



Model-Independent Constraints on Lepton-Flavor-Violating Top Decays

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Introduction

- LHC: huge top physics potential ($10^8 t$'s!).
- t unique, good place to look for New Physics (NP).
- Neutrino osc: Nature has lepton flavor violation.
- $t \rightarrow u(c)e^\pm\mu^\mp$: distinctive experimental signature.
- Are there experimental constraints on NP contributions to $t \rightarrow u(c)e^\pm\mu^\mp$?
- Take set of effective operators

$$\mathcal{L}_{eff} = \sum_i \frac{C_i}{\Lambda^{d-4}} \mathcal{O}_i + h.c.$$



to get constraints from B, K decays, $\mu \rightarrow e\gamma\dots$

Operators Contributing to $t \rightarrow u(c)e^\pm\mu^\mp$

- Take dimension-6, $SU(3) \times SU(2) \times U(1)$ -invariant op's.
- Separate into 2 classes, depending on T_L or t_R in operator.

Class One:

$$\begin{aligned}\mathcal{O}_{1,ijk} &= \bar{u}_R^i \gamma^\mu t_R \bar{L}_L^j \gamma_\mu L_L^k \\ \mathcal{O}_{2,ijk} &= \bar{u}_R^i \gamma^\mu t_R \bar{l}_R^j \gamma_\mu l_R^k \\ \mathcal{O}_{3,ijk} &= \epsilon^{ab} \bar{Q}_{La}^i t_R \bar{L}_{Lb}^j l_R^k \\ \mathcal{O}_{4,ijk} &= \epsilon^{ab} \bar{Q}_{La}^i \sigma^{\mu\nu} t_R \bar{L}_{Lb}^j \sigma_{\mu\nu} l_R^k\end{aligned}$$

$i = u, c; j, k = e\mu, \mu e$
 $a, b = SU(2)$ indices

Class Two:

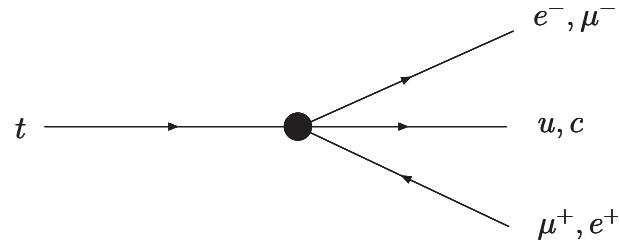
$$\begin{aligned}\mathcal{O}_{5,ijk} &= \bar{Q}_L^i \gamma^\mu T_L \bar{L}_L^j \gamma_\mu L_L^k \\ \mathcal{O}_{6,ijk} &= \bar{Q}_L^i \gamma^\mu T_L \bar{l}_R^j \gamma_\mu l_R^k \\ \mathcal{O}_{7,ijk} &= \epsilon^{ab} \bar{u}_R^i T_{La} \bar{l}_R^j L_{Lb}^k \\ \mathcal{O}_{8,ijk} &= \epsilon^{ab} \bar{u}_R^i \sigma^{\mu\nu} T_{La} \bar{l}_R^j \sigma_{\mu\nu} L_{Lb}^k\end{aligned}$$

Q_L, L_L, T_L : l.h. doublets
 u_R, l_R, t_r : r.h. singlets



Contributions of Op's to $t \rightarrow u(c)e^\pm\mu^\mp$

- Top decay for all op's proceeds via



- Assume can measure branching ratio of 10^{-7} .
- Taking $m_t = 170$ GeV and $|C_{n,ijk}| = 1$,

$$\begin{array}{lll} 2.1 & \text{TeV} & (n = 1, 2, 5, \text{ and } 6) \\ \Lambda \geq & 1.5 & \text{TeV} \quad (n = 3 \text{ and } 7) \\ & 4.0 & \text{TeV} \quad (n = 4 \text{ and } 8) \end{array}$$



- Results independent of flavor indices i, j, k .



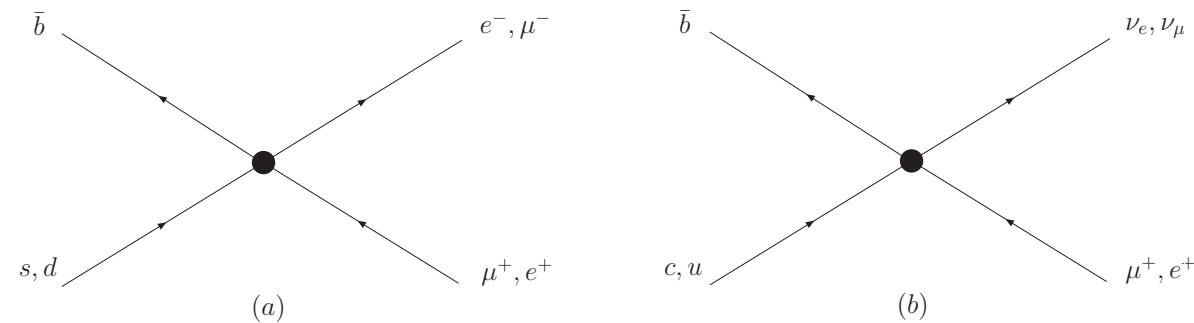
Constraints from B decays

Class Two op's contain T_L , include terms with b quarks.

