



NLO inclusive J/ψ photoproduction at large P_T at HERA and the EIC

Carlo Flore

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(virtual meeting)
13 Apr 2021**

CF, J.-P. Lansberg, H.-S. Shao, Y. Yedelkina, Phys. Lett. B 811 (2020) 135926

Quarkonium Production Model

Phys.Rept. 889 (2020) 1-106 & EPJC (2016) 76:107 for reviews

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 3. COLOR EVAPORATION MODEL:
based on **quark-hadron duality**;
only the invariant mass matters; semi-soft gluons emissions;
color-wise decorrelated $c\bar{c}$ prod. and hadr.

Leading P_T approximation of NLO $J/\psi + g$: NLO[★]

P. Artoisenet *et al.*, PRL 101 (2008) 152001, J.-P. Lansberg, EPJC 61 (2009) 693 & PLB 679 (2009) 340

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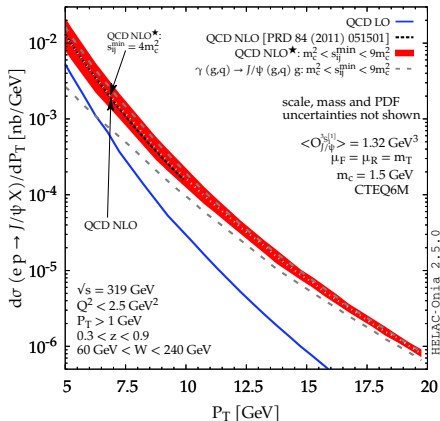
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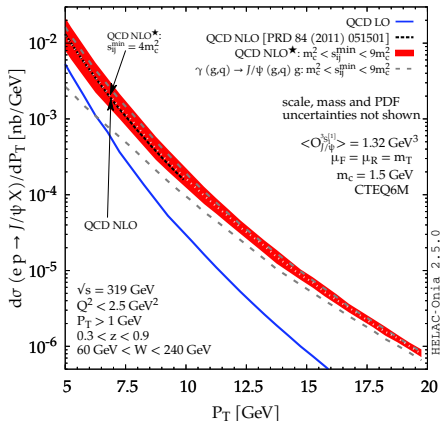
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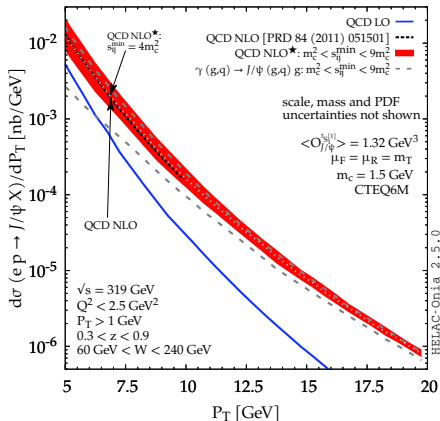


CSM QCD NLO from PRD 84 (2011) 051501

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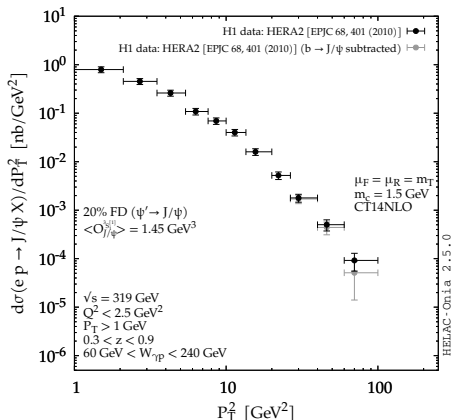
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Let's revisit HERA data!

