

6th of February 2015 — MNME 2015 (BNL, USA)

$K \rightarrow \pi l^+ l^-$  decays  
on the lattice

UNIVERSITY OF  
Southampton

Antonin J. Portelli  
(RBC-UKQCD collaborations)

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- ❖  $K^+ \rightarrow \pi^+ \ell^+ \ell^-$ : long-distance dominated
- ❖  $K_{L/S}^0 \rightarrow \pi^0 \ell^+ \ell^-$ : feature indirect / direct CP-violation interference

- ❖ Euclidean formulation
- ❖ Ultraviolet & infrared behaviour
- ❖ Preliminary lattice results
- ❖ Summary & perspectives

# Euclidean formulation

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# Minkowski amplitude

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$$\mathcal{A}_\mu^c(q^2) = \int d^4x \langle \pi^c(\mathbf{p}) | \mathbf{T}[J_\mu(0)H_W(x)] | K^c(\mathbf{k}) \rangle$$

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Spectral representation:

$$\begin{aligned} \mathcal{A}_\mu^c(q^2) = & i \int_0^{+\infty} dE \frac{\rho(E)}{2E} \frac{\langle \pi^c(\mathbf{p}) | J_\mu(0) | E, \mathbf{k} \rangle \langle E, \mathbf{k} | H_W(0) | K^c(\mathbf{k}) \rangle}{E_K(\mathbf{k}) - E + i\varepsilon} \\ & - i \int_0^{+\infty} dE \frac{\rho_S(E)}{2E} \frac{\langle \pi^c(\mathbf{p}) | H_W(0) | E, \mathbf{p} \rangle \langle E, \mathbf{p} | J_\mu(0) | K^c(\mathbf{k}) \rangle}{E - E_\pi(\mathbf{p}) + i\varepsilon} \end{aligned}$$

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# Euclidean correlation function

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$$\Gamma_{\mu}^{(4)c}(x, \mathbf{k}, \mathbf{p}) = \langle \phi_{\pi^c}(t_{\pi}, \mathbf{p}) \mathsf{T}[J_{\mu}(0) H_W(x)] \phi_{K^c}(t_K, \mathbf{k})^{\dagger} \rangle$$

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For  $-t_{\pi}, t_K \rightarrow +\infty$ :

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pion and kaon interpolating operators

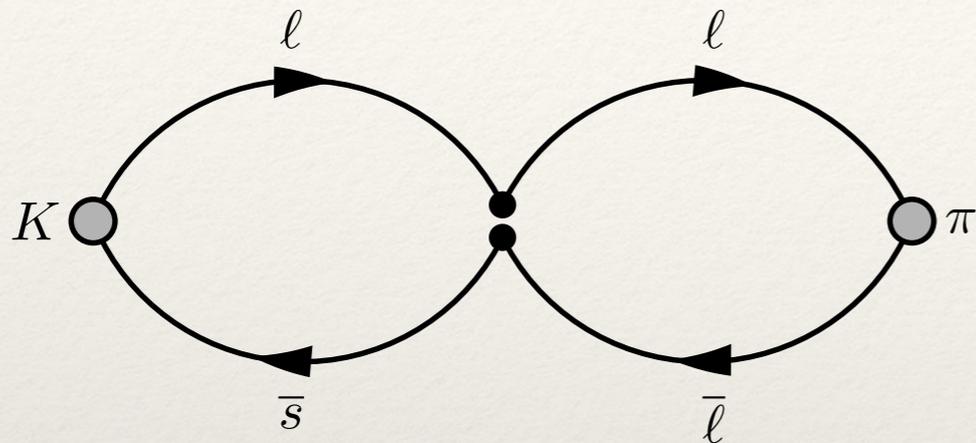
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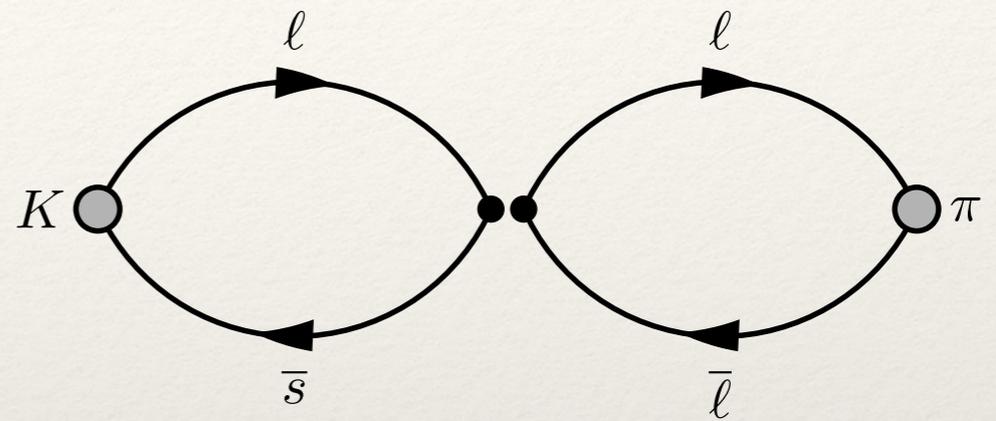
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$$\tilde{\Gamma}_{\mu}^{(4)c}(x, \mathbf{k}, \mathbf{p})$$

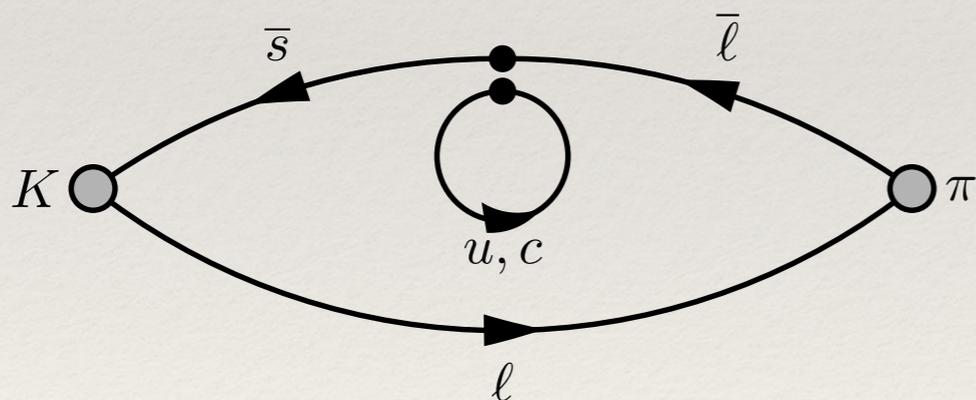
# Quark Wick contractions



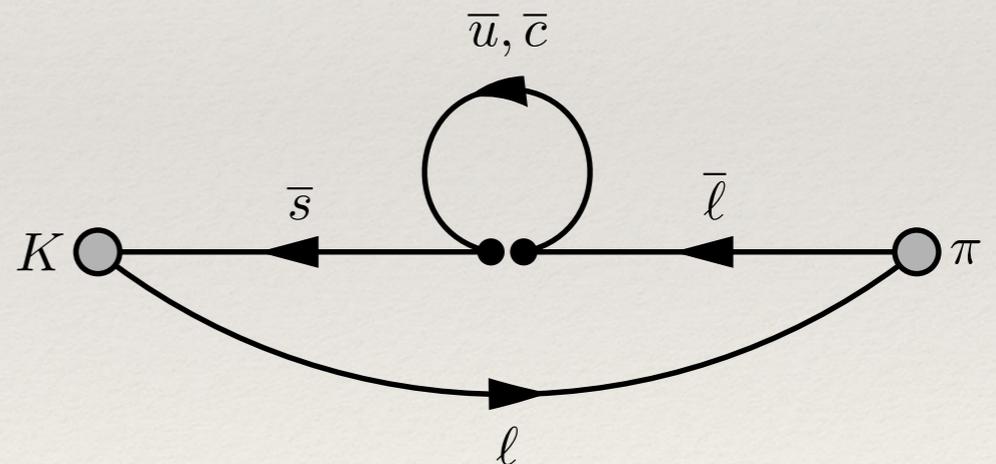
C: "Connected"