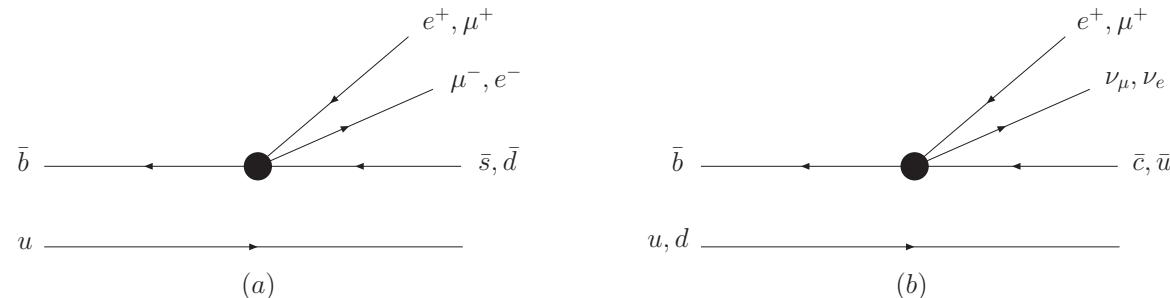
ex: $\mathcal{O}_{6,ijk} = \bar{Q}_L^i \gamma^\mu T_L \bar{l}_R^j \gamma_\mu l_R^k = (\bar{u}_L^i \gamma^\mu t_L + \bar{d}_L^{i'} \gamma^\mu b_L') \bar{l}_R^j \gamma_\mu l_R^k$

→ contribute at tree-level to B decay.

2-body



3-body



Constraints from 2-body B decays

The Class Two op's contribute to ($\langle 0 | \bar{d} \gamma^\mu \gamma_5 b | B^0(p) \rangle = i F_B p^\mu$):

Op's 5, 6: $B^0, B_s \rightarrow e^\pm \mu^\mp$, $\Gamma = \frac{1}{64\pi} \frac{|C_{n,ijk}|^2}{\Lambda^4} F_B^2 m_B m_\mu^2$
(helicity-suppressed)

Op 7: $B^+, B_c \rightarrow \ell^+ \nu$, $\Gamma = \frac{1}{64\pi} \frac{|C_{n,ijk}|^2}{\Lambda^4} F_B^2 \frac{m_B^5}{(m_b + m_{u(c)})^2}$

Op 8: 0 (tensor operator)

$$\text{Br}(B^0 \rightarrow e^\pm \mu^\mp) \leq 1.7 \times 10^{-7} (90\% \text{CL})$$

$$\rightarrow \frac{|C_{5(6),ujk}|}{\Lambda^2} \leq \frac{1}{(3.7 \text{ TeV})^2}, \quad \frac{|C_{5(6),cjk}|}{\Lambda^2} \leq \frac{1}{(1.8 \text{ TeV})^2}$$

$$\text{Br}(B^+ \rightarrow e^+ \nu) \leq 9.8 \times 10^{-6} (90\% \text{CL}) \rightarrow \frac{|C_{7,u\mu e}|}{\Lambda^2} \leq \frac{1}{(16 \text{ TeV})^2}$$

$$\text{Br}(B^+ \rightarrow \mu^+ \nu) \leq 1.7 \times 10^{-6} (90\% \text{CL}) \rightarrow \frac{|C_{7,u e \mu}|}{\Lambda^2} \leq \frac{1}{(10 \text{ TeV})^2}$$

Constraints from 3-body B decays

Similar to t decays: no helicity suppression, op 8 can contribute.

Op's 5,6: $\text{Br}(B^+ \rightarrow K^+ e^\pm \mu^\mp) \leq 9.1 \times 10^{-8}$,

Compare to inclusive via $B \rightarrow \pi \ell^+ \nu$.

$$\rightarrow \frac{|C_{5(6),cjk}|}{\Lambda^2} \leq \frac{1}{(16 \text{ TeV})^2}, \quad \frac{|C_{5(6),ujk}|}{\Lambda^2} \leq \frac{1}{(8 \text{ TeV})^2}$$

Op's 7,8: Take $2 \times$ exp. error as estimate of NP contribution.