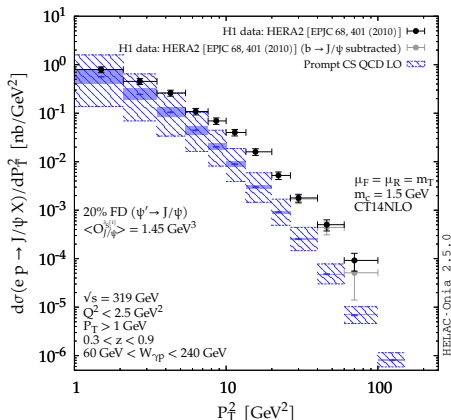
Revisiting HERA data (I)

CF, J.-P. Lansberg, H.-S. Shao, Y. Yedelkina, PLB 811 (2020) 135926



All the computations are done with **HELAC-Onia** [H.-S. Shao, CPC198 (2016) 238]. See also <https://nloaccess.in2p3.fr>

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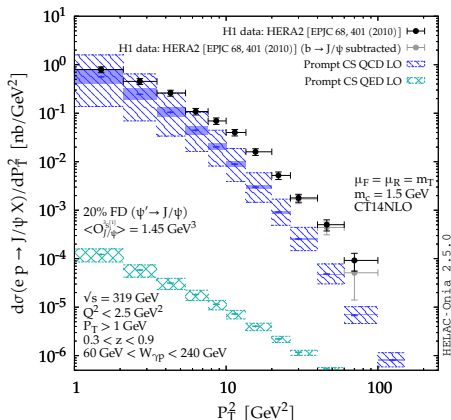


$$\gamma + g \rightarrow \psi + g @ \alpha\alpha_s^2$$

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[The quark and antiquark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.]

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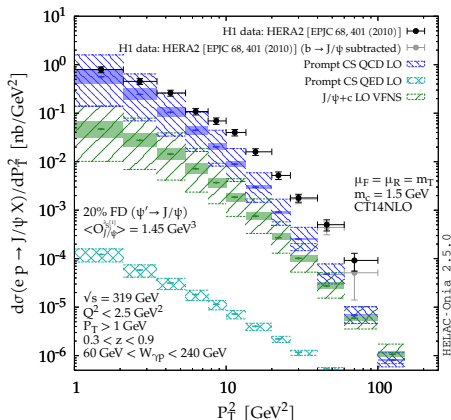


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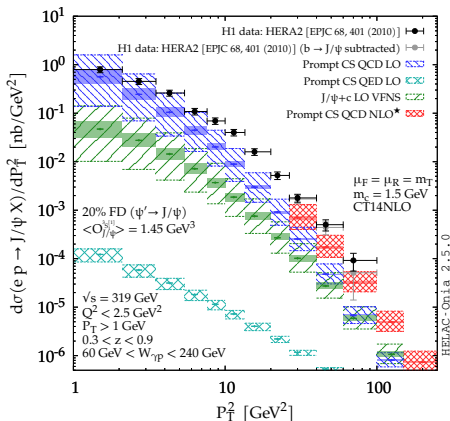


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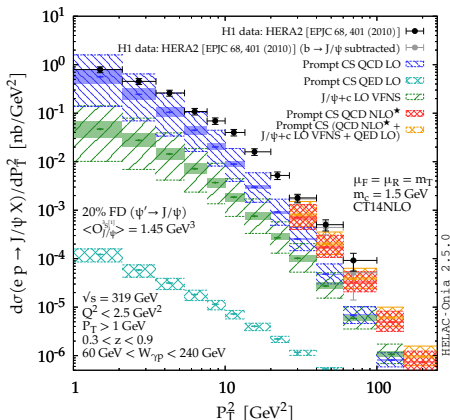
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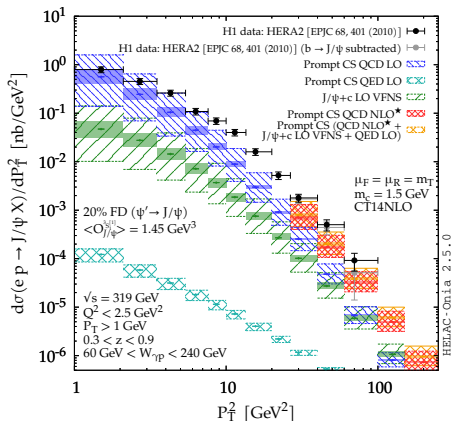
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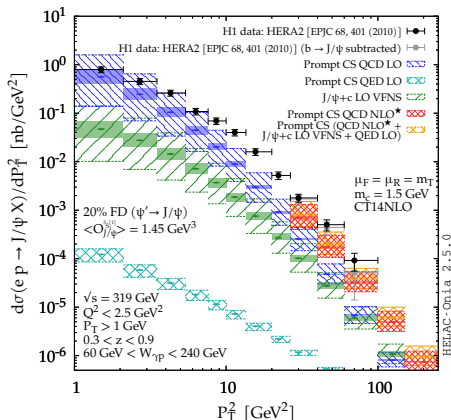
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- LO QCD works well at low P_T

