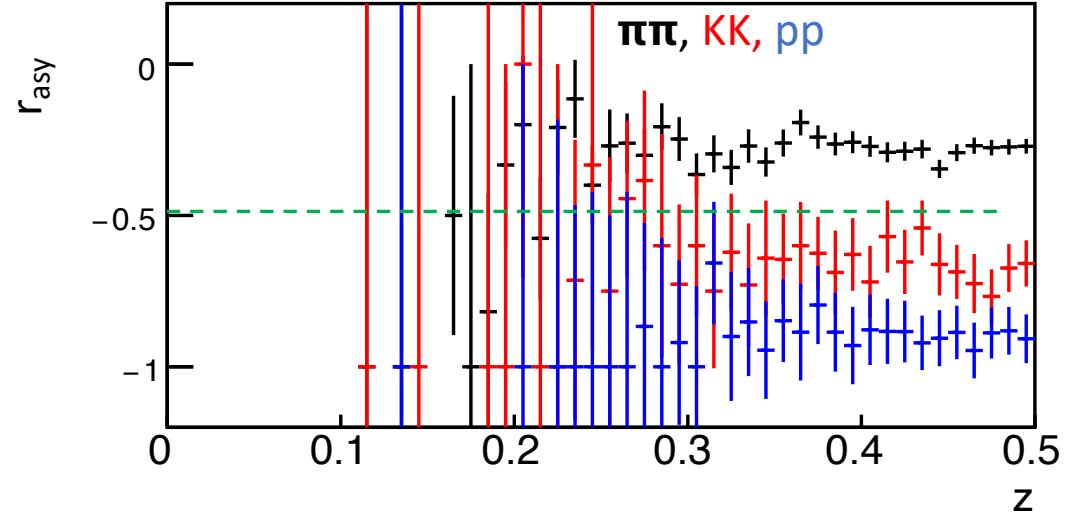
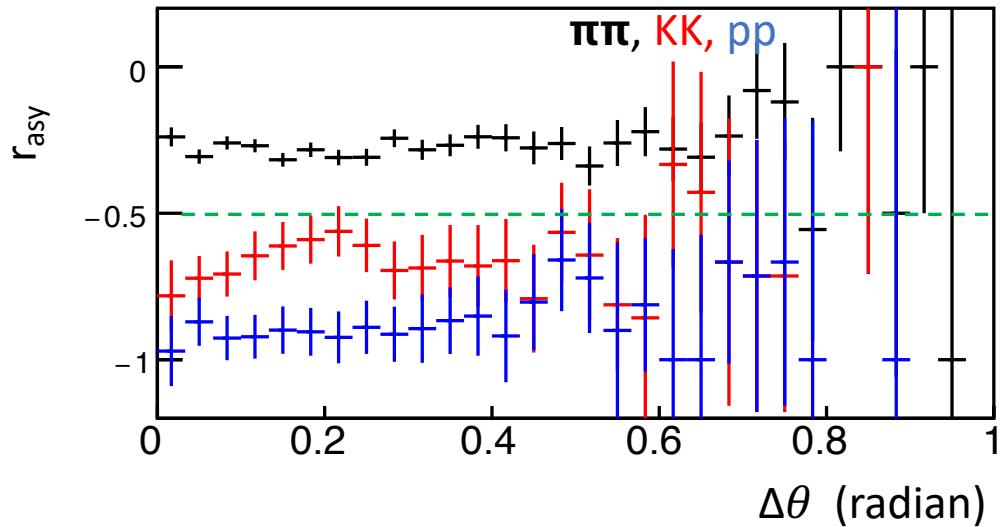
$r_{asy}$



$r_{asy}$



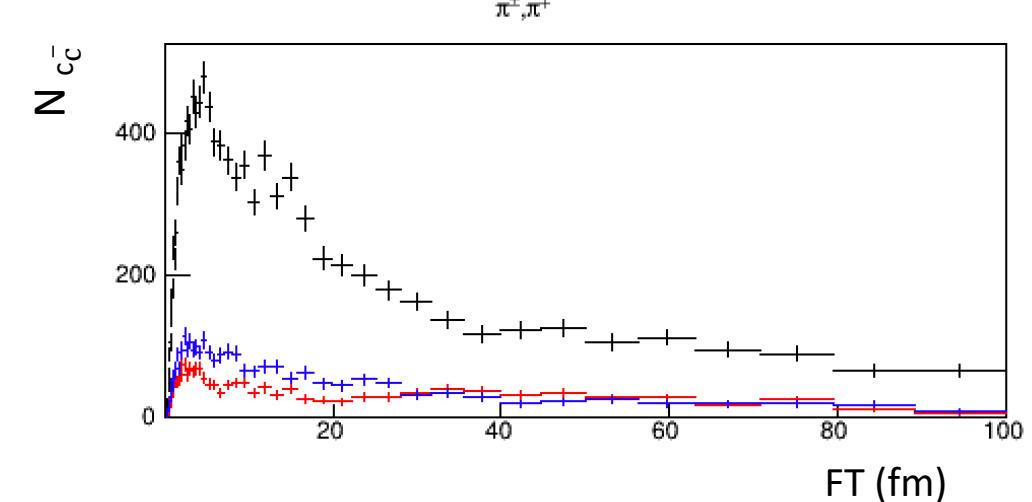
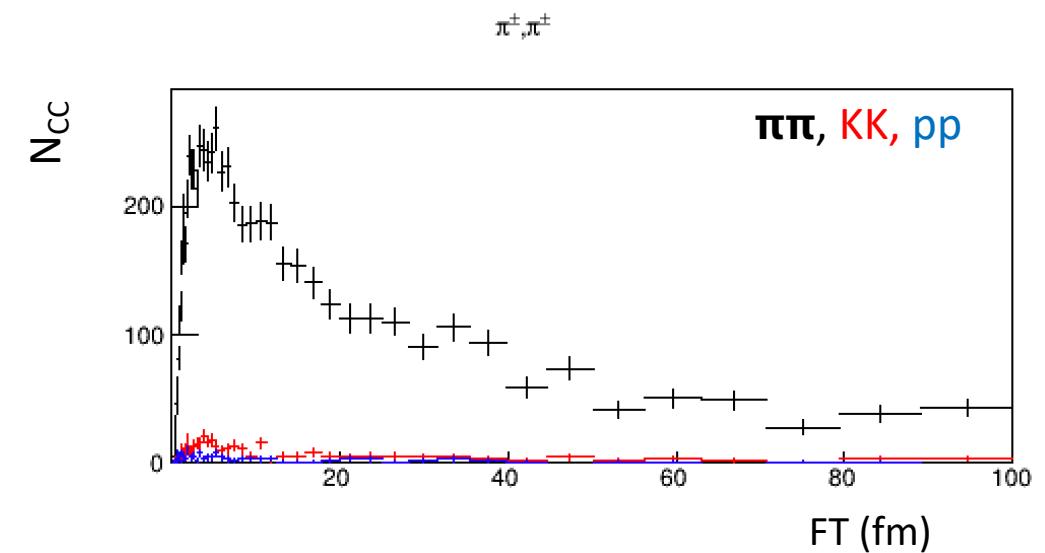
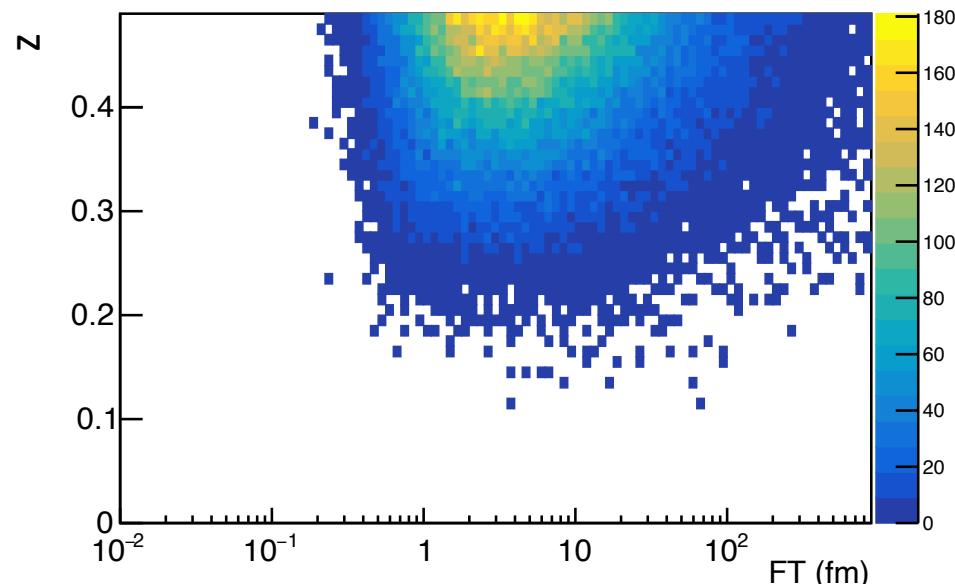
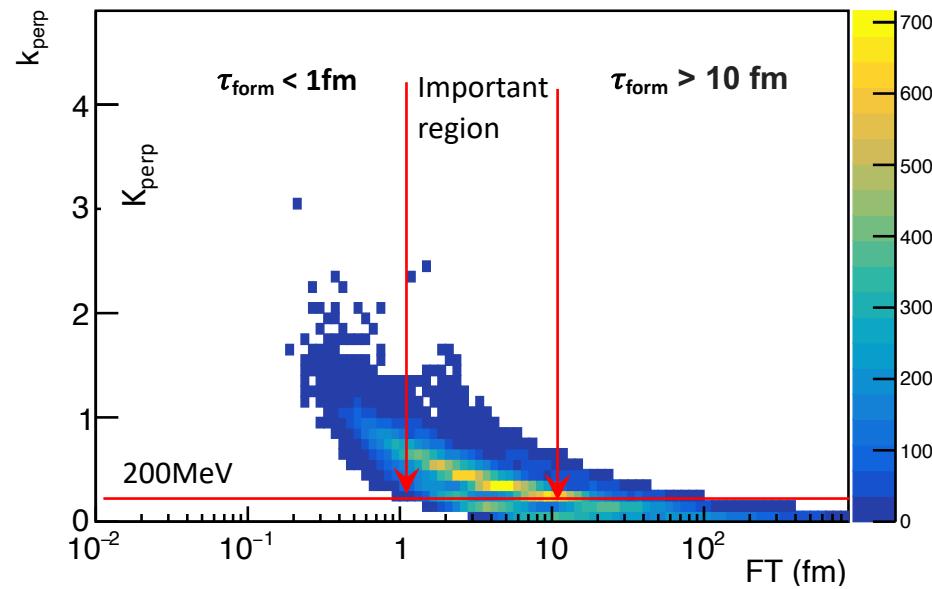
# Dependence on $\Delta\theta$ and z



