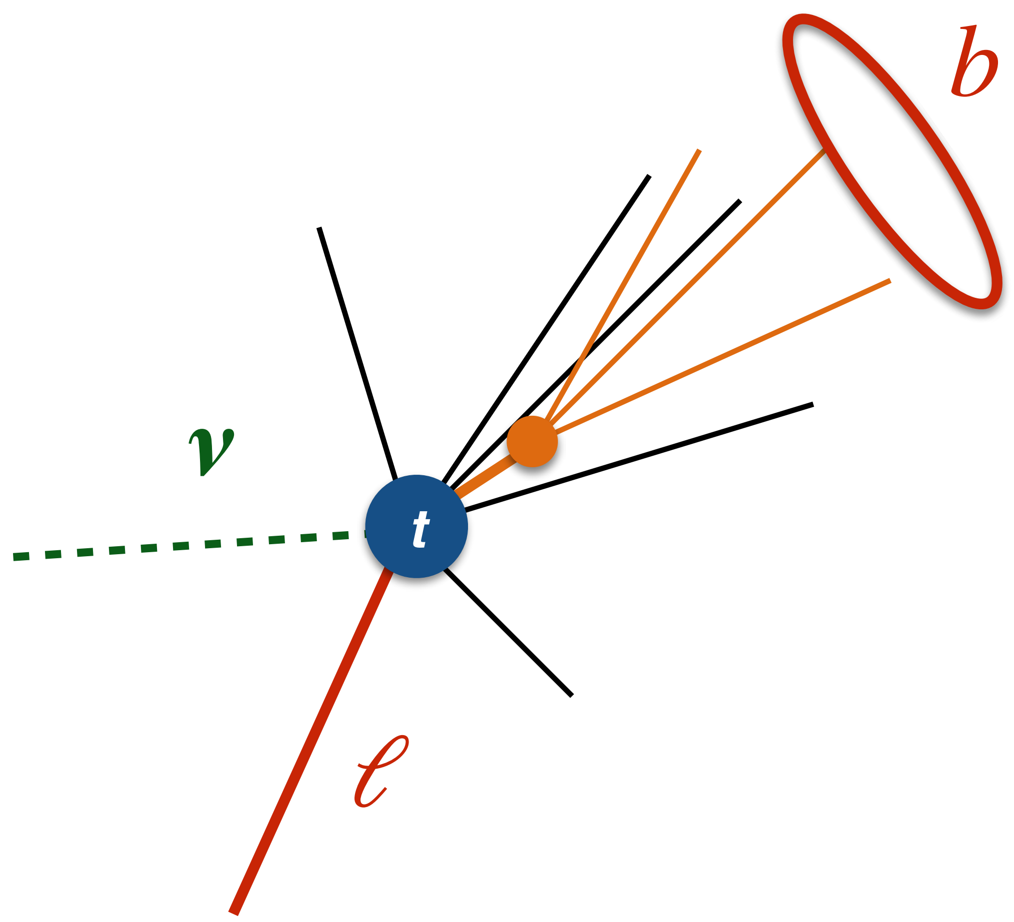
measuring b-fragmentation in $t\bar{t}$ events with ATLAS



Chris Pollard

on b-half of the ATLAS collaboration

DIS

2021 04 14

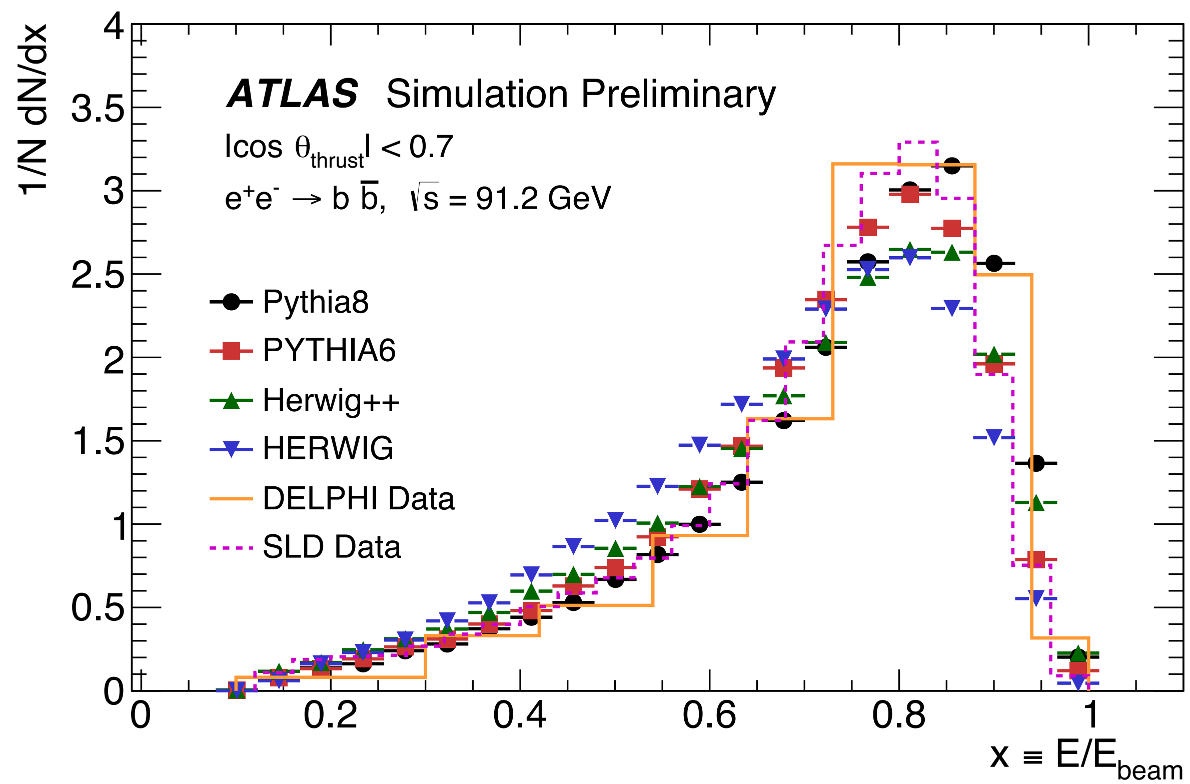
https://cds.cern.ch/record/2730444

ATLAS-CONF-2020-050

motivation I

the fragmentation of b-quarks into hadrons is of interest for many reasons.

- b-hadrons leave a ***striking experimental signature*** and…
- a ***unique correspondence*** to the originating b-quarks
  - ergo a ***precise probe of QCD***
- b-fragmentation currently ***tuned to e^+e^- data*** …
- then ***extrapolated to the LHC*** environment
  - to what degree is this correct?

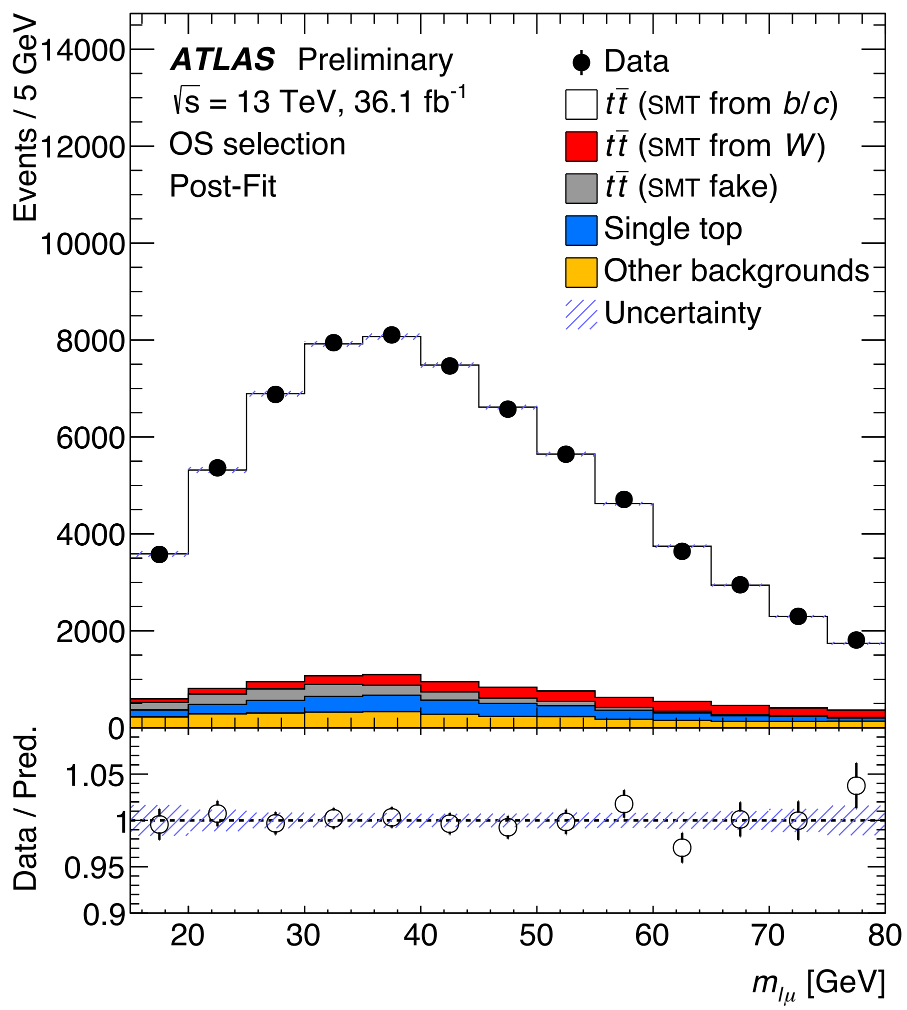


ATLAS-PHYS-PUB-2014-008

some tension between $e^+e^- \to bb$ measurements of b-fragmentation

parton-shower generators are also ***not in good agreement***

motivation II

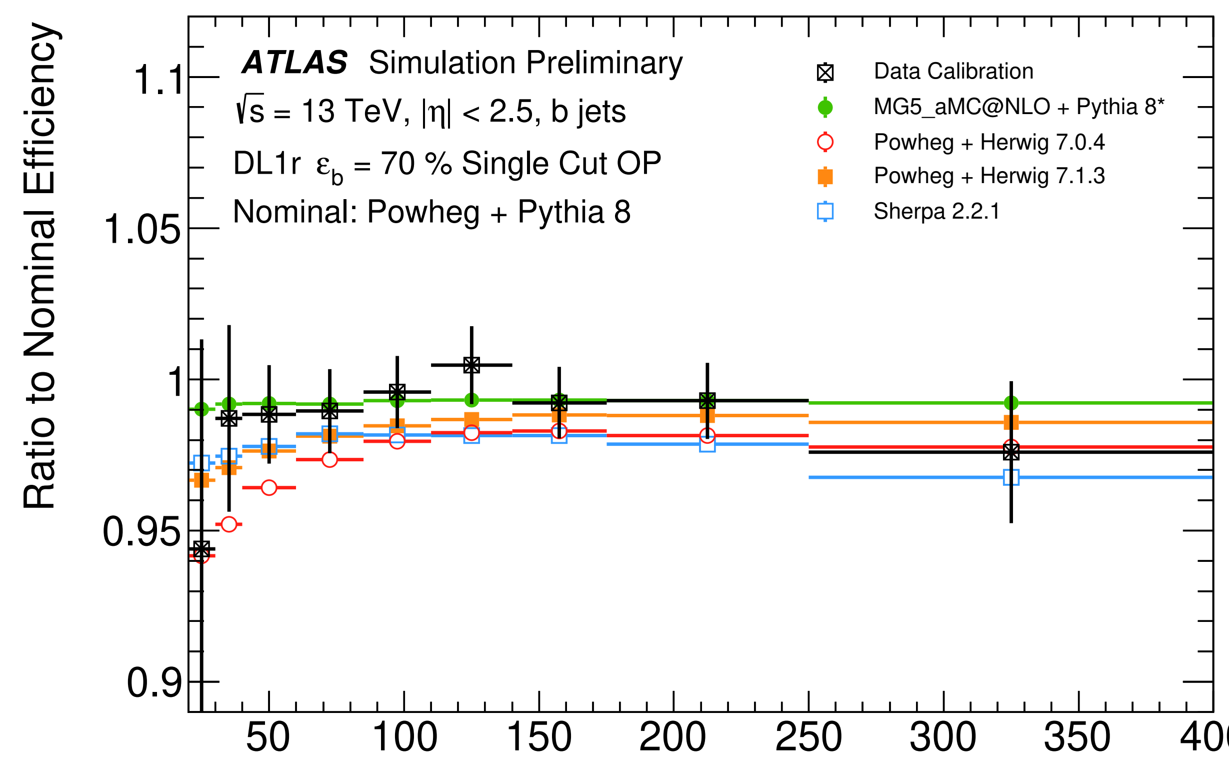


ATLAS-CONF-2019-046

$b \to \mu$ momentum transfer key for new top mass measurements, which rely on precise $b \to B$ prediction

critical for delivering the best LHC physics results with b-jets.

b-jet reconstruction and identification are very sensitive to b-fragmentation.