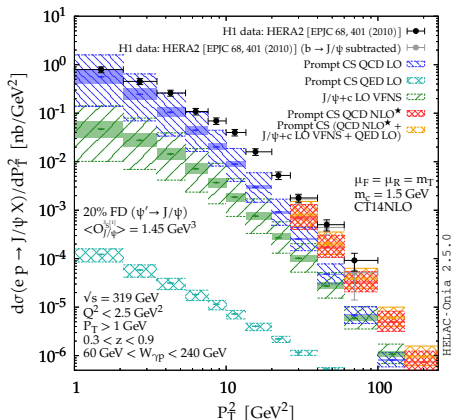
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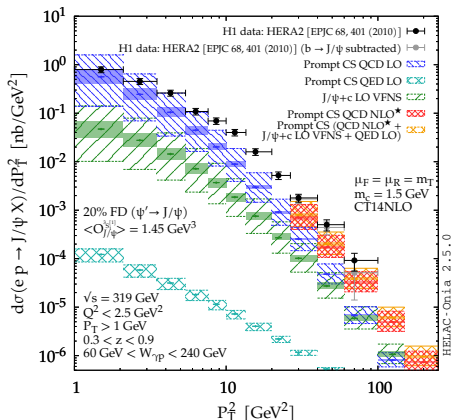
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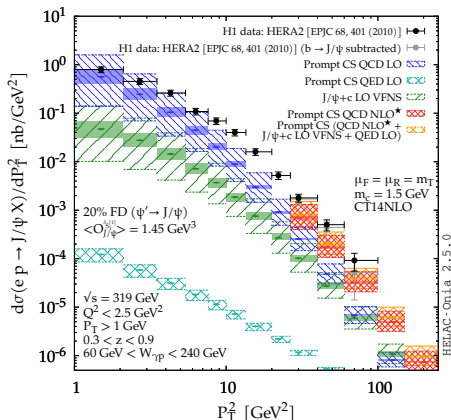
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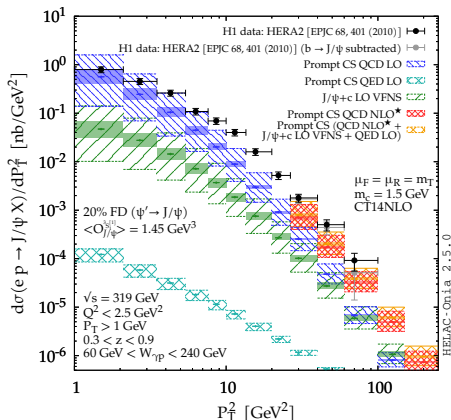