- $r_{\text{asy}}$  values remain almost flat opening angle and z with some little variations

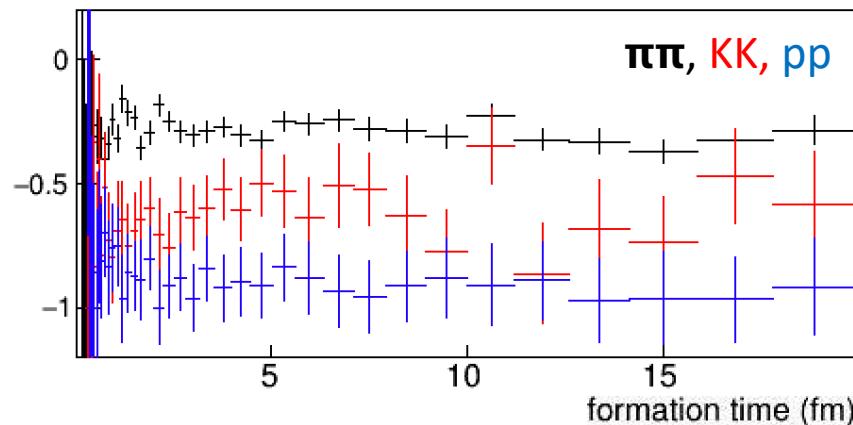
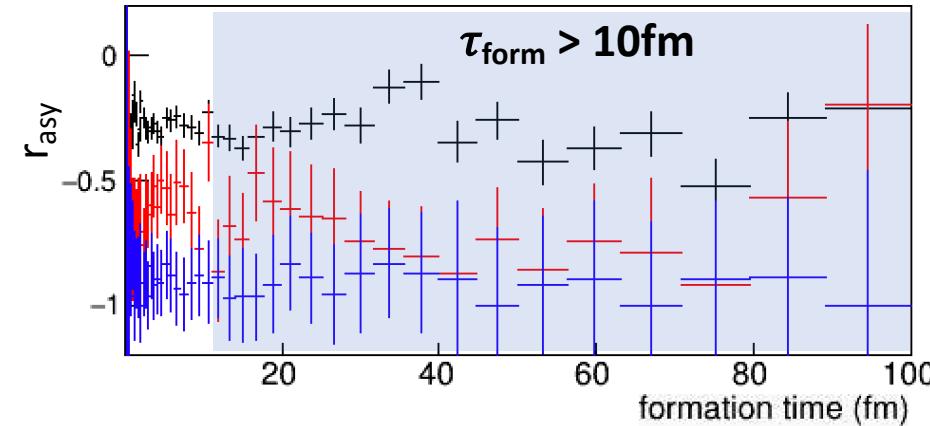
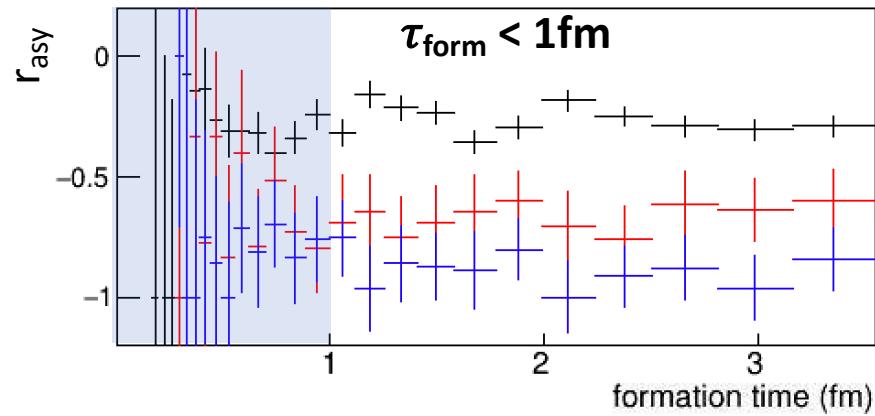
# Formation time

$$\text{Formation time, } \tau_{\text{form}} = [2z(1-z) P] / k_{\text{perp}}^2$$



# Formation time

$$\text{Formation time, } \tau_{\text{form}} = [2z(1-z) P] / k_{\text{perp}}^2$$

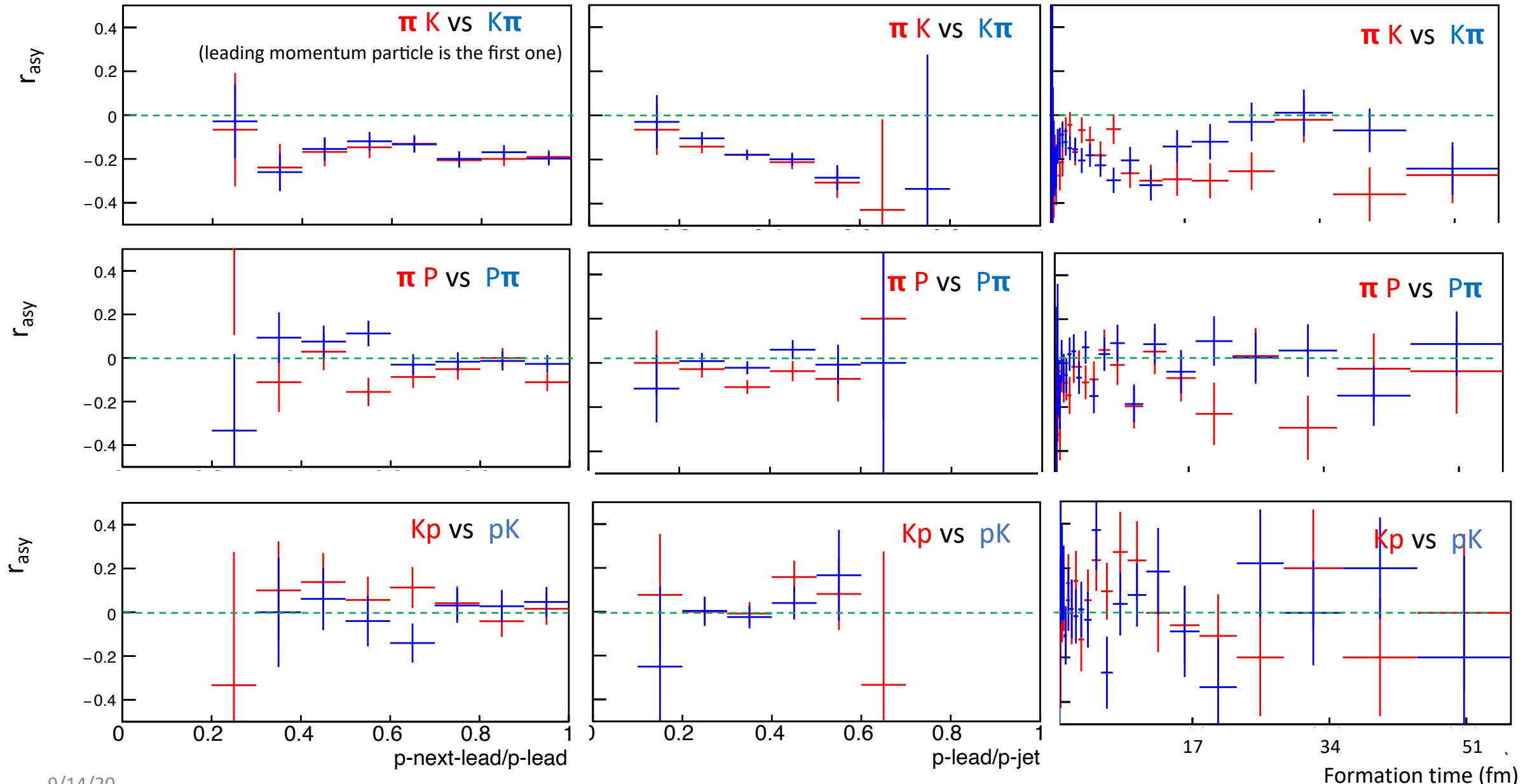


$\tau_{\text{form}} < 1 \text{ fm}$  : L and NL particles seem to separate after a very short time, which might decorrelate their hadronization.

