

# Combined Higgs boson measurements at the ATLAS experiment

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# Overview

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- ATLAS analyses target different Higgs boson production modes and decay channels
  - Many of measured quantities are **independent** of the Higgs boson production/decay
- Statistical combination of the most current measurements to get **best precision** and **limits** on important parameters in the Standard Model (SM)
  - **Updated** to full Run 2 dataset for some of the input analysis
  - Previous publication for 2015-2017 dataset [arXiv:1909.02845](https://arxiv.org/abs/1909.02845)

## Focus Today

Couplings - ([ATLAS-CONF-2020-027](#))

- Global signal strength
- Production mode cross sections (XS) and branching ratios (BR)
- Simplified template cross section (STXS) results
- Interpretations
  - $\kappa$ -framework
  - EFT([ATLAS-CONF-2020-053](#))

Differential - ([ATLAS-CONF-2019-032](#))

- $p_{T,H}$  combination for  $H \rightarrow ZZ^* \rightarrow 4\ell$  and  $H \rightarrow \gamma\gamma$

# Input to combined measurements

- So far  $H \rightarrow ZZ^* \rightarrow 4\ell$ ,  $H \rightarrow \gamma\gamma$ ,  $VH \rightarrow bb$  analyses used for STXS combination
- $H \rightarrow \mu\mu$  and  $H \rightarrow inv$  used in a subset of the  $\kappa$ -framework results

Analysis Decay Channel	Target Prod. Modes	$\mathcal{L}$ [fb <sup>-1</sup> ]	Link	Signal Strength	Prod. Modes Branching Ratios	STXS	$\kappa$ -interp	EFT
$H \rightarrow \gamma\gamma$	ggF,VBF,WH,ZH,t $\bar{t}$ H,tH	139	<a href="#">ATLAS-CONF-2020-026</a>	✓	✓	✓	✓	✓
$H \rightarrow ZZ^*$	ggF,VBF,WH,ZH,t $\bar{t}$ H(4 $\ell$ )	139	<a href="#">arXiv:2004.03447</a>	✓	✓	✓	✓	✓
	t $\bar{t}$ H excl. $H \rightarrow ZZ^* \rightarrow 4\ell$	36.1	<a href="#">arXiv:1712.08891</a> , <a href="#">arXiv:1806.00425</a>	✓	✓		✓	
$H \rightarrow WW^*$	ggF,VBF	36.1	<a href="#">arXiv:1808.09054</a>	✓	✓		✓	
	t $\bar{t}$ H		<a href="#">arXiv:1712.08891</a> , <a href="#">arXiv:1806.00425</a>					
$H \rightarrow \tau\tau$	ggF,VBF	36.1	<a href="#">arXiv:1811.08856</a>	✓	✓		✓	
	t $\bar{t}$ H		<a href="#">arXiv:1712.08891</a> , <a href="#">arXiv:1806.00425</a>					
$H \rightarrow b\bar{b}$	VBF	24.5 – 30.6	<a href="#">arXiv:1807.08639</a>	✓	✓		✓	
	WH,ZH	139	<a href="#">arXiv:2007.02873</a>	✓	✓	✓	✓	✓
	t $\bar{t}$ H	36.1	<a href="#">arXiv:1712.08895</a> , <a href="#">arXiv:1806.00425</a>	✓	✓		✓	
$H \rightarrow \mu\mu$	ggF,VBF,VH,t $\bar{t}$ H,tH	139	<a href="#">arXiv:2007.07830</a>				✓	
$H \rightarrow inv$	VBF	139	<a href="#">ATLAS-CONF-2020-008</a>				✓	

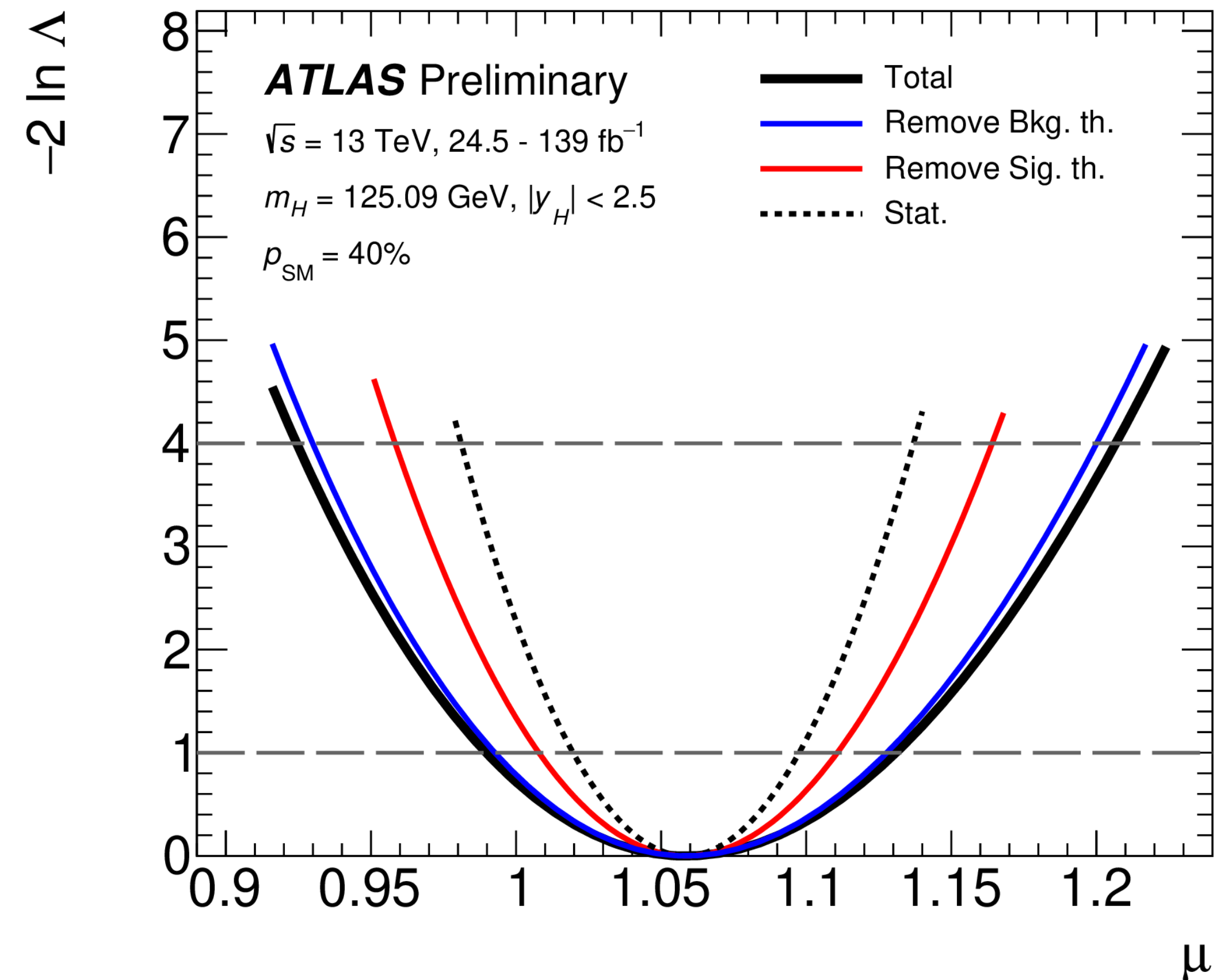
- Results of combination obtained from product of likelihood functions
- Systematics from same sources are correlated

# Global Signal Strength

- Simultaneous measurement of parameter  $\mu$  (where  $\mu_{if} = (\sigma_i \times \text{BR}_f)/(\sigma_i \times \text{BR}_f)_{\text{SM}}$ )

$$\mu = 1.06 \pm 0.07 = 1.06 \pm 0.04 \text{ (stat.)} \pm 0.03 \text{ (exp.)} {}^{+0.05}_{-0.04} \text{ (sig. th.)} \pm 0.02 \text{ (bkg. th.)}$$

- Systematic uncertainties **dominate** due to smaller statistical uncertainty from full Run-2 dataset
- Largest theory uncertainties are included
  - QCD scale ( $\sim 4\%$  on ggF)
  - PDF and  $\alpha_s$  uncertainty ( $\sim 3\%$  on ggF)
  - Note that ggF is  $\sim 90\%$  of signal and is predicted at N3LO
- Consistent with SM ( $\mu = 1$ ) with a p-value of 40%

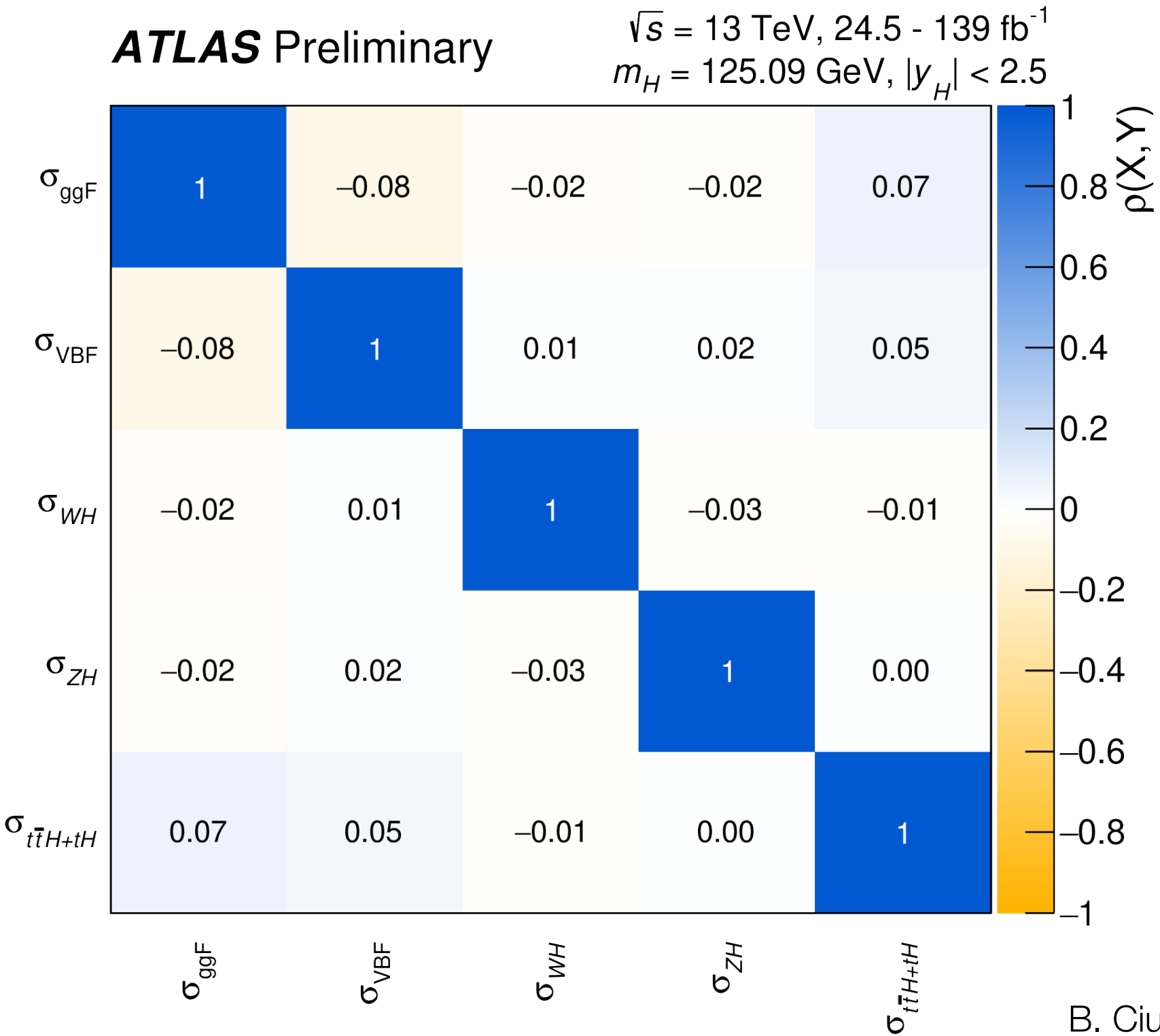
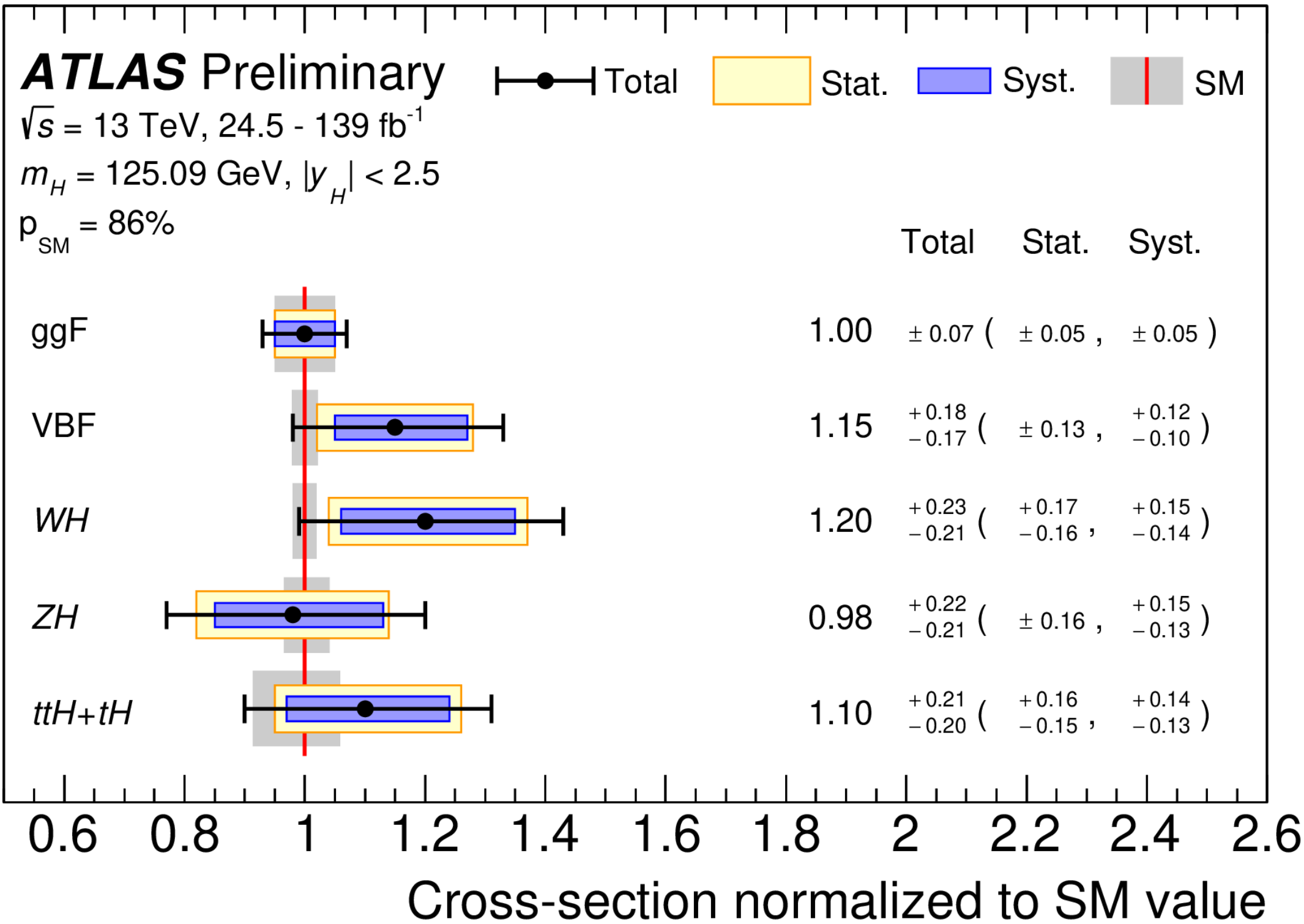




# Production Modes

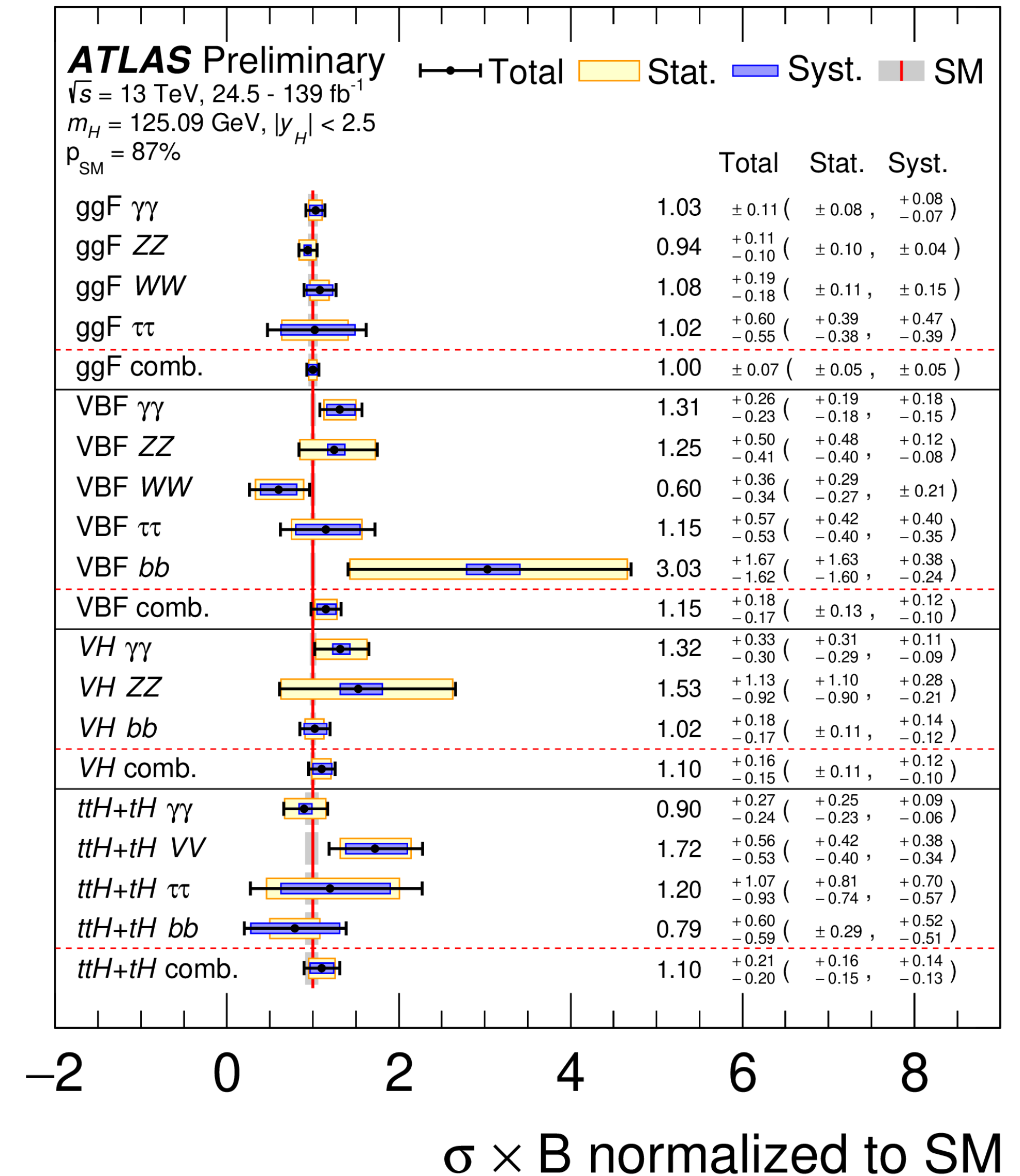
- Branching fractions fixed to SM expectations, cross sections are simultaneously determined in the fit to data
- Decrease in ggF and VBF correlation compared to previous publication ( $\sim 15\% \rightarrow 8\%$ )
  - Mainly from updated  $H \rightarrow ZZ^* \rightarrow 4\ell$  and  $H \rightarrow \gamma\gamma$  results
- **Above  $5\sigma$  observed** for all production modes
  - First **observation** of  $WH$  with observed (expected) significances of  $6.3\sigma$  ( $5.2\sigma$ )

Systematic uncertainty same/approaching statistical uncertainty



# Production and Decay

- Can also test the SM by measuring  $(\sigma \times B)_{\text{if}}$ 
  - ggF driven by  $H \rightarrow ZZ^* \rightarrow 4\ell$  and  $H \rightarrow \gamma\gamma$
  - VBF driven by  $H \rightarrow WW$ ,  $H \rightarrow \gamma\gamma$ , and  $H \rightarrow \tau\tau$
  - VH driven by  $H \rightarrow b\bar{b}$
- Statistically limited channels are fixed to the SM
  - ggF  $\rightarrow H \rightarrow b\bar{b}$
  - VH,  $H \rightarrow WW/\tau\tau$
- Relative fraction of ZH and WH fixed to the SM
- **Good agreement** with the SM (p-value of 87%)



