Parton Degrees of Freedom: Connected and Disconnected Sea Partons from CT18 Parametrization of PDFs

Tie-Jiun Hou
Northeastern University, China

Work with Jian Liang(UKY), Keh-Fei Liu(UKY), Meng-Shi Yan(PKU) and C.-P. Yuan(MSU)

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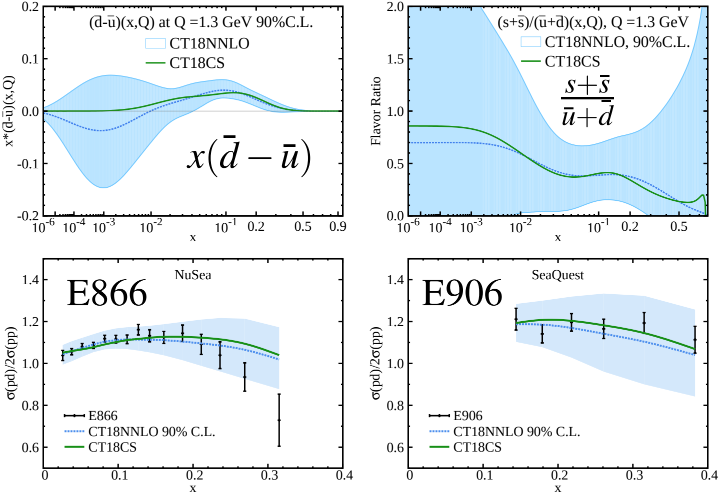
CT18CS Global PDF Analysis

An alternative parametrization of CT18 for $\bar{u}$ and $\bar{d}$ at $Q_0 = 1.3$ GeV scale with an input from Lattice QCD.



CT18

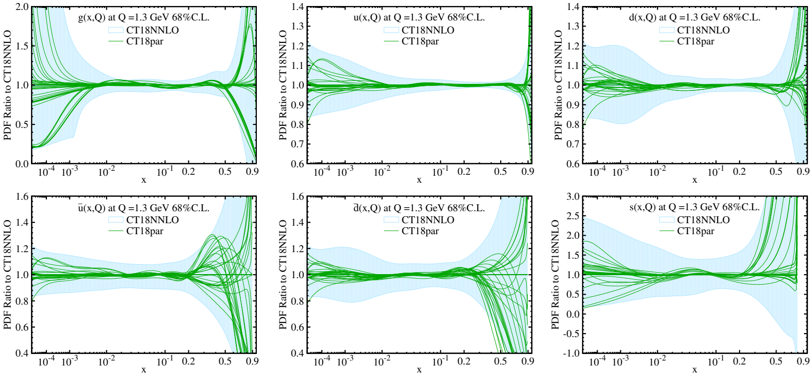
LHC data sets included in CT18 (1912.10053, T.- J. Hou *et al.*)

245	1505.07024	LHCb Z (W) muon rapidity at 7 TeV(applgrid)
246	1503.00963	LHCb 8 TeV Z rapidity (applgrid);
249	1603.01803	CMS W lepton asymmetry at 8 TeV (applgrid)
250	1511.08039	LHCb Z (W) muon rapidity at 8 TeV(applgrid)
253	1512.02192	ATLAS 7 TeV Z p_T (applgrid)
542	1406.0324	CMS incl. jet at 7 TeV with R=0.7 (fastNLO)
544	1410.8857	ATLAS incl. jet at 7 TeV with R=0.6 (applgrid)
545	1609.05331	CMS incl. jet at 8 TeV with R=0.7 (fastNLO)
573	1703.01630	CMS 8 TeV $t\bar{t}$ (p_T, y_t) double diff. distributions (fastNNLO)
580	1511.04716	ATLAS 8 TeV $t\bar{t}$ p_T and $m_{t\bar{t}}$ diff. distributions (fastNNLO)