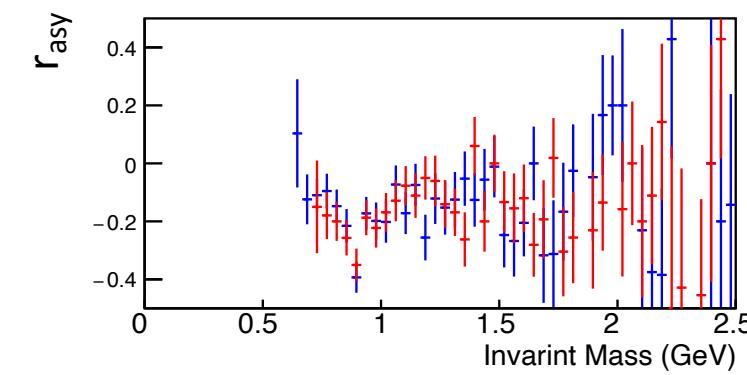
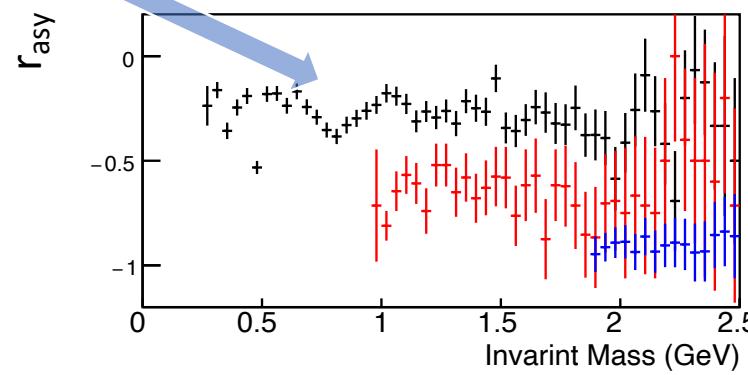
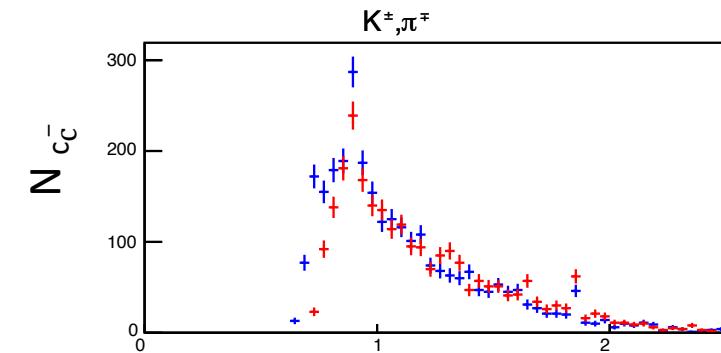
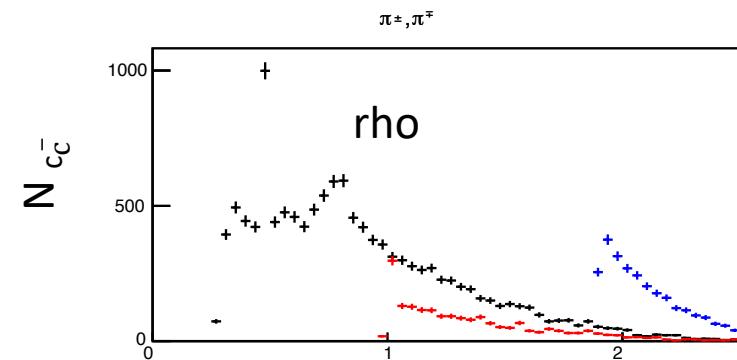
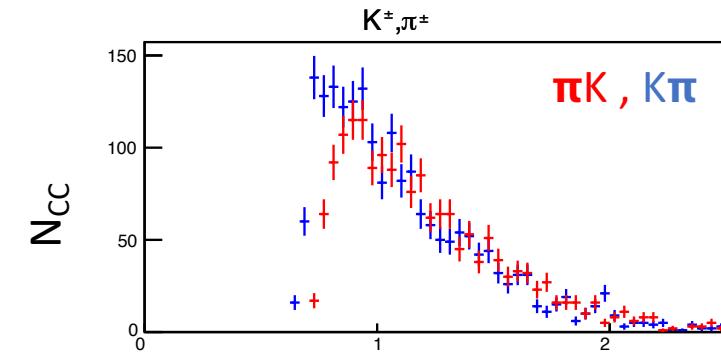
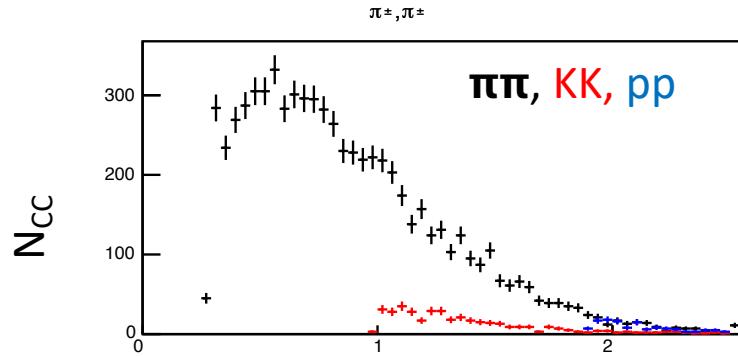
$\tau_{\text{form}} > 10 \text{ fm} (k_{\text{perp}} < 200 \text{ MeV})$  : nonperturbative transverse momenta in the jet, and we don't think that going to longer  $\tau_{\text{form}}$  or smaller  $k_{\text{perp}}$  leads to new dynamics

Important region to study in data  $\tau_{\text{form}} = \text{"a few fermi" and "a few dozen fermi"}, k_{\text{perp}} = \text{"a few GeV" to "several hundred MeV"}$

# Combination of different leading particles

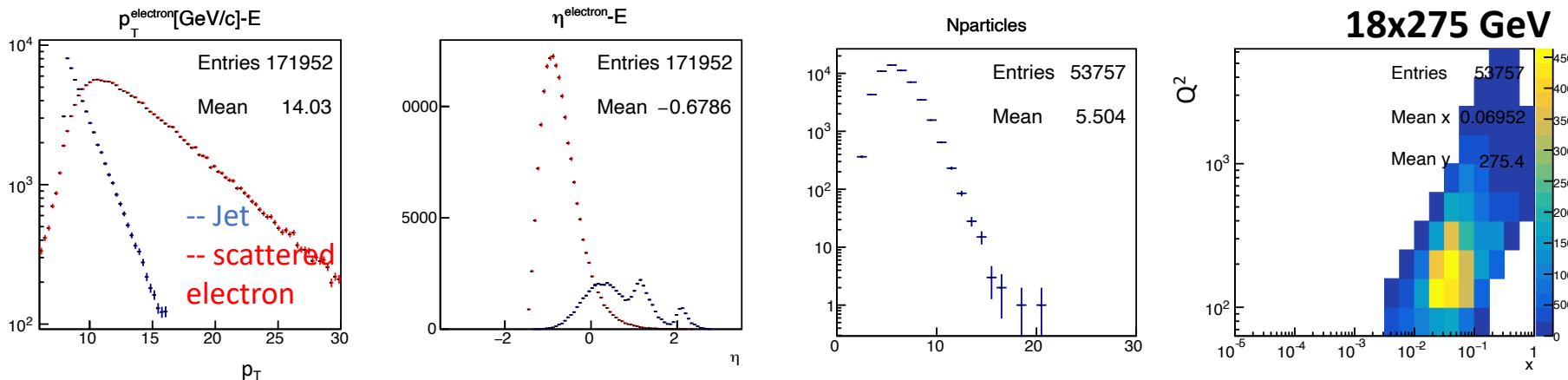


# Invariant mass

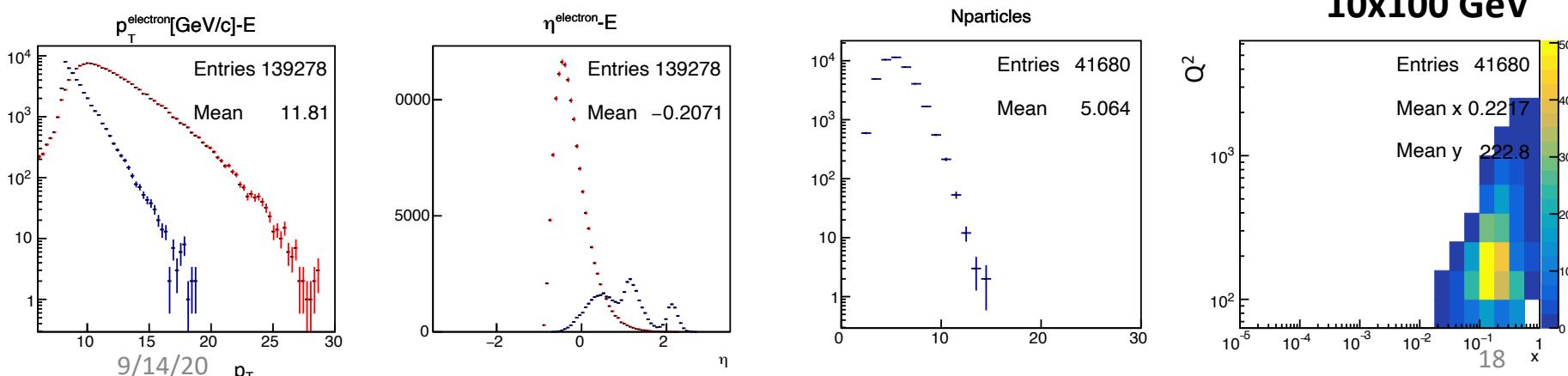
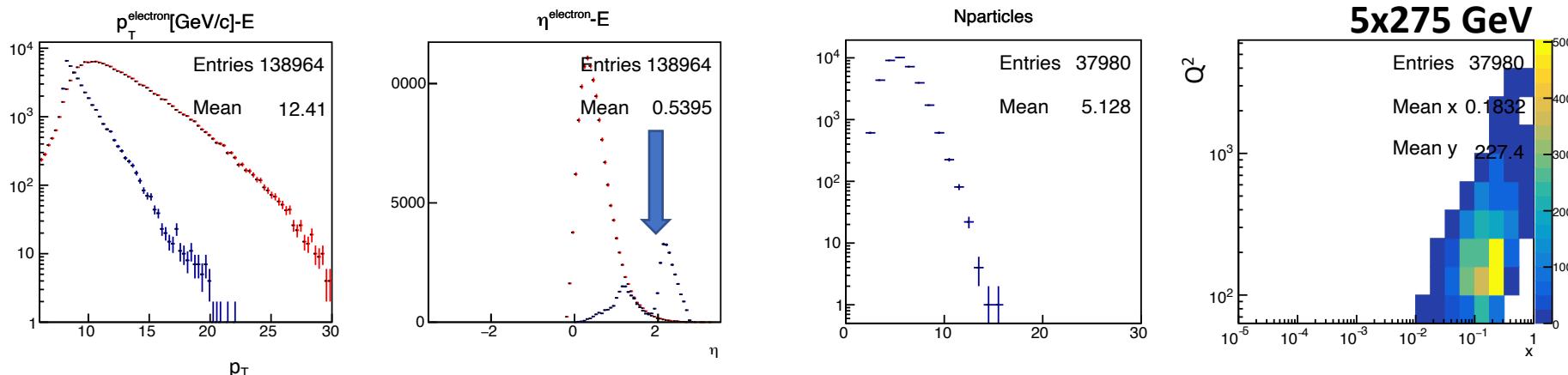


