



The Interplay of Baryon and Strangeness at the Beam Energy Scan

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Outline

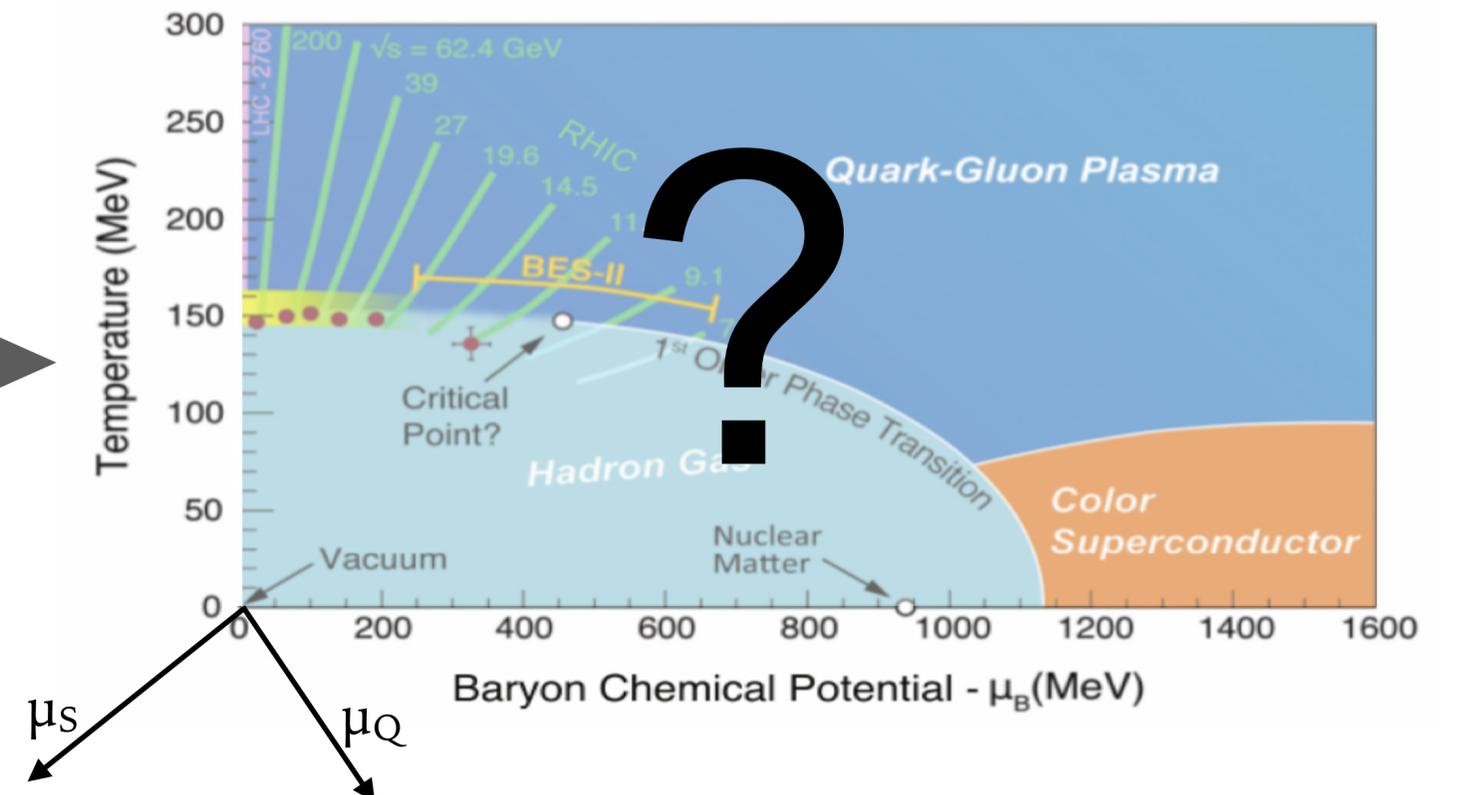
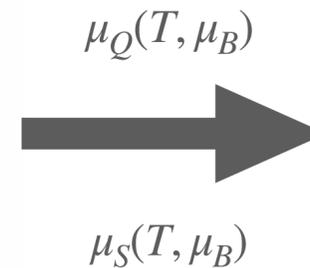
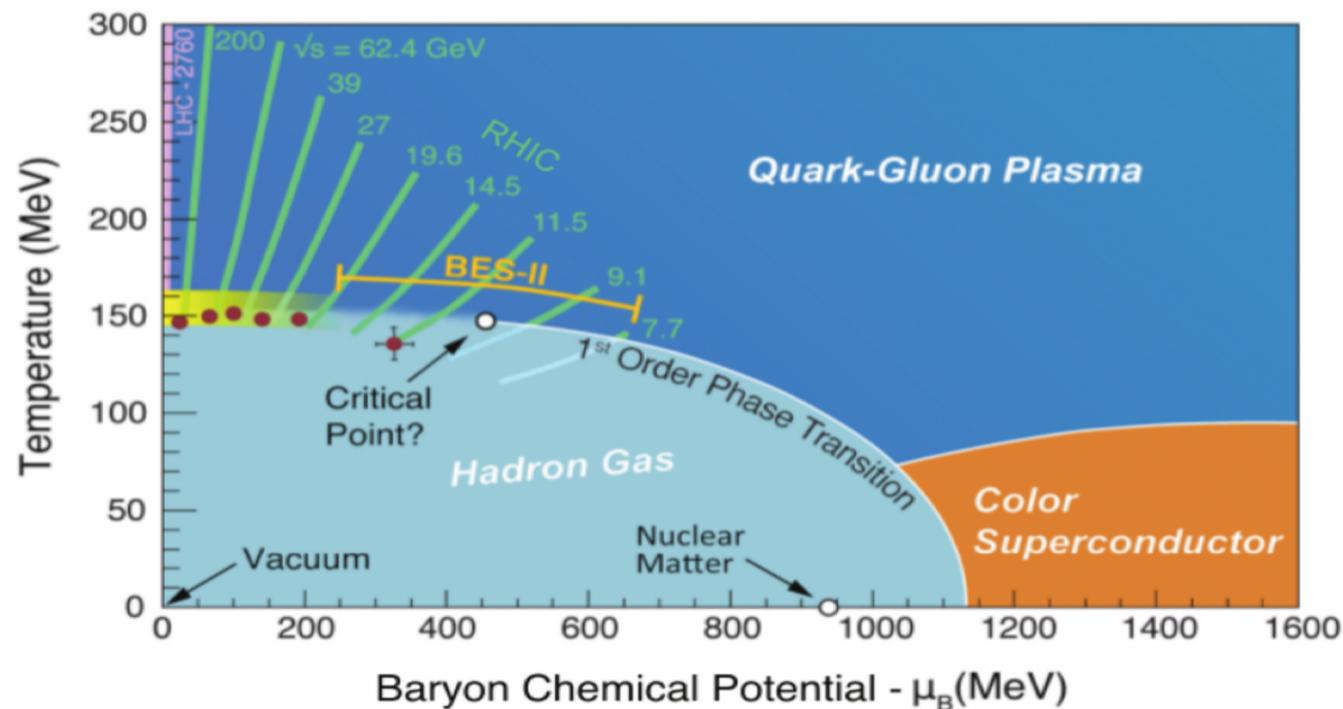
- I. Motivation for QCD modeling with all conserved charges
- II. Investigation of quark flavor hierarchy at chemical freeze-out
- III. Development of measurable cross-correlators of conserved charges
- IV. Recent model and theoretical approaches incorporating all conserved charges
- V. Conclusions

I. Motivation

QCD Phase Diagram

- ▶ The strongly interacting matter present in heavy-ion collisions carries a multitude of conserved quantum numbers: baryon number, strangeness and electric charge
 - ▶ This effects thermodynamics since each charge has an associated chemical potential

$$\frac{p}{T^4} = \frac{1}{VT^3} \ln Z(V, T, \mu_B, \mu_S, \mu_Q)$$



Lattice QCD Predictions

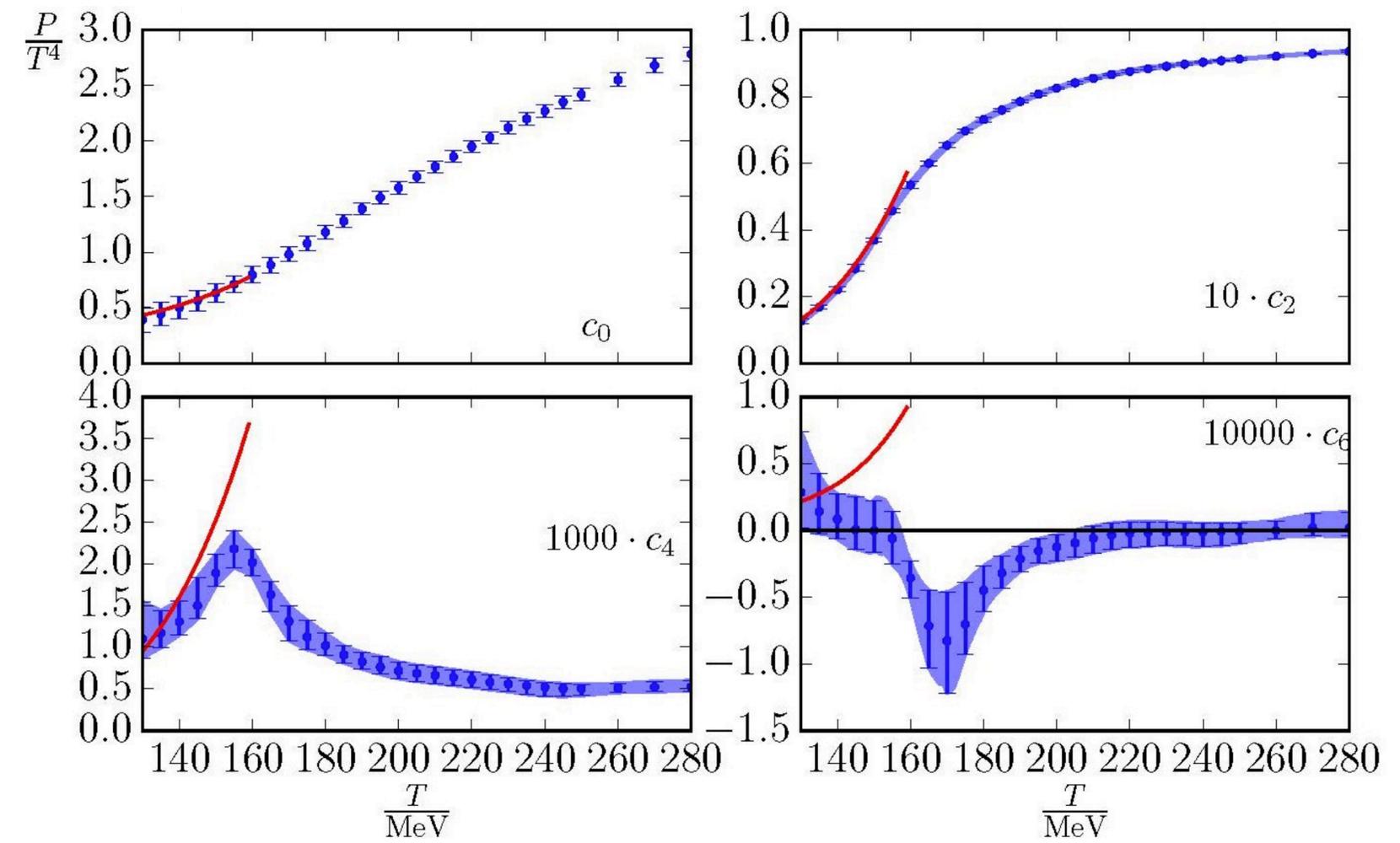
- The equation of state (EoS) for QCD has been calculated on the lattice under strangeness neutrality and fixed ratio of baryon number to electric charge, matching the heavy-ion situation

$$\frac{p(T, \mu_B)}{T^4} = \frac{p(T, 0)}{T^4} + \sum_{n=1}^{\infty} \frac{1}{(2n)!} \left. \frac{d^{2n}(p/T^4)}{d(\mu_B/T)^{2n}} \right|_{\mu_B=0} \left(\frac{\mu_B}{T} \right)^{2n}$$