Theory Calculation at NNLO

Obs.	Expt.	fast table	NLO code	K-factors	R,F scales
Inclusive jet	ATL 7 CMS 7/8	APPLgrid fastNLO	NLOJet++	NNLOJet	$\mathrm{p_T}, \mathrm{p_T^1}$
$\mathrm{p_T^Z}$	ATL 8	APPLgrid	MCFM	NNLOJet	$\sqrt{\mathrm{Q}^2 + \mathrm{p}_{\mathrm{T,Z}}^2}$
W/Z rapidity W asymmetry	LHCb 7/8 ATL 7 CMS 8	APPLgrid	MCFM/aMCfast	FEWZ/MCFM	$\mathrm{M_{W,Z}}$
DY (low,high mass)	ATL 7/8 CMS 8	APPLgrid	MCFM/aMCfast	FEWZ/MCFM	$\mathrm{Q_{ll}}$
$\mathrm{t\bar{t}}$	ATL 8 CMS 8	fastNNLO			$\frac{\mathrm{H_T}}{4}, \frac{\mathrm{m_T}}{2}$

Non-perturbative forms of PDFs in CT18 at $Q_0 = 1.3\mathrm{GeV}$



- CT18 – sample result of exploring various non-perturbative parametrization forms at $Q_0 = 1.3$ GeV.
- There is no data to constrain very large or very small x region.

Non-perturbative forms of PDFs in CT18 at $Q_0 = 1.3$ GeV

In CT18, 6 d.o.f of partons are parametrized at $Q_0 = 1.3$ GeV,

$$g, \quad u^v, \quad d^v, \quad \bar{u}, \quad \bar{d}, \quad s.$$

Where $\bar{s}(x) \equiv s(x)$ is assumed, and $u = u^v + \bar{u}$ and $d = d^v + \bar{d}$.
The functional form for the 6 parton flavors is

$$f^i(x, Q = Q_0) = a_0^i x^{a_1^i - 1}(1-x)^{a_2^i} P^i(x)$$

Where a_1^i term describes the behavior of PDFs as $x \to 0$, and a_2^i term describes the behavior of PDFs as $x \to 1$. The $P^i(x)$ is some choice of Berstein polynomial.
In total, there are 29 shape parameters used in CT18.
As a result, the full CT18 global fit yields $\chi^2 = 4292$, with a total of 3681 data points, and $\chi^2/N_{pt} = 1.17$.

Gottfried sum rule

Gottfried sum rule was originally obtain by assuming $\bar{u}$ and $\bar{d}$ to be the same, which leads to

$$S_G = \int_0^1 dx \frac{\left[F_2^p(x) - F_2^n(x)\right]}{x} = \frac{1}{3}, \quad \text{with } \bar{d}(x) \equiv \bar{u}(x)$$

New Muon Collaboration (NMC – PRL 66, 2712 (1991)), $\mu + p(n) \to \mu + X$, obtained $S_G = 0.240 \pm 0.016$ at $Q = 2$ GeV. The correct expression of sum rule is,

$$S_G = \frac{1}{3} - \frac{2}{3} \int_0^1 dx \left(\bar{d}(x) - \bar{u}(x)\right) + O(\alpha_s^2).$$

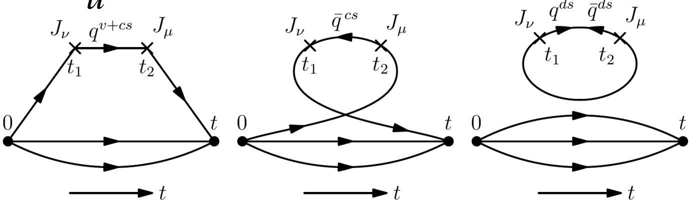
Hence, NMC data gives

$$\int_0^1 dx \left(\bar{d}(x) - \bar{u}(x)\right) = 0.140 \pm 0.024, \quad \text{at } Q_0 = 2.0 \text{ GeV}$$

What is the origin of $\bar{d} \neq \bar{u}$?