# STXS Scheme

- STXS scheme has bins of *exclusive phase space*

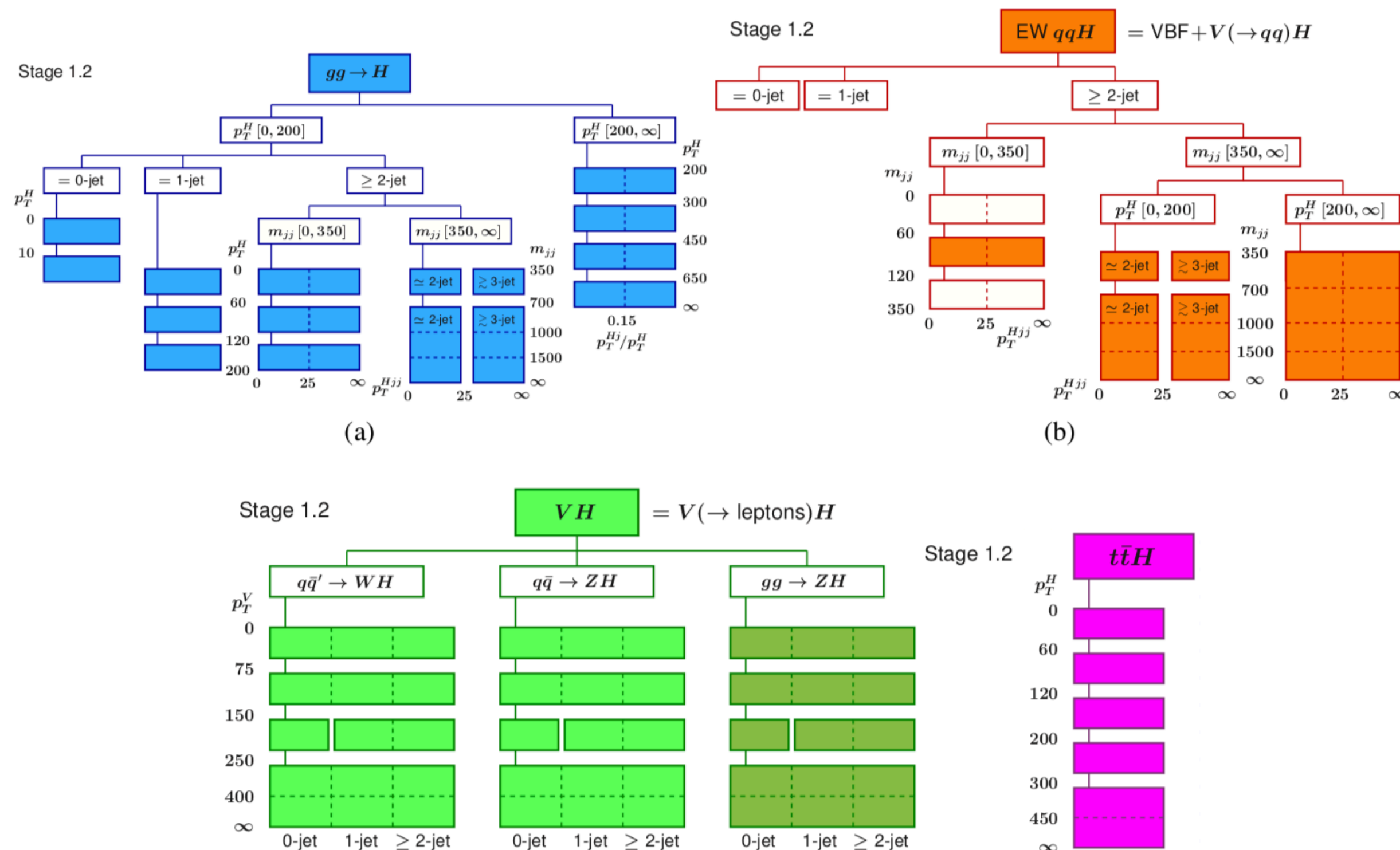
## STXS Motivation

- Reduce large theory uncertainties
- Bins are sensitive to BSM contributions
- Easier combination of different final states
- Full scheme measurement is statistically limited → merge bins

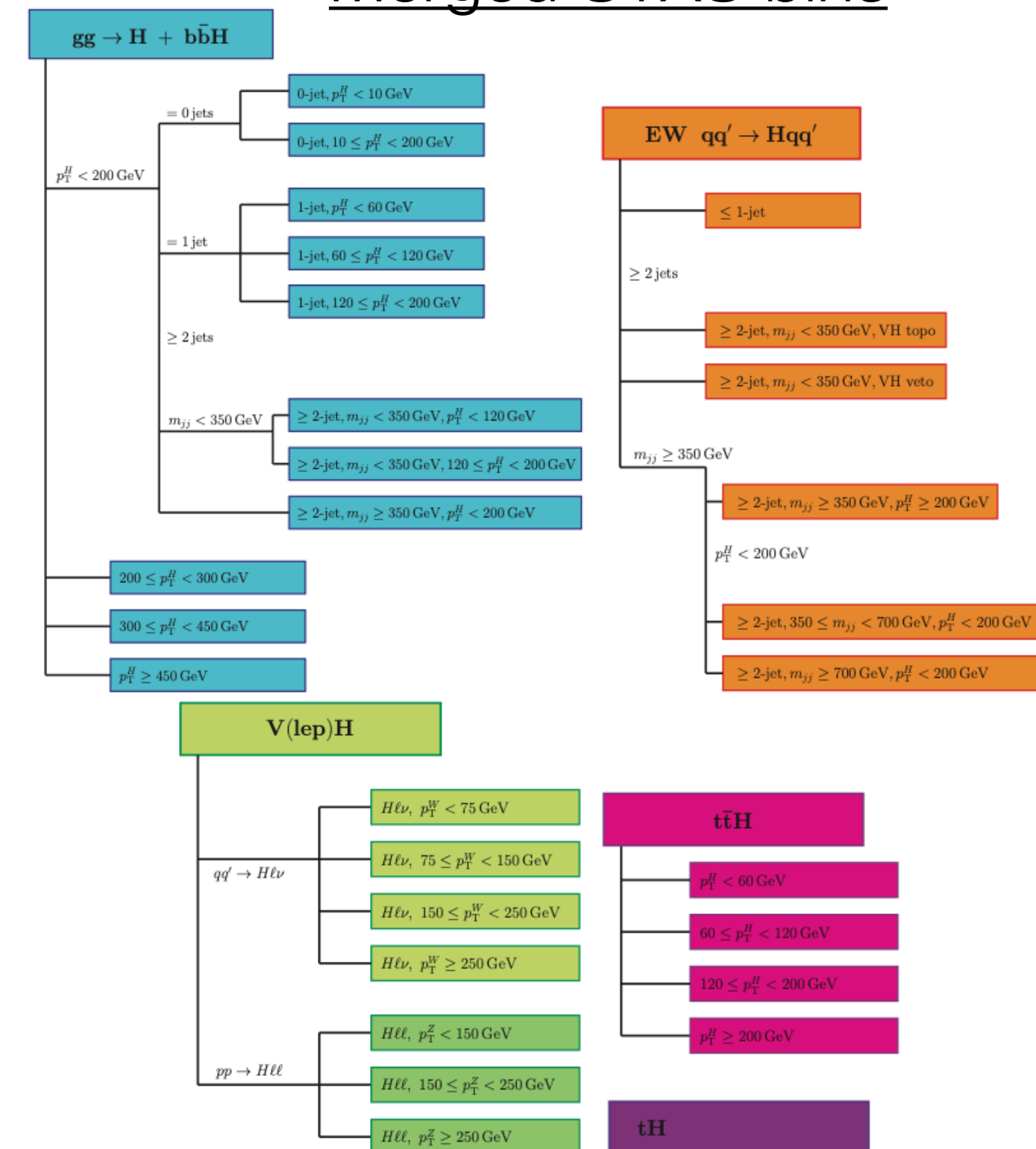
## Considerations for bin merging

- Merge highly correlated STXS bins
- Keep uncertainties < 100% with the exception of bins sensitive to BSM physics

## All STXS bins



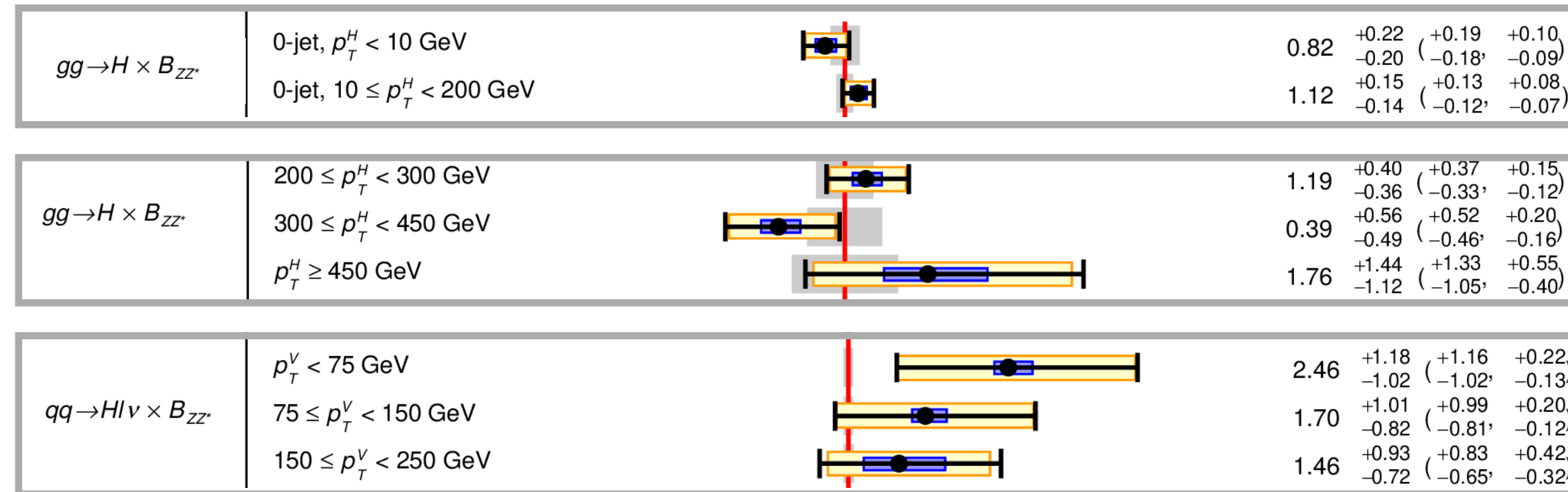
## Merged STXS bins



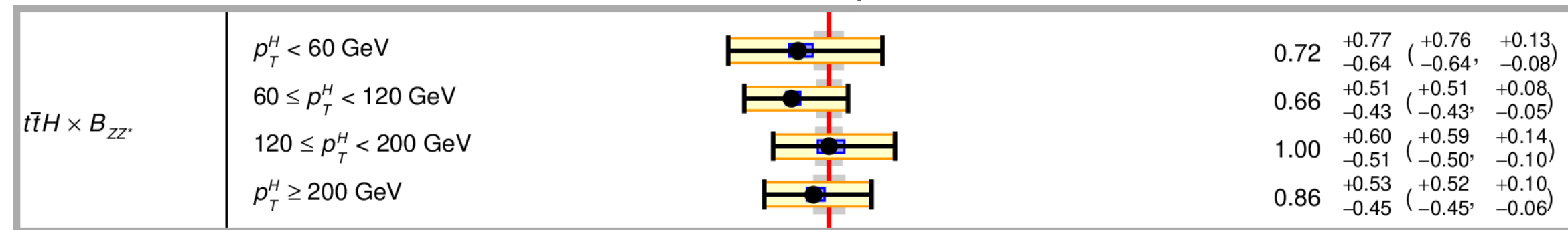


# STXS Results

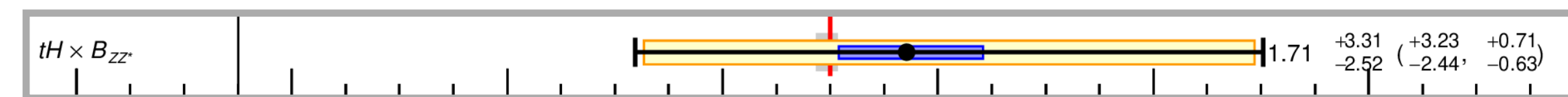
- Measure finer granularity of  $p_T$  bins



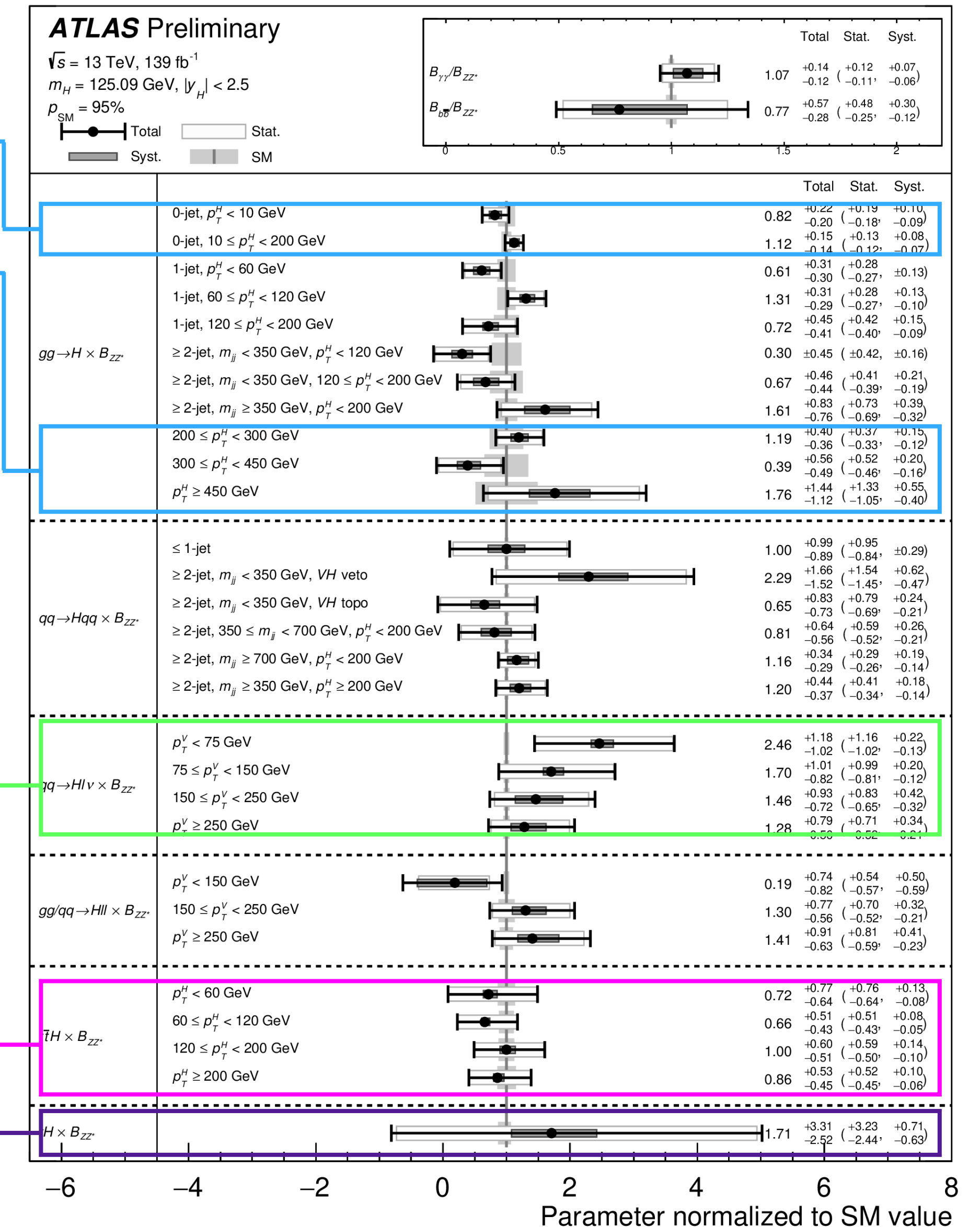
- Differential measurements for  $t\bar{t}H$  production



- First combined measurement of  $tH$



- Large correlation between VH and  $BR_{bb}$  (up to 83% in highest  $p_{TV}$  bin)
- Sensitivity for these come from  $VH \rightarrow bb$  and correlations will reduce when sensitivity from other production/decay channels improves
- Good agreement with the SM (p-value of 95%)

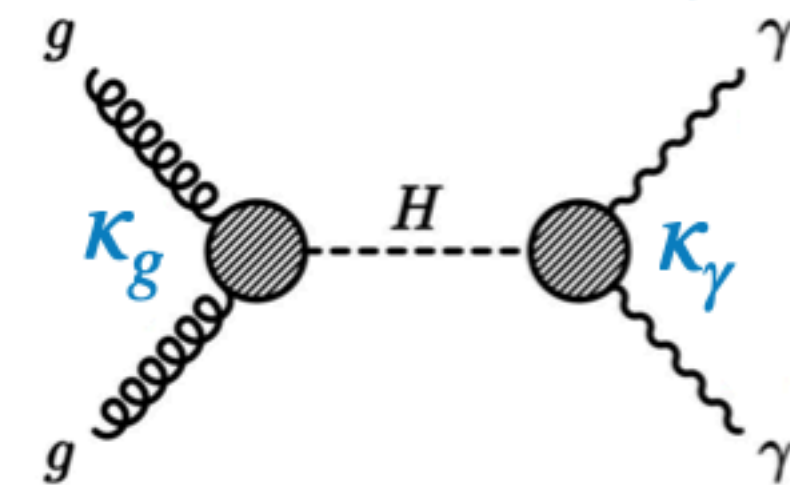


# Interpretations: K-Framework

- Multiplicative factors ( $\kappa_i$ ) factors applied to the SM XSs and BRs and can be used to constrain the couplings of the Higgs boson to particles

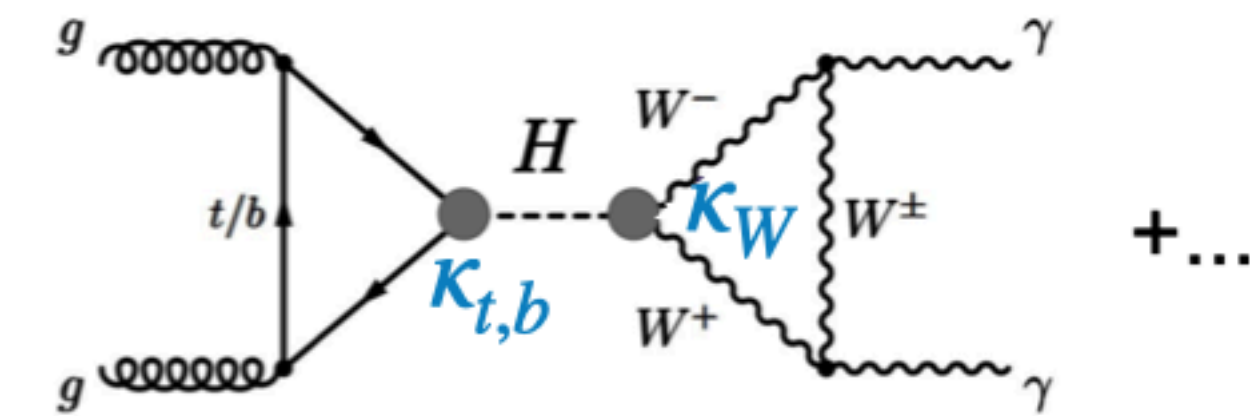
- Two ways to do this, for example in  $H \rightarrow \gamma\gamma$ :

Use effective  $H \rightarrow \gamma\gamma$  coupling



or

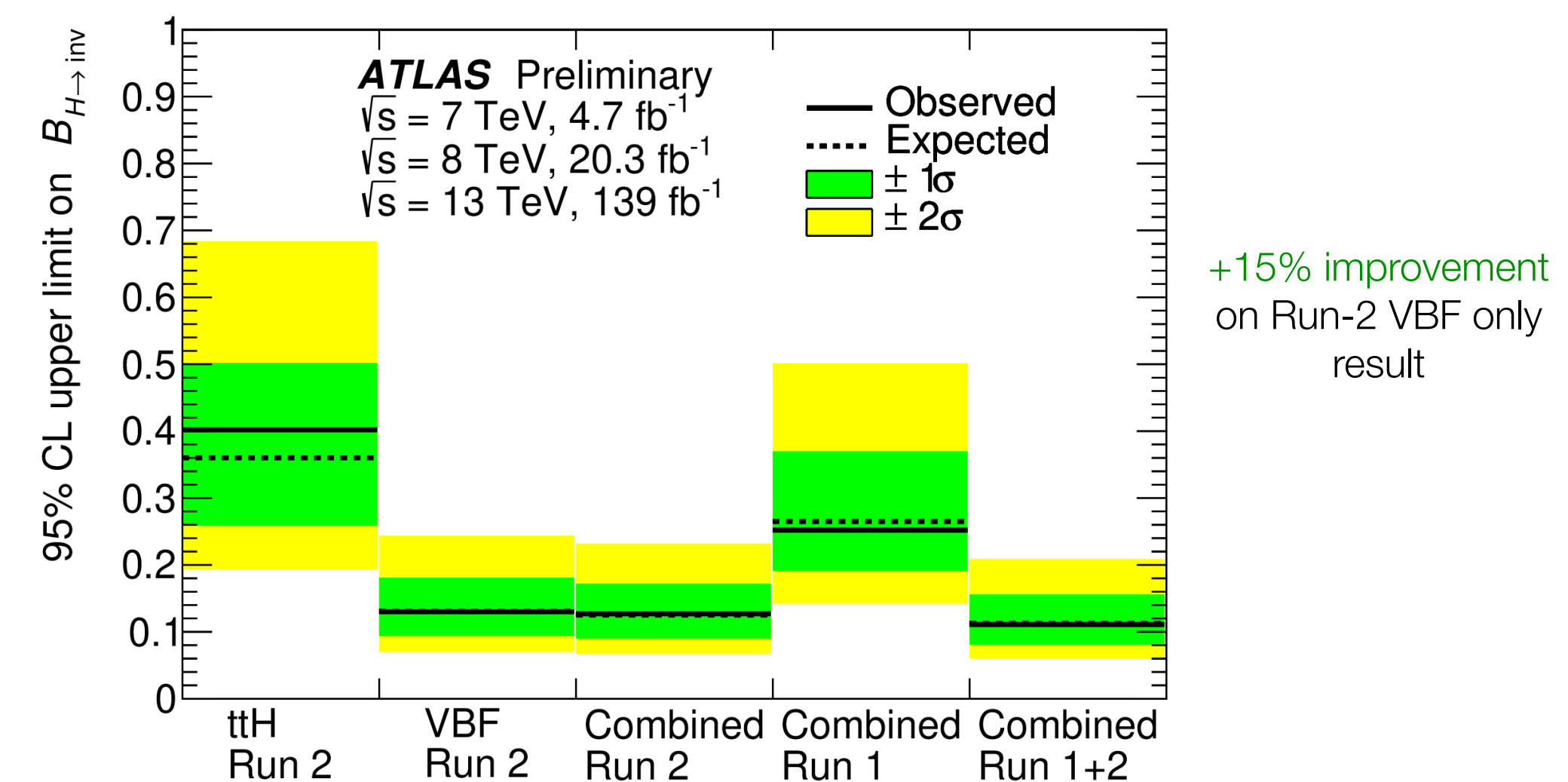
Use SM (e.g.  $\kappa_t$ ) in model



- Higgs width can also be affected by
  - Invisible decays ( $B_i$ ): Identified through  $E_T^{miss}$  (SM  $H \rightarrow ZZ^* \rightarrow 4\nu$ )
  - Undetected decays ( $B_u$ ): final states that analyses are not sensitive to (such as, light quarks, BSM particles without sizeable  $E_T^{miss}$ )

$$\Gamma_H(\kappa, B_i, B_u) = \frac{\sum_j B_j^{SM} \kappa_j^2}{1 - B_i - B_u} \Gamma_H^{SM}$$