$$= \sum_{n=0}^{\infty} c_{2n}(T) \left(\frac{\mu_B}{T} \right)^{2n}$$

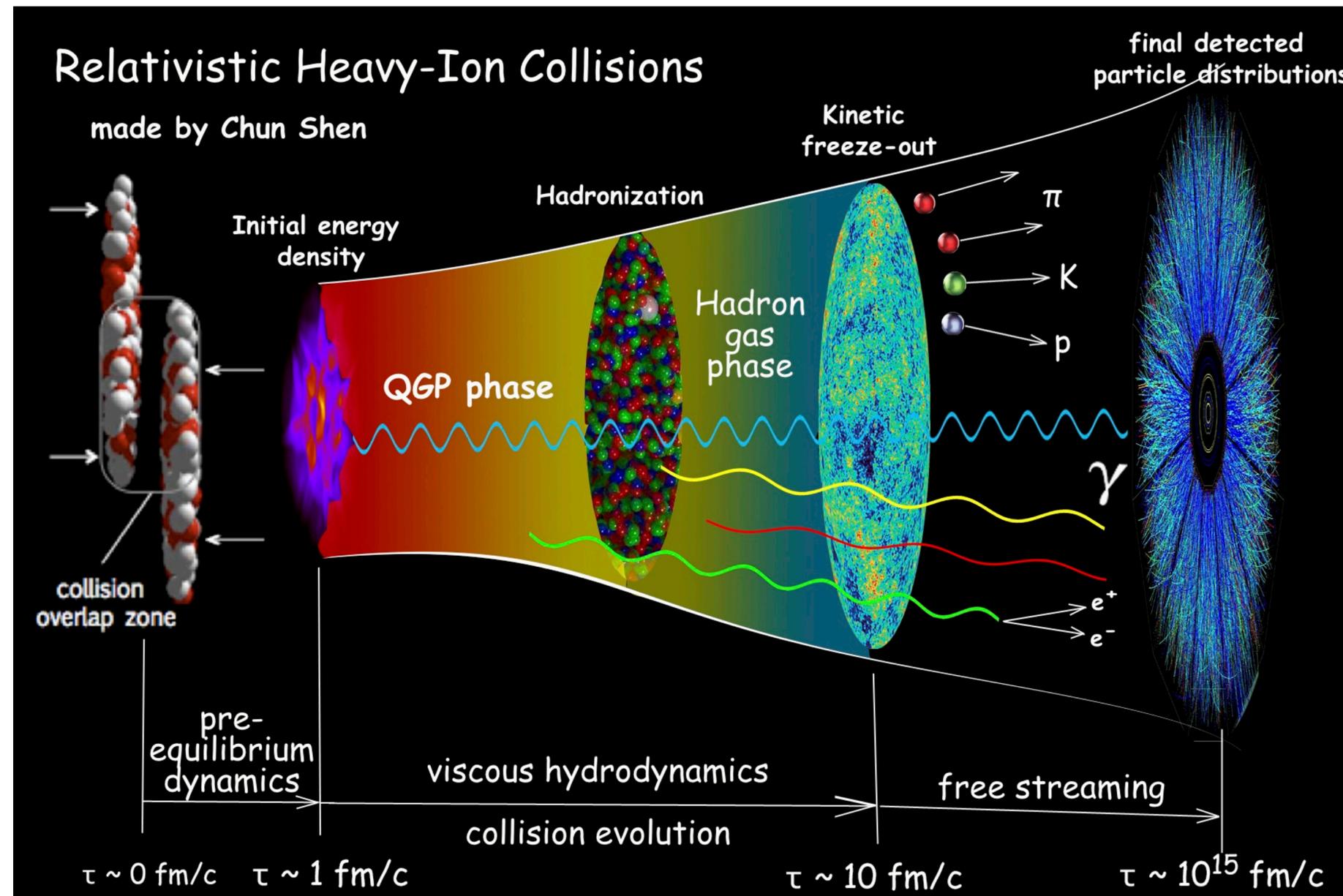
$$\langle n_s \rangle = 0$$

$$\langle n_Q \rangle = 0.4 \langle n_B \rangle$$



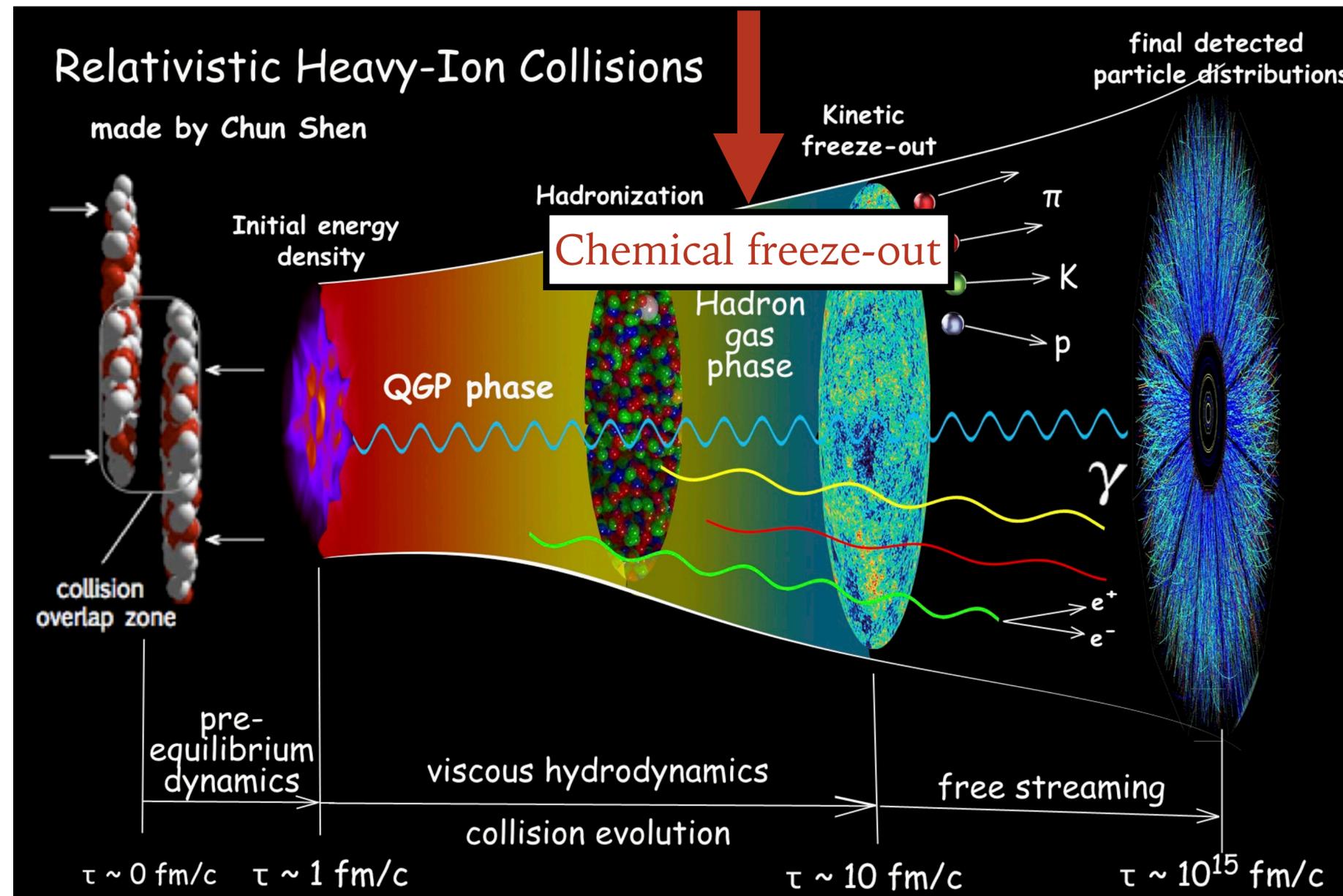
QCD Phenomenology

- All stages of heavy-ion collision modeling should seek to investigate the same slice of the phase diagram as the experiments



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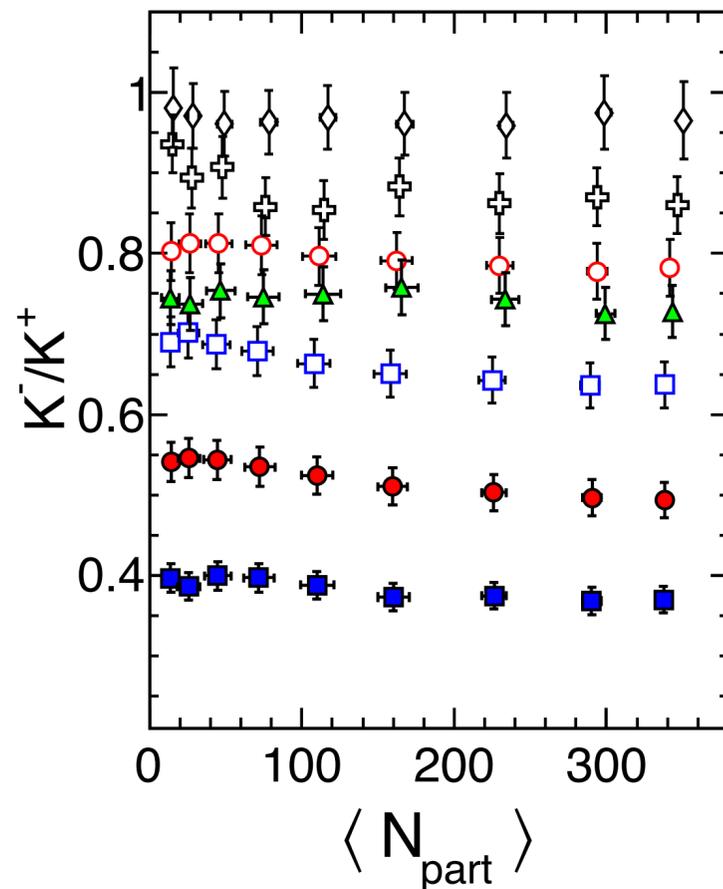


II. Flavor Hierarchy

Chemical Freeze-out Parameters

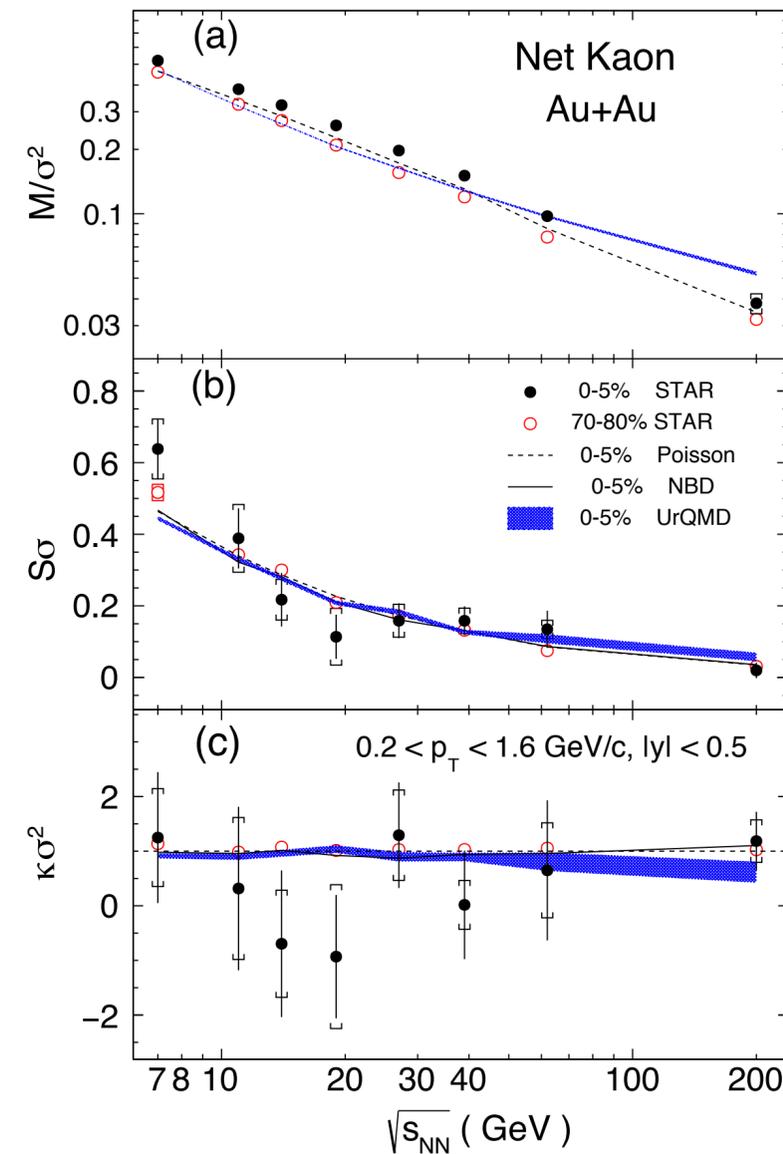
- Experimental particle yields and fluctuations tell us about the system at chemical freeze-out

Particle yields



- Ratios eliminate volume dependence

Fluctuations



- Moments can be calculated as derivatives of the pressure

Chemical Freeze-out Parameters

- Utilize Hadron Resonance Gas Model to perform thermal fits of yields and analyses of net-particle fluctuations including strangeness neutrality
- Investigate influence of number of states on freeze-out parameters

Particle yields



THERMAL-FIST: A package for heavy-ion collisions and hadronic equation of state

$$\frac{\chi^2}{N_{\text{dof}}} = \frac{1}{N_{\text{dof}}} \sum_{i=1}^N \frac{(N_i^{\text{exp}} - N_i^{\text{HRG}})^2}{\sigma_i^2}$$

$$\langle N_i \rangle = V n_i + V \sum_R \langle n_i \rangle_R n_R$$

Fluctuations

$$\chi_{lmn}^{BSQ} = \frac{\partial^{l+m+n} p / T^4}{\partial(\mu_B/T)^l \partial(\mu_S/T)^m \partial(\mu_Q/T)^n}$$

$$\text{mean : } M = \chi_1$$

$$\text{variance : } \sigma^2 = \chi_2$$

$$\text{skewness : } S = \chi_3 / \chi_2^{3/2}$$

$$\text{kurtosis : } \kappa = \chi_4 / \chi_2^2$$

$$\sigma^2 / M = \chi_2 / \chi_1$$

$$S \sigma^3 / M = \chi_3 / \chi_1$$

$$S \sigma = \chi_3 / \chi_2$$

$$\kappa \sigma^2 = \chi_4 / \chi_2$$

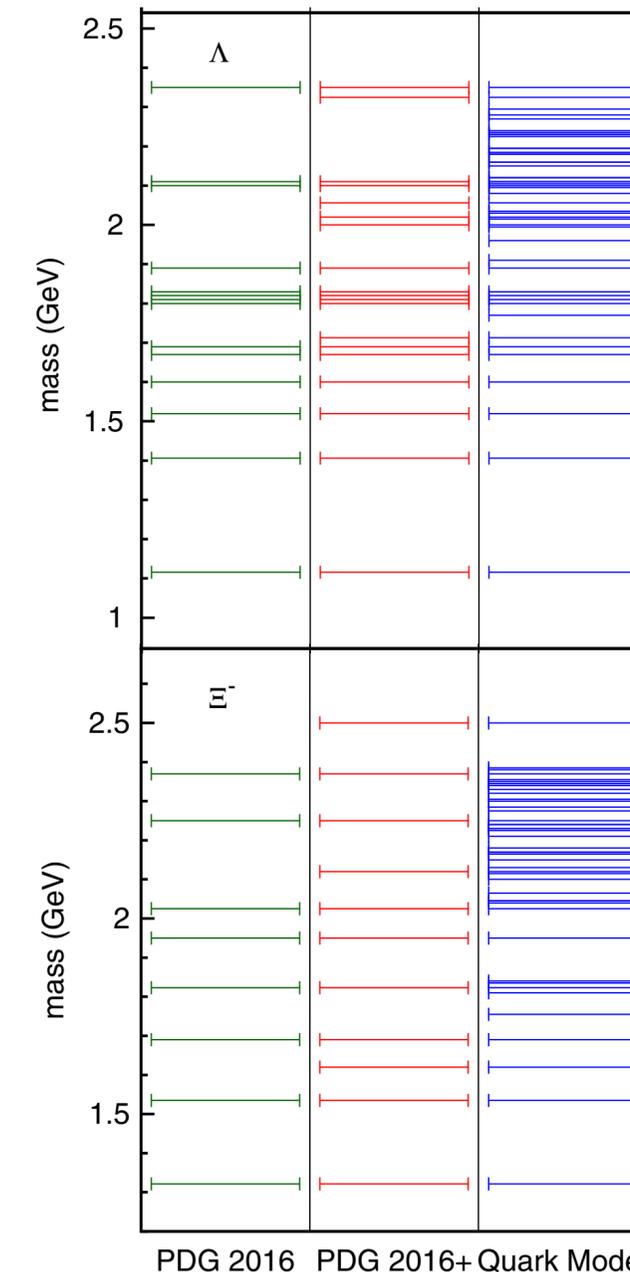
Hadronic States in HRG Model

- Pressure in HRG model depends on resonances included in the calculation:

$$\frac{P}{T^4} = \frac{1}{VT^3} \sum_i \ln Z_i(T, V, \vec{\mu})$$

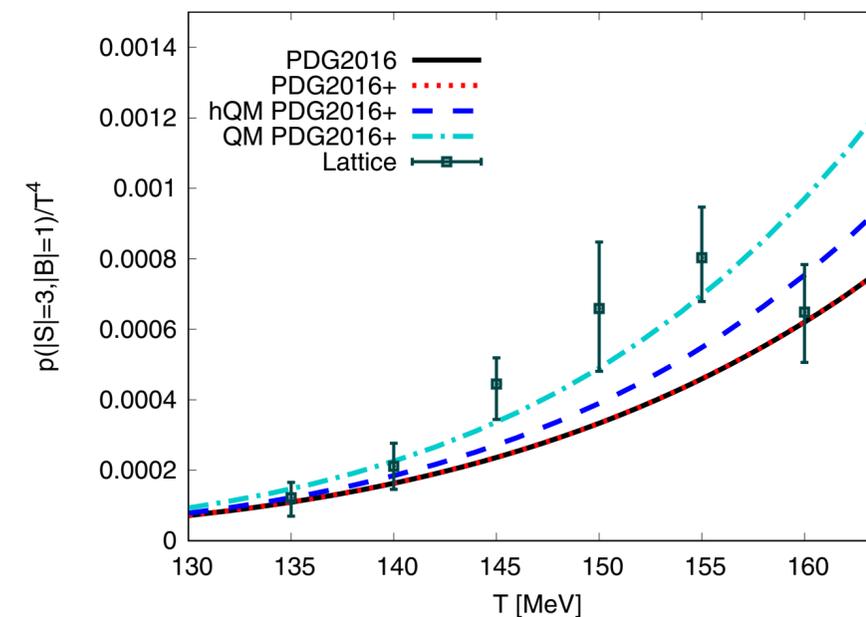
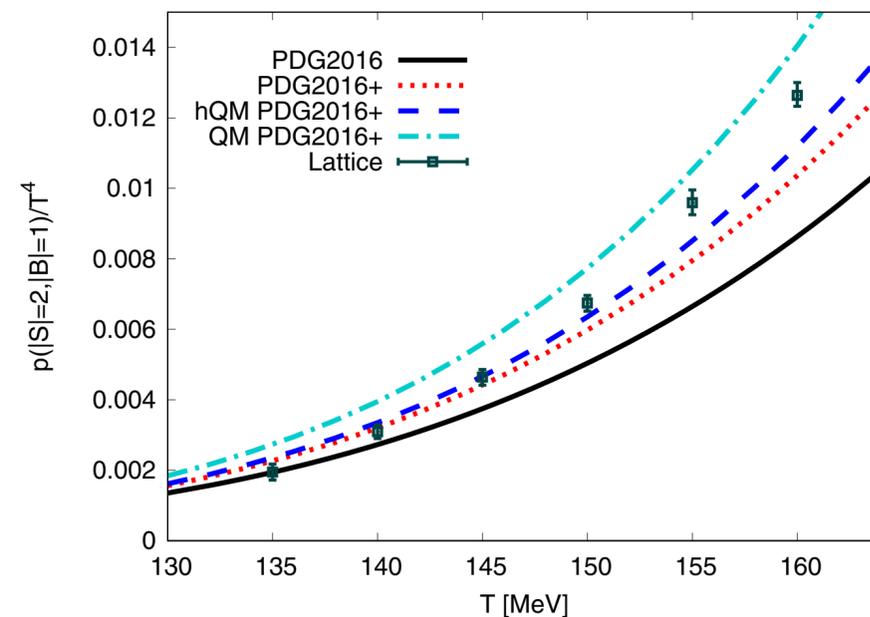
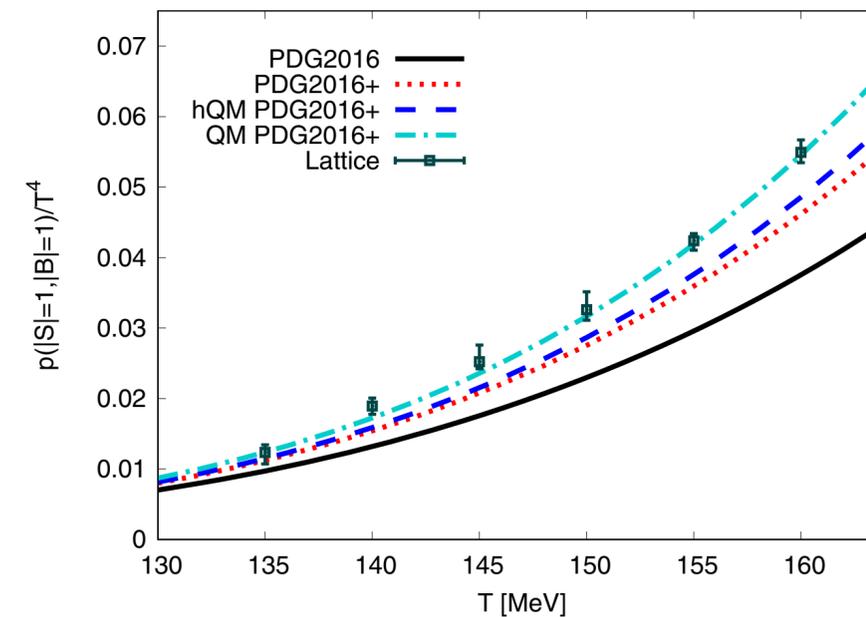
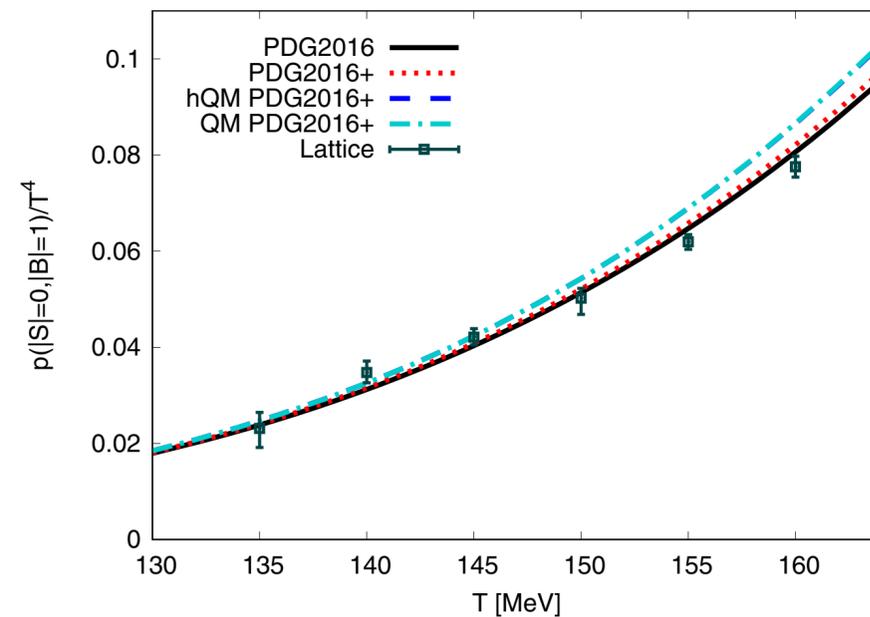
- PDG2005: 142 species
- PDG2016: 608 species
- PDG2016+: 738 species (all experimentally observed particles, i.e. *, **, ***, ****)
- QM: 1517 species (all states predicted by the Quark Model)

p	$1/2^+$	****
n	$1/2^+$	****
$N(1860)$	$5/2^+$	**
$N(1875)$	$3/2^-$	***
$\Delta(1232)$	$3/2^+$	****
$\Delta(1750)$	$1/2^+$	*



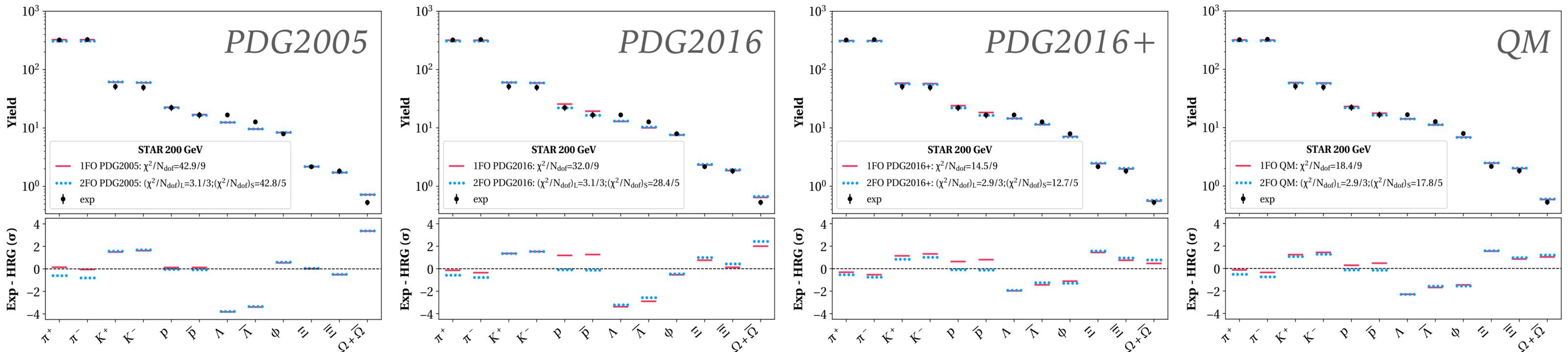
HRG Partial Pressures Confront LQCD

- Determine which particle list best matches Lattice QCD (LQCD) by breaking pressure into various sectors of baryon number and strangeness



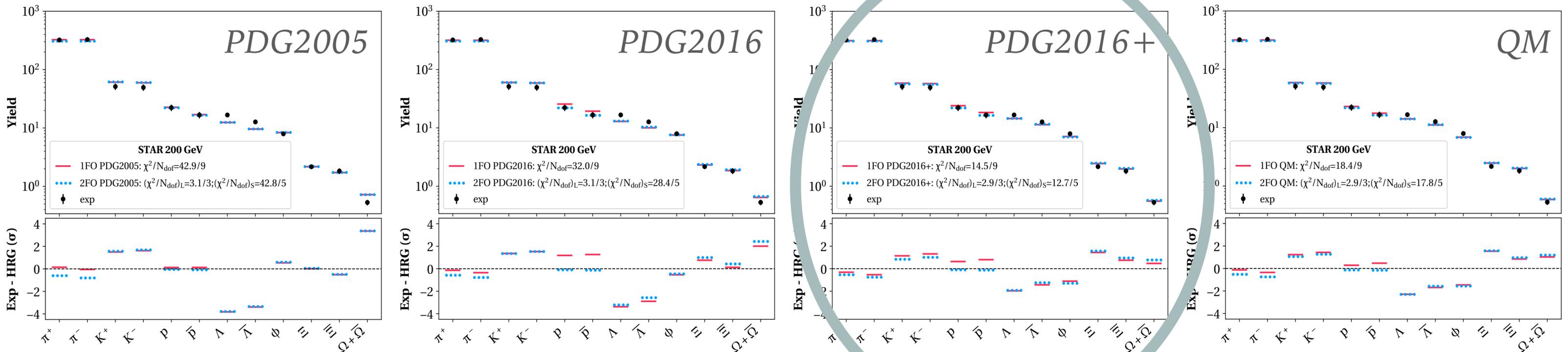
Thermal Fit Analyses

- Systematic study: chemical freeze-out dependence on number of resonances
- Flavor hierarchy investigation
 - Single freeze-out (1FO): simultaneous fit of all particles
 - Two flavor freeze-out (2FO): fit light and strange particles separately



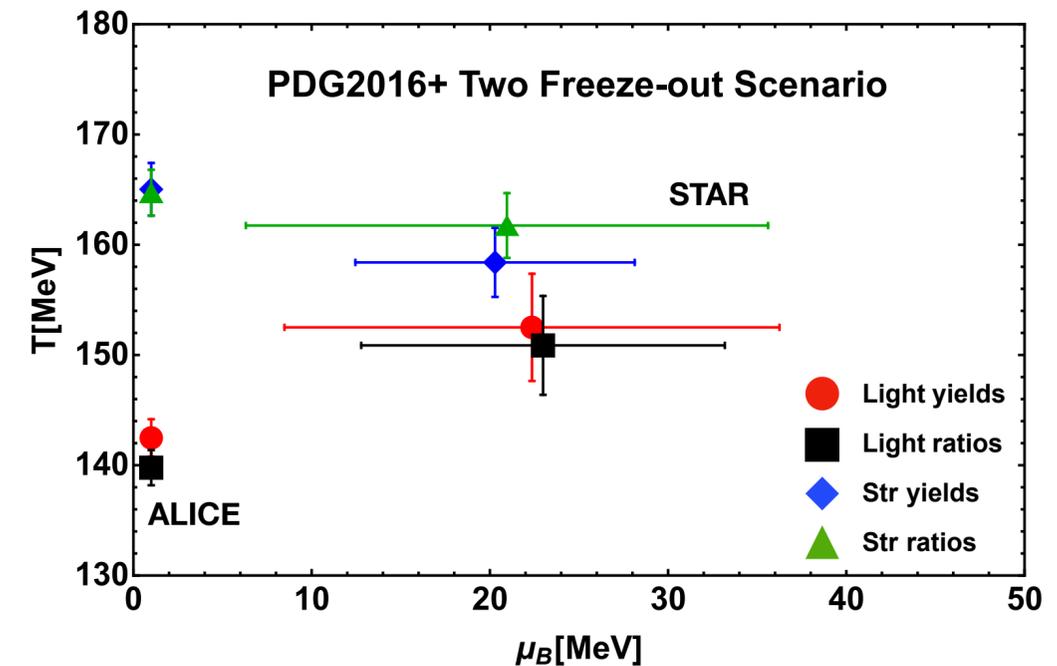
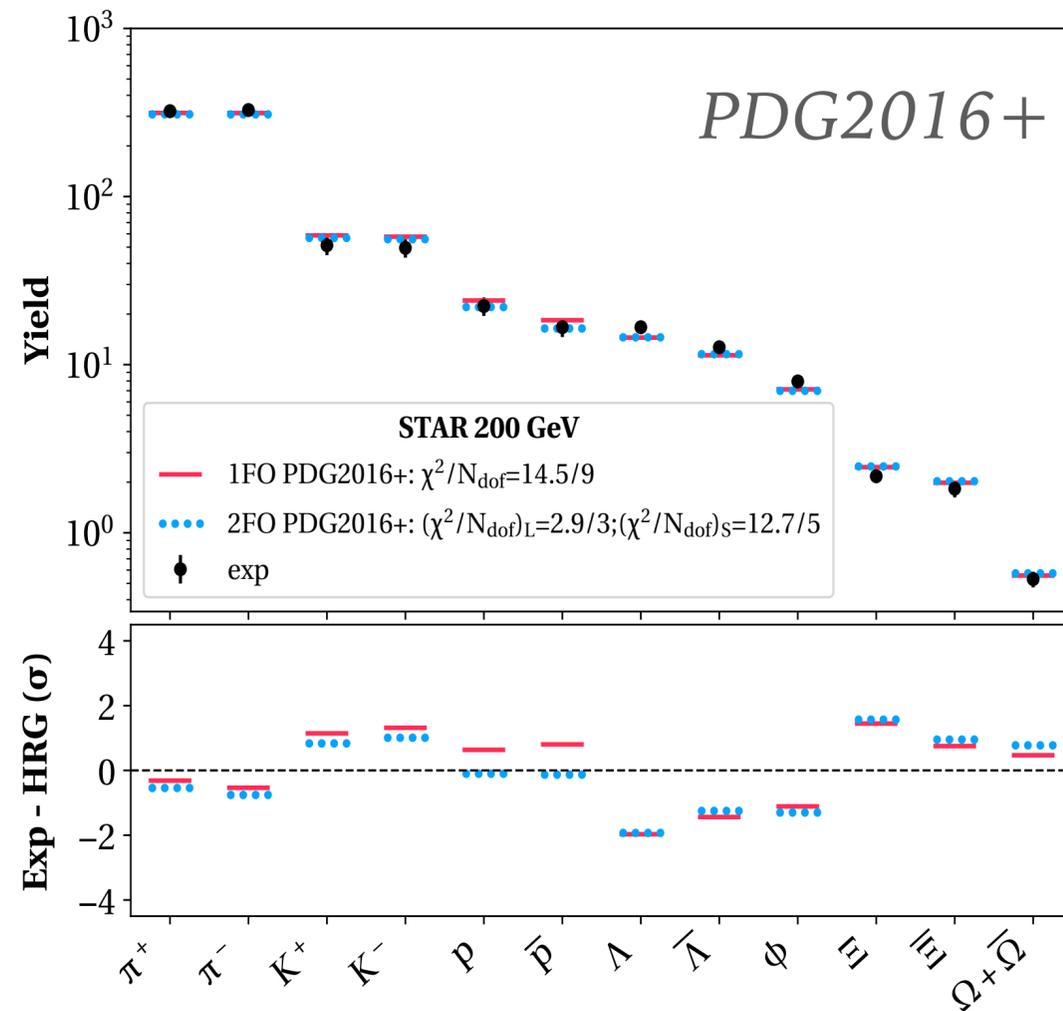
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Light v. Strangeness Freeze-out: Yields

- Unique chemical freeze-out parameters for light and strange even with many states
- Separating into light and strange particles produces an overall better fit