Motivated by Hadronic tensor in QCD path-integral formalism in Euclidian space.

u^{v+cs} $\bar{u}^{cs}$ u^{ds}, $\bar{u}^{ds}$



Connected sea(CS) Disconnected sea(DS)

$$
\begin{aligned}
u &= u^{v+cs} + u^{ds}, \quad & d &= d^{v+cs} + d^{ds} \\
\bar{u} &= \bar{u}^{cs} + \bar{u}^{ds}, \quad & \bar{d} &= \bar{d}^{cs} + \bar{d}^{ds}
\end{aligned}
$$

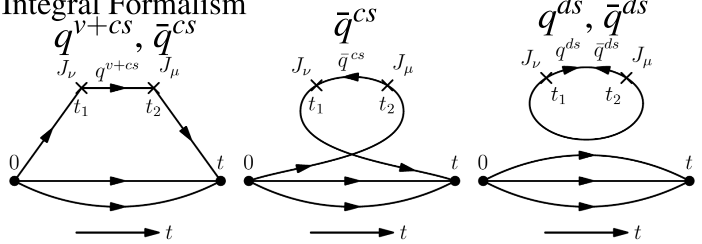
Define $u^v \equiv u^{v+cs} - \bar{u}^{cs}$, which is equivalent to defining $u^{cs} \equiv \bar{u}^{cs}$. Hence,

$$
\begin{aligned}
u - \bar{u} &\equiv (u^{v+cs} + u^{ds}) - (\bar{u}^{cs} + \bar{u}^{ds}) = u^v + (u^{ds} - \bar{u}^{ds}) \\
&\neq u^v, \quad \text{unless } u^{ds} = \bar{u}^{ds}
\end{aligned}
$$

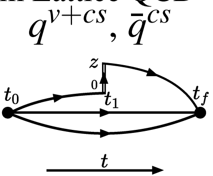
Similarly, $d^v \equiv d^{v+cs} - \bar{d}^{cs}$.

Hadronic tensor in Euclidean path-integral formalism versus Quasi-PDF from Lattice QCD

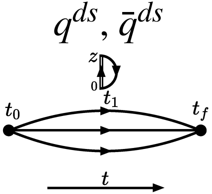
- Path-Integral Formalism



- Quasi-PDF from Lattice QCD

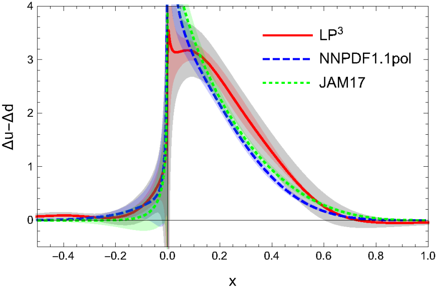


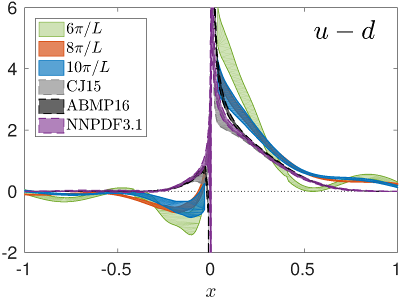
Connected Insertion(CI)



Disconnected Insertion(DI)

Quasi PDF results from LP3 and ETMC - connected insertion calculation





LP3 – H.W. Lin et al, PRL, arXiv:1807.07431

C. Alexandrou et al, PRL, arXiv:1803.02685

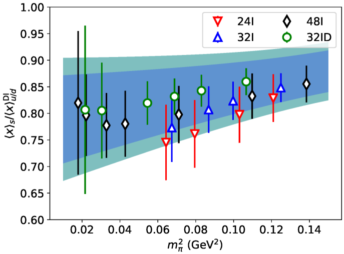
Where

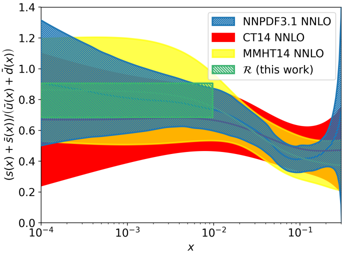
$$q(x>0) = q^{v+cs}(x), \qquad q(x<0) = -\bar{q}^{cs}$$

Parton degrees of freedom are the same as in hadronic tensor - K.F. Liu, PRD, arXiv:2007.15075

Lattice input to global fitting of PDFs

Lattice result from overlap on $N_f = 2+1$ DWF on 4 lattices, with one at physical pion mass (J. Liang et al., χQCD,PRD,arXiv:1901.07526)





$$\frac{1}{R} = \frac{\langle x \rangle_{s+\bar{s}}}{\langle x \rangle_{\bar{u}+\bar{d}}(DI)} (\text{at } 1.3 \text{ GeV}) = 0.822(69)(78)$$

This is the only Lattice data used in the CT18CS analysis.

