

Prompt photon production in pp scattering: high energy factorization approach

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Based on:

Phys. Rev. D **103**, 034013 (2021), 2010.00468

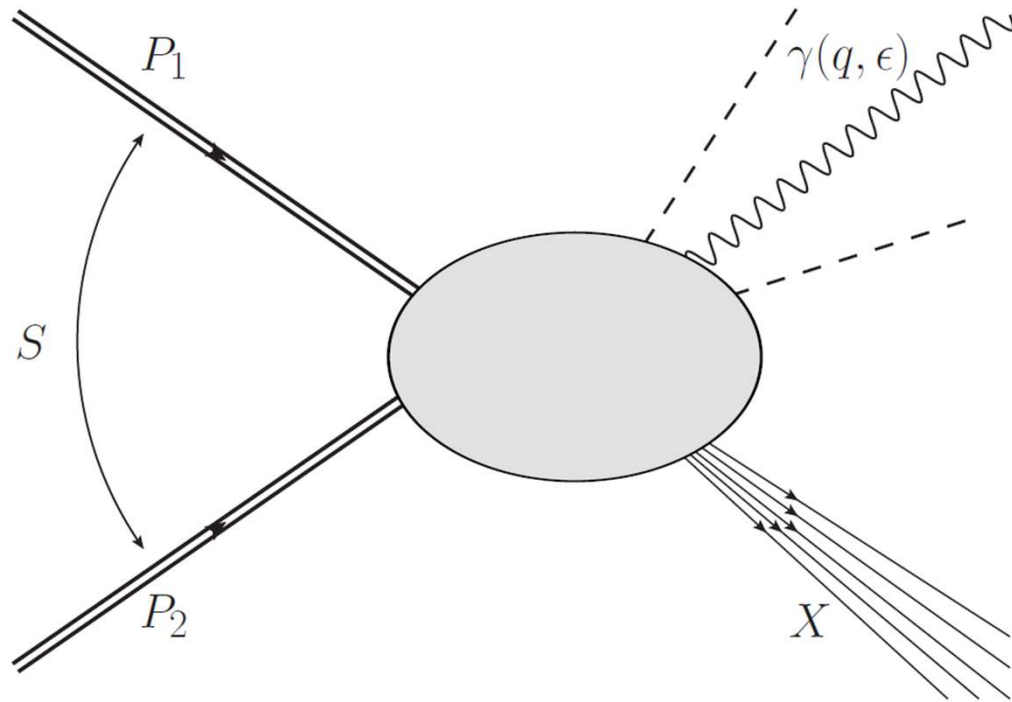


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Project is supported by National Science Centre,
grant no. 2019/32/C/ST2/00202.

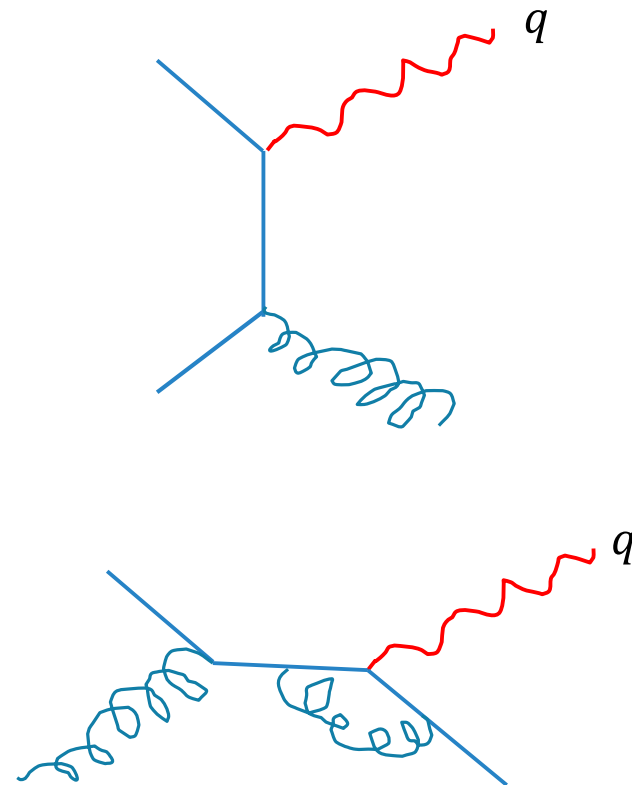
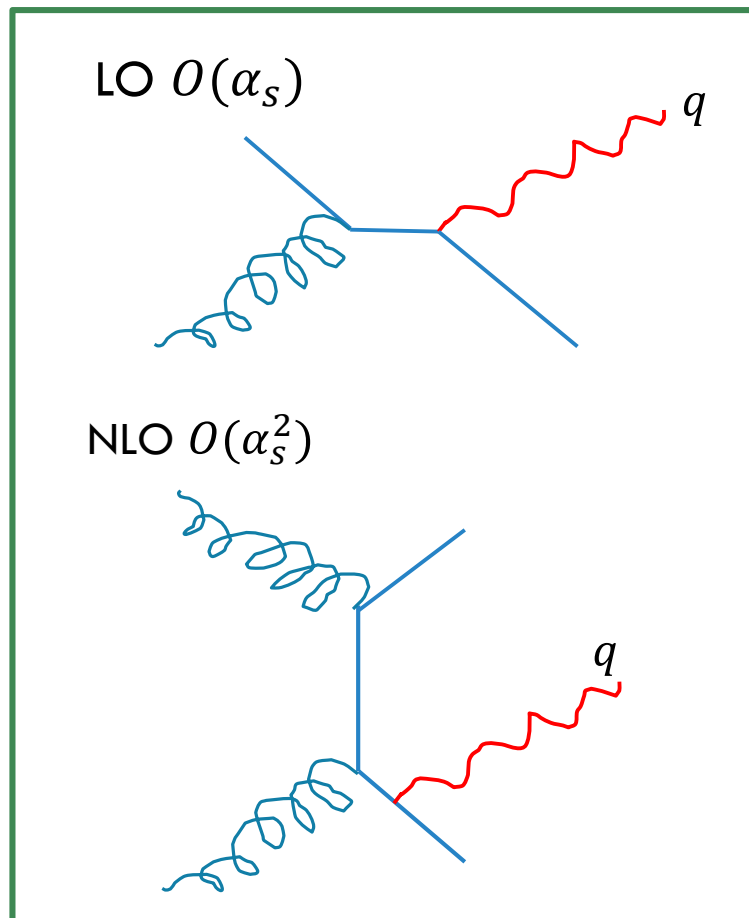
THE PROCESS

