

Transverse momentum dependent splitting functions in the Parton Branching method

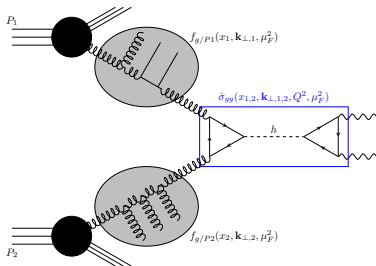
DIS 2021

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Kutak⁵, A. Lelek¹

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Motivation



- ▶ Collinear factorization is commonly used to describe hard production processes at LHC
 - ▶ Some classes of processes require more general scheme
 - ▶ Factorization in partonic cross-section and transverse momentum dependent PDFs (TMDs)
-
- ▶ TMDs much less known than PDFs at present → future experimental programs
 - ▶ TMDs from Parton Branching (PB) method: Can be used in Monte Carlo events generators → Widely applicable



Parton Branching Method

- $P_{ab}^R(z)$: (real emission part of) Splitting functions: Probability that a branching will happen
b: incoming parton, a: outgoing parton, z momentum fraction of parton a to b

- Sudakov form factor:

$$\Delta_a(\mu^2) = \exp \left(- \sum_b \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \int_0^{z_M} dz z P_{ba}^R(z) \right)$$

Interpretation: probability of an evolution without any resolvable branchings



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Iterative form of the PB evolution equations: [Hautmann, Jung, Lelek, Radescu, Zlebcik JHEP 01 (2018) 070, 1708.03279]

$$\begin{aligned} \tilde{\mathcal{A}}_a(x, \mathbf{k}_\perp, \mu^2) &= \Delta_a(\mu^2) \tilde{\mathcal{A}}_a(x, \mathbf{k}_\perp, \mu_0^2) + \sum_b \int \frac{d^2 \mu'_\perp}{\pi \mu'^2} \Theta(\mu^2 - \mu'^2) \Theta(\mu'^2 - \mu_0^2) \frac{\Delta_a(\mu^2)}{\Delta_a(\mu'^2)} \times \\ &\quad \times \int_x^{z_M} dz P_{ab}^R(z) \Delta_b(\mu'^2) \tilde{\mathcal{A}}_b\left(\frac{x}{z}, \mathbf{k}_\perp + (1-z)\mu'_\perp, \mu_0^2\right) + \dots \end{aligned}$$



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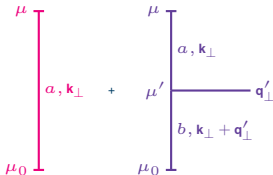
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$$\mathbf{q}'_\perp = (1 - z')\mu'$$



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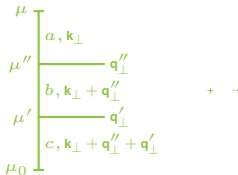
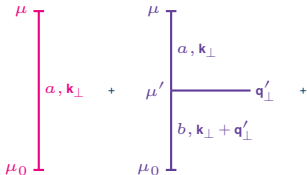
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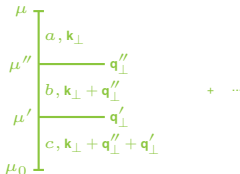
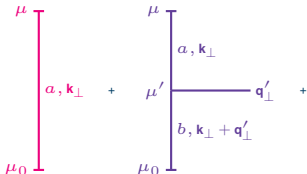
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PB calculates $\mathbf{k}_\perp$ from every branching:

$$\mathbf{k}_\perp = \mathbf{k}_{\perp,0} - \sum_i \mathbf{q}_{\perp,i}$$



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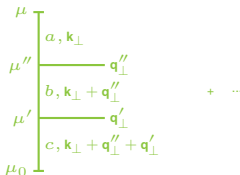
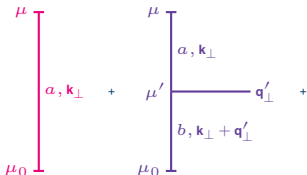
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Can be solved with MC methods.