Strategy for global analysis in CT18CS

In CT18CS, the non-perturbative PDFs parametrized at $Q_0 = 1.3$ GeV are :

$$g, \quad u^v, \quad d^v, \quad \bar{u}^{cs}, \quad \bar{d}^{cs}, \quad s^{ds}.$$

In this analysis,

- Disconnected Sea (DS) components:
  - Similar to CT18 fit, we assume $\bar{s}(x) = s(x)$. Hence, $s^{ds}(x) = \bar{s}^{ds}(x)$.
  - Likewise, we also assume $u^{ds}(x) = \bar{u}^{ds}(x)$ and $d^{ds}(x) = \bar{d}^{ds}(x)$ for simplicity.
  - Assuming isospin symmetry for u and d quark PDF

  This leads to

$$u^{ds} = \bar{u}^{ds} = d^{ds} = \bar{d}^{ds} = Rs = R\bar{s},$$

With $1/R = 0.822$ at $Q_0 = 1.3$ GeV.

Parton degrees of freedom at $Q_0 = 1.3$ GeV

- Connected Sea (CS) components:
  We define $u^{cs} \equiv \bar{u}^{cs}$ and $d^{cs} \equiv \bar{d}^{cs}$. They will be separately determined by the global fit, though with the same a_1 and a_2 components.

The physical parton degrees of freedom used in CT18CS are then:

In CT18

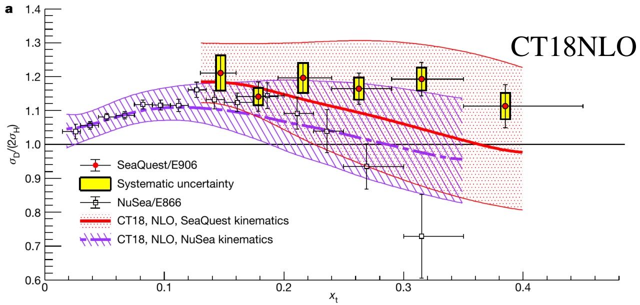
$$
\begin{aligned}
g &= g_{par} \\
u^v &= u^v_{par} \\
d^v &= d^v_{par} \\
\bar{u} &= \bar{u}^{cs} + \bar{u}^{ds} = \bar{u}_{par} + R s_{par} \\
\bar{d} &= \bar{d}^{cs} + \bar{d}^{ds} = \bar{d}_{par} + R s_{par} \\
s &= \bar{s} = s_{par}
\end{aligned}
$$

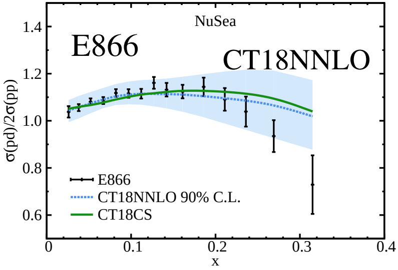
In CT18CS

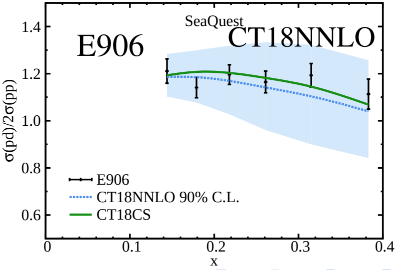
Both CT18 and CT18CS have 6 independent non-perturbative PDF functions, hence 6 parton degrees of freedom, at Q_0 scale.
The PDFs $f(x,Q)$ at $Q > Q_0$ are obtained by applying DGLAP evolution equations, as in CT18.