W: "Wing"

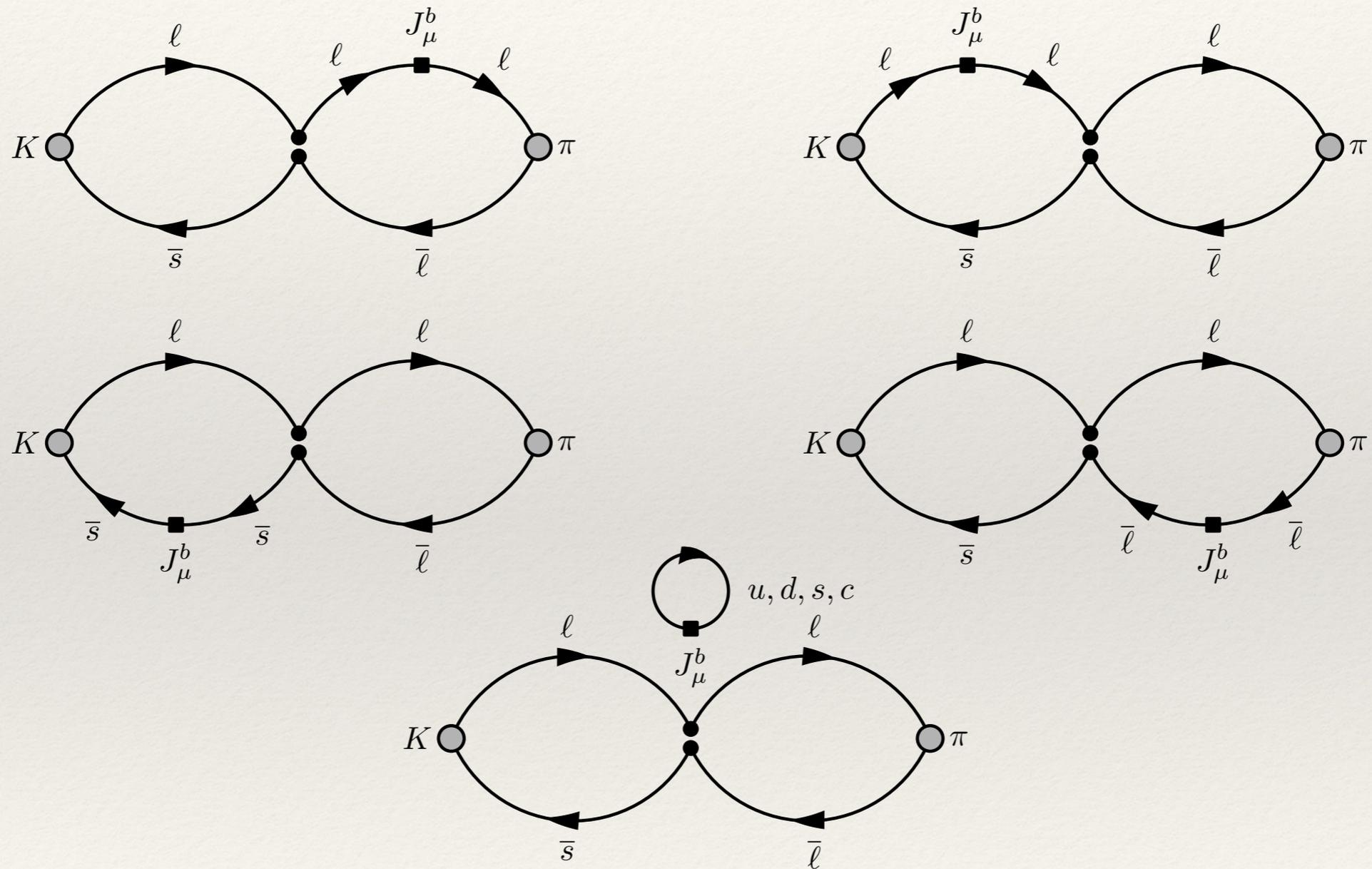


E: "Eye"



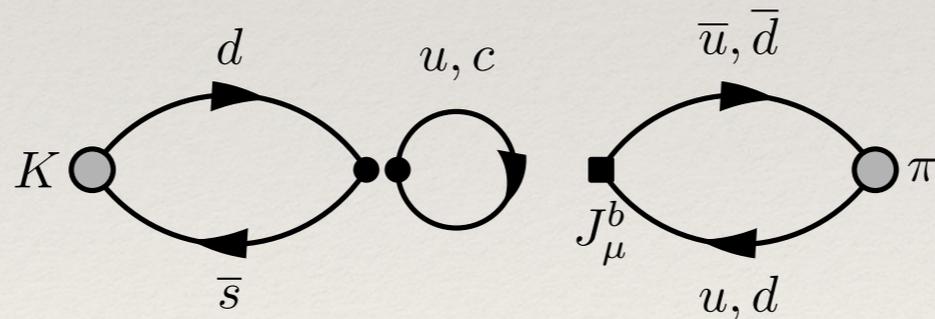
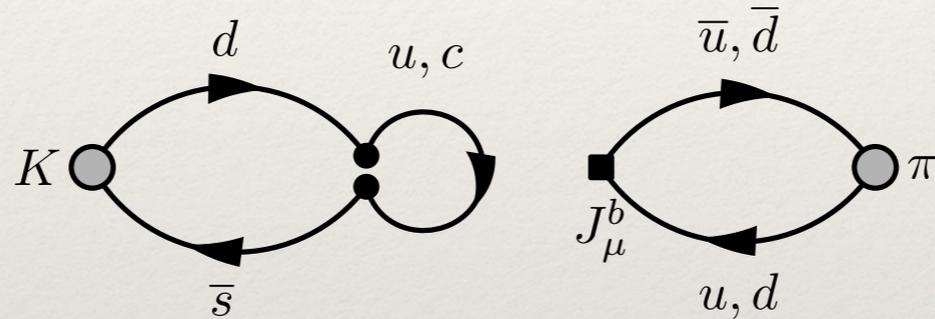
S: "Saucer"

# Quark Wick contractions



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Neutral case additional diagrams:



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# Euclidean spectral representation

