

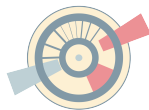
Event-dependent jet vetoes: a new prospect for multilepton searches ¹

DIS 2021, SUNY Stony Brook, Zoom Campus

Richard Ruiz

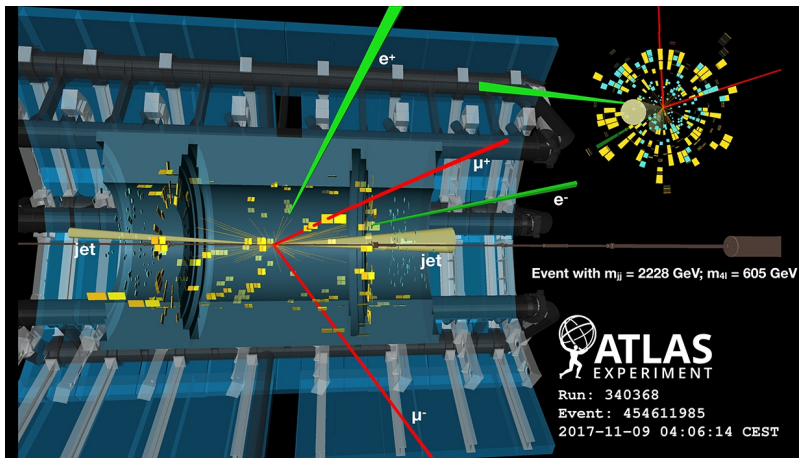
Institute of Nuclear Physics – Polish Academy of Science (IFJ PAN)

14 April 2021



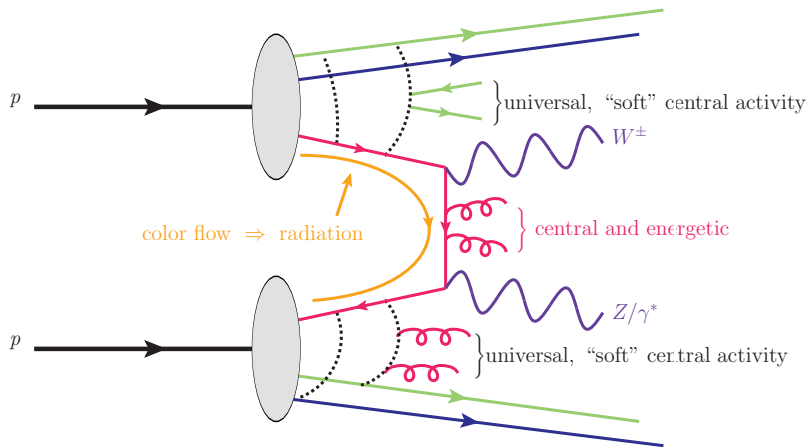
¹ Pascoli, RR, Weiland [PLB ('18), 1805.09335], JHEP ('19), [1812.08750]; w/ Fuks, et al. PRD ('19) [1901.09937] =

What do high- p_T LHC collisions look like?



In high-momentum transfer (Q^2) events...

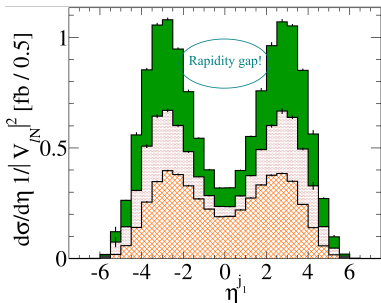
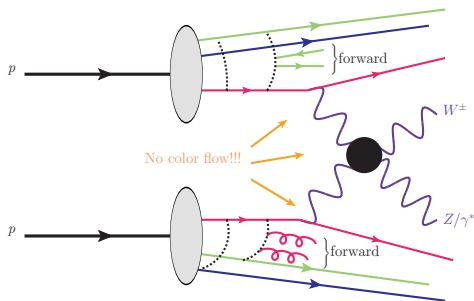
flow of QCD color charges is *everywhere!*



From (classical) electrodynamics, accelerated dipoles $\Rightarrow$ real radiation
- **Notably**, while radiation existence is universal, emission pattern is not

In color-singlet processes, there is an absence of central color flow
 $\Rightarrow$ absence of central, high- p_T jets

[“Rapidity gaps” Dokshitzer, Khoze, Troyan ('86)]



Basis for Central Jet Veto

[Barger, et al, PRD ('91) + PLB ('95); Bjorken, PRD ('94)]:

- Reject events with any jet satisfying $p_T^j > 25 - 30$ GeV, $|\eta^j| < 2 - 3$
 - After applying meaningful jet PID, equivalent to keeping only 0-jet bin
- CJVs are crucial to precision Higgs and EW studies
 - Discrepancies between SM and Run I data inspired an entire community

Revolution in Jet Veto Resummation

$$d\sigma(X + 0j) = \underbrace{f \otimes f \otimes \Delta}_{\text{LL, NLL, NNLL, etc.}} \otimes \underbrace{d\hat{\sigma}}_{\text{LO, NLO, etc.}} + \underbrace{\text{(matching scheme)}}_{\text{removes dbl counting}}$$

- Jet veto factorization thms

[Stewart, et al ('09); Becher, et al ('12); etc.]