E906/SeaQuest

$$\frac{\sigma(pd)}{2\sigma(pp)}$$



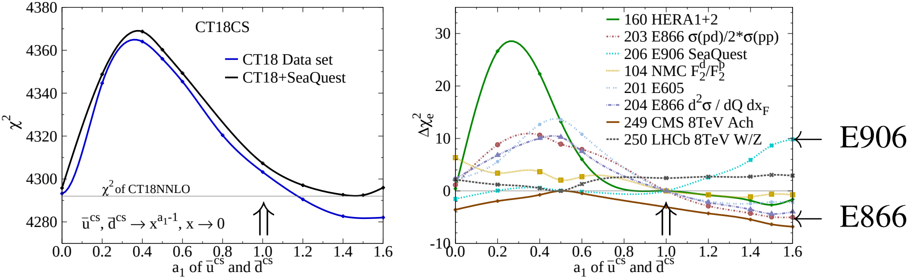




Small-x behavior of CT18CS PDFs

The non-perturbative PDF functions are chosen so that

- $\bar{d}/\bar{u} \xrightarrow{x\to 0} 1.$ $\Leftarrow$ Isospin symmetry, similar to CT18
- $\bar{u}^{ds}, \bar{d}^{ds}, \bar{s}^{ds} \xrightarrow{x\to 0} x^{-1}.$ $\Leftarrow$ Similar to CT18



- We scan the a_1 parameter of $\bar{u}^{cs}$ and $\bar{d}^{cs}$, and choose $a_1 = 1.0$ as CT18CS, where most of data are well fitted.
- The SeaQuest data was not included in the global fit.
- CT18CS has total $\chi^2 = 4299$ for $N_{pt} = 3681$. CT18 has total $\chi^2 = 4292$, only lower by 7 units.

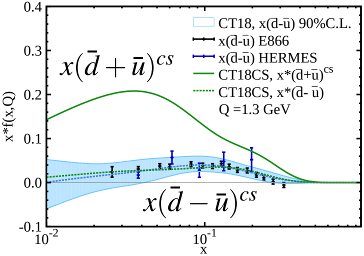
CT18CS PDFs

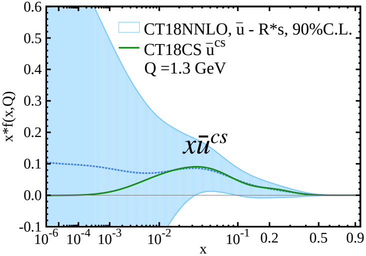
With the input of $\bar{u}^{ds} = \bar{d}^{ds} = Rs^{ds}$ from lattice QCD, and considering the scenario of small-x behavior, we obtain the CT18CS at $Q_0 = 1.3$ GeV scale.

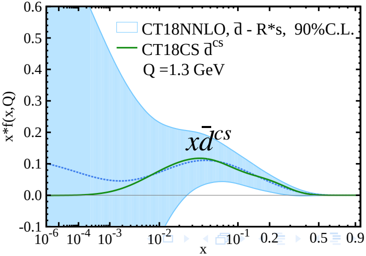
$$\frac{1}{R} \equiv \frac{\langle x \rangle_{s+\bar{s}}}{\langle x \rangle_{\bar{u}+\bar{d}}(DI)}$$