Resonance  
decorrelate the  
asymmetry

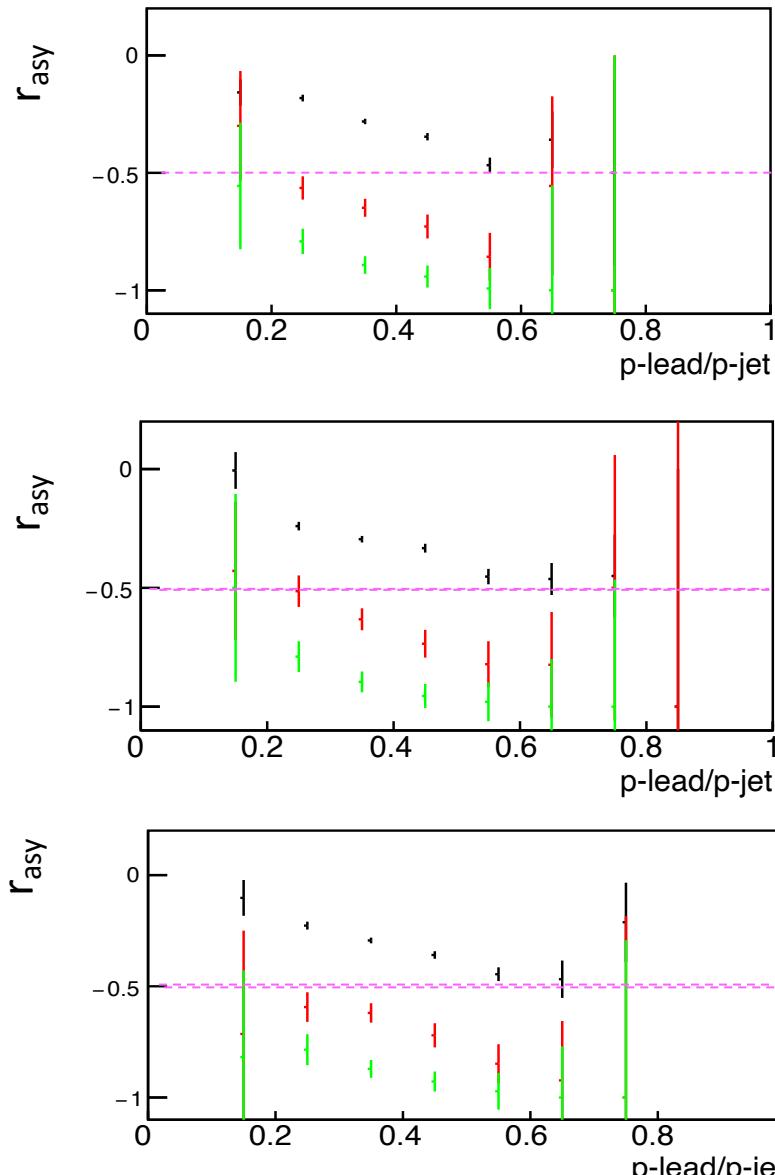
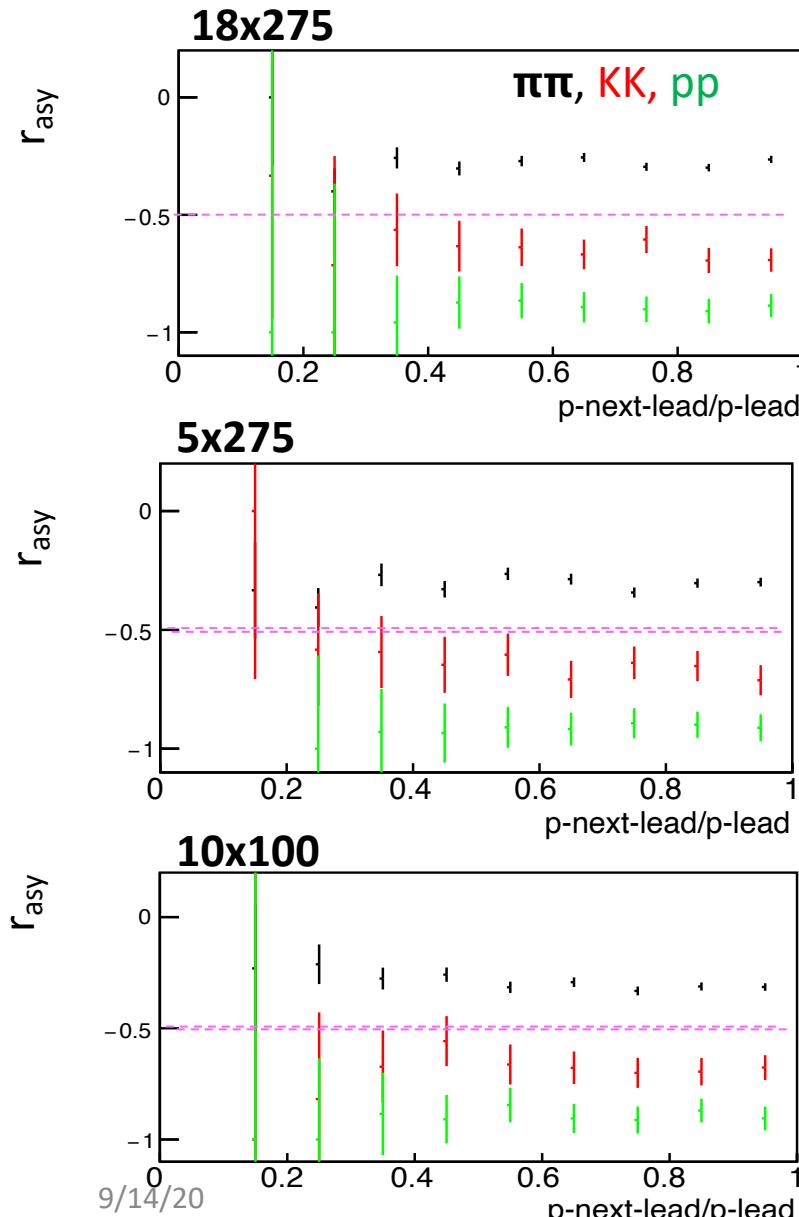
# Acceptance at various collision energies



Most asymmetric –  
acceptance shift to  
Forward rapidity :  
5x275 GeV



# At different collision energies



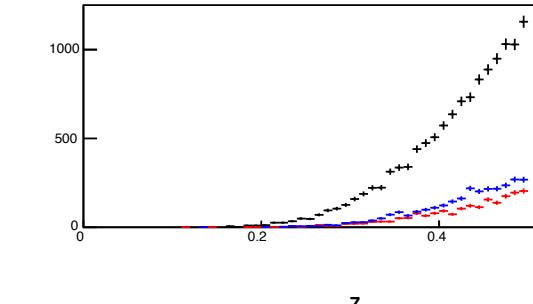
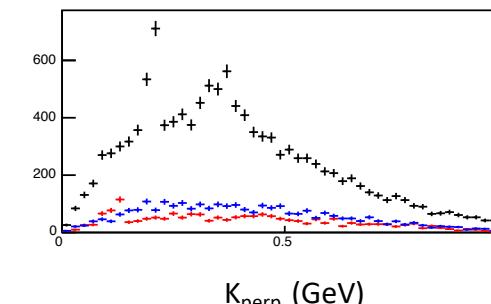
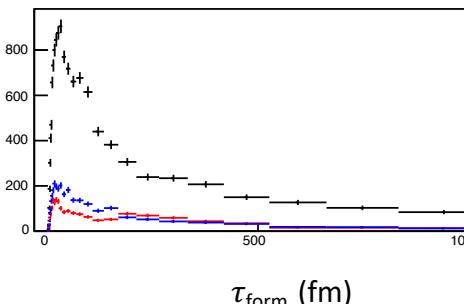
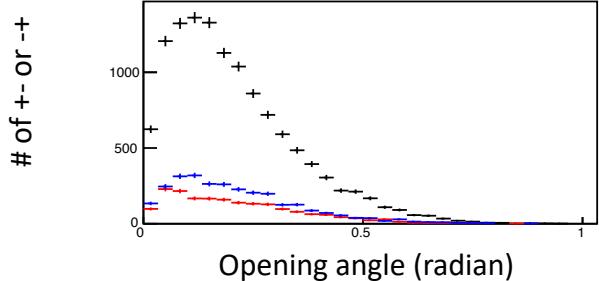
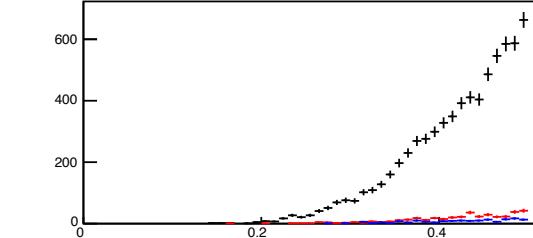
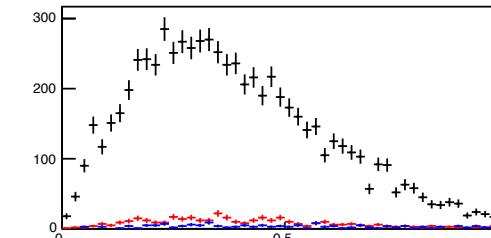
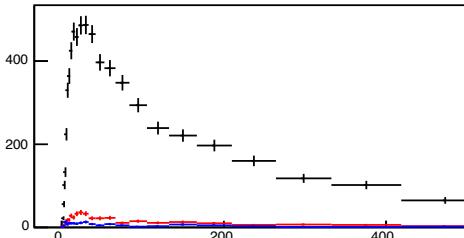
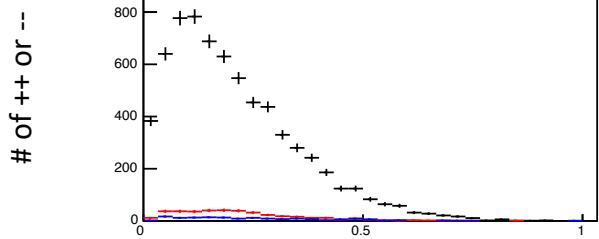
- At various center of mass energies  $r_{\text{asy}}$  looks similar

