

Lattice study of the phase structure of the strong Yukawa model

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The collaboration

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- Japan
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Kei-Ichi Nagai.
- Taiwan
 - National Chiao-Tung University, Hsinchu
C.-J.David Lin, Kenji Ogawa.
 - National Taiwan University, Taipei
George W.-S. Hou, Bastian Knippschild, Brian Smigielski (→ U.S.).

Motivation

Heavy fermions beyond SM3?

- Not much is known for strong (*non-perturbative*) Yukawa theory.
- Heavy extra generation of fermions may
 - enhance CP violation.

G.W.S. Hou, 2008

- offer an alternative way to break EW symmetry dynamically and induces bound states to unitarise WW scattering.

B. Holdom, 2007

- UV stabilise the SM.

P.Q. Hung, C. Xiong, 2009

Outline

- Goals, general issues and recent developments.
- Simulation setup.
- The phase structure.
- Exploratory numerical studies.
 - VEV.
 - Susceptibility and critical exponents.
 - Binder's cumulant.
- Future plan.

Targets for the bare strong-Yukawa regime

- The nature of the phase transitions.
 - ⇒ Connection to the continuum world (next slide).
- Possible bound states.
 - ⇒ Computation of the spectrum.
- Possible new mechanism for dynamical symmetry breaking.
 - ⇒ Heavy scalar with fermion condensate?

General issues and strategy

- The triviality (Landau-pole) problem.
 - ⇒ Non-trivial to take the lattice spacing to zero.
- Look for 2nd-order phase transitions via "scanning simulations".
 - ⇒ $\xi \rightarrow \infty$.
- Problem: Finite-volume effects.
 - ⇒ Phase transitions are washed out.
 - ⇒ Severe near the critical points since $L = \hat{L}a$.
- Chiral fermions required.

New ingredients in current work

- Previous studies (*circa* 1990):
Lee, Shigemitsu, Shrock; Bock *et al.*,...
 - Use fermions without exact chiral symmetry.
⇒ Ambiguity in defining chiral fermions.
 - Small ($\sim 8^3 \times 16$) volumes and no $L \rightarrow \infty$ limit taken.
- Current new-generation simulations:
 - Use the overlap fermion (exact chiral symmetry).
 - Several large volumes and $L \rightarrow \infty$ limit taken.
⇒ Test finite-size scaling behaviour.
⇒ Determine the order of the phase transition.

Reminder: Notation for scalar field theory

- The discretised Euclidean scalar action ($a = 1$)

$$S_\varphi = - \sum_{x,\mu} \varphi_x^\alpha \varphi_{x+\hat{\mu}}^\alpha + \sum_x \left[\frac{1}{2} (2d + m_0^2) \varphi_x^\alpha \varphi_x^\alpha + \frac{1}{4} \lambda_0 (\varphi_x^\alpha \varphi_x^\alpha)^2 \right].$$

- $\varphi = \sqrt{2\kappa}\phi$, $m_0^2 = \frac{1-2\hat{\lambda}}{\kappa} - 2d$, $\lambda_0 = \frac{\hat{\lambda}}{\kappa^2}$

$$S_\phi = -2\kappa \sum_{x,\mu} \phi_x^\alpha \phi_{x+\hat{\mu}}^\alpha + \sum_x [\phi_x^\alpha \phi_x^\alpha + \hat{\lambda}(\phi_x^\alpha \phi_x^\alpha - 1)^2],$$

$$Z_\phi = \int \prod_{x,\alpha} d\phi_x^\alpha \exp(-S_\phi) = \int \prod_{x,\alpha} d\mu(\phi_x^\alpha) \exp \left(2\kappa \sum_{x,\mu} \phi_x^\alpha \phi_{x+\hat{\mu}}^\alpha \right),$$