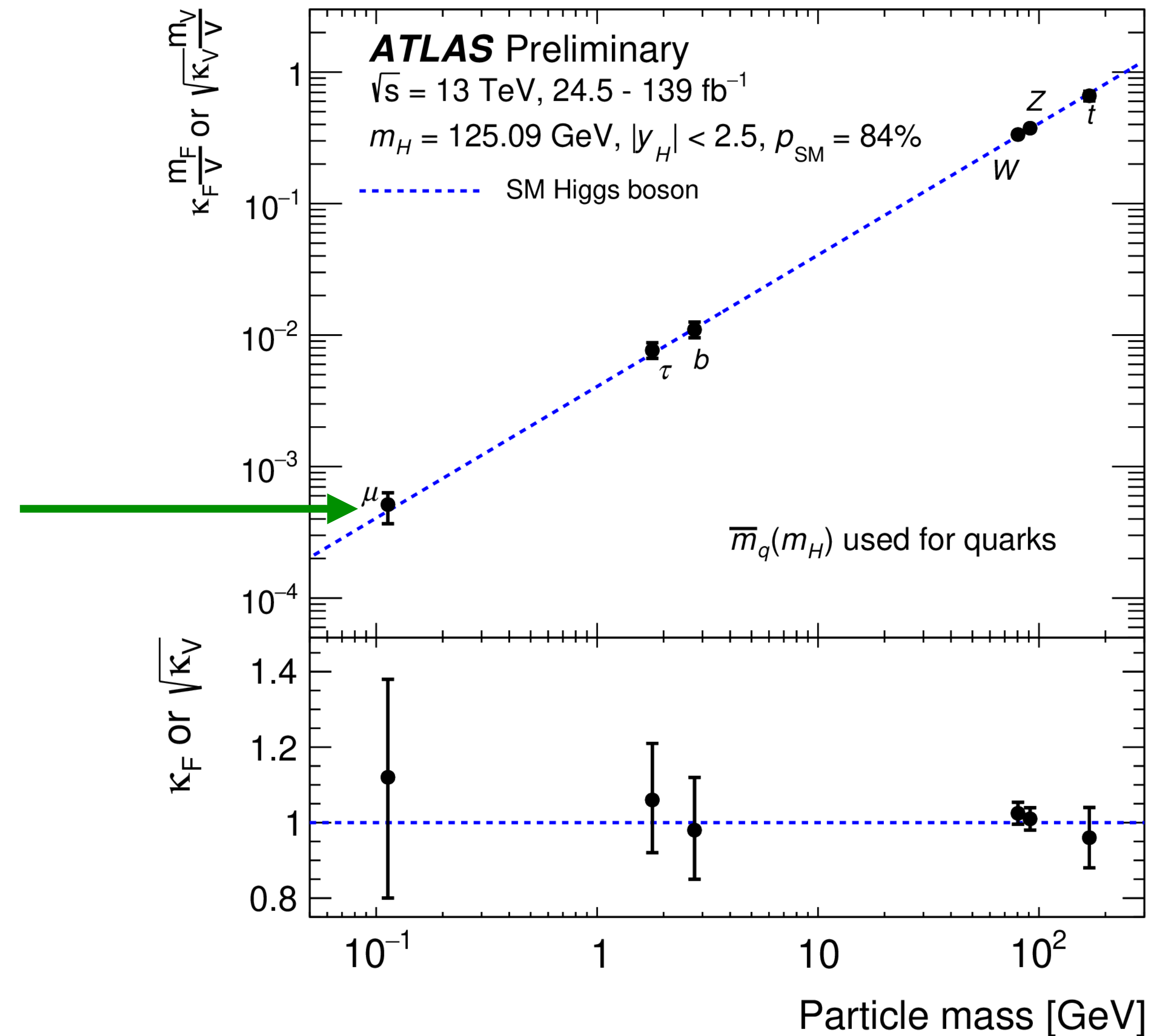
Not covered here: **new results** (with interpretations) of  $H \rightarrow \text{inv}$  combination available — [ATLAS-CONF-2020-052](https://atlas.cern/ATLAS-CONF-2020-052)





# Measurement of Yukawa couplings

- Assuming no new physics affects loop-induced processes or Higgs boson total width
  - Fit  $\kappa$  for W, Z, t, b,  $\tau$  and  $\mu$
- $\kappa_\mu$  is no longer a limit but a measurement with full Run-2 dataset using latest  $H \rightarrow \mu\mu$  result
- Good agreement with SM expectation through three orders of magnitude of particle mass

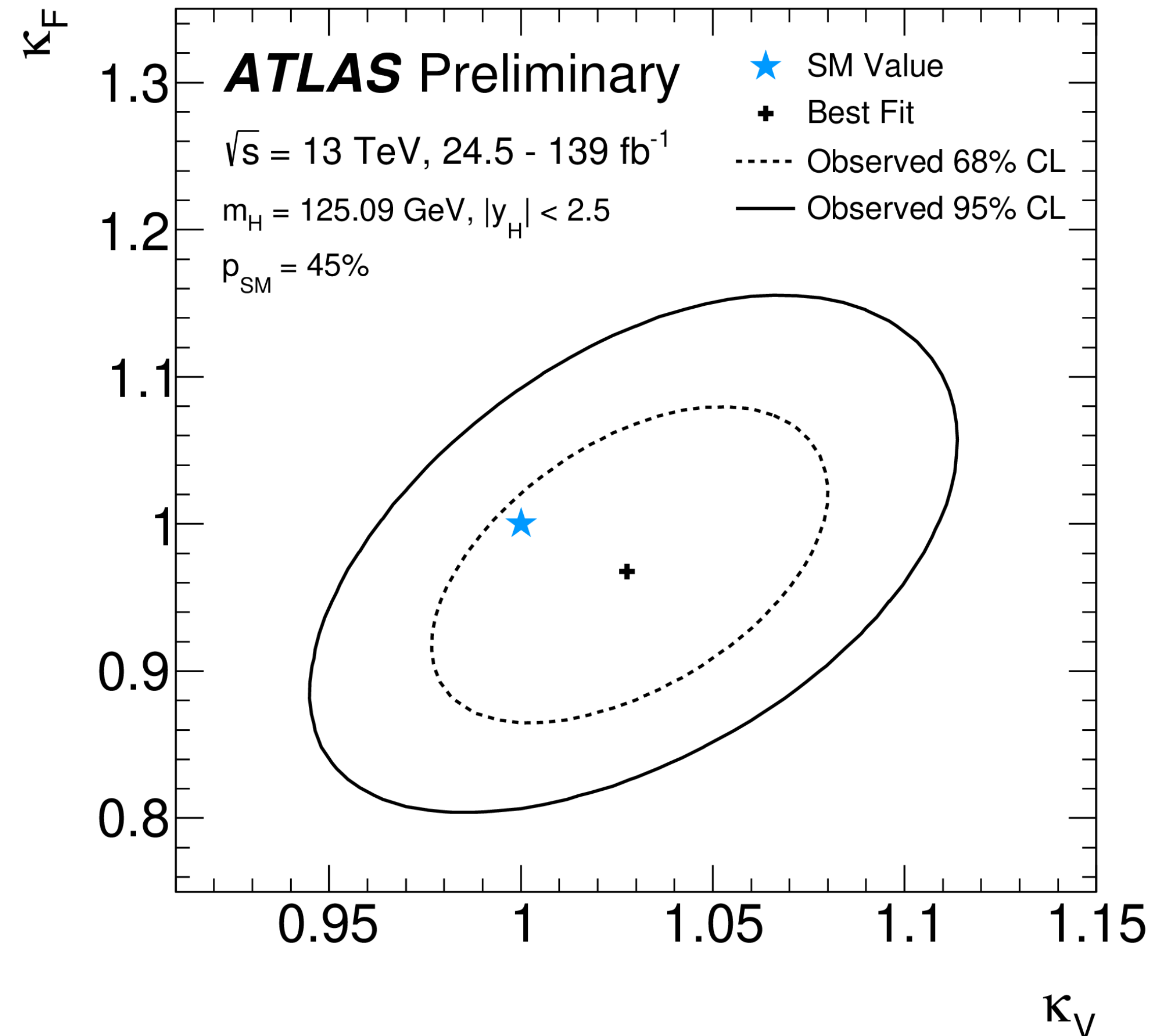


# Fermion and gauge boson couplings: $\kappa_V$ and $\kappa_F$

- Study the universal coupling-strength of:
  - Gauge bosons  $\kappa_V = \kappa_W = \kappa_Z$
  - Fermions  $\kappa_F = \kappa_t = \kappa_b = \kappa_\mu = \kappa_\tau$
- Assume no invisible or undetected decays and that  $\kappa_V, \kappa_F > 0$
- Best-fit values:

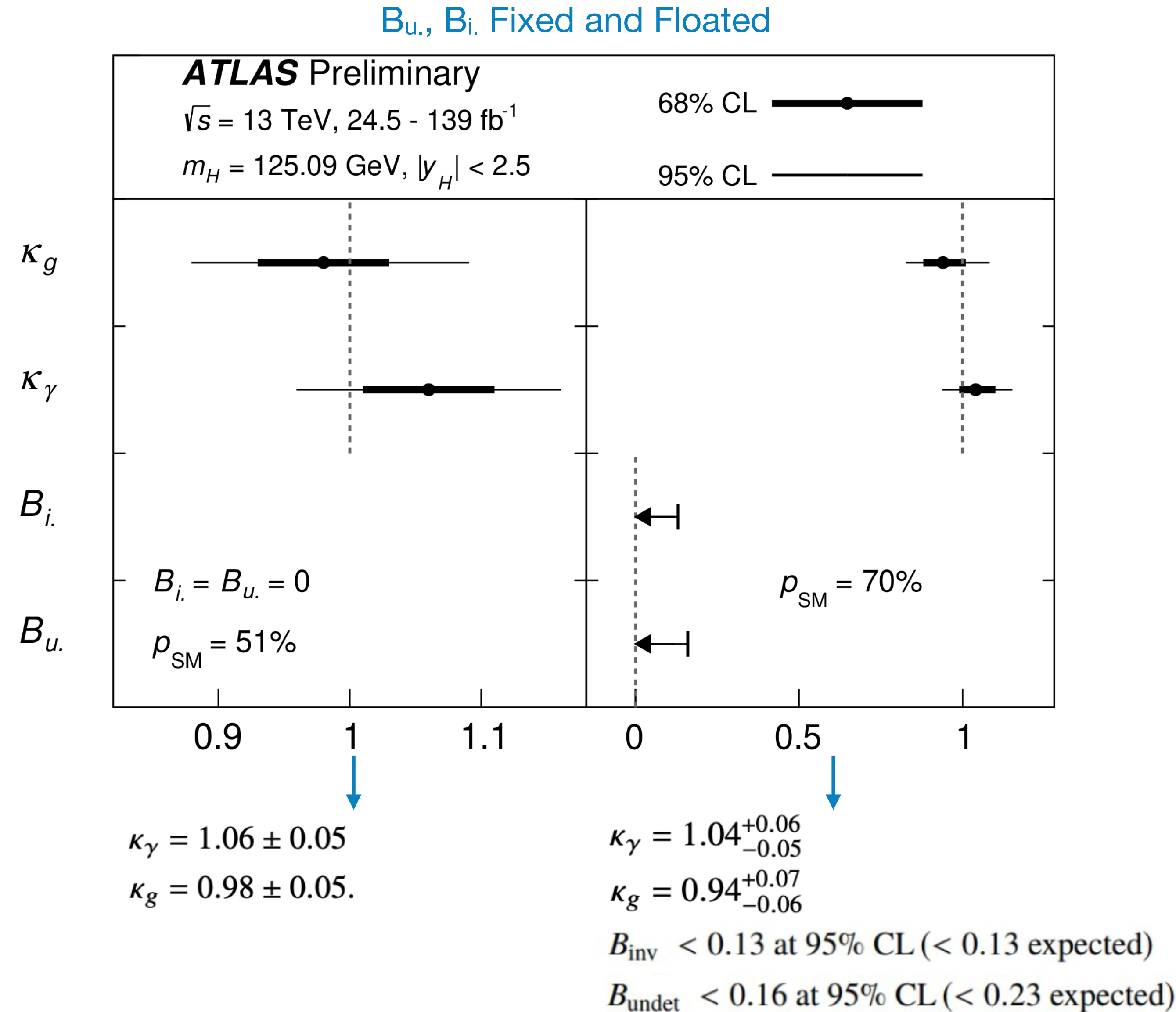
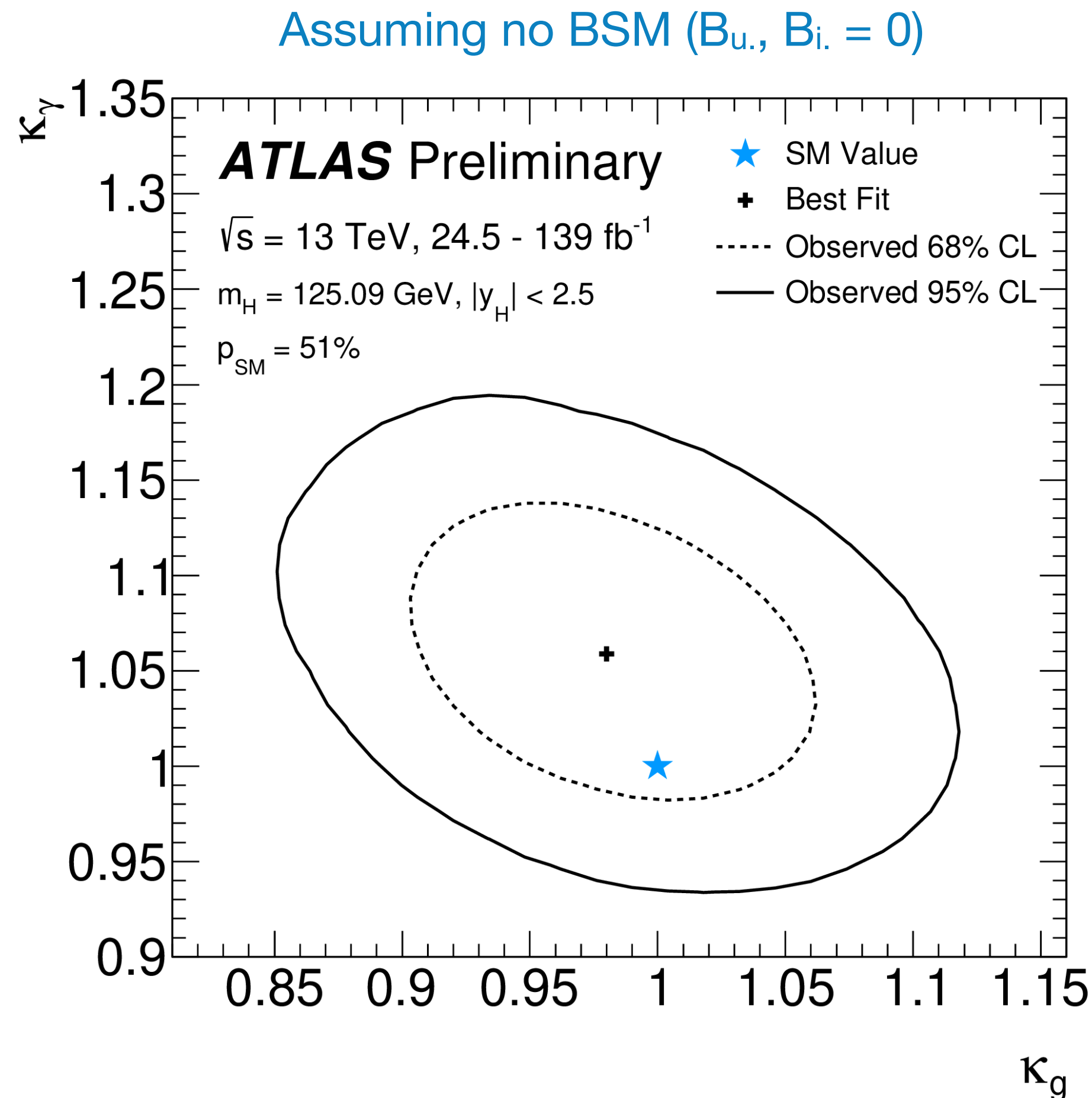
$$\kappa_V = 1.03 \pm 0.03$$

$$\kappa_F = 0.97 \pm 0.07$$



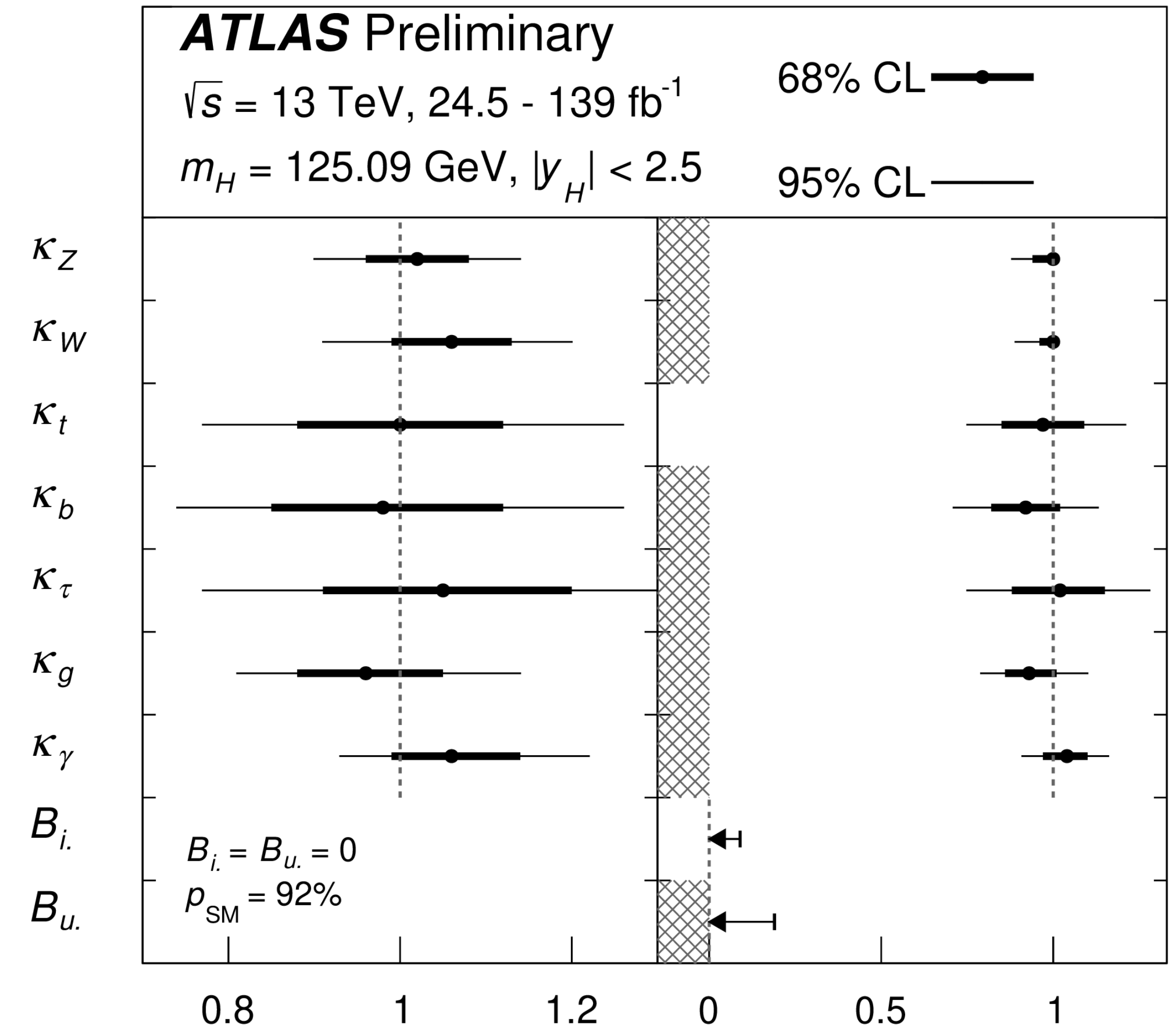
# Probe BSM contributions in loops and decays: $\kappa_g$ and $\kappa_\gamma$

- Effective couplings  $\kappa_g$  and  $\kappa_\gamma$  are **sensitive to new physics in loops**
- Investigate two assumptions ( $B_i$ ,  $B_u$  are fixed or floated)



# Generic parametrization

- Effective coupling-strength modifiers,  $\kappa_g$  and  $\kappa_\gamma$ , for loop processes
- Two scenarios considered —  $B_i$ ,  $B_u$  fixed and floated
  - Assume  $|\kappa_V| < 1$  to constrain undetected BR
- Negative  $\kappa_t$  **excluded** at a p-value of 0.0037 (or  $2.7\sigma$ ) assuming no BSM contributions to total width



## $\kappa$ -framework Limitations

- Inherently leading order
- Missing sensitivity to differential information



Motivates use of frameworks such as **EFT** or models such as 2HDM, MSSM (backup slides) among others



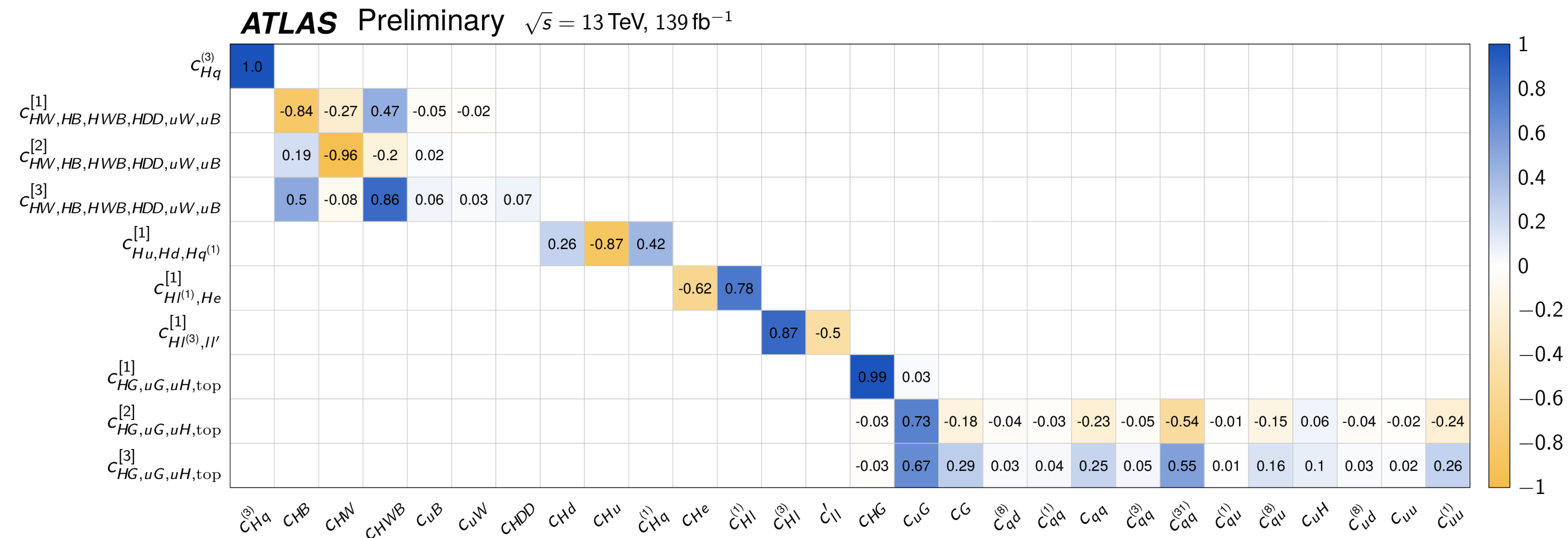
# Interpretations: SM Effective Field Theory

