

Particle production at all x : from proton-nucleus collisions to DIS

(toward a unified evolution equation including both large and small x)

Jamal Jalilian-Marian

Baruch College and City University of New York Graduate Center

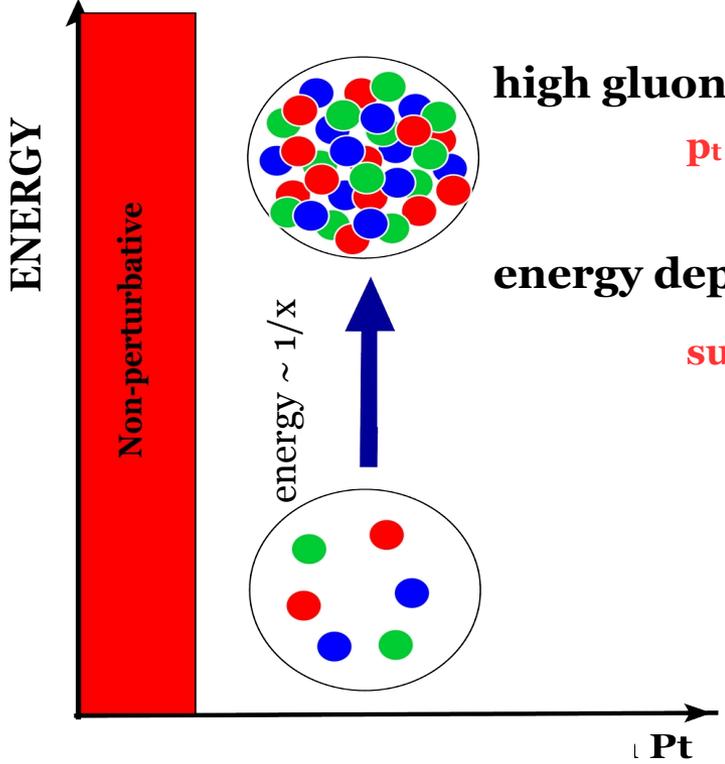
based on:

PRD102 (2020) 1, 014008

PRD99 (2019) 1, 014043

and work in progress

QCD at high energy/small x: gluon saturation



high gluon density: Eikonal multiple scattering
 p_t broadening (generic to multiple scattering)

energy dependence: x-evolution via JIMWLK/BK
 suppression of spectra/away side peaks

$$Q_s^2(x, b_t, A) \sim A^{1/3} \left(\frac{1}{x}\right)^{0.3}$$

$$Q_s^2(x = 3 \times 10^{-4}) \sim 1 \text{ GeV}^2$$

for a proton target (quarks)

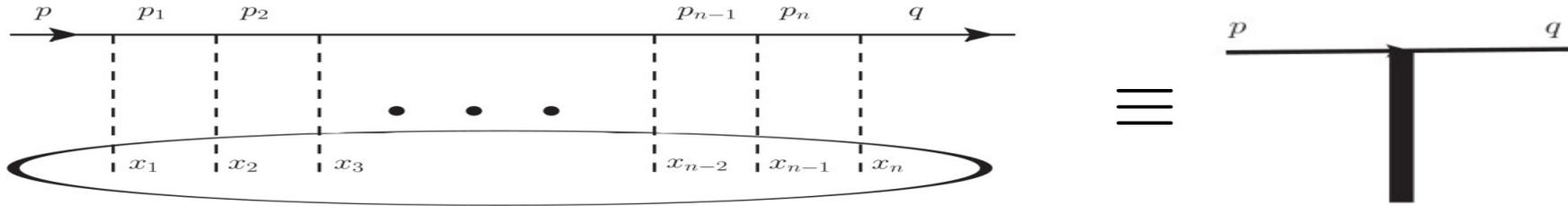
a framework for multi-particle production in QCD at small x/low p_t