$$q_\gamma = (q_\gamma^+, q_\gamma^-, \mathbf{q}_T) = (x_F \sqrt{S}, \mathbf{q}_T^2 / x_F \sqrt{S}, \mathbf{q}_T)$$



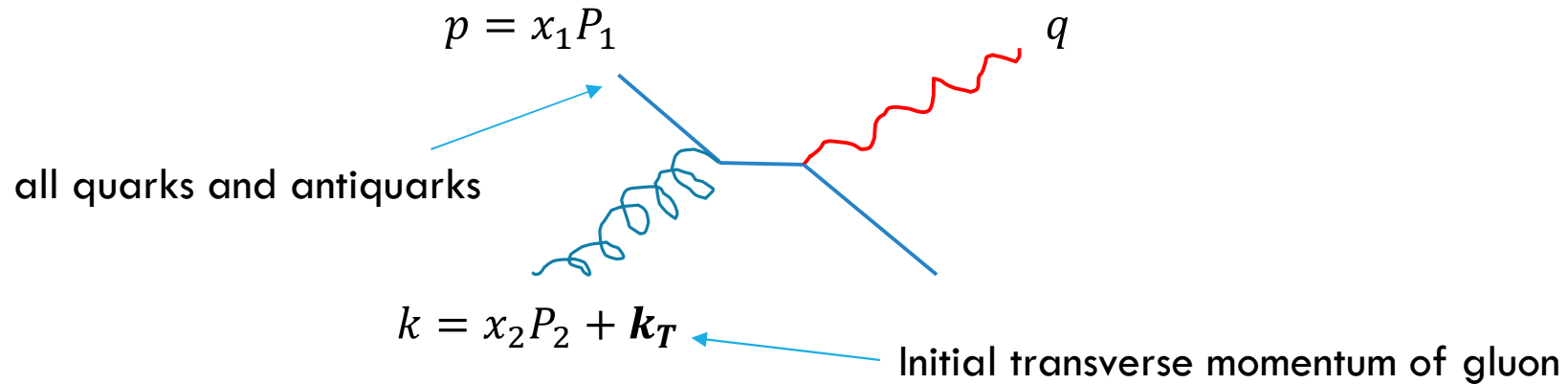
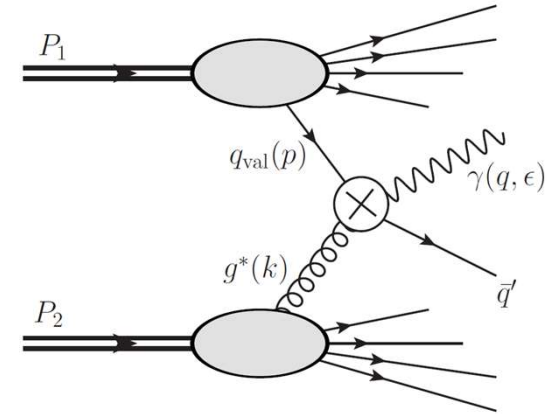
HIGH ENERGY APPROXIMATION

$E \gg \text{other scales}$ \longrightarrow Gluon distribution dominates



APPROACH qg^* (LO)