- $pp \rightarrow h/Z + 0j$

[Becher, et al ('12,'13); Banfi, et al ('12); Stewart, et al ('13)]

- $pp \rightarrow WW + 0j$ [Meade, et al ('14); Jaiswal, et al

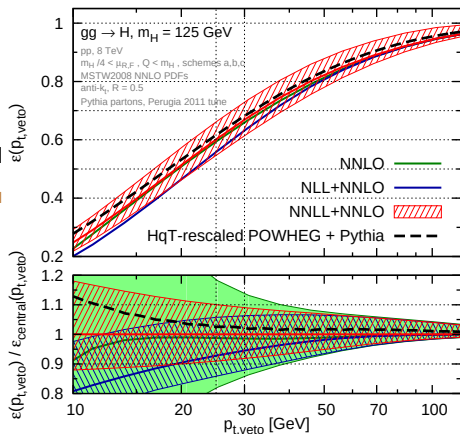
('14); Monni, et al ('14); Becher, et al ('14)]

- Analytic resummation vs (N)NLO vs NLO+PS(LL)

[Monni, et al ('14); Fuk, RR, et al ('17,'18,'19); Pascoli, RR, et al ('18); Jager, et al ('18)]

- Rapidity window-dependence

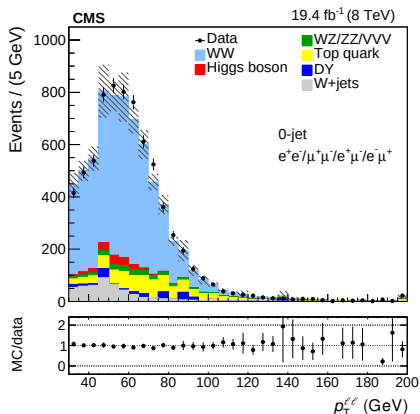
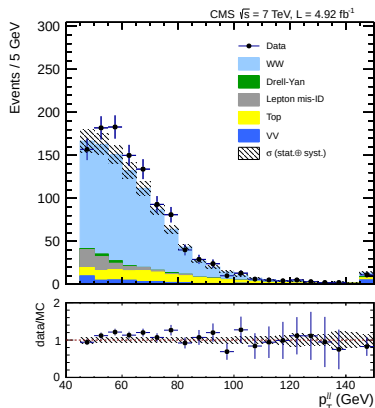
[Gangal, et al ('14,'17); Michel, et al ('18)]



Banfi, et al [[1203.5773](#), [1206.4998](#)]

Standard Model vs Jet Vetoes

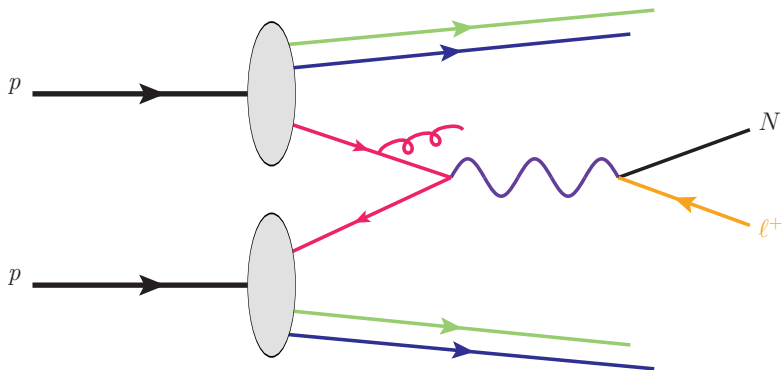
Plotted: p_T of the $(\ell_1^+ \ell_2^-)$ -system in $pp \rightarrow W^+ W^- + 0j \rightarrow \ell_1^+ \ell_2^- \nu \bar{\nu} + 0j$



Summary: Past discrepancies resolved by accounting for (resumming) **ultra soft** and **ultra forward** gluon emission

jet vetoes in searches for new physics

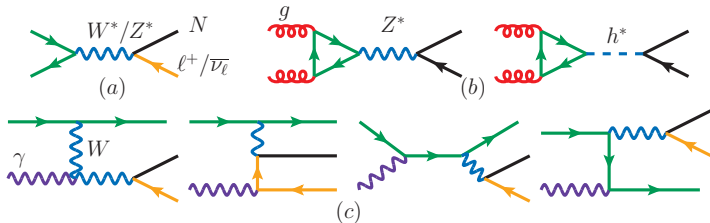
Case Study: Heavy Neutrinos at the LHC²



²Pascoli, RR, Weiland [PLB ('18), 1805.09335], JHEP ('19), 1812.08750]