# Summary

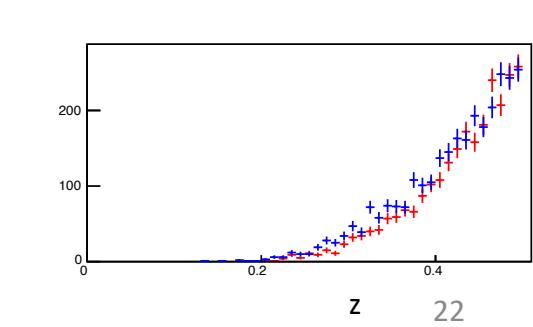
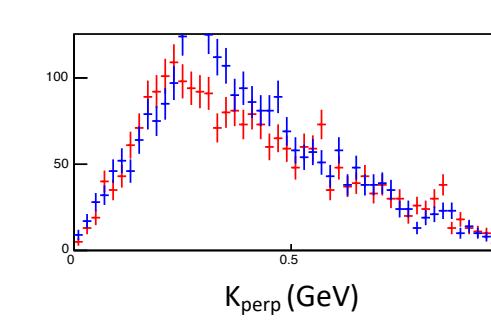
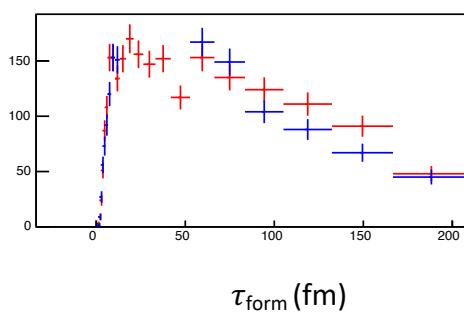
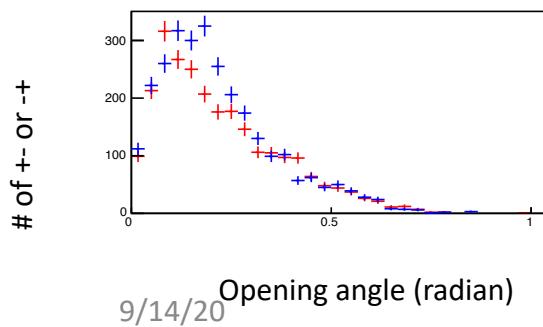
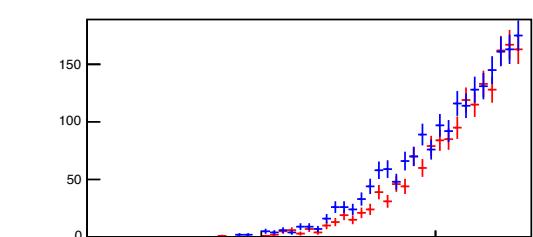
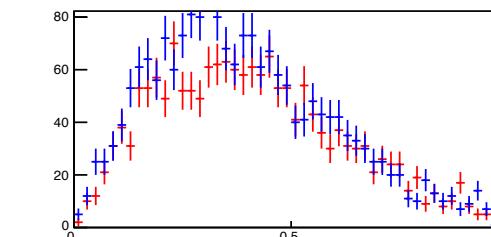
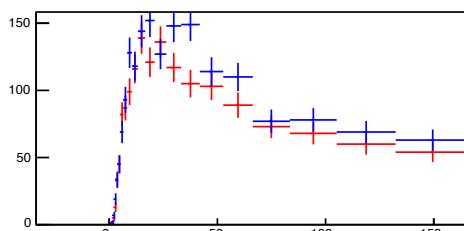
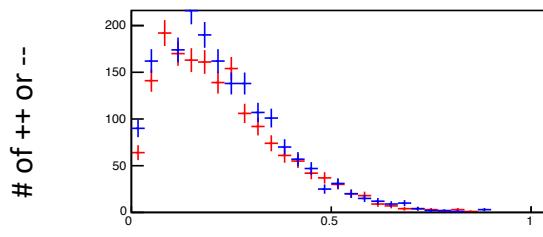
- Correlations in momentum, charge and flavor of leading particles in jet carry information of non perturbative aspect of jet fragmentation
- At EIC this can be measured with high momentum PID capabilities
- The embedded correlations in PYTHIA is studied at different set of collision energies for pions, kaons and protons
- We plan to study :
  - ✓ Bringing neutral particle in the correlations
  - ✓ other fragmentation models using HERWIG
  - ✓ Inclusion of jet charge in our study
- We would like to include the studies in Yellow Report activities

# backup

# $\pi\pi$ , KK, pp (18x275 GeV)

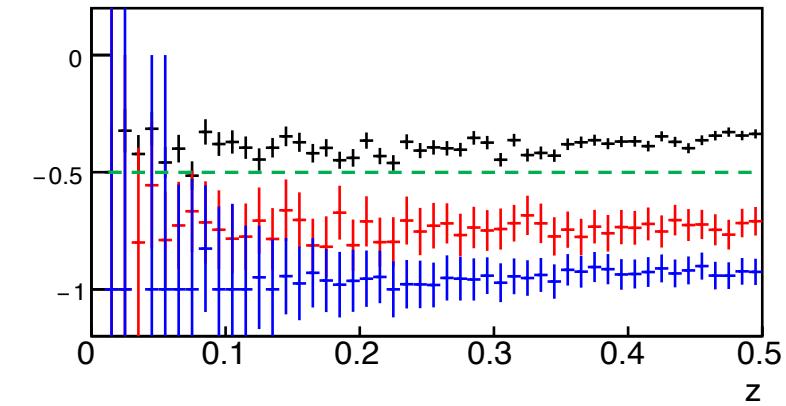
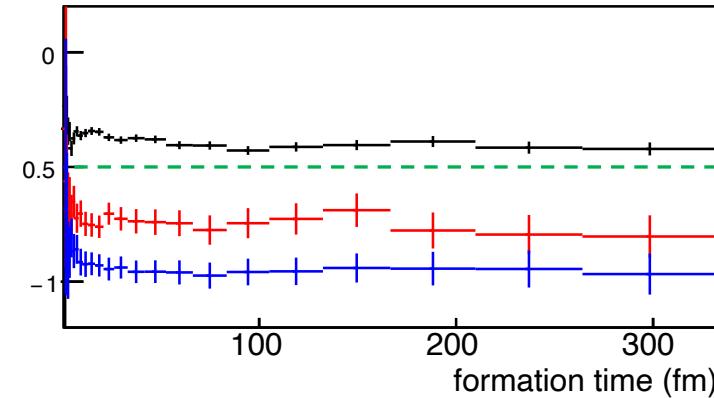
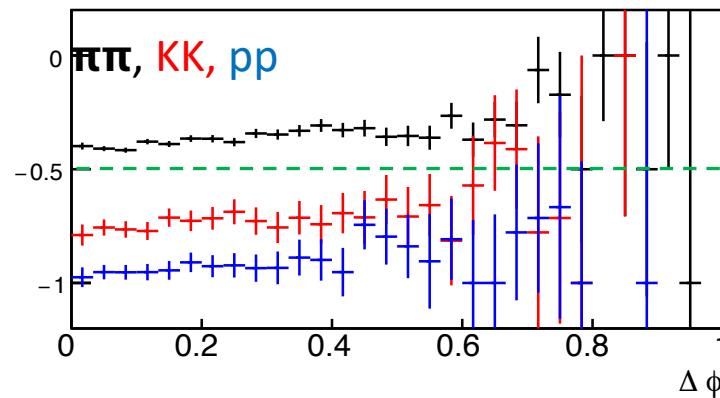


## $\pi K$ vs $K\pi$



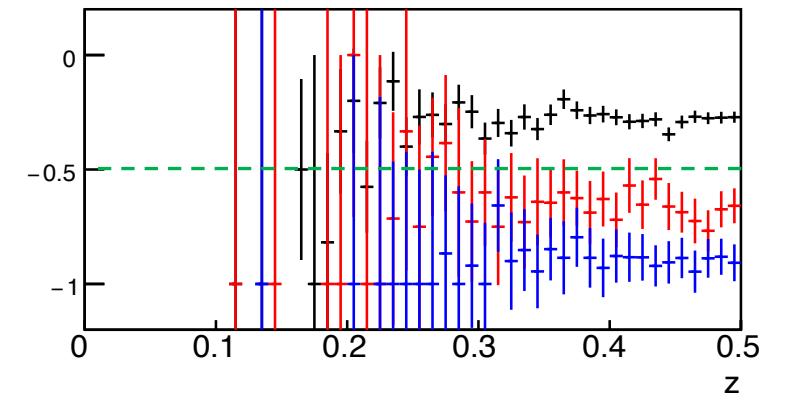
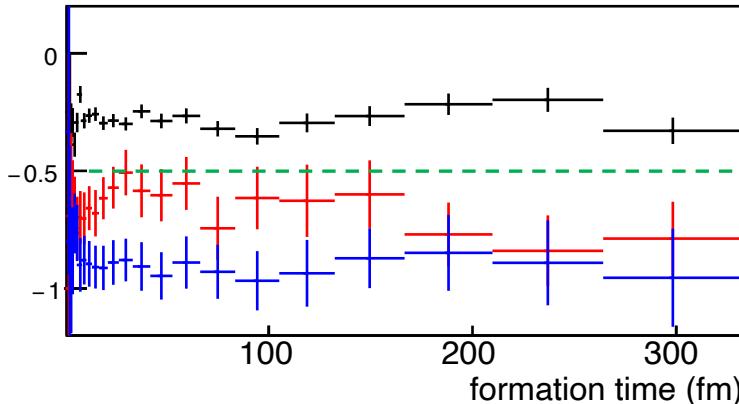
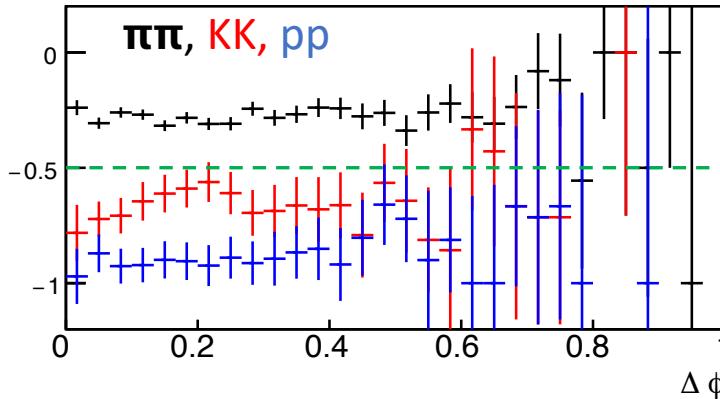
$$r = \frac{(\# of + or --) - (\# of \pm or \mp)}{\text{total pairs}}$$