ATLAS-PHYS-PUB-2020-009

analysis strategy

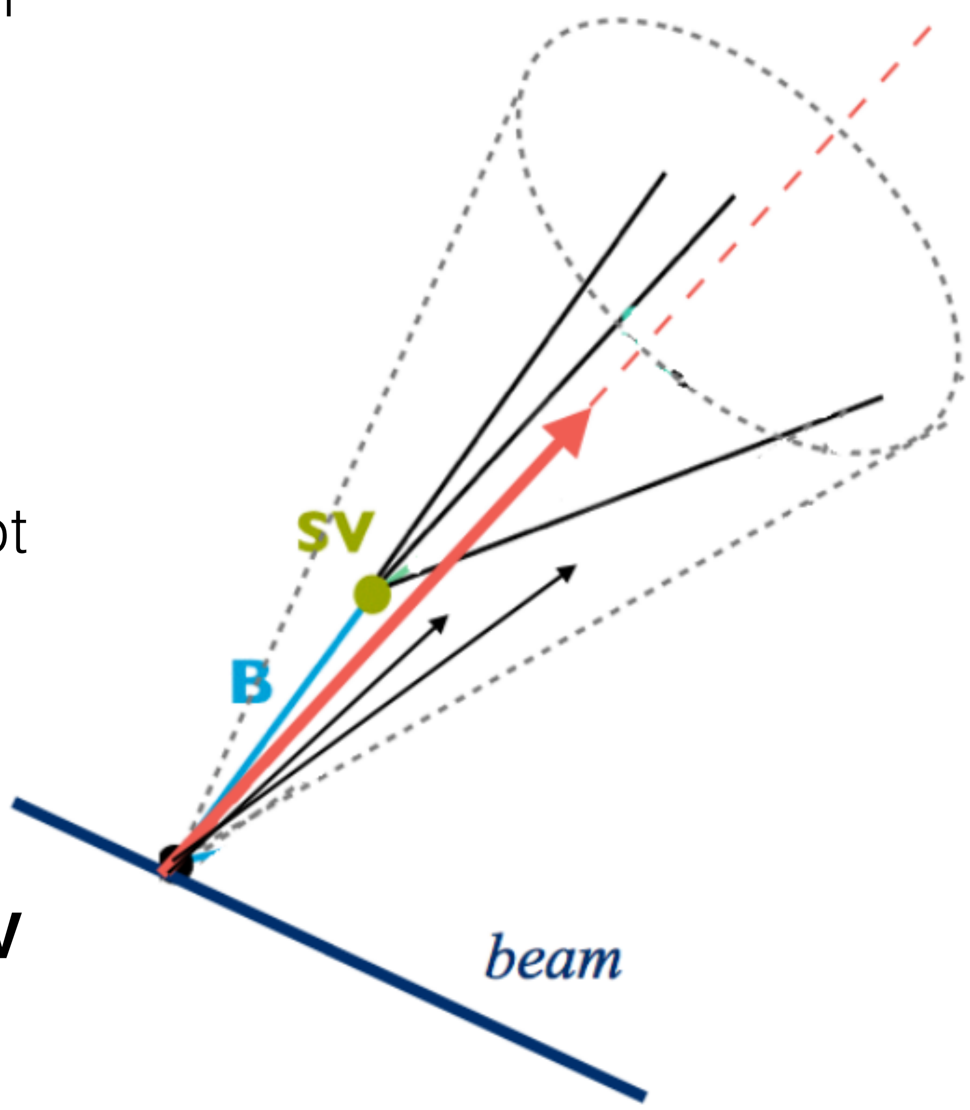
1. obtain sample of b-jets from $t\bar{t} \rightarrow e\mu\nu\nu bb$ events

2. identify b-hadron candidates using ATLAS single secondary vertex finder (SSVF)

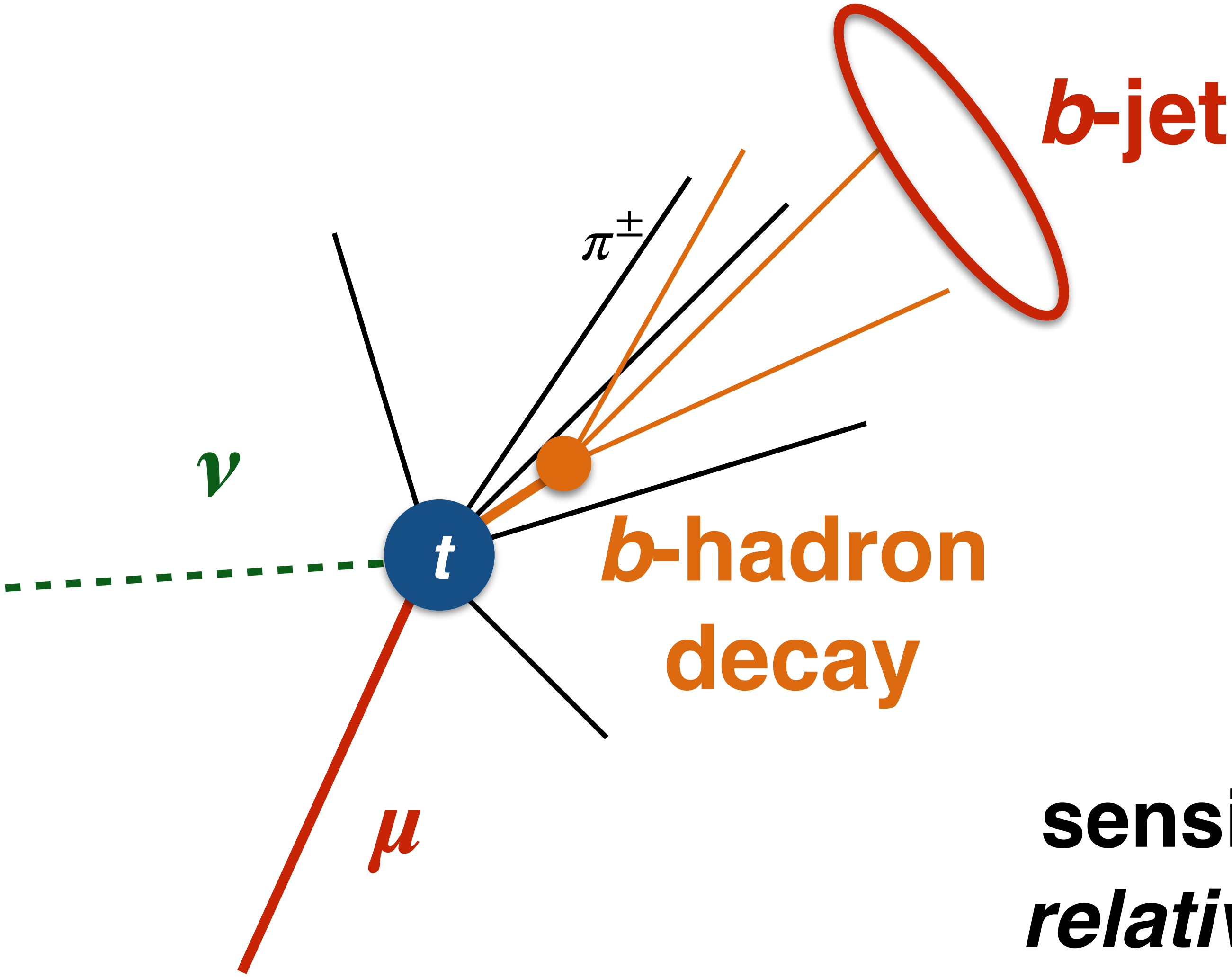
3. measure properties of the b-hadron sensitive to the modeling of b-quark fragmentation, unfolded to particle level

fiducial definition

- $t\bar{t}$ events (single top is treated as background for ease of interpretation)
- 1 μ, 1 e
- exactly 2 b-jets with $|\eta| < 2.1$, $p_T > 30$ GeV, $\Delta R(jet, jet) > 0.5$
  - jet of stable particles, excluding neutrinos and prompt leptons
  - 1 weakly-decaying b-hadron ($p_T > 5$ GeV) with ΔR (jet, hadron) < 0.3
- **sum over stable charged particles with $p_T > 500$ MeV**
  - b-hadron decay products make up $\vec{p}_b^{\,\mathrm{ch}}$
  - particles ghost-associated to jet make up $\vec{p}_{\mathrm{jet}}^{\,\mathrm{ch}}$



observables I



$$z^{\text{ch}}_{\text{L,b}} = \frac{\vec{p}^{\,\text{ch}}_{\text{b}} \cdot \vec{p}^{\,\text{ch}}_{\text{jet}}}{|p^{\text{ch}}_{\text{jet}}|^2},$$

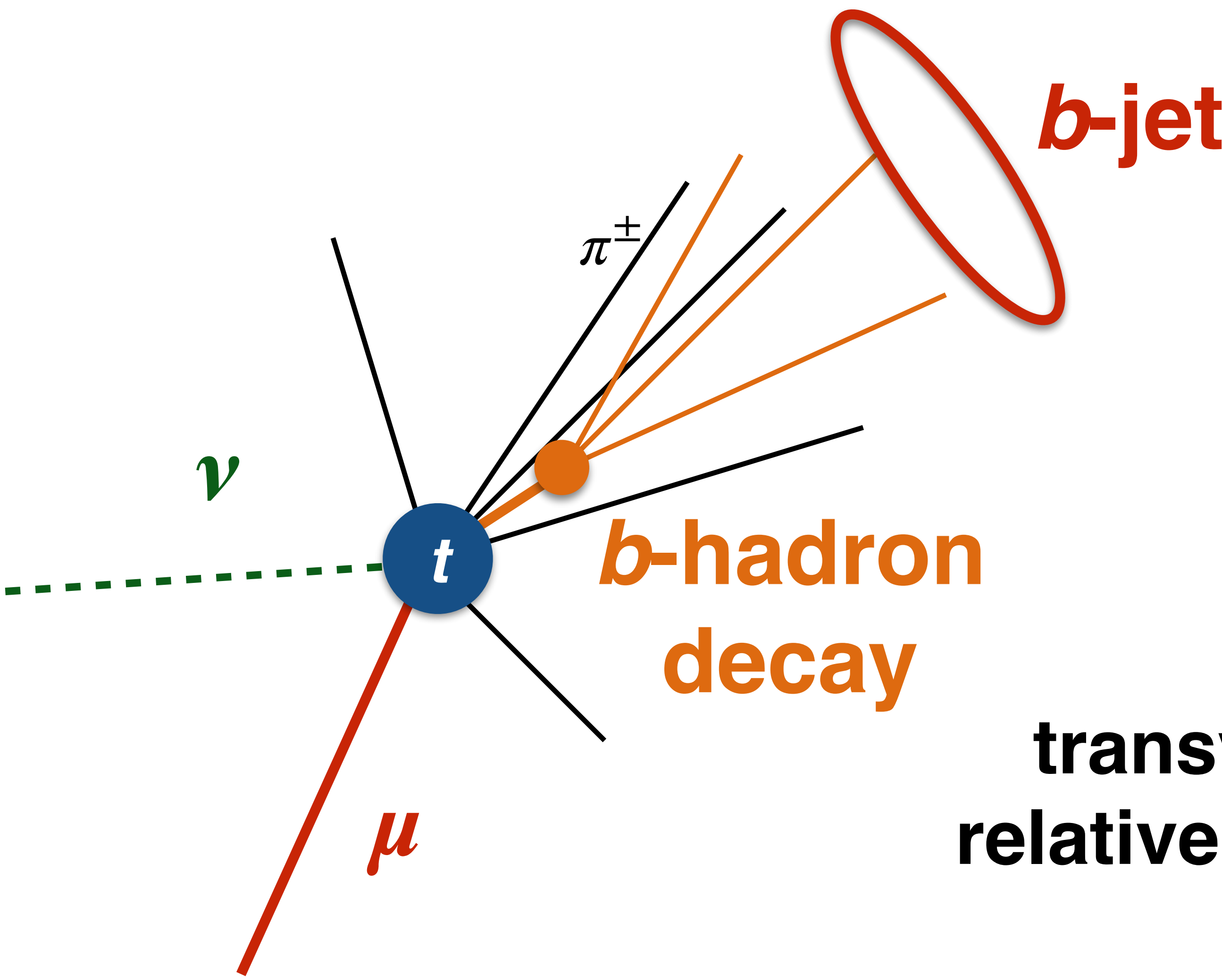
$$z^{\text{ch}}_{\text{T,b}} = \frac{p^{\text{ch}}_{\text{T,b}}}{p^{\text{ch}}_{\text{T,jet}}},$$

sensitive to momentum of b-hadron *relative to the local hadronic activity,*

e.g. the b-hadron momentum compared to the b-quark momentum

and related model parameters: r_B, α_S^{FSR}

observables II



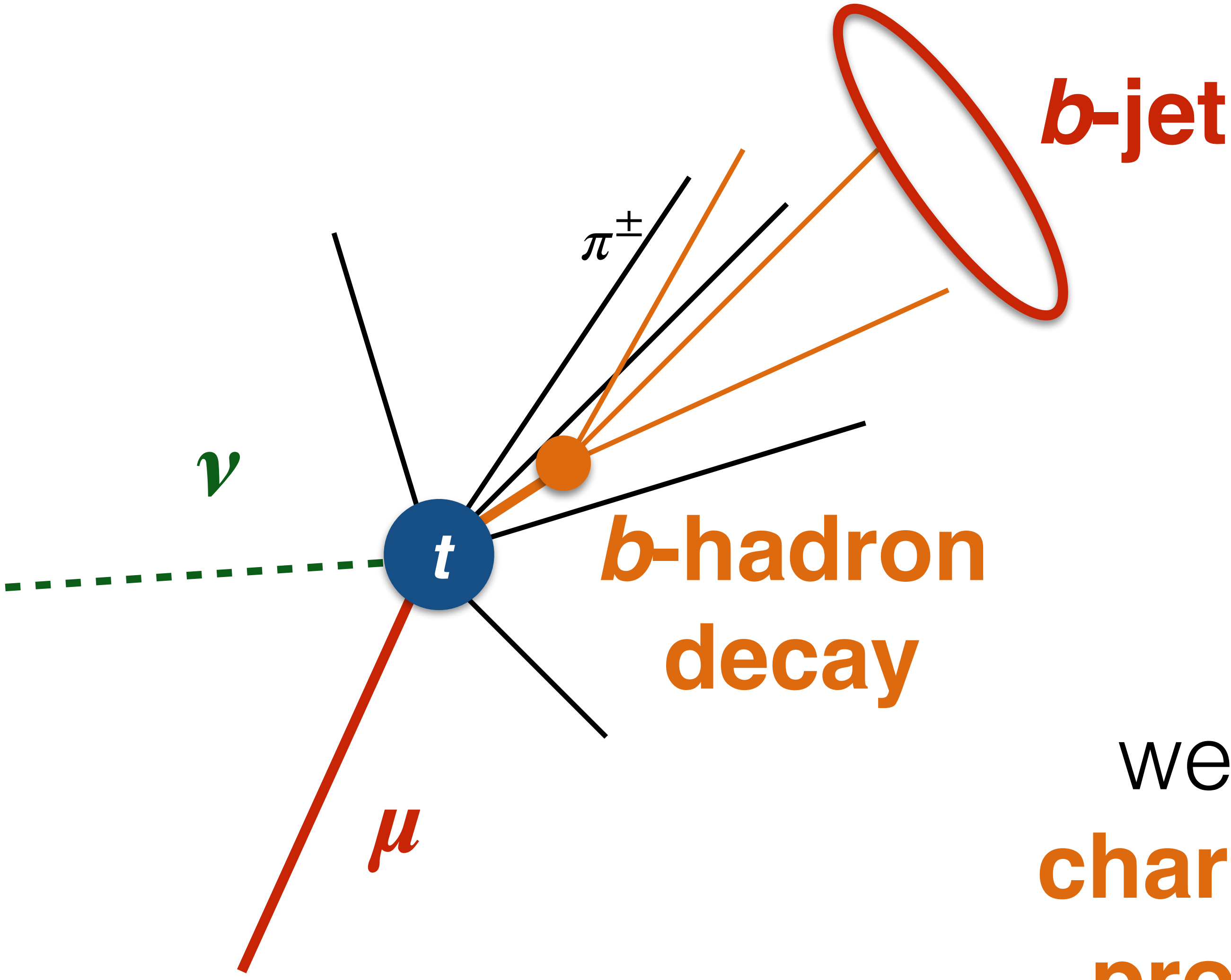
$$\rho = \frac{2p_{\mathrm{T,b}}^{\mathrm{ch}}}{p_{\mathrm{T}}^{e} + p_{\mathrm{T}}^{\mu}}$$

transverse momentum of *b*-hadron relative to the average pt of the leptons in the event

sensitive to choice of e.g. $\alpha_S(FSR)$ but ***not*** $\alpha_S(ISR)$

sensitivity also to wide-angle radiation in top-quark decays

observables III



we also unfold the **number of charged, stable b-hadron decay products with $p_T > 500$ MeV**.