- Model independent way to parametrize effects of BSM theories
- Parametrized as higher dimensional operators that respect SM symmetry → cross-sections measurements used to constrain Wilson coefficient ( $c_i$ ) at fixed scale ( $\Lambda = 1 \text{ TeV}$ )

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_j \frac{b_j}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots$$

- Consider  $d = 6$  operators using the Warsaw basis

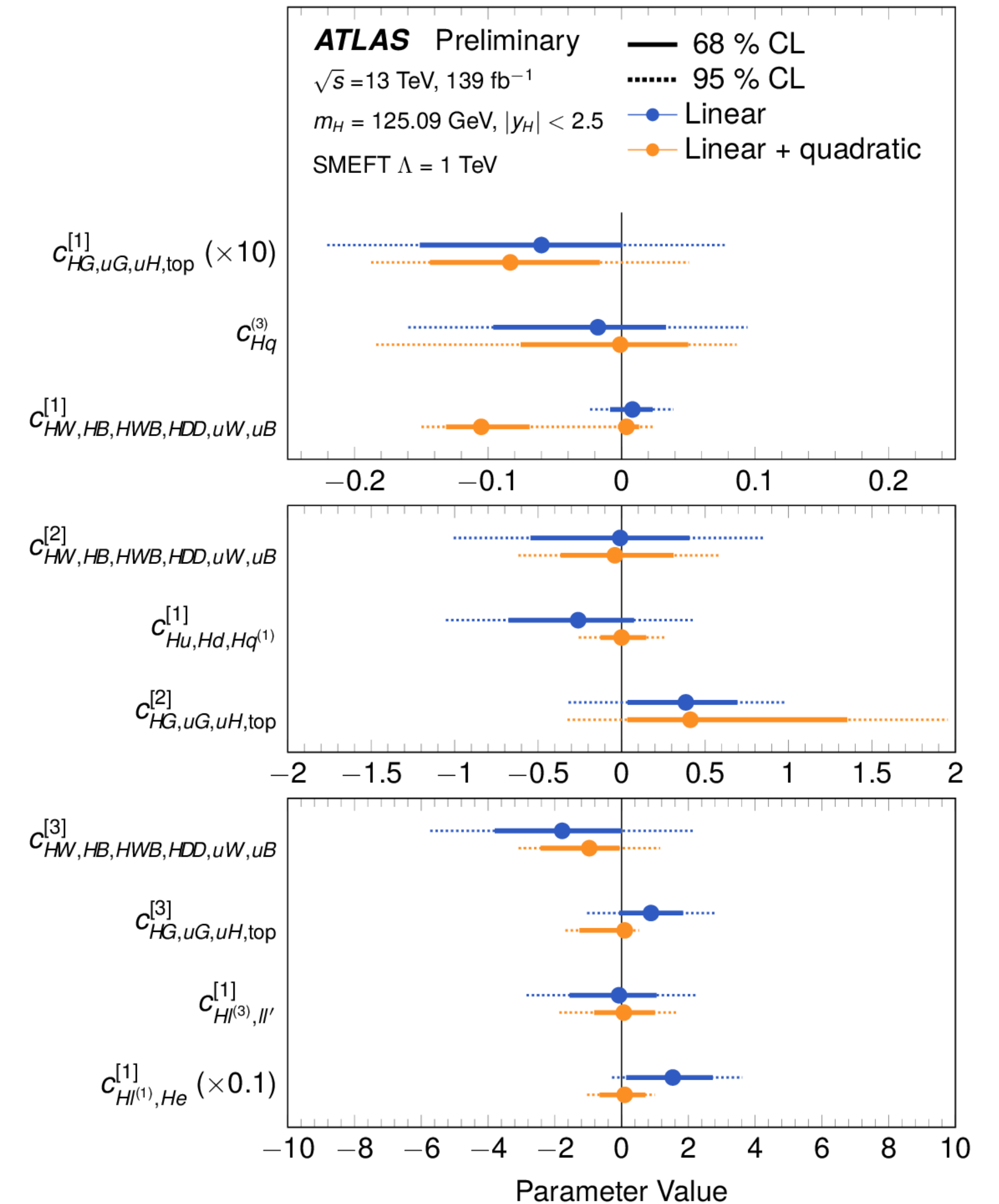
- Interpretations using **full** STXS scheme
- Probe as many  $c_i$  as possible
- Eigenvector decomposition to determine combinations of SMEFT operators analysis is sensitive to





# SM Effective Field Theory

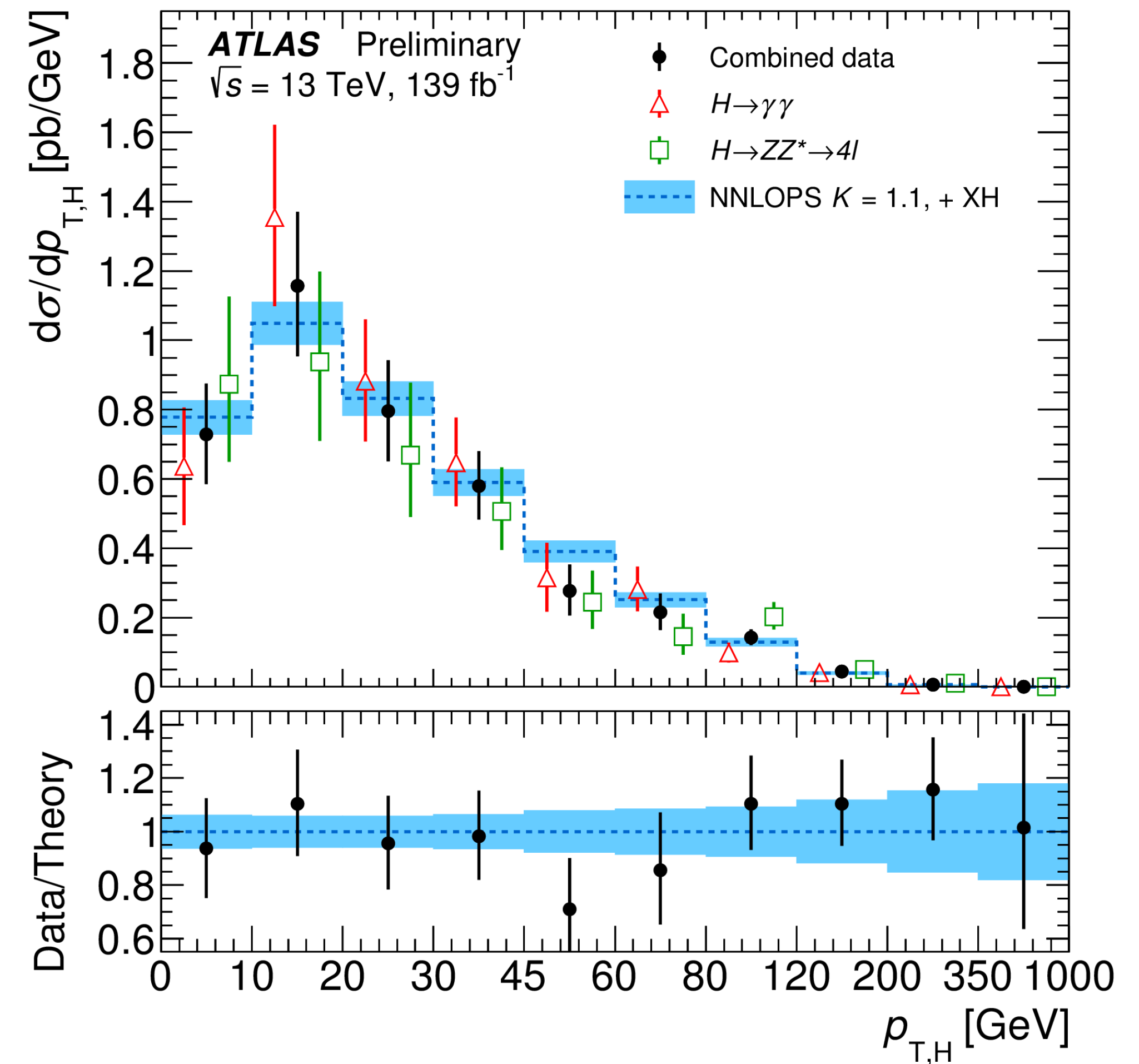
- Matrix element squared from d=6 operators have linear ( $c_j/\Lambda^{-2}$ ) and quadratic ( $c_j c_k/\Lambda^{-4}$ ) terms
- Parametrize ( $\sigma \times B$ ) using linear or linear+quadratic terms
  - For quadratic terms, interference with d = 8 operators not considered
- Including **quadratic terms** shows non-negligible influence
- All results compatible with the SM
- For more details on combination: [ATLAS-CONF-2020-053](#)



# Differential combination: $p_{T,H}$

- Measuring differential XS probes the SM and is largely model independent
- Deviations in  $p_{T,H}$  distribution can indicate BSM physics such as:
  - Modified light quark coupling at low  $p_{T,H}$
  - New BSM particles in ggF loop at high  $p_{T,H}$
- Input results:
 

Analysis	Decay Channel	Prod. Modes	$\mathcal{L}$ [fb $^{-1}$ ]	Link
	$H \rightarrow \gamma\gamma$	All	139	<a href="#">ATLAS-CONF-2019-029</a>
	$H \rightarrow ZZ^*$	All	139	<a href="#">arXiv:1504.05833</a>
- P-value of 78%, compatible with SM
- More results to be added, detailed interpretations on combined to be performed
  - Other observables to be included in combination



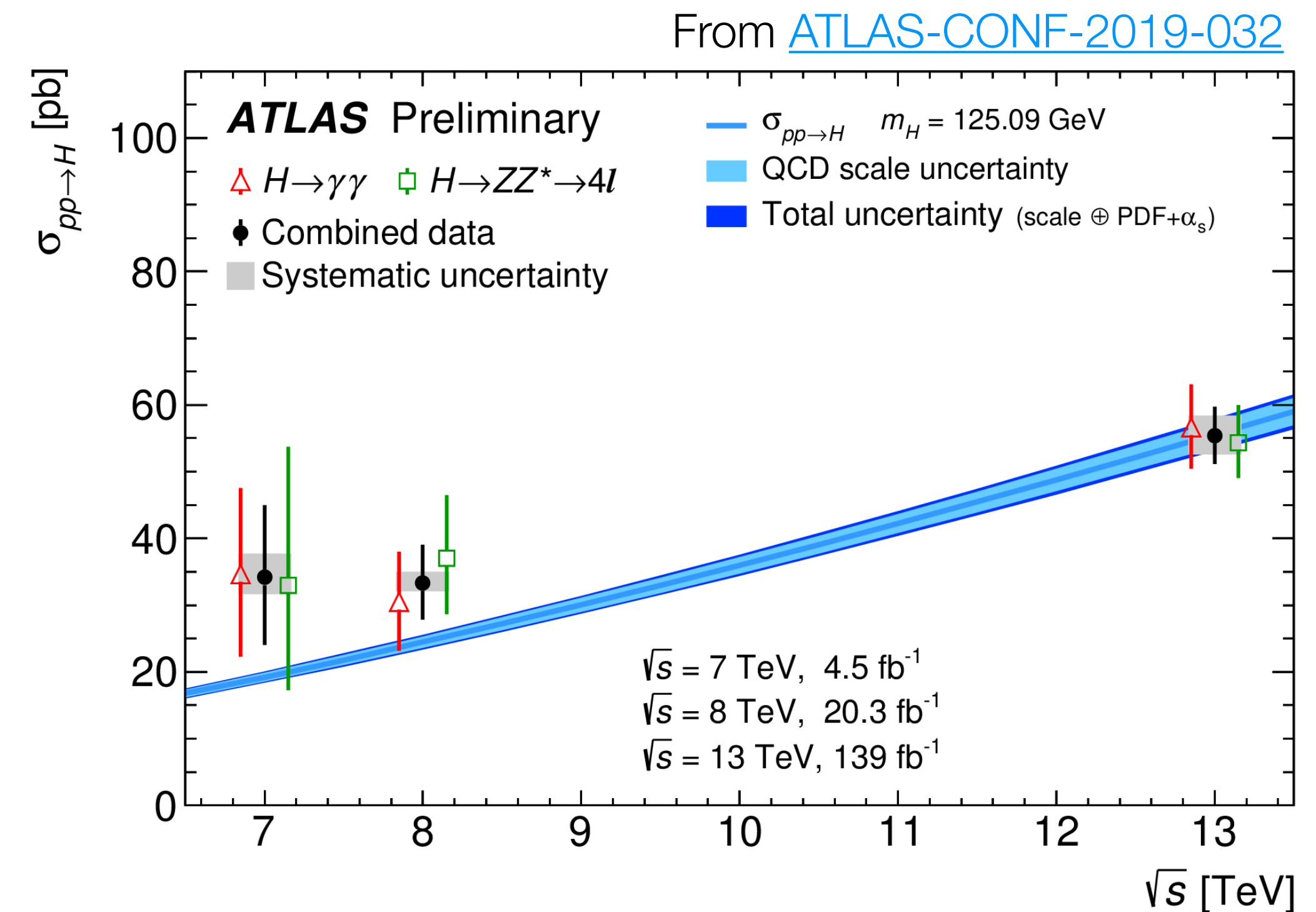
# Summary

## Couplings

- Extended to include full Run-2 results for several channels!
- Signal strength, production modes and branching ratios combined measurements made with **latest available** results from analyses
- Extended phase space measured in STXS!
- Interpretations of XS and BRs show **consistency with SM**

## Differential

- $p_{T,H}$  combination for  $H \rightarrow ZZ^* \rightarrow 4\ell$  and  $H \rightarrow \gamma\gamma$  show good agreement with SM



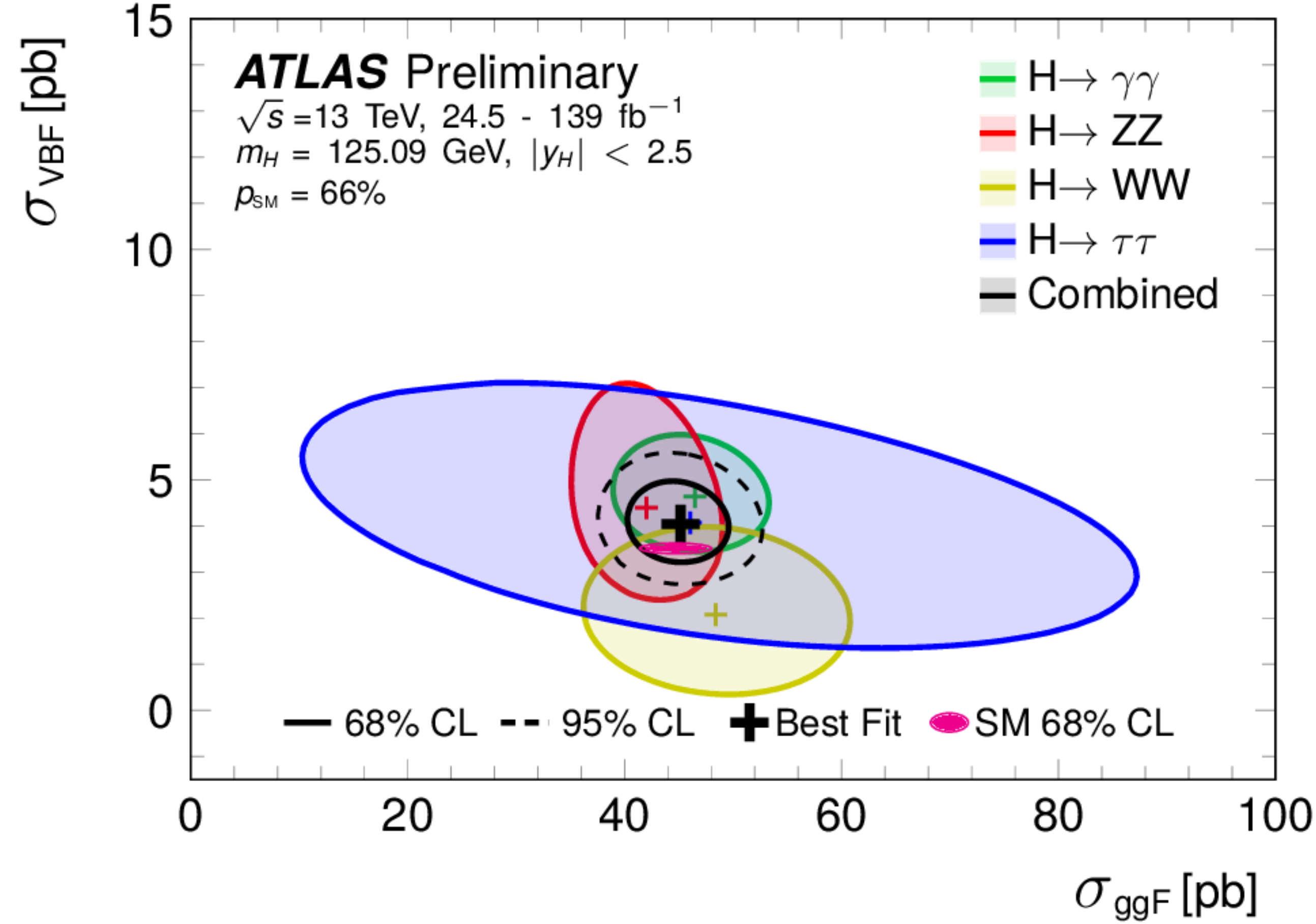
## Outlook

- More **full Run-2 results** to be added to combinations soon!
  - Most recent  $H \rightarrow WW$  ([ATLAS-CONF-2021-014](#))
  - Increased precision in several regions of phase space with addition of analyses
- Combination work on other observables to come soon too!

Backup

# Production mode measurements: Additional information

- Observed likelihood contours for VBF and ggF XS



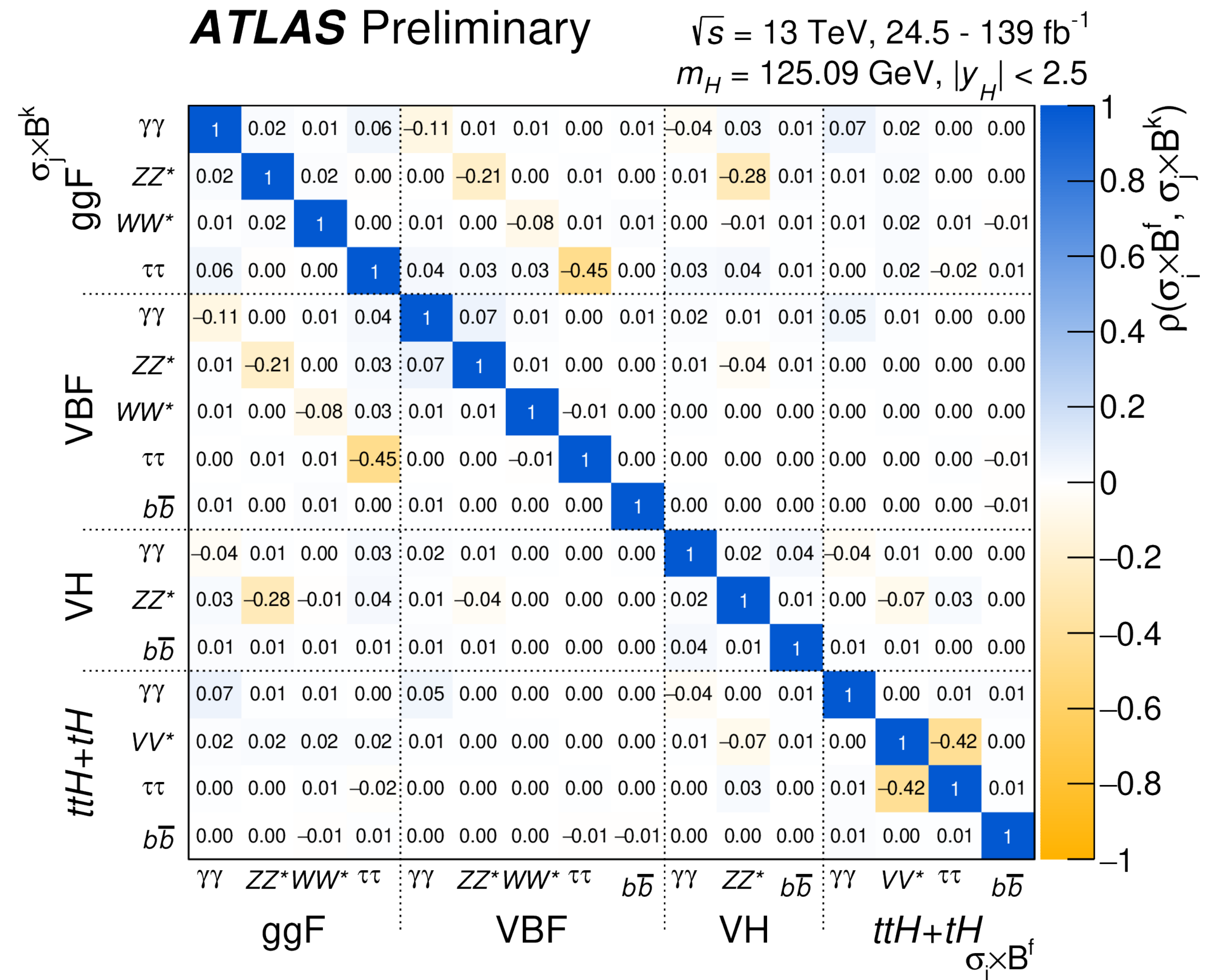
- Decomposition of uncertainty into components
- Systematic uncertainty approaching/same as statistical uncertainty