Heavy N can be produced in a variety of ways in pp collisions

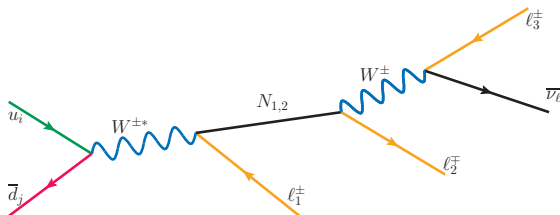
Review on ν mass models at colliders, Y. Cai, T. Li, T. Han, RR [1711.02180]



DY + $W\gamma$ fusion dominant mechanisms for heavy N at LHC

Alva, Han, RR [1411.7305]; Degrande, Mattelaer, RR, Turner [1602.06957]

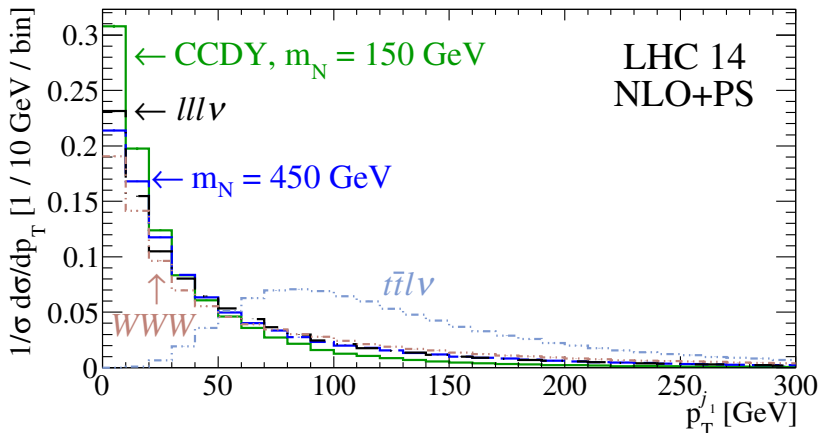
Focus on the final state: $pp \rightarrow 3\ell$



Jet vetoes tries to discriminate jet activity in signal and background.

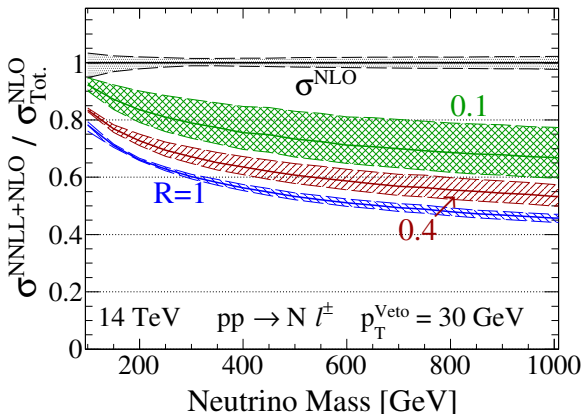
Consider $pp \rightarrow 3e/\mu + X$ at NLO+PS(LL)+anti- k_T ($R = 1$)

Plotted: $p_T^{j_1}$ for CC DY signal + SM Bkg



Separation according to Born color structure clear

Plotted: veto efficiency³ $\varepsilon = \sigma^{\text{NLO+NNLL}}(p_T^{\text{Veto}}) / \sigma_{\text{Tot.}}^{\text{NLO}}$ for $pp \rightarrow W'_{\text{SSM}}$



($R=\text{jet size}$)

- For fixed p_T^{Veto} threshold, signal survival crashes for increasing mass!⁴
- **Question:** Is a veto even useful when $p_T^{\text{Veto}} > 30 - 40$ GeV?

³ a la automated resummation with mg5amc@nlo+scet by Becher, et al [1412.8408] (more on this later!)

⁴ Ditto for sleptons, Tackmann, et al [1603.03052]; and W'/Z' , Fuks, RR [1701.05263]

Reenvisioning Jet Vetoes:

Event-Based⁵ Jet Vetoes⁶

⁵aka: “dynamic jet veto,” “phase space-dependent jet veto”, or “safe jet veto”

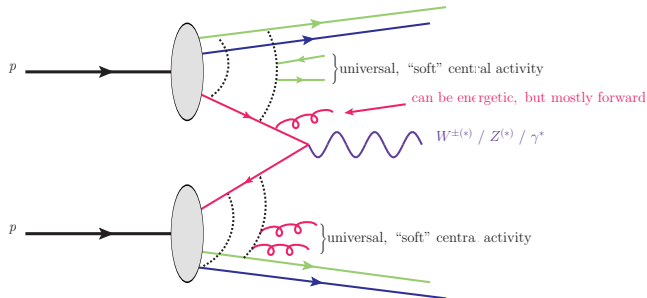
⁶**Early literature:** Bjorken (dynamic rapidity gap) [PRD ('93)]; Denner, et al (regulator trick) [0906.1656]; Companario, et al ($x = E_T / \sum_k E_T^k$ fraction) [1410.4840] ; **more recent work** by DESY+NIKEF+Mainz (rapidity-dependent vetoes) [1506.0544]

an idea:⁷: on an event-by-event basis,
allow an event's kinematics to set p_T^{Veto}