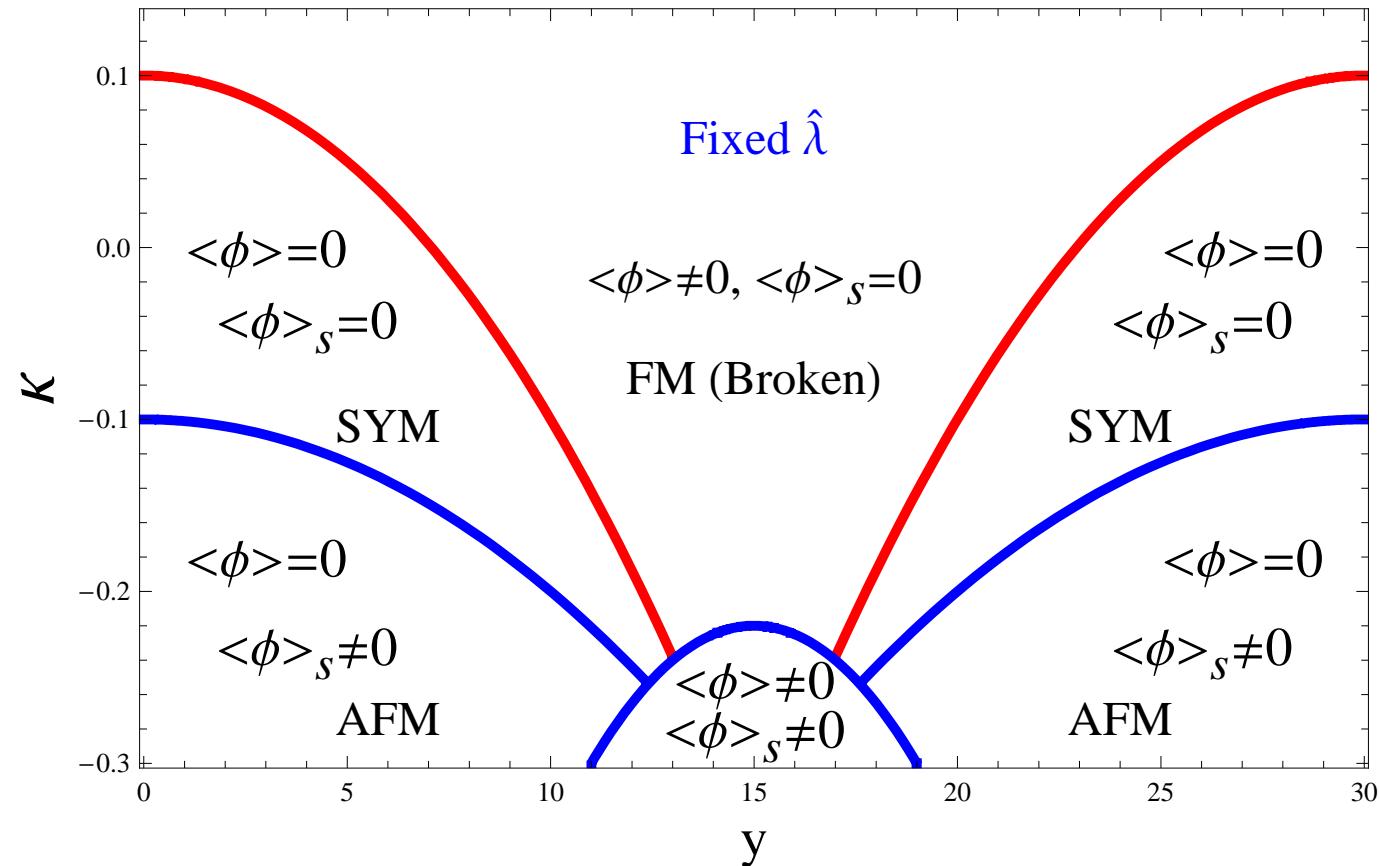
$$d\mu(\phi_x^\alpha) = d\phi_x^\alpha \exp [-\phi_x^\alpha \phi_x^\alpha - \hat{\lambda}(\phi_x^\alpha \phi_x^\alpha - 1)^2].$$

- “staggered symmetry”: $\kappa \rightarrow -\kappa$ and $\phi_x^\alpha \rightarrow (-1)^{x_1+x_2+\dots+x_d} \phi_x^\alpha$.

