

# Jet Substructure at EIC

Mriganka Mouli Mondal  
Stony Brook University

## **In Collaboration with :**

Stony Brook University, Stony Brook

Theorists : Yang-Ting Chien, George Sterman

Experimentalists : Abhay Deshpande, Roli Esha, Henry Klest, Meng Jia Ou, Jinlong Zhang

University of California, Riverside

Experimentalists : Miguel Arratia, Sean Preins

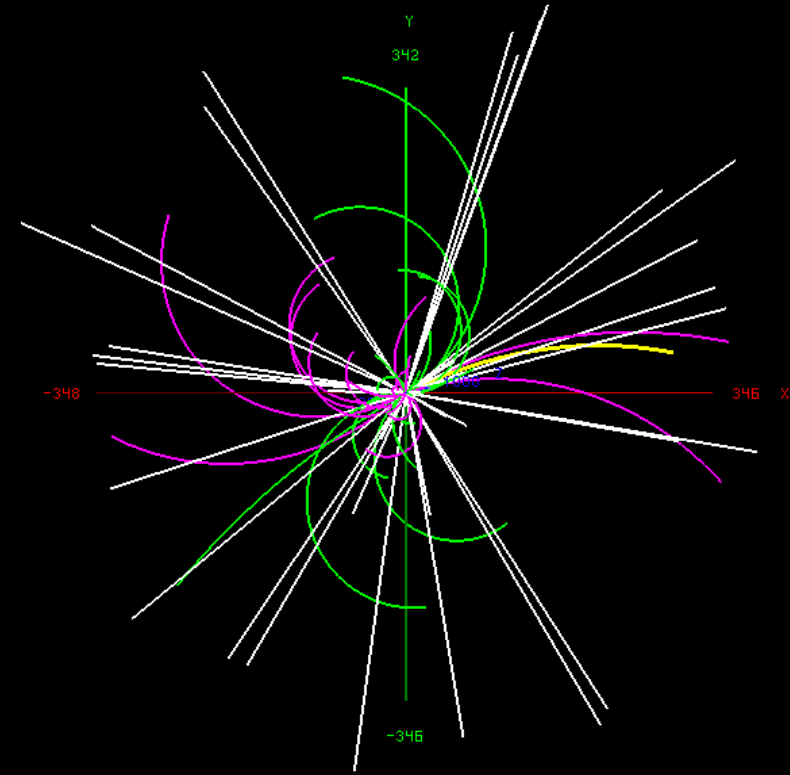
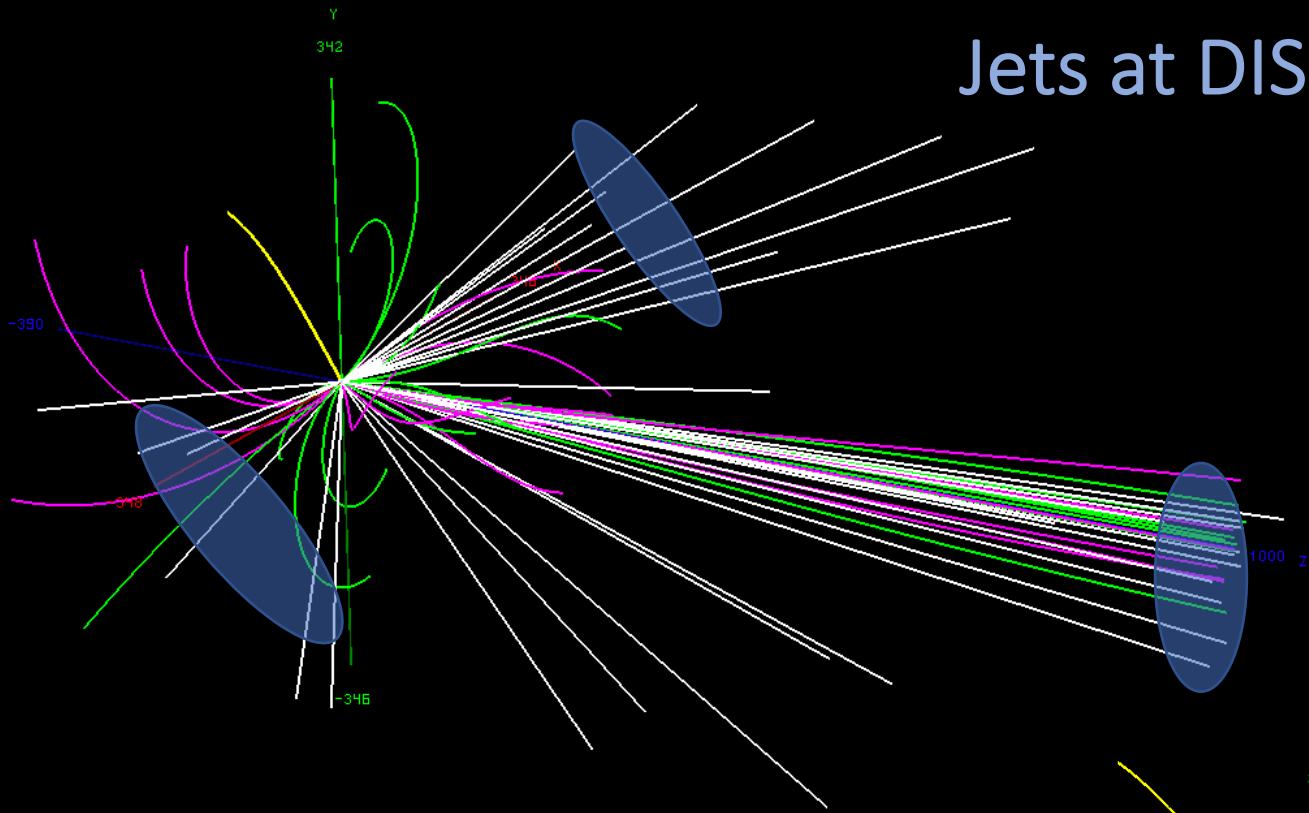
Wayne State University, Detroit

Experimentalist : Raghav Kunnawalkam Elayavalli

# Introduction

- Jets and their structure
  - Correlations in momentum, charge and flavor : leading and next to leading particles in a jet
  - access to the dynamics of fragmentation and color entanglement in QCD
- Observable
  - Charge asymmetry and connection to dihedron fragmentation function
- Pythia event studies on the observable and PID limits
  - Acceptance of Jets with with limits in PID
  - Charge asymmetries for two different PID limits in central region
- Summary

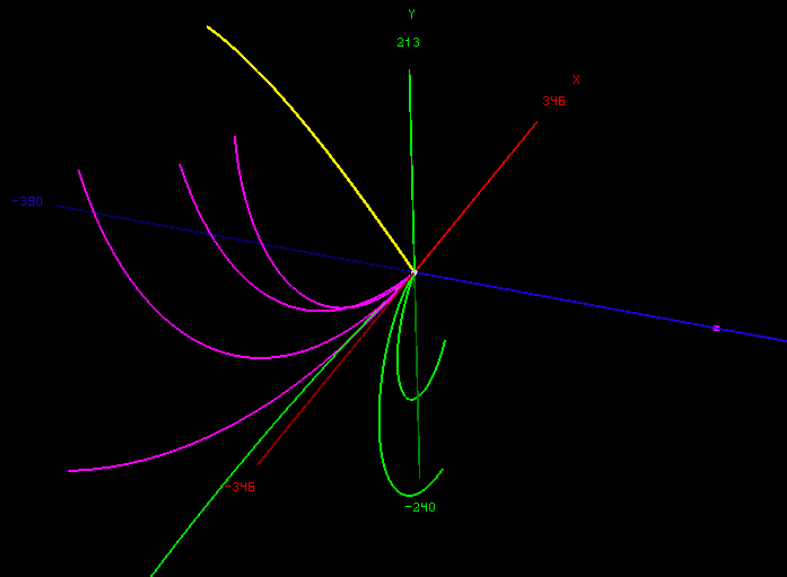
# Jets at DIS



Leading momentum particles in the jet –  $K, \bar{p}$

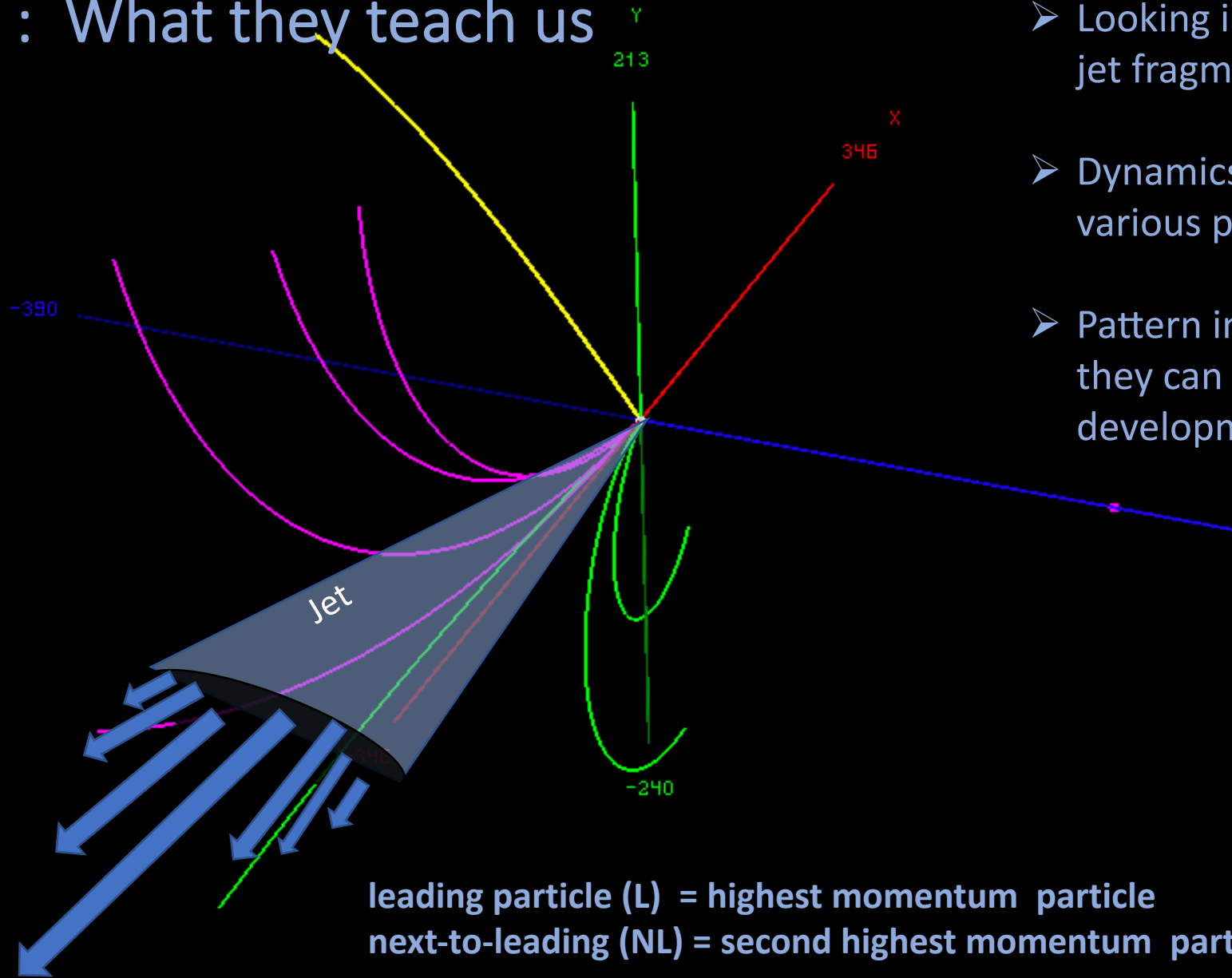
Particles in Jet :