Separation between light and strange freeze-out temperatures:

$$T_{l,STAR} \simeq 150 \text{ MeV}$$

$$T_{s,STAR} \simeq 160 \text{ MeV}$$

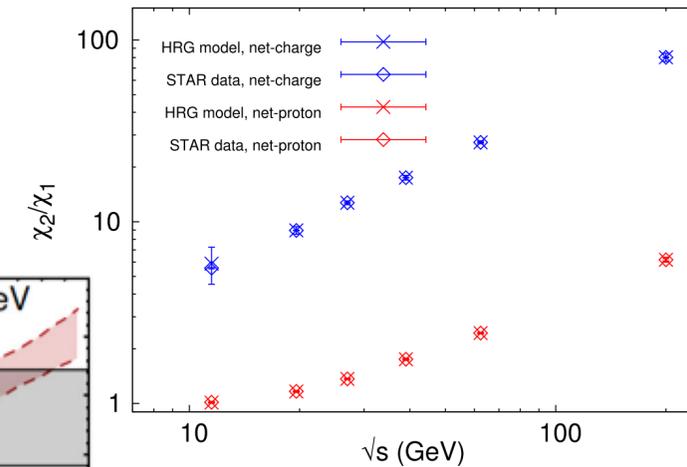
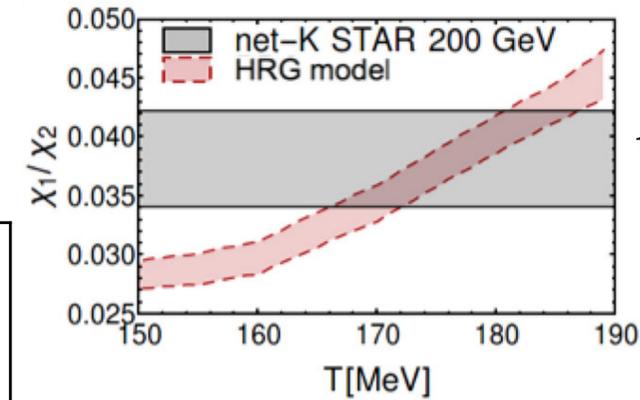
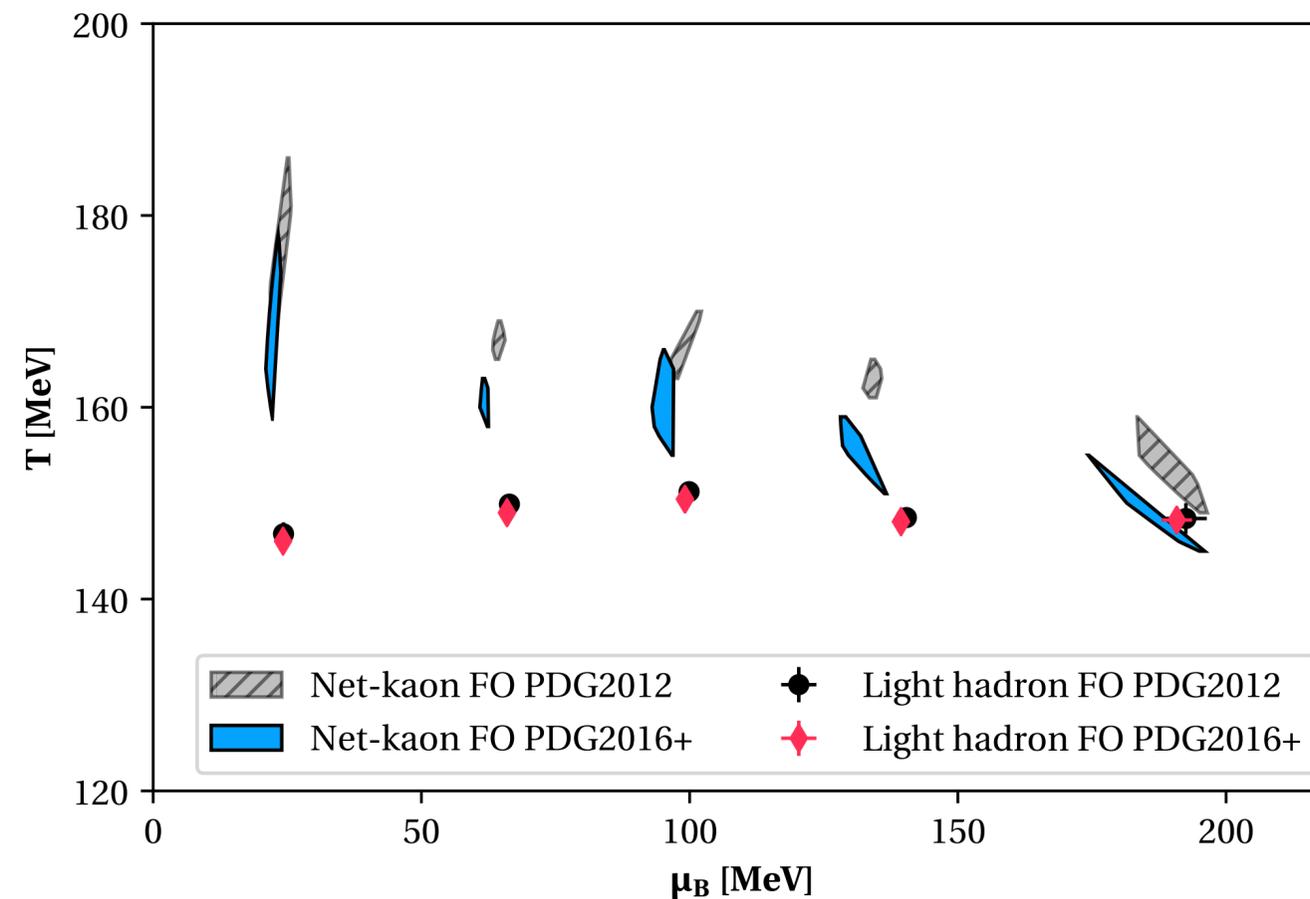
See also: J. Steinheimer et al, PRL (2013)
 F. Beccatini et al, PRL (2013)
 A. Bazavov et al, PRL (2014)

G. Aarts et al, PRD (2019)
 A. Andronic et al, Phys.Lett. B (2019)

P. Alba, JS et al, PRC (2020)
 R. Bellwied et al, arXiv: 2009.14781

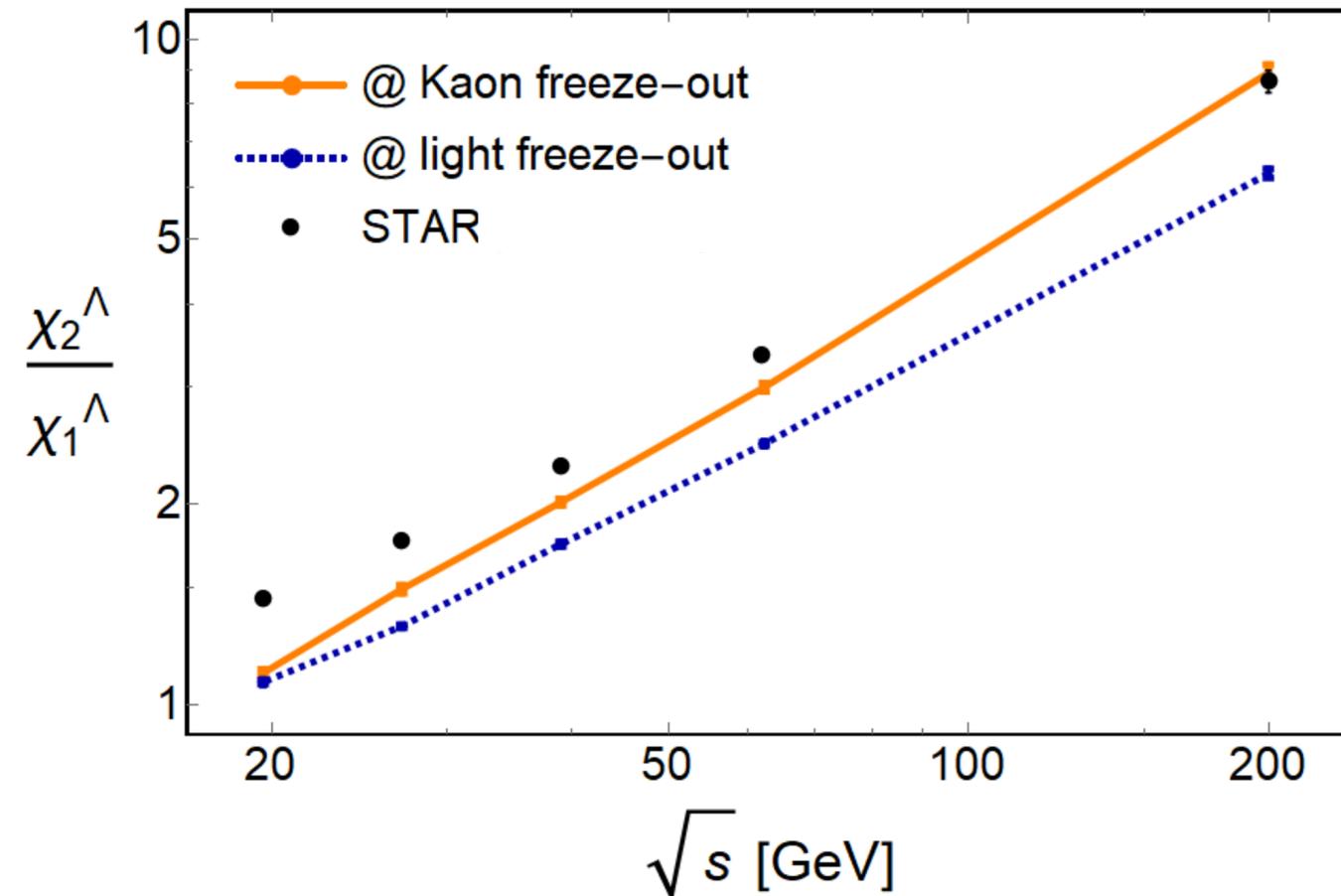
Light v. Strangeness Freeze-out: Fluctuations

- Freeze-out parameters calculated via net-particle fluctuations are *different* for p and K
 - Light freeze-out determined by combined fit of χ_1^p/χ_2^p and χ_1^Q/χ_2^Q
 - Calculate χ_1^K/χ_2^K along the Lattice QCD isentropes



Light v. Strangeness Freeze-out: Fluctuations

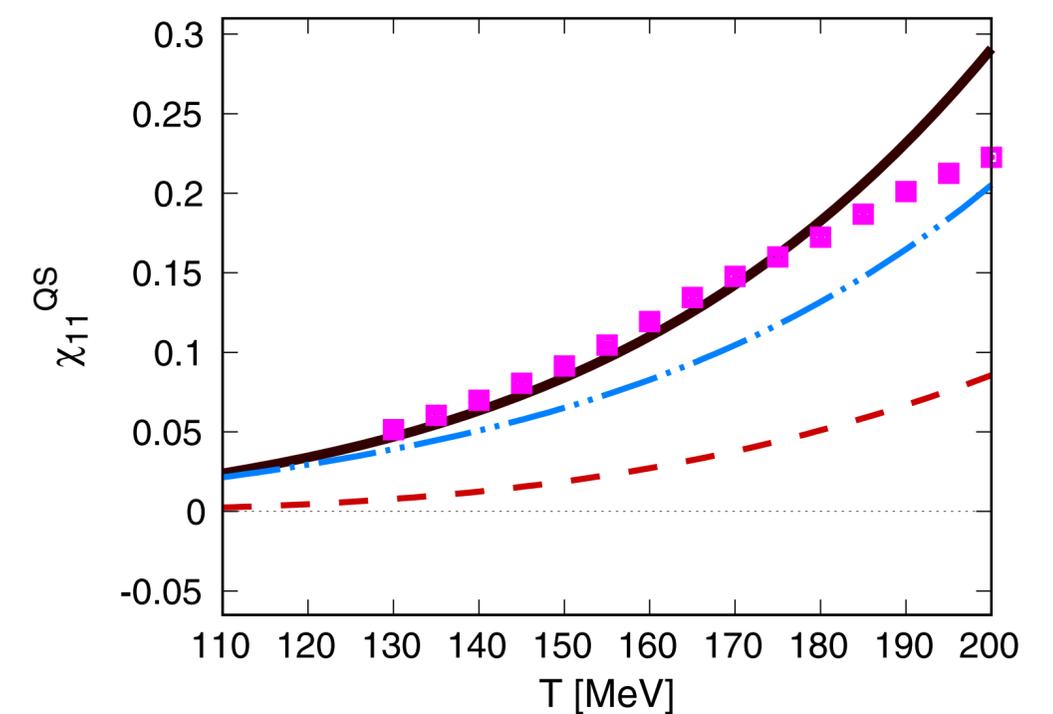
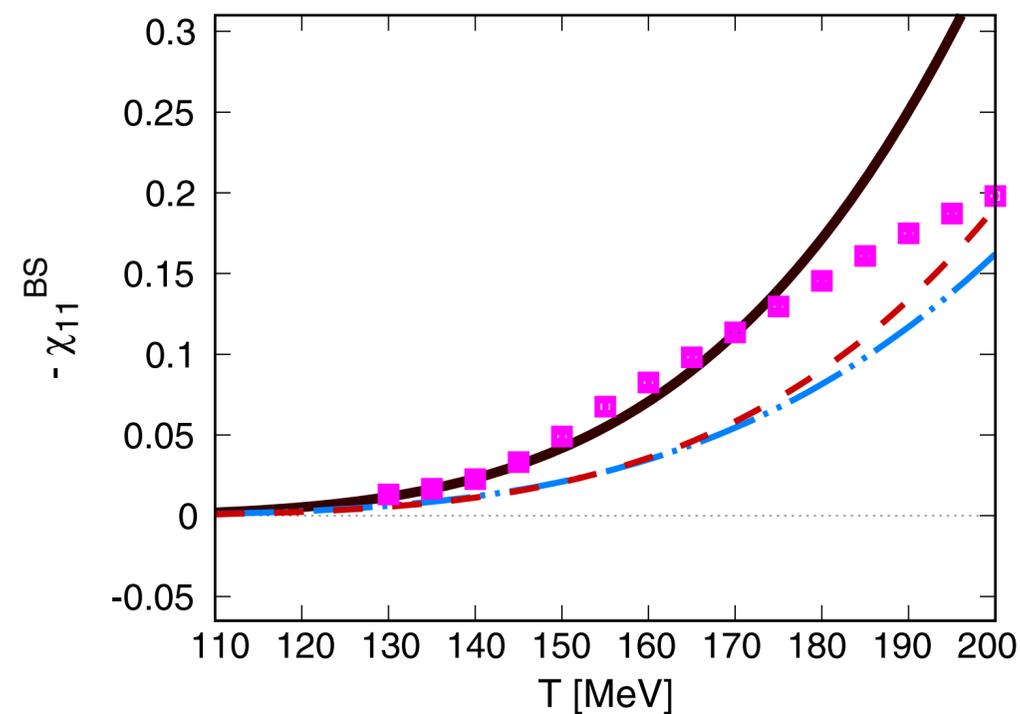
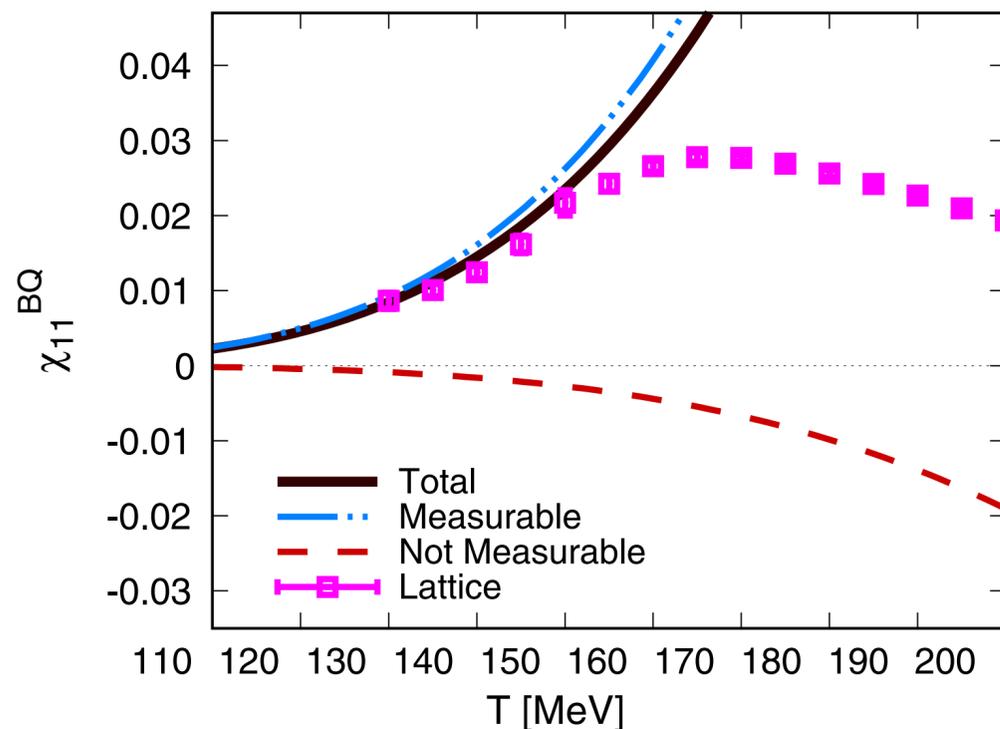
- Freeze-out curves for light particles and net-kaons compared to net- Λ fluctuations shows this strange baryon prefers to freeze-out with strange particles