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Integrated correlator on a finite time interval  $[-T_a, T_b]$ :

$$\begin{aligned} \int d^3\mathbf{x} \int_{-T_a}^{T_b} dt \tilde{\Gamma}_{\mu}^{(4)c}(t, \mathbf{x}, \mathbf{k}, \mathbf{p}) &= - \int_0^{+\infty} dE \frac{\rho(E)}{2E} \frac{\langle \pi^c(\mathbf{p}) | J_{\mu}(0) | E, \mathbf{k} \rangle \langle E, \mathbf{k} | H_W(0) | K^c(\mathbf{k}) \rangle}{E_K(\mathbf{k}) - E} \\ &\quad \times (1 - e^{[E_K(\mathbf{k}) - E]T_a}) \\ &+ \int_0^{+\infty} dE \frac{\rho_S(E)}{2E} \frac{\langle \pi^c(\mathbf{p}) | H_W(0) | E, \mathbf{p} \rangle \langle E, \mathbf{p} | J_{\mu}(0) | K^c(\mathbf{k}) \rangle}{E - E_{\pi}(\mathbf{p})} \\ &\quad \times (1 - e^{-[E - E_{\pi}(\mathbf{p})]T_b}) \end{aligned}$$

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- ❖ growing exponential for  $E < E_K(\mathbf{k})$

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- ❖ growing exponential for  $E < E_K(\mathbf{k})$
- ❖ need to be removed to obtain the Minkowski amplitude
- ❖ generated by 1, 2 and 3-pion intermediate states

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# Removing the single-pion divergence

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1. Reconstruct the divergent single-pion term by computing  $J_\mu$  and  $H_W$  matrix elements for  $\pi \rightarrow \pi\gamma^*$  and  $K \rightarrow \pi$  transitions

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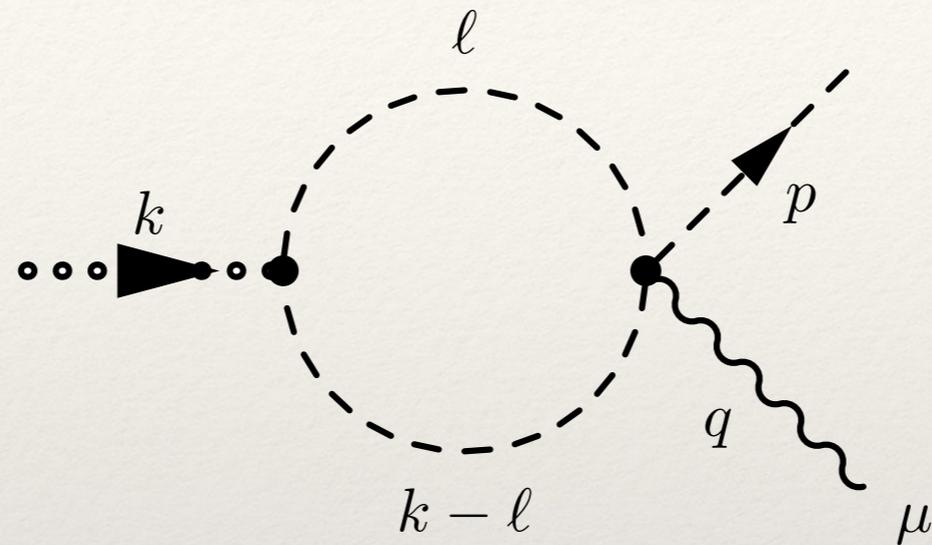
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1. Reconstruct the divergent single-pion term by computing  $J_\mu$  and  $H_W$  matrix elements for  $\pi \rightarrow \pi\gamma^*$  and  $K \rightarrow \pi$  transitions
2. One can show that the physical amplitude is invariant under  $H_W \mapsto H_W + c_S \bar{s}d$ ,  $c_S$  can be tuned to cancel the  $K \rightarrow \pi$  matrix element

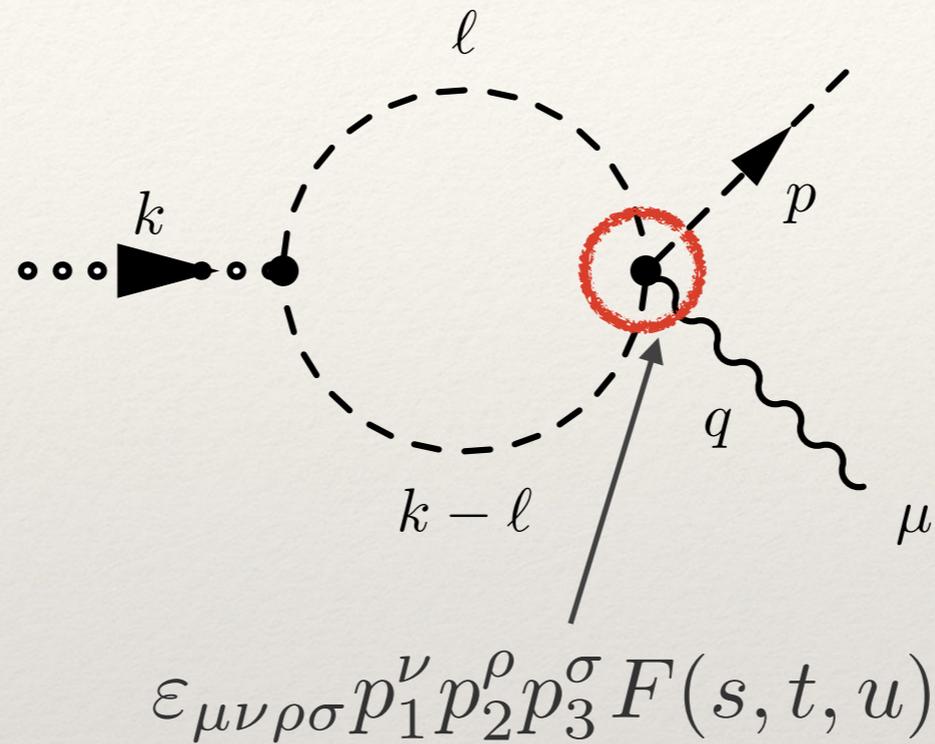
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# Two-pion intermediate states