### NO momentum cuts



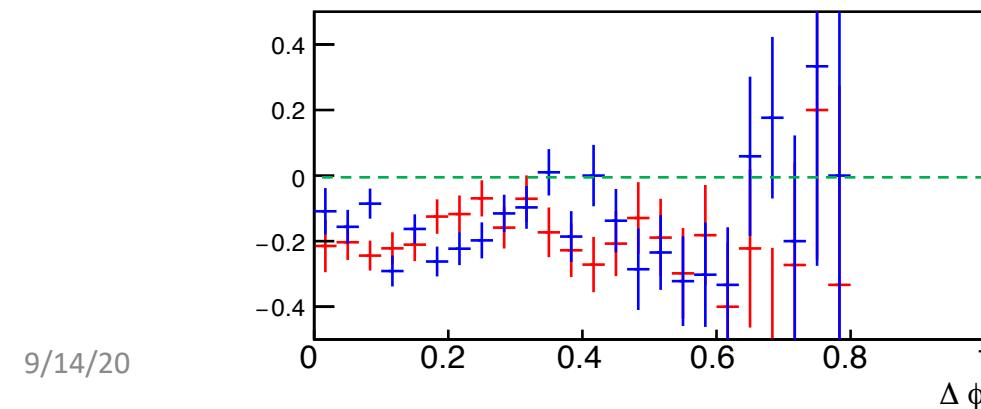
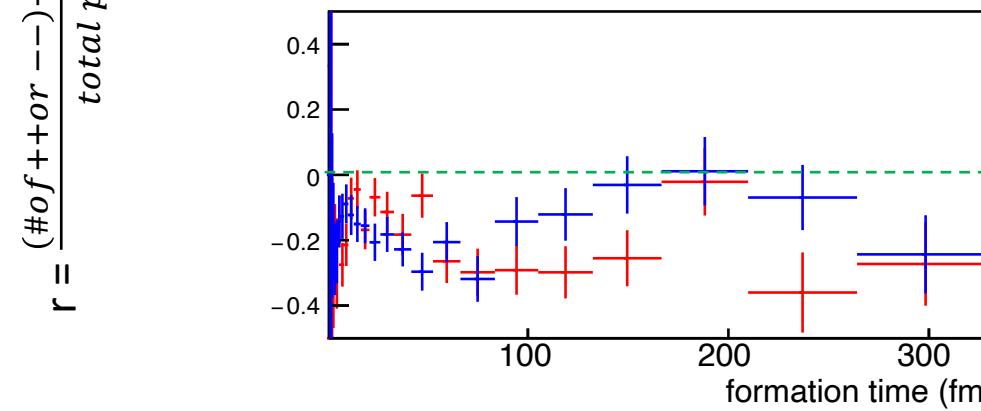
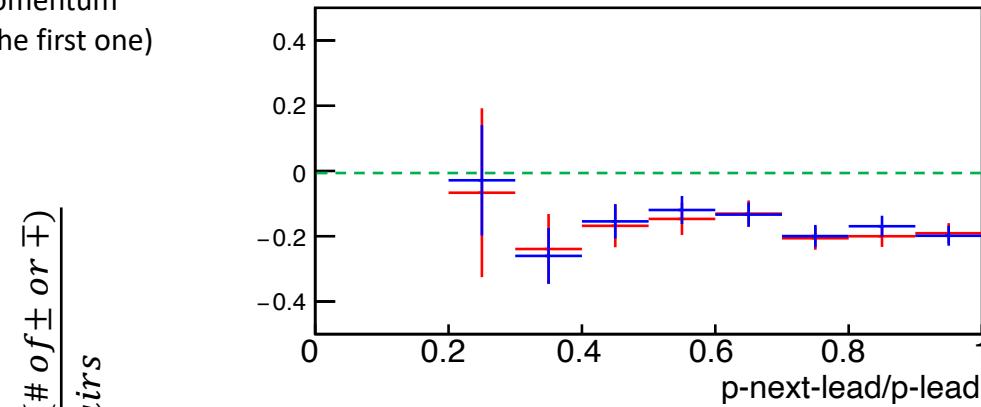
$$r = \frac{(\# of + or --) - (\# of \pm or \mp)}{\text{total pairs}}$$

### After momentum cuts

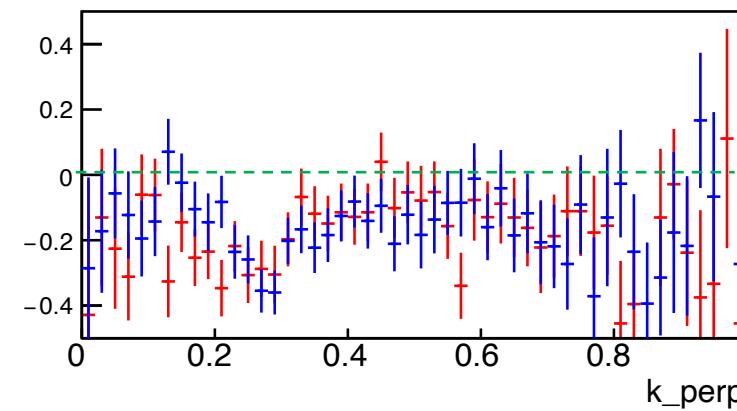
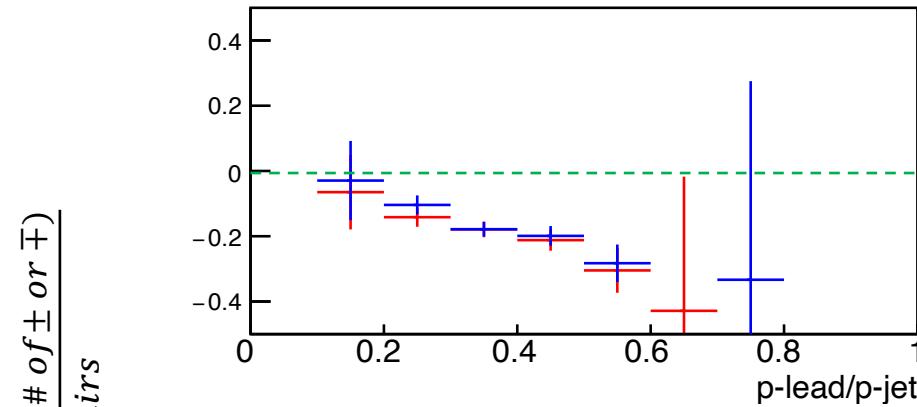


# $\pi K$ vs $K\pi$ - with p cut

(leading momentum particle is the first one)