III. Cross-correlators

Off-diagonal Correlators of Conserved Charges

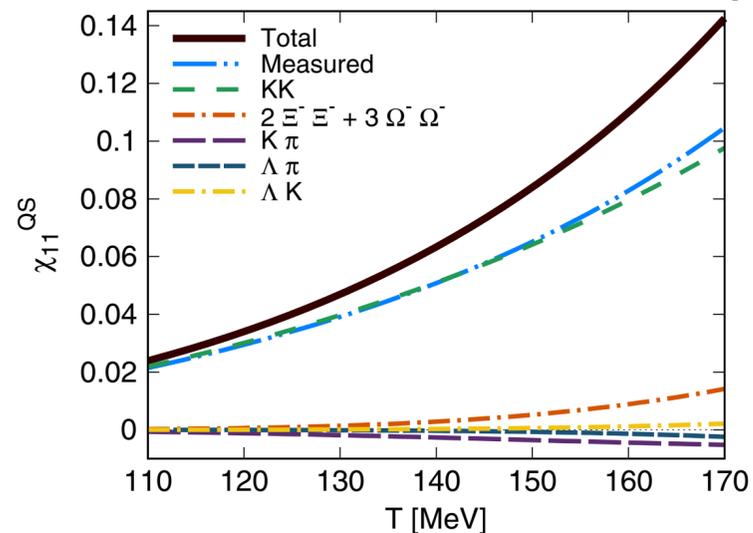
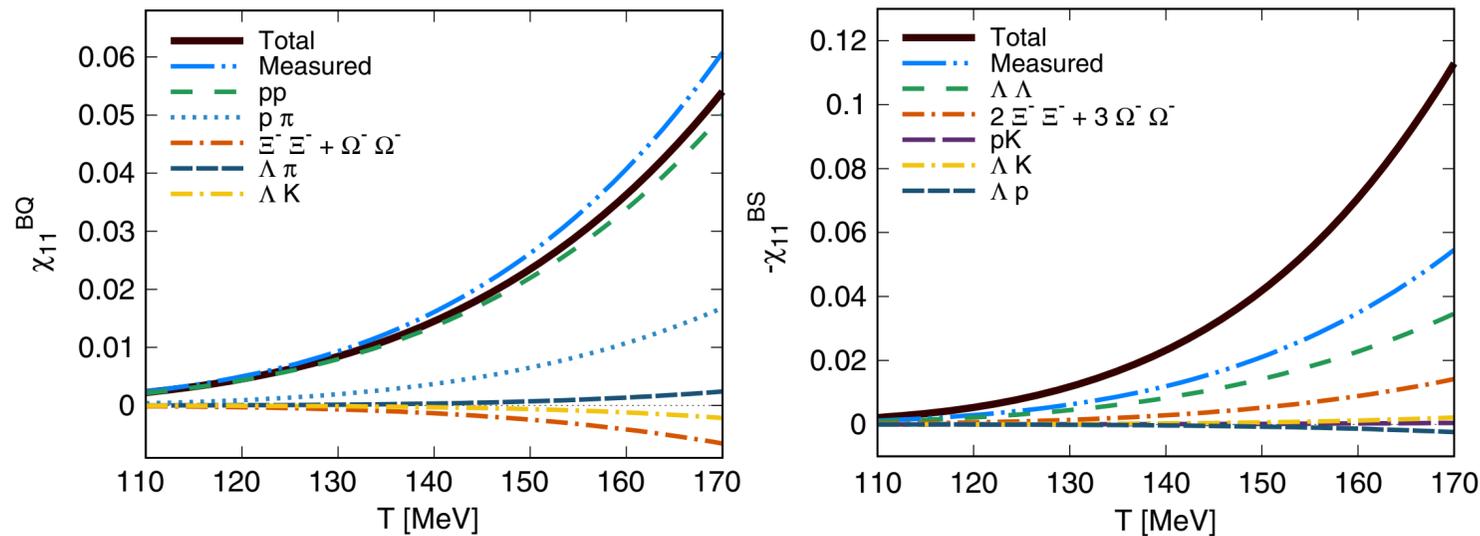
- The contribution of individual species to correlators on the lattice can be determined by comparing to the HRG model
- Main goal: develop proxies for correlators of conserved charges that can be measured experimentally



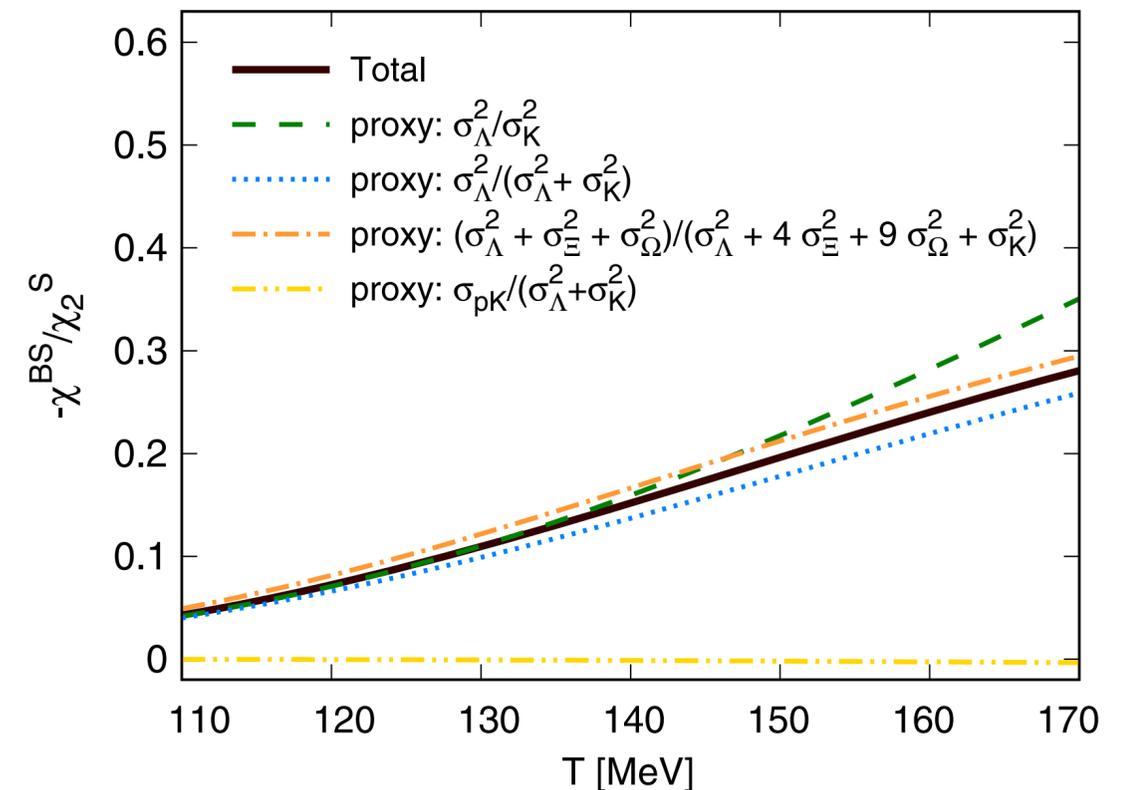
Measurable: π^+ , π^- , K^+ , K^- , p , \bar{p} , Λ , $\bar{\Lambda}$, Ξ^- , Ξ^+ , Ω^- , Ω^+

Off-diagonal Correlators of Conserved Charges

- In reality, there are only a few terms that makeup the overwhelming majority of a correlator
- Determine proxy by constructing ratios with the most dominate species

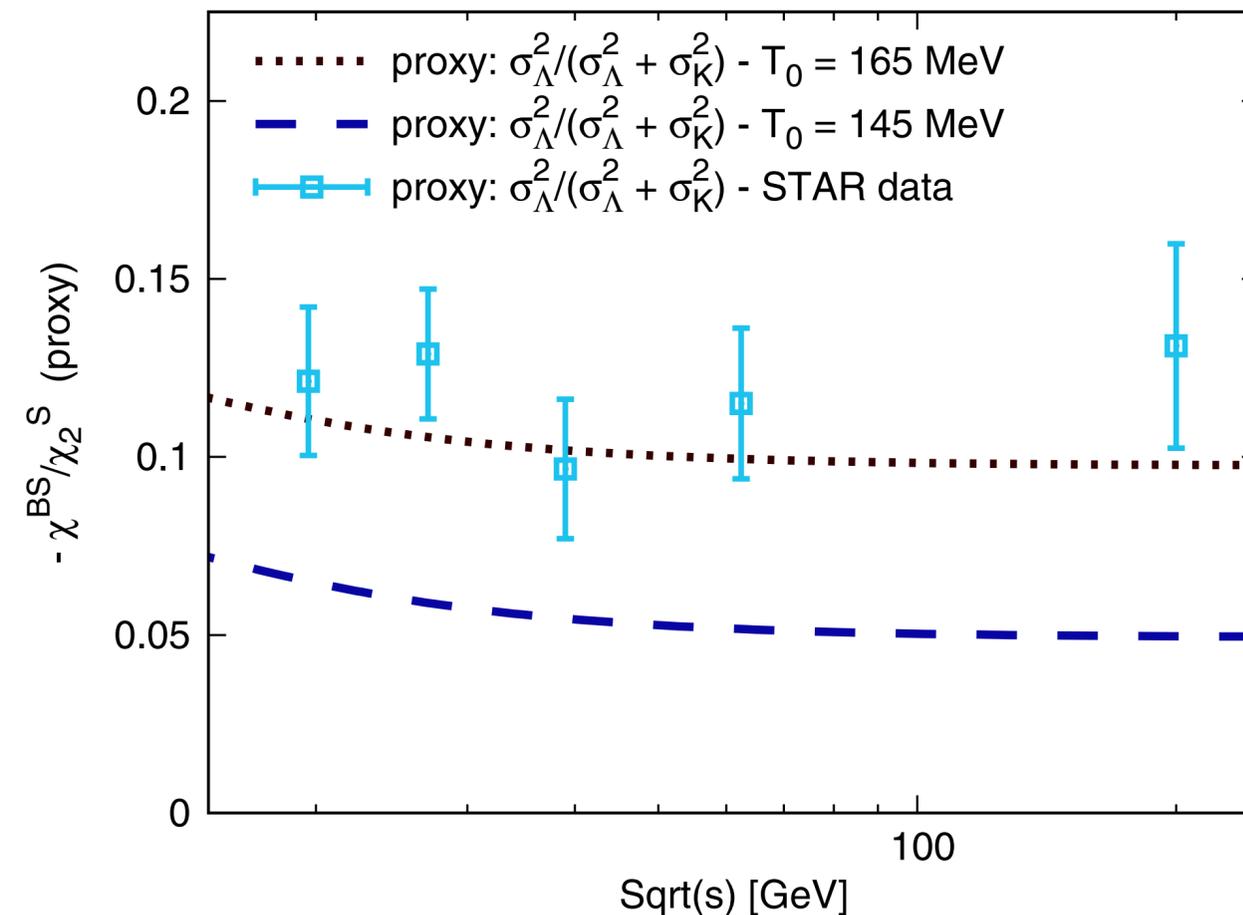


Proxy to measure **baryon-strangeness correlator**



Cross-correlators at Freeze-out

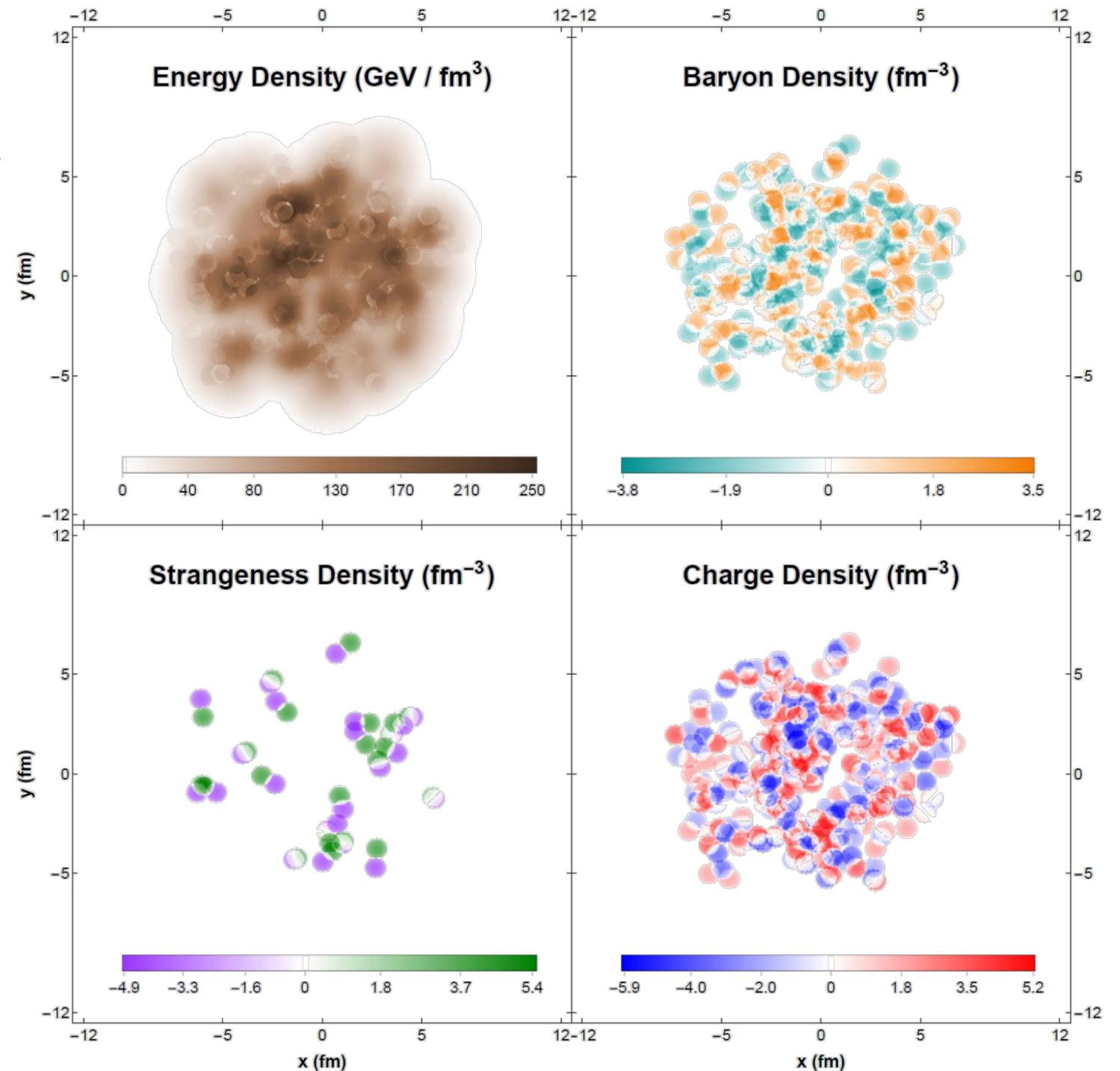
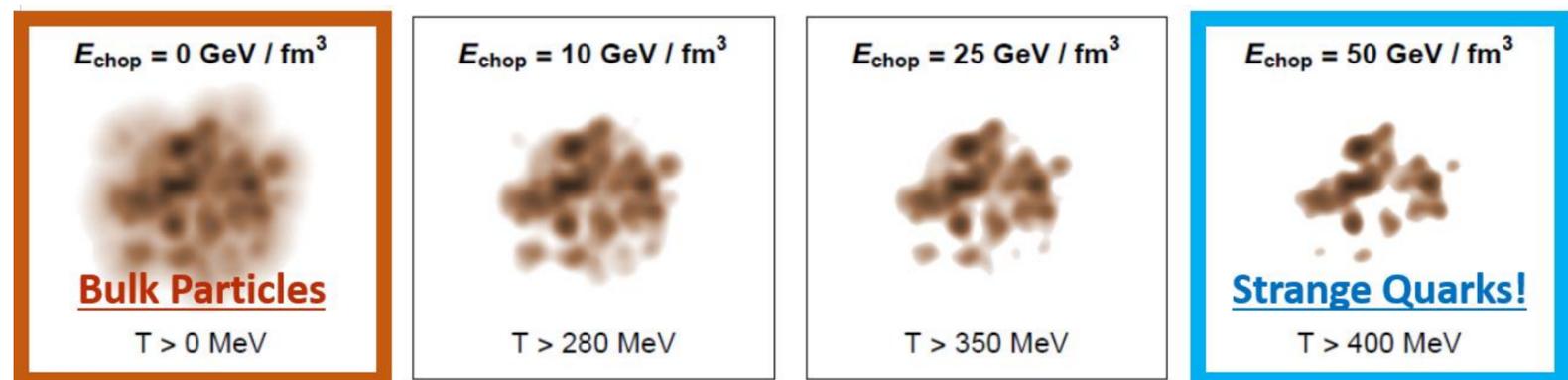
- ▶ The baryon-strangeness correlator can test the flavor hierarchy
 - ▶ Lines show parametrization of proxy on upper and lower bounds of crossover
 - ▶ Experimental data favor higher temperature