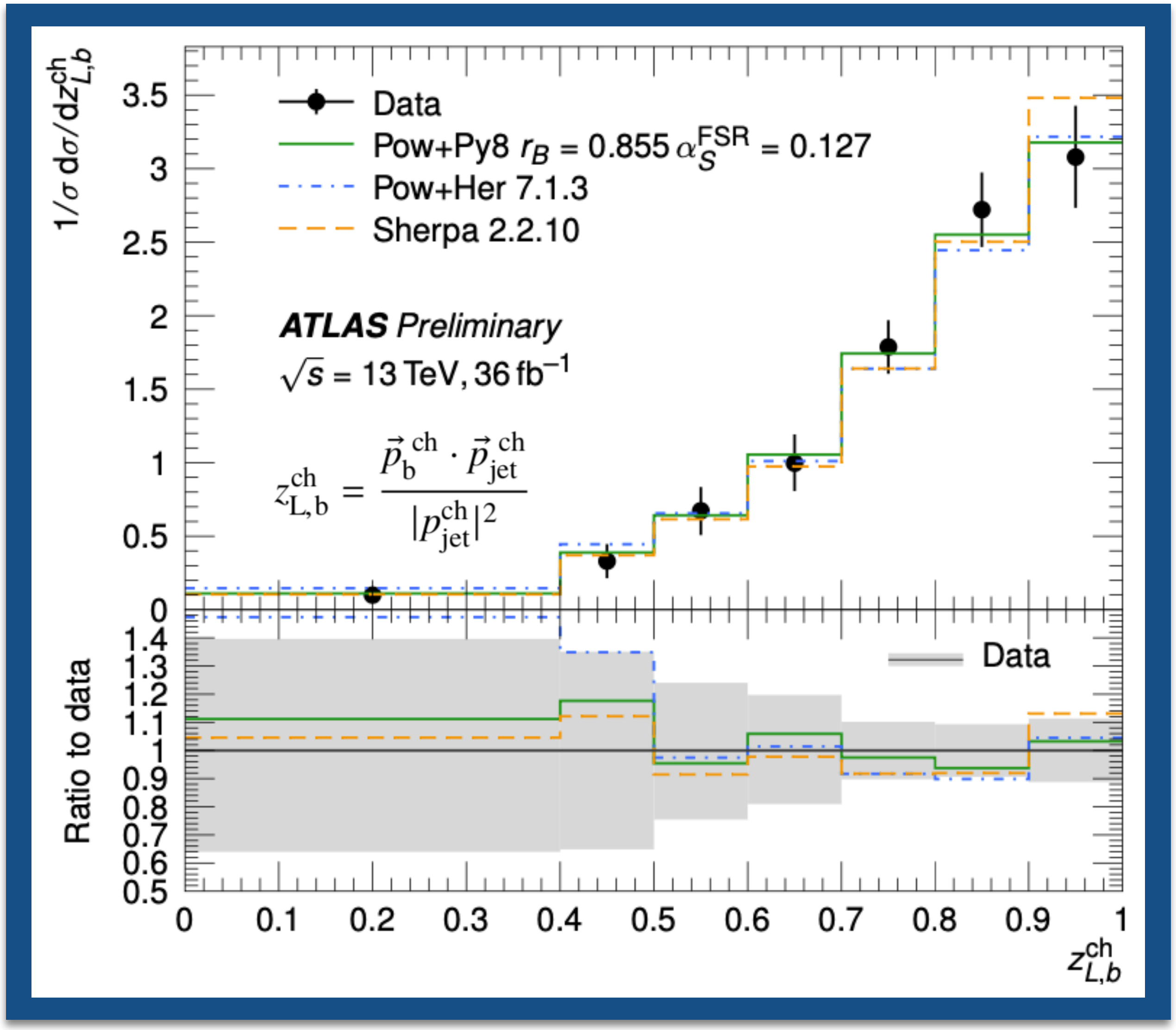
sensitive to the of b-hadron species production rates

systematic uncertainties

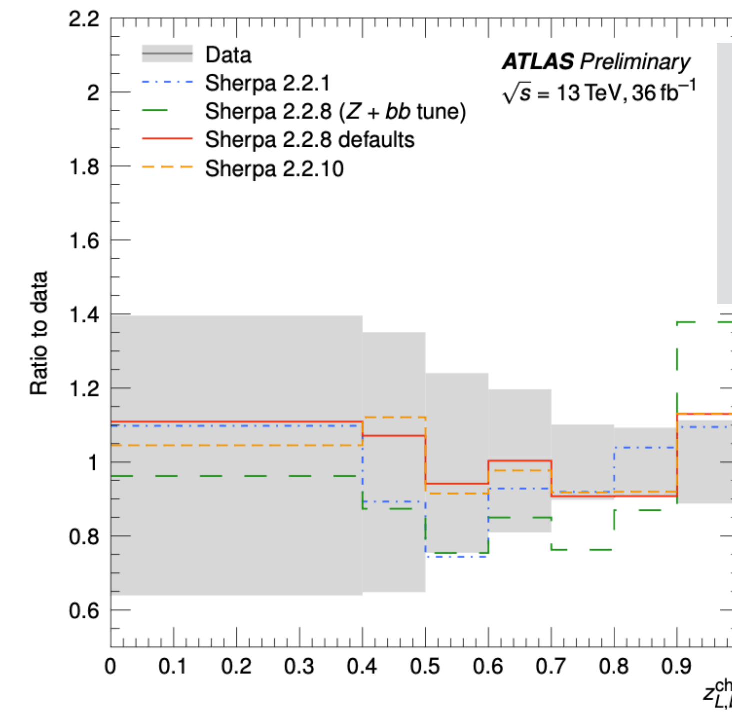
the dominant uncertainties in the unfolding model come from

- tracking efficiency and impact-parameter resolutions
- $t\bar{t}$ modeling:
    - Pythia vs Herwig
    - α_S^{FSR} tune variations
    - B production fractions
    - b-jet p_T mismodeling
- pileup modelling

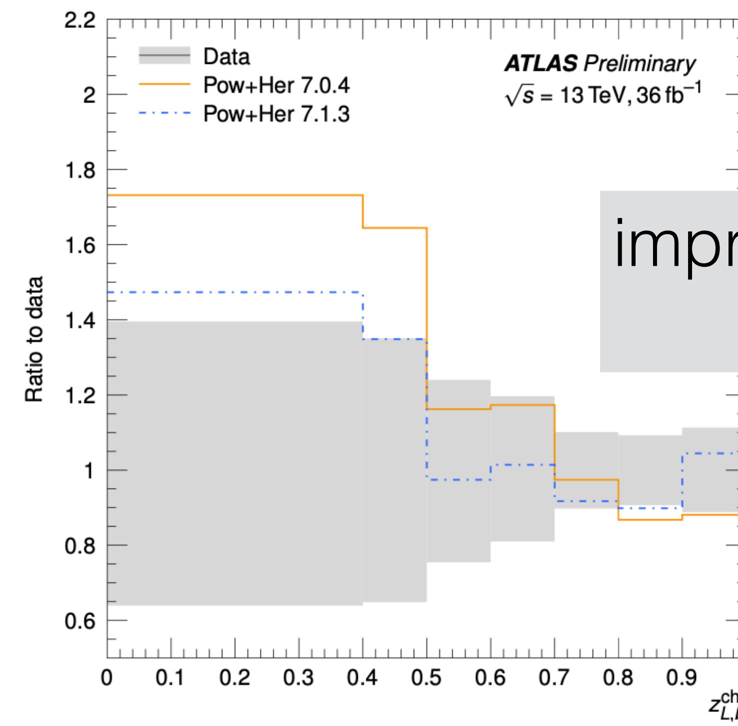
fiducial data/MC comparisons I



recent generator tunes
to $e^+e^- \rightarrow Z \rightarrow bb$
do a reasonable job.

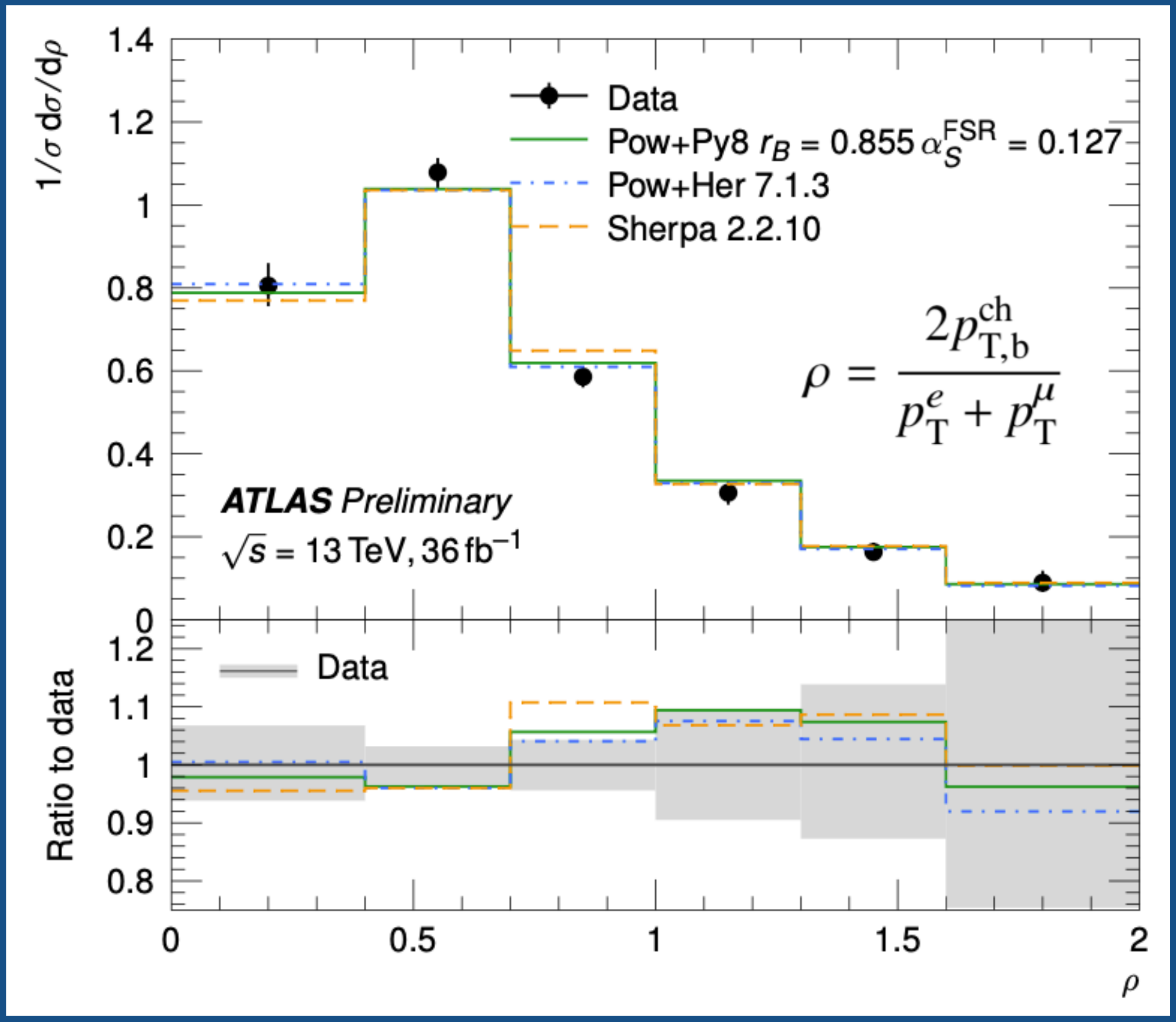


Sherpa: 2.2.10 with default tune looks good.