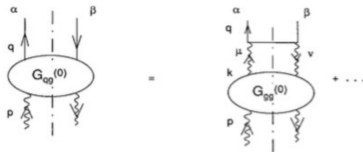
TMD Splitting functions

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- ▶ Concept from high-energy factorization [Catani, Hautmann NPB427 (1994) 475524, hep-ph/9405388]
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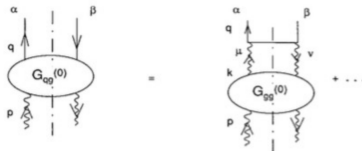
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► $\tilde{P}_{qg}(z, \mathbf{k}_{\perp}, \mathbf{q}_{\perp})$ originally calculated

► Recently other splitting functions calculated [Gutliar, Hentschinski, Kutak JHEP 01 (2016) 181, 1511.08439],

[Hentschinski, Kusina, Kutak, Serino EPJC 78 (2018) 174, 1711.04587]





TMD Splitting functions in PB

TMD Splitting functions: $\tilde{P}_{ab}^R(z, \mathbf{k}'_{\perp}, \mu'_{\perp}) = \frac{d^2 \tilde{\mathbf{k}}_{\perp}}{\tilde{\mathbf{k}}_{\perp}^2} \frac{\mu'_{\perp}{}^2}{d^2 \mu'_{\perp}} \tilde{P}_{ab}^R(z, \mathbf{k}'_{\perp}, \tilde{\mathbf{k}}_{\perp})^1$

- | | | |
|--|---|--|
| ▶ Defined within 2GI-Kernel, Jacobian included in redefinition | $\left. \begin{array}{l} k, a \\ k', b \end{array} \right \frac{q}{\quad}$ | $\begin{aligned} \mu'_{\perp} &= \frac{q_{\perp}}{1-z} \\ k' &= k + q \end{aligned}$ |
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In limit $\mathbf{k}'_{\perp} \rightarrow 0$

$$\bar{P}_{qq} = T_R(z^2 + (1-z)^2)$$

$$\bar{P}_{gq} = C_F \frac{1+(1-z)^2}{z}$$

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In DGLAP limit $\rightarrow$ DGLAP Splitting functions

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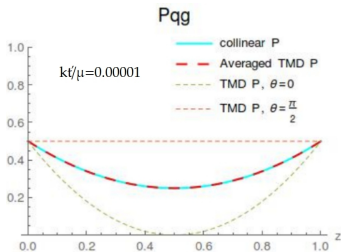
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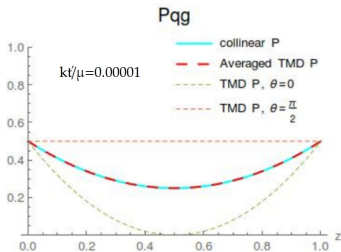
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In DGLAP limit $\rightarrow$ DGLAP Splitting functions

All $\tilde{P}_{ab}^R(z, \mathbf{k}'_{\perp}, \mu'_{\perp})$ are positive definite

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$k_{\perp}$ -dependent Sudakov form factor

Resolvable branchings:

$$P_{ab}(z) \rightarrow \tilde{P}_{ab}(z, \mathbf{k}_{\perp}, \mu'_{\perp})$$



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Sudakov form factor: probability of an evolution without any resolvable branchings

- Definition of $\mathbf{k}_{\perp}$ -dependent Sudakov form factor
- Should sum over all possible splittings $\rightarrow$ Angular averaged splitting functions

$$\Delta_a(\mu^2) = \exp \left(- \sum_b \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \int_0^{z_M} dz z P_{ba}^R(z) \right) \rightarrow$$

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$k_{\perp}$ -dependent Sudakov form factor

Resolvable branchings:

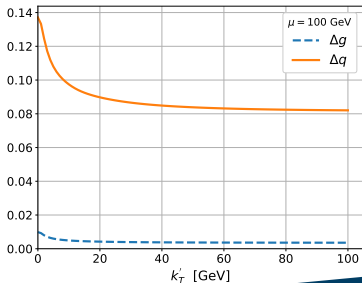
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The Sudakov form factor decrease with higher $k_{\perp}$



Studies on TMD P in evolution equations

New evolution equation:

$$\begin{aligned}\tilde{\mathcal{A}}_a(x, \mathbf{k}_\perp, \mu^2) = & \Delta_a(\mu^2, \mathbf{k}_\perp) \tilde{\mathcal{A}}_a(x, \mathbf{k}_\perp, \mu_0^2) + \sum_b \int \frac{d^2 \mu'_\perp}{\pi \mu'^2} \frac{\Delta_a(\mu'^2, \mathbf{k}_\perp)}{\Delta_a(\mu'^2, \mathbf{k}_\perp)} \Theta(\mu^2 - \mu'^2) \Theta(\mu'^2 - \mu_0^2) \times \\ & \times \int_x^{z_M} dz \tilde{P}_{ab}^R(z, \mathbf{k}_\perp + (1-z)\mu'_\perp, \mu'_\perp) \tilde{\mathcal{A}}_b\left(\frac{x}{z}, \mathbf{k}_\perp + (1-z)\mu'_\perp, \mu'^2\right)\end{aligned}$$