Process ( $ y_H  < 2.5$ )	Value [pb]	Uncertainty [pb]						SM pred. [pb]
		Total	Stat.	Syst.	Exp.	Sig. Th.	Bkg. Th.	
ggF	44.7	$\pm 3.1$	$\pm 2.2$	$\pm 2.2$	$+1.8$ $-1.7$	$+1.0$ $-0.9$	$+0.9$ $-0.7$	$44.7 \pm 2.2$
VBF	4.0	$\pm 0.6$	$\pm 0.5$	$\pm 0.4$	$+0.3$ $-0.2$	$\pm 0.3$	$\pm 0.1$	$3.51^{+0.08}_{-0.07}$
WH	1.45	$+0.28$ $-0.25$	$+0.20$ $-0.19$	$+0.18$ $-0.17$	$+0.13$ $-0.12$	$+0.08$ $-0.06$	$+0.10$ $-0.09$	$1.204 \pm 0.024$
ZH	0.78	$+0.18$ $-0.17$	$\pm 0.13$	$+0.12$ $-0.10$	$+0.08$ $-0.07$	$+0.07$ $-0.05$	$\pm 0.06$	$0.797^{+0.033}_{-0.026}$
$t\bar{t}H + tH$	0.64	$\pm 0.12$	$\pm 0.09$	$\pm 0.08$	$+0.06$ $-0.05$	$+0.03$ $-0.02$	$\pm 0.05$	$0.59^{+0.03}_{-0.05}$



# Production mode and branching ratios correlation matrix

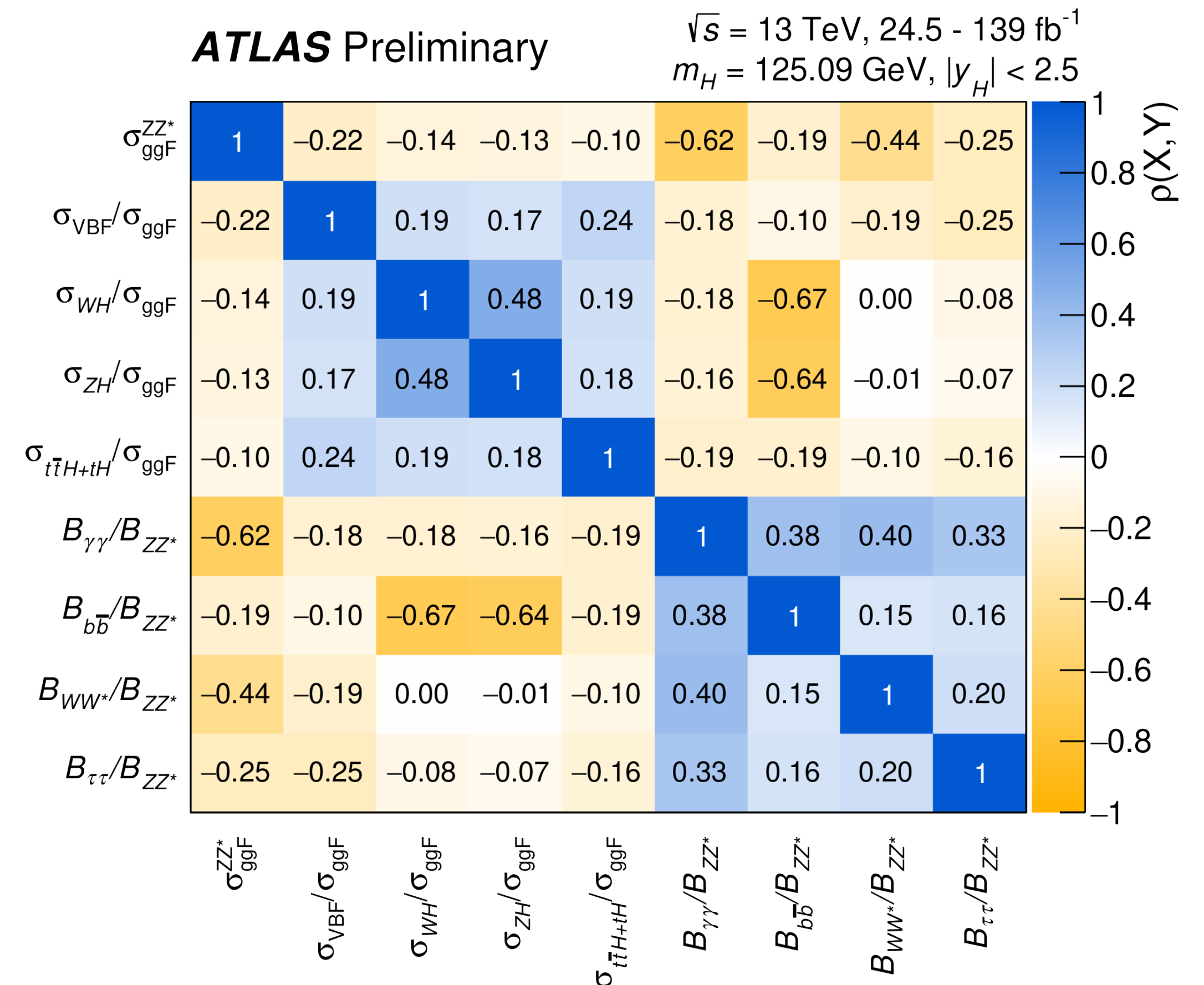
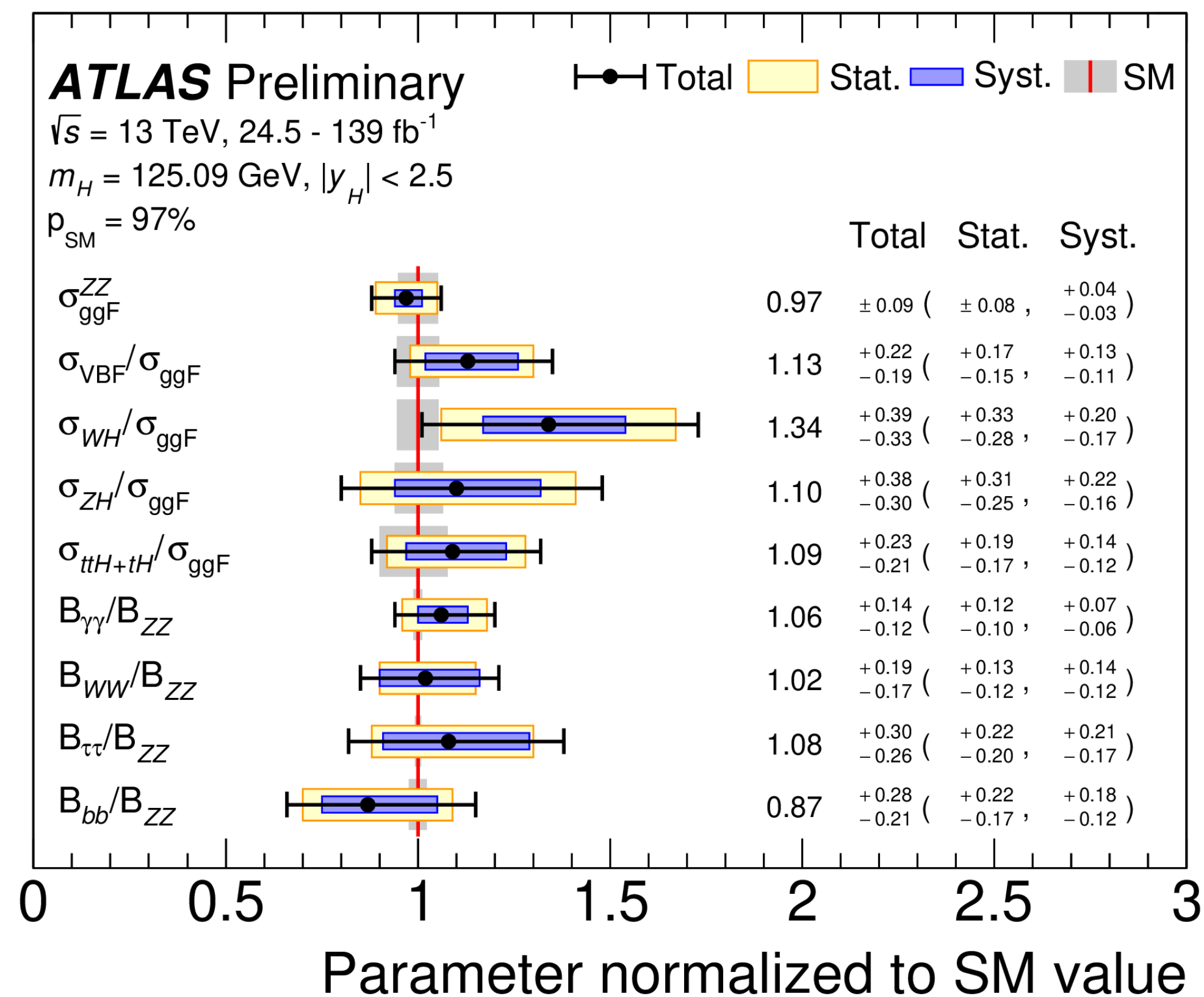
- Can also test the SM by measuring  $(\sigma \times B)_{\text{if}}$
- Statistically limited channels are fixed to the SM
  - $ggF \rightarrow H \rightarrow b\bar{b}$
  - $VH, H \rightarrow WW/\tau\tau$
- Relative fraction of ZH and WH fixed to the SM
- **Good agreement** with the SM (p-value of 87%)



# Production mode and branching ratios: ratio model

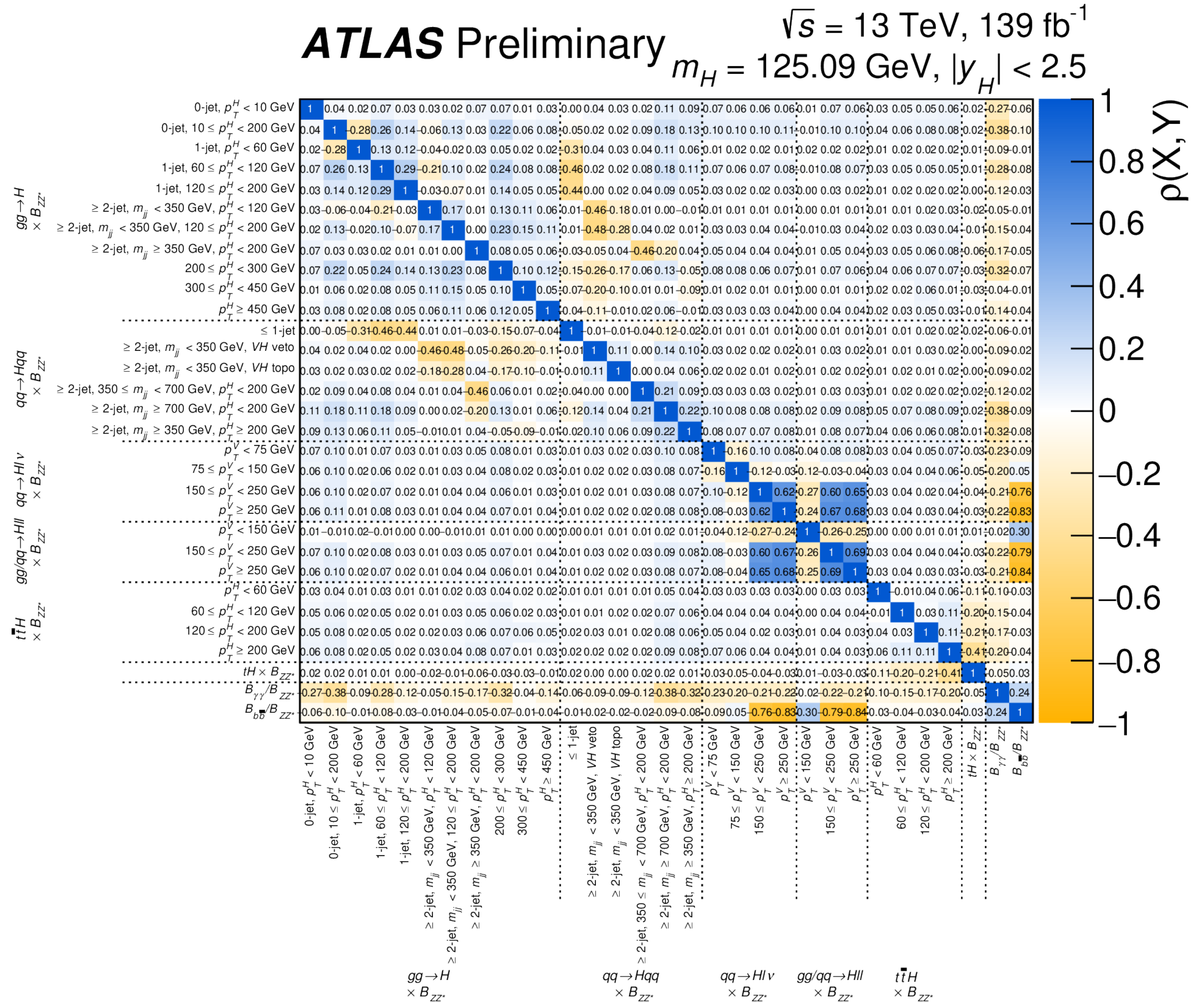
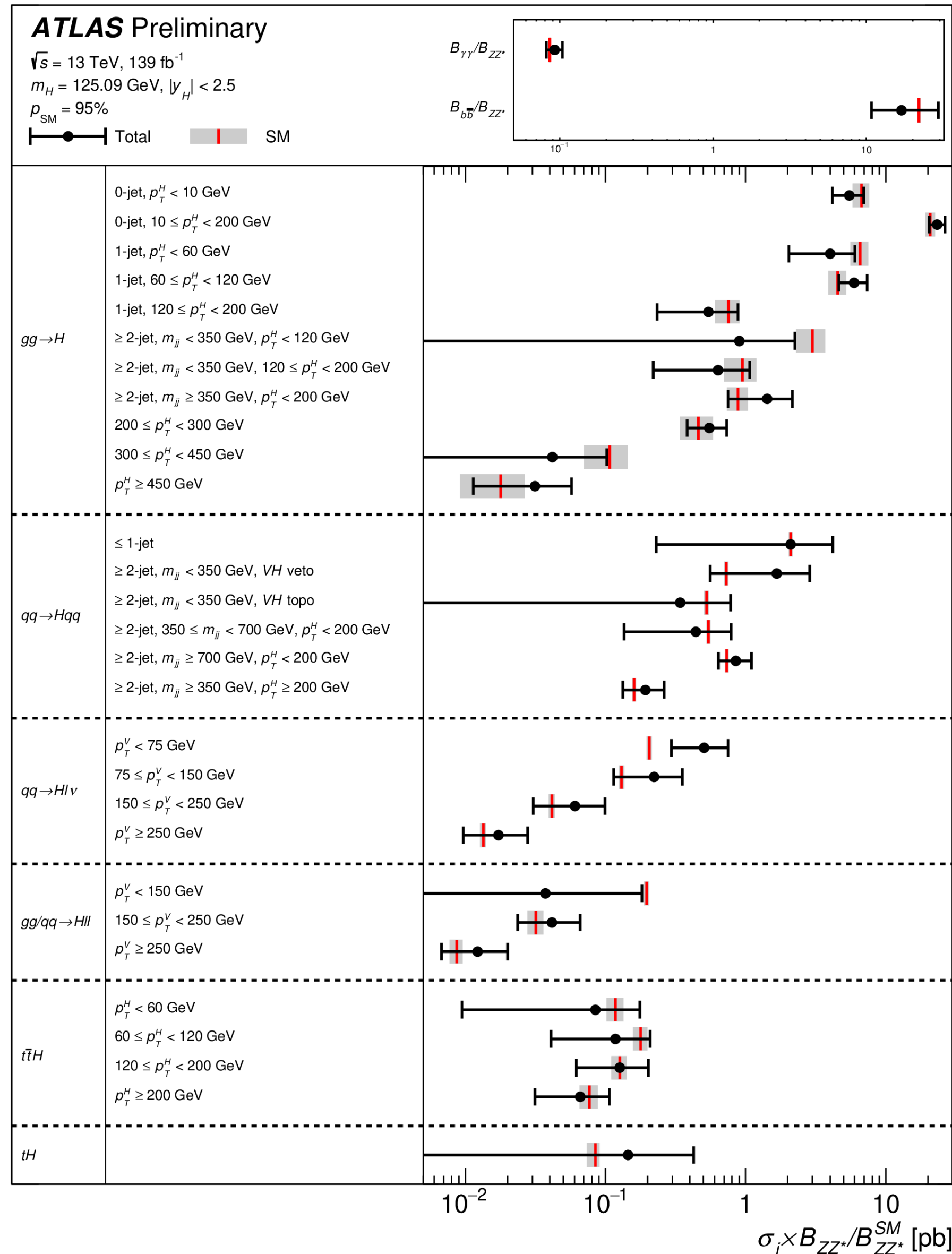
- Can also measure  $(\sigma \times B)_{if}$  with respect to a reference process
- $gg \rightarrow H \rightarrow 4\ell$  chosen due to its small experimental uncertainties
  - Systematics cancel in the ratio!

$$(\sigma \times B)_{if} = \sigma_{ggF}^{ZZ} \cdot \left( \frac{\sigma_i}{\sigma_{ggF}} \right) \cdot \left( \frac{B_f}{B_{ZZ}} \right)$$





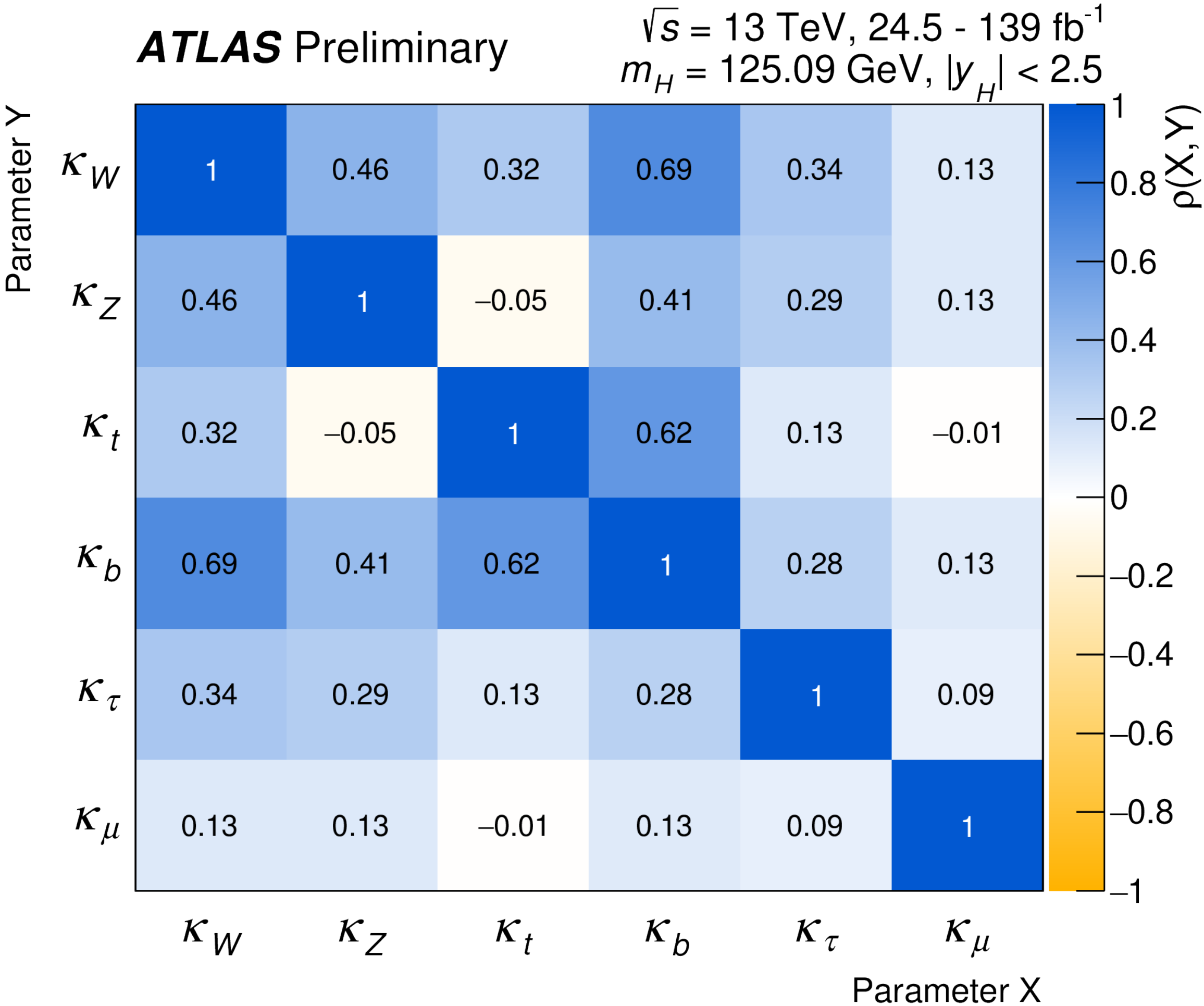
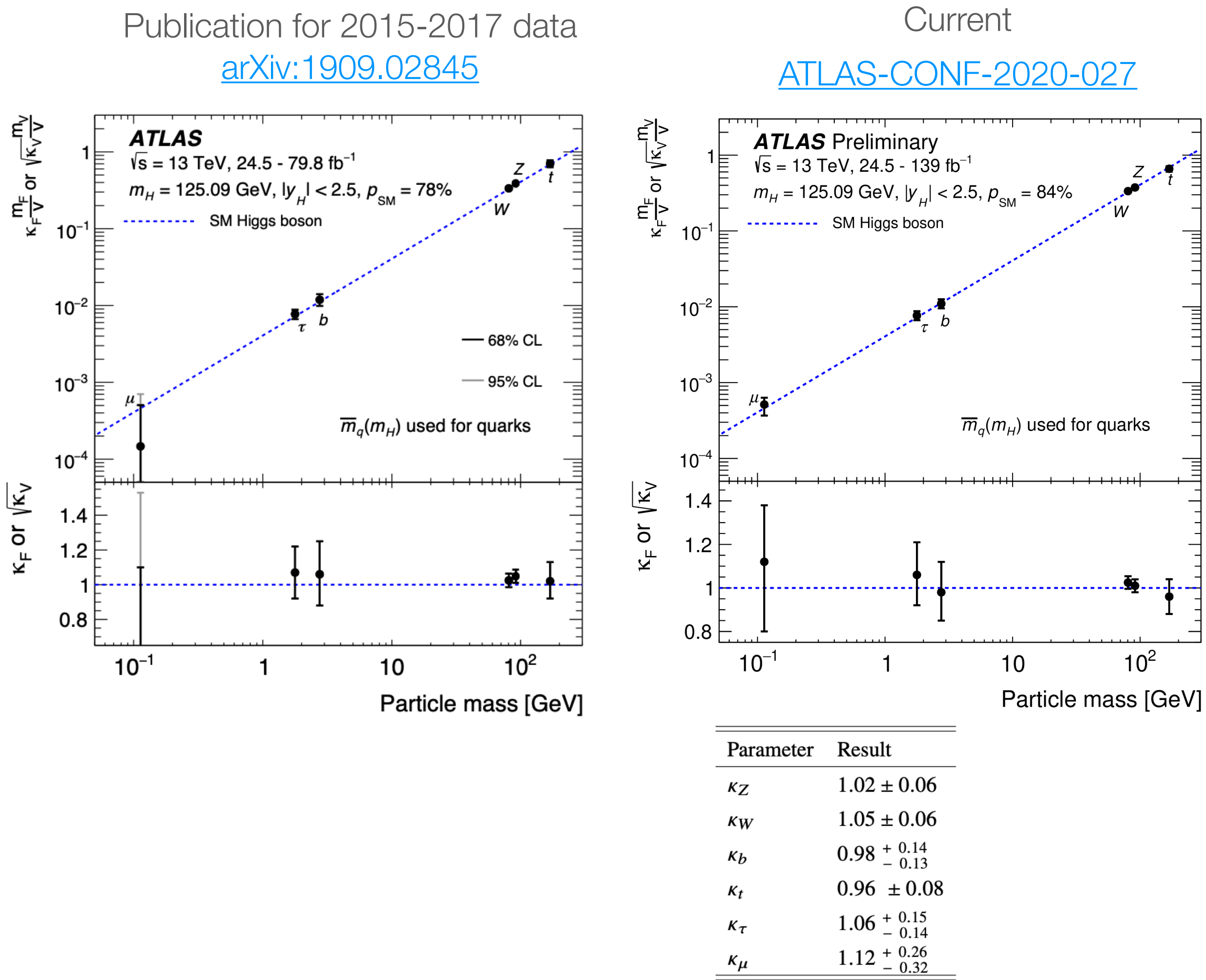
# Additional results: STXS