improvements in newer versions of H7

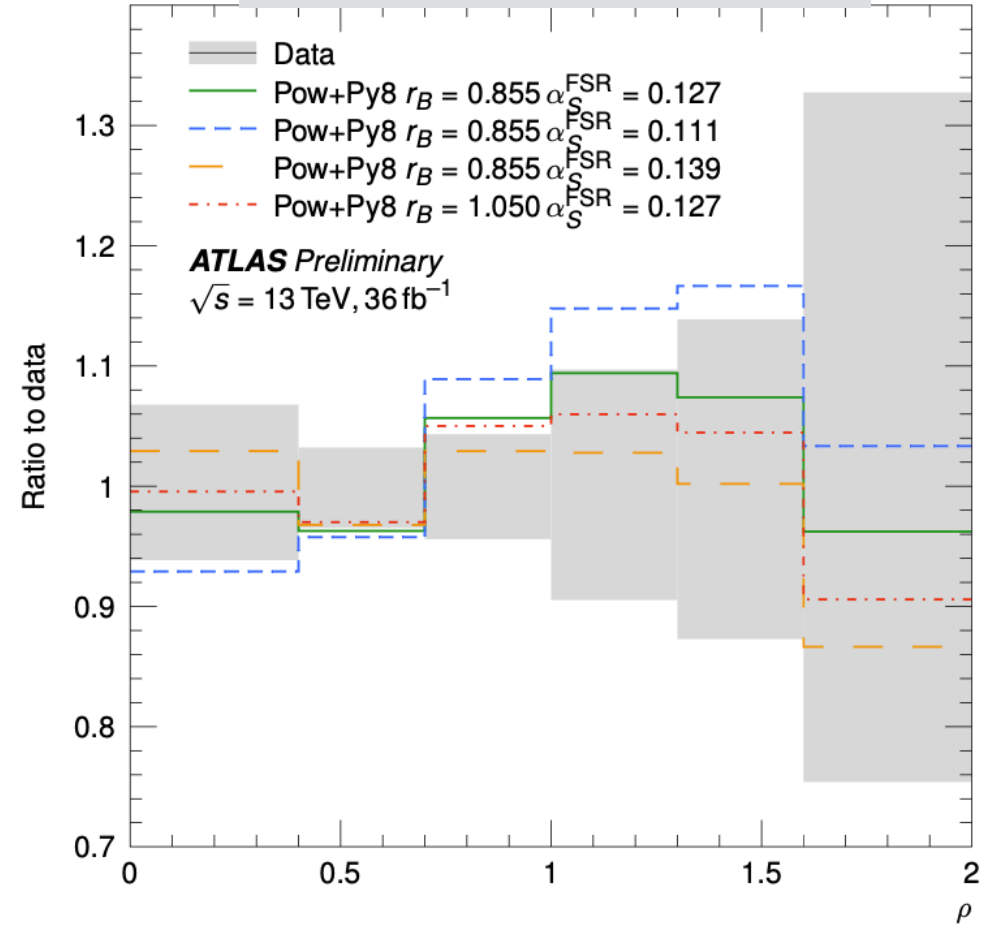
fiducial data/MC comparisons II



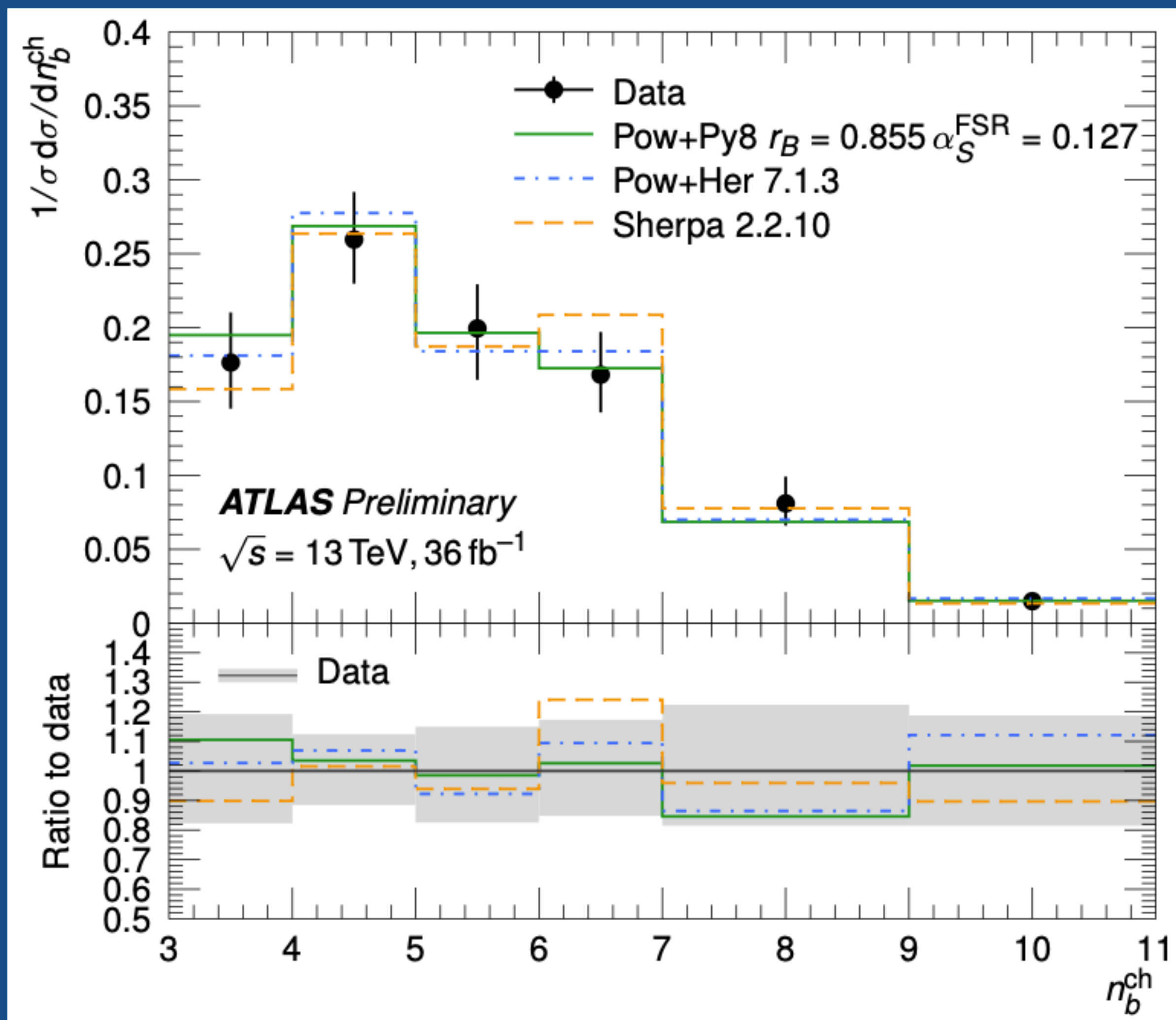
$\sim 1\sigma$ tensions
with recent generators
→ further investigation.

large impact from
A14 α_S^{FSR} variations

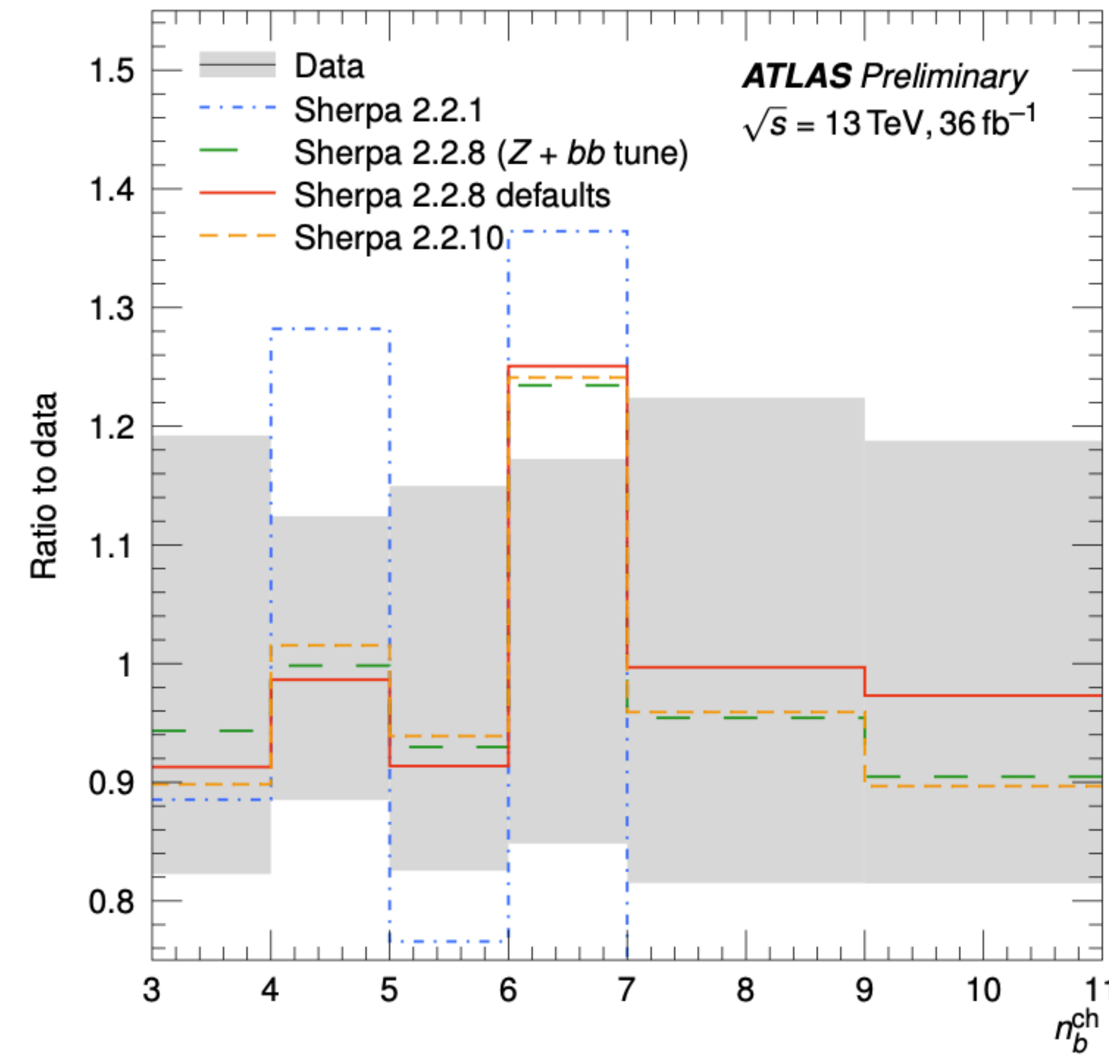
no sensitivity to α_S^{ISR}



fiducial data/MC comparisons III



poor choice of heavy baryon enhancement in ATLAS's Sherpa 2.2.1 setup has now been fixed.

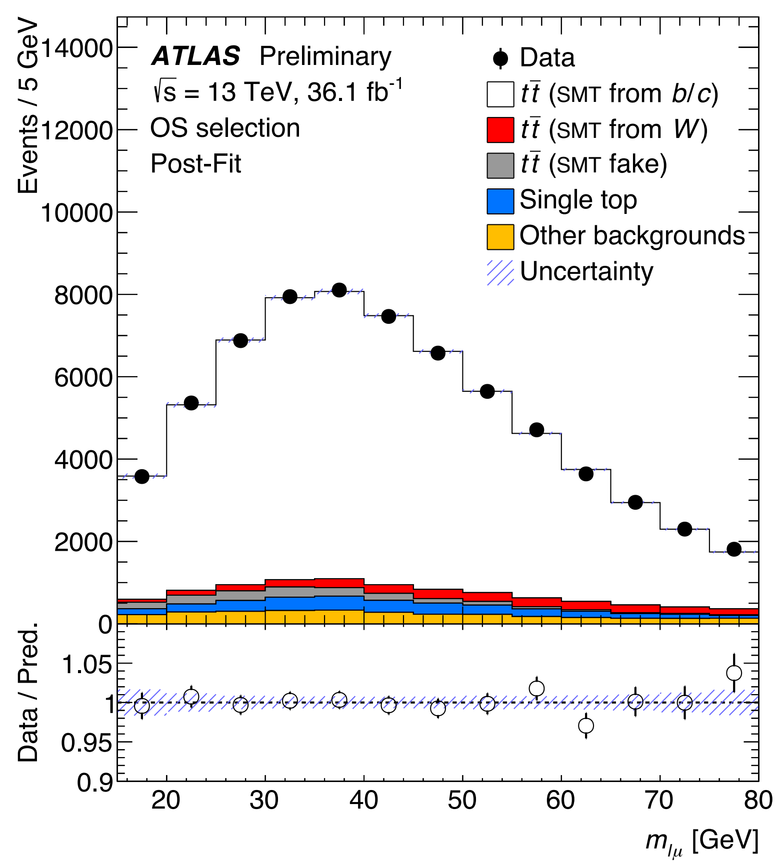


Py8/Her7 + EvtGen and Sherpa 2.2.10 all perform admirably.

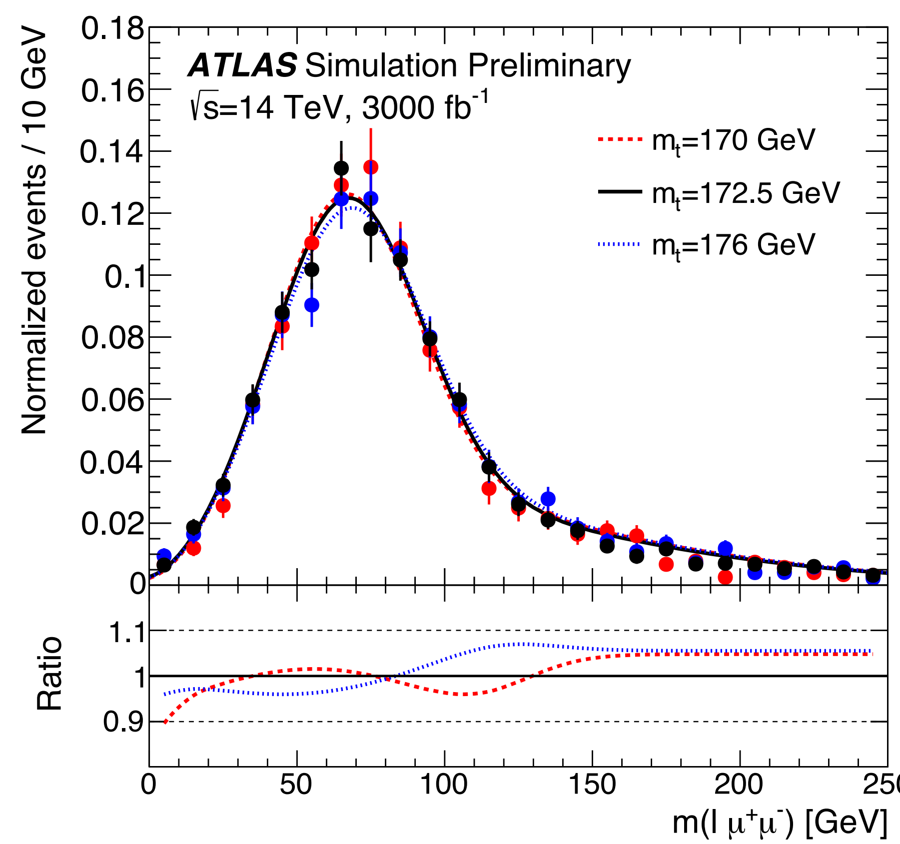
impact: top mass

at our current precision, **generators *well-tuned* to $e^+e^- \to bb$ data yield good predictions in $t\bar{t}$ events and top-quark decays**.