- Shadowing/Nuclear modification factor*
- Azimuthal angular correlations (photon-jet,...)*
- Long range rapidity correlations (ridge,...)*
- Initial conditions for hydro*
- Thermalization (?)*

$$x \leq 0.01$$

$$\alpha_s \ln(x_v/x) \sim 1$$

CGC: eikonal approximation (tree level)



$$i\mathcal{M}(p, q) = 2\pi\delta(p^+ - q^+) \bar{u}(q) \not{n} \int d^2x_t e^{-i(q_t - p_t) \cdot x_t} [V(x_t) - 1] u(p)$$

with $V(x_t) \equiv \hat{P} \exp \left\{ ig \int_{-\infty}^{+\infty} dx^+ S_a^-(x^+, x_t) t_a \right\}$

scattering from small x gluons of the target
can cause only a small angle deflection

Dipole: DIS, proton-nucleus collisions
 x dependence from JIMWLK/BK evolution equation

$$\langle \text{Tr} V(x_\perp) V^\dagger(y_\perp) \rangle$$

toward precision at small x :

NLO corrections:

Chirilli+Xiao+Yuan, PRL (2012)

Balitsky+Chirilli, PRD88 (2013)

.....

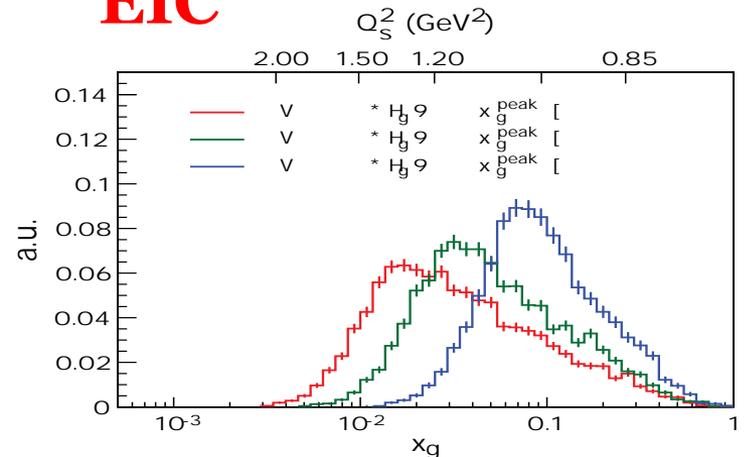
sub-eikonal corrections:

Kovchegov+Pitonyak+Sievert, JHEP (2017)

Agostini+Altinoluk+Armesto, EPJC (2019)

.....