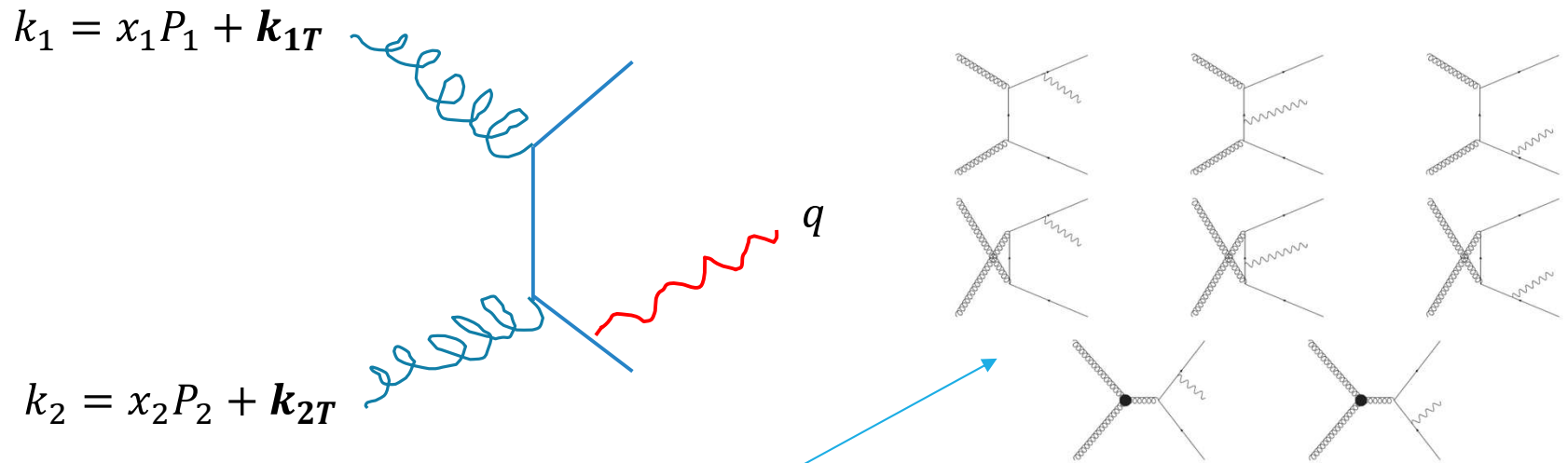
k_T -factorization approach (or rather „hybrid“):



$$\frac{d\sigma^\gamma}{dy d^2\mathbf{q}_T} = \frac{\alpha_{\text{em}}}{3\pi} \int_{x_F}^1 \frac{dz}{z} \frac{x_F}{z} \sum_{i \in \{f, \bar{f}\}} e_i^2 q_i\left(\frac{x_F}{z}, \mu_F\right) \\ \times \int \frac{d^2\mathbf{k}_T}{k_T^2} \alpha_s F_g(x_g, k_T, \mu_F) \frac{[1 + (1 - z)^2] z^2 \mathbf{k}_T^2}{\mathbf{q}_T^2 (\mathbf{q}_T - z\mathbf{k}_T)^2} + (y \rightarrow -y)$$

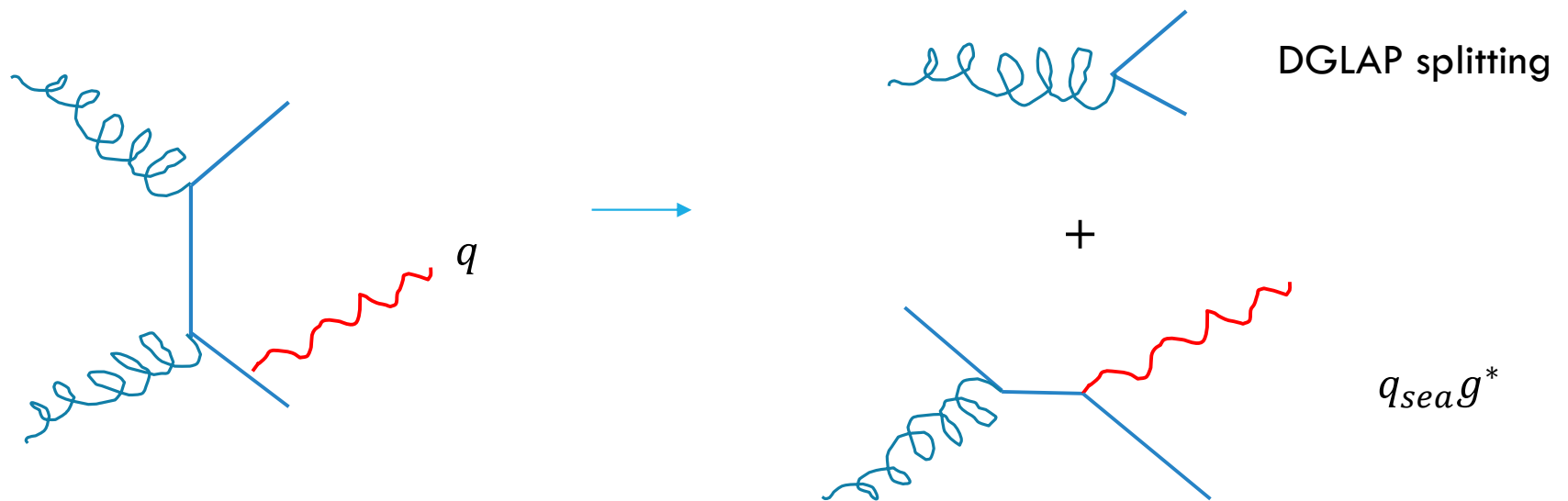
APPROACH $g^* g^*$ („NLO“)

k_T -factorization approach:



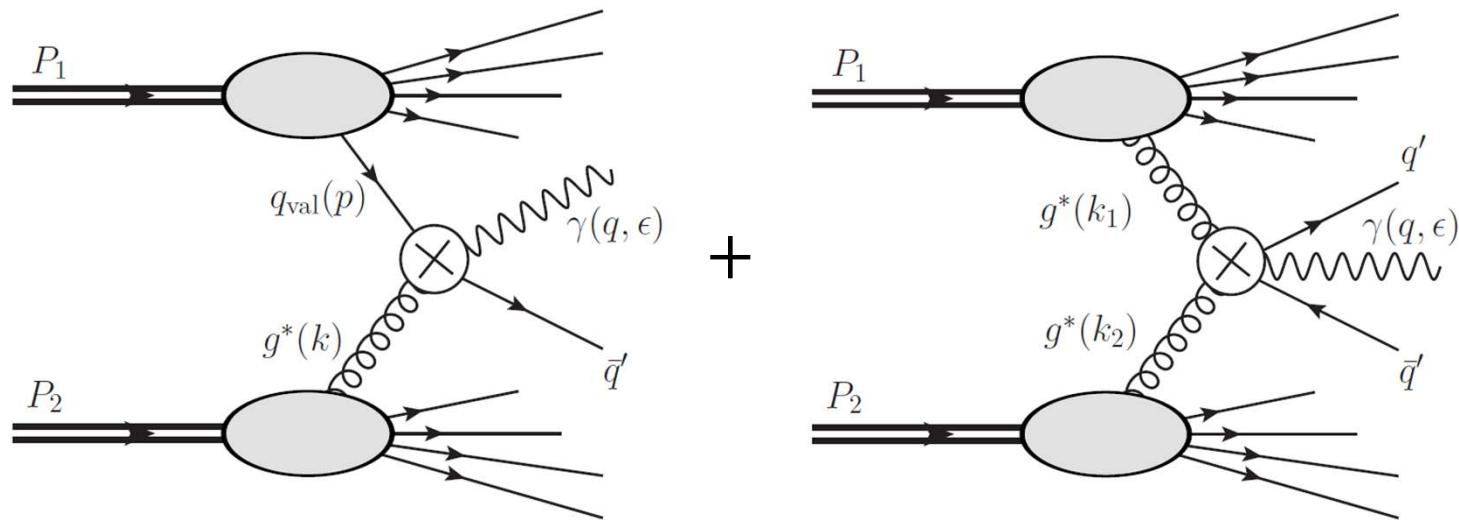
$$d\sigma_{\sigma}^{(g^* g^* \rightarrow q \bar{q} \gamma)} = \int dx_1 \int \frac{d^2 \mathbf{k}_{1T}}{\pi k_{1T}^2} F_g(x_1, k_{1T}, \mu_F) \int dx_2 \int \frac{d^2 \mathbf{k}_{2T}}{\pi k_{2T}^2} F_g(x_2, k_{2T}, \mu_F) \\ \times \frac{(2\pi)^4}{2S} \mathcal{H}_{\sigma} dPS_3(k_1 + k_2 \rightarrow p_3 + p_4 + q),$$

APPROACH $g^* g^*$ („NLO“)



We need to add the missing $q_{valence} g^*$ contribution

APPROACH $g^* g^* + q_{val} g^*$



$$d\sigma^\gamma = d\sigma^{(q_{val} g^* \rightarrow q \gamma)} + d\sigma^{(g^* g^* \rightarrow q \bar{q} \gamma)}$$

UNINTEGRATED GLUON DISTRIBUTIONS (uPDF)

$$d\sigma_{\sigma}^{(g^*g^* \rightarrow q\bar{q}\gamma)} = \int dx_1 \int \frac{d^2\mathbf{k}_{1T}}{\pi k_{1T}^2} F_g(x_1, k_{1T}, \mu_F) \int dx_2 \int \frac{d^2\mathbf{k}_{2T}}{\pi k_{2T}^2} F_g(x_2, k_{2T}, \mu_F) \\ \times \frac{(2\pi)^4}{2S} \mathcal{H}_{\sigma} dPS_3(k_1 + k_2 \rightarrow p_3 + p_4 + q),$$

g^*g^* channel

$$\frac{d\sigma^{\gamma}}{dy d^2\mathbf{q}_T} = \frac{\alpha_{\text{em}}}{3\pi} \int_{x_F}^1 \frac{dz}{z} \frac{x_F}{z} \sum_{i \in \{f, \bar{f}\}} e_i^2 q_i\left(\frac{x_F}{z}, \mu_F\right) \\ \times \int \frac{d^2\mathbf{k}_T}{k_T^2} \alpha_s F_g(x_g, k_T, \mu_F) \frac{[1 + (1-z)^2] z^2 \mathbf{k}_T^2}{\mathbf{q}_T^2 (\mathbf{q}_T - z\mathbf{k}_T)^2} + (y \rightarrow -y)$$