- we now have unfolded data to guide our choice of related uncertainties.
- precise b-fragmentation → smaller uncertainties on the b-quark to b-jet transfer
- a better understanding of $t \to b^{\text{quark}} \to b^{\text{hadron}}$ crucial for fully-leptonic template mass extractions ($m_{\ell\mu}$ or $m_{\ell,J/\psi}$)
    - joint fragmentation+mass measurements in leptonic modes **should be investigated**



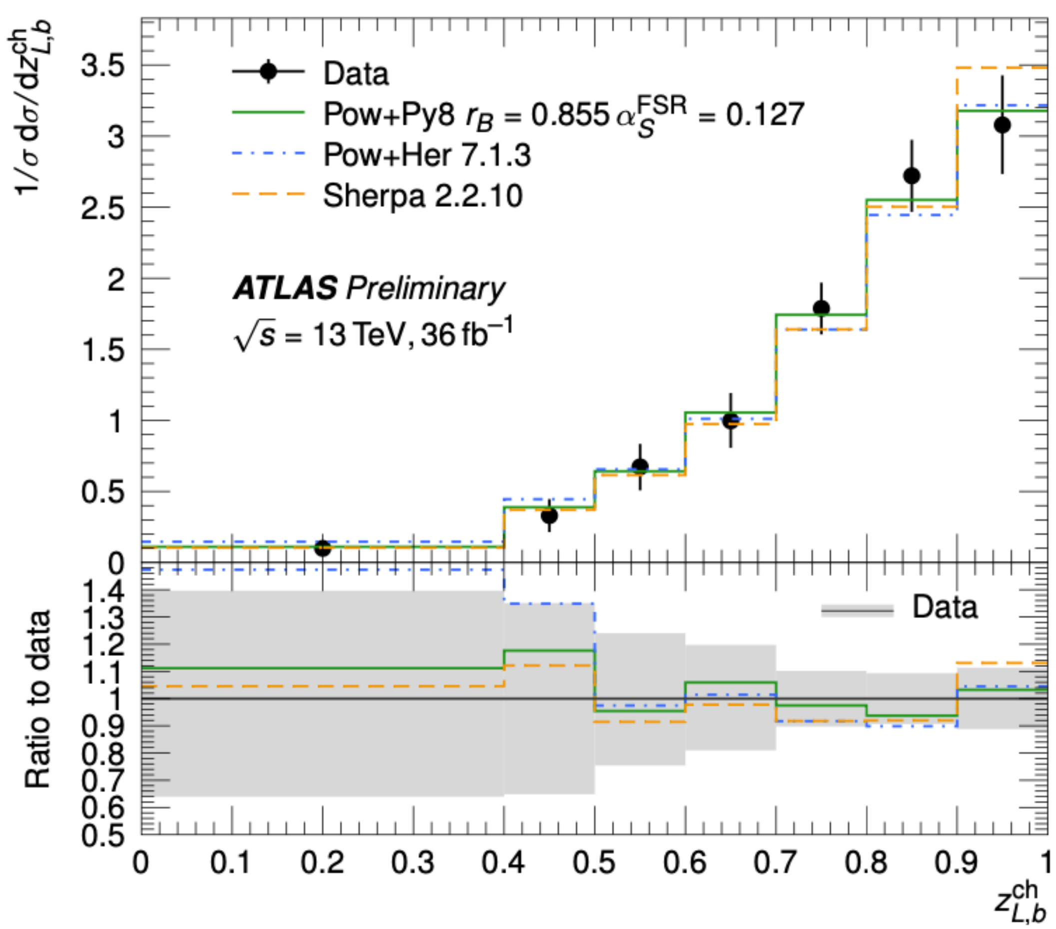
ATLAS-CONF-2019-046

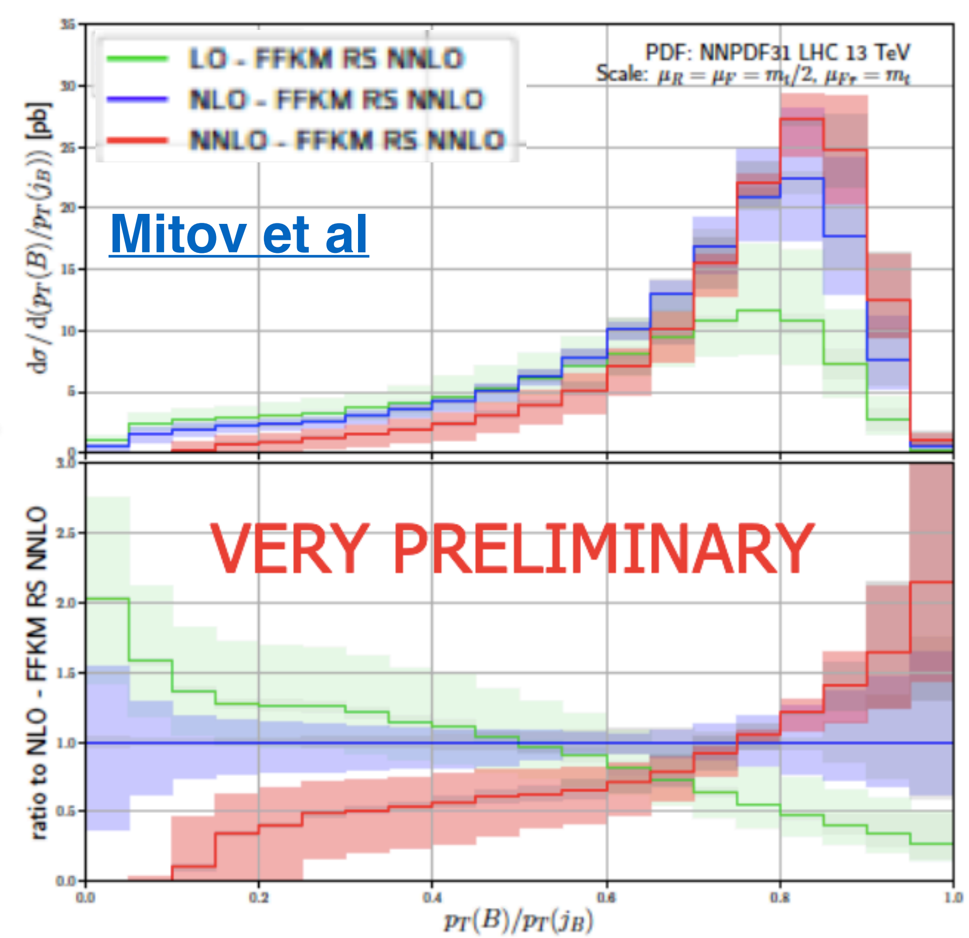


ATLAS-CONF-2018-042

status of ~analytic calculations

- Mitov et al. showed the first quasi-analytic calculations of similar distributions at TOP2020.
- will be very interesting to compare to unfolded data
- challenges for comparison:
  - calculations are for production of exclusive B-hadron species
  - jets are defined at parton-level in calculation
  - ATLAS measurement is to stable, charged particles





summary and outlook

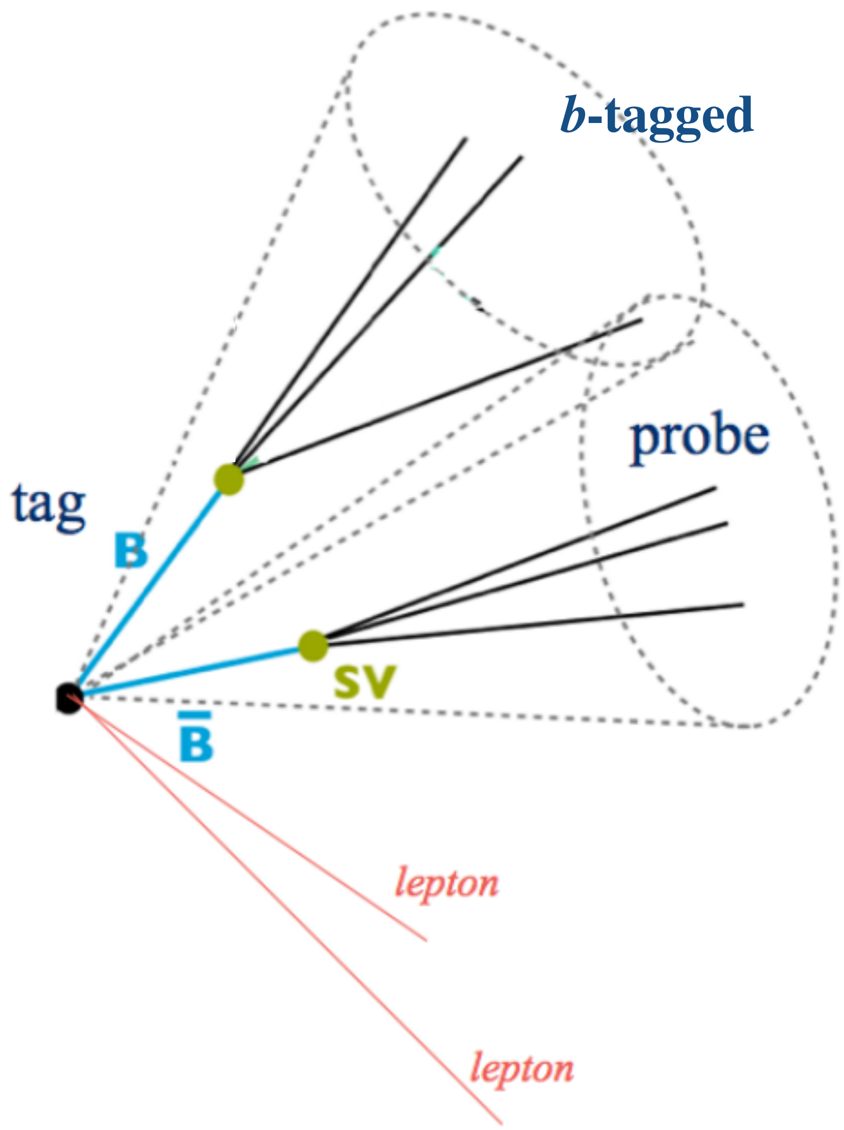
- several differential observables sensitive to b-quark fragmentation have been measured in $tt \rightarrow e\mu bb$ events with the 2015+2016 ATLAS data
- comparisons to modern MC generator predictions are made
- we will provide Rivet routine + HEPData entries for the unfolded data.
- many improvements possible: more precise measurements may be necessary for best “direct” top-quark mass extraction
- first ~analytic calculations of b-quark fragmentation in top-quark decays are on the horizon for comparison!

bonus

overview: detector-level event selection

obtain sample of b-jets from $t\bar{t} \rightarrow e\mu\nu\nu bb$ events

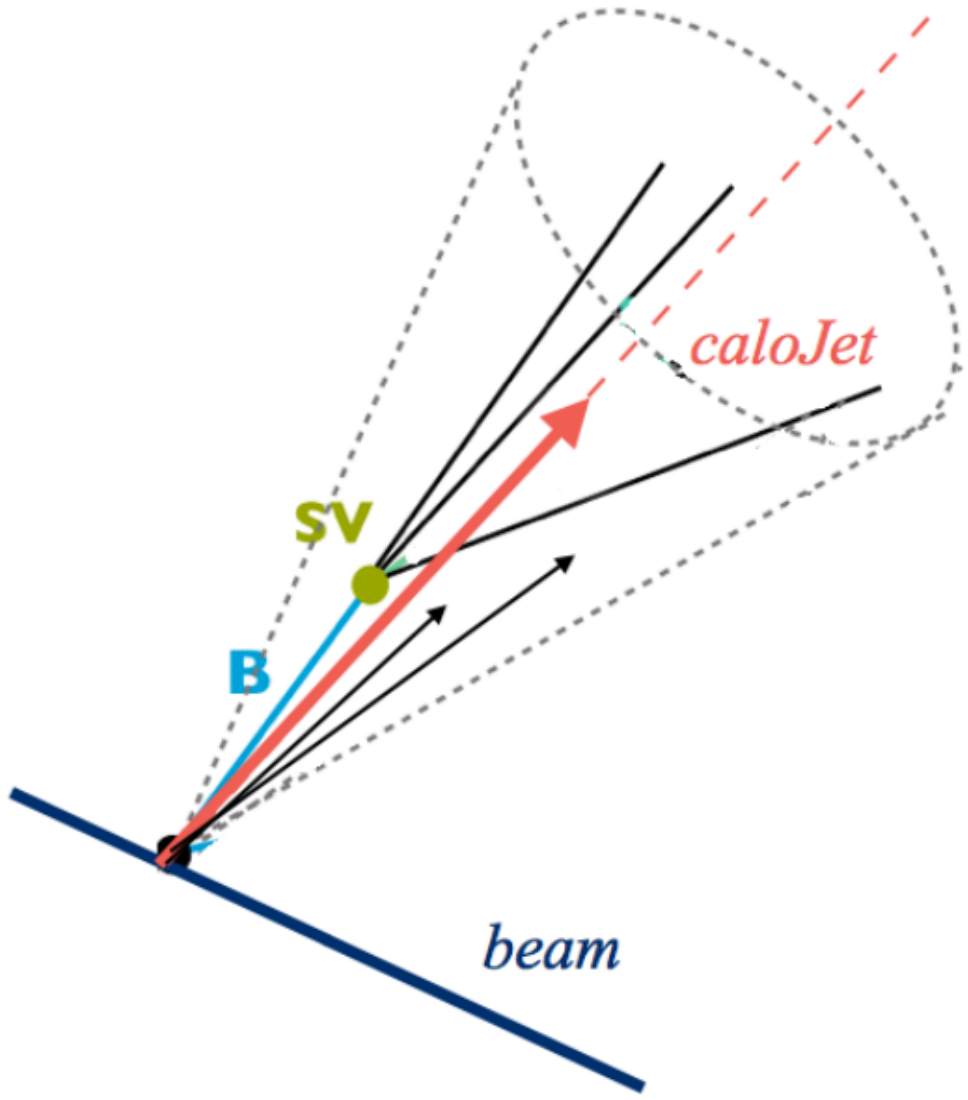
- require ==1 e and ==1 μ with opposite charge
- ==2 anti-k_t $R = 0.4$ calorimeter jets
  - $|\eta| < 2.5$, $p_T > 30$ GeV
  - at least **one jet must be *b*-tagged** ("the tag")
  - we measure the other jet ("the probe")
  - if both jets are b-tagged, **both jets are measured**



b-hadron reconstruction

identify b-hadron candidates using the ATLAS single secondary vertex finder (SSVF)

- we require a reconstructed SSVF vertex with ≥ 3 tracks for "probe" jets
  - **95+% b-jet purity after this cut**
- the SSVF vertex **only has tracking information**
- track constituents of the SV make up the b-hadron candidate.
- tracks ghost-associated to the jet make up the jet charged component.



detector-level yields: predicted and observed

	events with $e\mu jj$ (≥ 1 b-tag)		probe-jets	
process	predicted yields			
fiducial $t\bar{t}$	–		44000	± 12000
non-fiducial $t\bar{t}$	–		6700	± 1900
total $t\bar{t}$	76000	± 18000	51000	± 12000
single top	4400	± 1500	1580	± 600
Z+jets	125	± 45	13.0	± 5.1
diboson	90	± 34	9.7	± 3.9
total non-$t\bar{t}$	4600	± 1600	1600	± 600
b-jets	–		52200	± 12000
c-jets	–		180	± 60
other jets	–		250	± 70
total prediction	81000	± 18000	53000	± 12000
	observed yields			
data	88511		57476	

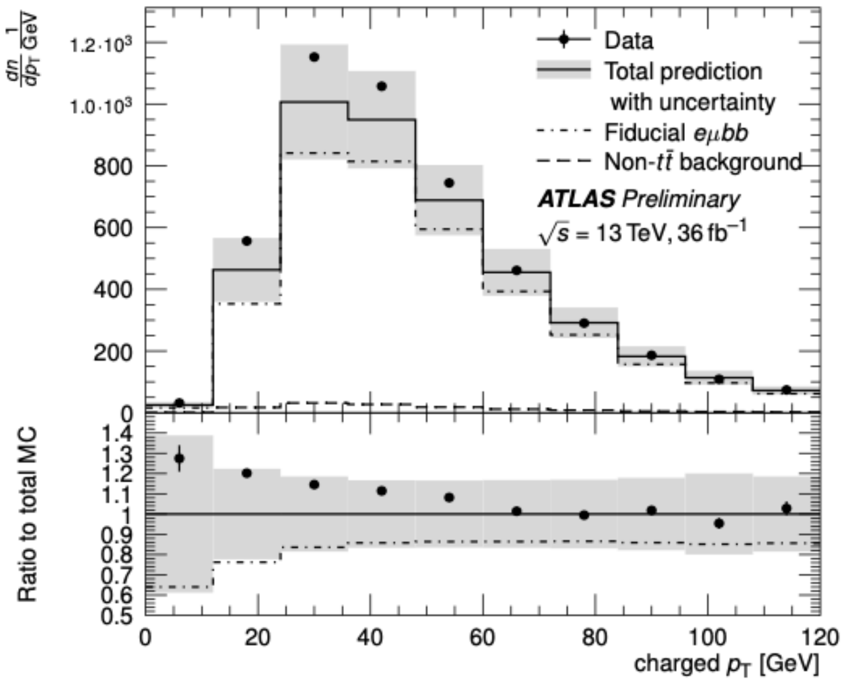
Table 1: Summary of expected and observed event and probe-jet yields.