EIC

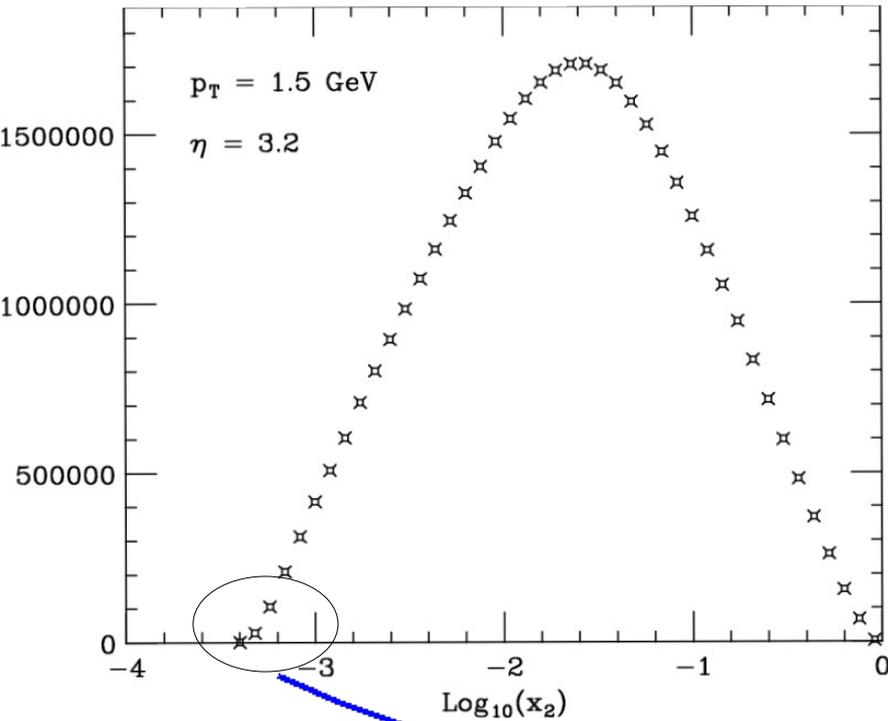


Aschenauer et al. ArXiv:1708.01527

single inclusive pion production in pp at RHIC

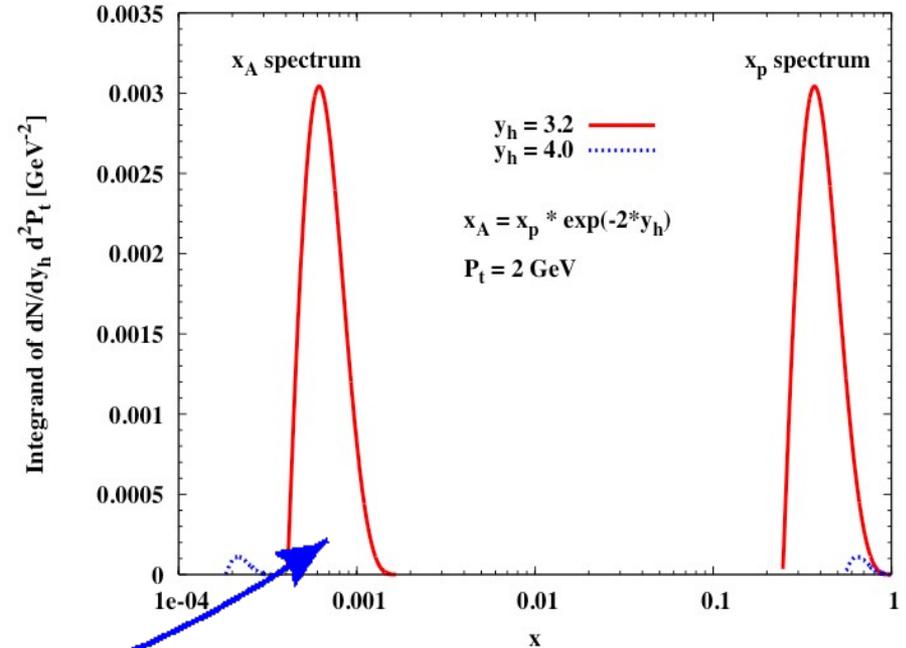
collinear factorization

GSV, PLB603 (2004) 173-183



CGC

DHJ, NPA765 (2006) 57-70

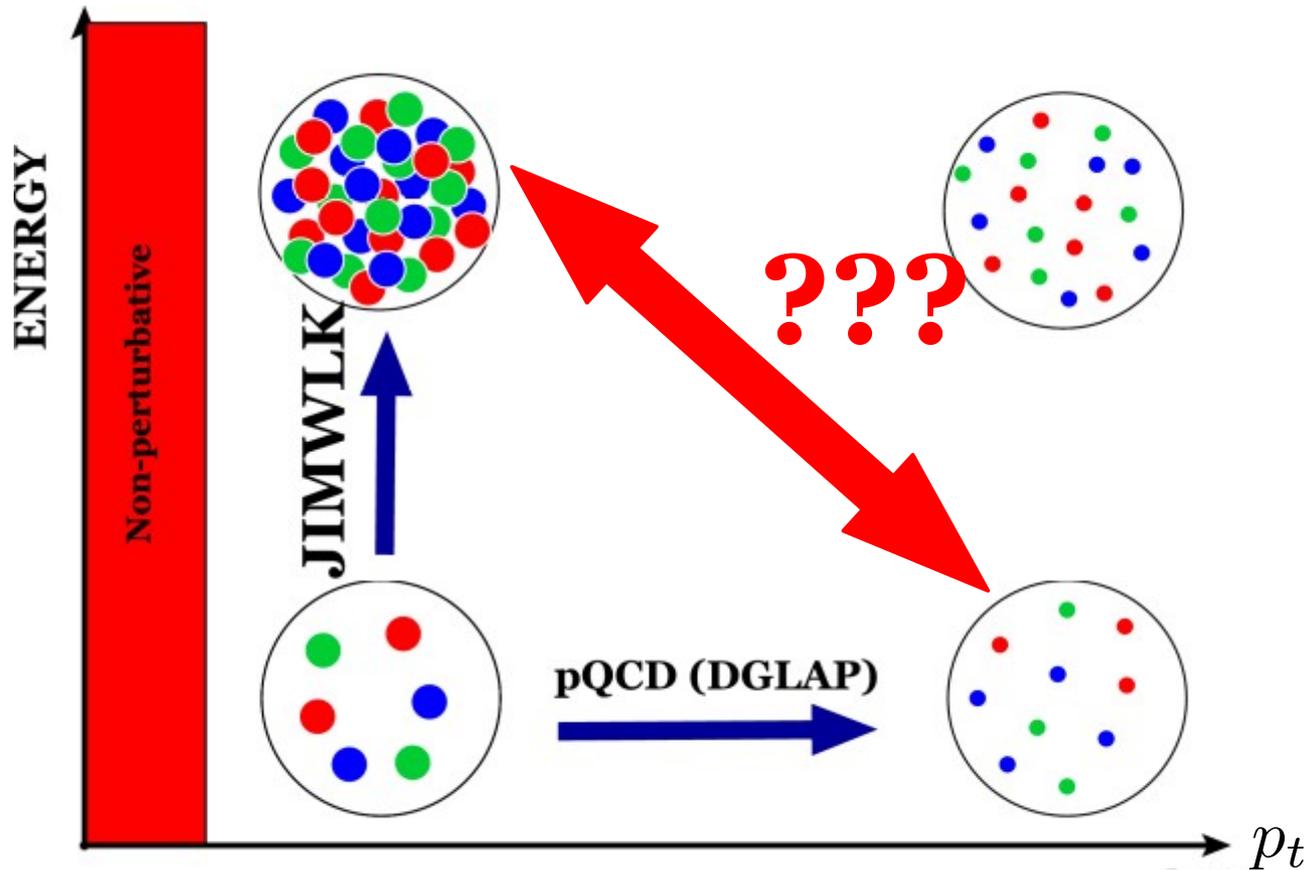


$$\int_{x_{\min}}^1 dx x G(x, Q^2) \dots \dots \rightarrow x_{\min} G(x_{\min}, Q^2) \dots$$