Fermions and the Yukawa couplings

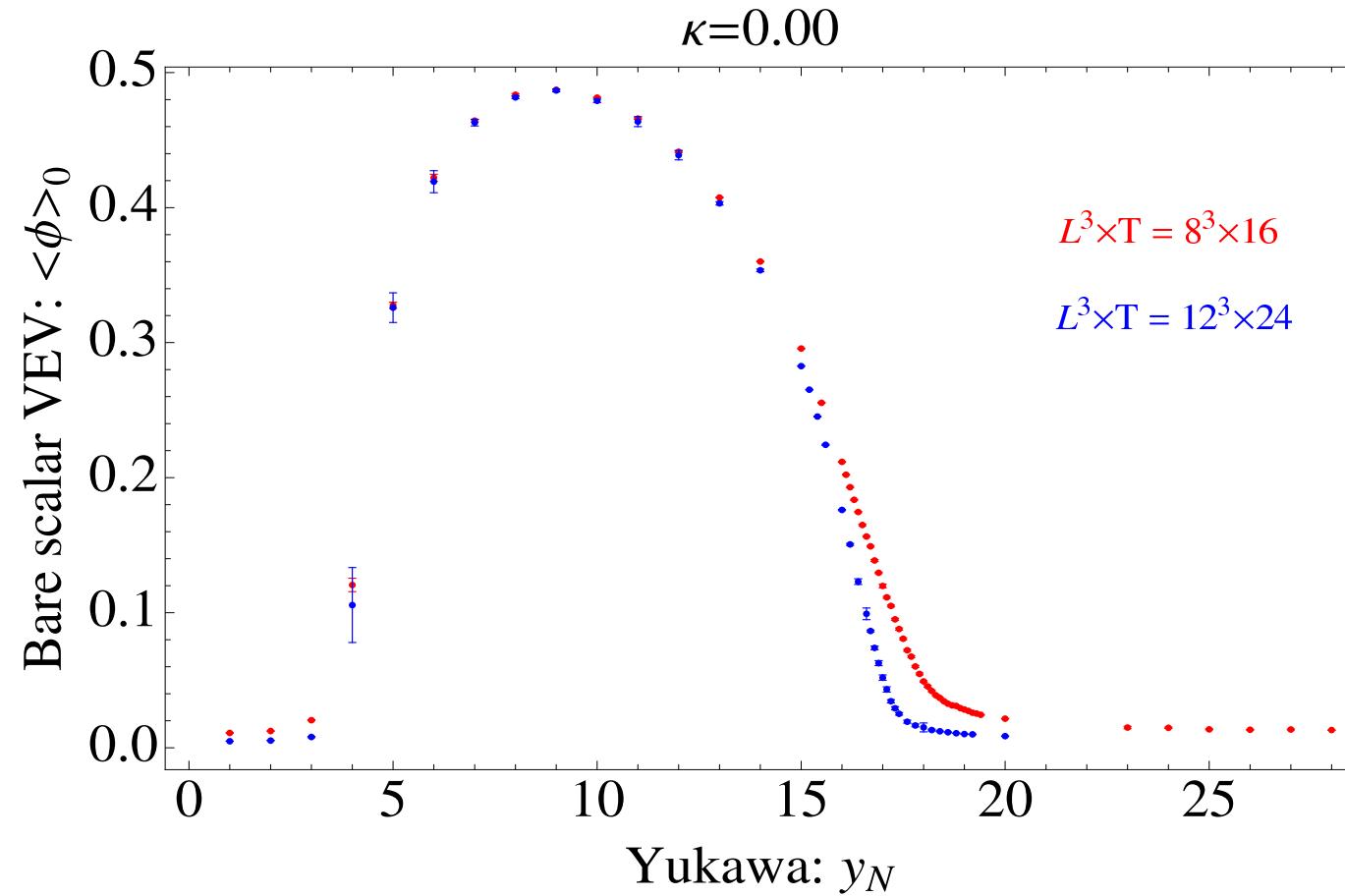
- Use the overlap Dirac operator with exact lattice chiral symmetry.
- The Yukawa terms $S_{HY} = \sum_x \textcolor{blue}{y}(\bar{t}_x, \bar{b}_x)_L \Phi_x b_{x,R} + \textcolor{blue}{y}(\bar{t}_x, \bar{b}_x)_L \tilde{\Phi}_x t_{x,R} + \text{h.c.}$.
 - Φ is a complex scalar doublet and $\tilde{\Phi} = i\tau_2 \Phi^*$.
- Results presented in this talk are from $8^3 \times 16$, $12^3 \times 24$ and $16^3 \times 32$.