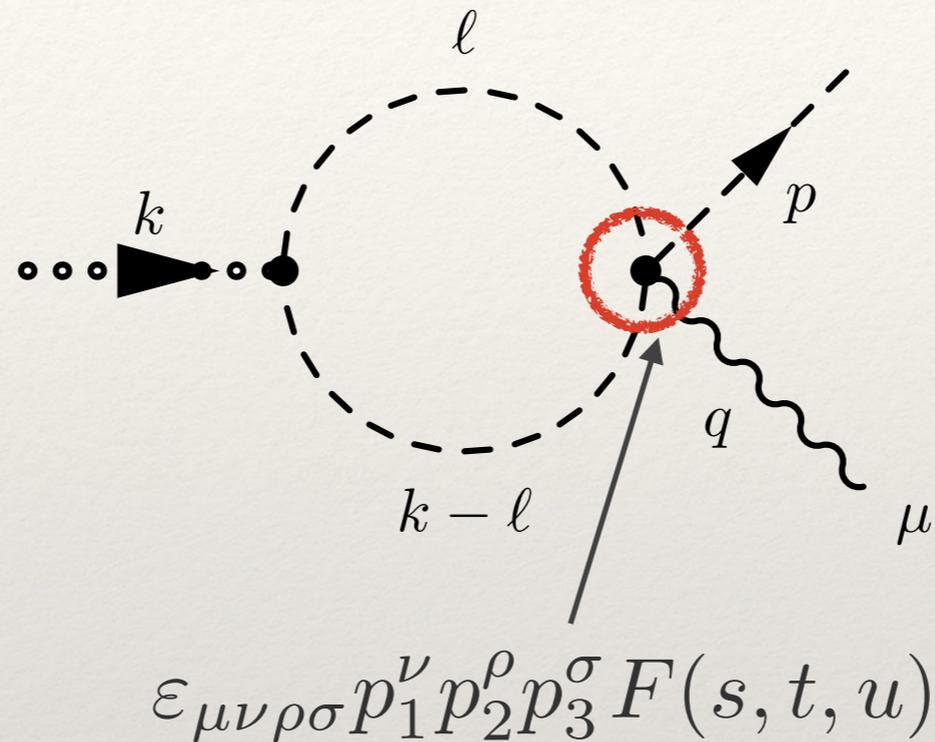
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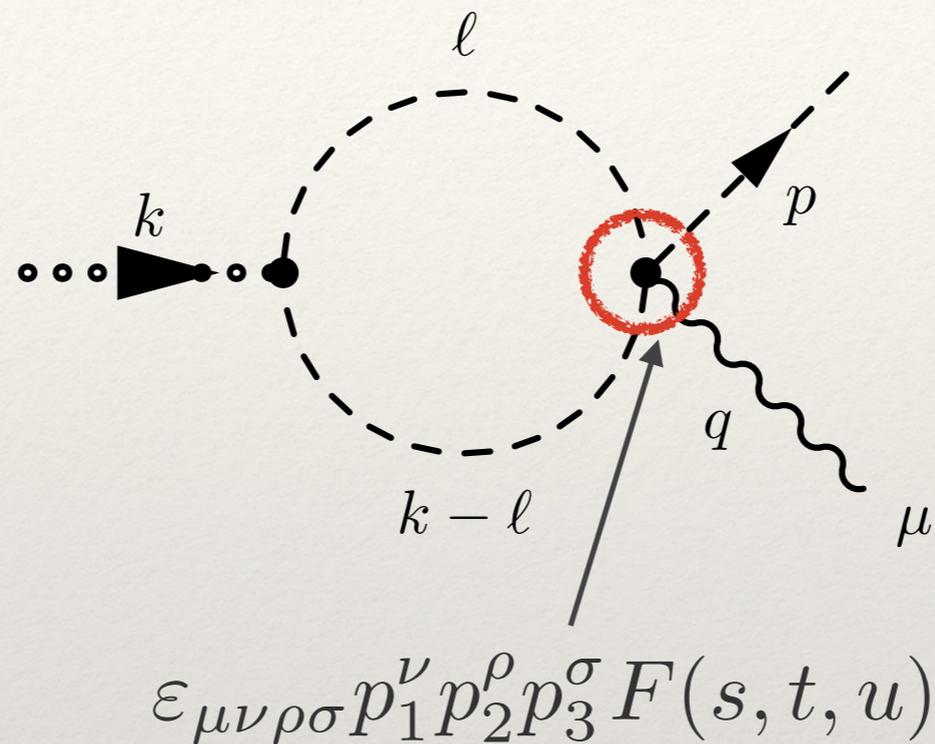


# Two-pion intermediate states



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**No two-pion intermediate state**

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- ❖ Only a problem for pion masses less than  $\sim 165$  MeV

# Ultraviolet & infrared behaviour

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# Individual operator renormalisation

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- ❖ The vector current is conserved and does not need renormalisation

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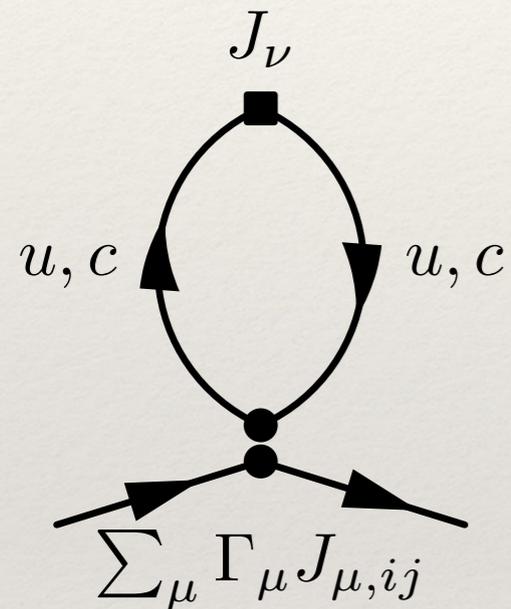
# Individual operator renormalisation

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- ❖ The vector current is conserved and does not need renormalisation
- ❖ The renormalisation of the weak hamiltonian is also known and is much more simple with chiral fermions (*cf. e.g.* [Z. Bai, *et al.* PRL, 113(1), p. 112003, 2014]).

# Short distance operator product

UV divergences may appear in loops between  $J_\mu$  and  $H_W$  :

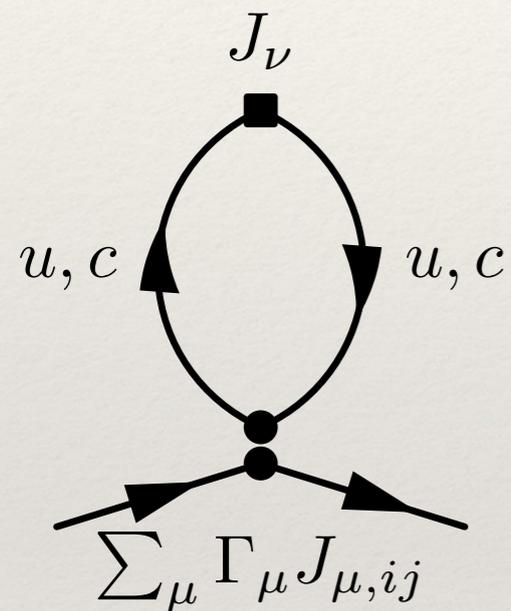


$$= \sum_\nu \Gamma_\nu [\Pi_{\mu\nu,ij}^u(q) - \Pi_{\mu\nu,ij}^c(q)]$$

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Same divergence structure than HVP

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OPE with lattice regularisation:

$$\Pi_{\mu\nu,ij}^f(q) \underset{a \rightarrow 0}{=} C_{\mu\nu,ij}^1 + C_{\mu\nu,ij}^{\bar{f}f} \langle m_f \bar{f} f \rangle + \dots$$

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- ❖ GIM subtraction cancels mass independent divergences