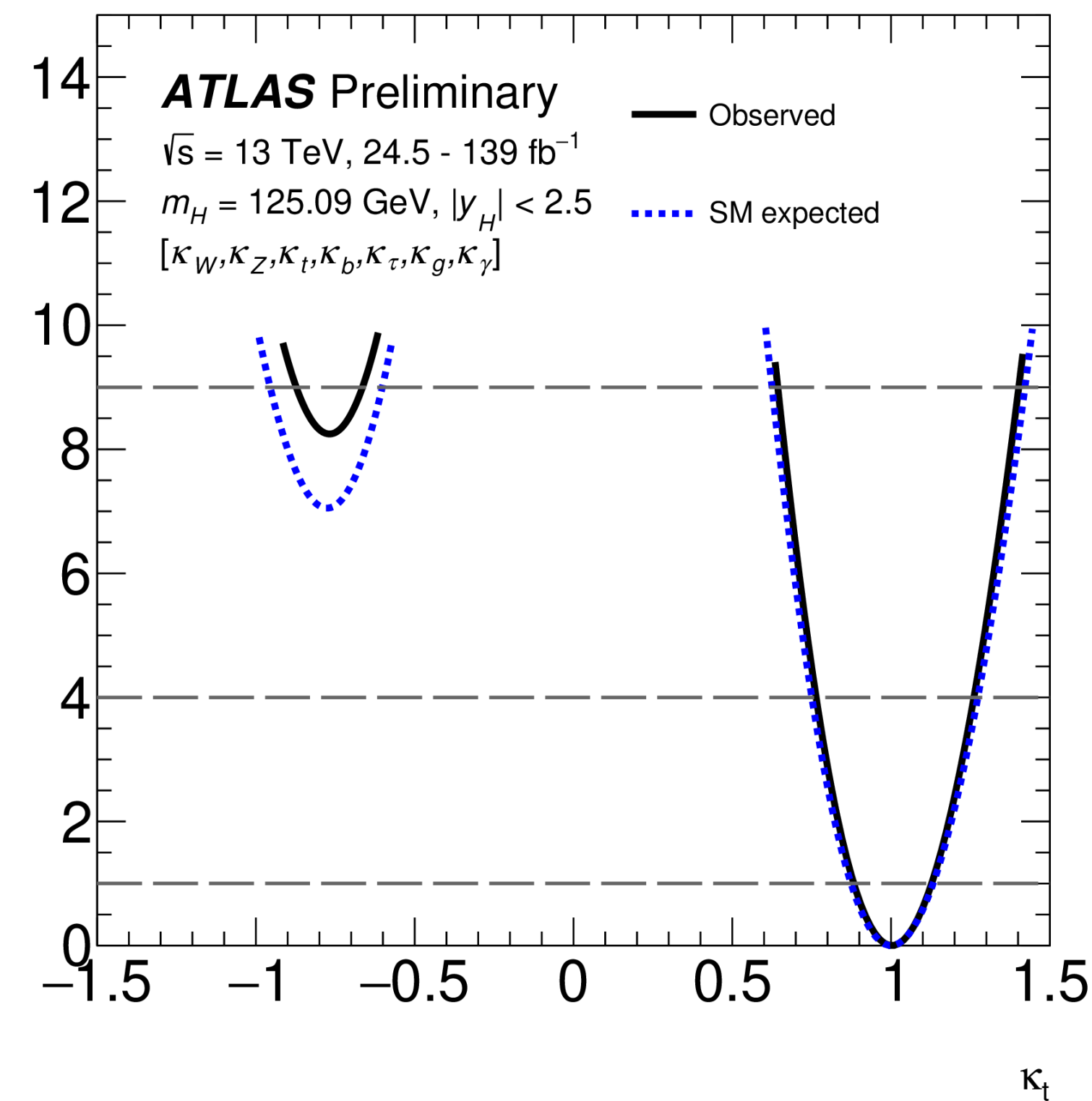
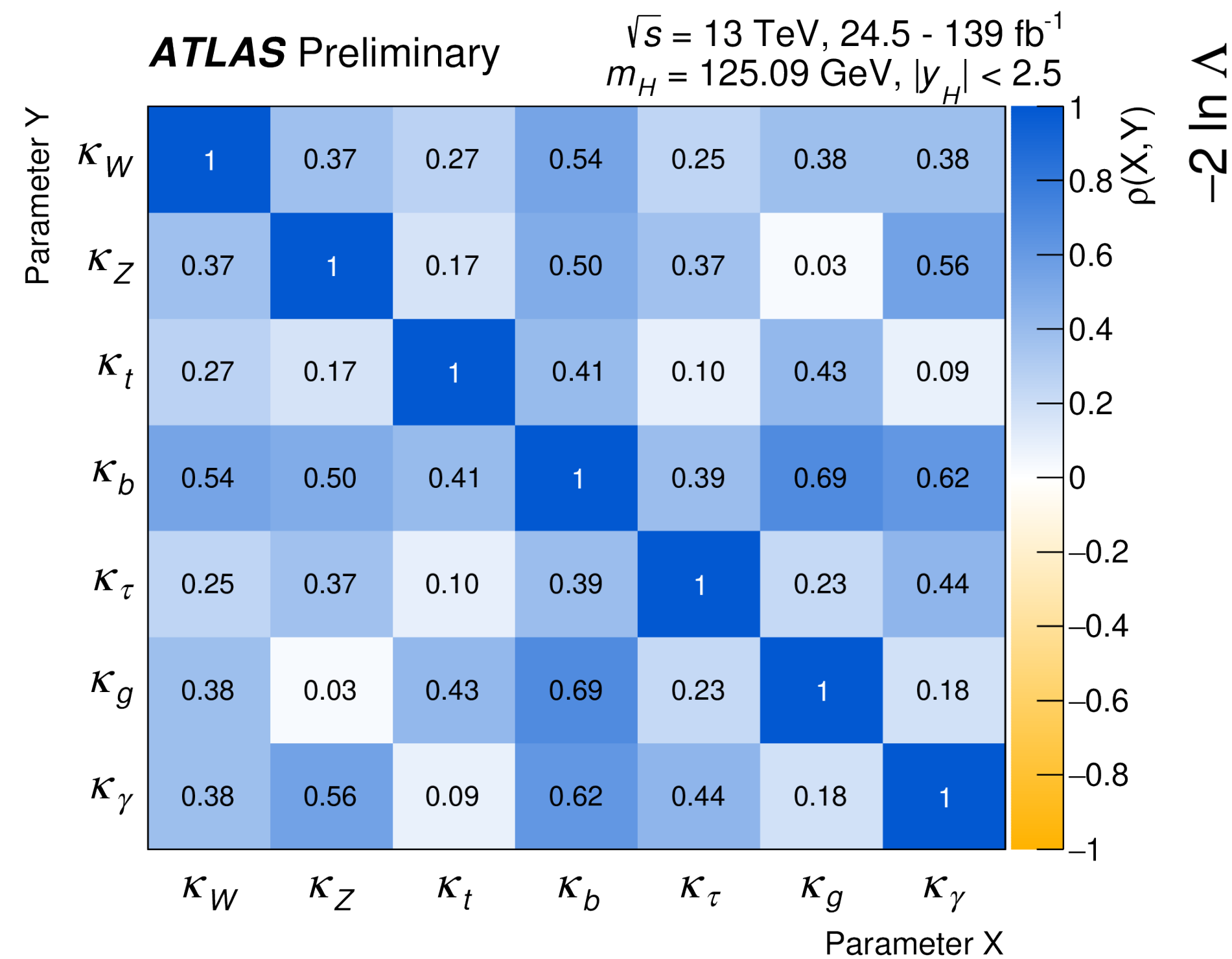
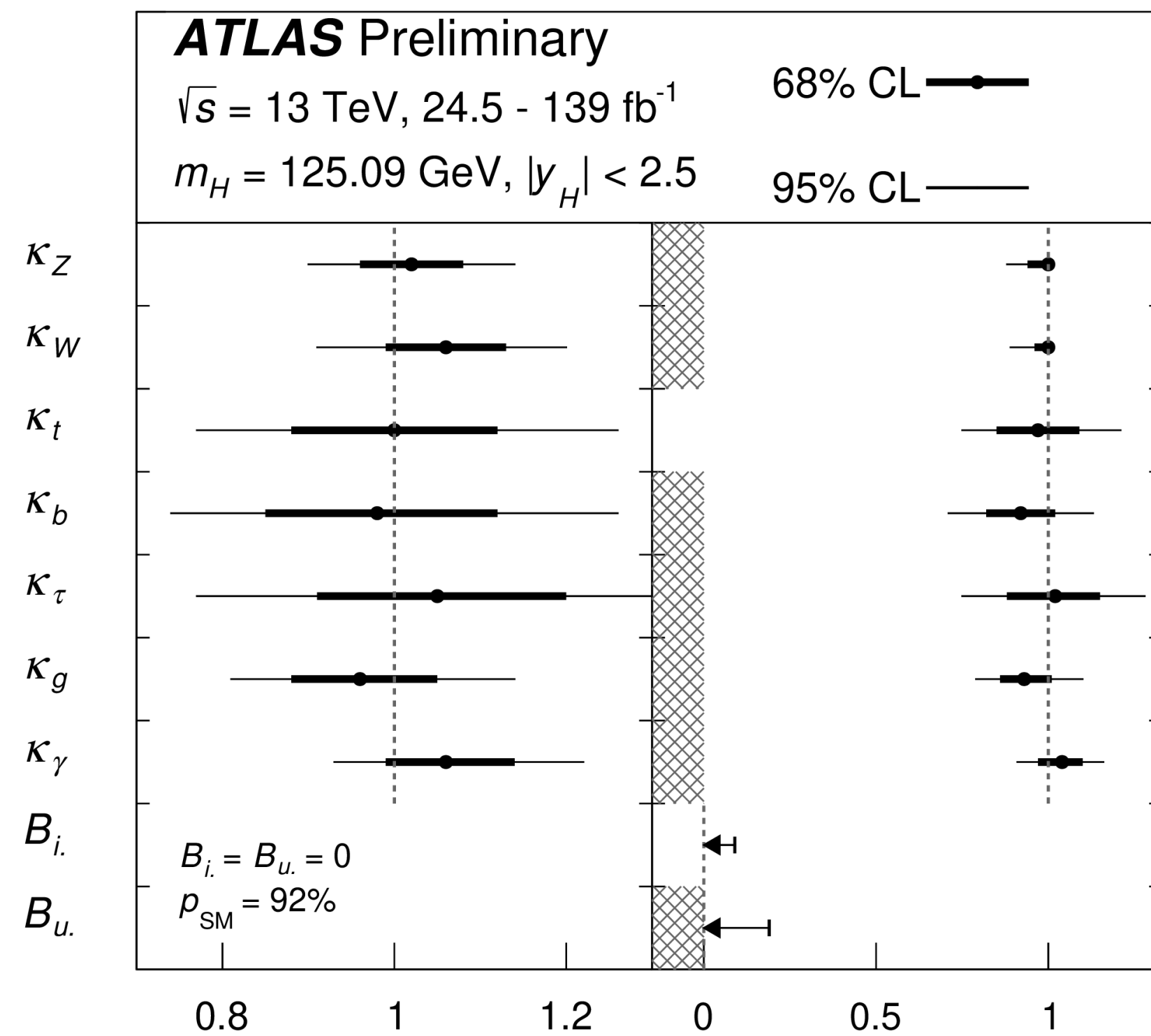
# Generic parametrization

- Excluding BSM couplings, fit for  $\kappa$  for W, Z, t, b,  $\tau$  and  $\mu$



# Generic parametrization: float effective couplings

- Effective coupling-strength modifiers,  $\kappa_g$  and  $\kappa_\gamma$ , for loop processes
- Negative  $\kappa_t$  **excluded** at  $2.9\sigma$  ( $2.7\sigma$  expected) assuming no BSM contributions to total width

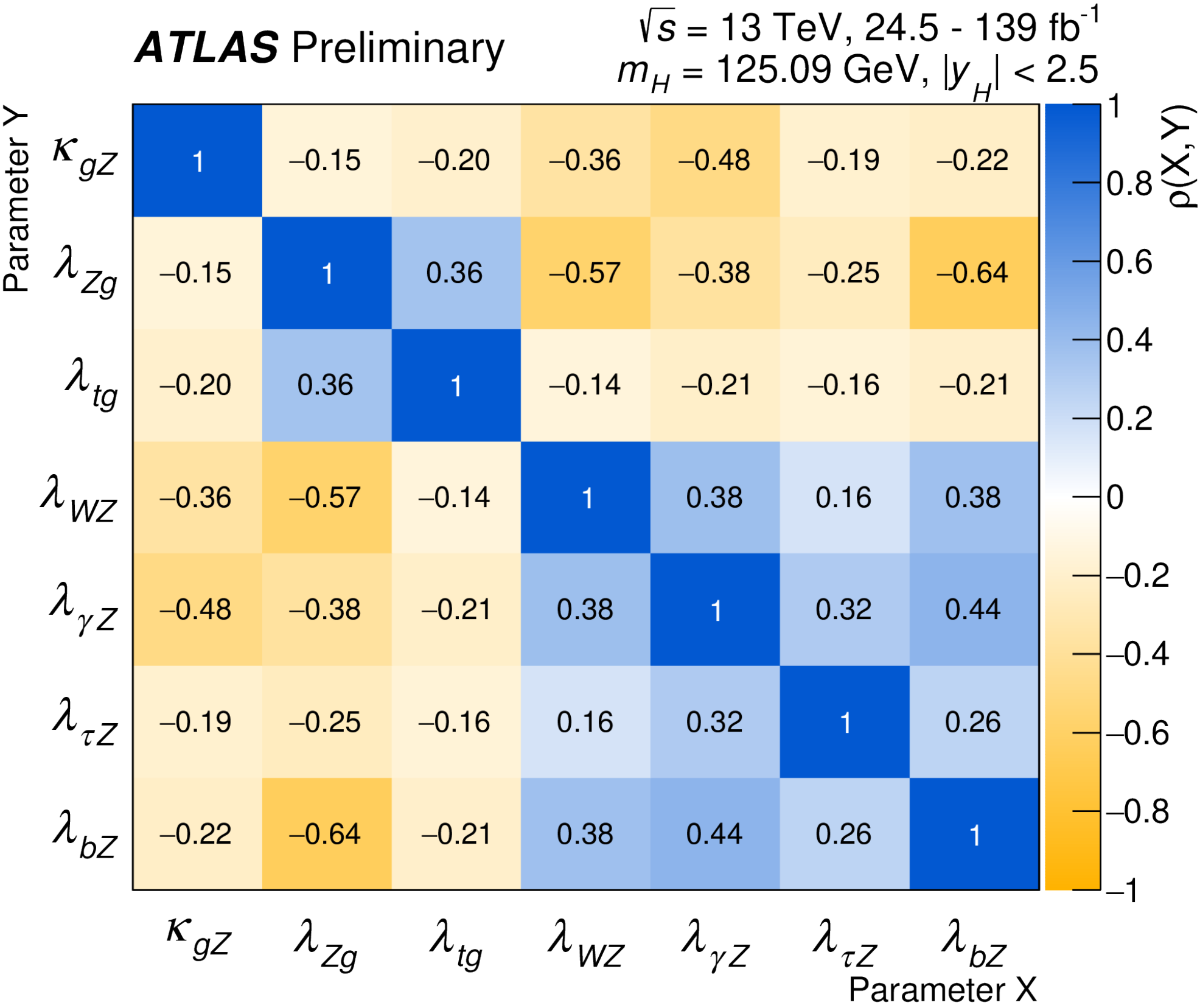




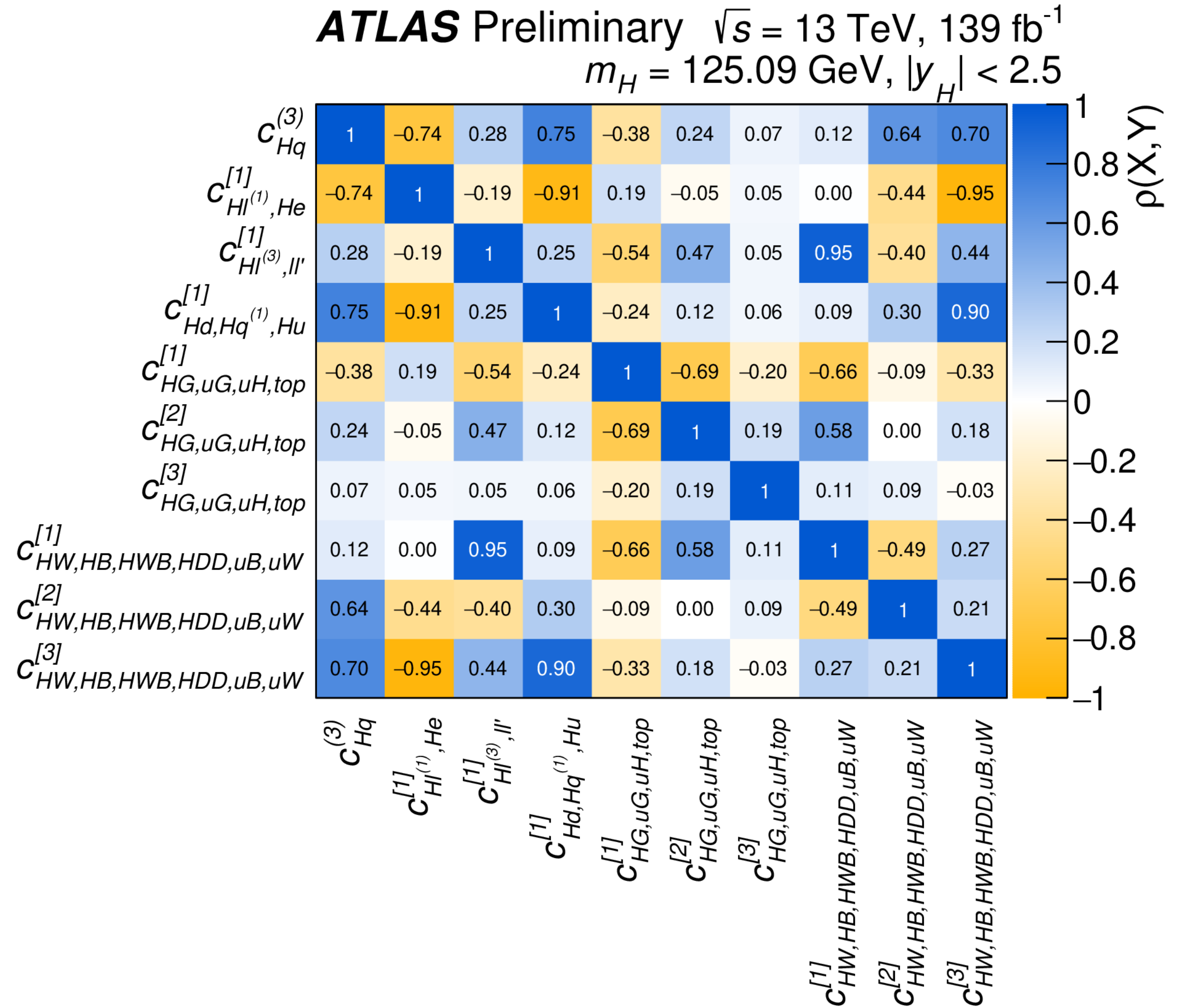
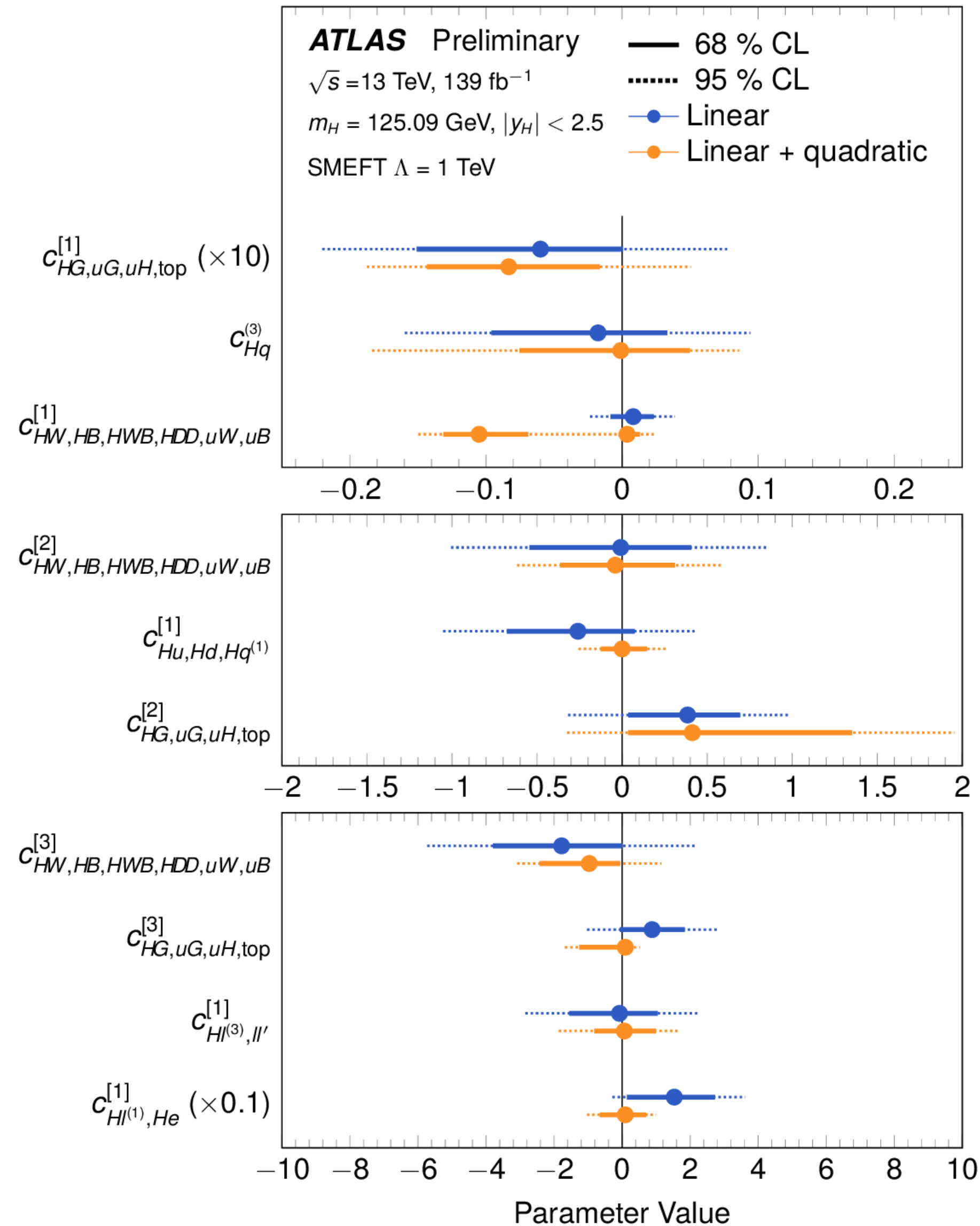
# Generic parametrization using ratios

