

EIC Physics Hadronization - Sec. 7.4

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$$\mathcal{P}_{A/I}(z, P_{\mathrm{T}}) = \frac{1}{2z(2\pi)^3} \int \mathrm{d}x^- \,\mathrm{d}^2 x_{\mathrm{T}} \,\mathrm{e}^{\imath k^+ x^- - \imath k_{\mathrm{T}} \cdot x_{\mathrm{T}}} \qquad \text{environment}$$

$$\times \frac{1}{3} \operatorname{tr}_{\text{color}} \frac{1}{2} \operatorname{tr}_{\text{Dirac}} \left\{ \gamma^+ \langle 0 | \psi(x) a_A(P^+, 0) a_A(P^+, 0) \bar{\psi}(0) | 0 \rangle \right\}$$



String fragmentation



Light Flavor Fragmentation





Diagram showing the three elements of pQCD factorization: parton distribution functions f a,b (x), partonic cross sections σ a+b--c , and fragmentation functions D h c (z).

- Data from electron-positron annihilation mostly con-strain the singlet combination of the FFs. Proton-proton collisions primarily constrain gluon FFs
- Production of light mesons in SIDIS primary channel for the differentiation between the fragmentation of light quarks and anti-quarks. High sensitivity to the separation of quark flavors



E. Aschenauer et al. 2019

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Global DSS analyses with the inclusion of EIC pseudo-data, at a c.m.s. energy \sqrt{s} =140 GeV

Polarized Lambda Fragmentation



 X_{F}

General questions of polarized fragmentation

A. Metz et al. 2016

Spin transfer to Λ - both longitudinal and transverse (strange quark helicity or transversity) origin of Λ Statistical Projection origin of $\overline{\Lambda}$ 0.08 Δ. 6000 18x275 10 fb⁻¹ 6000-18x275 $Q^2 > 1$ 0.01 < y < 0.95 0.06 $p_{-}^{p/\pi} > 0.1 \text{ GeV}$ lη < 3.5 4000 S 4000 S Λ 0.04 s Λ SU(6) s (40% reco eff.) di-quark A DIS di-quark HF quarks HF quarks ⊼ SU(6) 0.02 feed-down feed-down 2000 ⊼ DIS 2000 0 0 0 -0.5 0.2 0.4 0.6 0.8 -1 0 0.5 0.5 0 -0.5 -1 0 X_{F} Feynman-x

Projection of longitudinal spin transfer for A and A from proton beam at 18×275 GeV at EIC. The two righthand panels show the origin of the reconstructed Λ/Λ . In the current fragmentation re-gion a significant fraction originates from feed-down. A dominant part of the feed-down component is contributed by $\Sigma 0 \rightarrow \Lambda y$.

T-odd fragmentation-polarization fragmentation function DL1T

Callos et al. 2020



Projected Λ polarization using the extraction for the highest energy configuration. The projected uncertainty on Σ0 polarization is also shown. No feed-down which most likely reduces the magnitude of the asymmetries.

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Partial wave decomposition of polarized and unpolarized di-hadron FFs including TMDs



 $F_{UT,T}^{P_{d',m}\sin((1+m)\phi_{k}-m\phi_{k_{\perp}}-\phi_{S})} = -\mathcal{I}\left[\frac{|k_{T}|}{M}\cos(\phi_{k}+m\phi_{p}-(1+m)\phi_{h})(f_{1T}^{\perp}D_{1}^{|\ell',m\rangle} + \mathrm{sign}[m]g_{1T}G_{1}^{|\ell',m\rangle})\right]$

 $F_{UT}^{P_{\ell,m}\sin((1-m)\phi_h+m\phi_{R_\perp}+\phi_S)} = -\mathcal{I}\left[\frac{|p_T|}{M}\cos((1-m)(\phi_p-\phi_h))h_1H_1^{\perp|\ell,m\rangle}\right],$

Di-hadron FFs are more powerful than single-hadron FFs, due to the additional degrees of freedom. This allows FFs to exist that do not have a single-hadron analog *S. Gliske et al. 2015*



Projections nine partial waves contributing at twist-2 to A_{UT} using L=10 fb-1of data at 5×41. The labels on the figure indicate the m,l state and which PDF and FF the

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More sensitive at lower center of mass energies

 $F_{UT,L}^{P_{\mathscr{E},m}\sin((1+m)\phi_h-m\phi_{R_\perp}-\phi_S)}=0.$

Hadronization in nuclei





A. Accardi et al. 2009

- The space-time picture of hadronization
- Competing physics explanations based on energy loss and absorption

Light hadron measurements cannot differentiate between competing mechanisms



Attempt to parametrize nFFs assuming universality.

 A way of estimating what the differences might be coming form underlying physics
 Effect of 10 fb-1 EIC data

Sassot et al. 2009



Heavy meson tomography

 $R^h_{eA}(p_T,\eta,z) =$



Normalized by inclusive large radius jet production. To LO equivalent inclusive normalization. Idea is to eliminate nPDF effects

Heavy flavor can differentiate between energy loss and absorption models. Allows to develop theory further

The larger CM energies imply partonic interactions

Help constrain the transport properties of nuclear matter:

Z. Liu et al. 2020

2.0

$$2\frac{\mu^2}{\lambda g} = 0.12\frac{GeV^2}{fm} \quad (vary \times 2,/2)$$





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It is beneficial to look at forward rapidities and lower CM energies

https://indico.bnl.gov/event/9273/

 $N^n(p_T,\eta,z)$

 $N^{h}(p_T,\eta,z)$

Heavy meson reconstruction

- PYTHIA simulation for 10 GeV electron and 100 GeV proton collisions • with integrated luminosity: 10 fb^{-1} , 500 pb^{-1} for eA.
- X. Li et al. 2020 **Reconstructed D-meson and B-meson mass distributions** ullet