IV. BQS Modeling

BQS Initial Conditions

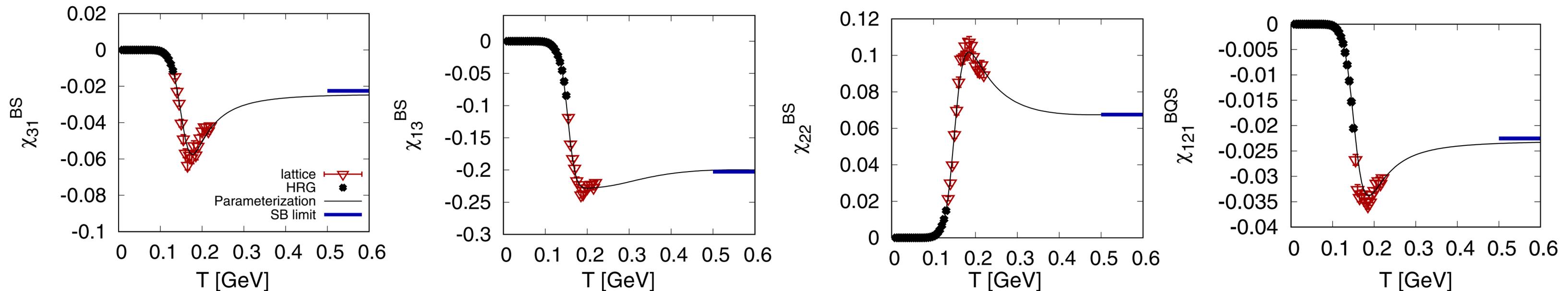
- ▶ Initial densities for each conserved charge can be constructed by re-sampling the initial energy density to map gluon density to $q\bar{q}$ production
- ▶ Baryon and charge densities largely mirror energy density
- ▶ Strangeness density is related to distribution of hot spots



BQS EoS

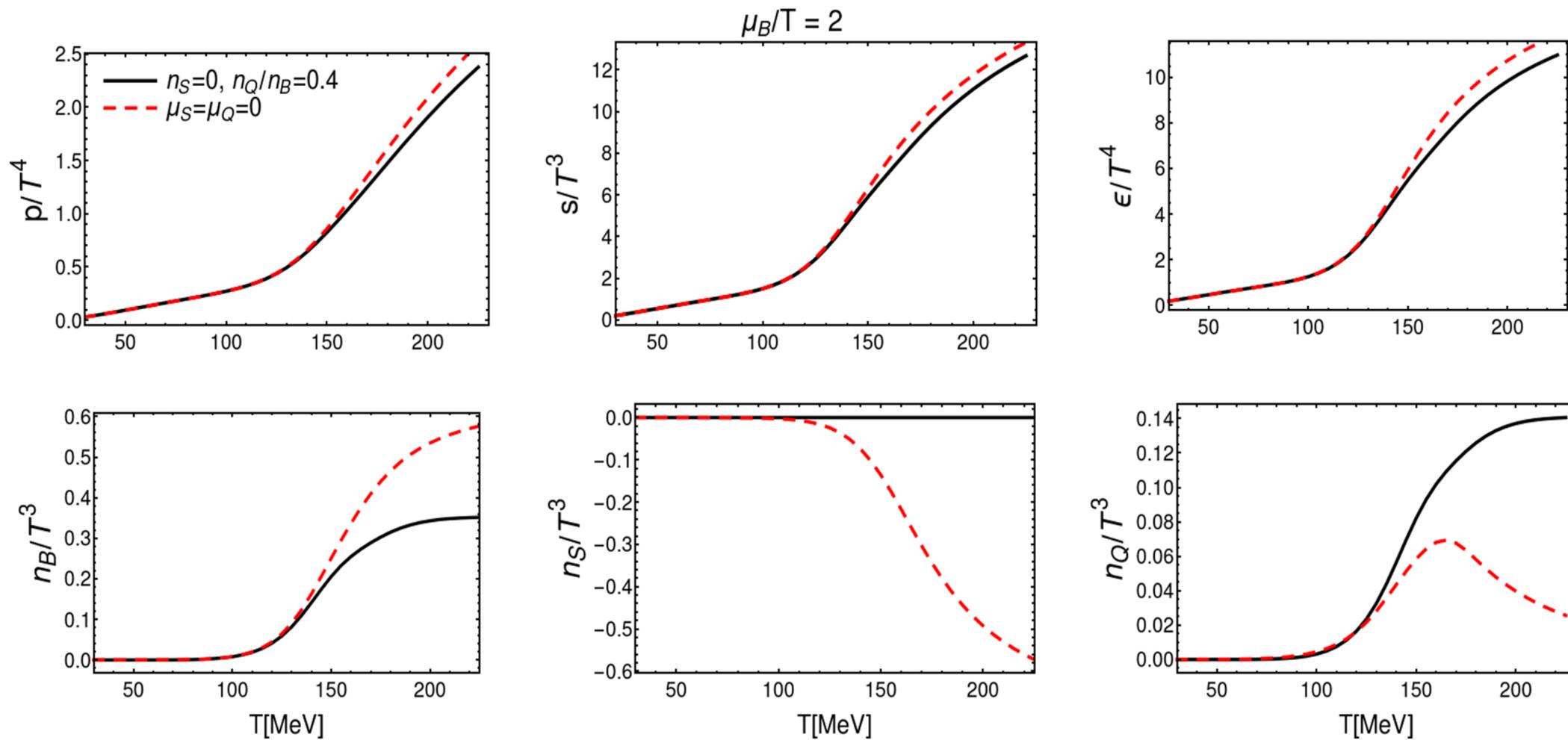
- Reconstruct the QCD equation of state from all diagonal and off-diagonal Taylor expansion coefficients up to $\mathcal{O}(\mu_B^4)$

$$\frac{P(T, \mu_B, \mu_Q, \mu_S)}{T^4} = \sum_{i,j,k} \frac{1}{i!j!k!} \chi_{ijk}^{BQS}(T) \left(\frac{\mu_B}{T}\right)^j \left(\frac{\mu_Q}{T}\right)^k \left(\frac{\mu_S}{T}\right)^i$$

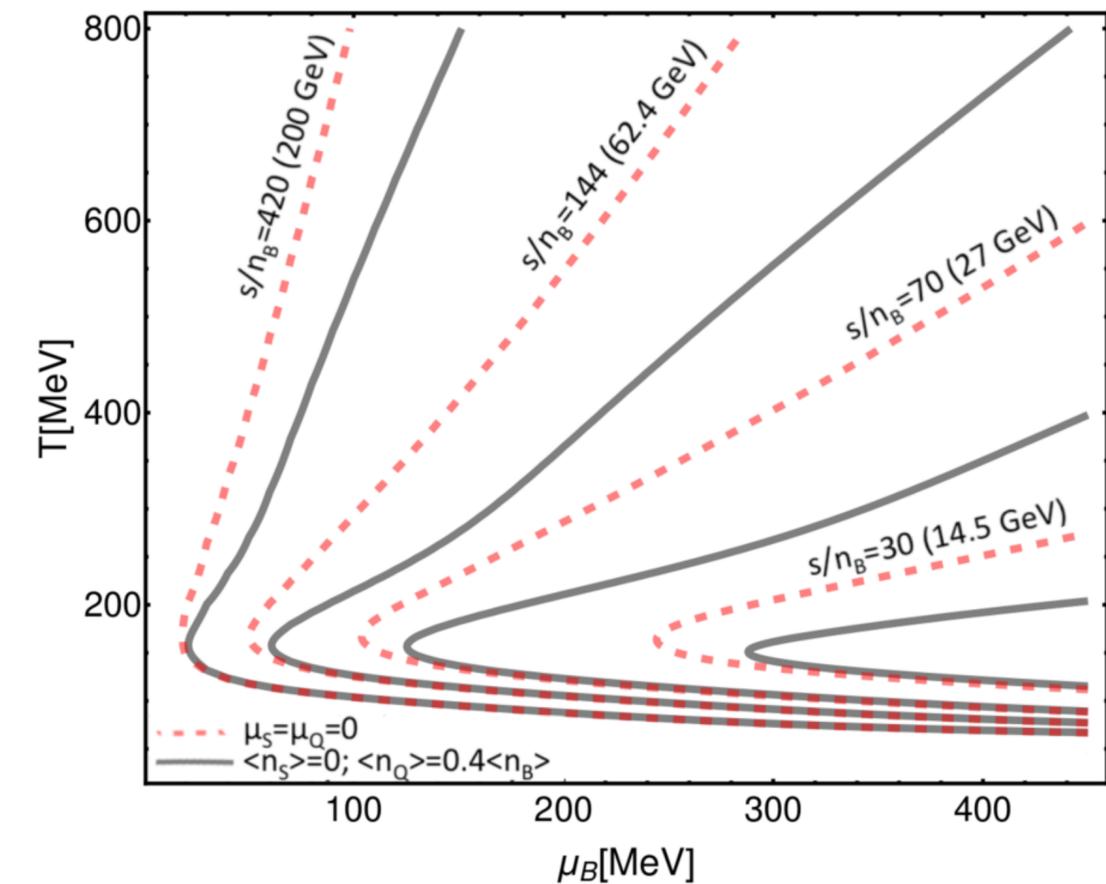


BQS EoS

- Charge conservation conditions can be applied in BQS EoS

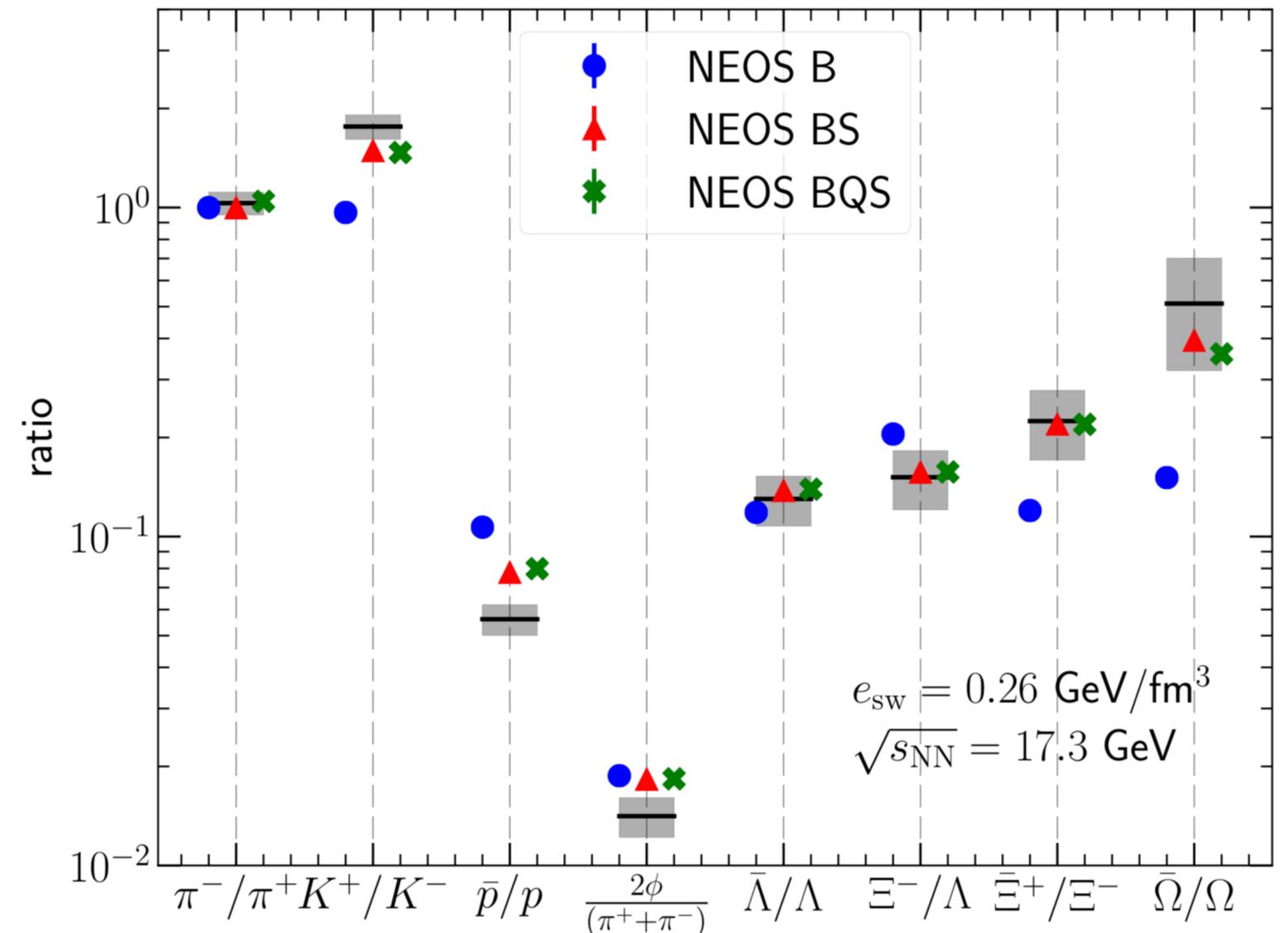


Isentropic trajectories stress importance of BQS modeling for heavy-ion phenomenology



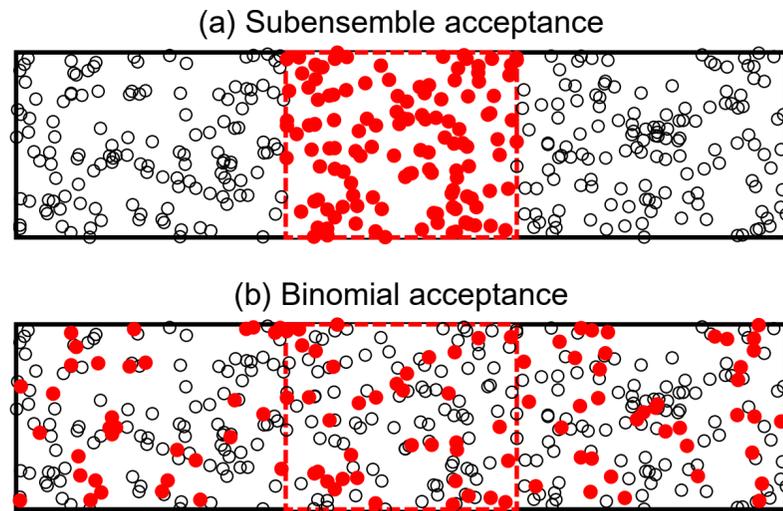
BQS EoS

- ▶ Full dynamical modeling with successive additions of charge conservation conditions increases agreement of predicted particle ratios with experimental data
 - ▶ Including only baryon number is not sufficient
 - ▶ Strange baryons especially affected
 - ▶ Small difference after adding Q