⁷ Pascoli, RR, Weiland [PLB ('18), [1805.09335]

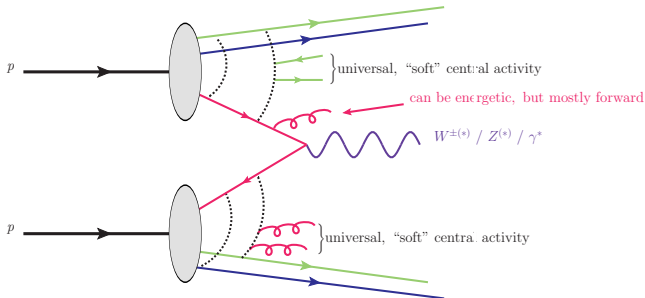
Leptons in High-Mass Drell-Yan and VBS

For leptons in high-mass DY: $p_T^{\ell_k}, S_T \sim M_{V^*} \gg p_T^j \sim \text{Sudakov shoulder}$



Leptons in High-Mass Drell-Yan and VBS

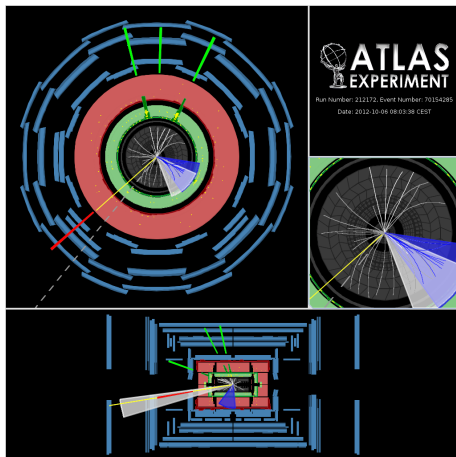
For leptons in high-mass DY: $p_T^{\ell_k}, S_T \sim M_{V^*} \gg p_T^j \sim \text{Sudakov shoulder}$



An idea: on event-by-event basis, set $p_T^{\text{Veto}} = p_T^{\ell_1}$

- High-mass, DY- and VBS-type processes pass by construction
- Since $(M_{V^*}/p_T^{\text{Veto}}) \sim 1$, jet veto logarithms not inherently large
- What about background processes, e.g., top quarks?

Top Quark Background vs Event-Based Jet Vetoes



$pp \rightarrow t\bar{t}Z \rightarrow 1\mu + 3e + 2j_b + \cancel{E}_T$
 candidate event [1509.05276]

Textbook kinematics:

- $m_{ee} = 93 \text{ GeV}$
- $\cancel{E}_T = 57 \text{ GeV}$

Typically,

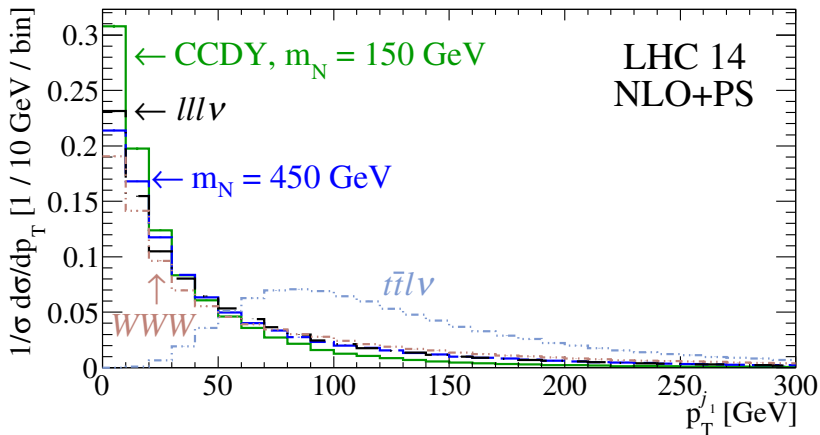
- $p_T^{e3} \sim \frac{M_Z}{2} \sim 45 \text{ GeV}$
 - $p_T^{e1} \sim \frac{m_t}{4} \left(1 + \frac{M_W^2}{m_t^2}\right) \sim 50 \text{ GeV}$
 - $p_T^{b1} \sim \frac{m_t}{2} \left(1 - \frac{M_W^2}{m_t^2}\right) \sim 60 \text{ GeV}$
- $p_T^{b1} > p_T^{\ell_k} \implies \text{event vetoed!}$

Setting $p_T^{\text{Veto}} = p_T^{\ell_1}$ can **reduce** top background **without** b -jet tagging!
 - Diboson and triboson backgrounds also reduced! (see backup)

Jet vetoes tries to discriminate jet activity in signal and background.