qg^* channel

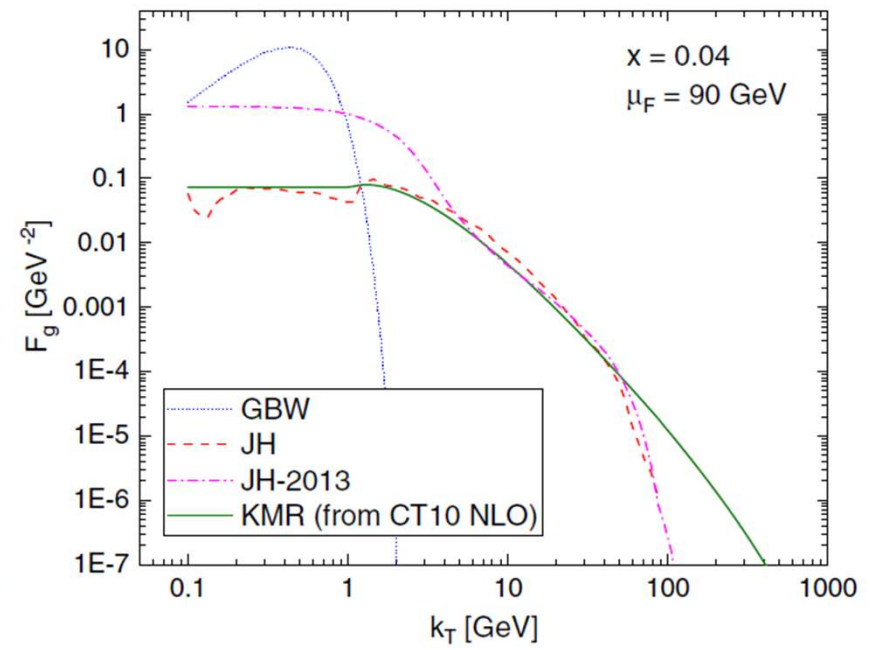
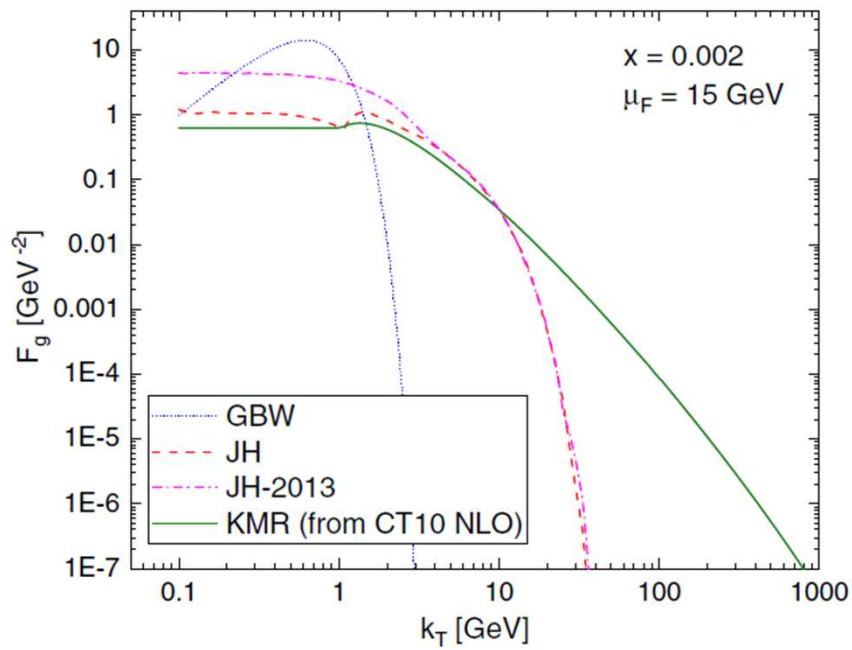
$F_g(x, k_T, \mu_F)$ is nonperturbative quantity and some models are needed.

$$\alpha_s F_g(x, k_T, \mu) = \frac{3\sigma_0}{4\pi^2} \frac{k_T^2}{Q_s^2} \exp(-k_T^2/Q_s^2) \times \left(\frac{1-x}{1-0.01} \right)^7$$

We will use:

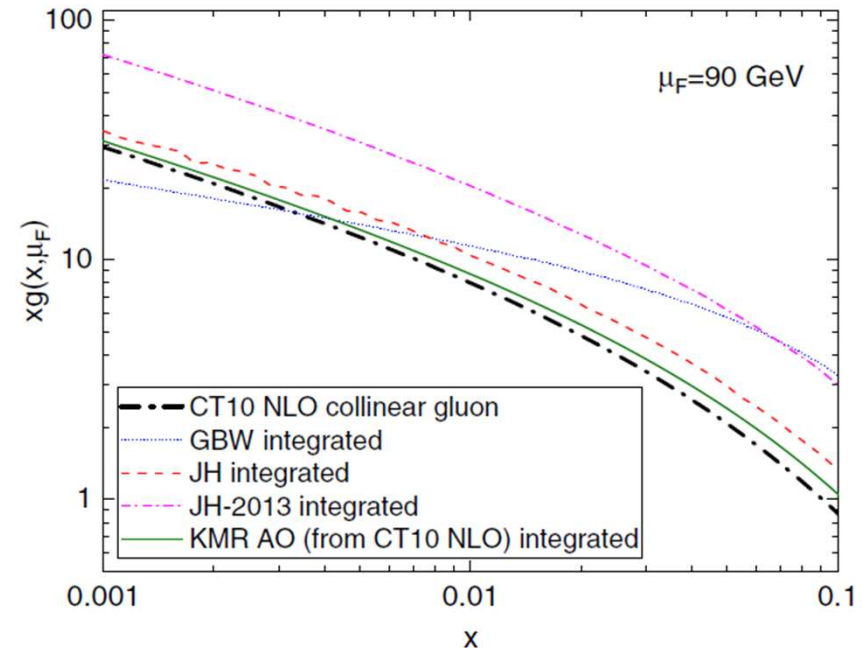
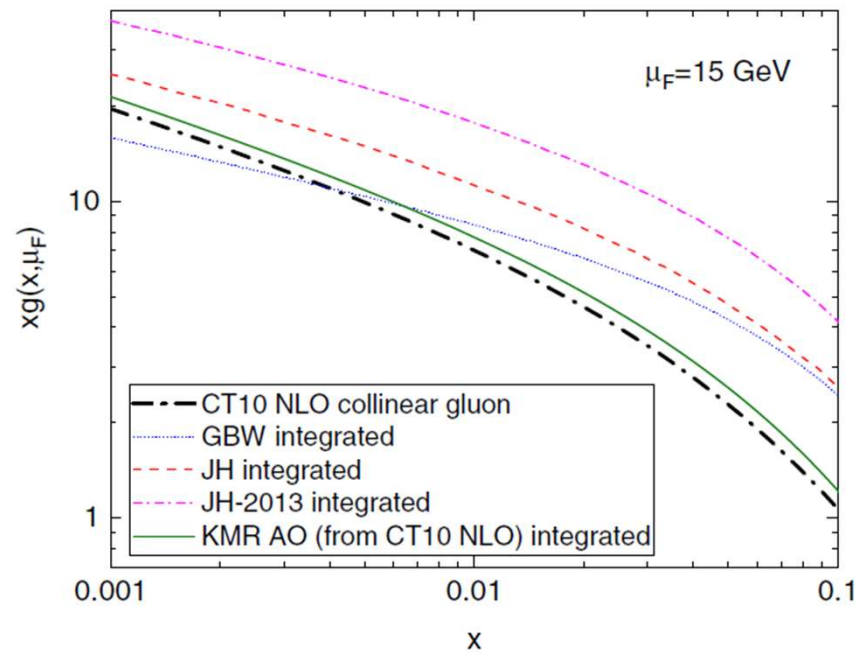
- 1) Golec-Biernat – Wusthoff (GBW) model (modified at large x)
- 2) CCFM solution by Jung and Hansson (JH) or Jung and Hautmann (JH-2013)
- 3) Kimber–Martin–Ryskin (KMR) uPDF

PLOTS OF u PDFs




PLOTS OF INTEGRATED u PDFs

$$\int_0^{\mu_F^2} dk_T^2 F_g(x, k_T, \mu_F) = xg(x, \mu_F)$$



DATA FOR ISOLATED PHOTON

We shall use the following data sets:

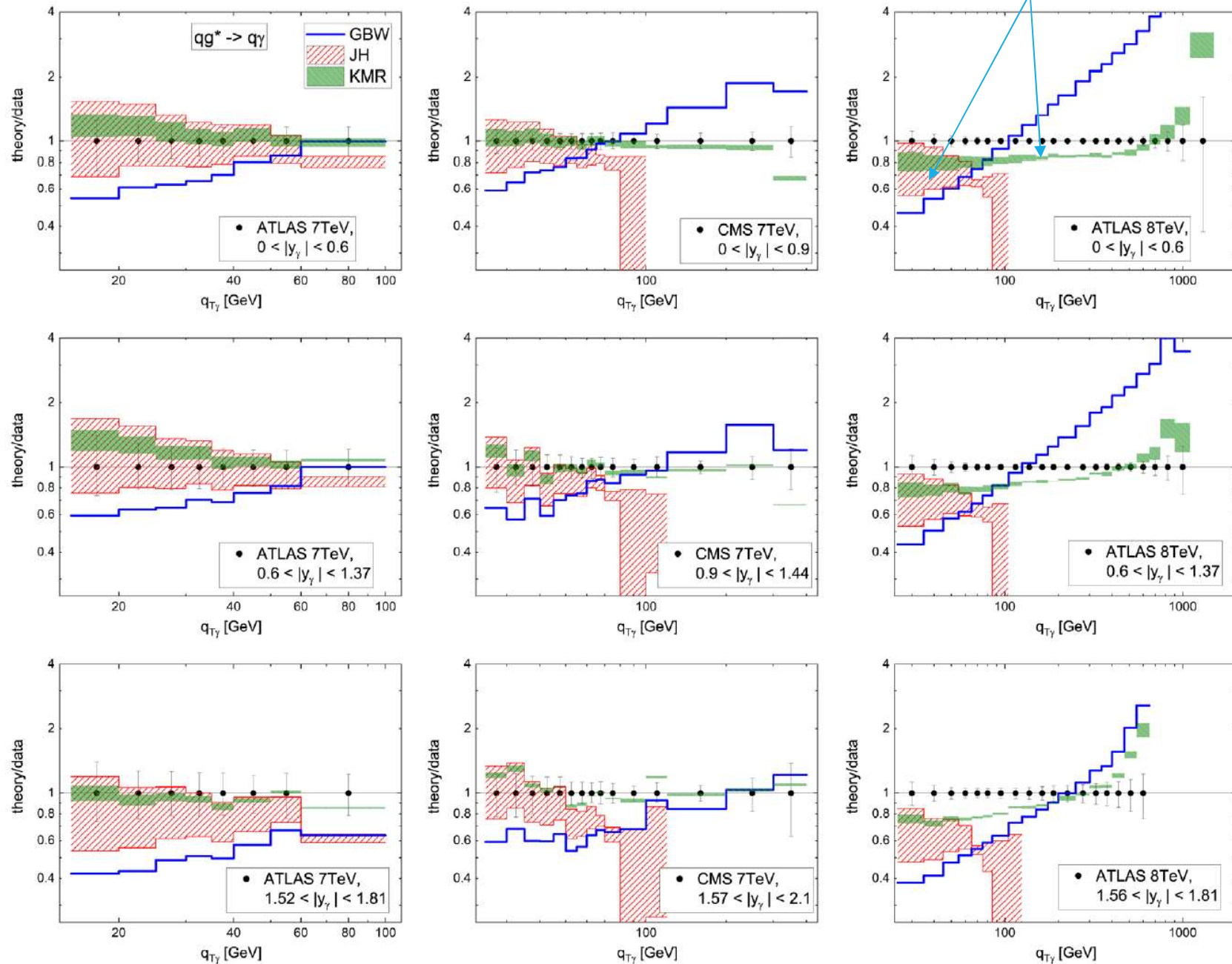
$$x \approx 0.002$$


- ATLAS data at $\sqrt{S} = 7$ TeV with $15 \text{ GeV} < q_T < 100 \text{ GeV}$ (ATLAS@7TeV)
- CMS data at $\sqrt{S} = 7$ TeV with $25 \text{ GeV} < q_T < 400 \text{ GeV}$ (CMS@7TeV)
- ATLAS data at $\sqrt{S} = 8$ TeV with $25 \text{ GeV} < q_T < 1500 \text{ GeV}$ (ATLAS@8TeV)

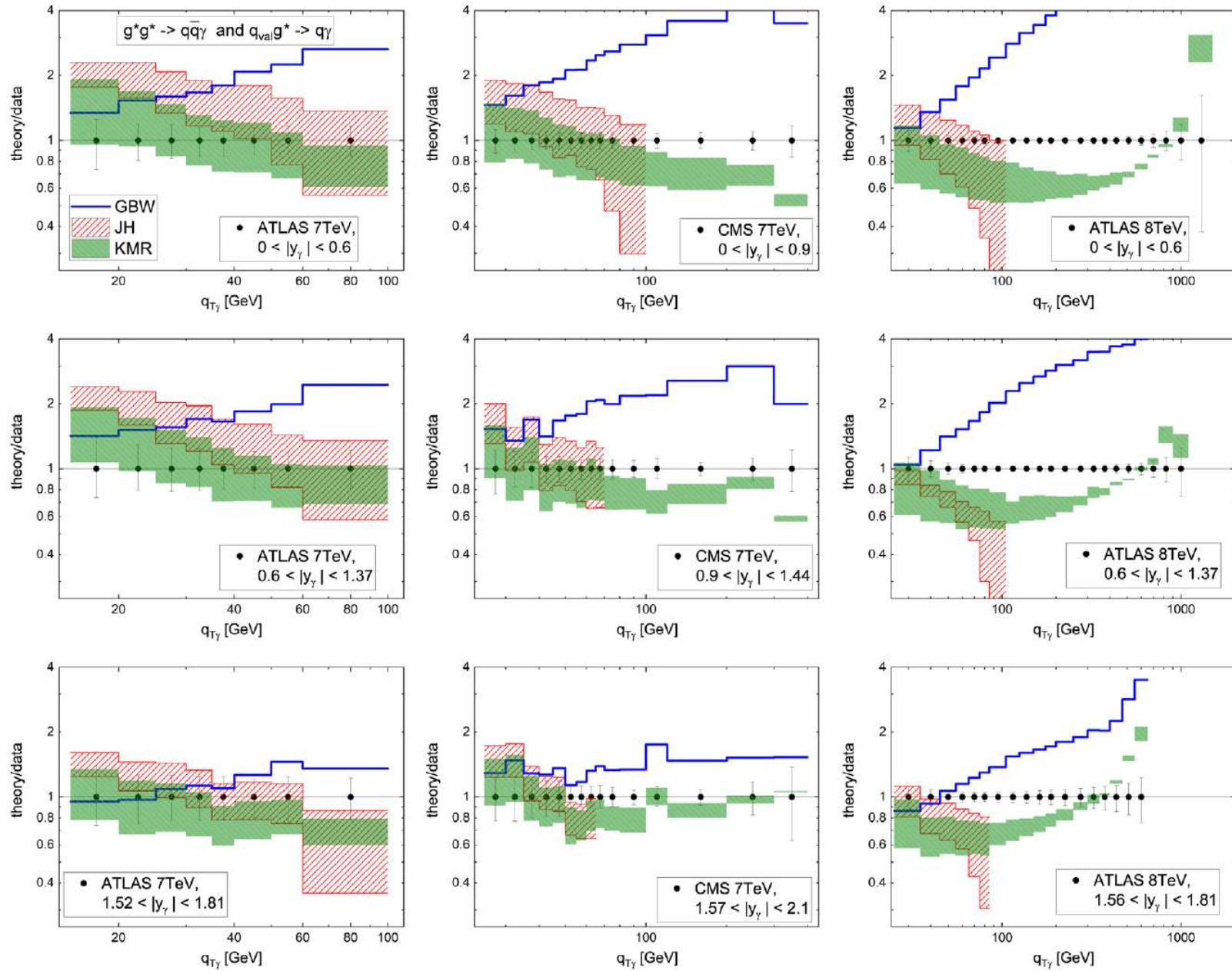

$$x \approx 0.2$$

These data can test the k_T -factorization framework at different values of x .