| Px    | Py    | Pz    | PID   |
|-------|-------|-------|-------|
| -7.64 | -4.41 | -4.21 | 321   |
| -2.06 | -1.89 | -1.40 | -2212 |
| -1.44 | -0.87 | -0.69 | -211  |
| -1.07 | -0.08 | 0.04  | 2212  |
| -0.58 | -0.45 | -0.53 | -211  |
| -0.70 | -0.39 | -0.39 | -211  |
| -0.65 | -0.52 | -0.04 | 211   |



Scattered-electron  
 +ve particles  
 -ve particles  
 Neutral particles

# Jets : What they teach us



- Looking in the nonperturbative aspect of jet fragmentation
- Dynamics that led to fragmentation of various particle species in certain ways
- Pattern in charge and flavor separation : they can inform in future theoretical development

leading particle (L) = highest momentum particle  
next-to-leading (NL) = second highest momentum particle

# Momentum-charge correlations

□ Leading particle (L) and next-to-leading (NL) are both pions generated in two pictures

i) “**random**” picture : L is fixed and NL is random and both L and NL pions are charged

$$N_{C\bar{C}}^{\text{random}} = N_{CC}^{\text{random}} = \frac{N^{\text{random}}}{2}$$

$C\bar{C}$  indicates opposite charges  
 $CC$  same charge

ii) “**alternating**” picture : perturbative shower gives  $q_L$  followed by  $\bar{q}'_{NL}$ , which form pions by sharing a soft pair:

$$q_L + \bar{q}_{NL} \rightarrow q_L + (\bar{q}_s + q_s) + \bar{q}'_{NL} \rightarrow \pi(q_L, \bar{q}_s) + \pi(q_s, \bar{q}'_{NL})$$

$$N_{C\bar{C}}^{\text{alternating}} = N_{CC}^{\text{alternating}} \quad \text{and} \quad N_{CC}^{\text{alternating}} = 0$$

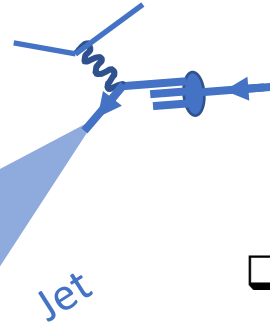
□ **The observable :**

$$r_{\text{asy}} \equiv \frac{N_{CC} - N_{C\bar{C}}}{N_{CC} + N_{C\bar{C}}} = \frac{1 - a}{2} - \left( \frac{1 - a}{2} + a \right) = -a$$

- ✓ provided every event results from one of these two processes, with no interference
- ✓ percentage of “alternating” =  $a$ ; and percentage of random events =  $1 - a$

➤  $r_{\text{asy}}$  is a measurement of the fraction of hadronizations that are “string-like”, alternating between quark and antiquark (classical picture)

18x275 GeV



$$r_{\text{asy}} \equiv \frac{N_{CC} - N_{C\bar{C}}}{N_{CC} + N_{C\bar{C}}}$$

next-lead  
leading

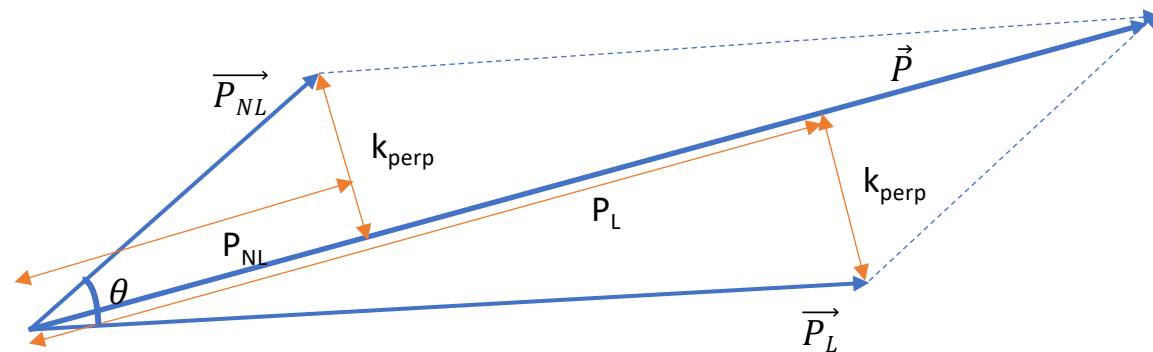
❑ Construct  $r_{\text{asy}}$  with particle compositions with various parameters

- ✓ Leading particle ( $\pi$ ) and next leading particles ( $\pi/K/p$ )
- ✓ Leading particle ( $K$ ) and next leading particles ( $K/\pi/p$ )
- ✓ Leading particle ( $p$ ) and next leading particles ( $p/\pi/K$ )

❑ Looking via

- ✓ Momentum-next lead particle/momentum of leading particle
- ✓ Faction of jet momentum carried by leading particle
- ✓ Angle between the leading and next to leading particles ( $\Delta\theta$ )
- ✓ relative transverse momentum ( $k_{\text{perp}}$ )
- ✓ pair invariant mass
- ✓ Formation time :  $[2z(1-z) P] / k_{\text{perp}}^2$

PYTHIA-6 :  
 $Q_2 > 65 \text{ GeV}$   
 Jet Reconstruction :  
 anti-kt  $R = 0.7$   
 Jet  $p_t > 8 \text{ GeV}$   
 pt-tracks  $> 0.2 \text{ GeV}$   
 track  $|\eta| < 3.5$   
 Jet  $|\eta| < 2.8$



$$z = P_{NL} / (P_{NL} + P_L)$$

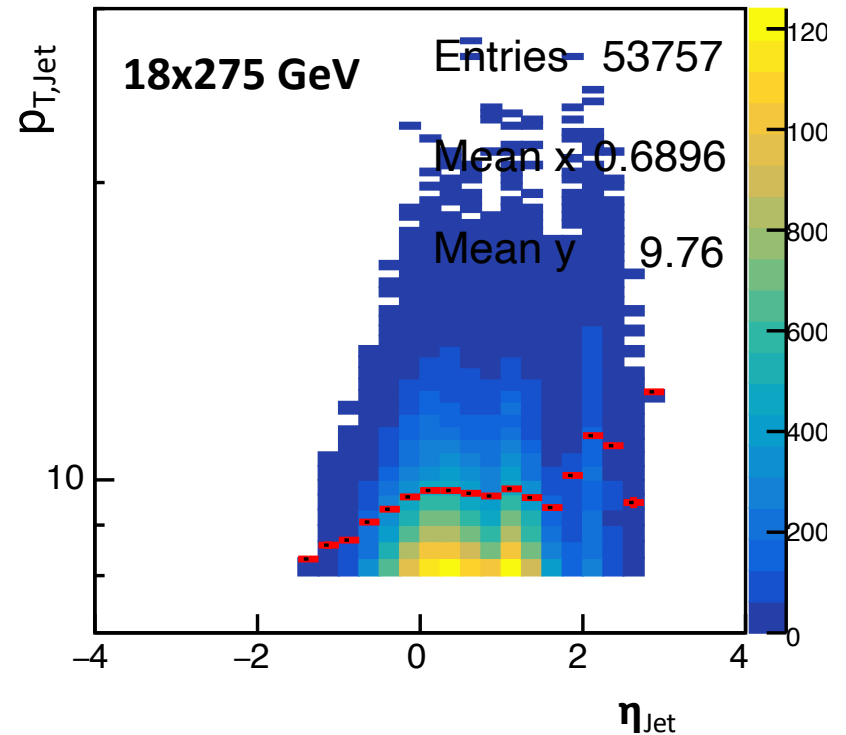
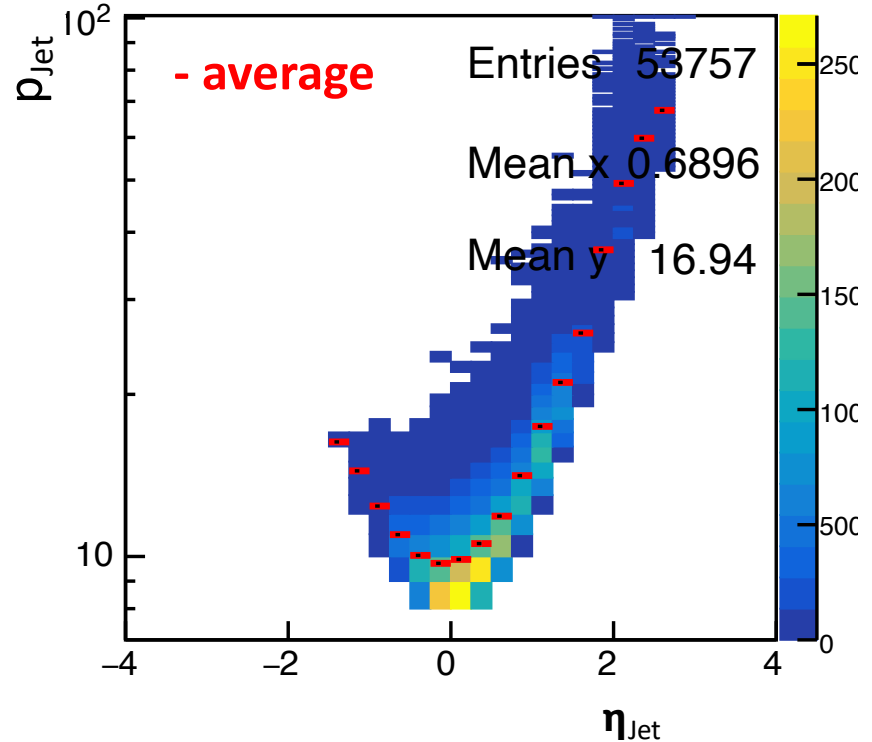
$$P_L = (1-z)P$$