Studies on evolution using same starting distribution:

$$\tilde{\mathcal{A}}_a(x, \mathbf{k}_{\perp,0}, \mu_0^2) = \tilde{f}_a(x, \mu_0^2) \times \frac{1}{q_s^2 \pi} \exp\left(-\frac{\mathbf{k}_{\perp,0}^2}{q_s^2}\right), \quad \begin{array}{l} \tilde{f}_a(x, \mu_0^2): \text{HERAPDF20_LO_EIG,} \\ \mu_0 = 1.3784 \text{ GeV, } q_s = 0.5 \text{ GeV} \end{array}$$



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$$\tilde{\mathcal{A}}_a(x, \mathbf{k}_{\perp,0}, \mu_0^2) = \tilde{f}_a(x, \mu_0^2) \times \frac{1}{q_s^2 \pi} \exp\left(-\frac{\mathbf{k}_{\perp,0}^2}{q_s^2}\right), \quad \begin{array}{l} \tilde{f}_a(x, \mu_0^2): \text{HERAPDF20_LO_EIG,} \\ \mu_0 = 1.3784 \text{ GeV, } q_s = 0.5 \text{ GeV} \end{array}$$

Momentum of proton is **conserved** in equation with **TMD P** and **TMD Sudakov**
(and for col. P with col. Sudakov: the regular PB method):



Studies on TMD P in evolution equations

New evolution equation:

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$\sum_a \int_{10^{-5}}^1 dx \tilde{f}_a(\mu^2, x) :$	$\frac{\mu \text{ (GeV)}}{100}$	col. P	TMD P col. Sud.	TMD P TMD Sud.
		0.992	1.129	0.991

Collinear PDFs or integrated TMD: $\tilde{f}_a(x, \mu^2) = \int d^2 \mathbf{k}_\perp \tilde{\mathcal{A}}_a(x, \mathbf{k}_\perp, \mu^2)$

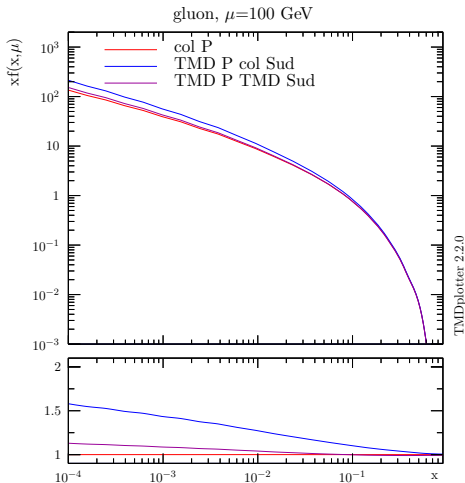


Integrated TMDs

We studied effects of TMD Splitting functions on the evolution.

No fits has been done yet:

- Not yet ready for phenomenology





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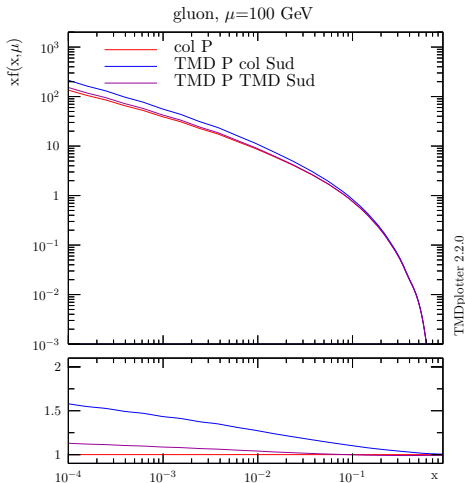
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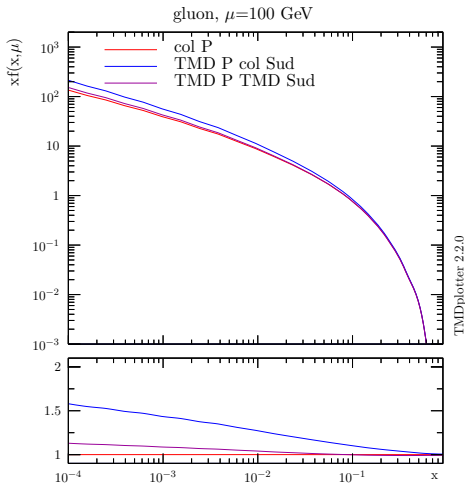
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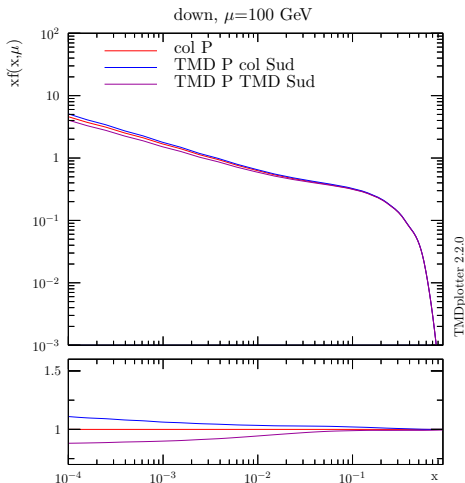
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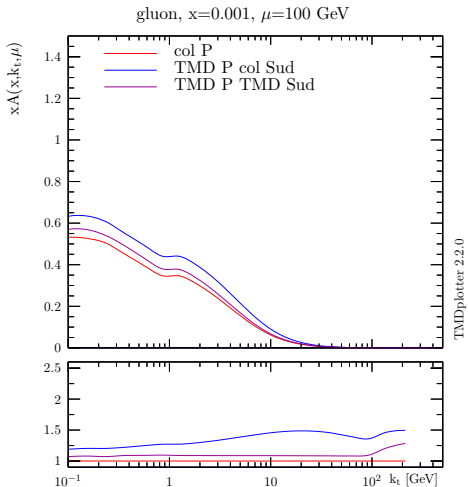
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TMDs vs $k_{\perp}$

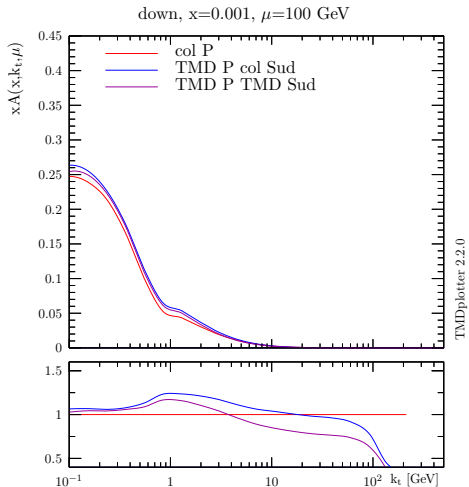
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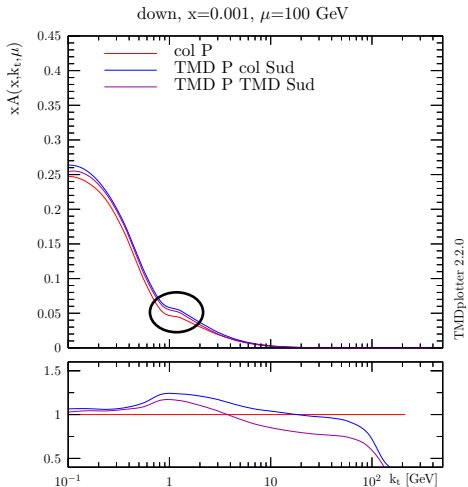




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Kink distributions $\rightarrow$ non-perturbative input





Non-perturbative input

Dynamical resolution scale:

[Hautmann, Keersmaekers, Lelek, van Kampen NuclPhysB

(2019) 114795,1908.08524]

$$z_M = 1 - \frac{q_0}{\mu'} \rightarrow q_\perp > q_0$$



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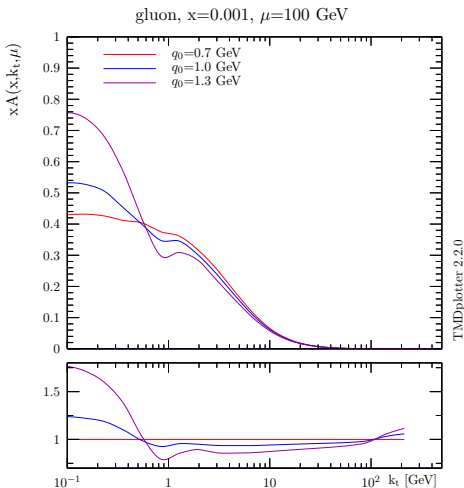
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Kink due to matching intrinsic $k_{\perp,0}$ and evolution