RESULTS FOR qg^*



RESULTS FOR $g^*g^* + q_{val}g^*$



SUMMARY

- We calculate cross-section for prompt photon production using k_T -factorization framework.
- We have tested three well-know transverse momentum dependent gluon distributions (GBW, CCFM, KMR)
- Good description of data was obtained, even at high q_T .
- Surprisingly, the g^*g^* channel does not improve the description of data.

THANK YOU!

Backup slides

GOLEC-BIERNAT – WUSTHOFF (GBW) GLUON DISTRIBUTION

$$\alpha_s F_g(x, k_T, \mu) = \frac{3\sigma_0}{4\pi^2} \frac{k_T^2}{Q_s^2} \exp(-k_T^2/Q_s^2)$$

where the the saturation scale $Q_s^2 = (x/x_0)^{-\lambda} \text{ GeV}^2$.

The parameters σ_0, x_0, λ were fitted to data with $x < 0.01$.

We introduce additional factor which modifies uPDF at high x :

$$\alpha_s F_g(x, k_T, \mu) = \frac{3\sigma_0}{4\pi^2} \frac{k_T^2}{Q_s^2} \exp(-k_T^2/Q_s^2) \times \left(\frac{1-x}{1-0.01} \right)^7$$

CATANI–CIAFALONI–FIORANI– MARCHESINI (CCFM) EQUATION

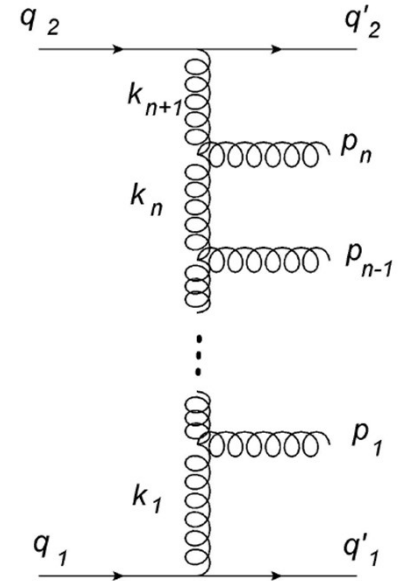
CCFM formalism is one the most popular modification of small- x evolution equation (BFKL) for $x \rightarrow 1$.

θ_i - emission angle of gluon p_i w.r.t. z axis defined by hadrons' momenta P_1 and P_2

CCFM requires angular ordering: $\theta_i > \theta_{i-1}$.

$$f_g(x, k_T, Q) = f_g^0(x, k_T, Q_0) + \int \frac{d^2 \vec{q}}{\pi q^2} \int_x^1 \frac{dz}{z} \theta(Q - zq) \theta(q - Q_0) \frac{\alpha_s(q)}{2\pi} \Delta_S(Q, zq) \\ \times (2N_c) \left[\frac{\Delta_{NS}(k_T, q, z)}{z} + \frac{\theta(1 - z - Q_0/q)}{1 - z} \right] f_g\left(\frac{x}{z}, |\vec{k}_T + (1 - z)\vec{q}|, q\right)$$

We use the solution of CCFM by Jung and Hansson (2003), hep-ph/0309009



KIMBER—MARTIN—RYSKIN (KMR)

The main idea is to obtain uPDFs from collinear PDF (DGLAP equation). This is often done by:

$$f_a(x, k_\perp, Q) = \frac{\partial}{\partial \ln k_\perp^2} [T_a(Q, k_\perp) D_a(x, k_\perp)]$$

where Sudakov formfactor is $T_a(Q, k_\perp) = \exp \left\{ - \int_{k_\perp^2}^{Q^2} \frac{dp_\perp^2}{p_\perp^2} \sum_{a'} \int_0^{1-\Delta} dz z P_{a'a}(z, p_\perp) \right\}$

$$\Delta(k_T, Q) = \frac{k_T}{k_T + Q} \quad (\text{angular ordering version of KMR})$$

It was shown by K. Golec-Biernat and A. Staśto (1803.06246[hep-ph]) that better prescription is:

$$F_g(x, k_T, Q) \equiv \frac{T_a(Q, k_T)}{k_T^2} \sum_{a' \in \{f, \bar{f}, g\}} \int_x^{1-\Delta(k_T, Q)} \frac{dz}{z} P_{ga'}(z, k_T) D_{a'}\left(\frac{x}{z}, k_T\right)$$

COMPARISON BETWEEN CHANNELS (FOR KMR)

$$\sigma(qg^*) = \sigma(q_{sea}g^*) + \sigma(q_{val}g^*)$$