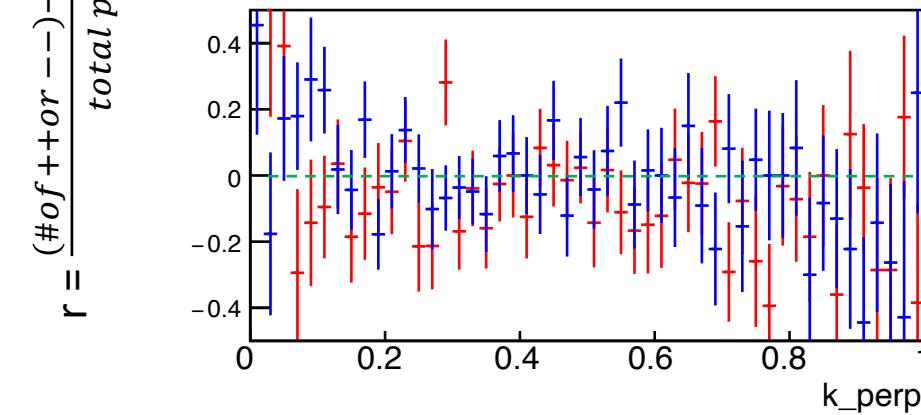
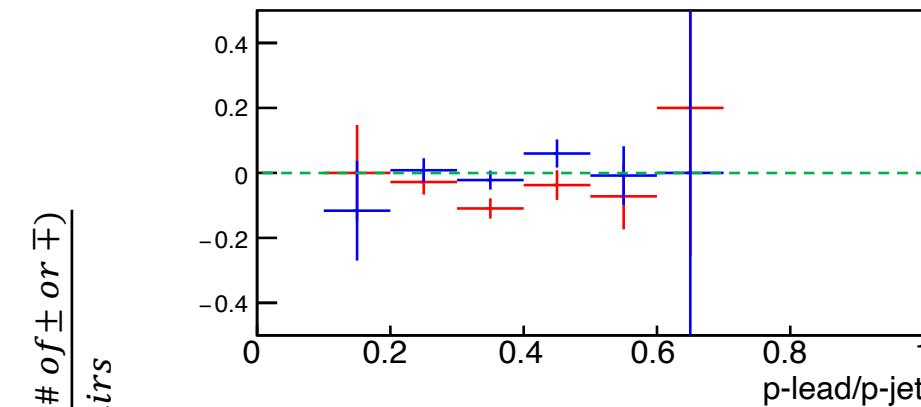
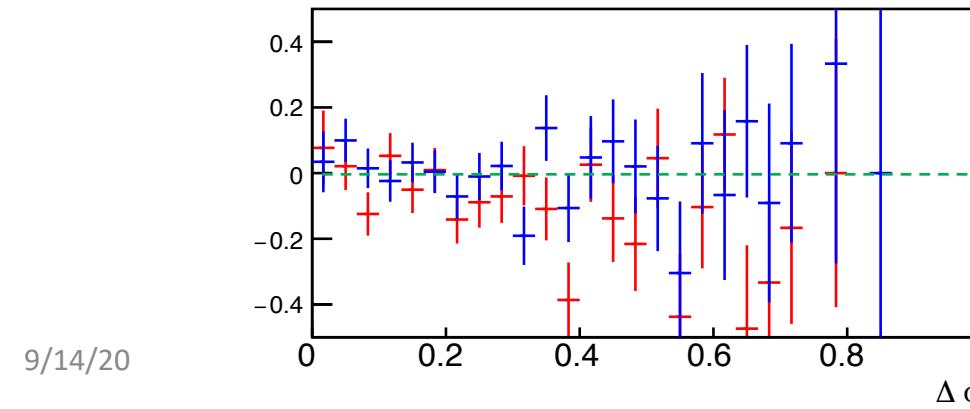
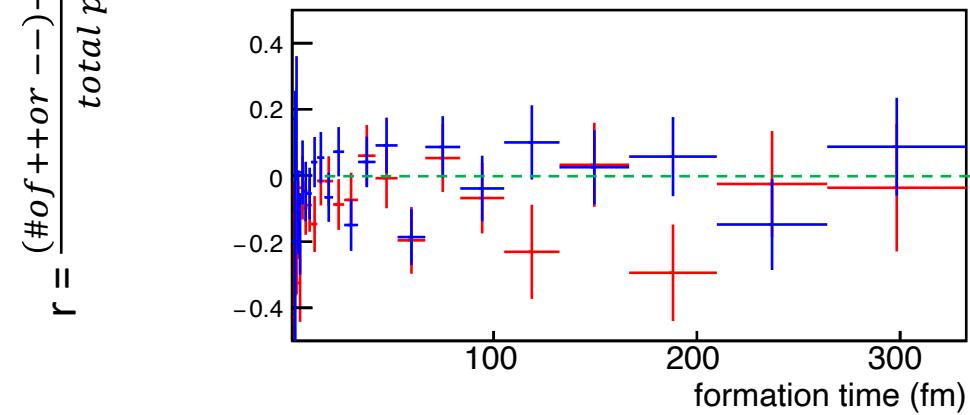
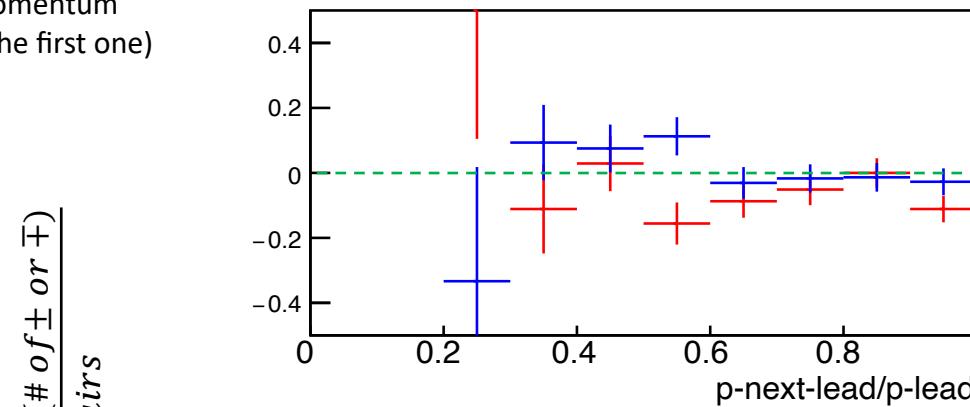
9/14/20



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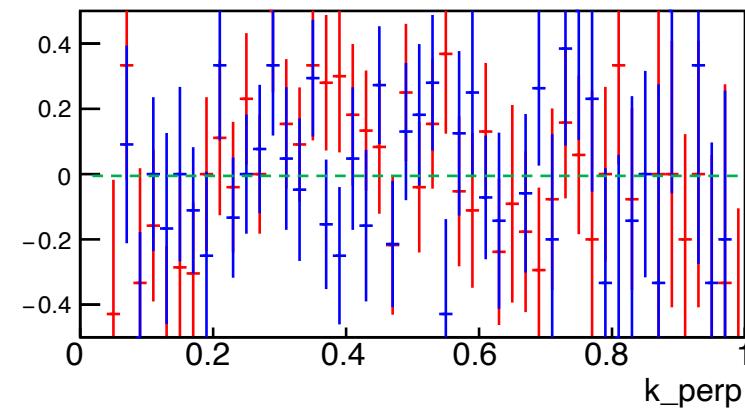
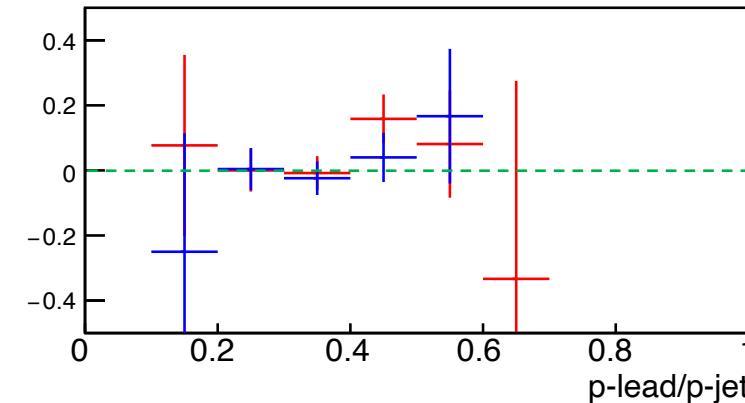
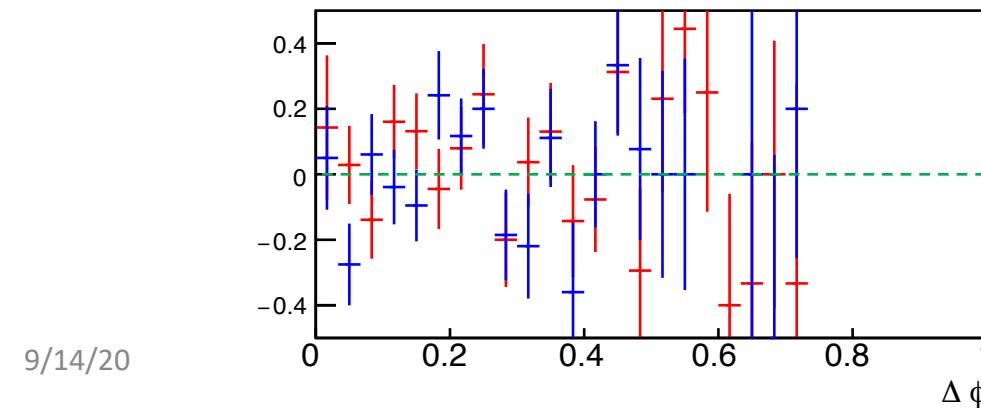
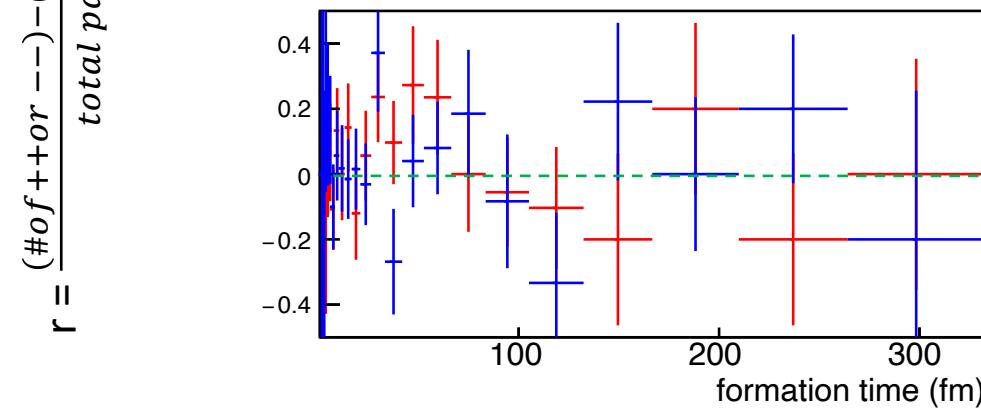
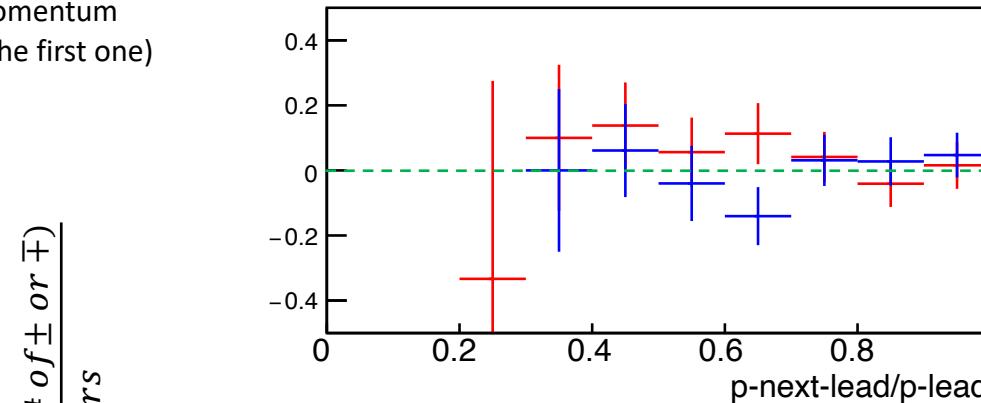
# $\pi P$ vs $P\pi$ - with p cut