- Express coupling-strength ( $\kappa_t, \kappa_b, \kappa_\tau, \kappa_W, \kappa_Z$ ) and effective loop ( $\kappa_g, \kappa_\gamma$ ) scale factors as ratios to  $gg \rightarrow H \rightarrow ZZ^*$ 
  - Ratios are independent of assumptions on total width
- $\lambda_{WZ}$  should be 1
  - required by SU(2) custodial symmetry and
  - $\rho$  measurements at LEP and Tevatron
- $\lambda_{\gamma Z}$  sensitive to charge particle contribution in loops
- $\lambda_{tg}$  sensitive to new coloured particles in ggF loop

Parameter	Definition in terms of $\kappa$ modifiers	Result
$\kappa_{gZ}$	$\kappa_g \kappa_Z / \kappa_H$	$0.98 \pm 0.05$
$\lambda_{tg}$	$\kappa_t / \kappa_g$	$1.04 \pm 0.12$
$\lambda_{Zg}$	$\kappa_Z / \kappa_g$	$1.06^{+0.12}_{-0.11}$
$\lambda_{WZ}$	$\kappa_W / \kappa_Z$	$1.04^{+0.08}_{-0.07}$
$\lambda_{\gamma Z}$	$\kappa_\gamma / \kappa_Z$	$1.04^{+0.07}_{-0.06}$
$\lambda_{\tau Z}$	$\kappa_\tau / \kappa_Z$	$1.04 \pm 0.13$
$\lambda_{bZ}$	$\kappa_b / \kappa_Z$	$0.96^{+0.12}_{-0.11}$



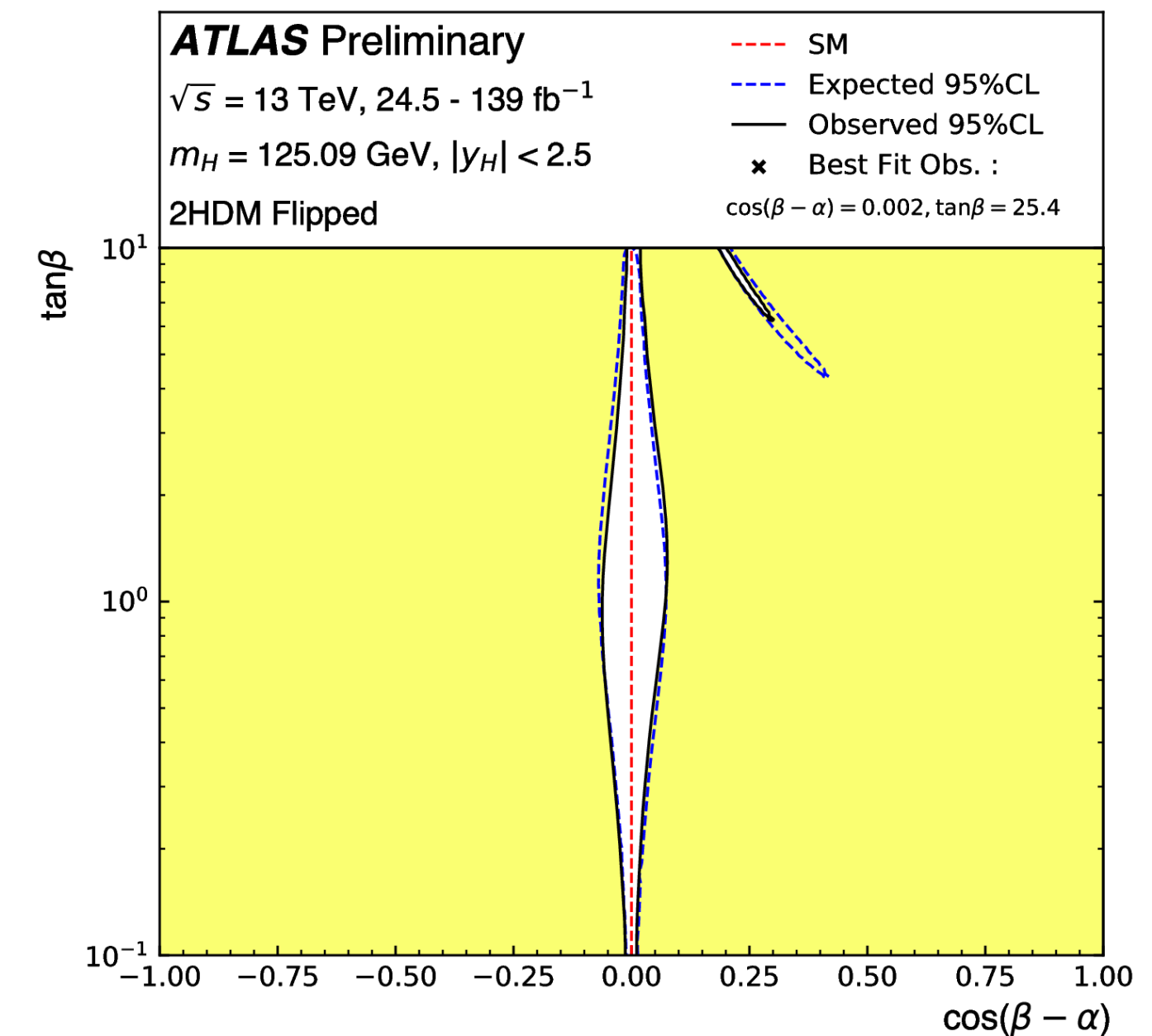
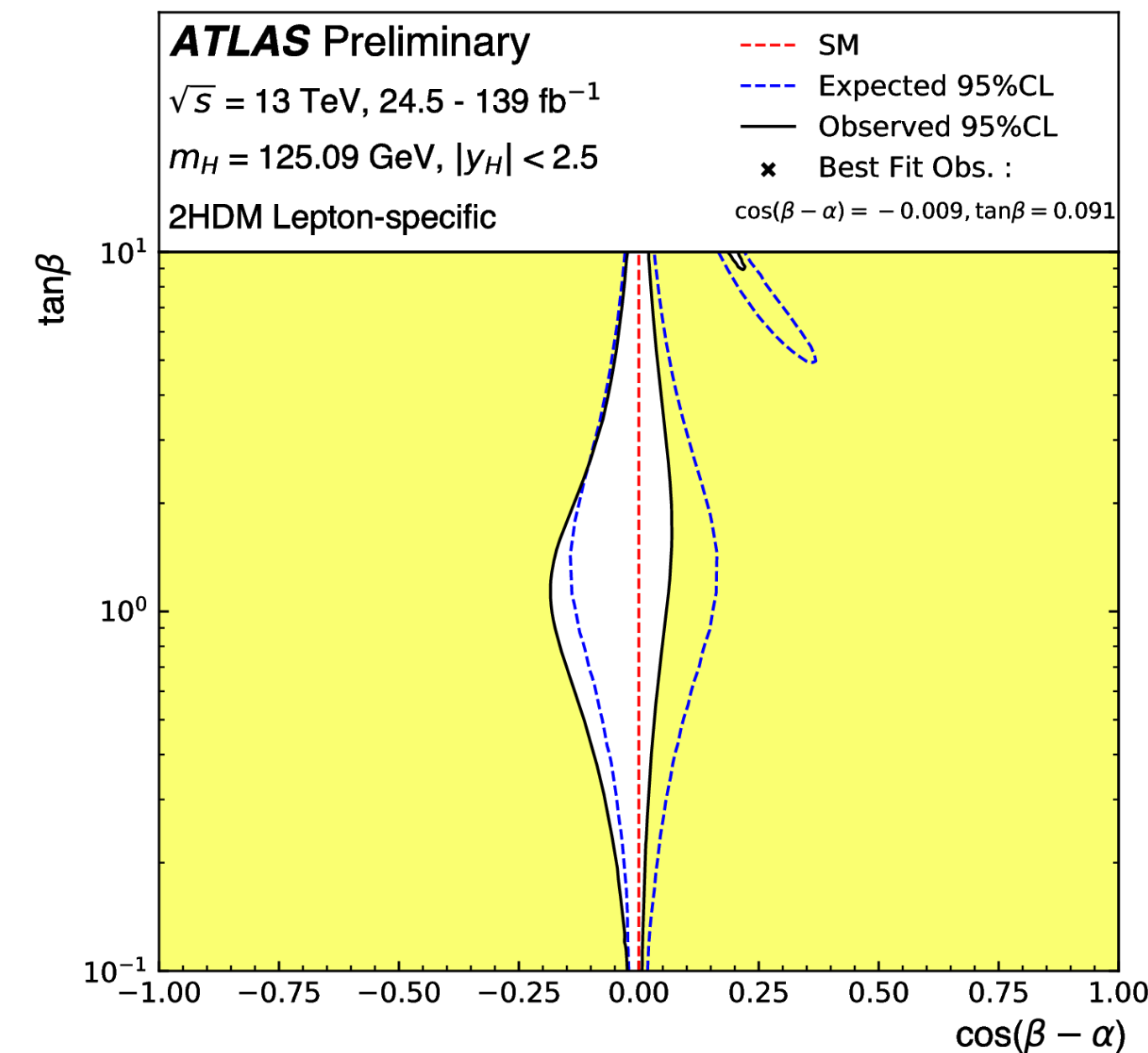
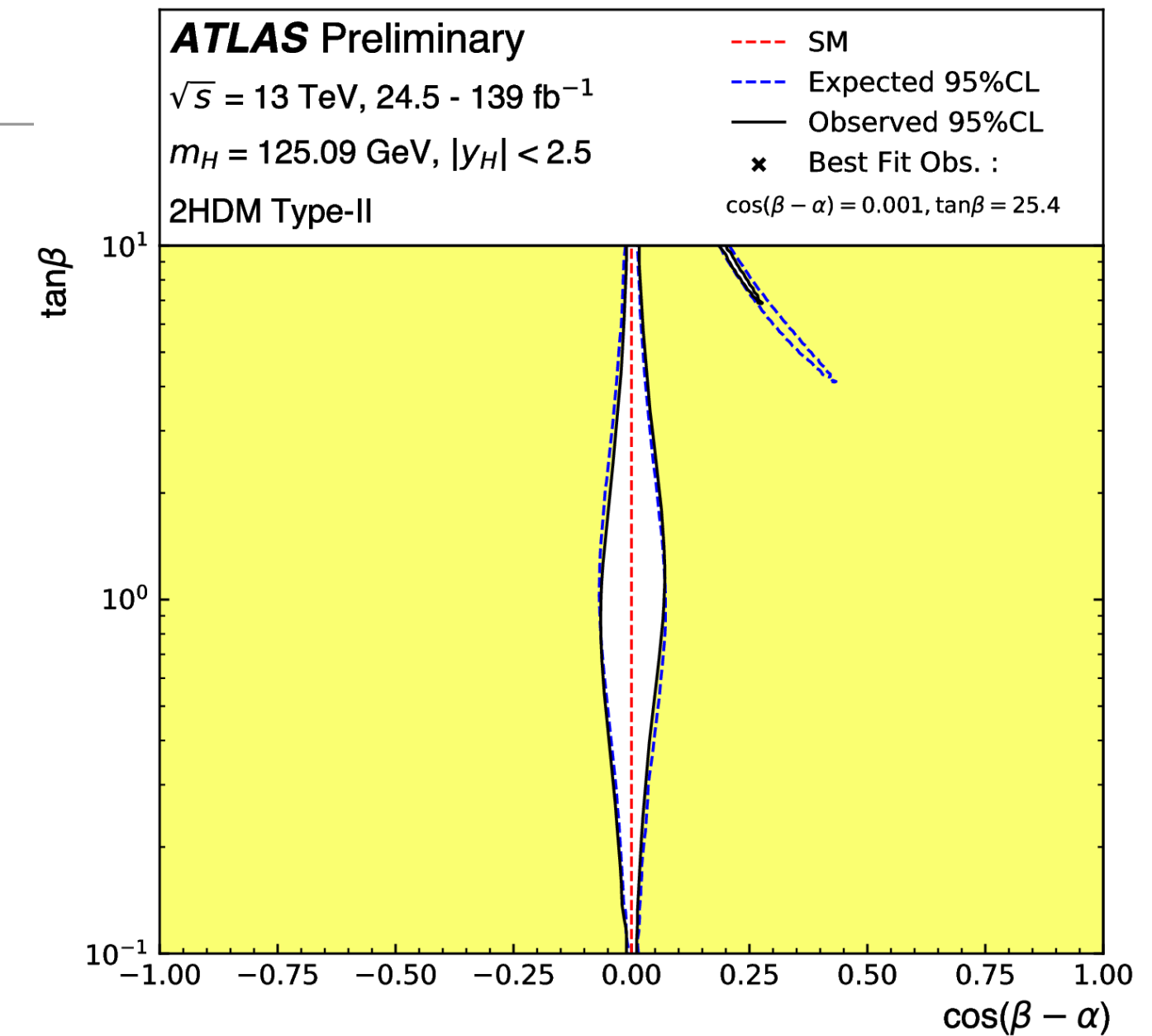
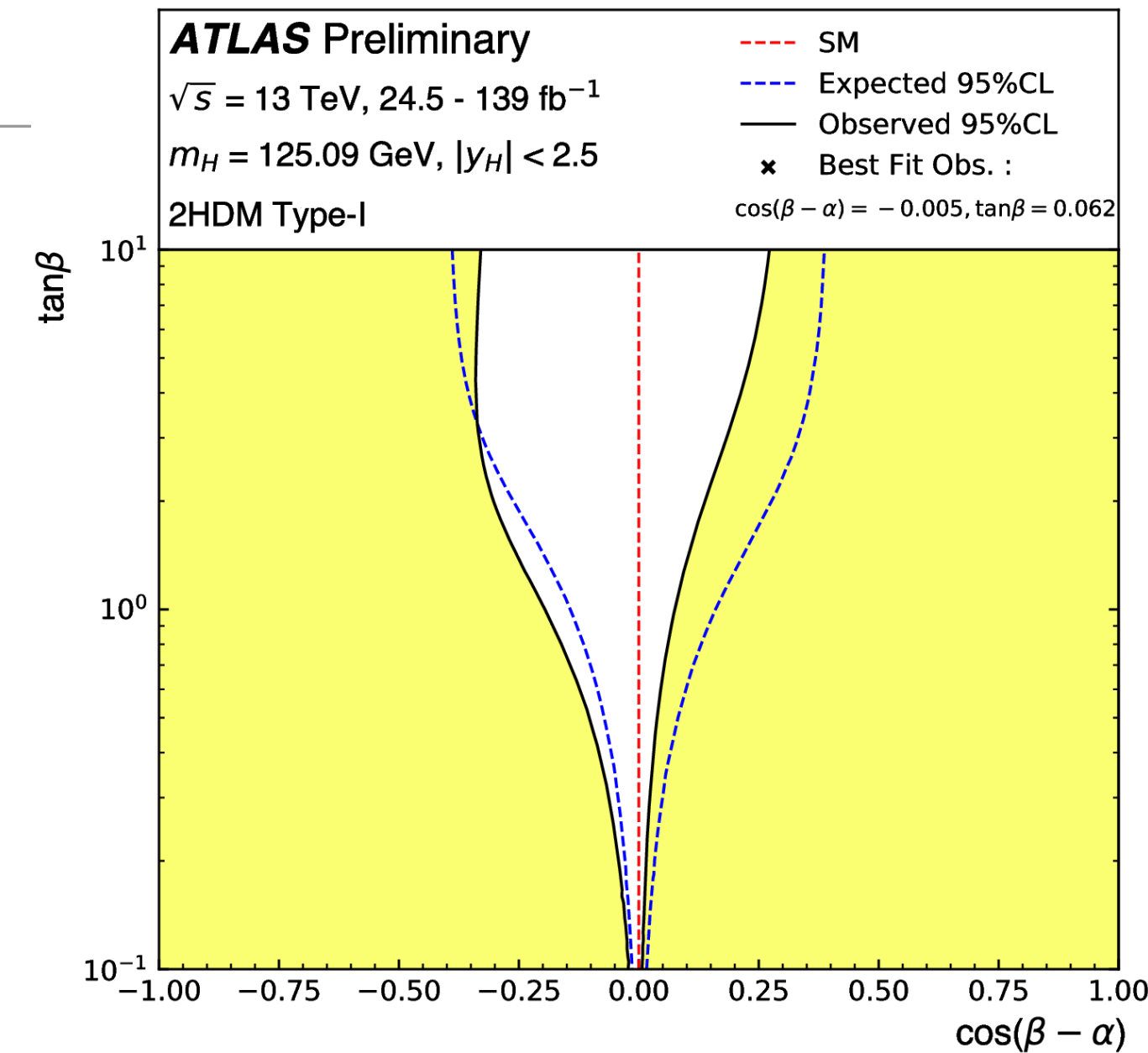
# SMEFT Correlation matrix



# 2HDM

- Can interpret results in 2HDM
  - Reparametrization of  $\kappa_Z$ ,  $\kappa_b$ ,  $\kappa_W$ ,  $\kappa_t$ ,  $\kappa_\tau$ , and  $\kappa_\mu$
- Use coupling to each SM particle, assume no BSM contribution to total width