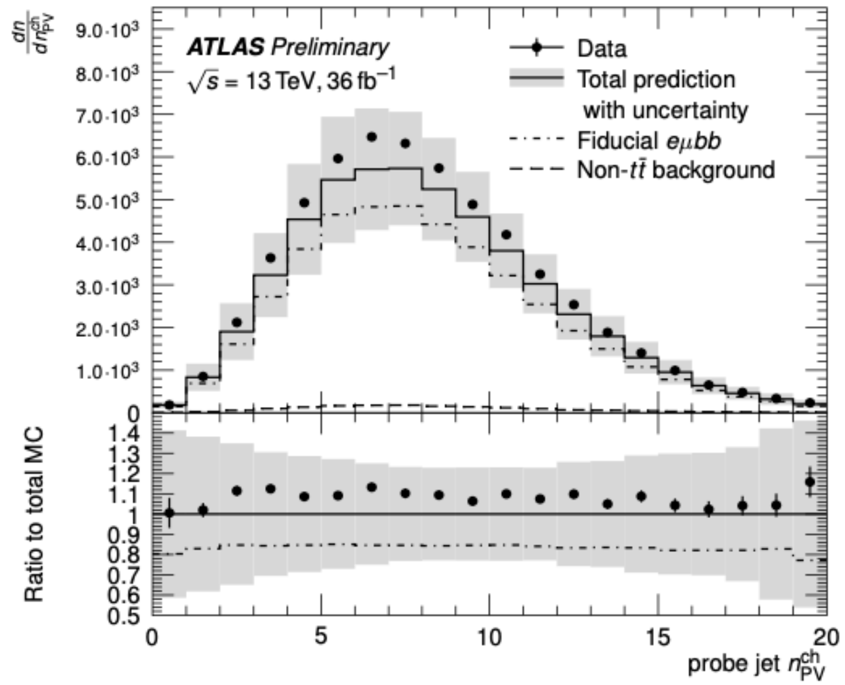
tt dominant

high expected *b*-jet purity

7.5% data / MC discrepancy

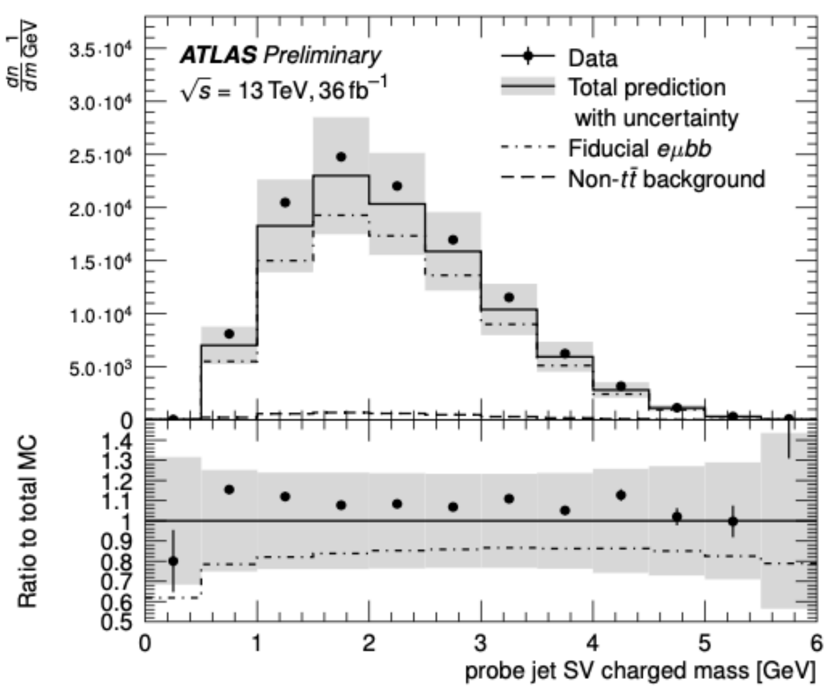
detector-level control plots

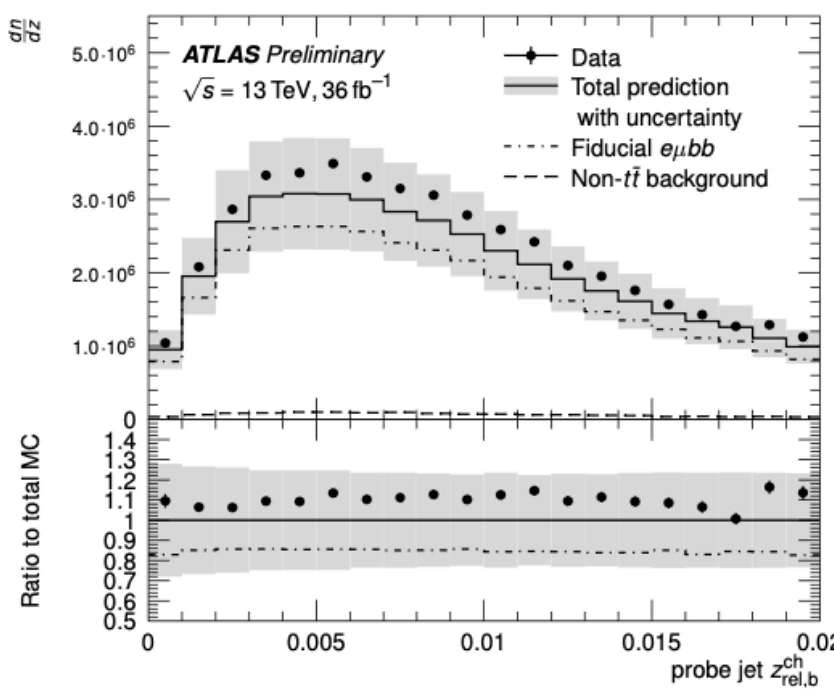




charged component of jet p_T has similar slope to jet p_T (data vs MC taken as uncertainty)

other relevant PV and SV quantities are well-modeled





SV mass and relative momentum transverse to to jet axis,

$$z^{ch}_{rel,b} = (\vec{p}_b^{\,ch} \times \vec{p}_{jet}^{\,ch}) / |\vec{p}_{jet}^{\,ch}|^2$$

are also shown.

unfolding (FBU)

- **based on the "Fully Bayesian Unfolding" technique (arxiv: 1201.4612)**
- we use the following likelihood to unfold the detector effects:

$$\mathcal{L}(d|\sigma,\Lambda) = \prod_{i\,\in\,\mathrm{recobins}} \mathrm{Poiss}(d_i|x_i(\sigma,\Lambda))$$

$$x_i(\sigma,\Lambda) = L(\Lambda)\times(b_i(\Lambda)+M_{ij}(\Lambda)\,\sigma_j)$$

- where $\overrightarrow{x}$ is the total number of predicted events in detector-level bins, L is the luminosity, $\overrightarrow{b}$ are predicted background events, M_{ij} is probability that a particle-level jet in bin i is reconstructed in detector-level bin j, *and* Λ are the systematic nuisance parameters.
- **we marginalize over nuisance parameters through MCMC sampling to build the marginal posteriors on $\overrightarrow{\sigma}$, the parameters-of-interest.**

unfolding procedure: background and response matrices

- to calculate the background predictions and response matrices at a given point in parameter space, we take linear combinations of the background predictions for nominal and the 1-sigma variation corresponding to a systematic uncertainty:

$$M_{ij}^{\mathrm{p}\to\mathrm{d}}(\vec{p}) = M_{ij,0}^{\mathrm{p}\to\mathrm{d}} + \sum_{k\in\mathrm{systematics}} p_k * (M_{ij,k}^{\mathrm{p}\to\mathrm{d}} - M_{ij,0}^{\mathrm{p}\to\mathrm{d}})$$

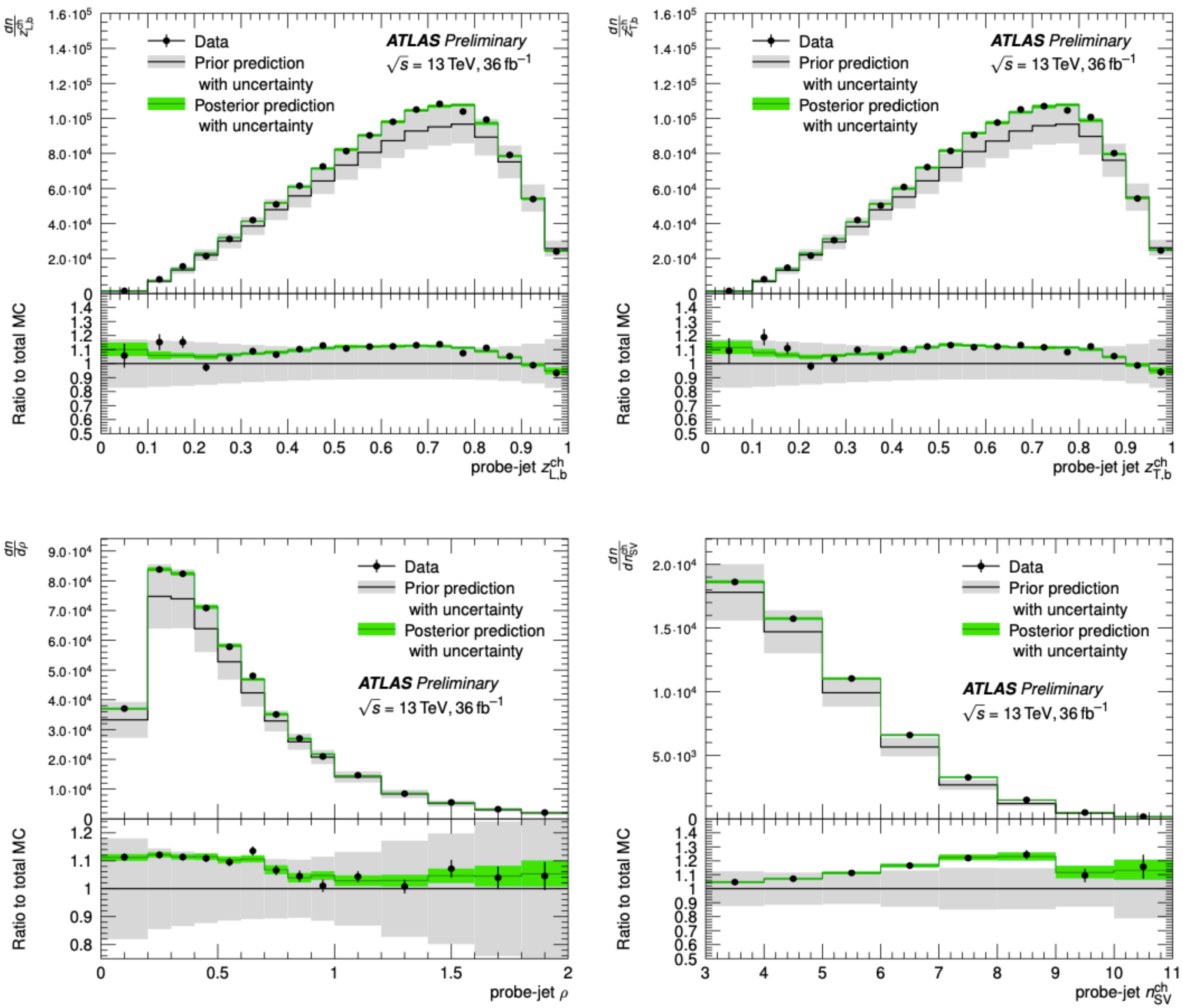
$$b_i(\vec{p}) = b_{i,0} + \sum_{k\in\mathrm{systematics}} p_k * (b_{i,k} - b_{i,0})$$

where the $\overrightarrow{p}$ are the nuisance parameters at a particular point in parameter space.

- in this way we can coherently and smoothly vary the response matrices and background predictions corresponding to systematic variations.