$$P_{NL} = zP$$

# Acceptance

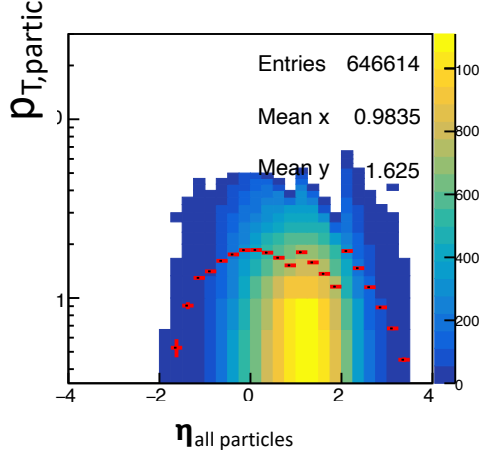
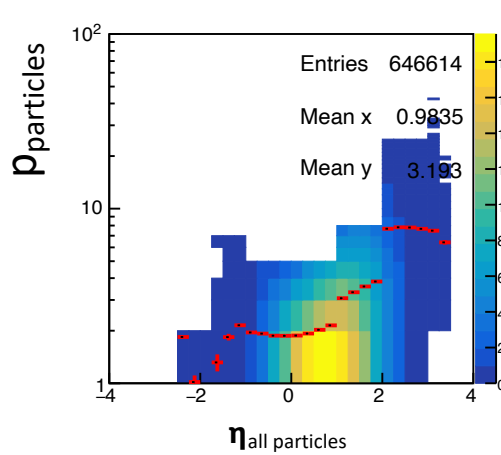
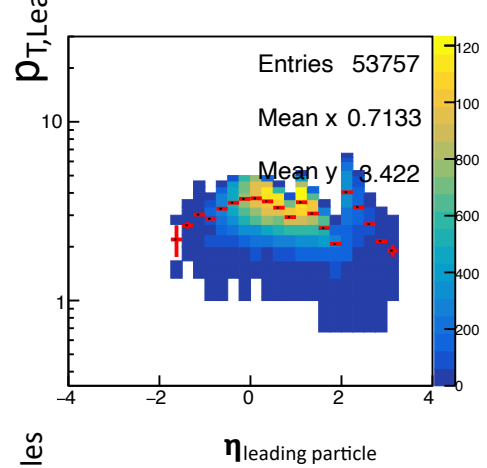
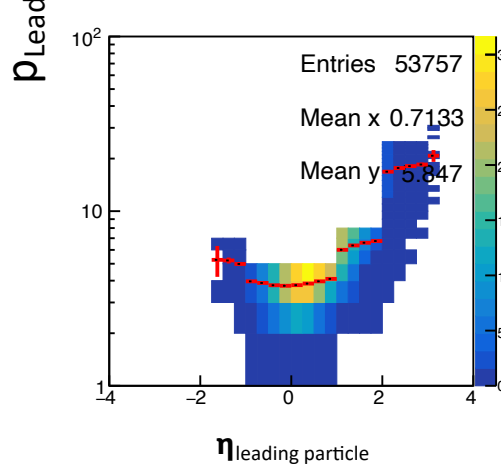
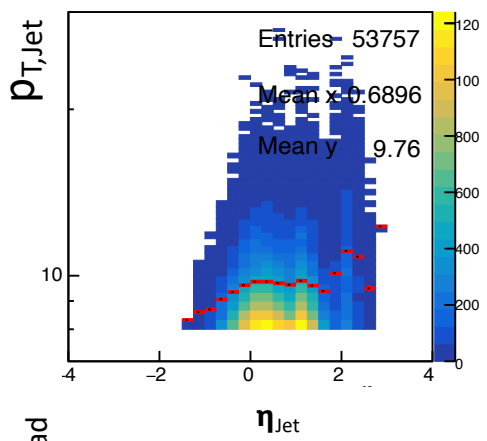
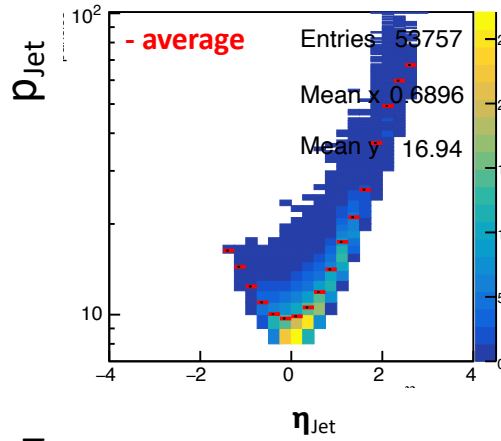
Momentum cuts on tracks for approximate PID requirement at EIC at different  $\eta$  regions (currently we made one hard limit at certain momentum)

| $\eta$ - range | Momentum cut ( $\text{GeV}/c^2$ ) |
|----------------|-----------------------------------|
| -3.5 to -1.0   | 7                                 |
| -1.0 to 1.0    | 5                                 |
| 1.0 to 2.0     | 8                                 |
| 2.0 to 2.0     | 25                                |
| 3.0 to 3.5     | 45                                |



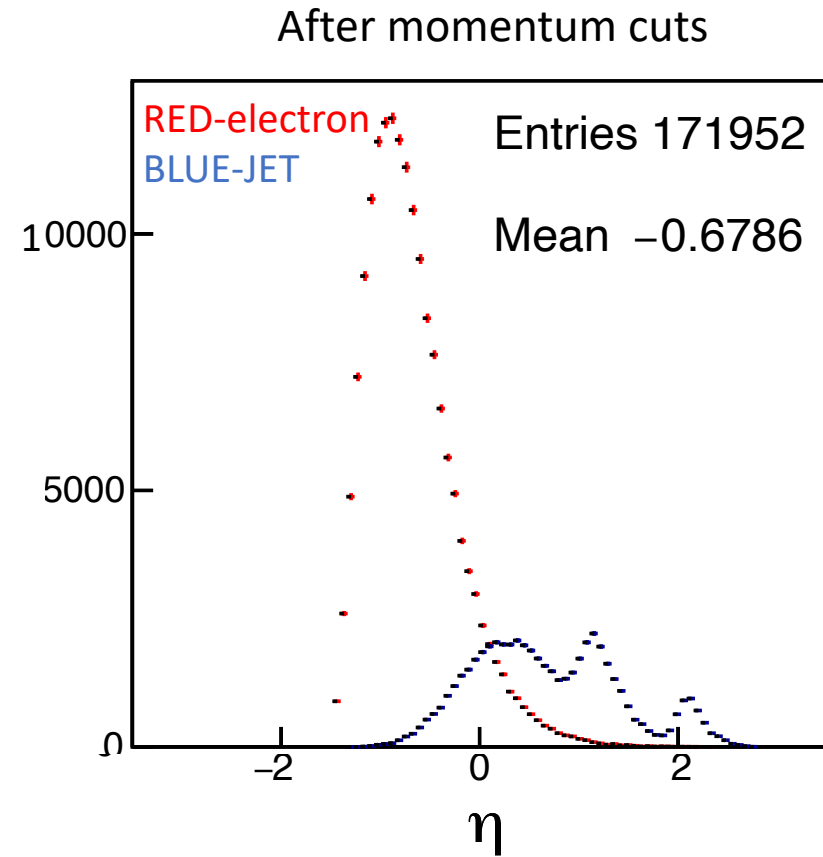
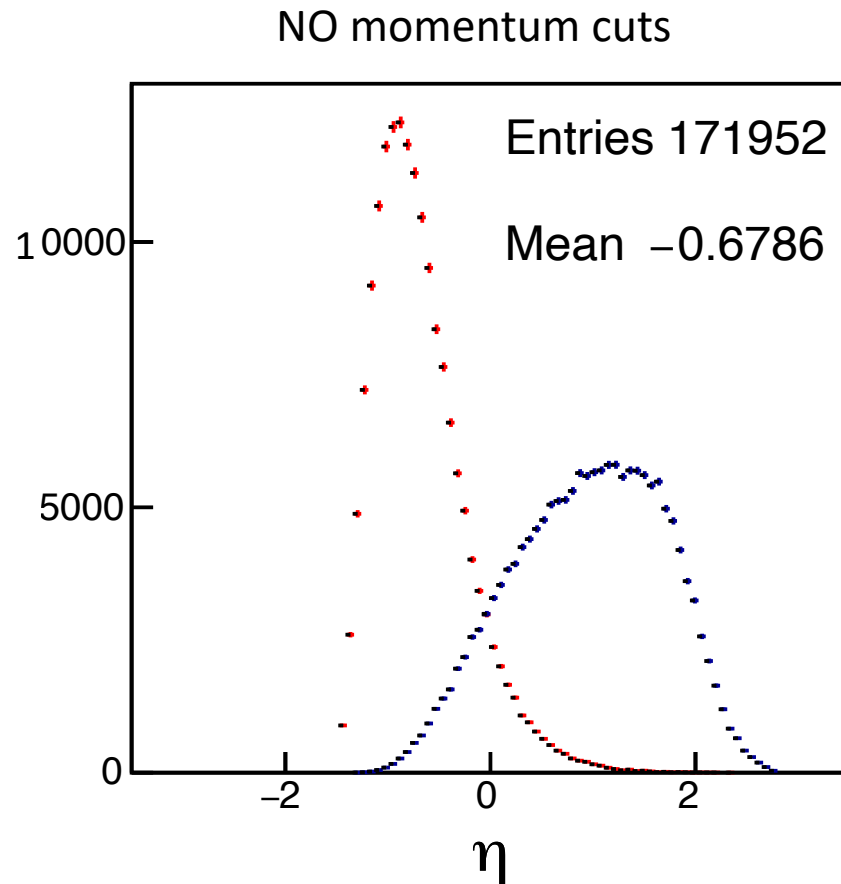
# Acceptance

| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -3.5 to -1.0   | 7                                  |
| -1.0 to 1.0    | 5                                  |
| 1.0 to 2.0     | 8                                  |
| 2.0 to 2.0     | 25                                 |
| 3.0 to 3.5     | 45                                 |