Phase diagram of the H-Y model (qualitative)



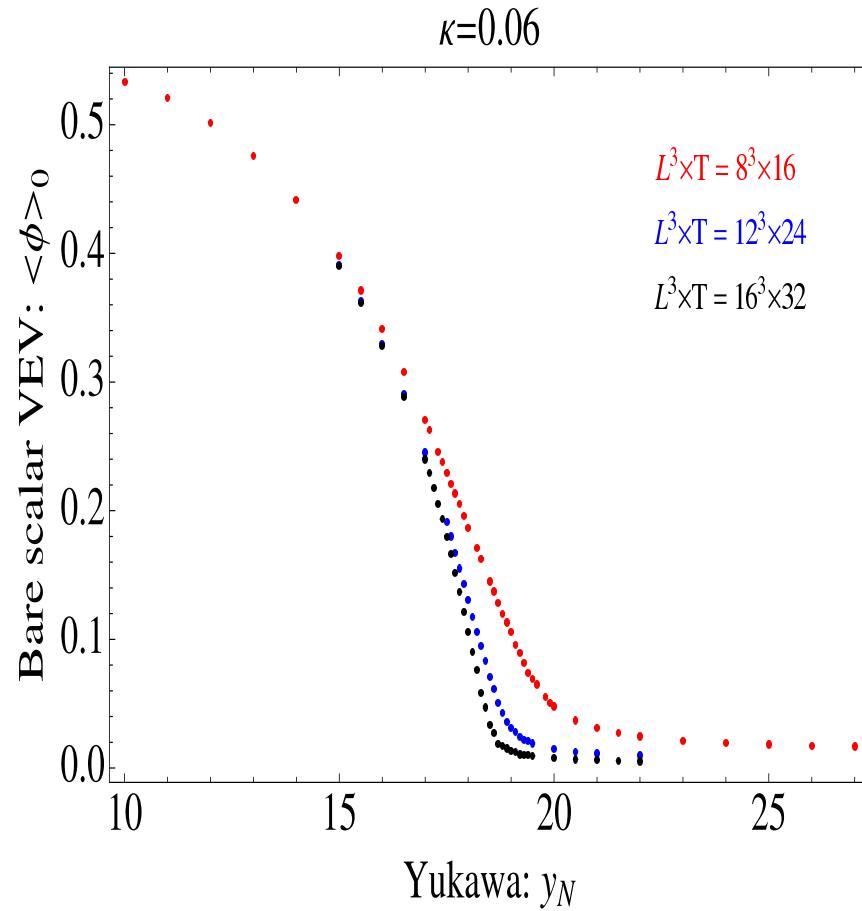
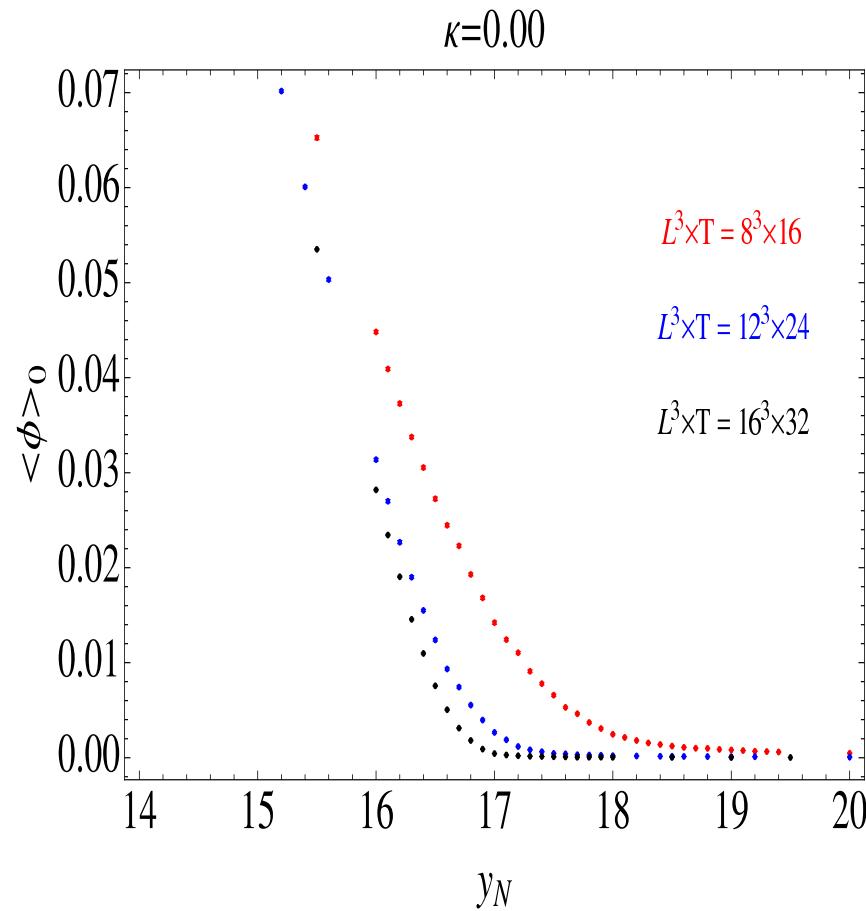
- * From earlier work using Wilson fermions.
- ⇒ Controversy from staggered-fermion calculations.