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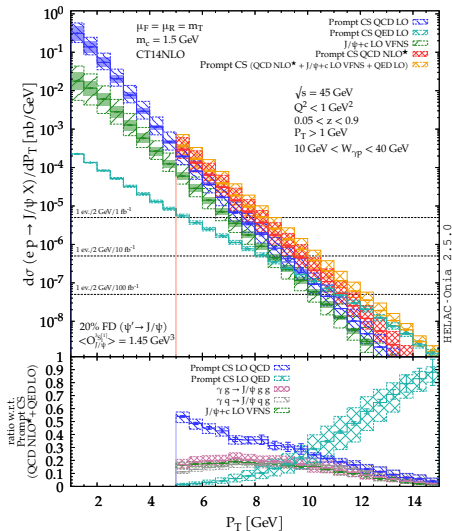


The CSM up to $\alpha\alpha_s^3$ reproduces J/ψ photoproduction at HERA

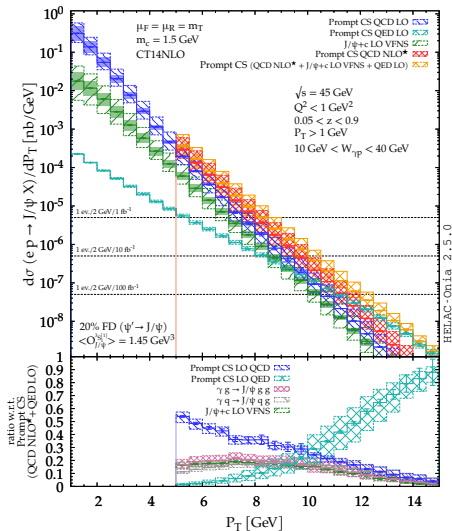
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→ we will restrict to CSM for our EIC predictions

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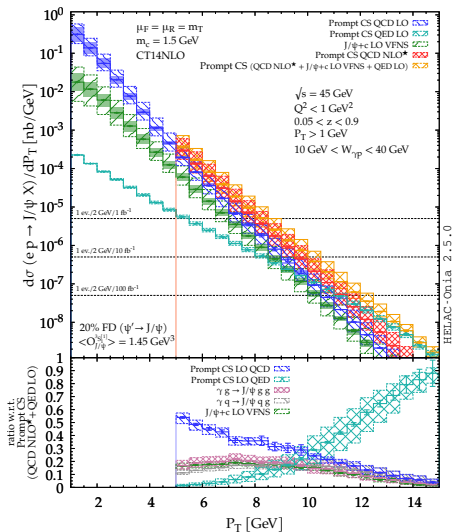


- At $\sqrt{s_{ep}} = 45 \text{ GeV}$, one gets into valence region



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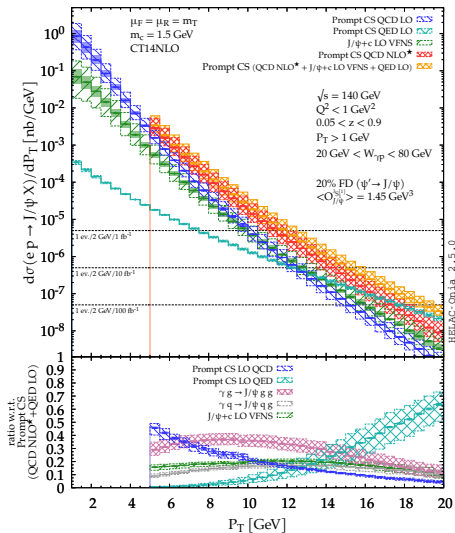
- At $\sqrt{s_{ep}} = 45 \text{ GeV}$, one gets into **valence region**
- Yield steeply falling with P_T
- Yield can be measured **up to** $P_T \sim 11 \text{ GeV}$ with $\mathcal{L} = 100 \text{ fb}^{-1}$
[using both ee and $\mu\mu$ decay channels and $\epsilon_{J/\psi} \simeq 80\%$]