which kinematics are we in?



QCD kinematic phase space



unifying saturation with high p_t (large x) physics?

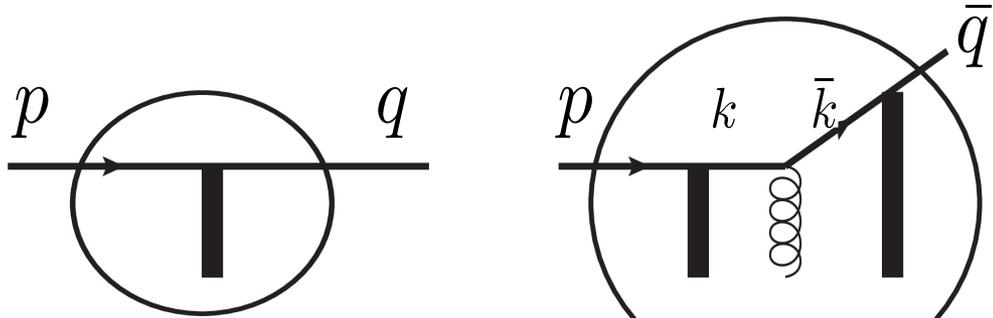
kinematics of saturation: where is saturation applicable?
*jet physics, high p_t and forward-backward correlations,
spin physics, early time e -loss in heavy ion collisions,*