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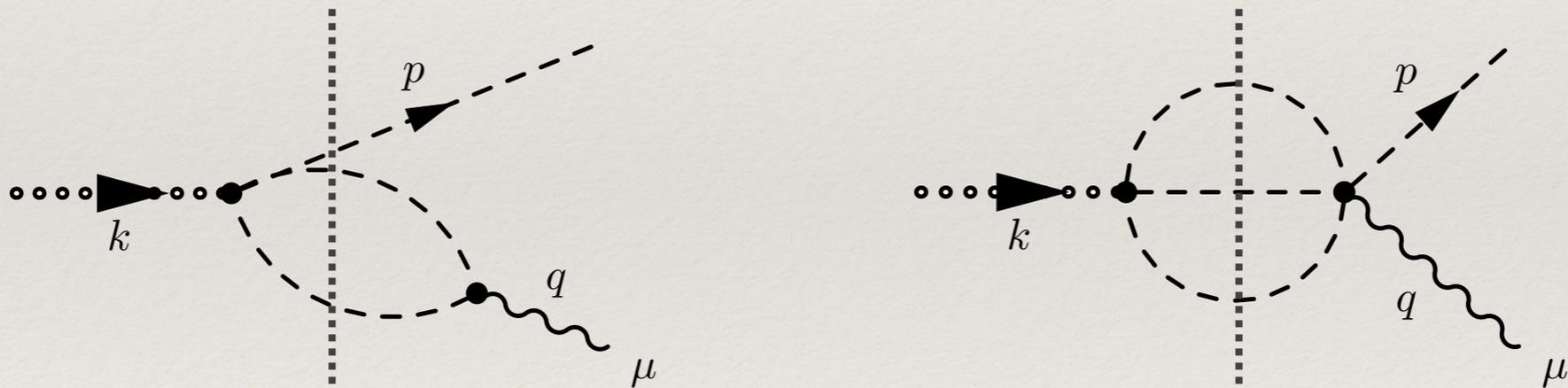
# Finite-size effects

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(*cf. e.g. S. Sharpe's talk yesterday*)

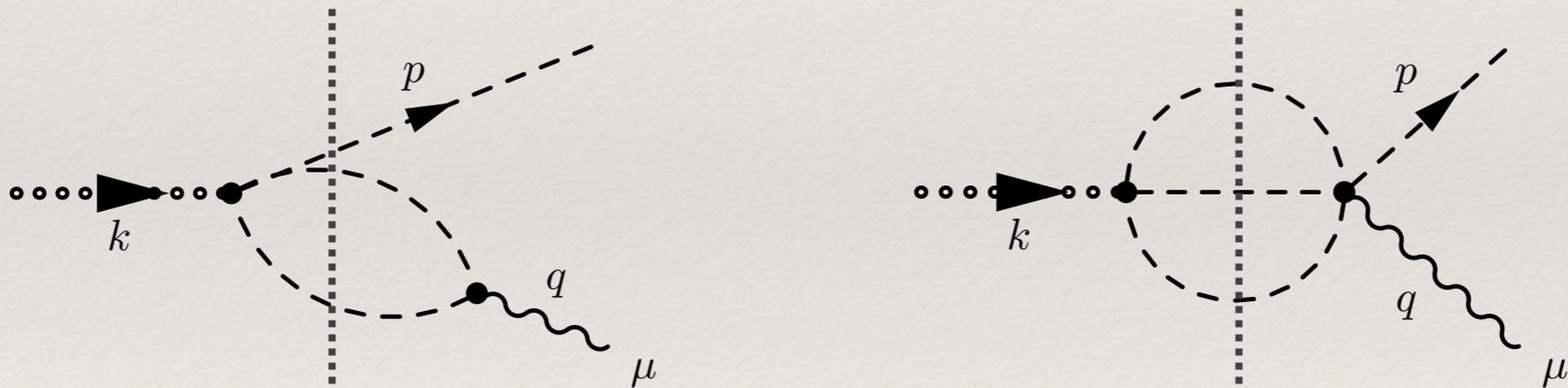
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- ❖ All other finite-size effects: exponentially suppressed