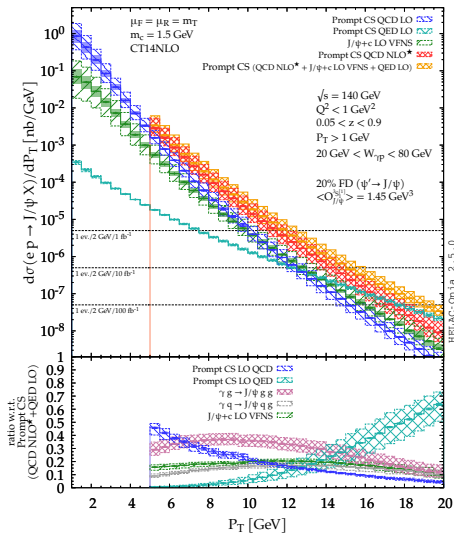
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- QED** contribution **leading** at the largest reachable P_T
- $\gamma + q$ fusion contributes more than 30% for $P_T > 8 \text{ GeV}$

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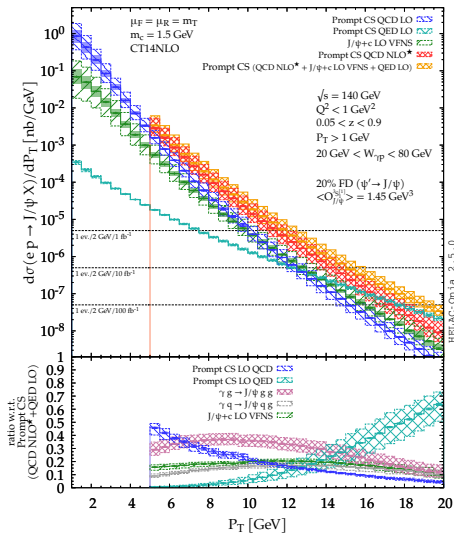


- At $\sqrt{s_{ep}} = 140 \text{ GeV}$ larger P_T range, up to $\sim 18 \text{ GeV}$



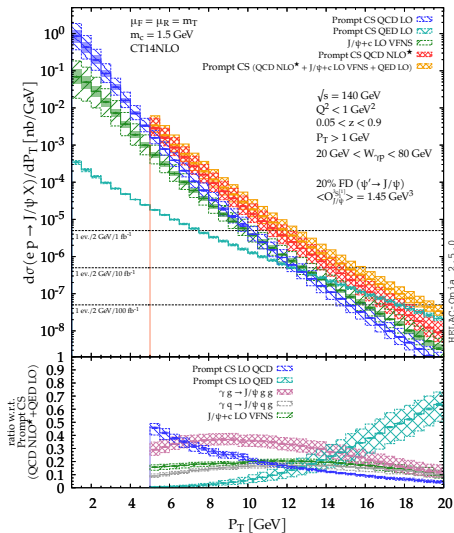
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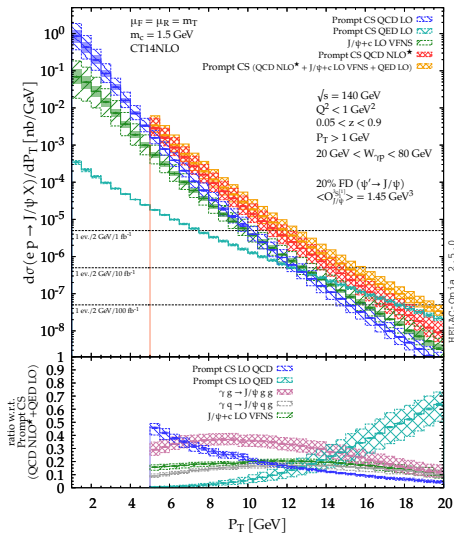
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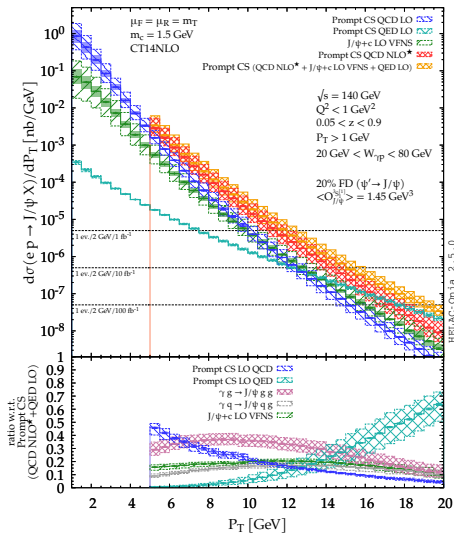
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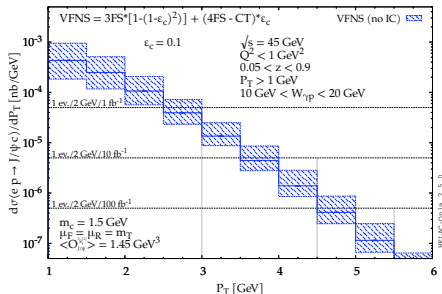
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- We expect the $d\sigma$ to vanish when $E_{\text{jet}_2}^{J/\psi \text{ rest fr.}} \rightarrow 0$