Evidence of a symmetric phase at large y



Consistent with recent results in P. Gerhold and K. Jansen, 2007.

The bare scalar vev at large Y



Finite-size scaling of susceptibility

- Susceptibility: $\chi = V_4 (\langle \phi^2 \rangle - \langle \phi \rangle \langle \phi \rangle)$.
- The scaling behaviour from solving the RGE,
 - Universal function $\chi L_s^{-\gamma/\nu} \sim g(\tilde{t} L_s^{1/\nu})$, where $\tilde{t} = (y/y_{\text{crit}} - 1)$.
 - critical exponents γ and ν .
 - Modelling the scaling violation from

M. Fisher and M. Barber, 1972

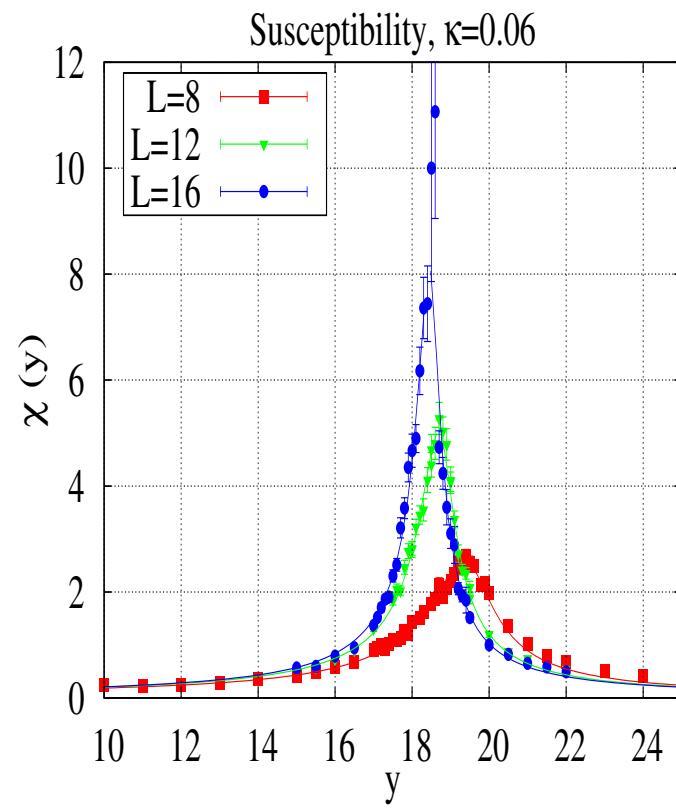
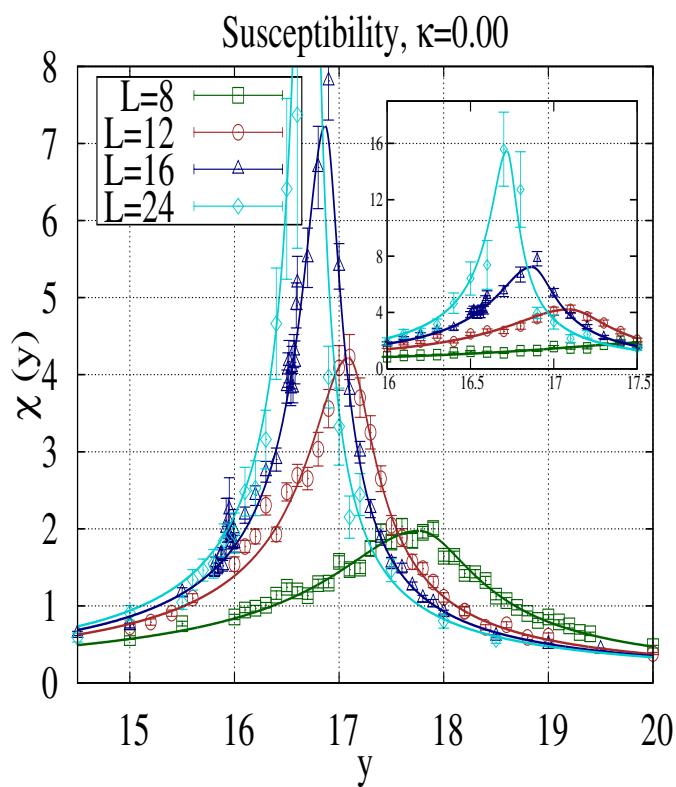
$$\Rightarrow \chi L_s^{-\gamma/\nu} \sim g(t L_s^{1/\nu}), \text{ where } t = (y/(y_{\text{crit}} - A_4/L_s^b) - 1).$$

- Fit all the data to the (partly empirical) function at fixed κ

K. Jansen and P. Seuferling, 1990

$$\chi = A_1 \left\{ L_s^{-2/\nu} + A_{2,3} (y - y_{\text{crit}} - A_4/L_s^b)^2 \right\}^{-\gamma/2}.$$