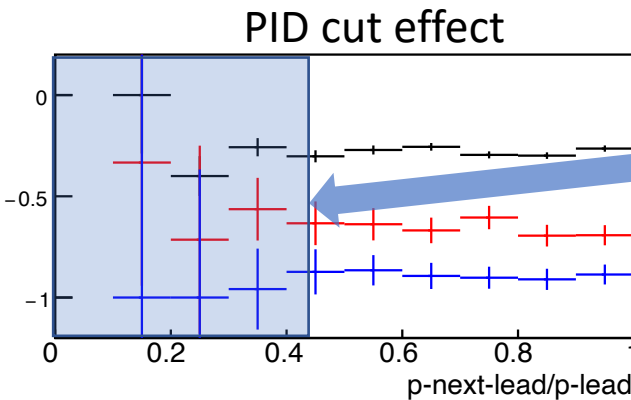
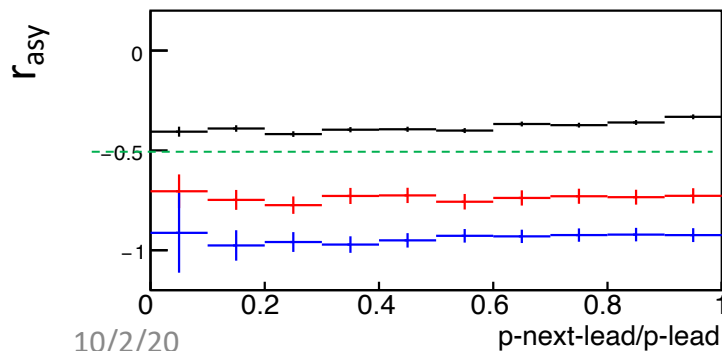
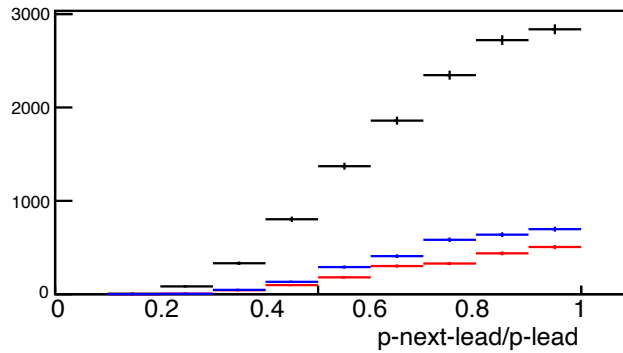
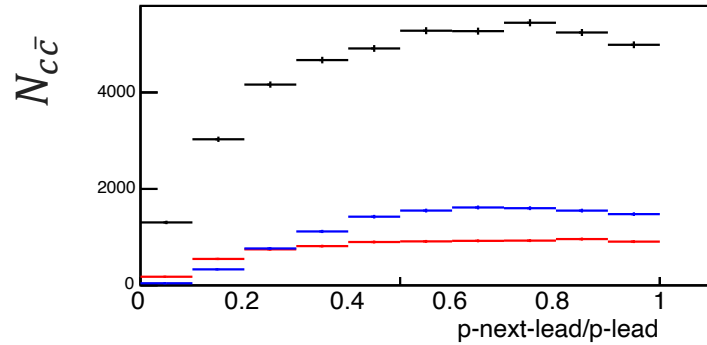
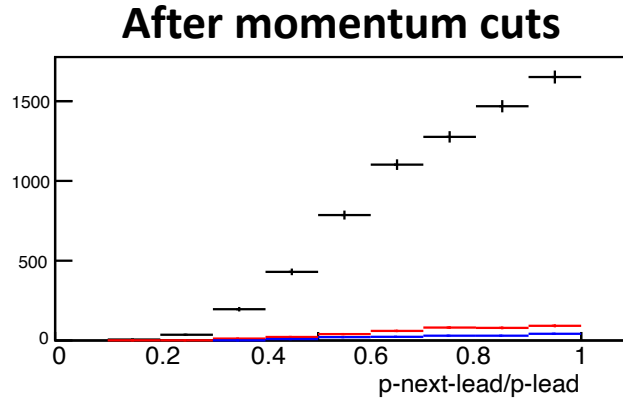
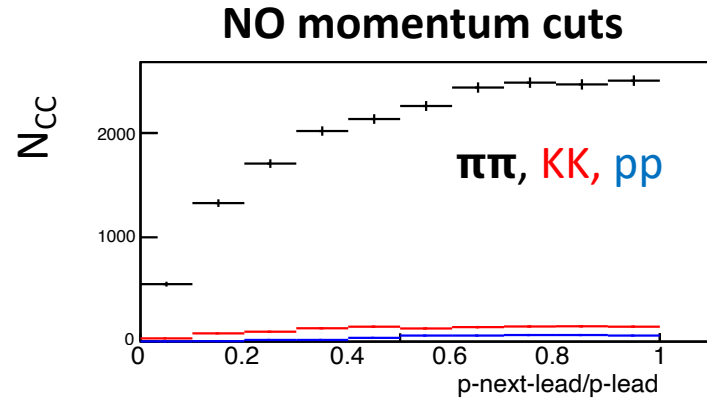


# The effect of momentum cuts in pseudorapidity acceptance of jets



PID limit has huge impact on jet acceptance when leading particle's PID is performed

# Correlation with two leading particles



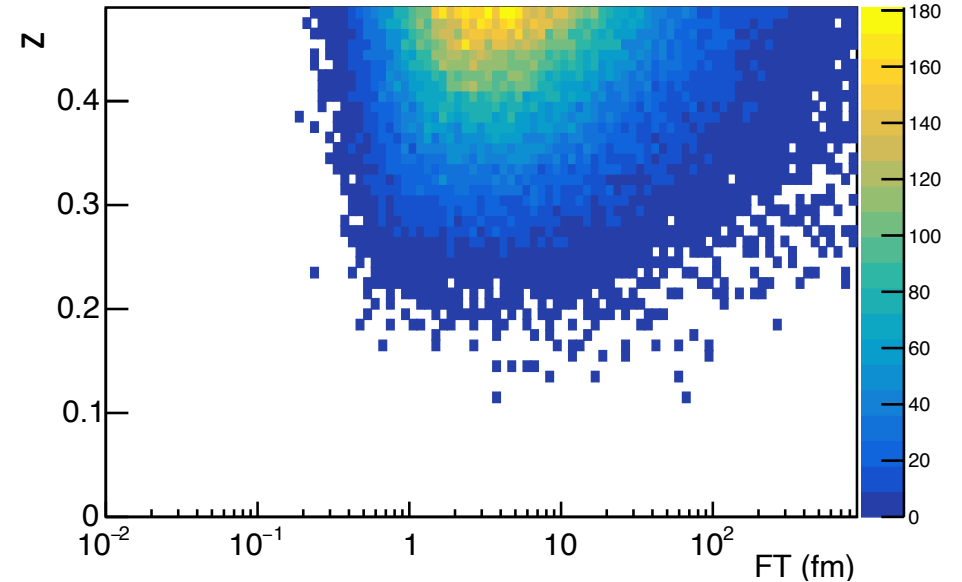
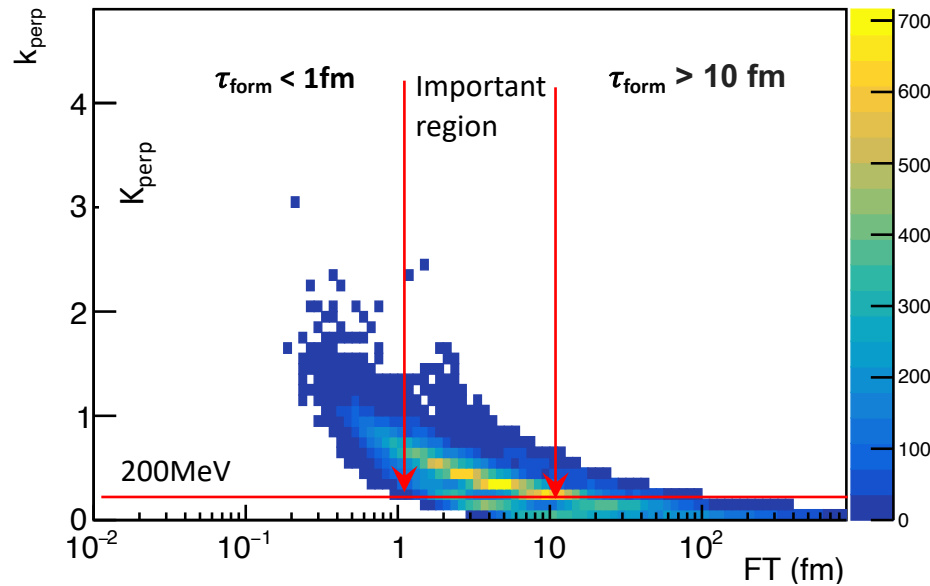
$$r_{asy} \equiv \frac{N_{CC} - N_{C\bar{C}}}{N_{CC} + N_{C\bar{C}}}$$

- The strength of correlations are different for pions, kaons and proton

The transition region from soft proton emission to hard is very important

# Formation time

$$\text{Formation time, } \tau_{\text{form}} = [2z(1-z) P] / k_{\text{perp}}^2$$



$\tau_{\text{form}} < 1\text{ fm}$  : L and NL particles seem to separate after a very short time, which might decorrelate their hadronization