J/ψ + charm associated production at the EIC

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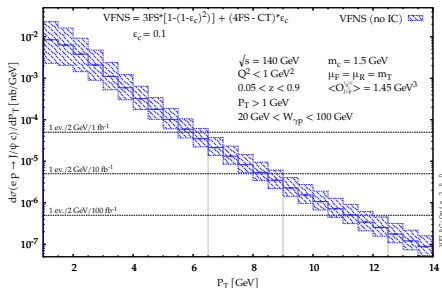
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$$d\sigma^{\text{VFNS}} = d\sigma^{\text{3FS}} [1 - (1 - \varepsilon_c)^2] + (d\sigma^{\text{4FS}} - d\sigma^{\text{CT}}) \varepsilon_c$$

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- But it is clearly observable if $\varepsilon_c = 0.1$ with **$\mathcal{O}(500, 50, 5)$ events** for **$\mathcal{L} = (100, 10, 1) \text{ fb}^{-1}$**

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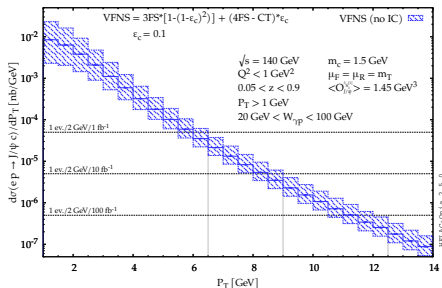
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- At $\sqrt{s_{\text{sep}}} = 140 \text{ GeV}$, P_T range up to 10 GeV with **up to thousands of events with $\mathcal{L} = 100 \text{ fb}^{-1}$**
- Could be observed via **charm jet**

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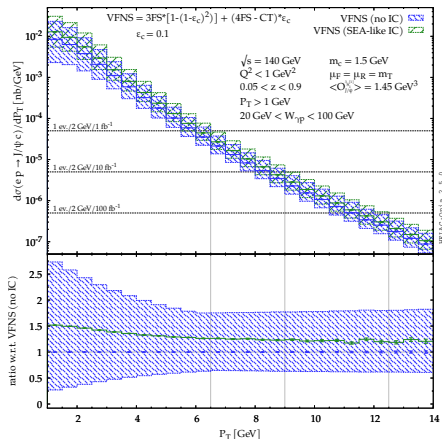
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J/ψ + charm associated production at the EIC