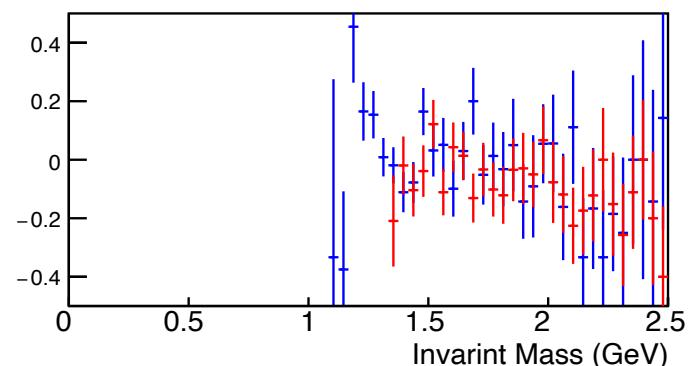
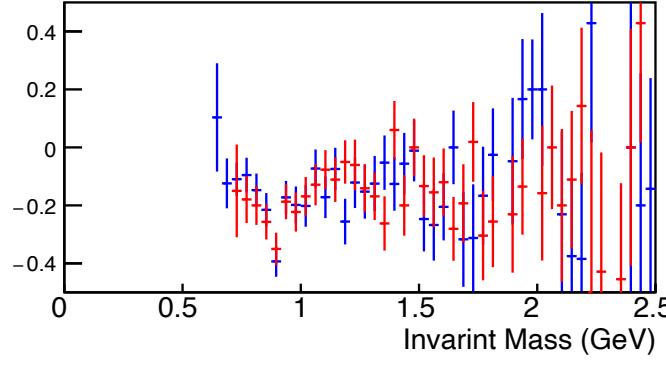
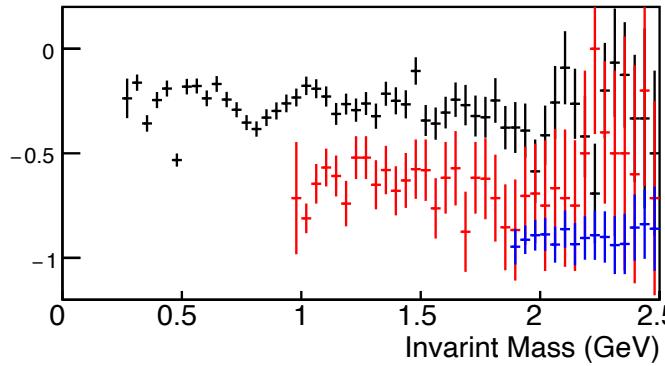
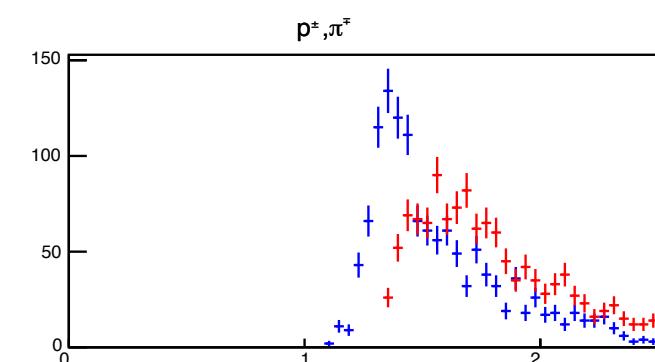
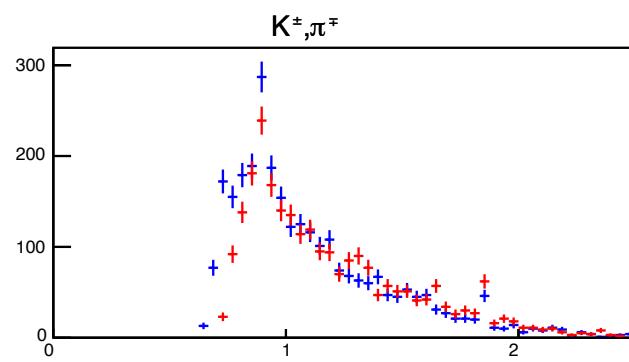
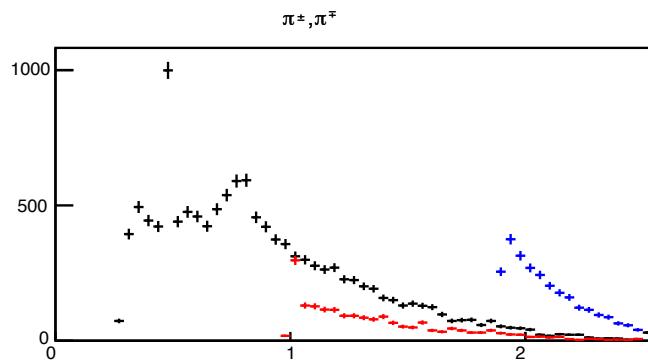
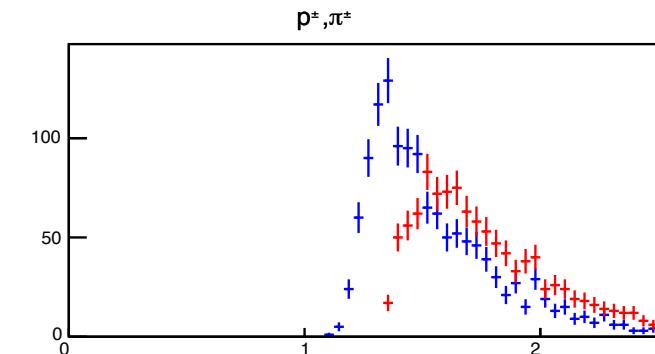
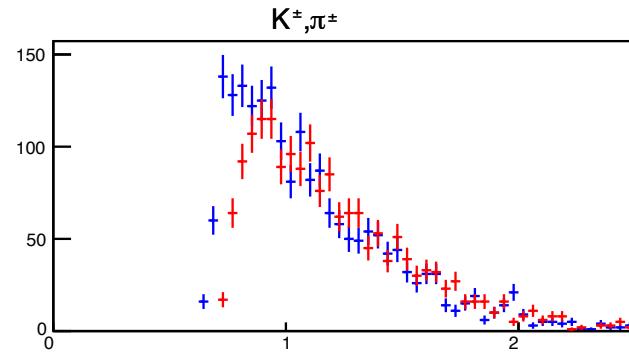
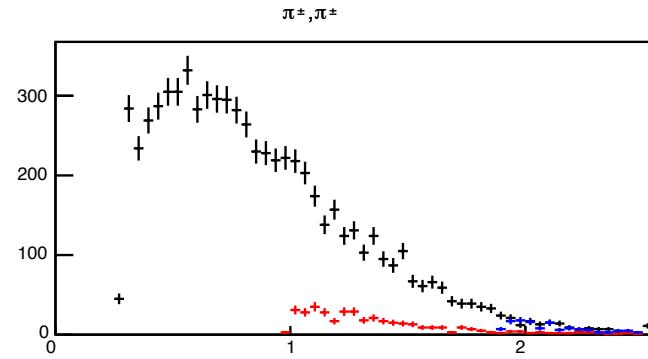
(leading momentum particle is the first one)



# K<sub>p</sub> vs pK - with p cut

(leading momentum particle is the first one)





$Q^2 = 65 \text{ GeV}^2$

Maximum Cross section =  $2.2 \times 10^4 \text{ mb}$

$2.2 \times 10^{-4} \text{ mb}$

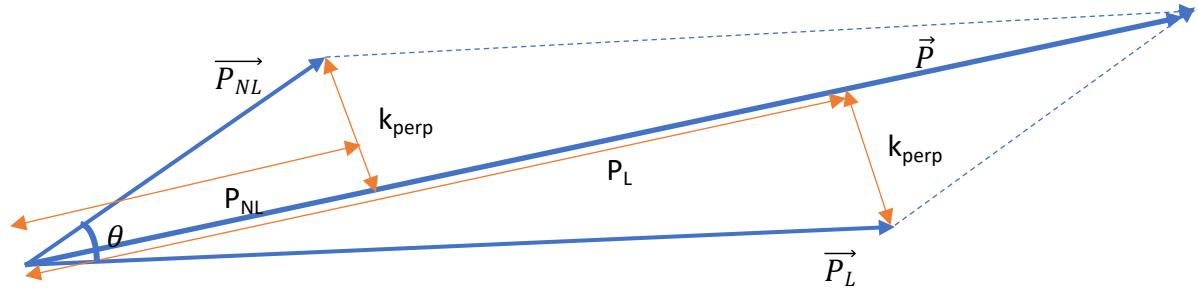
$2.2 \times 10^8 \text{ fb}$

Integrated luminosity =  $10 \text{ fb}^{-1}$

=  $2.2 \times 10^9 \text{ events}$

= 2200 M events

# Formation time



$$z = \frac{\vec{P}_{NL}}{\vec{P}_{NL} + \vec{P}_L}$$

$$\vec{P}_L = (1-z)\vec{P}$$

$$\vec{P}_{NL} = z\vec{P}$$

$$\text{formation time} = [2z(1-z) \vec{P}] / k_{\text{perp}}^2$$

What is the interpretation of  $\tau_{\text{form}}$ ?

It's the inverse of the "energy deficit" in the lab frame if we imagine that two massless particles with 3-momenta  $\vec{P}_L$  and  $\vec{P}_{NL}$  emerged from the decay of an off-shell massless particle with total momentum  $\vec{P}$ .

Long formation times should correspond to particles that "stayed together" for a long time before separating.