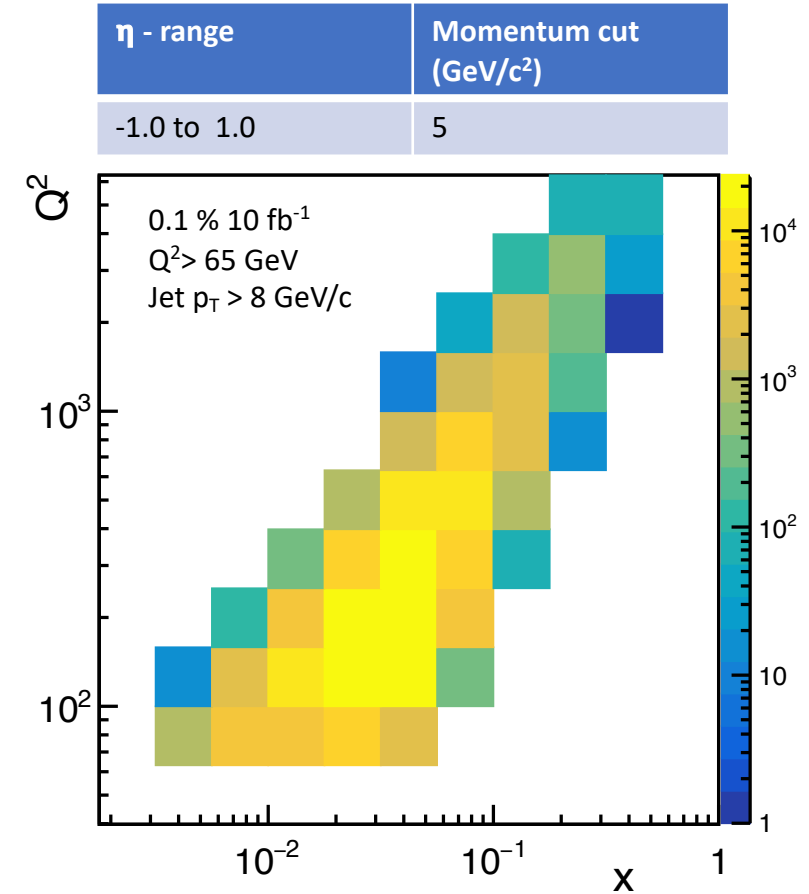
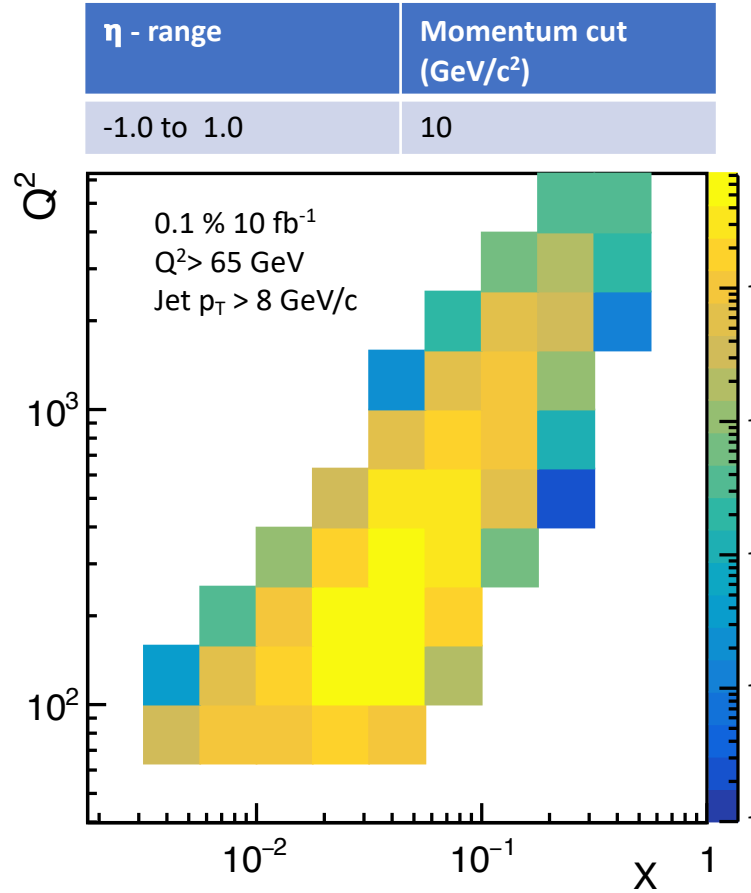
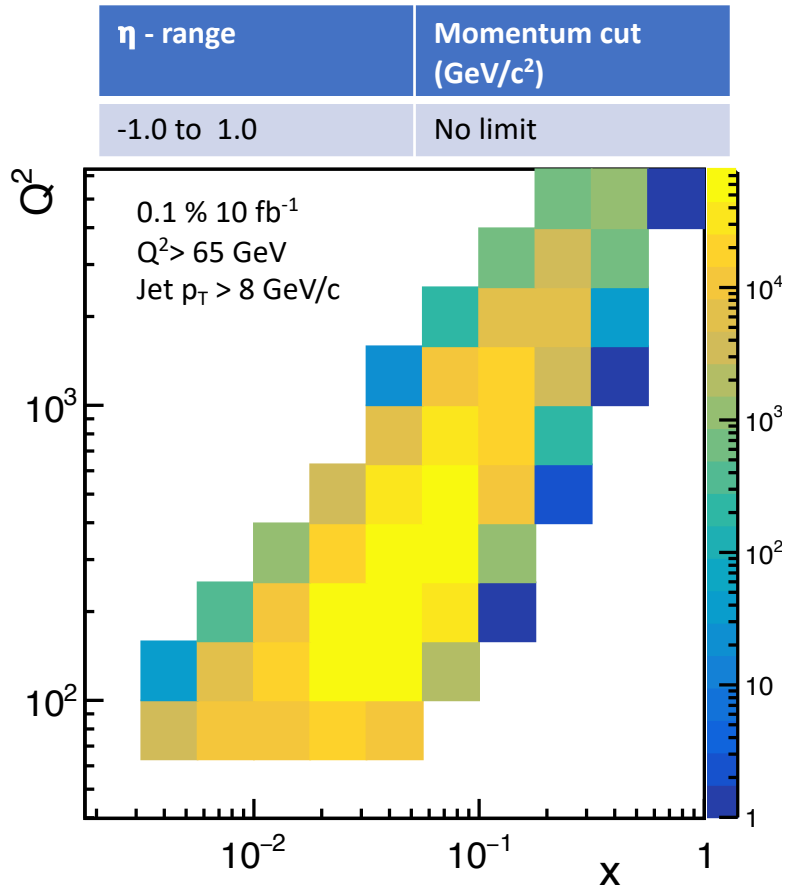
$\tau_{\text{form}} > 10\text{ fm}$  ( $k_{\text{perp}} < 200\text{ MeV}$ ) : nonperturbative transverse momenta in the jet, and we don't think that going to longer  $\tau_{\text{form}}$  or smaller  $k_{\text{perp}}$  leads to new dynamics

Important region to study in data  $\tau_{\text{form}} =$  "a few fermi" and "a few dozen fermi",  $k_{\text{perp}} =$  "a few GeV" to "several hundred MeV"

## Different PID limits at Barrel region : a comparative study

| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -1.0 to 1.0    | No limit                           |
| -1.0 to 1.0    | 10                                 |
| -1.0 to 1.0    | 5                                  |

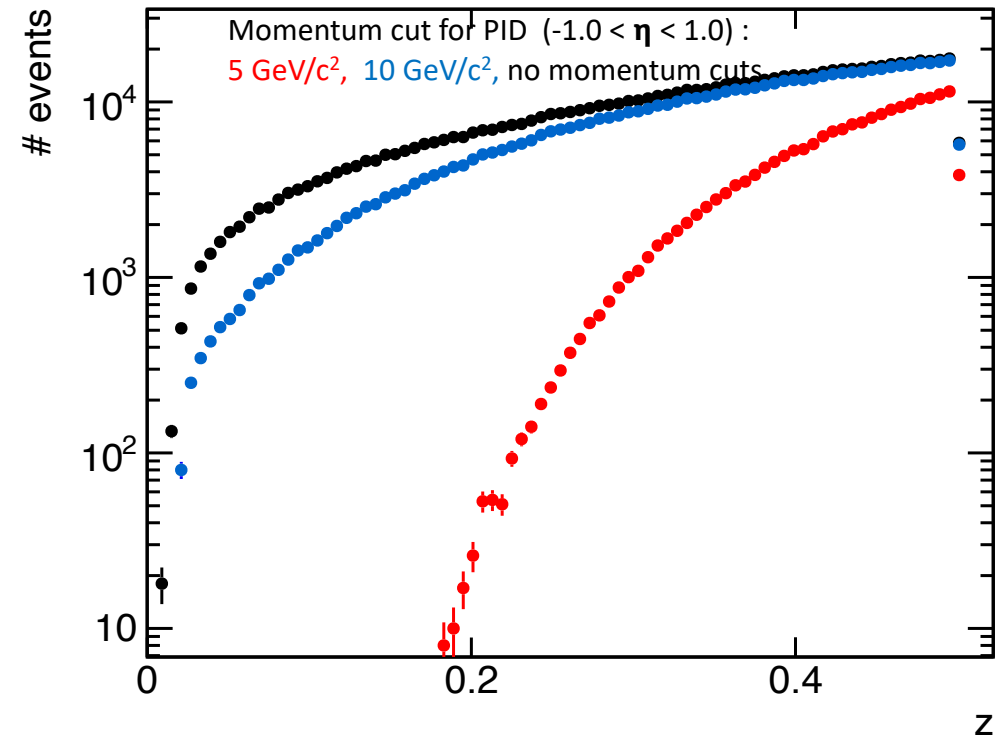
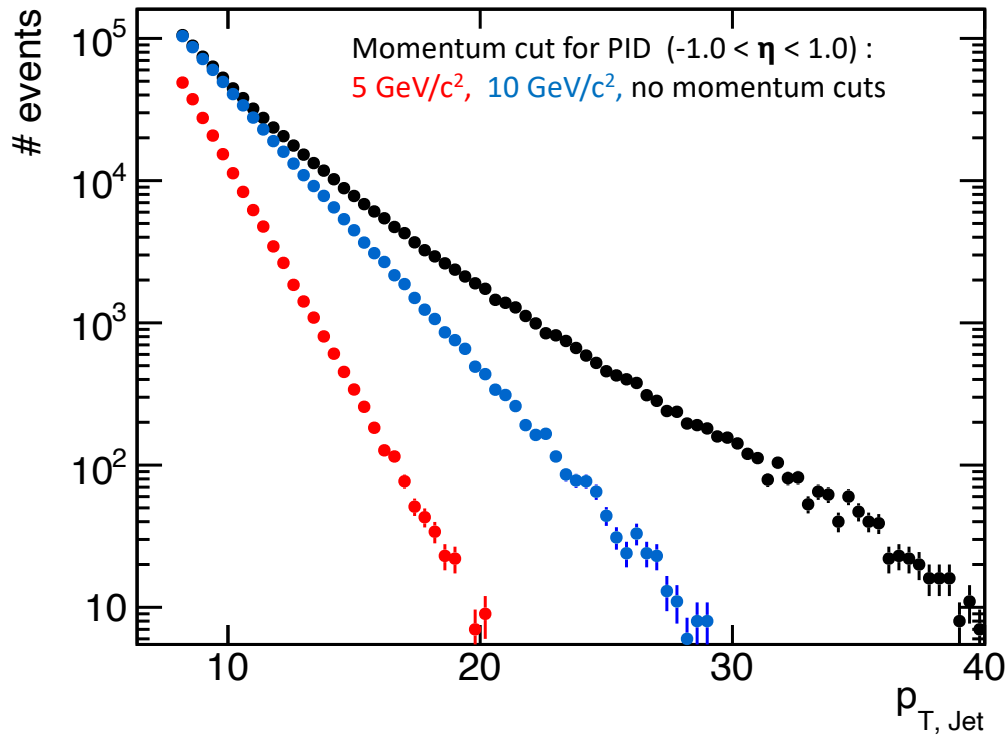
# Acceptance in $x$ - $Q^2$