$$= 0.822$$

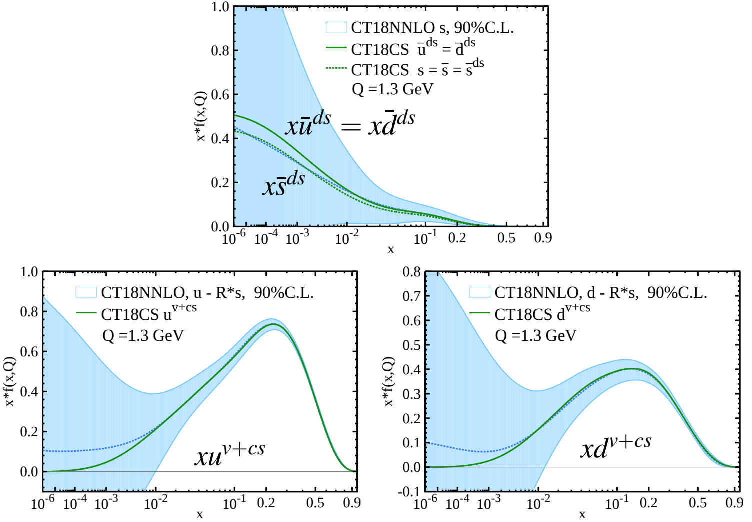
at 1.3 GeV



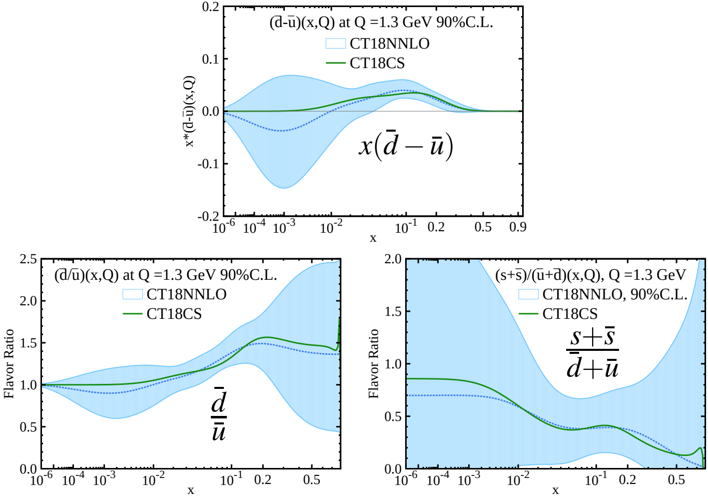




CT18CS PDFs


CT18NNLO s, 90%C.L.
CT18CS $\bar{u}^{ds} = \bar{d}^{ds}$
CT18CS $s = \bar{s} = \bar{s}^{ds}$
Q =1.3 GeV
$x\bar{u}^{ds} = x\bar{d}^{ds}$
$x\bar{s}^{ds}$
x*f(x,Q)
x
CT18NNLO, u - R*s, 90%C.L.
CT18CS u^{v+cs}
Q =1.3 GeV
xu^{v+cs}
CT18NNLO, d - R*s, 90%C.L.
CT18CS d^{v+cs}
Q =1.3 GeV
xd^{v+cs}

CT18CS PDFs



(d̄-ū)(x,Q) at Q =1.3 GeV 90%C.L.
CT18NNLO
CT18CS
$x(\bar{d}-\bar{u})$
x*(d̄-ū)(x,Q)
x
(d̄/ū)(x,Q) at Q =1.3 GeV 90%C.L.
CT18NNLO
CT18CS
$\frac{\bar{d}}{\bar{u}}$
Flavor Ratio
x
(s+s̄)/(ū+d̄)(x,Q), Q =1.3 GeV
CT18NNLO, 90%C.L.
CT18CS
$\frac{s+\bar{s}}{\bar{d}+\bar{u}}$
Flavor Ratio
x

The $\langle x \rangle$ Moment of CT18CS at 1.3 GeV

The $\langle x \rangle$ moment of CT18CS at 1.3 GeV:

u^{v+cs}	d^{v+cs}	$\bar{u}^{cs}$	$\bar{d}^{cs}$	$2u^{ds}(2d^{ds})$	$2s^{ds}$	gluon
0.335	0.157	0.013	0.021	0.038	0.026	0.386

More direct comparison between global analysis and lattice calculation can be done one by one, instead of being limited to $u-d$ and s.