Consider $pp \rightarrow 3e/\mu + X$ at NLO+PS(LL)+anti- k_T ($R = 1$)

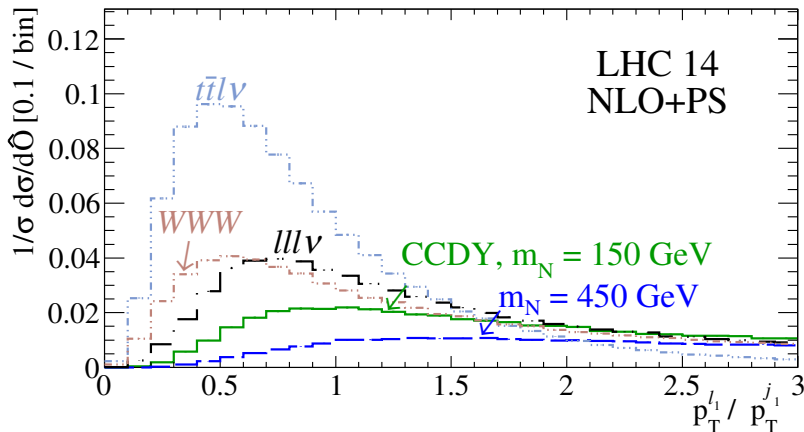
Plotted: $p_T^{j_1}$ for CC DY signal + SM Bkg



Separation according to **Born color structure** clear

Event-based jet vetoes discriminate against hadronic and leptonic activities

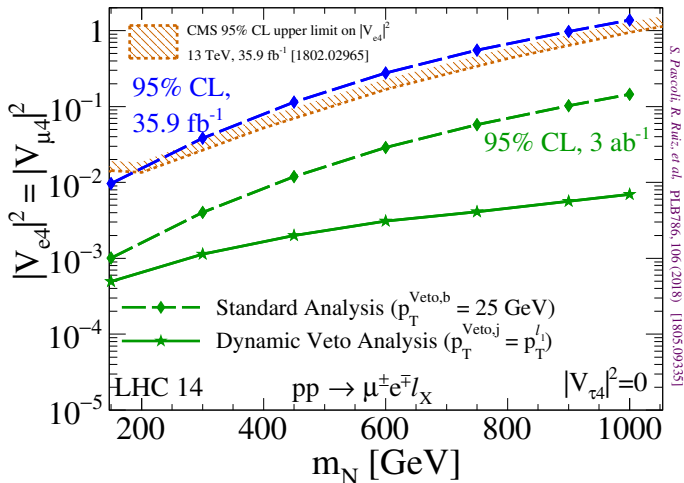
Plotted: $r_{j1}^{\ell_1} = p_T^{\ell_1} / p_T^{j_1}$ for CC DY signal + SM Bkg



Similar behavior for other ratios, e.g., S_T/H_T , $p_T^{\ell_k}/H_T$

- Not universal, e.g. degenerate mass limit in $\tilde{\mu} \rightarrow \mu \tilde{\chi}^0$ [1901.09937]

Plot: LHC 14 sensitivity to **active-sterile neutrino mixing** (coupling²) vs **heavy N mass** (m_N), in search for $pp \rightarrow \mu^\pm e^\mp \ell_X + X$ ($\ell_X = e, \mu, \tau_h$)



Improved sensitivity up to $10 - 11\times$ with $\mathcal{L} = 3 \text{ ab}^{-1}$.

See also [1812.08750] for various lepton flavor permutations, uncertainty plots, etc

Summary

- ① The science of jet vetoes is very mature
 - ▶ Precision sufficient for both indirect and direct BSM searches
 - ▶ Uncertainties extracted at NLO+PS(LL) may be large but reliable (consistent with higher accuracy results)
 - ▶ High precision resummation necessary if $p_T^{\text{Veto}} \ll Q$
- ② Building jet veto thresholds on event-by-event basis, e.g., $p_T^{\text{Veto}} = S_T$, appears very promising:
 - ▶ New schemes reveals $> 90 - 95\%$ veto acceptance for signal with little-to-no dependence on BSM mass scales [1805.09335, 1901.09937]
 - ▶ Irreducible backgrounds, e.g., $t\bar{t}$ and multi-boson particularly impacted
 - ▶ Setting $p_T^{\text{Veto}} \sim Q \implies$ milder jet veto logs $\implies$ less need for high-precision resummation [1812.08750]
- ③ No single veto scheme is optimal for all analyses (have fun, explore!)
 - ▶ Dedicated and general-purpose MC tools are available and in dev.
 - ▶ BSM UFO model files with ingredients for NLO in QCD available feynrules.irmp.ucl.ac.be/wiki/NLOModels

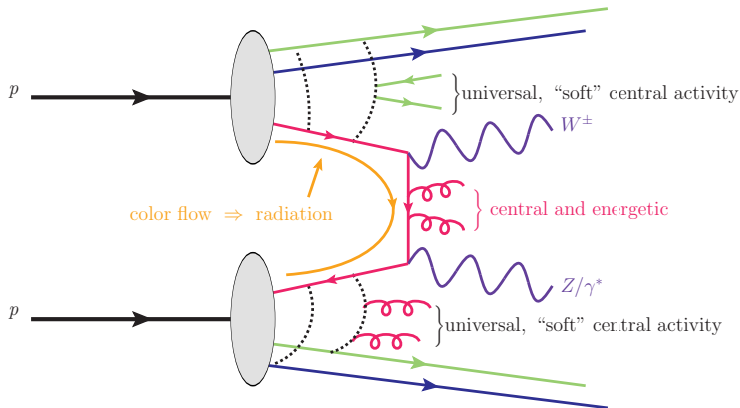


Thank you.