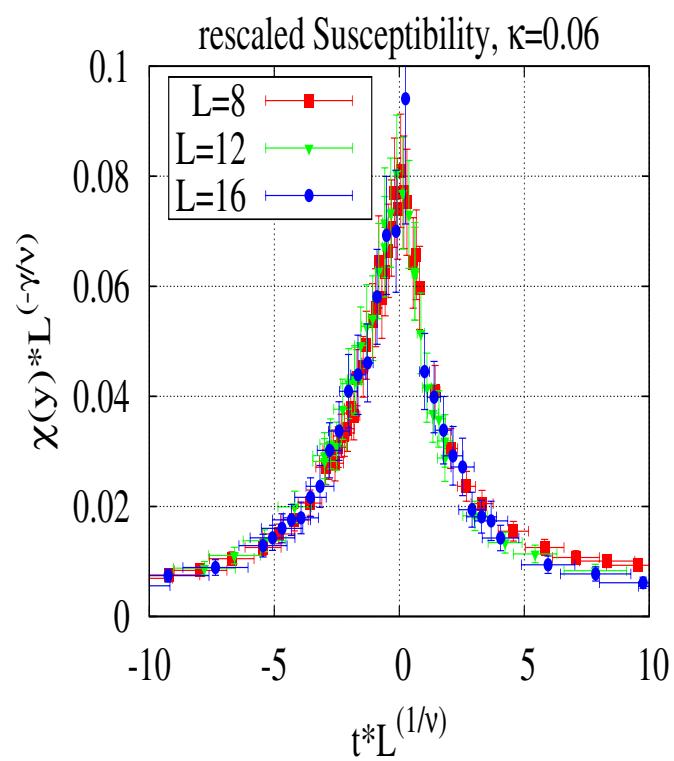
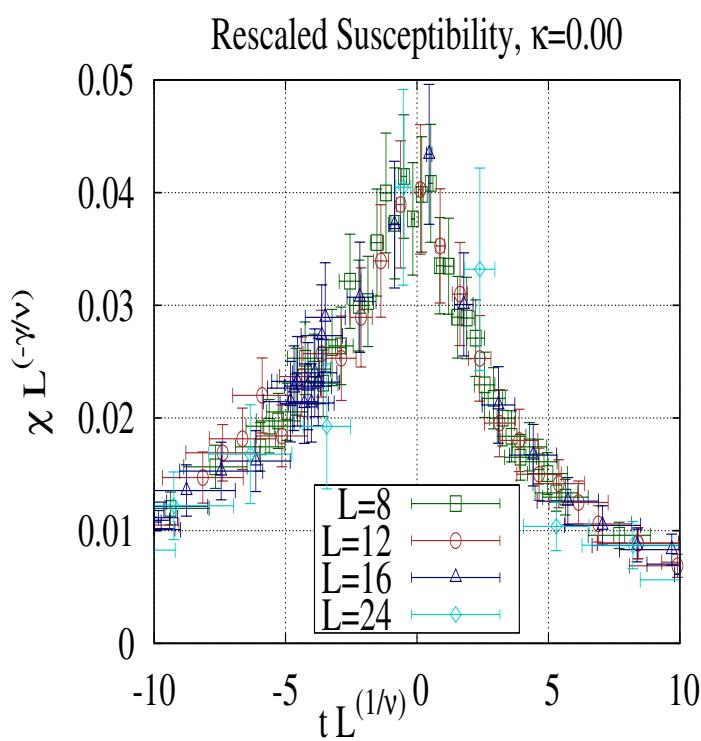
Finite-size fit of susceptibility