- ▶ $\mathbf{k}_\perp = \mathbf{k}_{\perp,0} - \sum_i \mathbf{q}_{\perp,i} \rightarrow$
 $k_\perp$ from evolution starts to build
 around q_0
- ▶ Many branchings smooth out kinks





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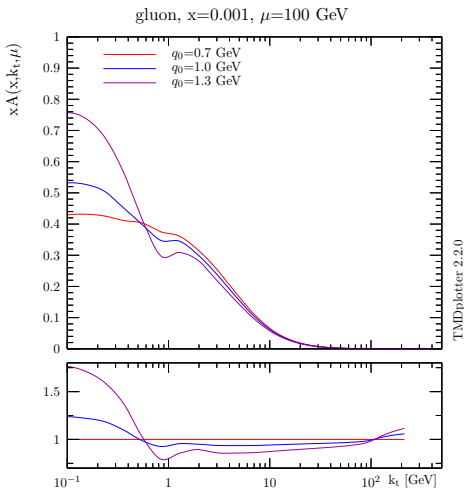
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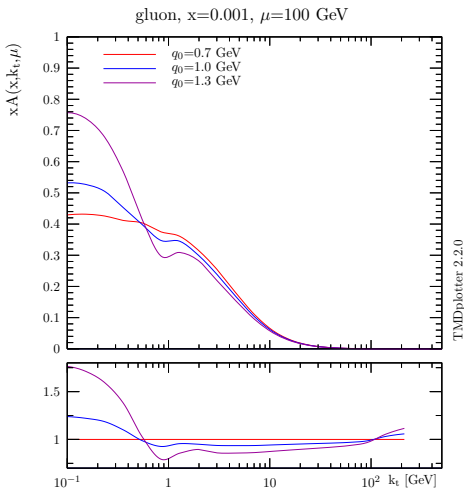
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Result shown here: collinear P

TMD P: similar results



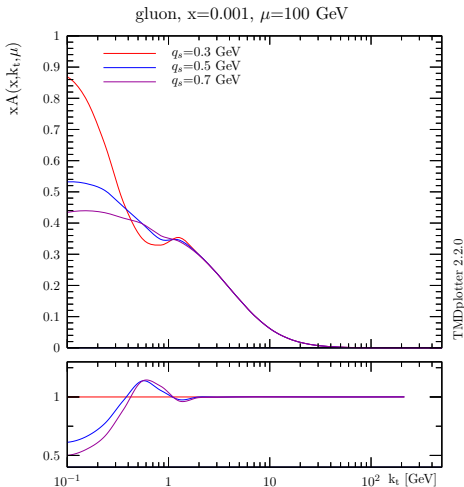


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When q_s close to q_0 : smoother

Large $k_{\perp}$ -region unaffected by q_s
when $\mu \gg \mu_0$





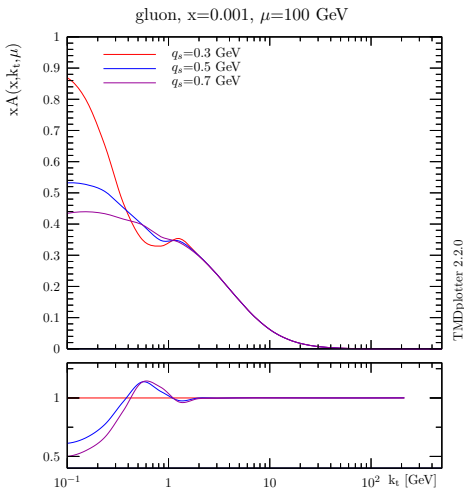
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Summary and outlook

- ▶ I presented a parton branching algorithm for space-like parton evolution with $k_{\perp}$ -dependent splitting functions
- ▶ The splitting functions are a (positive-definite) $k_T \neq 0$ continuation of the LO DGLAP splitting functions originally obtained from high-energy factorization
- ▶ $k_{\perp}$ -dependent splittings affect both real emission and Sudakov form factors
- ▶ Our evolution equations conserve the momentum of the proton
- ▶ New code is working and produces both collinear and TMD parton distributions - paper in preparation
- ▶ Prospects to do phenomenology:
 - ▶ Perform fits to DIS and DY data to determine nonperturbative TMDs
 - ▶ Use them to make PB-TMD predictions for LHC and EIC processes including for the first time the effects of TMD splittings

Thank you!