Diboson and Triboson Processes

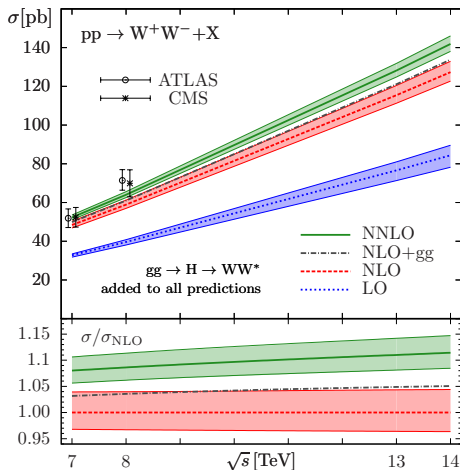
Di- / triboson processes contain large amounts of hadronic activity

- Significant s - and t -channel cancelations for Born-like configurations⁸
- Nontrivial color flow, despite naïve, color-singlet nature



⁸ i.e., Radiation zeros (super interesting!) Mikaelian ('78); +Brown, et al ('78-'79); Zhu ('80); Brodsky, et al ('82,'83);

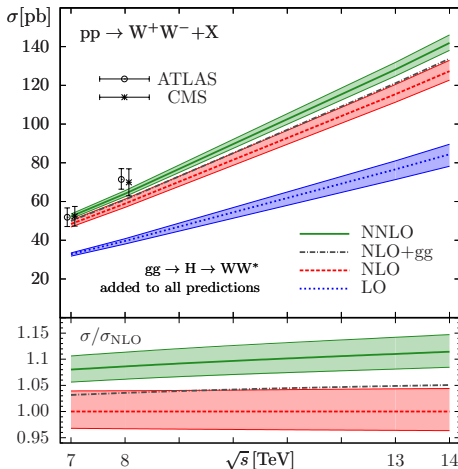
Irreducible Bkgs vs Event-Based Jet Vetoes



Inclusive diboson production driven by $pp \rightarrow VV'j + X$

- Manifests as large K -factor in $pp \rightarrow W^+W^-$ at NNLO [1408.5243]
- $p_T^\ell \sim \frac{M_V}{2} + \frac{p_T^j}{2n_b}$ ($n_b = \#$ bosons)
- $\Rightarrow$ large $p_T^j > p_T^\ell$ tail

Irreducible Bkgs vs Event-Based Jet Vetoes



Inclusive diboson production driven by $pp \rightarrow VV'j + X$

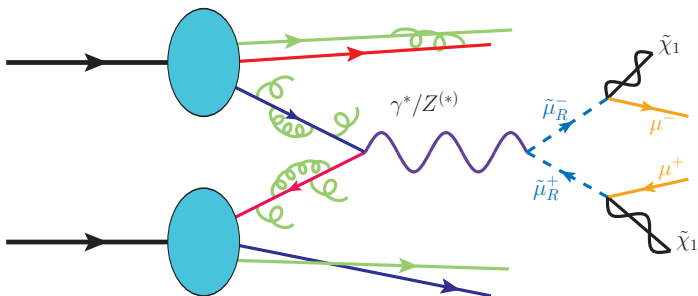
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- $\Rightarrow$ large $p_T^j > p_T^\ell$ tail

Jets mistagged/IDed as τ_h/e , i.e., “fake leptons,” major irreducible bkg:

- E.g., $pp \rightarrow W^+W^-j$
- Subleading jets more likely to be mis-IDed than leading jet

Setting $p_T^{\text{Veto}} = p_T^{\ell_1}$ can **reduce** irreducible backgrounds

Case Study II: sleptons and generalizing event-based jet veto definitions⁹



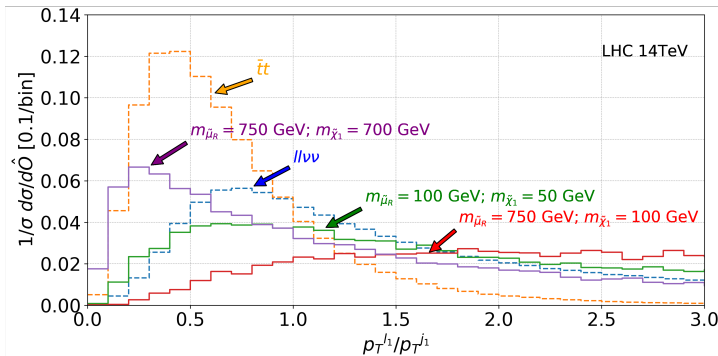
⁹ Fuks, Nordstrom, RR, Williamson [1901.09937]