Fit range: $y = 15.0 \sim 20.0$

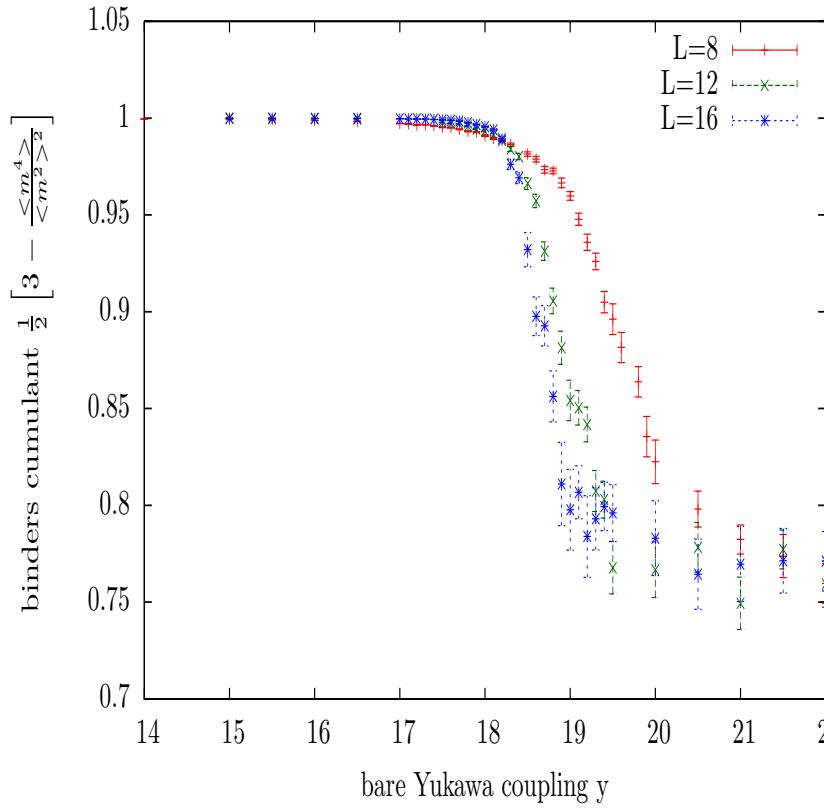
$y = 14.0 \sim 22.0$

Finite-size scaling of susceptibility

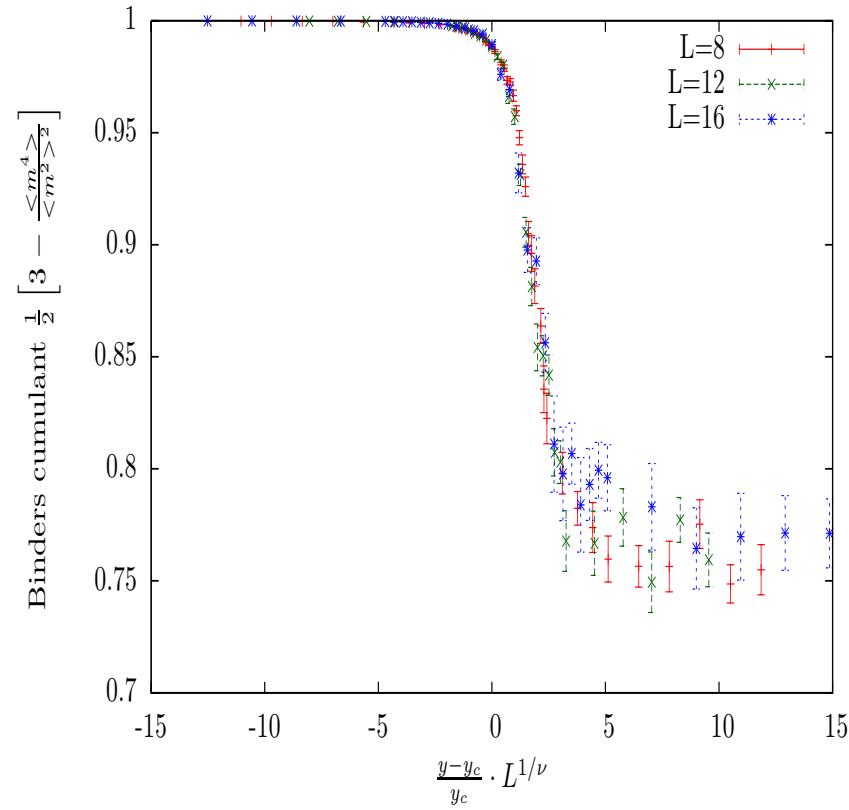


Finite-size scaling of Binder's cumulant

unscaled Binders cumulant for $\kappa = 0.06$



Rescaled Binders cumulant, $\nu = 0.650000, y_c = 18.200000$



Probing the phase structure using susceptibility

	$\kappa = 0.00$	$\kappa = 0.06$	O(4) scalar model
y_{crit}	16.57 ± 0.06	18.11 ± 0.06	N/A
γ	0.97 ± 0.02	1.08 ± 0.01	1
ν	0.52 ± 0.01	0.66 ± 0.02	0.5
b	2.18 ± 0.09	2.04 ± 0.20	?

- Quoted errors are statistical.
- Estimate systematics by changing the fit range in y .
- Systematic effects
 - y_{crit} is very stable.
 - γ can change by $\sim 2\%$.
 - ν can vary by $\sim 8\%$. \Rightarrow Different from O(4) scalar model?

Outlook

- Improving results by
 - running more at large lattices, $24^3 \times 48$. (finishing soon.)
 - using more sophisticated procedure to investigate the susceptibility.
 - studying the details of the scaling behaviour of Binder's cummulant.
- More information:
 - Compute three renormalised couplings to “trade” with κ , $\hat{\lambda}$ and y .
 - Study the spectrum in the strong Yukawa regime.

A lot more to do and to understand.