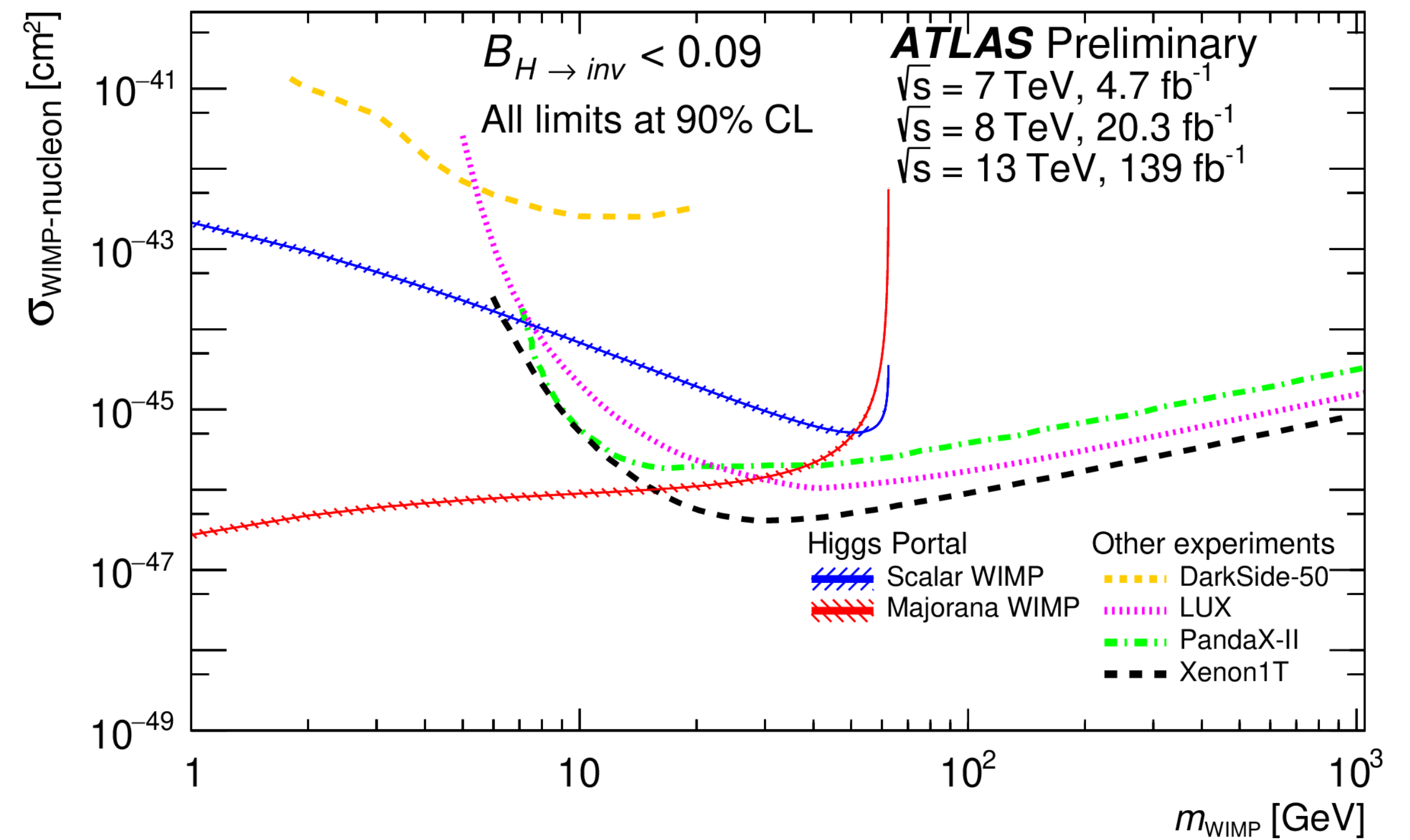
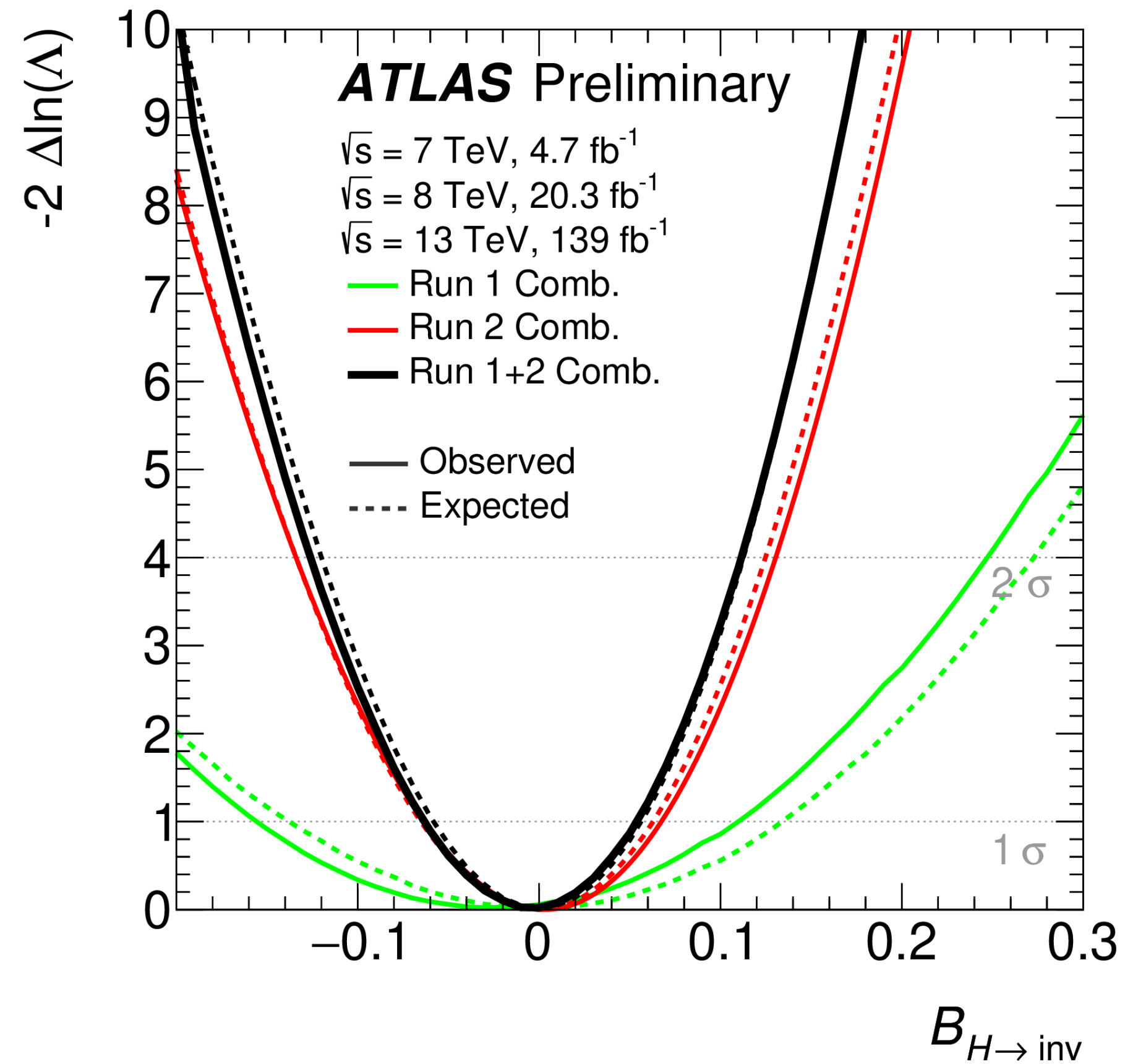
- Type I: One Higgs doublet couples to vector bosons, other couples to fermions
- Type II: One Higgs doublet couples to up-type quarks and other to down-type quarks and charged leptons
- Lepton-specific: The Higgs bosons have the same coupling to quarks as in Type I and to charged leptons as in Type II
- Flipped: The Higgs bosons have the same couplings to quarks as in Type II model and to charged leptons as in Type I





# B-Inv

- Provides a complementary search space to direct-search experiments



# SM EFT diagrams

Coefficient	Operator	Example process
$c_{HDD}$	$(H^\dagger D^\mu H)^* (H^\dagger D_\mu H)$	
$c_{HG}$	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	
$c_{HB}$	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	
$c_{HW}$	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	
$c_{HWB}$	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	
$c_{eH}$	$(H^\dagger H)(\bar{l}_p e_r H)$	
$c_{Hl}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$	
$c_{Hl}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$	
$c_{He}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{e}_p \gamma^\mu e_r)$	
$c_{Hq}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$	
$c_{Hq}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{q}_p \tau^I \gamma^\mu q_r)$	
$c_{Hu}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{u}_p \gamma^\mu u_r)$	
$c_{Hd}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$	

- Most relevant EFT operators affecting Higgs boson production and decay in the considered phase space with examples

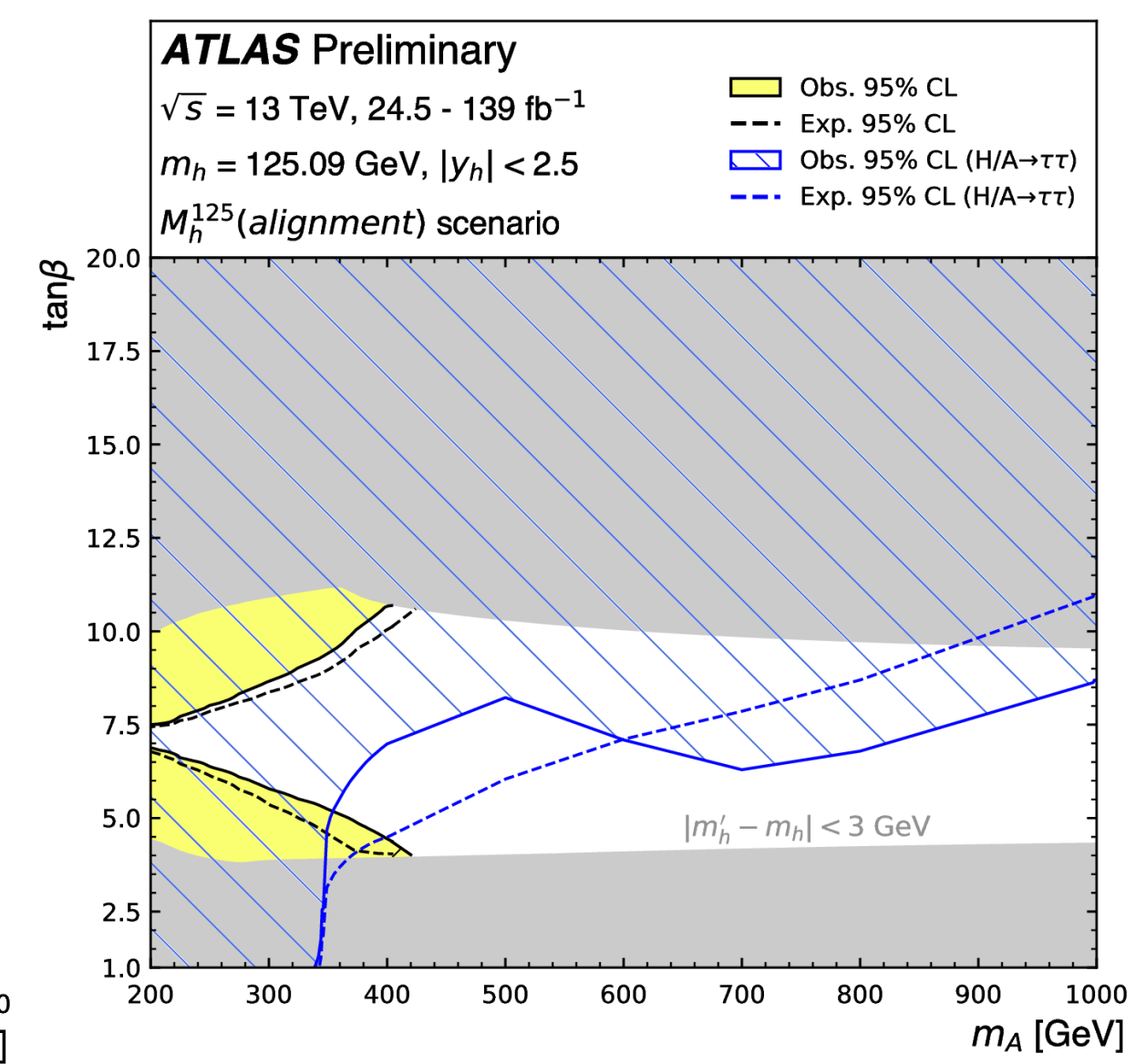
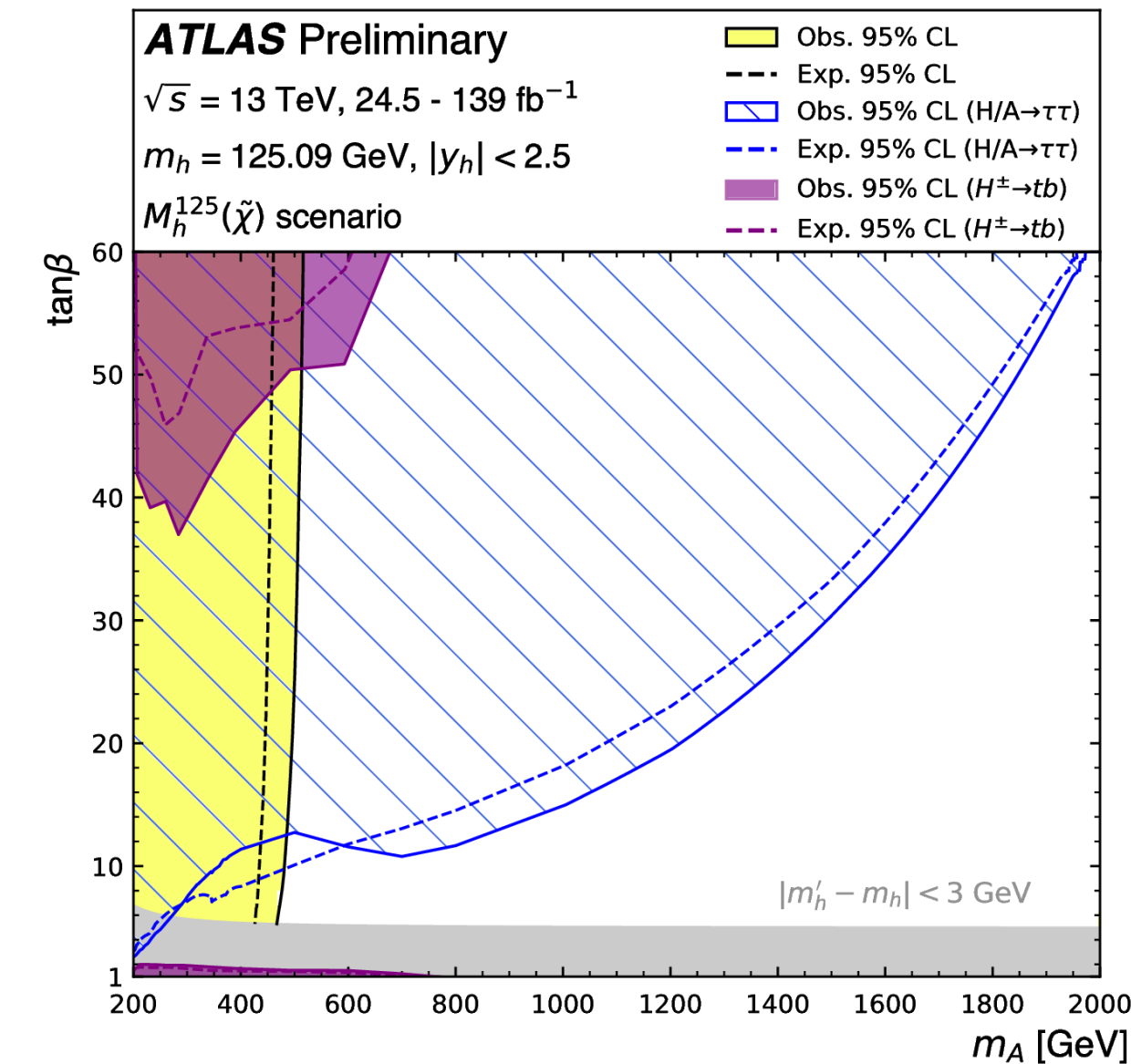
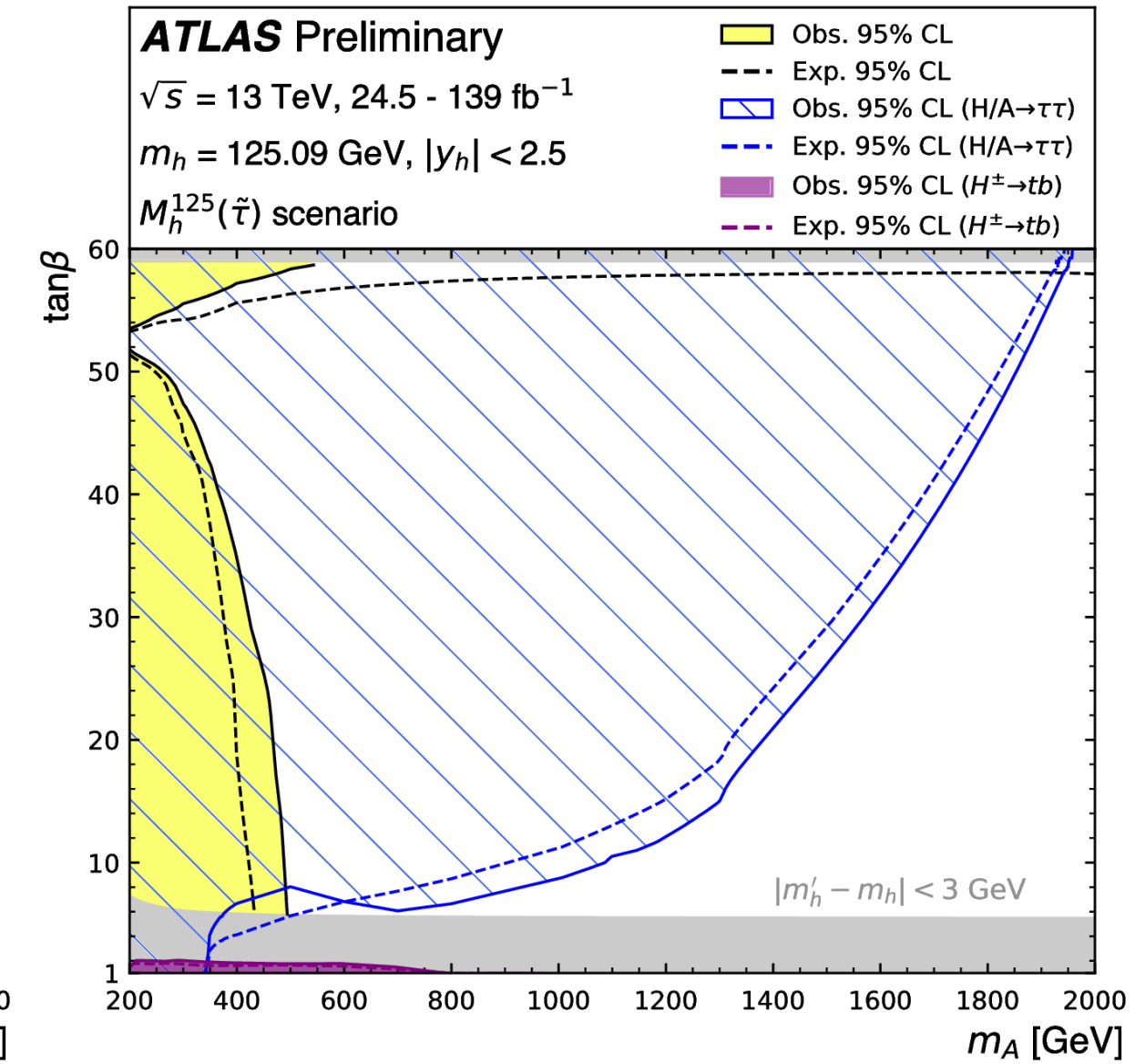
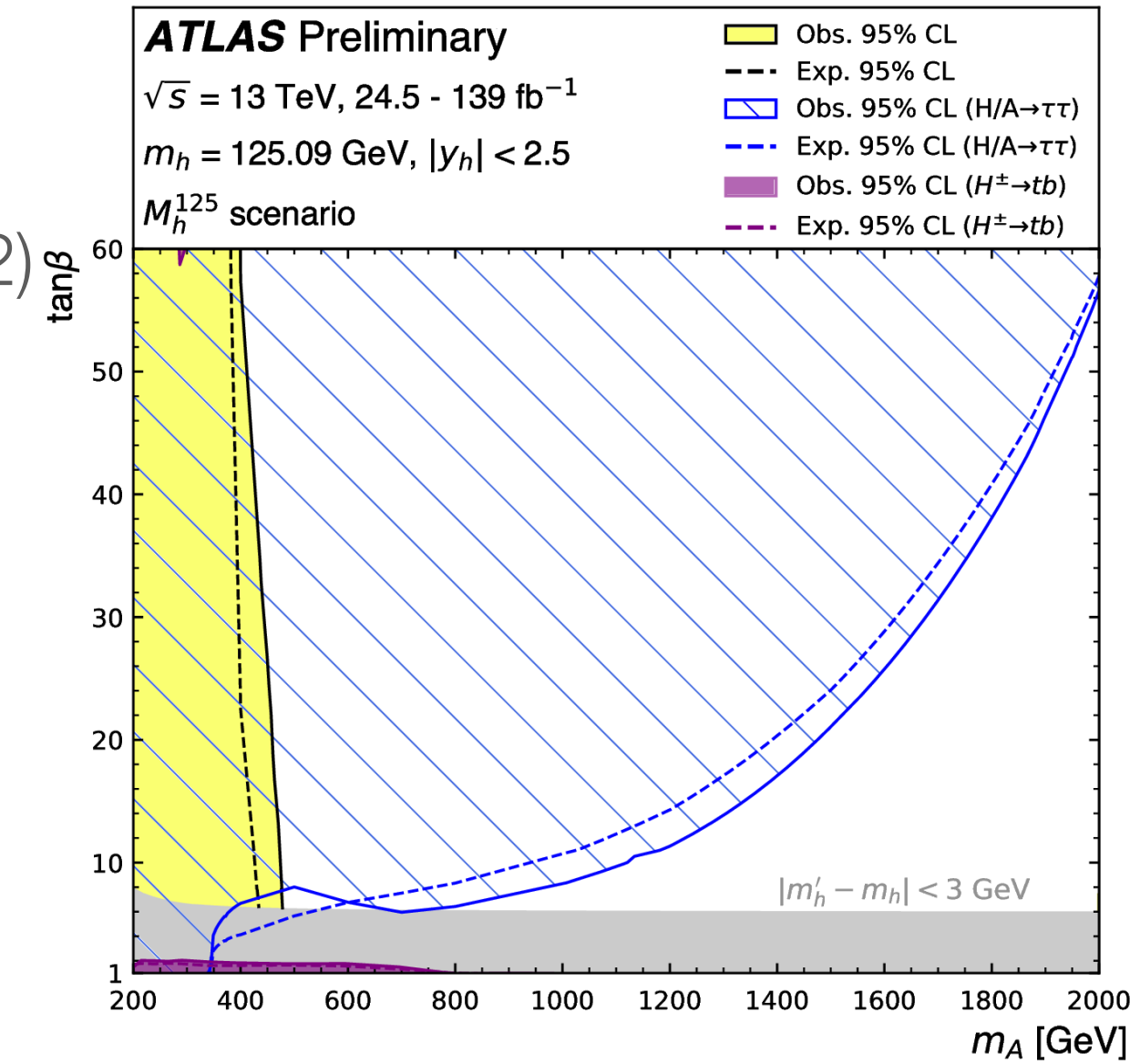
Coefficient	Operator	Example process
$c_{uG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{H} G_{\mu\nu}^A$	
$c_{uW}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$	
$c_{uB}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$	
$c_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{q}_r \gamma^\mu q_s)$	
$c_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	
$c_{qq}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{q}_r \gamma^\mu q_s)$	
$c_{qq}^{(31)}$	$(\bar{q}_p \gamma_\mu \tau^I q_t)(\bar{q}_r \gamma^\mu \tau^I q_s)$	
$c_{uu}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	
$c_{uu}^{(1)}$	$(\bar{u}_p \gamma_\mu u_t)(\bar{u}_r \gamma^\mu u_s)$	
$c_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{u}_r \gamma^\mu u_s)$	
$c_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	
$c_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$	
$c_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$	
$c_G$	$f^{ABC} G_\mu^A G_\nu^B G_\rho^C$	



# Minimal supersymmetric extension of SM (MSSM)

- Supersymmetry is theoretically motivated to address unanswered questions in the SM
- Minimal SUSY (MSSM) predicts partners for SM states and two SU(2) doubles which both acquire a VEV,  $v_1$  and  $v_2$  with ratio  $\beta = v_2/v_1$
- Over 100 parameters in MSSM but under certain assumptions can reduce the parameter space
- Observed Higgs boson in light CP-even Higgs boson  $h$  of MSSM

- $M_h^{125}$ : Superparticles are so heavy that production/decay of MSSM Higgs bosons are only mildly affected
- $M_h^{125}(\tilde{\tau})$ : Light staus and gaugino-like charginos and neutrinos
- $M_h^{125}(\tilde{\chi})$ : Chargino & neutralinos relatively light with significant higgsino-gaugino mixing
- $M_h^{125}(\text{alignment})$ : For a given  $\tan(\beta)$ , one of the two neutral CP-even scalars has SM-like couplings independently of the mass spectrum of the remaining Higgs bosons



# Minimal supersymmetric extension of SM (MSSM)

5.  $M_{h,EFT}^{125}$  scenario: characterized by a flexible mass scale of the superpartners. In 1-4 sfermions are tried to TeV scale. To reopen  $\tan(\beta) < 5$  the sfermion mass scale is adjusted dynamically to achieve a 125 GeV Higgs boson. All super particles chosen to be heavy.
6.  $M_{h,EFT}^{125}(\tilde{\chi})$  scenario: Similar to  $M_{h,EFT}^{125}$  but features light neutralinos and charginos that significantly alter phenomenology of Higgs boson. SUSY scale adjusted at every parameter point to achieve 125 GeV Higgs