Leptons vs Hadrons Redux

Consider the signature $pp \rightarrow \mu^+ \mu^- + X$ at NLO+PS(LL)

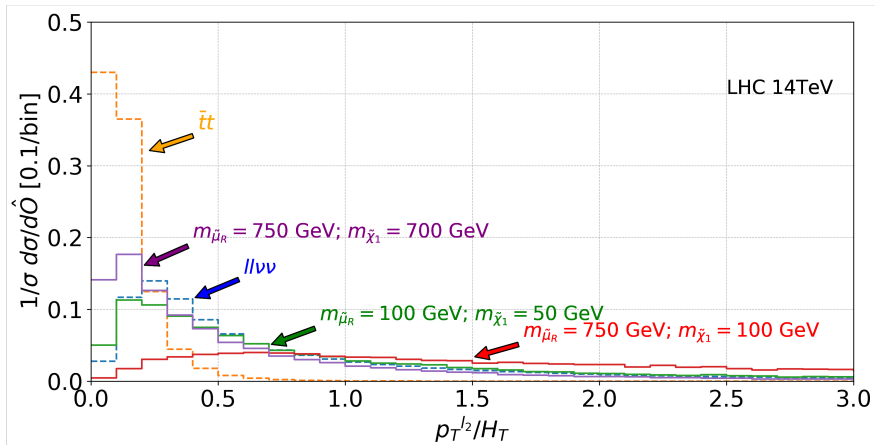
- Generator cuts + $p_T^\ell > 15$ GeV + $|\eta^\ell| < 2.4$ + anti- k_T ($R = 1$)

Plotted: $r_{j1}^{\ell_1} = p_T^{\ell_1} / p_T^{j1}$



Poorer S/B separation than heavy neutrino case, esp. for $m_{\tilde{\mu}} \sim m_{\tilde{\chi}}$

Plotted: $r_{H_T}^{\ell_2} = p_T^{\ell_2}/H_T$, $H_T = \sum_{k \in \{clusters\}} E_T^k$



Turns out $p_T^{\text{Veto}} = p_T^{\ell_1}$ is not best ratio :)

Nevertheless, separation between $t\bar{t}$ from $V^{(*)}V^{(*)}$ is clear :)

Had "fun" looking into uncertainties :)