quark scattering: beyond small x approximation

large x partons of target can cause a large-angle deflection of the quark

target gluon field

$$\mathcal{A}^\mu = \mathbf{S}^\mu + \mathbf{A}^\mu$$

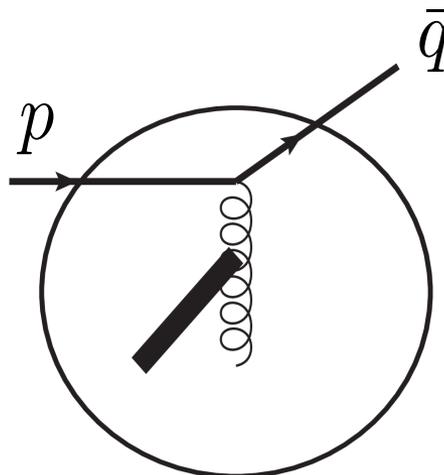


eik

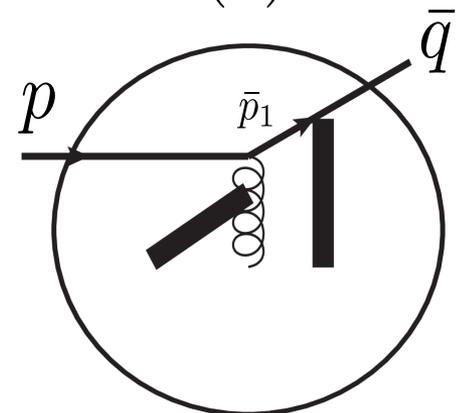
(1)

single scattering from large x gluons of target

$$\mathbf{A}^\mu = (\mathcal{A}^\mu - \mathbf{S}^\mu)$$



(2)



(3)

multiple scatterings from small x gluons of target

\mathbf{S}^μ

soft (eikonal) limit: $i\mathcal{M} \longrightarrow i\mathcal{M}_{eik}$

use spinor helicity formalism: helicity amplitudes

Including large x gluons of the target leads to:

longitudinal double spin asymmetries (ALL)

baryon transport (beam rapidity loss),

toward one-loop corrections: leading log evolution

gluon radiation

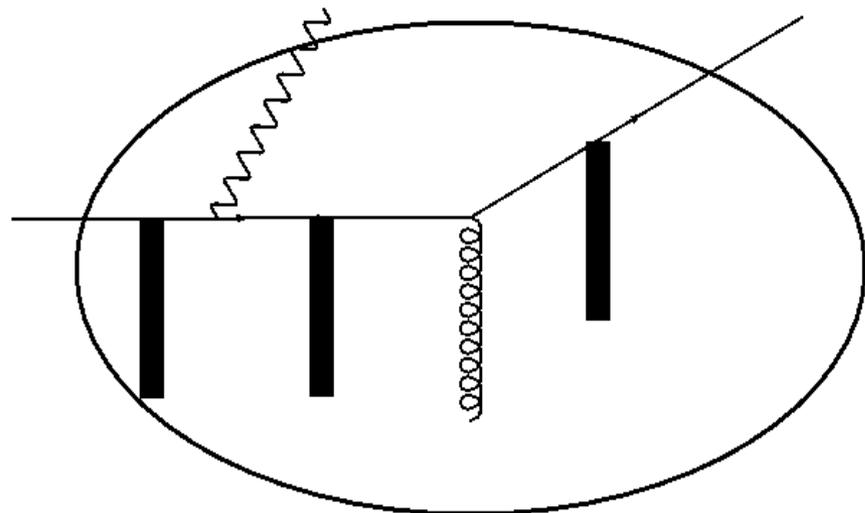
related problem: photon radiation

photon-jet correlations:

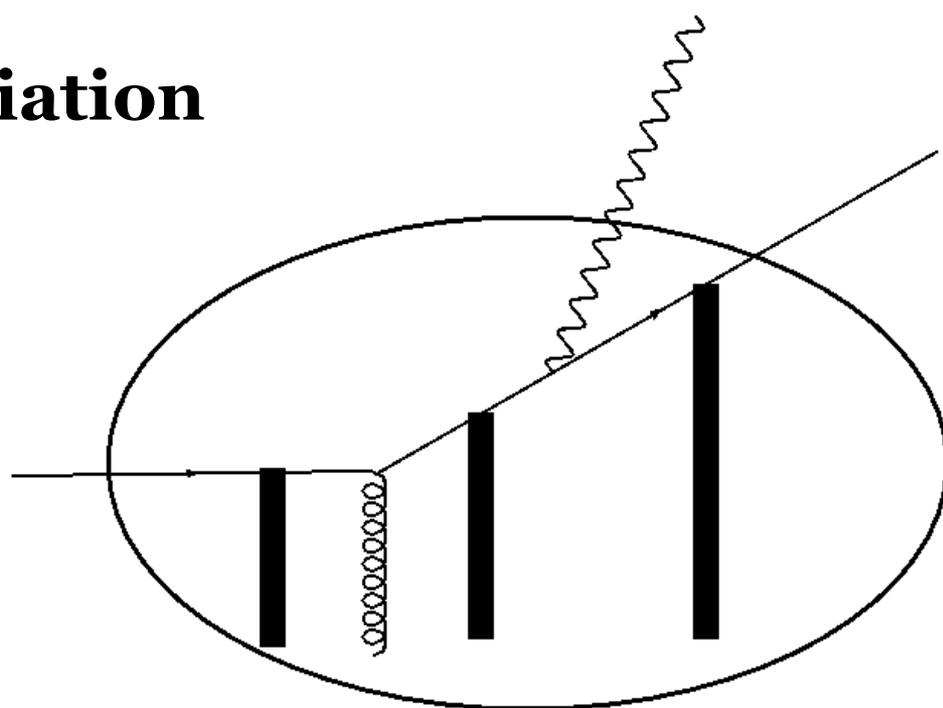
azimuthal angular correlations from low to high p_t