Covers wider  $x$ - $Q^2$  phase with high momentum PID

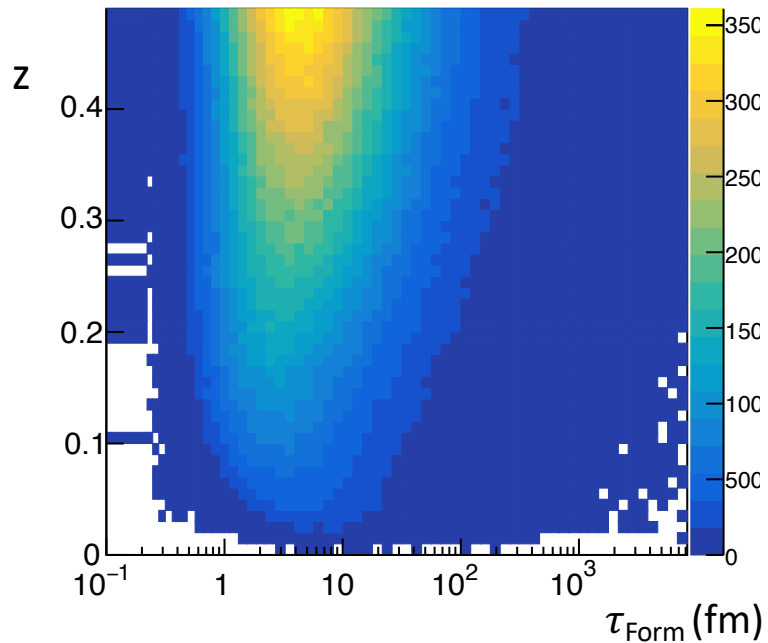
# $p_{T,Jet}$ and $z$ reach for different PID limits at Barrel region

$$z = P_{NL}/(P_{NL}+P_L)$$

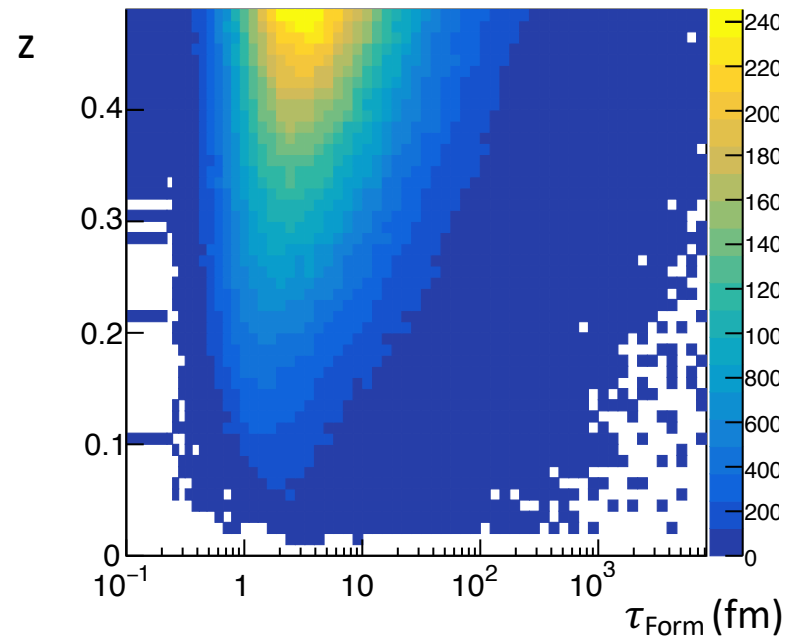


- PID up to 10 GeV/c<sup>2</sup> impose significant limitation in small  $z$  and high  $p_{T,Jet}$  reach
- However, increasing PID performance up to 10 GeV/c<sup>2</sup> would :
  - ✓ significantly cover small  $z$
  - ✓ enhance  $p_{T,Jet}$  limit of measurement

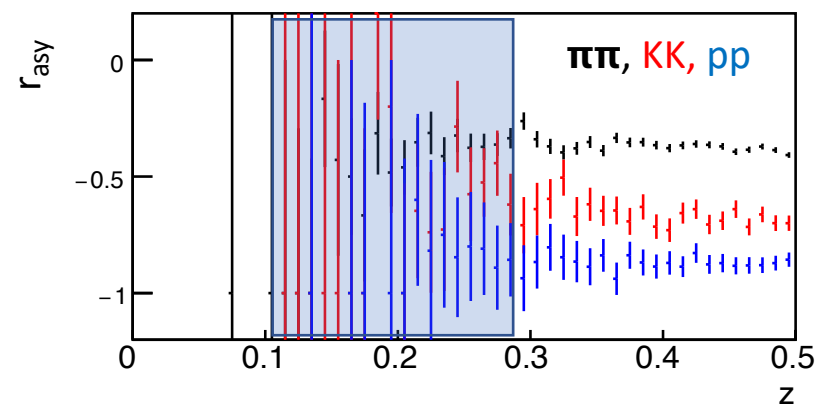
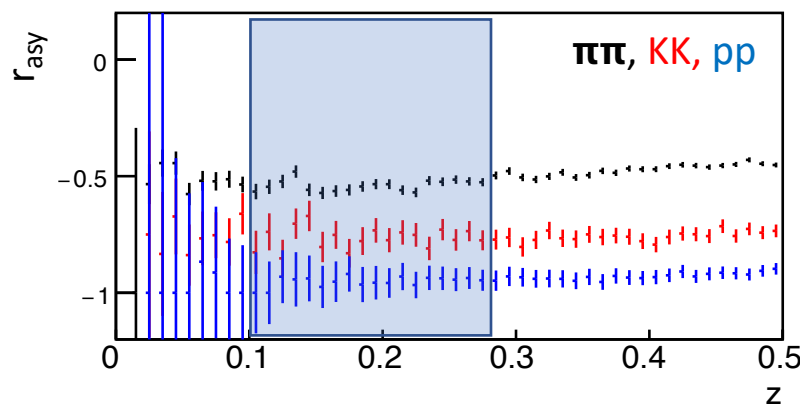
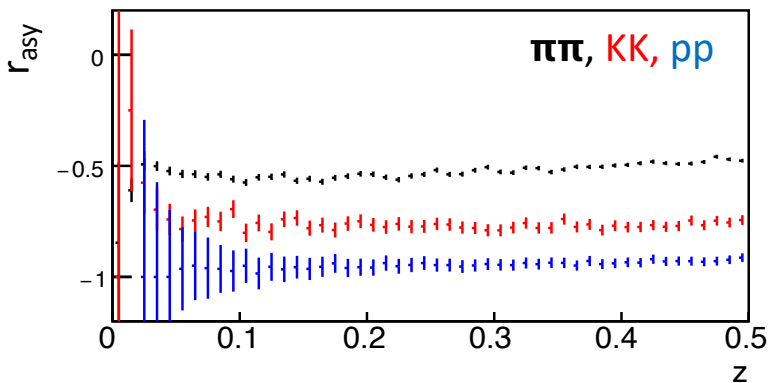
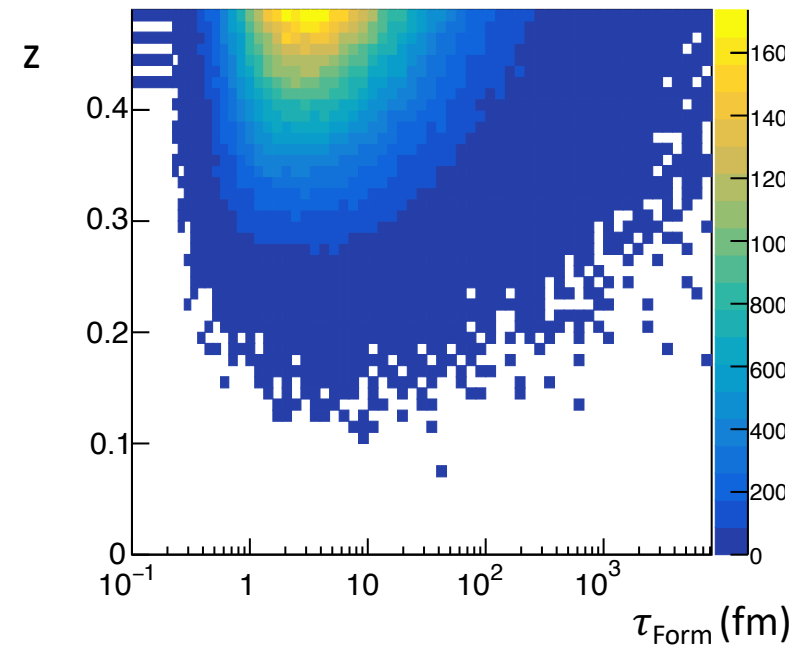
| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -1.0 to 1.0    | No limit                           |



| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -1.0 to 1.0    | 10                                 |