# Preliminary lattice results

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# Lattice setup

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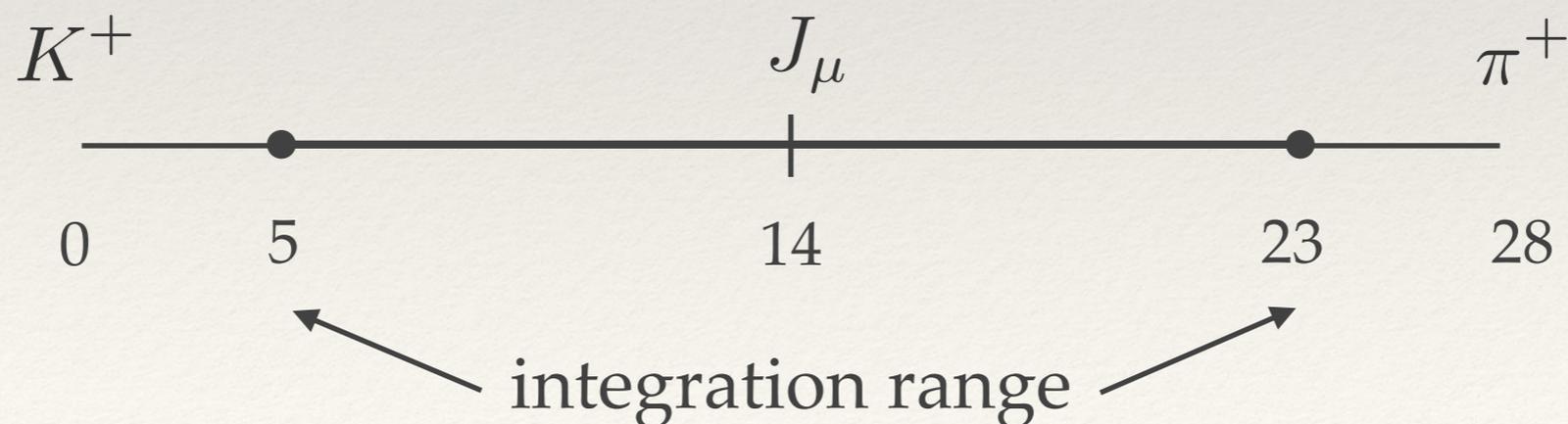
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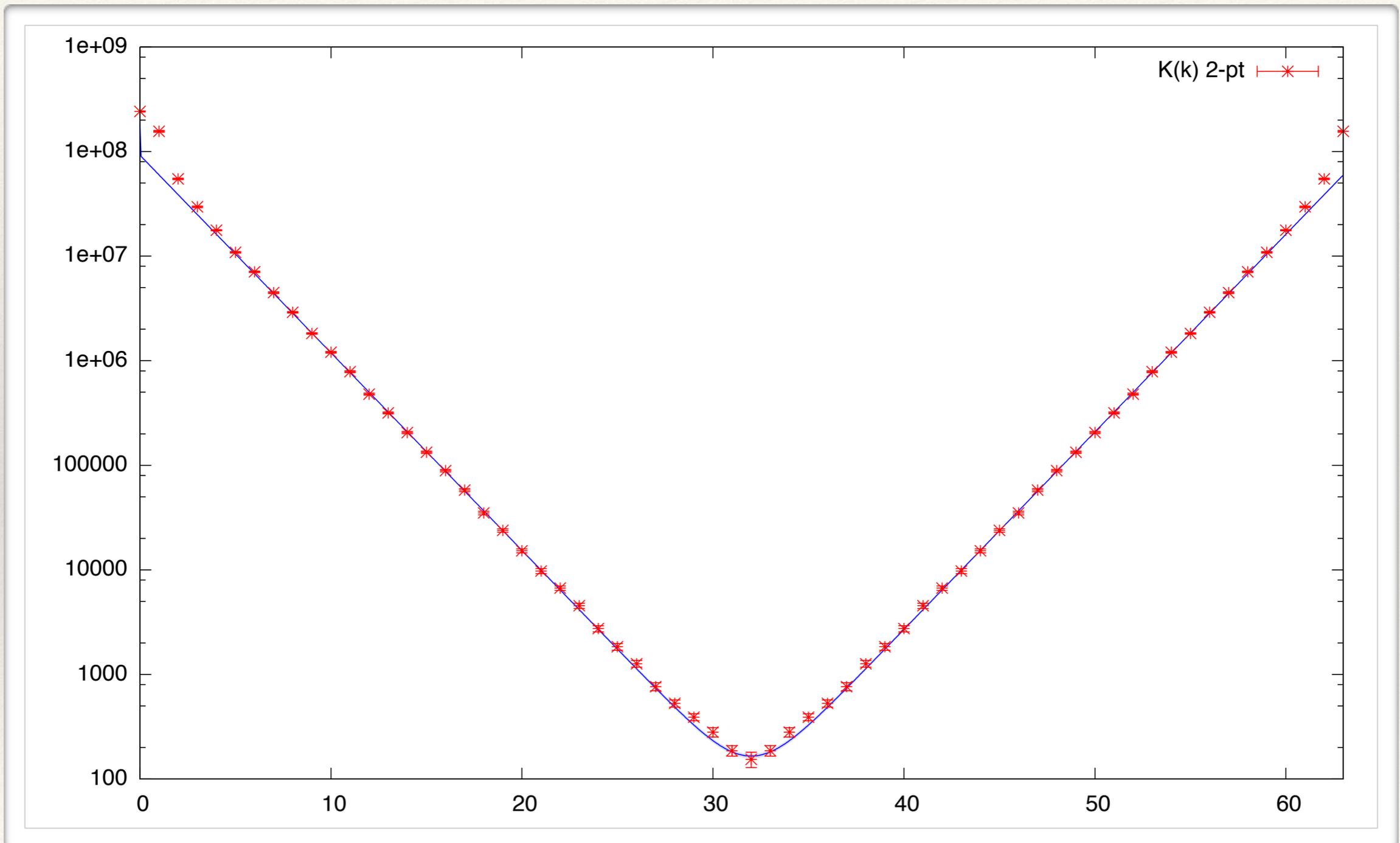
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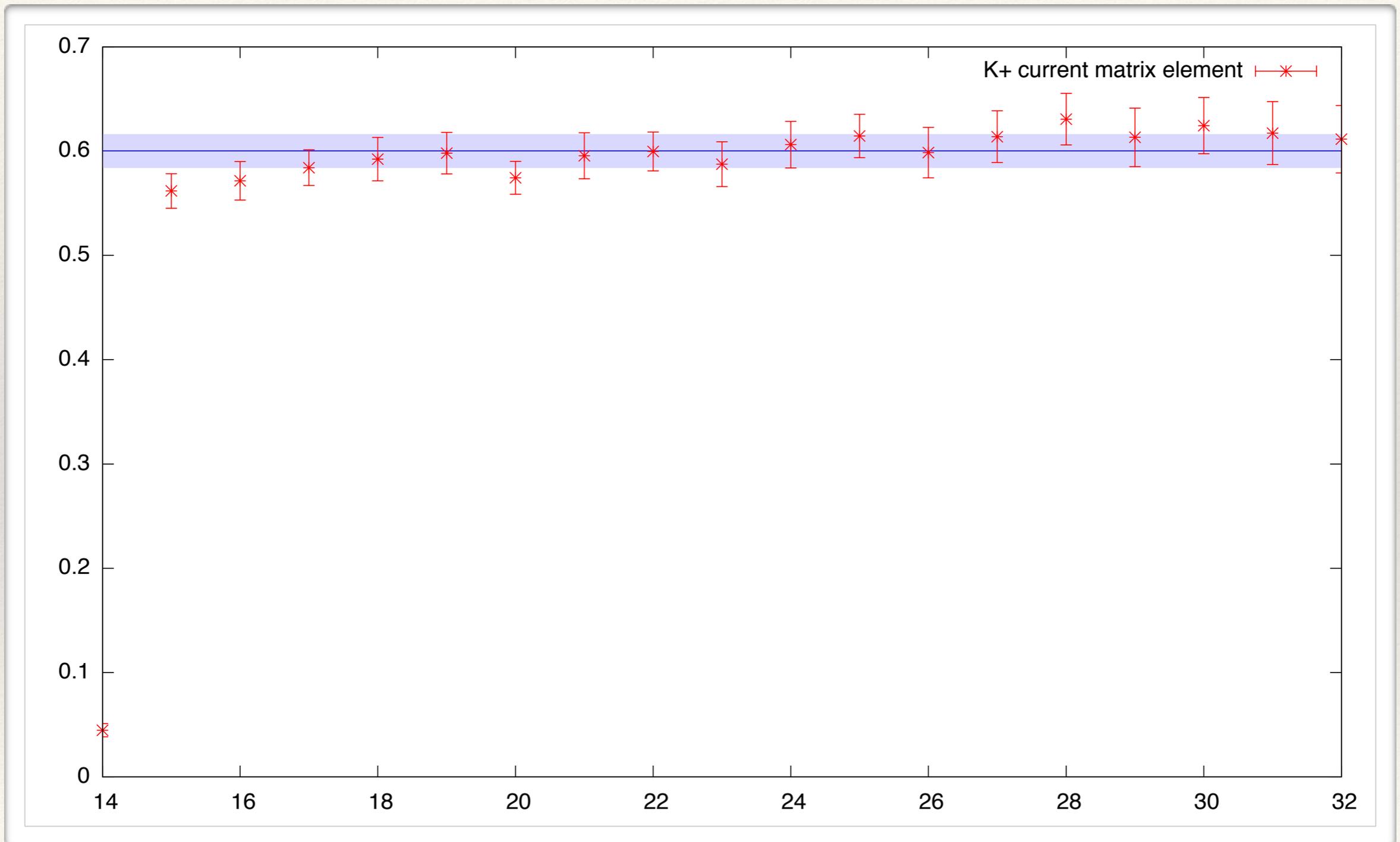
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- ❖ gauge fixed wall sources, sequential current insertion