forward-backward rapidity correlations

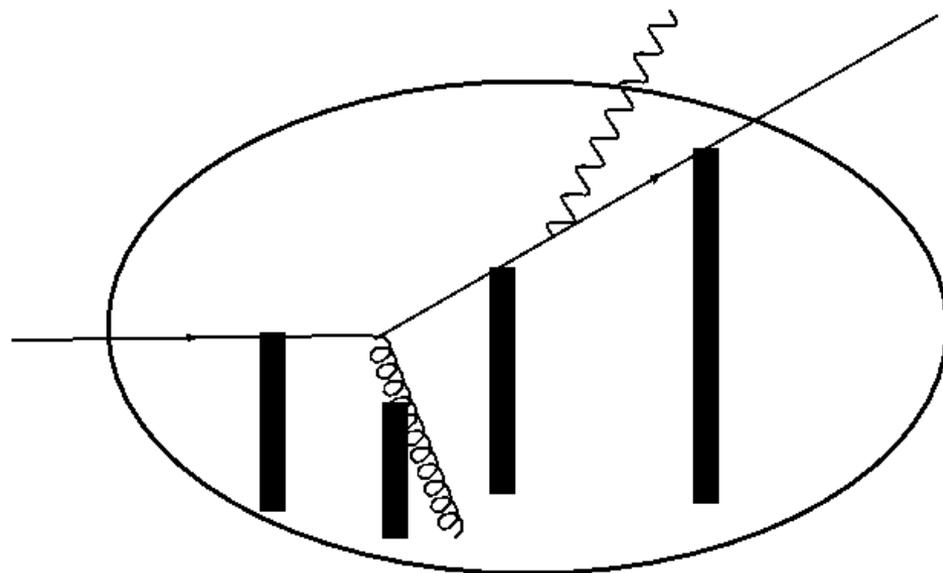
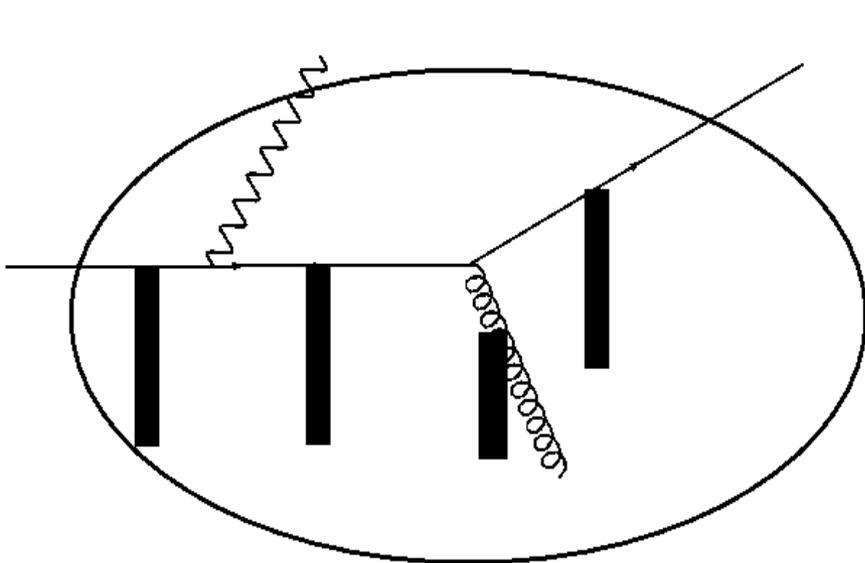
photon radiation



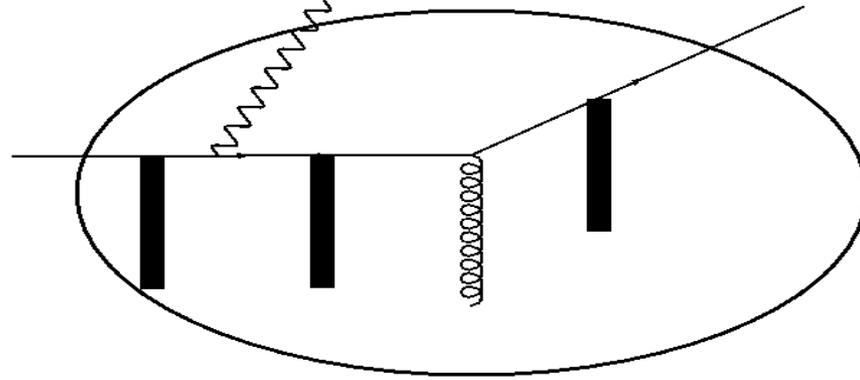
before hard scattering



after hard scattering



photon radiation: helicity amplitudes



$$\mathcal{N}_{1-1} = \bar{u}(\bar{q}) \frac{\not{\epsilon}}{2\bar{n} \cdot \bar{q}} \not{A}(x) \frac{\not{\epsilon} \not{l} \not{k}_1 \not{\epsilon}}{2n \cdot p 2n \cdot (p-l) 2n \cdot (p-l)} u(p)$$

$$\mathcal{N}_{1-2} = \bar{u}(\bar{q}) \frac{\not{\epsilon}}{2\bar{n} \cdot \bar{q}} \not{A}(x) \frac{\not{\epsilon} \not{l} \not{k}_1 \not{\epsilon}}{2n \cdot p 2n \cdot (p-l)} u(p)$$

$$\mathcal{N}_{1-1}^{++} = (\mathcal{N}_{1-1}^{--})^* = -\sqrt{\frac{n \cdot p}{n \cdot (p-l)}} \frac{[n \cdot l k_{2\perp} \cdot \epsilon_{\perp}^* - n \cdot (p-l) l_{\perp} \cdot \epsilon_{\perp}^*]}{n \cdot l n \cdot (p-l)} \langle \bar{k}_1^+ | \not{A}(x) | k_3^+ \rangle$$