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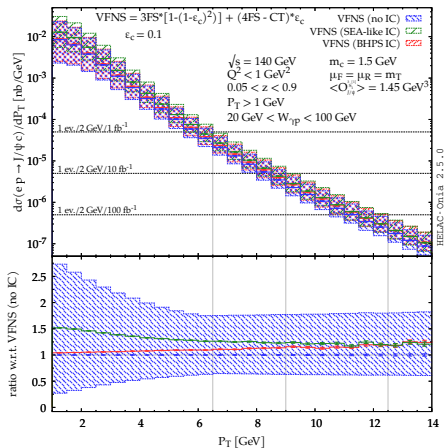
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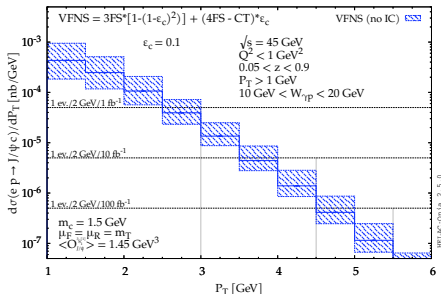
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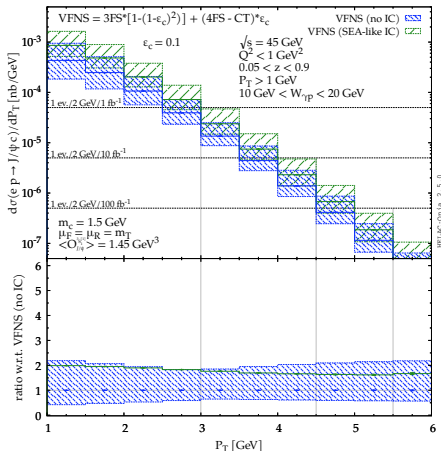
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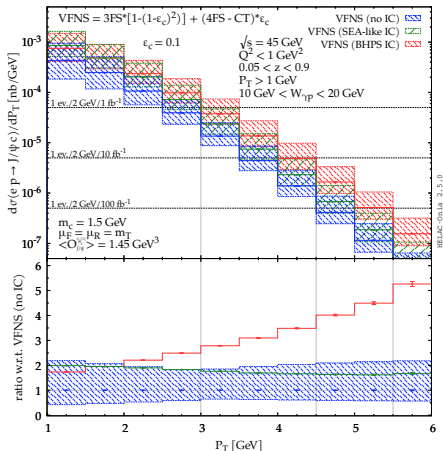
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- $\sqrt{s_{ep}} = 45$ GeV:
 - $\gamma + q$ QED contribution [new!] leading at high P_T
 - J/ψ +charm sensitive to charm PDFs

Thank you

Backup

Kinematics and cross section

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- $s_{ep} = (P_e + P_p)^2 = 4E_e E_p$; $s_{\gamma p} = W_{\gamma p}^2 = (P_\gamma + P_p)^2$; $P_\gamma = x_\gamma P_e$, so $s_{\gamma p} = x_\gamma s_{ep}$
- $z = \frac{P_Q \cdot P_p}{P_\gamma \cdot P_p}$: fraction of the photon energy taken by the J/ψ in the proton rest frame
- cross section:

$$\frac{d\sigma}{dz dP_T} = \int_{x_\gamma^{\min}}^1 dx_\gamma \frac{2x_a P_T f_{\gamma/e}(x_\gamma, Q_{\max}^2) f_{a/p}(x_a(x_\gamma), \mu_F)}{z(1-z)} \\ \times \frac{1}{16\pi s^2} |\mathcal{M}(\gamma + a \rightarrow Q + k)|^2,$$

where $x_a = \frac{M_T^2 - m_Q^2 z}{x_\gamma s_{ep} z(1-z)}$ and $x_\gamma^{\min} = \frac{M_T^2 - m_Q^2 z}{s_{ep} z(1-z)}$