- we're wandering around the space of possible response matrices and background predictions available given our understanding of the predictions.

detector-level data/MC comparison before/after unfolding



detector-level observables
before (prior) and
after (posterior) unfolding

unfolding model is capable of describing observed data

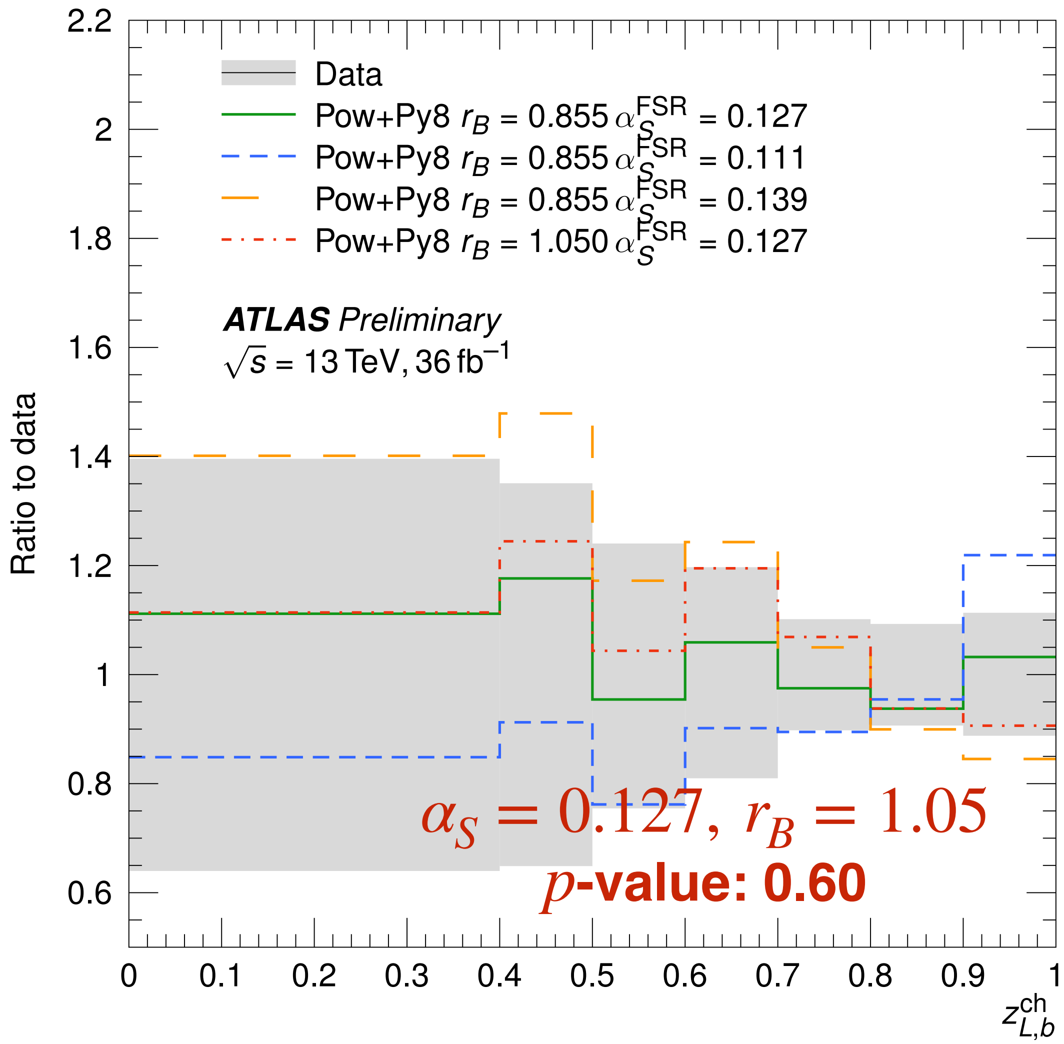
parameter anti-correlations cause reduction in total uncertainty

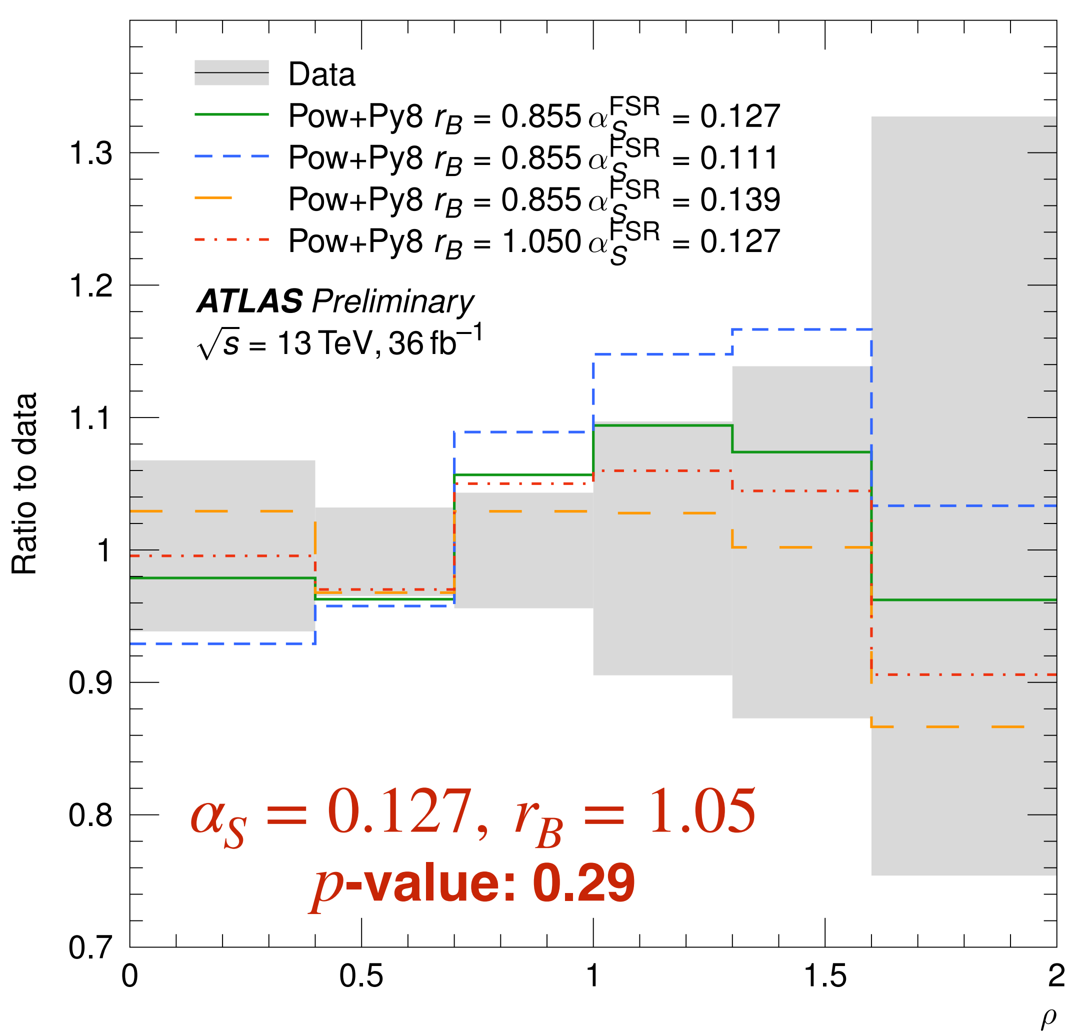
posterior uncertainty closely follows data stat uncertainty

impact: top mass

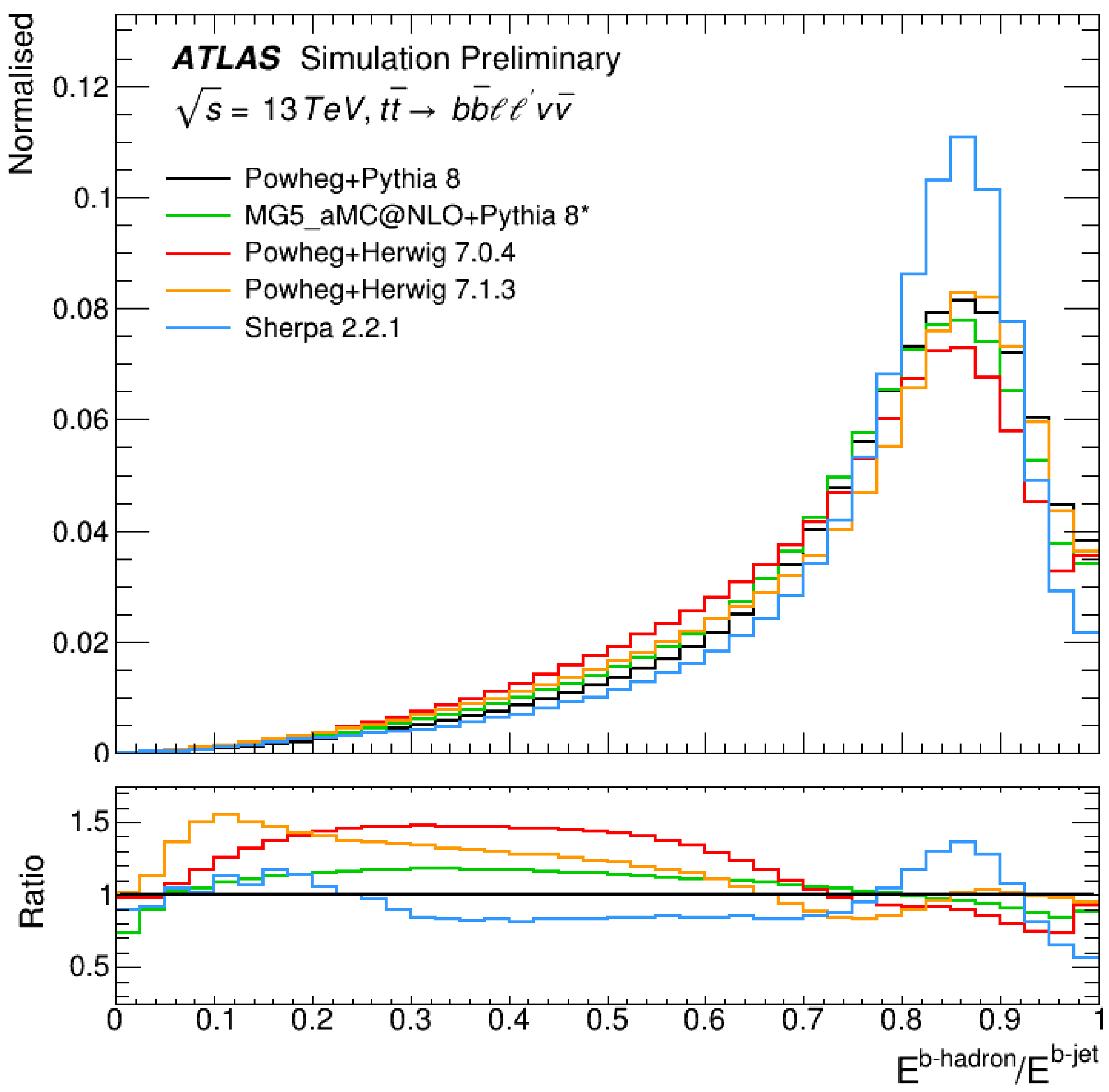
ATLAS retuned the Lund-Bowler r_B parameter to LEP data with the A14 α_S

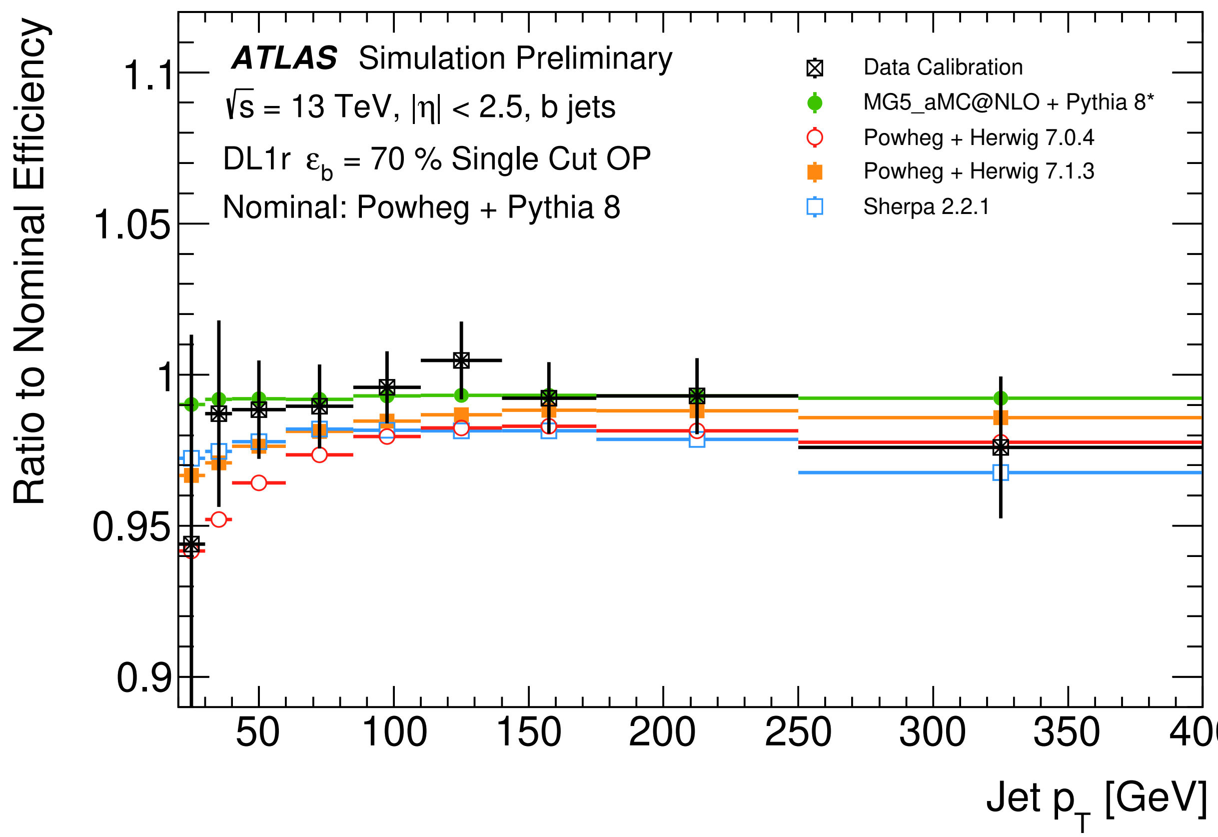
- used in preliminary mass measurement ATLAS-CONF-2020-046
- this combination of $\alpha_S = 0.127,\ r_B = 1.05$ agrees with observed data to current precision. **(dash-dotted red curve)**