	$\mu = 2.0$ GeV		$\mu = 1.3$ GeV
	CT18	Lattice	CT18CS
$\langle x \rangle_{u^+-d^+}$	0.156(7)	0.184(32)	0.172
$\langle x \rangle_{s^+}$	0.033(9)	0.092(41)	0.026

$$u^+ - d^+ = (u+\bar{u}) - (d+\bar{d}) = (u^{v+cs} + u^{ds} + \bar{u}^{cs} + \bar{u}^{ds}) - (d^{v+cs} + d^{ds} + \bar{d}^{cs} + \bar{d}^{ds})$$

$$\xrightarrow{CT18CS} (u^{v+cs} - d^{v+cs}) + (\bar{u}^{cs} - \bar{d}^{cs})$$

$$s^+ = s + \bar{s} = s^{ds} + \bar{s}^{ds} \xrightarrow{CT18CS} 2s^{ds}$$

Allow direct comparison between lattice calculations and global analysis for each parton degree of freedom.

Summary

- With the input from lattice QCD, $\frac{1}{R} \equiv \frac{\langle x \rangle_{s+\bar{s}}}{\langle x \rangle_{\bar{u}+\bar{d}}(DI)}$, we consider global analysis with the connected sea parton degrees of freedom taken into account, which are respondsible for $\bar{u} \neq \bar{d}$, as suggested by data.

- The result of global analysis, the CT18CS, is found to be compatible with CT18.

- The CT18CS allows to provide direct comparison between lattice calculations and global analysis for each parton degree of freedom.

Hadronic tensor in Euclidean path-integral formalism

- DIS in Minkowski space $\quad \frac{d^2\sigma}{dE'd\omega} = \frac{\alpha^2}{q^4}\frac{E'}{E} l^{\mu\nu} W_{\mu\nu}$

$$W_{\mu\nu}(\vec{q},\vec{p},\nu) = \frac{1}{\pi}\mathrm{Im}T_{\mu\nu} = \langle N(\vec{p})| \int \frac{d^4x}{4\pi} e^{iq\cdot x} J_\mu(x) J_\nu(0)|N(\vec{p})\rangle_{spin\ ave.}$$

$$= \frac{1}{2}\sum_n \int \prod_{i=1}^n \left[\frac{d^3p_i}{(2\pi)^3 2E_{pi}}\right](2\pi)^3\delta^4(p_n - p - q)\langle N(\vec{p})|J_\mu(0)|n\rangle\langle n|J_\nu(0)|N(\vec{p})\rangle_{spin\ ave.}$$

- Euclidean path-integral
(K.F. Liu and S.J. Dong, PRL 72, 1790 (1994), K.F. Liu, PRD 62, 074501 (2000))

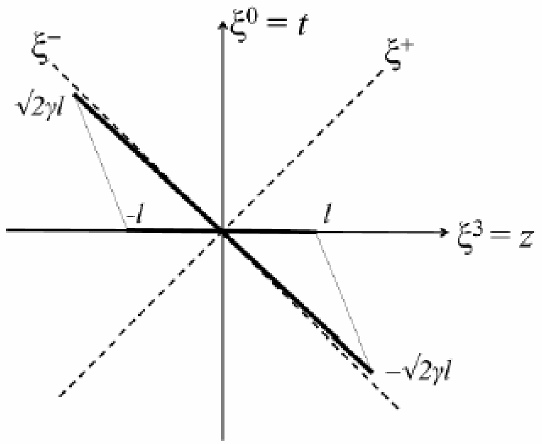
$$\widetilde{W}_{\mu\nu}(\vec{q},\vec{p},\tau)$$
$$= \frac{1}{4\pi}\sum_n \frac{2m_N}{E_n}\delta(\vec{p}_n - \vec{p} - \vec{q}) < N(\vec{p})|J_\mu|n >< n|J_\nu|N(\vec{p}) >_{spin\ ave.}\ e^{-(E_n - E_p)\tau}$$
$$= < N(\vec{p})|\sum_{\vec{x}} \frac{e^{-i\vec{q}\cdot\vec{x}}}{4\pi} J_\mu(\vec{x},\tau) J_\nu(0,0)|N(\vec{p}) >_{spin\ ave.}$$

Inverse problem
(Laplace transform)

$$D(\tau) = \int K(\tau,\nu)\rho(\nu)d\nu,$$
$$D(\tau) = \widetilde{W}_{\mu\nu}(\tau), \quad K(\tau,\nu) = e^{-\nu\tau}, \quad \rho(\nu) = W_{\mu\nu}(q^2,\nu)$$

PDF on Lattice: Quasi-PDF

Problem: No light-cone direction on Euclidean lattice.