- WW distribution

$$f_{\gamma/e}(x_\gamma, Q_{\max}^2) = \frac{\alpha}{2\pi} \left[\frac{1 + (1 - x_\gamma)^2}{x_\gamma} \ln \frac{Q_{\max}^2}{Q_{\min}^2(x_\gamma)} + 2m_e^2 x_\gamma \left(\frac{1}{Q_{\max}^2} - \frac{1}{Q_{\min}^2(x_\gamma)} \right) \right]$$

where $Q_{\min}^2(x_\gamma) = m_e^2 x_\gamma^2 / (1 - x_\gamma)$

VFNS treatment of J/ψ +charm yield

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J/ψ +charm production follows from

- $\gamma + g \rightarrow J/\psi + c + \bar{c} @ \alpha\alpha_S^3$ (3FS)

- $\gamma + \{c, \bar{c}\} \rightarrow J/\psi + \{c, \bar{c}\} @ \alpha\alpha_S^2$ (4FS)

$$d\sigma_{\gamma c \rightarrow J/\psi + c} = \frac{1}{2(\hat{s} - m_c^2)} dx_c f_{c/p}(x_c, \mu_F^2) \times \overline{|\mathcal{M}_{\gamma c \rightarrow J/\psi c}|^2} d\Phi(p_\gamma, p_c \rightarrow P_Q, p'_c),$$

with

$$f_{c/p}(x_c, \mu_F^2) = \tilde{f}_{c/p}^{(1)}(x_c, \mu_F^2) + \mathcal{O}(\alpha_S^2),$$

where

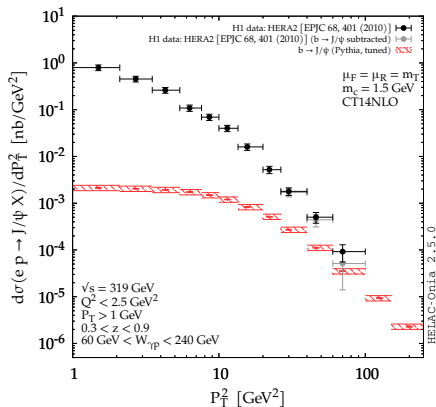
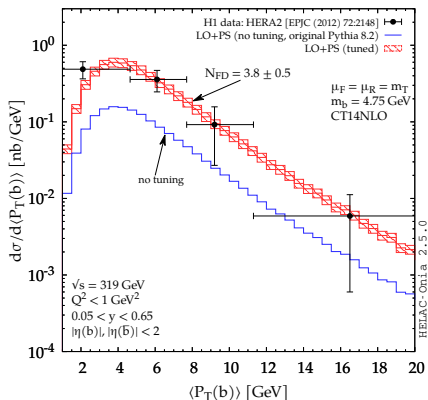
$$\tilde{f}_{c/p}^{(1)}(x_c, \mu_F^2) = \frac{\alpha_S}{2\pi} \log\left(\frac{\mu_F^2}{m_c^2}\right) \int_{x_c}^1 \frac{dz}{z} P_{qg}(z) f_{g/p}\left(\frac{x_c}{z}, \mu_F^2\right)$$

with AP splitting function $P_{qg}(z) = \frac{1}{2} [z^2 + (1-z)^2]$. Overlap CT to be subtracted from 3FS:

$$d\sigma_{\text{CT}, \gamma c \rightarrow J/\psi + c} = \frac{1}{2(\hat{s} - m_c^2)} dx_c \tilde{f}_{c/p}^{(1)}(x_c, \mu_F^2) \overline{|\mathcal{M}_{\gamma c \rightarrow J/\psi c}|^2} d\Phi(p_\gamma, p_c \rightarrow P_Q, p'_c).$$

$b \rightarrow J/\psi$ feed-down

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H1 $b\bar{b}$ production from EPJC 72 (2012) 2148 [Note: $\langle P_T(b) \rangle = \sqrt{(P_{T,b}^2 + P_{T,\bar{b}}^2)/2}$]

$N_{FD} = 3.8 \pm 0.5$ estimated via χ^2 -minimisation

$b \rightarrow J/\psi$ tuned with PYTHIA 8.2