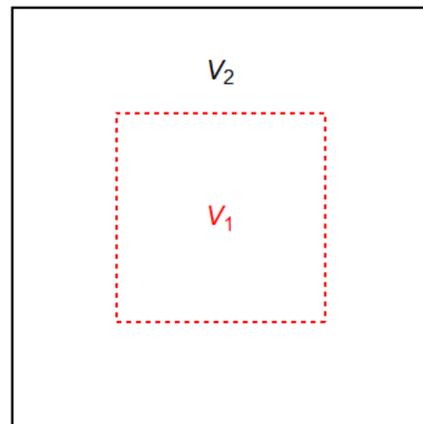


BQS Conserved Fluctuations

- Using the subensemble acceptance method, the conservation effects can be studied for all QCD charges



$$\alpha \equiv V_1/V \equiv \Delta Y_{\text{acc}}/\Delta Y_{4\pi} \lesssim 0.2$$

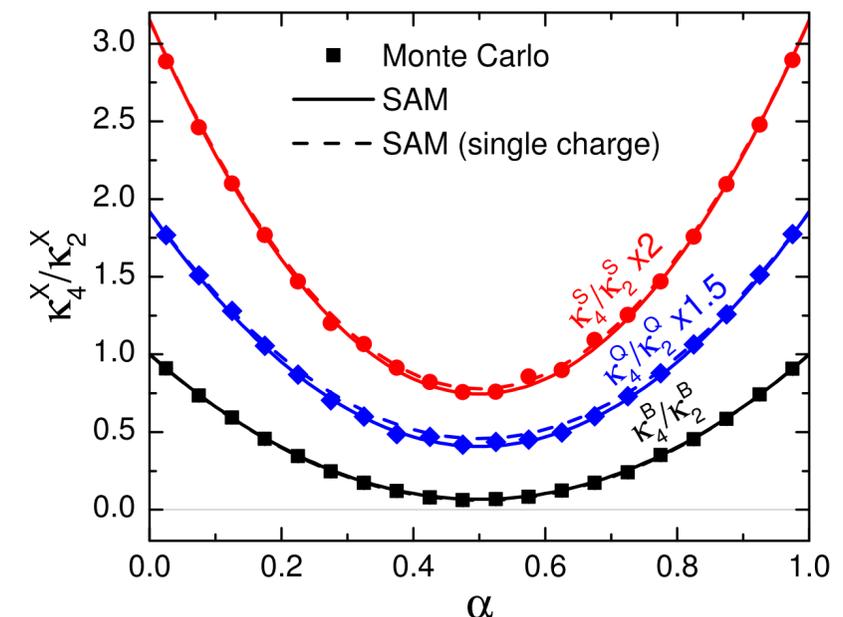
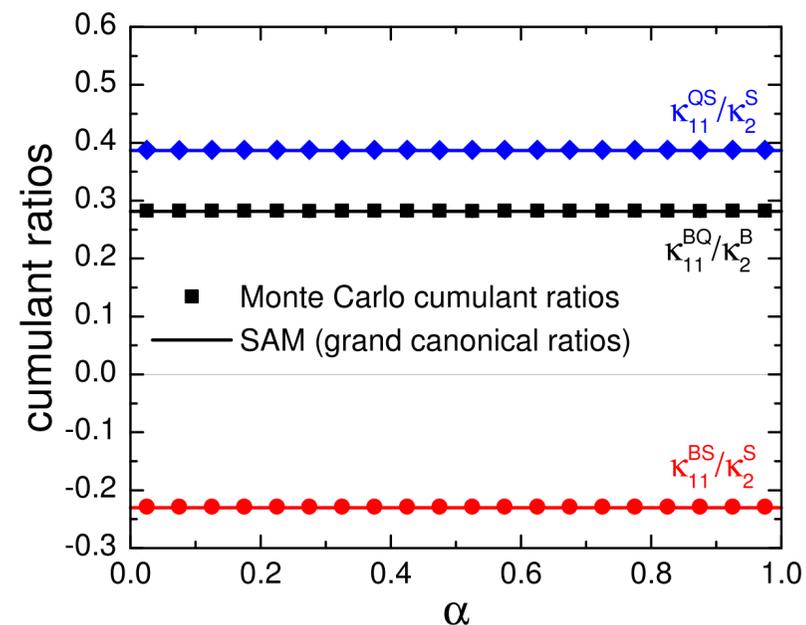


$$\beta \equiv 1 - \alpha$$

$$\text{B conserved: } \kappa_4[B^1] = \alpha VT^3 \beta \left[(1 - 3\alpha\beta) \chi_4^B - 3\alpha\beta \frac{(\chi_3^B)^2}{\chi_2^B} \right]$$

$$\text{BQS conserved: } \kappa_4[B^1] = \alpha VT^3 \beta \left[(1 - 3\alpha\beta) \chi_4^B - \frac{3\alpha\beta}{D[\hat{\chi}_2]} \times \left\{ (\chi_3^B)^2 [\chi_2^Q \chi_2^S - (\chi_{11}^{QS})^2] + (\chi_{21}^{BQ})^2 [\chi_2^B \chi_2^S - (\chi_{11}^{BS})^2] + (\chi_{21}^{BS})^2 [\chi_2^B \chi_2^Q - (\chi_{11}^{BQ})^2] - 2\chi_3^B \chi_{21}^{BQ} (\chi_2^S \chi_{11}^{BQ} - \chi_{11}^{BS} \chi_{11}^{QS}) - 2\chi_3^B \chi_{21}^{BS} (\chi_2^Q \chi_{11}^{BS} - \chi_{11}^{BQ} \chi_{11}^{QS}) \right\} \right]$$

$$D[\hat{\chi}_2] = \chi_2^B \chi_2^Q \chi_2^S + 2\chi_{11}^{BQ} \chi_{11}^{BS} \chi_{11}^{QS} - \chi_2^B (\chi_{11}^{QS})^2 - \chi_2^Q (\chi_{11}^{BS})^2 - \chi_2^S (\chi_{11}^{BQ})^2$$



V. Conclusions

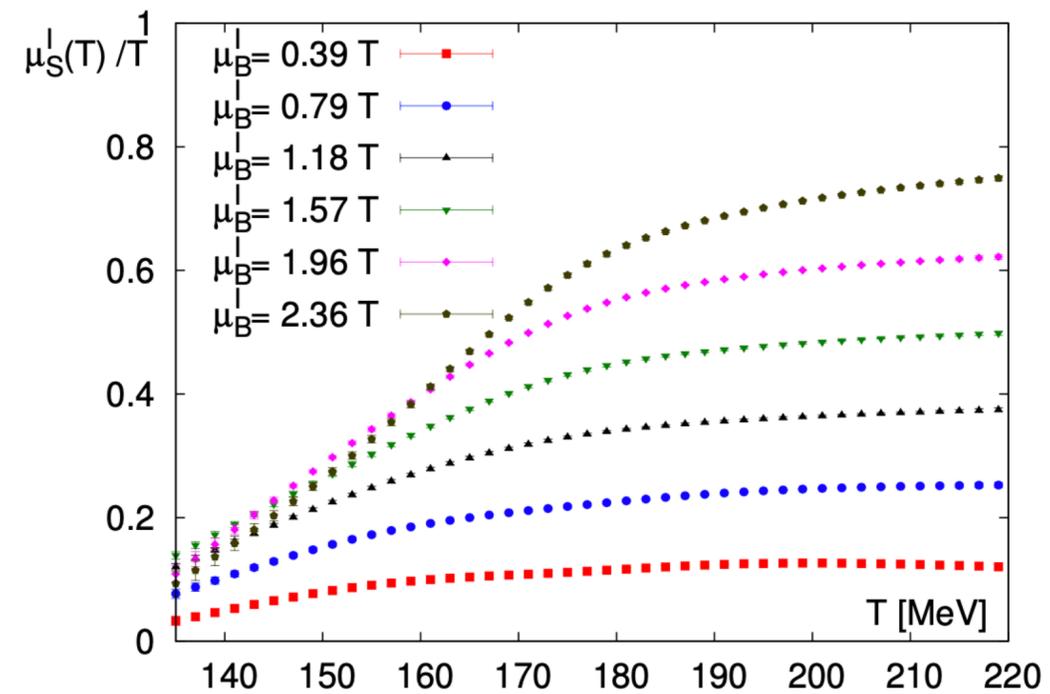
Conclusions

- ▶ Evidence for a light v. strange flavor hierarchy at chemical freeze-out from:
 - ▶ Thermal fits with separate light and strange freeze-out parameters,
 - ▶ Net-K and net- Λ fluctuations yield higher freeze-out temperatures
- ▶ Many models can now incorporate all BQS conserved charges:
 - ▶ ICCING provides initial densities for all charges
 - ▶ BQS EoS provides strangeness neutrality, matching experimental situation
 - ▶ Charge conservation effects for BQS can now be corrected for any EoS
 - ▶ ...more exciting developments out there!

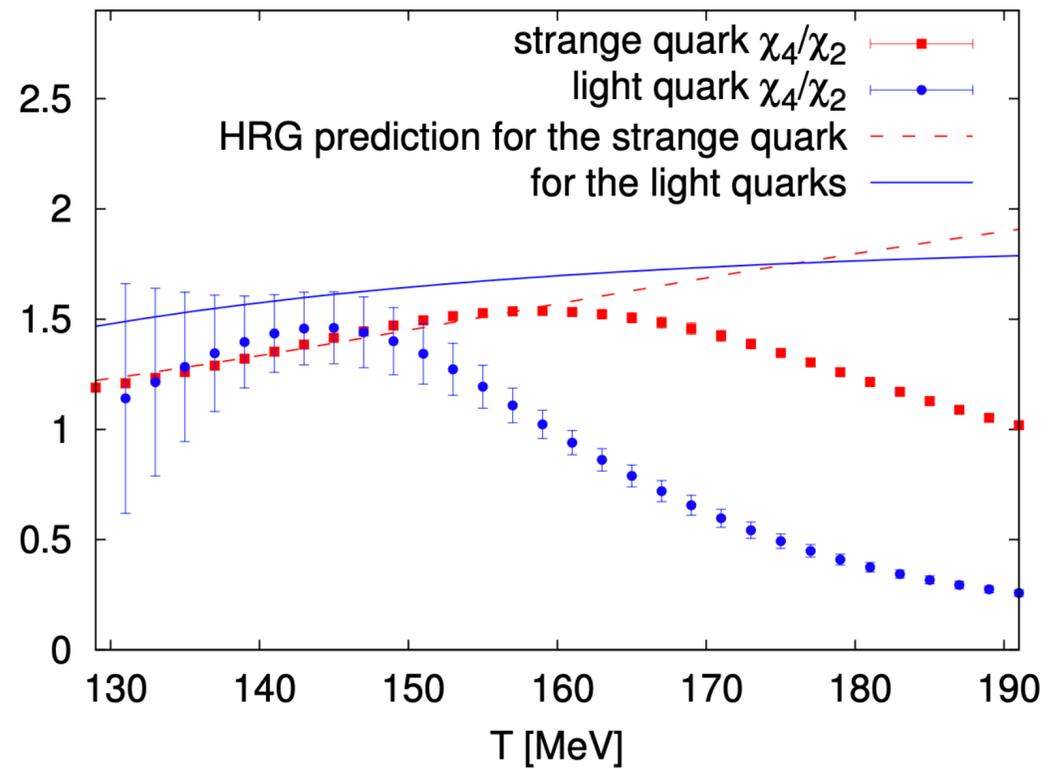
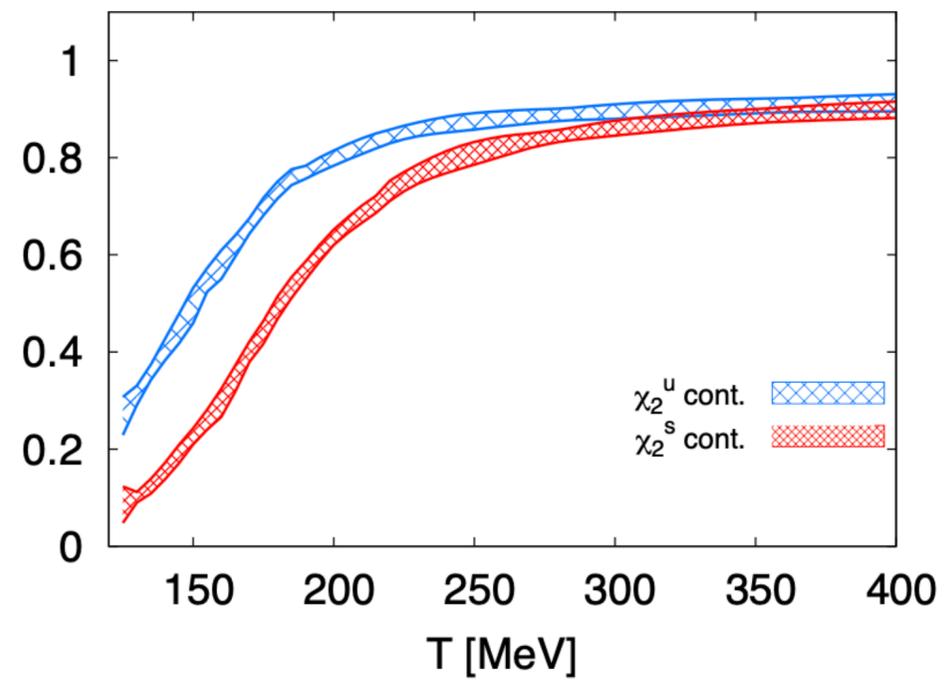