Cut / Channel	$\sigma(\ell\ell\nu\nu)$ [fb]			$\sigma(\ell\ell\ell\nu)$ [fb]			$\sigma(t\bar{t})$ [fb]		
	Incl.	Incl.+MPI	FxFx+MPI	Incl.	Incl.+MPI	FxFx+MPI	Incl.	Incl.+MPI	FxFx+MPI
Common Analysis Requirements									
Generator	10200 ^{+4.1%} _{-4.8%}	10200 ^{+4.1%} _{-4.8%}	10300 ^{+4.5%} _{-5.1%}	1600 ^{+5.5%} _{-6.5%}	1600 ^{+5.5%} _{-6.5%}	1680 ^{+5.5%} _{-6.2%}	85800 ^{+9.2%} _{-10.4%}	85800 ^{+9.2%} _{-10.4%}	91000 ^{+11.9%} _{-11.7%}
Dimuon Selection	850 ^{+4.1%} _(8.4%)	850 ^{+4.5%} _(8.4%)	840 ^{+4.9%} _(8.4%)	140 ^{+6.3%} _(8.7%)	140 ^{+6.3%} _(8.7%)	140 ^{+6.1%} _(8.6%)	6400 ^{+12%} _(7.5%)	6400 ^{+12%} _(7.5%)	6800 ^{+12%} _(7.5%)
+ $m_{\ell\ell}$ Requirements	590 ^{+4.2%} _(69%)	590 ^{+4.5%} _(69%)	580 ^{+5.3%} _(69%)	40 ^{+7.0%} _(28%)	40 ^{+6.9%} _(28%)	40 ^{+6.3%} _(29%)	5000 ^{+12%} _(77%)	5000 ^{+12%} _(77%)	5300 ^{+12%} _(78%)
+Minimum M_{T2}	2.9 ^{+6.6%} _(0.49%)	3.0 ^{+4.6%} _(0.51%)	2.4 ^{+9.0%} _(0.42%)	0.74 ^{+9.6%} _(1.9%)	0.76 ^{+7.2%} _(1.9%)	0.56 ^{+7.0%} _(1.4%)	8.8 ^{+13%} _(0.18%)	9.2 ^{+13%} _(0.18%)	8.6 ^{+12%} _(0.16%)
Benchmark Static Jet Veto Analysis Requirements									
+ $p_T^{\ell_1} > 50$ GeV, $p_T^{\ell_2} > 20$ GeV	2.7 ^{+9.0%} _(94%)	2.8 ^{+4.6%} _(94%)	2.3 ^{+12%} _(96%)	0.70 ^{+12%} _(95%)	0.71 ^{+7.4%} _(94%)	0.53 ^{+7.8%} _(95%)	7.7 ^{+13%} _(87%)	8.3 ^{+15%} _(90%)	7.8 ^{+12%} _(91%)
+Static Jet Veto	1.5 ^{+13%} _(54%)	1.5 ^{+5.6%} _(53%)	1.6 ^{+13%} _(70%)	0.27 ^{+11%} _(39%)	0.26 ^{+9.3%} _(36%)	0.31 ^{+7.8%} _(60%)	0.033 ^{+240%} _(0.43%)	0.049 ^{+66%} _(0.60%)	0.15 ^{+68%} _(1.9%)
Dynamic Jet Veto Analysis Requirements									
$p_T^{\text{Veto}} = p_T^{\ell_1}$	2.6 ^{+6.7%} _(86%)	2.5 ^{+7.0%} _(85%)	2.4 ^{+10%} _(77%)	0.58 ^{+11%} _(77%)	0.58 ^{+11%} _(77%)	0.55 ^{+7.3%} _(99%)	3.8 ^{+12%} _(43%)	3.9 ^{+12%} _(43%)	4.5 ^{+32%} _(57%)
$p_T^{\text{Veto}} = p_T^{\ell_2}$	2.0 ^{+4.4%} _(67%)	1.9 ^{+7.6%} _(68%)	2.3 ^{+11%} _(91%)	0.41 ^{+15%} _(55%)	0.43 ^{+9.1%} _(57%)	0.50 ^{+7.1%} _(87%)	0.84 ^{+25%} _(9.0%)	0.95 ^{+22%} _(10%)	0.98 ^{+57%} _(12%)
$p_T^{\text{Veto}} = S_T$	2.8 ^{+6.2%} _(94%)	2.7 ^{+8.2%} _(94%)	2.4 ^{+8.9%} _(90%)	0.68 ^{+8.9%} _(90%)	0.67 ^{+9.4%} _(90%)	0.56 ^{+6.9%} _(90%)	5.8 ^{+19%} _(62%)	7.1 ^{+13%} _(71%)	6.6 ^{+23%} _(89%)
$H_T^{\text{Veto}} = p_T^{\ell_1}$	2.2 ^{+7.4%} _(73%)	1.9 ^{+10%} _(67%)	2.2 ^{+9.2%} _(87%)	0.46 ^{+9.4%} _(60%)	0.41 ^{+16%} _(55%)	0.46 ^{+7.3%} _(84%)	0.40 ^{+23%} _(4.2%)	0.49 ^{+25%} _(4.9%)	0.54 ^{+22%} _(7.1%)
$H_T^{\text{Veto}} = p_T^{\ell_2}$	1.5 ^{+4.9%} _(51%)	1.5 ^{+5.3%} _(43%)	1.5 ^{+7.7%} _(60%)	0.31 ^{+8.3%} _(41%)	0.25 ^{+8.9%} _(34%)	0.33 ^{+8.6%} _(58%)	0.049 ^{+69%} _(0.53%)	0.016 ^{+200%} _(0.17%)	0.12 ^{+62%} _(1.5%)
$H_T^{\text{Veto}} = S_T$	2.5 ^{+4.6%} _(84%)	2.4 ^{+8.3%} _(82%)	2.3 ^{+9.6%} _(97%)	0.55 ^{+13%} _(74%)	0.53 ^{+11%} _(71%)	0.54 ^{+7.0%} _(95%)	1.4 ^{+23%} _(16%)	1.1 ^{+44%} _(11%)	1.9 ^{+31%} _(23%)

Table 2 The cross section [fb] with uncertainties [$^{+ \%}_{- \%}$] and cut efficiency [%] of the selection cuts in table 1 for the dominant SM backgrounds, when modeled at the inclusive NLO+PS level without MPI (Incl.), with MPI (Incl.+MPI), and FxFx-merging with MPI (FxFx+MPI). Uncertainties are obtained by adding the renormalization and factorization scale envelope with the shower scale envelope and statistical uncertainty in quadrature. At the generator-level, statistical confidence corresponds to 5-10 M events for each sample and shower variation.

As a benchmark, we considered CMS's Run II search [1806.05264]

- Changes: relaxed leading $p_T^{\ell_1}$ cut and swapped out $p_T^{\text{Veto}} = 25$ GeV
- + FxFx1j merging, MPI/UE tuning

Plot: Sensitivity change using $H_T^{\text{Veto}} = p_T^{\ell_2}$ in $(m_{\tilde{\chi}}, m_{\tilde{\mu}})$ -space

darker = improvement

- Improved sensitivity
to large mass splitting

- Worse for small
mass splitting

- Other ratios show
qualitatively opposite
behavior [1901.09937]