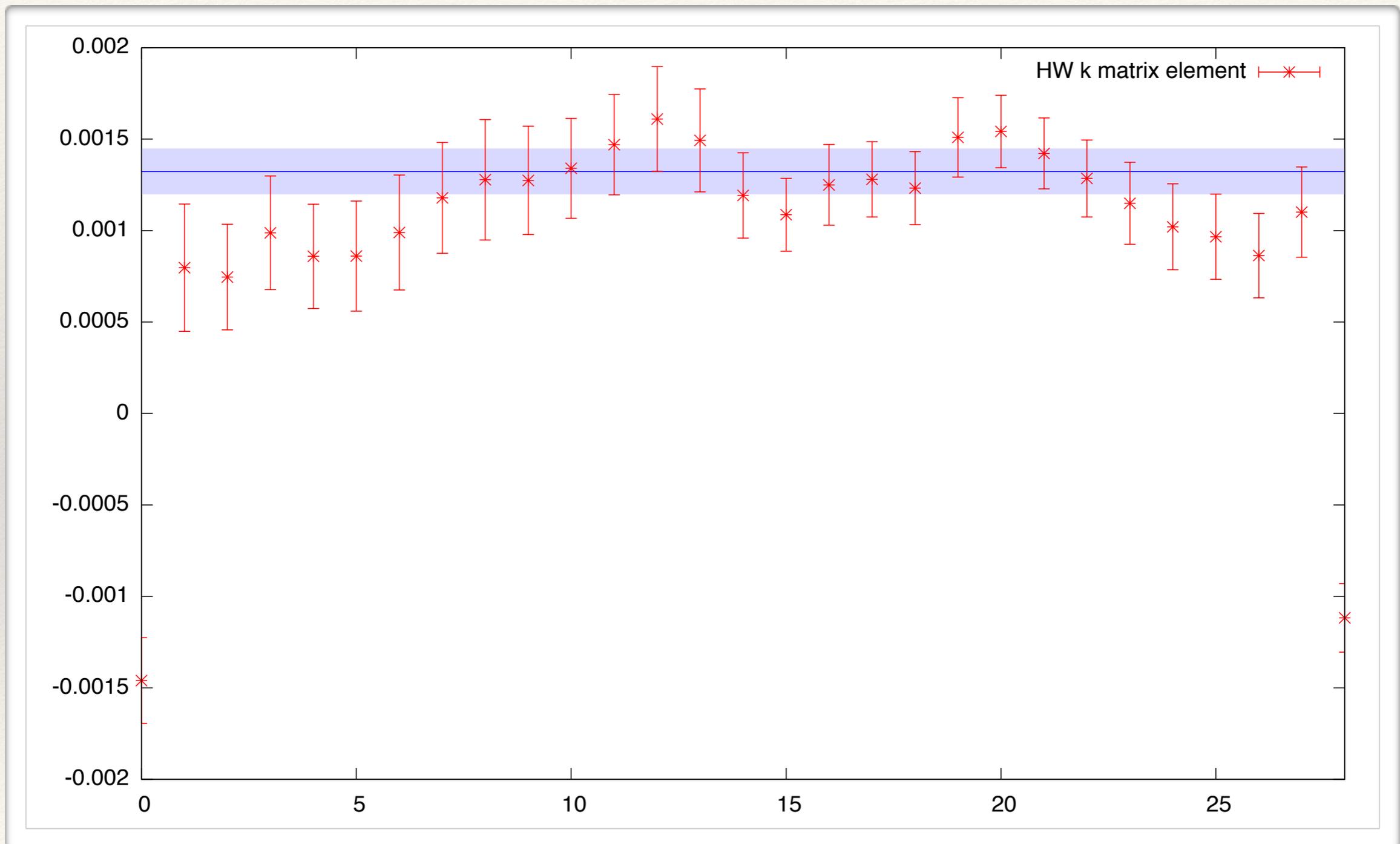
# 2-point function fit



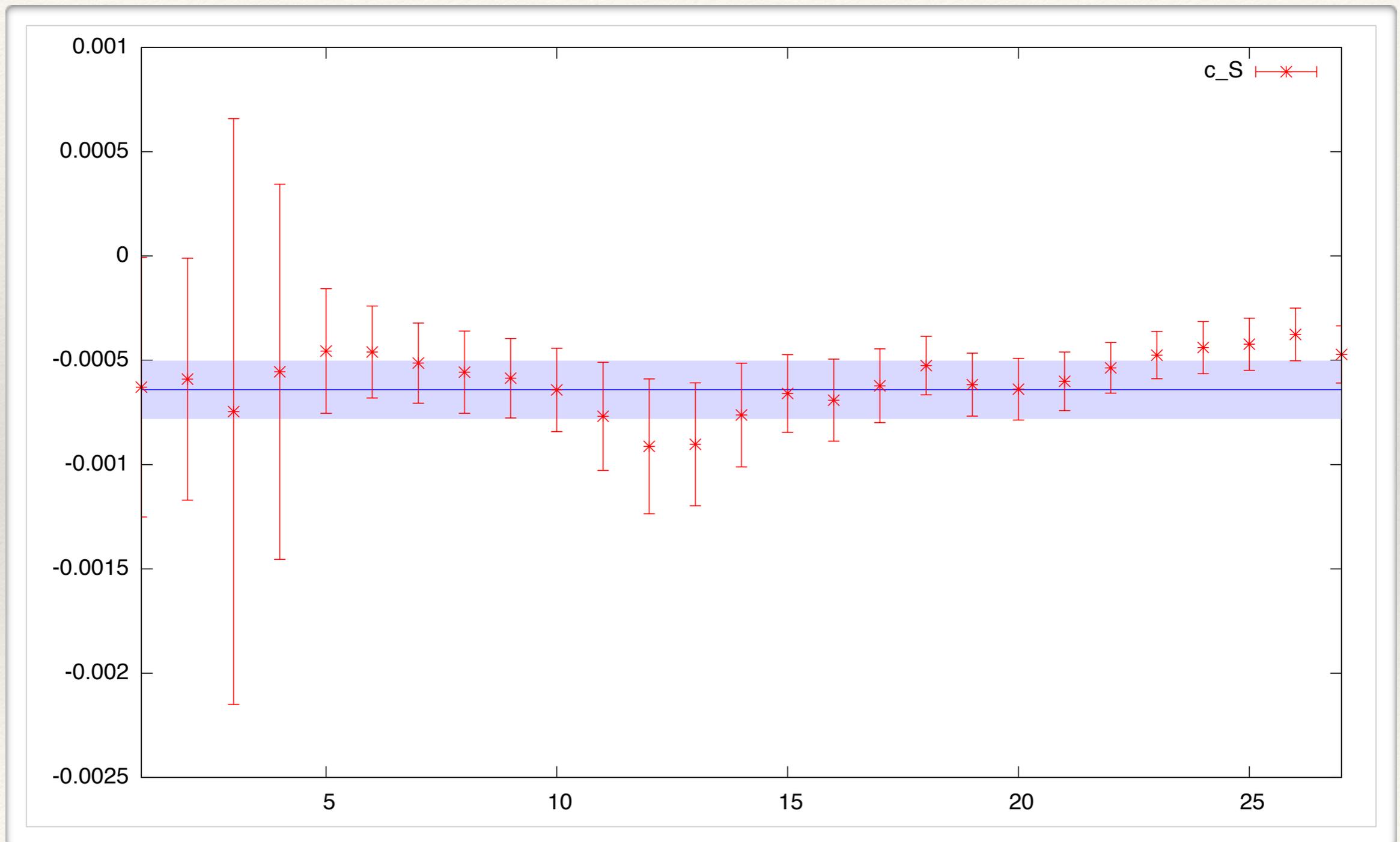
# EM current matrix element



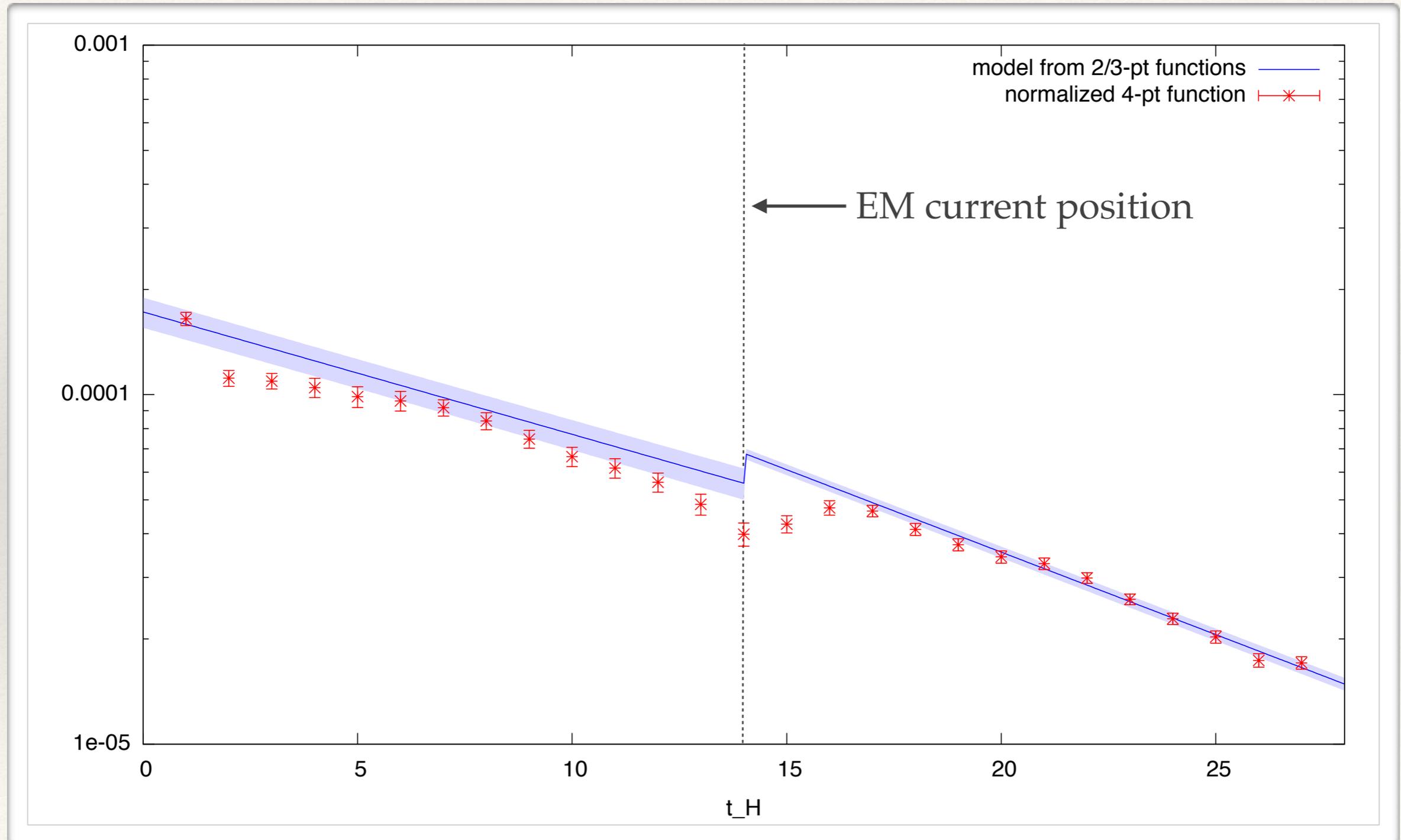
# Weak Hamiltonian matrix element



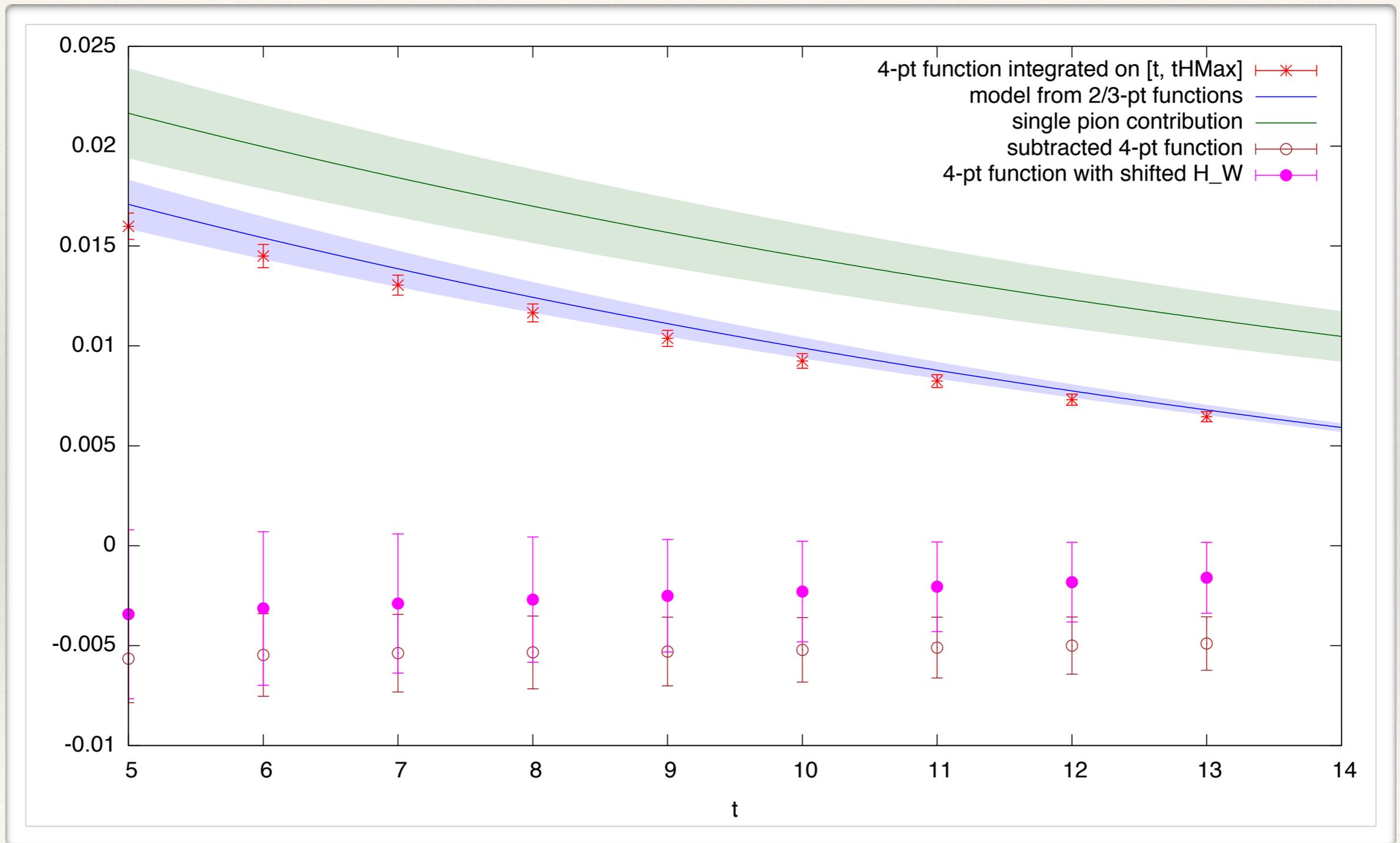
# $c_S$ determination



# Rare kaon decay correlation function



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# Summary & outlook

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- ❖ Preliminary lattice calculations agree with theory

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# Outlook

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- ❖ Aim at lighter quark masses
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*Thank you!*