Back-up Slides

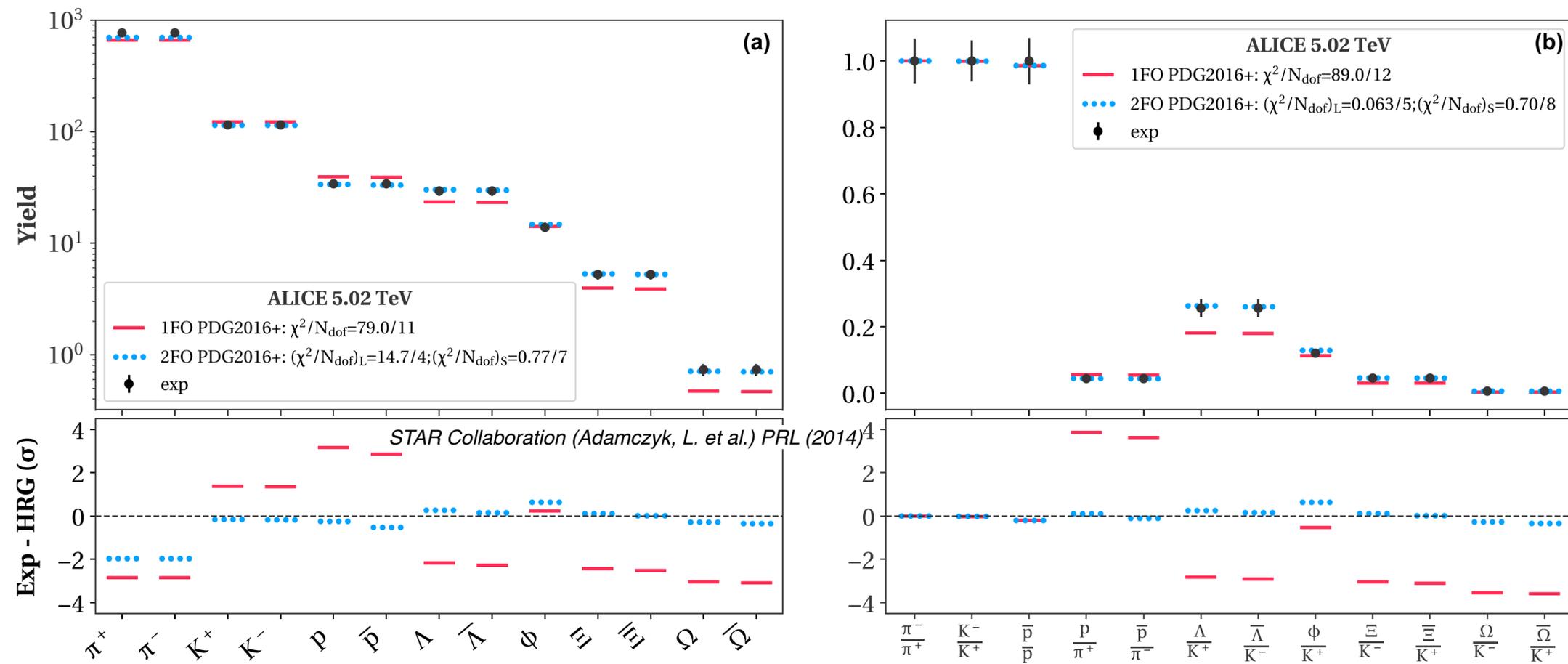
Strangeness neutrality on the lattice



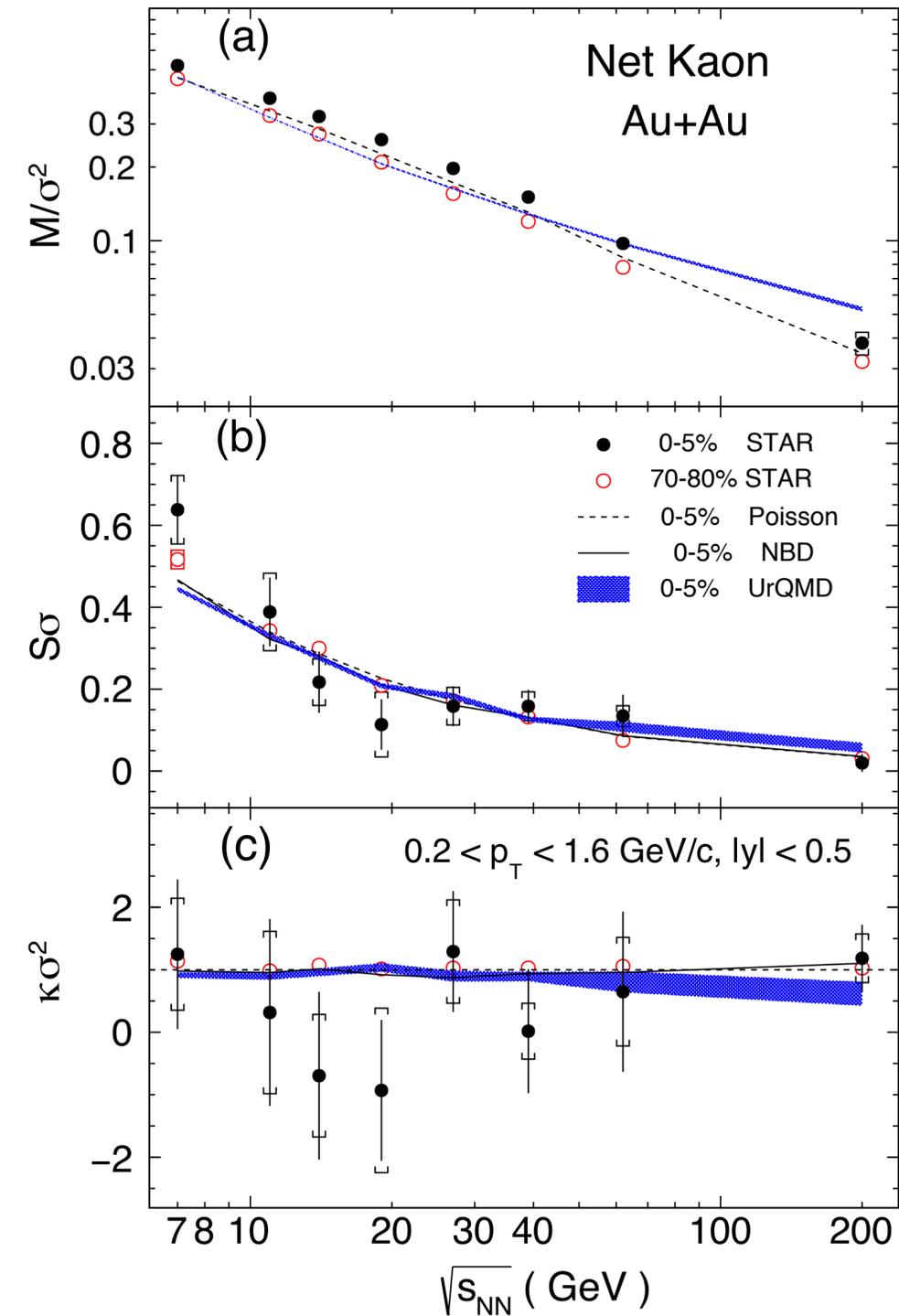
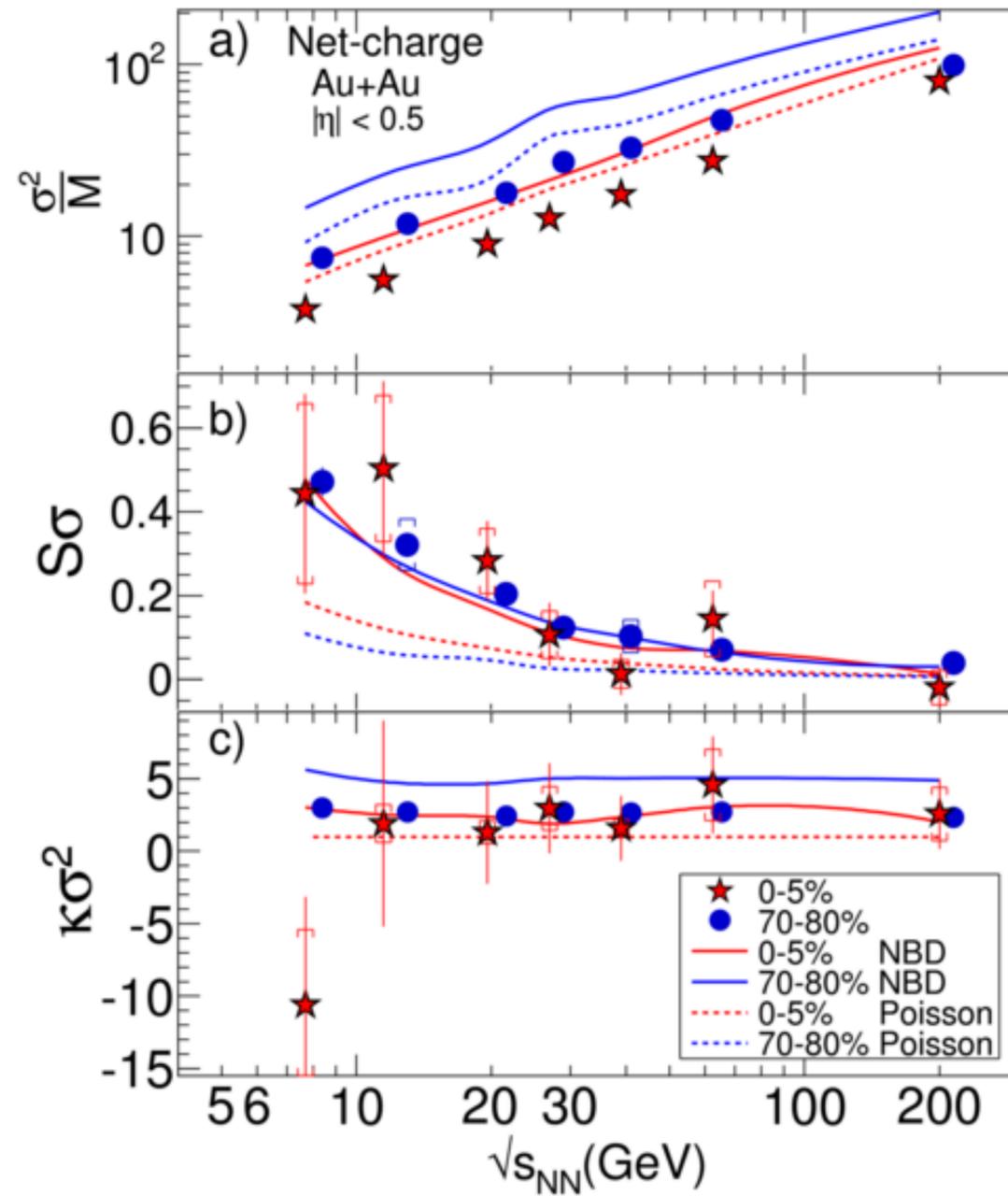
Indications of Flavor Hierarchy



Thermal Fits for ALICE

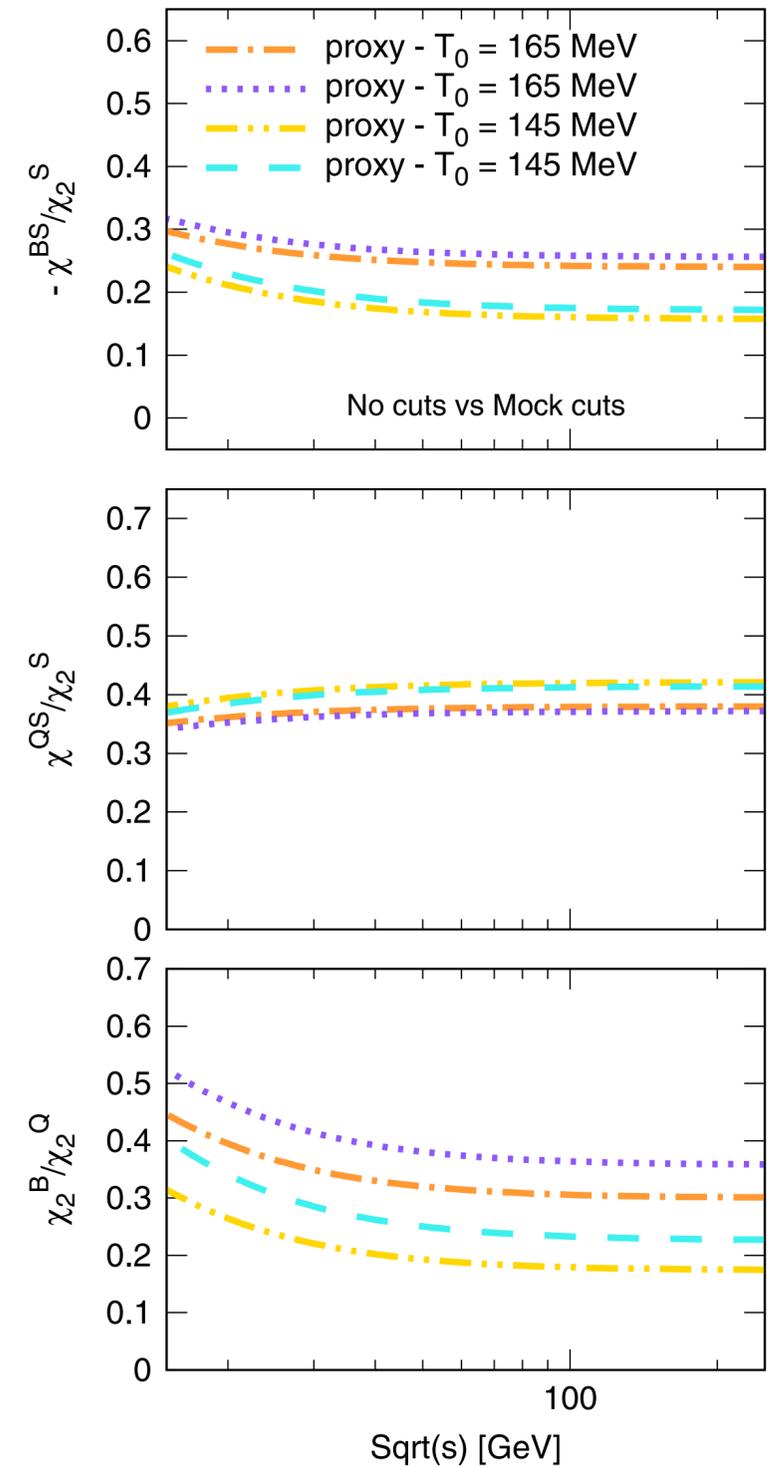
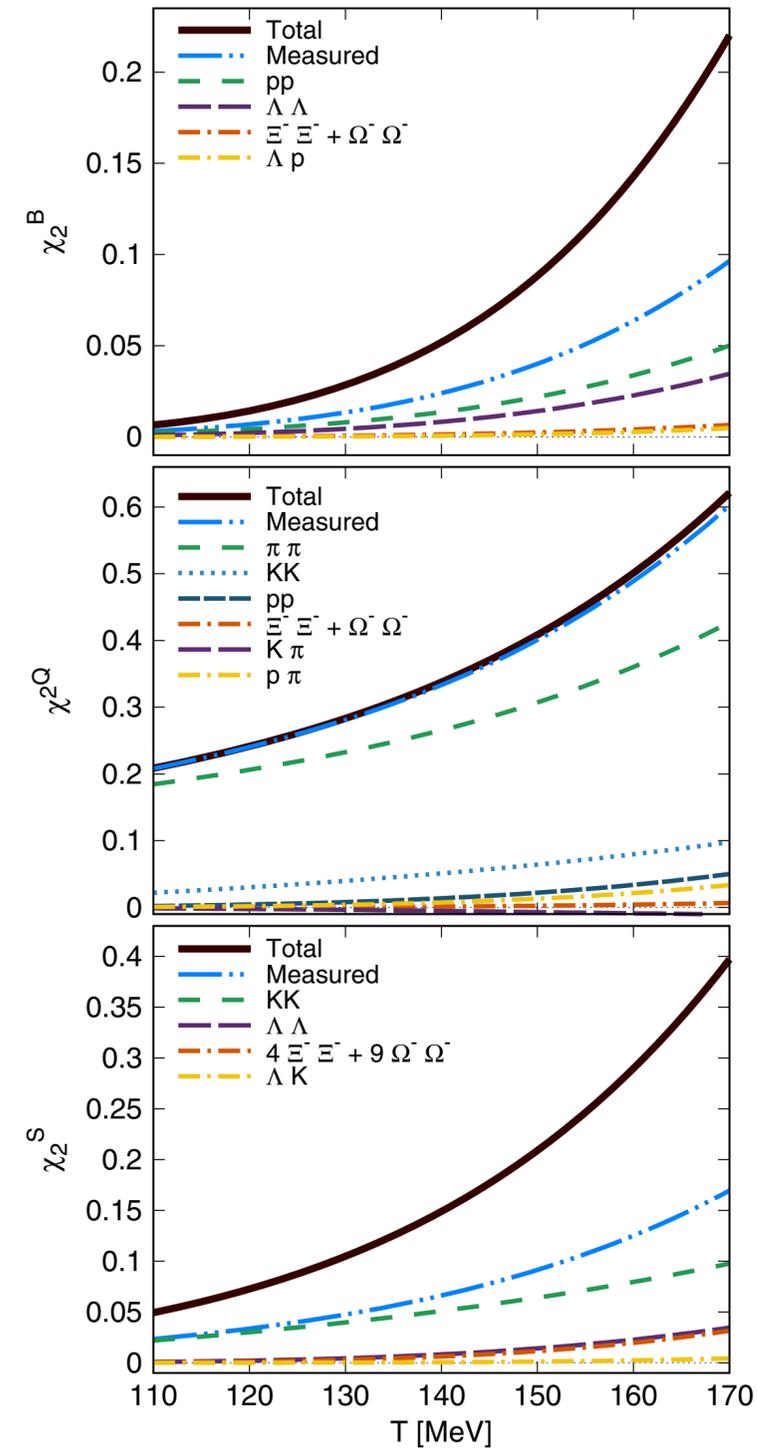


Net-particle Fluctuations from STAR



Cross-correlators

- Diagonal correlators
- Effect of acceptance cuts on the proxies



ICCING Sampling Algorithm