Alternative approach:
spatial separation
$\downarrow$ *boost*
light-like separation

Credit: *Ji, PRL.2013*

Quasi-PDF:

$$\tilde{q}(x, \mu_R, P^z) = \int \frac{dz}{4\pi} e^{-ix\cdot zP^z} \boxed{\langle N|\bar{\psi}(z)\Gamma W(z,0)\psi(0)|N\rangle}$$

Can be calculated on lattice

Large Momentum Effective Theory (LaMET) [Ji,PRL.2013]:

$$\tilde{q}(x, \mu_R, P_z) = \int_{-1}^{1} \frac{dy}{|y|} C(\frac{x}{y}, \frac{\mu_R}{\mu}, \frac{\mu}{yP^z}) q(y, \mu) + O\left(\frac{\Lambda_{QCD}^2}{P_z^2}, \frac{m_N^2}{P_z^2}\right)$$

Strange quark and antiquark Quasi-PDFs

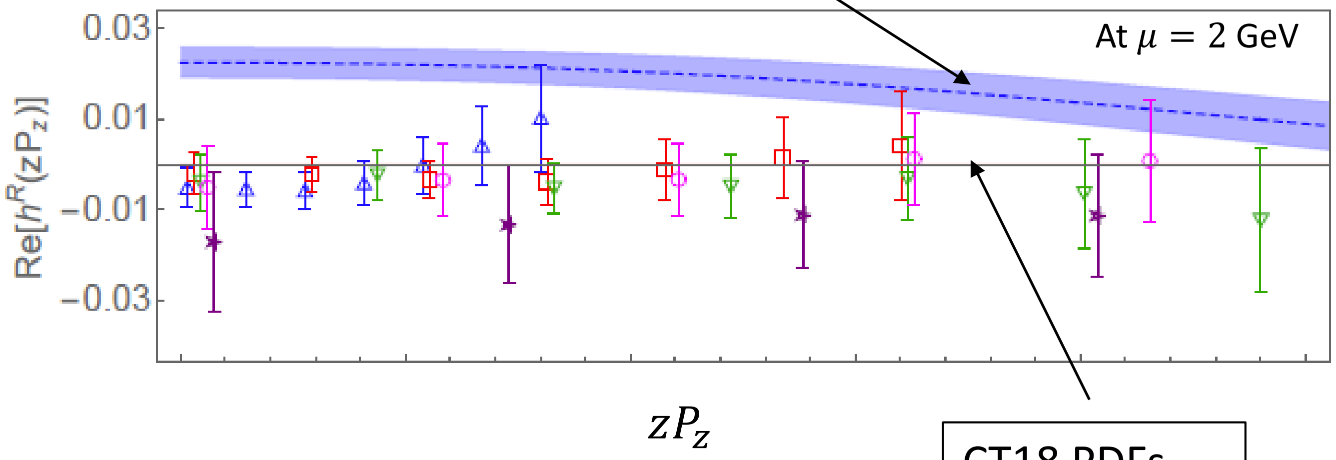
- Not a valence quark of the nucleon.
- Only disconnected contribution $\langle N|\bar{\psi}(z)\Gamma W(z,0)\psi(0)|N\rangle$ on Lattice.

Disconnected insertion diagrams are more difficult to calculate on Lattice. The first attempt was done in arXiv: 2005.01124 (by Rui Zhang, Huey-Wen Lin and Boram Yoon)

NNPDF3.1 PDFs

Ball, et.al, EPJC 77 (2017)

$$\mathrm{Re}[h(z)] \propto \int dx\,(s(x)-\bar{s}(x))\cos(xzP_z)$$



$s(x)=\bar{s}(x)$ is a good approximation,
to be consistent with Quasi-PDF calculation.

CT18 PDFs
$s(x)=\bar{s}(x)$

Hou, et.al, PRD 103 (2021)