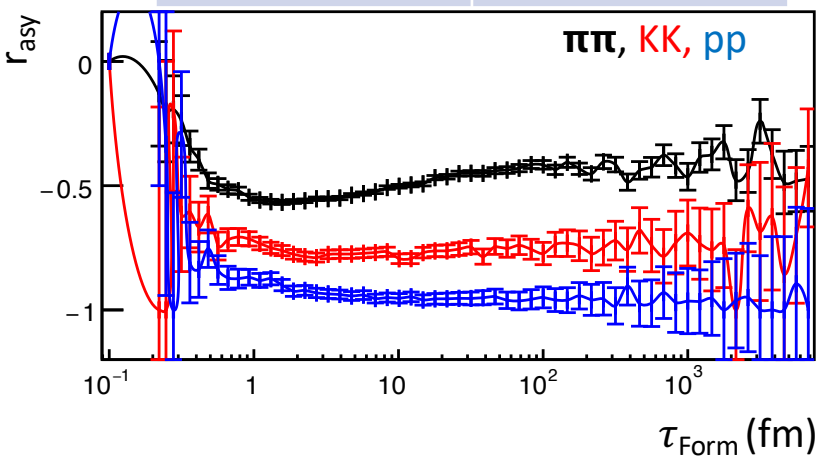
| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -1.0 to 1.0    | 5                                  |



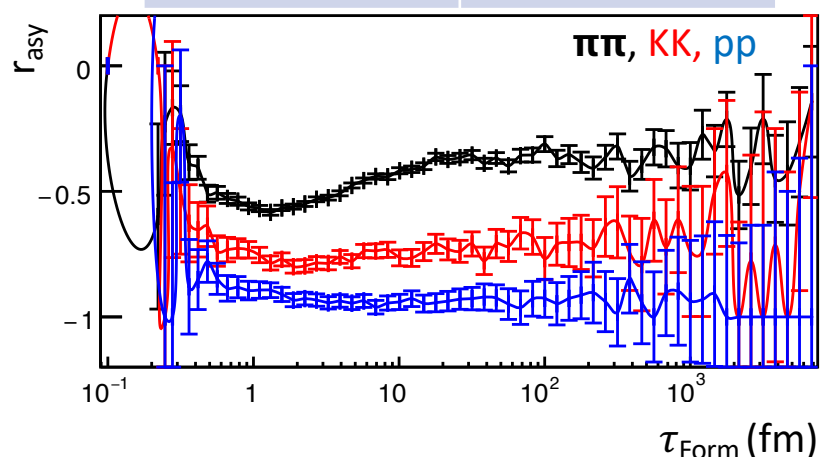
10/2/20

Better measurements at small z region with high momentum PID

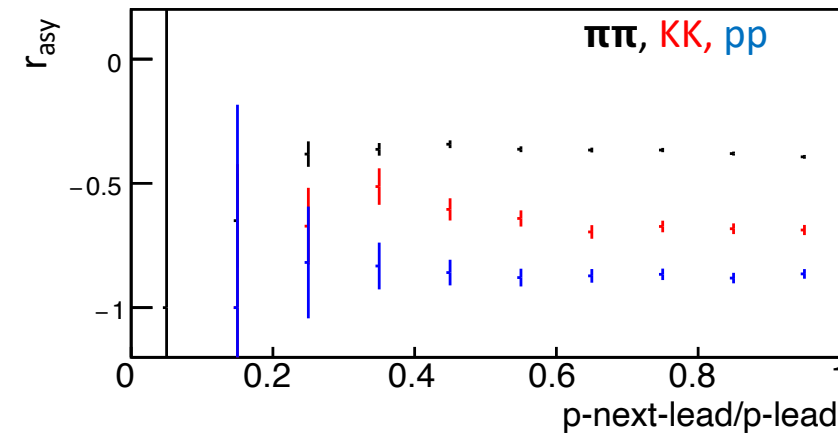
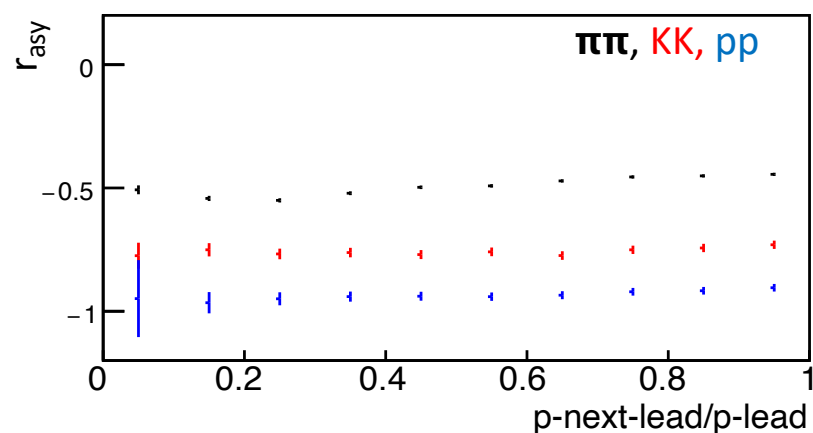
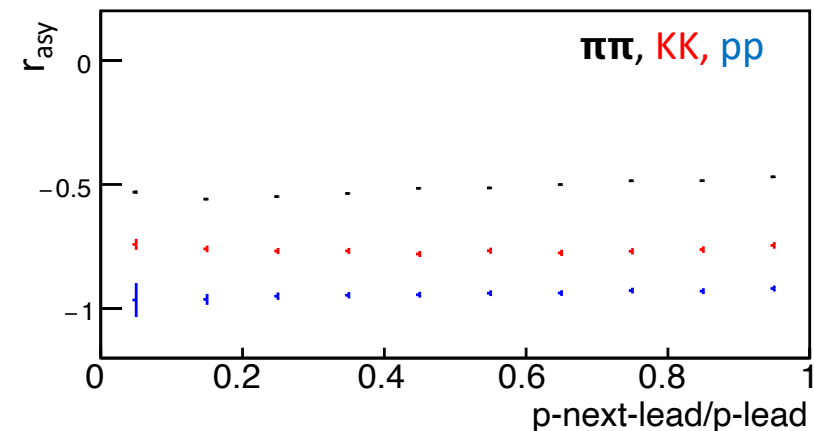
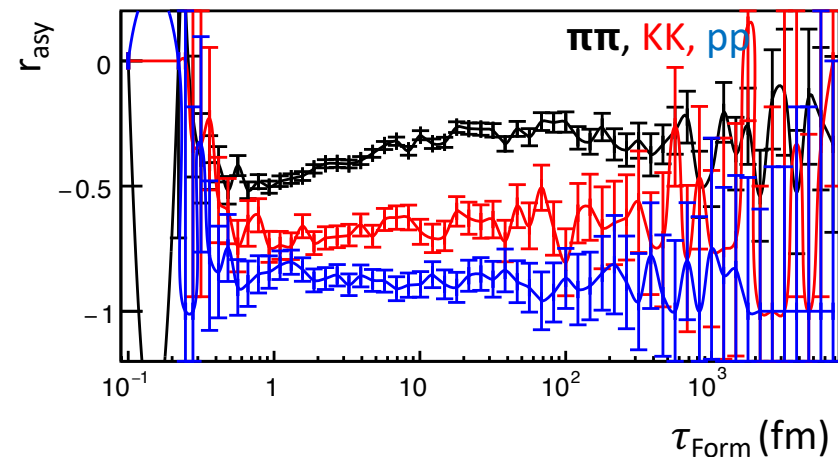
| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -1.0 to 1.0    | No limit                           |



| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -1.0 to 1.0    | 10                                 |



| $\eta$ - range | Momentum cut (GeV/c <sup>2</sup> ) |
|----------------|------------------------------------|
| -1.0 to 1.0    | 5                                  |



- increasing PID performance up to 10 GeV/c<sup>2</sup> would significantly make better measurement in small z and small p-net-lead/p-lead region



# Summary

- Correlations in momentum, charge and flavor of leading particles in jet carry information of non perturbative aspect of jet fragmentation
- At EIC this can be measured with high momentum PID capabilities compared to the experiments in HERA
- PID with 5 GeV/c momentum limit keeps much of the important region inaccessible
- Enhancing momentum reach 10 GeV/c is very useful in reaching the very important inaccessible region
  - ✓ Covers wider  $x$ - $Q^2$  phase space
  - ✓ Access smaller  $z$  region
  - ✓ enhance  $p_{T,Jet}$  limit of measurement
- Such studies going in very forward region

backup

# Measurements of $r_{\text{asy}}$ and expressing in terms of di-hadron fragmentation functions

- **Measurements of  $r_{\text{asy}}$ :**
  - differentially in fractions  $z_L$  and  $z_{NL}$  in a jet,
  - “transverse” kinematic variables:
    - ✓ relative transverse momentum
    - ✓ pair invariant mass
    - ✓ pair formation time
    - ✓ including polarization where applicable

## □ $r_{\text{asy}}$ and its connection with generalized di-hadron fragmentation functions

generalized di-hadron fragmentation functions for any hadrons  $h_1, h_2$ :  $D_{h_L, h_{NL}}^>(z_L z_{NL})$