$$\mathcal{N}_{1-2}^{++} = (\mathcal{N}_{1-2}^{--})^* = -\sqrt{\frac{n \cdot p}{n \cdot (p-l)}} \langle \bar{k}_1^+ | \not{A}(x) | n^+ \rangle$$

$$\mathcal{N}_{1-1}^{+-} = (\mathcal{N}_{1-1}^{-+})^* = -\sqrt{\frac{n \cdot p}{n \cdot (p-l)}} \frac{[n \cdot p l_{\perp} \cdot \epsilon_{\perp} - n \cdot l k_{1\perp} \cdot \epsilon_{\perp}]}{n \cdot p n \cdot l} \langle \bar{k}_1^+ | \not{A}(x) | k_3^+ \rangle$$

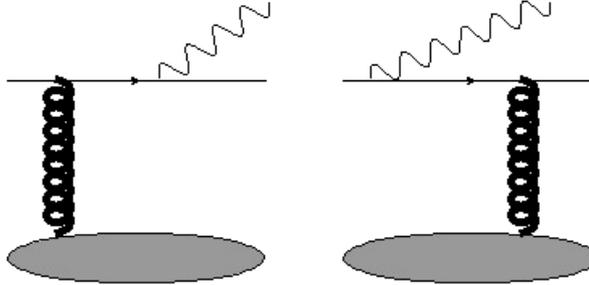
$$\mathcal{N}_{1-2}^{+-} = \mathcal{N}_{1-2}^{-+} = 0$$

from pA to DIS: crossing symmetry

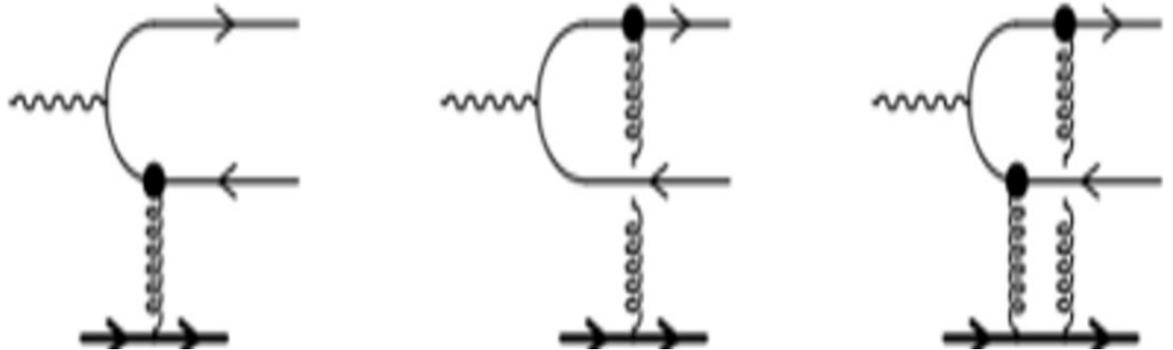
FG+JJM, hep-ph/0211363

recall at small x

$$q T \rightarrow q \gamma^* X$$



$$\gamma^* T \rightarrow q \bar{q}$$



can do the same crossing here to get to dijets in DIS

SUMMARY

CGC is a systematic approach to high energy collisions

strong hints from RHIC, LHC,...

toward precision: NLO, sub-eikonal corrections, ...

CGC breaks down at large x (high p_t)

a significant part of EIC/RHIC/LHC phase space is at large x

transition from large x physics to CGC (kinematics?)

Toward a unified formalism:

particle production in both small and large p_t kinematics

two-particle correlations: from forward-forward to forward-backward

one-loop correction: both collinear and CGC factorization limits

need to clarify/understand: gauge invariance, initial conditions,

beyond eikonal approximation:

large x partons of target can cause a large-angle deflection of the quark