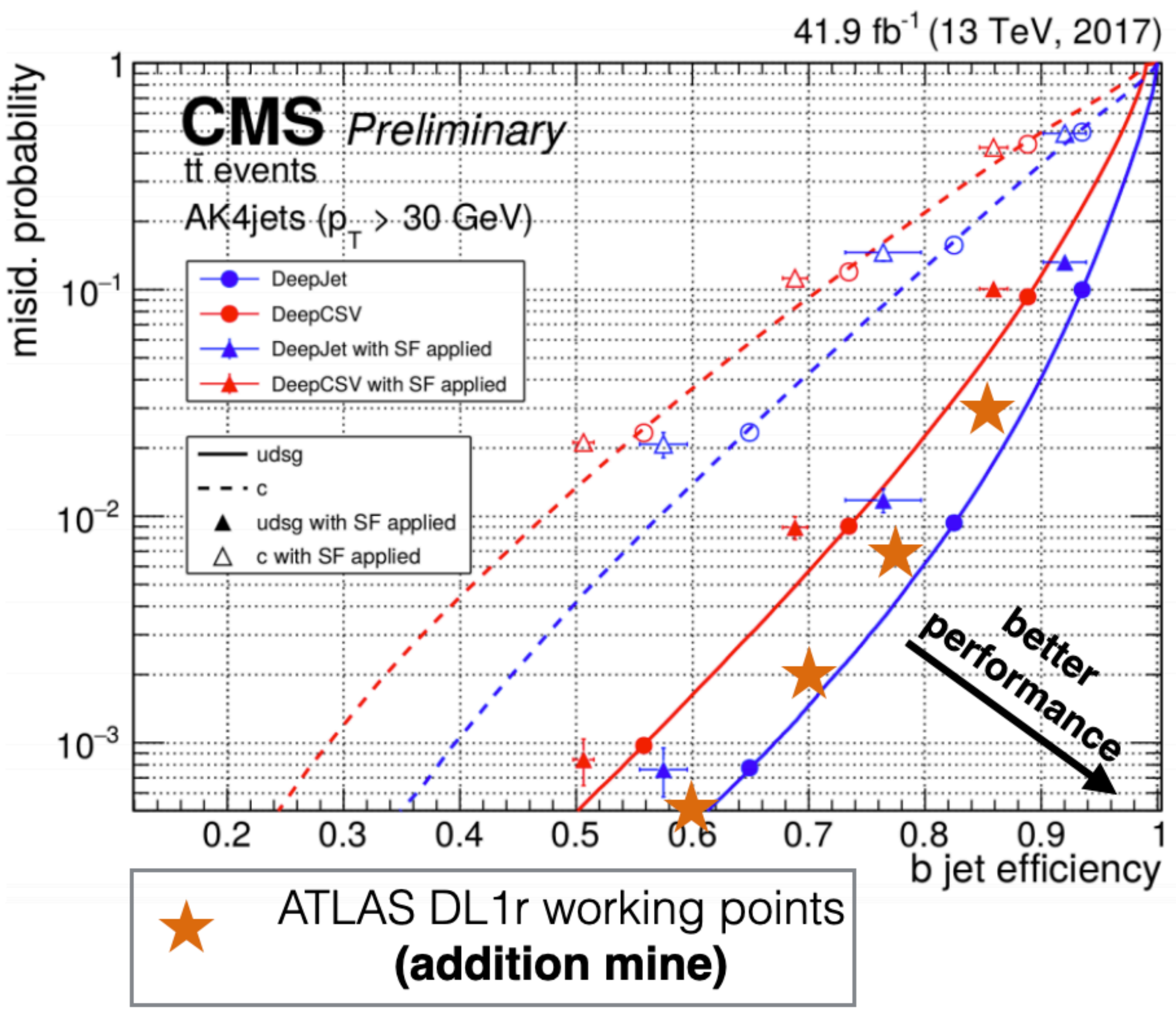
impact on b-tagging

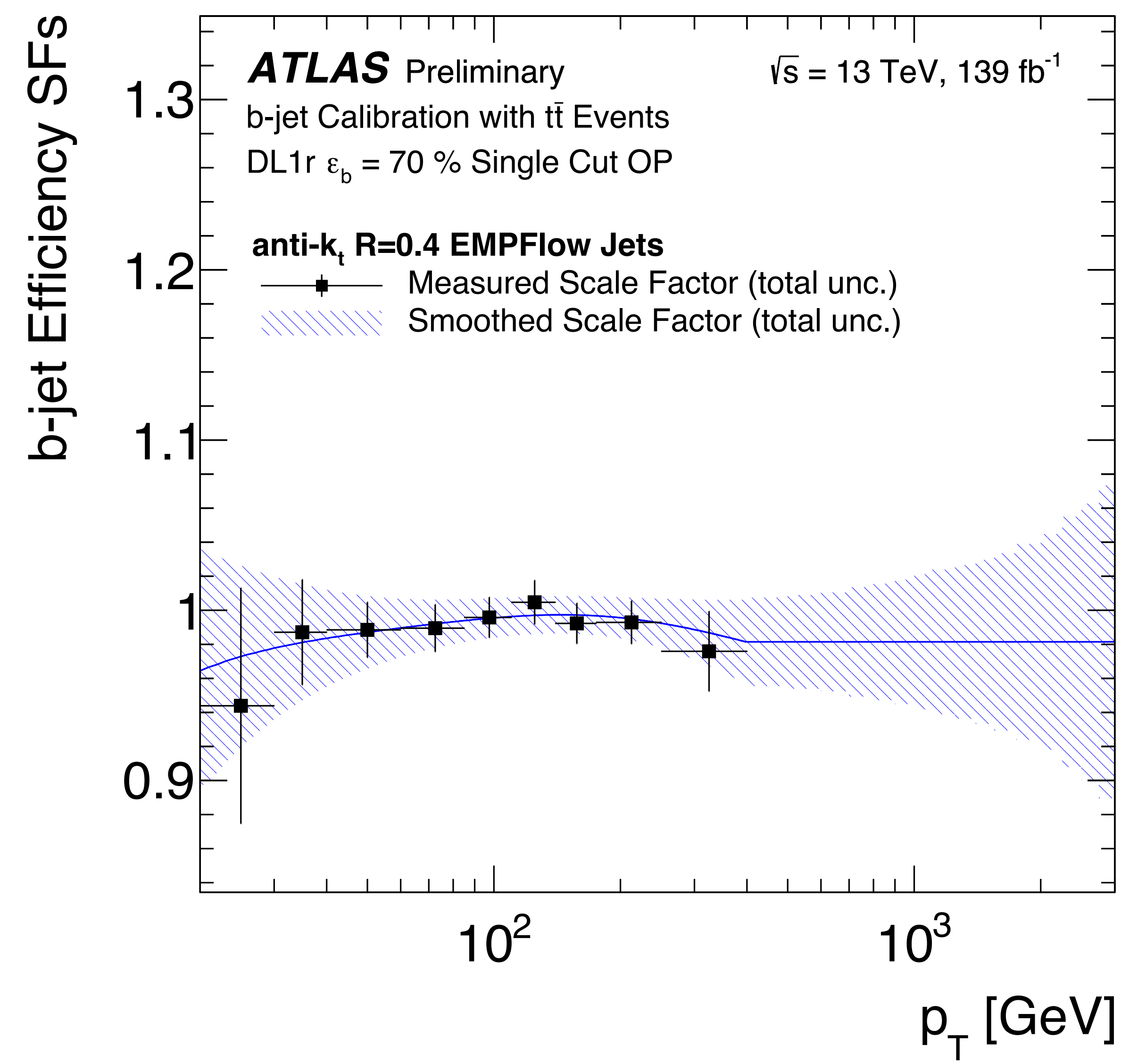




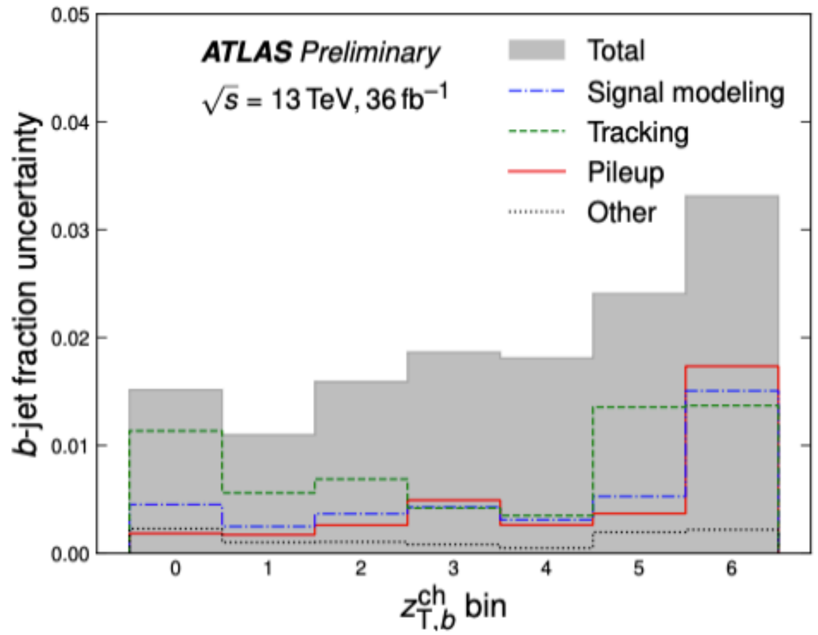
ATLAS-PHYS-PUB-2020-009

b-tagging SFs

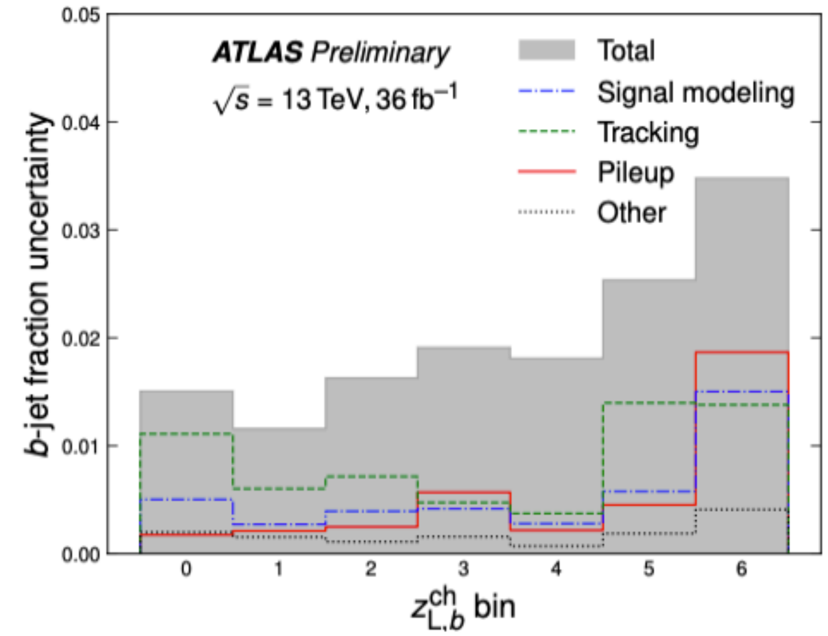




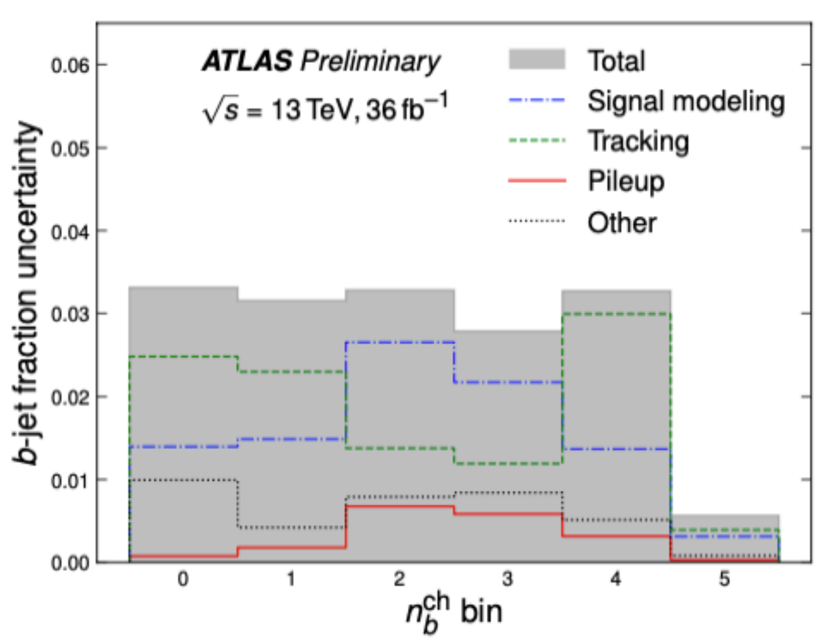
uncertainty breakdowns



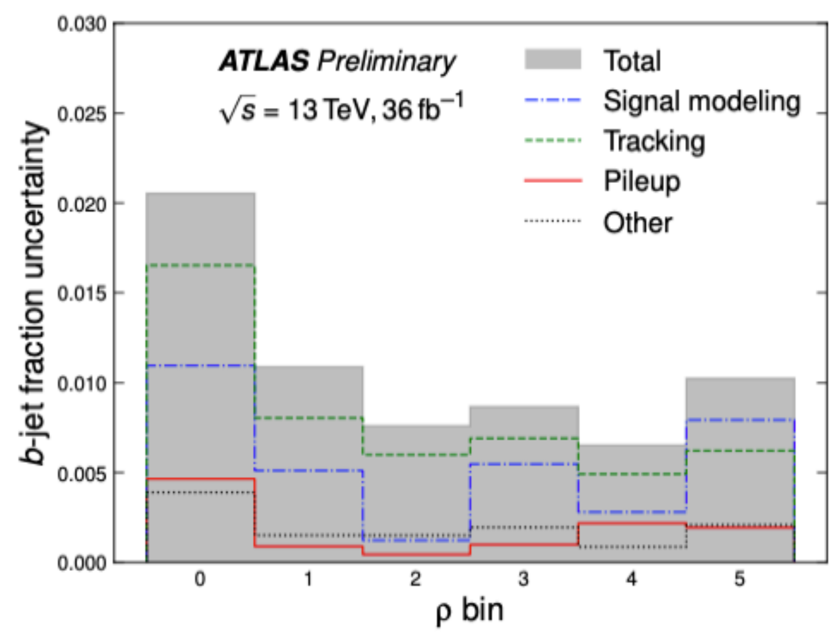
(a) $z^{ch}_{T,b}$



(b) $z^{ch}_{L,b}$



(c) n^{ch}_{b}



(d) ρ

p-values

generator configuration	$z^{\mathrm{ch}}_{\mathrm{T},b}$	$z^{\mathrm{ch}}_{\mathrm{L},b}$	n^{ch}_{b}	ρ
Powheg+Pythia8 A14 $r_B = 0.855$ $\alpha_S^{\mathrm{FSR}} = 0.127$	0.50	0.98	0.94	0.20
Powheg+Pythia8 A14 $r_B = 0.855$ $\alpha_S^{\mathrm{FSR}} = 0.139$	0.13	0.37	0.95	0.31
Powheg+Pythia8 A14 $r_B = 0.855$ $\alpha_S^{\mathrm{FSR}} = 0.111$	0.16	0.62	0.95	0.07
Powheg+Pythia8 A14 $r_B = 1.050$ $\alpha_S^{\mathrm{FSR}} = 0.127$	0.18	0.60	0.94	0.29
Powheg+Herwig7.0.4	0.22	0.50	0.95	0.12
Powheg+Herwig7.1.3	0.42	0.84	0.92	0.18
Sherpa2.2.1	0.14	0.45	0.01	0.52
Sherpa2.2.8 ($Z + bb$ tune)	0.004	0.02	0.45	0.24
Sherpa2.2.8	0.32	0.79	0.49	0.11
Sherpa2.2.10	0.27	0.84	0.32	0.08

Table 3: Summary of p-values for various MC generator configurations based on the observed data distributions