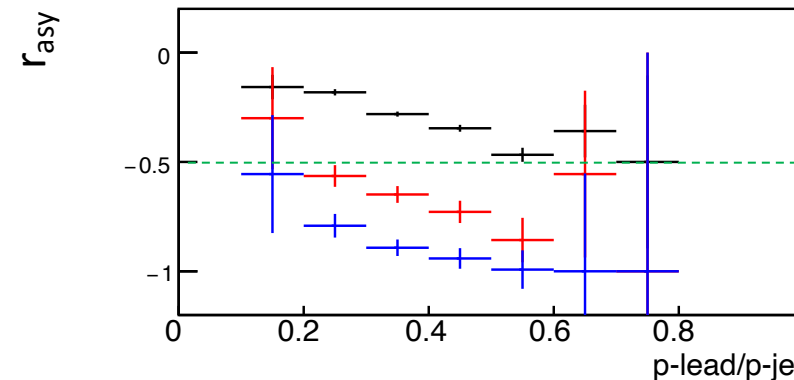
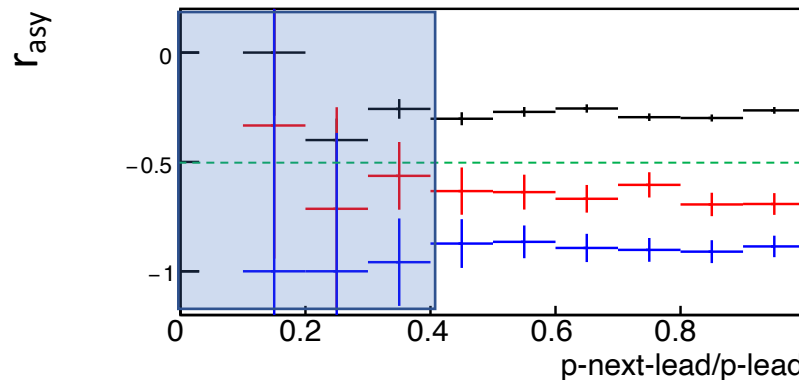
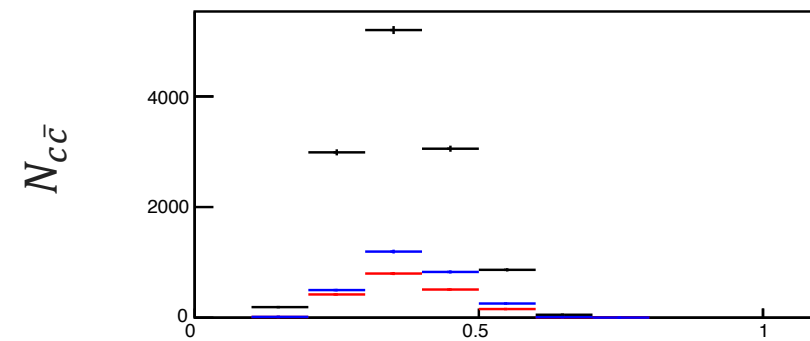
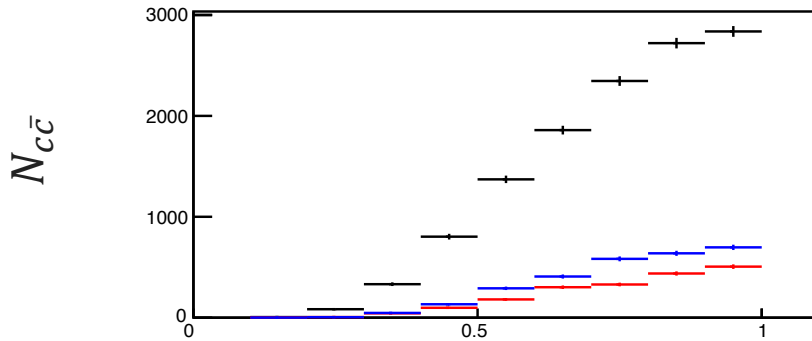
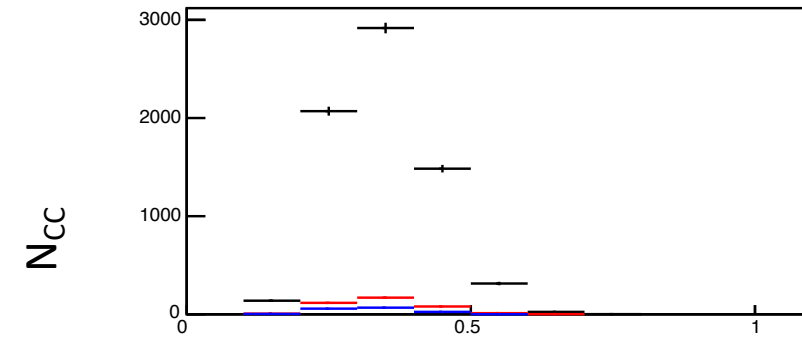
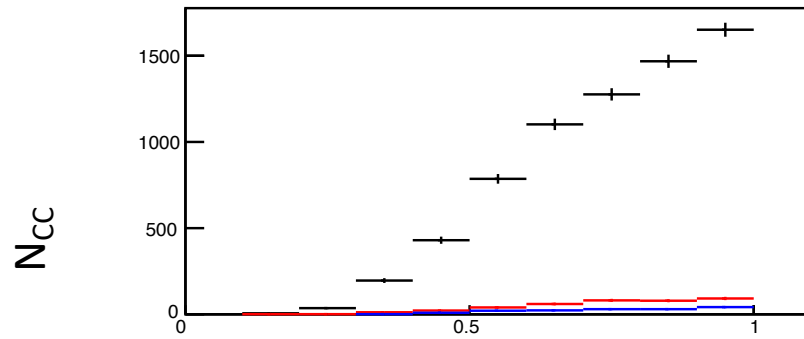
$$N_{h_L, h_{NL}}^> = \int_0^1 dz_L \int^{\min(z_L, 1-z_L)} dz_{NL} D_{h_1, h_2}^>(z_L, z_{NL}, Q)$$

When  $z_L$  and  $z_{NL}$  are large enough, this is the usual di-hadron distribution  $D^>(x_1, x_2, Q) = D(x_1, x_2, Q)$  when  $x_2 > 1 - x_1 - x_2$

$$r_S = \frac{\sum_{h_1, h_2 \in S} Q_{h_1} Q_{h_2} N_{h_1, h_2}^>}{\sum_{h_1, h_2 \in S} |Q_{h_1} Q_{h_2}| N_{h_1, h_2}^>}$$

## After momentum cuts

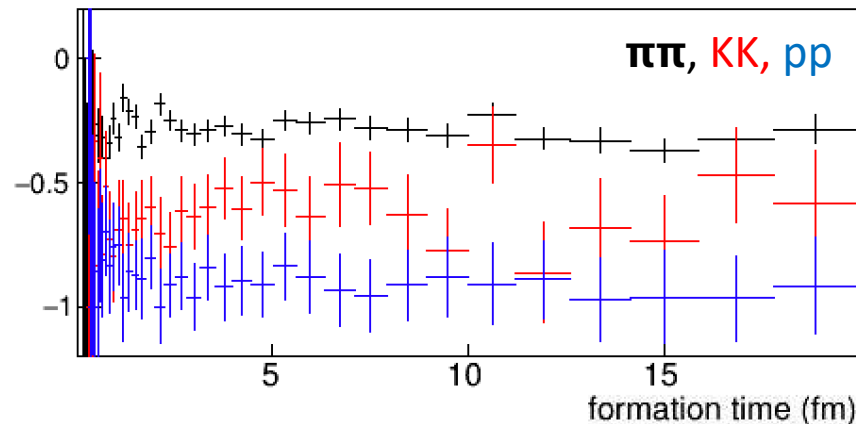
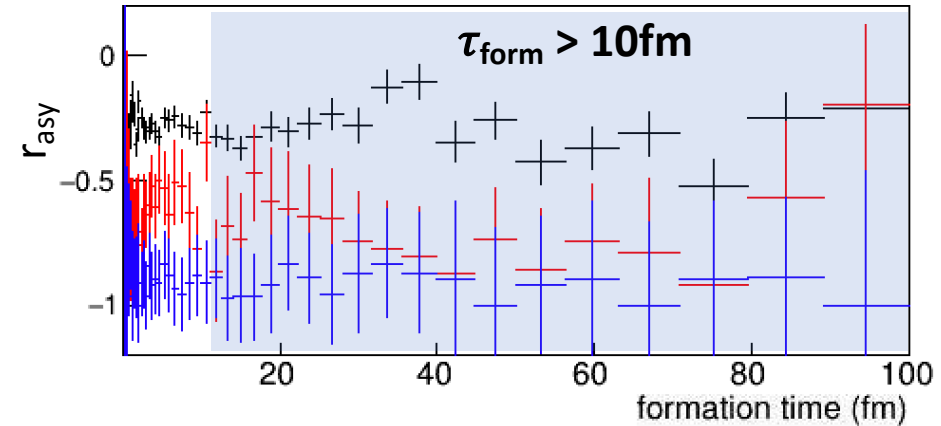
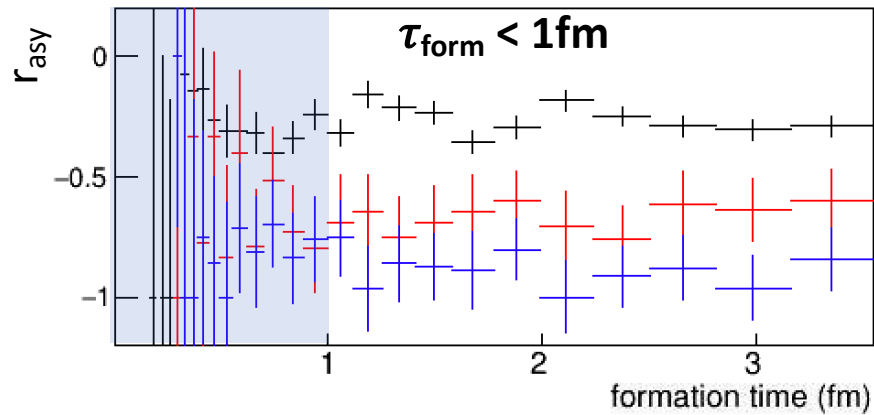
- The strength of correlations are different decrease as an acceptance effect
- Acceptance changes in certain regions due to momentum cuts



The transition region from soft patron emission to hard is very important

# Formation time

$$\text{Formation time, } \tau_{\text{form}} = [2z(1-z) P] / k_{\text{perp}}^2$$



$\tau_{\text{form}} < 1\text{fm}$  : L and NL particles seem to separate after a very short time, which might decorrelate their hadronization.

$\tau_{\text{form}} > 10\text{ fm}$  ( $k_{\text{perp}} < 200\text{ MeV}$ ) : nonperturbative transverse momenta in the jet, and we don't think that going to longer  $\tau_{\text{form}}$  or smaller  $k_{\text{perp}}$  leads to new dynamics

Important region to study in data  $\tau_{\text{form}} =$  "a few fermi" and "a few dozen fermi",  $k_{\text{perp}} =$  "a few GeV" to "several hundred MeV"