Reconstructed D[±]







Projected hadron $R_{_{eA}}$ vs $z_{_{h}}$

It is beneficial to look at forward rapidities and lower CM energies

Together with D meson back-to-back correlations to X. Dong et al. 2020 study TMD – basis for out tracking request

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Target Fragmentation





Fragmentation at large Feynman-x $e+p \rightarrow e'+X+h(x_F,p_T)$

Limited knowledge from HERA. Quantum numbers matter in the large x_F region



- x-dependence of target fragmentation: qualitative changes of the x_F distributions of p and n depending on the x of the removed parton: ∝(1-x_F) at x>0.2; constant in x_F at x~0.2; ∝1/x_F at x <<0.1
- Spin dependence and polarization transfer: Fragmentation into self-analyzing Λ baryons, use of the Collins asymmetry
- Quark vs. gluon fracture functions: how the hadronization process changes depending on whether a quark or gluon is removed from the nuclear wavefunctiopn
- Correlations of target and current fragmentation: correlations could be revealed in back-to-back pion correlations with p_T≈0.5 GeV and moderate rapidity separations∆η≈4

An important requirement is continuous coverage in x_F from~1 down to~0.1, without gaps between the central (η <4) and forward detectors

New Particle Production Mechanisms • Los Al

Pomerons – quantum numbers of the vacuum J^{PC}=0⁺⁺

Predominantly produce vector mesons in the final state

- Reggeons summed meson trajectories carry a wider range of quantum numbers
- Odderons three gluon color singlet exchanges

Mandatory to explain the difference between p-p and pbarp scattering

Forward Backward P. R ? Q^2 Q^2 forward backward hard L, TTDA GPD GPD 0 θ (CM)

TOTEM collaboration may have observed the Odderon (2017)

At the EIC $- f_2(1270)$ via photon-Odderon fusion

U-channel exclusive vector mesons -omega

$$\frac{d\sigma}{dt} = A(s/1\text{GeV})^B \exp(-Ct)$$

$$A \approx 18 \mu b/\text{GeV}^2 \text{ and } B = -1.92 \text{ C} \approx 10 \text{ GeV}^{-2}$$

$$\frac{d\sigma}{du} = A(s/1\text{GeV})^B \exp(-Cu)$$

$$A \approx 4 4 \mu b/\text{GeV}^2 \text{ and } B = -1.92 \text{ C} \approx 21 \text{ GeV}^{-2}$$

The decay $\omega \rightarrow \pi^0 \gamma$ could be studied with farforward calorimetry.

U-channel electroproduction - π 0, π ±, ρ , η , η and ϕ

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 $\frac{d\sigma}{dt}$

Quarkonia and exotics

 $\mathcal{L}_{\mathrm{NRQCD}_G} = \mathcal{L}_{\mathrm{NRQCD}} + \mathcal{L}_{Q-G/C}(\psi, A_{G/C}^{\mu,a}) + \mathcal{L}_{g-G/C}(A_s^{\mu,b}, A_{G/C}^{\mu,a}) + \psi \leftrightarrow \chi$



Illustrative examples of quarkonium production mechanism in ep and eA colliders: (a) Direct photo/lepto-production, (b) resolved-photon quarkonium production,(c) exclusive quarkonium production, and (d) heavy quark pair production and subsequent Glauber/Coulomb gluon exchanges with nuclear matter

New exotic state structure studies at the EIC

Use eA collisions – nucleus as a filter to differentiate between tightly bound (quark) and molecular states



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M. Durham et al. 2020



Y. Makris *et al. 2020*

Constrain LDMEs

Understand TMD production of quarkona (and shape functions)

Develop EFTs of quarkonia in matter (indications from HI collisons)

Open quantum systems approach at EIC



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Spectroscopy



M. Albaladejo et al. 2020

Tetraquark states X, Z ...; Pentaquark states P ...

- Photoproduction through photon-Pomeron fusion lead predominantly to J^{PC}=1⁻⁻ states like the J/ψ, etc., so is only sensitive to exotic with those quantum numbers
- Photon-Reggeon fusion leads to states with a wider range of spin, parity and even charge



Z⁺ photoproduction rapidity distributions and integrated cross section predictions for fixed-spin exchange, valid at low energies (center), and for Regge exchange, valid at high energies (right) JLab energies - u-channel production, Expect for EIC S. Kein *et al.* 2019 t-channel production

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Jets



 The idea behind inclusive jets was of course to minimize hadronization effects

At the EIC jet energies are relatively low and hadronization effects are expected to play a significant role

Hadronization is important when we look inside jets: shapes, fragmentation functions, angularities

Heavy flavor jet tagging and sub-structure





L. Cunqueiro et al. 2020

$$\tau_a \equiv \tau_a^{pp} \equiv \frac{1}{p_T} \sum_{i \in J} p_T^i \left(\Delta \mathcal{R}_{iJ} \right)^{2-a}$$

 $E_{\rm e}{\rm :}~10~GeV$ $E_{\rm p}{\rm :}~100~GeV$ Integrated Lumi: 10 fb^-1

Have been studies experimentally

Can be used to differentiate between light and heavy flavor jets

P. Wong *et al. 2020*

Energy loss of jets in e+A is important but in a different section

In place of conclusions



The chapter on hadronization covers an impressive number of topics

Thanks to all the PWG members who contributed to chapter 7.4 !