$\text{Br}(B \rightarrow X_u \ell^+ \nu) = 2.33 \pm .22 \times 10^{-3}$, $\text{Br}(B^- \rightarrow X_c e^+ \nu) = 10.8 \pm 0.4 \%$:

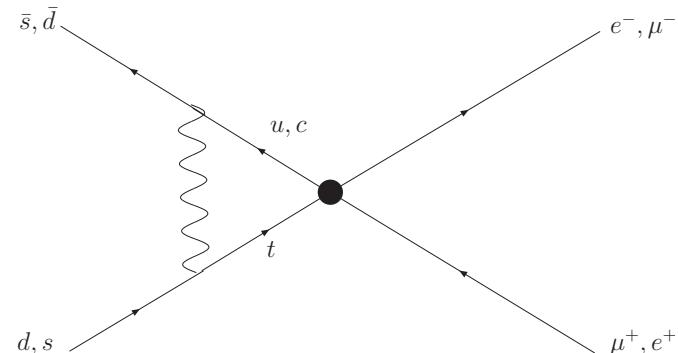
$$\begin{aligned} \rightarrow \frac{|C_{7,ujk}|}{\Lambda^2} &\leq \frac{1}{(3 \text{ TeV})^2}, & \frac{|C_{7,cjk}|}{\Lambda^2} &\leq \frac{1}{(1 \text{ TeV})^2} \\ \frac{|C_{8,ujk}|}{\Lambda^2} &\leq \frac{1}{(7 \text{ TeV})^2}, & \frac{|C_{8,cjk}|}{\Lambda^2} &\leq \frac{1}{(3 \text{ TeV})^2} \end{aligned}$$



Constraints from K decays

Again, **Class Two** op's give (CKM-suppressed) contributions at tree level to $K_L \rightarrow e^\pm \mu^\mp$ (just like B case).

Op's **of both classes** can contribute at one loop.
i.e.,



Using

$\text{Br}(K_L \rightarrow e^\pm \mu^\mp) \leq 4.7 \times 10^{-12}$:
(3-body K decays weaker)

$$\frac{|C_{5(6),ujk}|}{\Lambda^2} < \frac{1}{(60 \text{ TeV})^2}$$

$$\frac{|C_{5(6),cjk}|}{\Lambda^2} < \frac{1}{(40 \text{ TeV})^2}$$

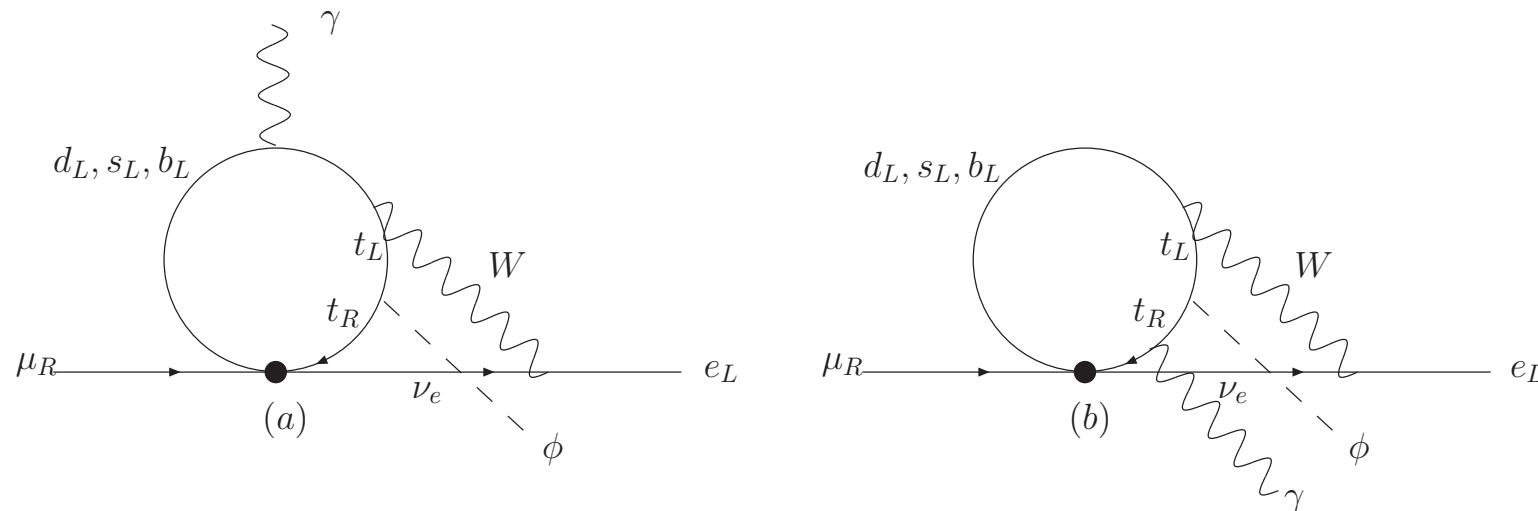
$$\frac{|C_{4,ujk}|}{\Lambda^2} \lesssim \frac{1}{(1 \text{ TeV})^2}$$



$\mu \rightarrow e\gamma$ and $\mu \rightarrow e$ Conversion



Ops 3, 4: $\mu \rightarrow e\gamma$ via GIM-sup'd ($\sim m_h^2/m_W^2$) 2-loop diags:



$$\text{Br}(\mu \rightarrow e\gamma) < 1.2 \times 10^{-11} \rightarrow \frac{|C_{3,cjk}|}{\Lambda^2} \lesssim \frac{1}{(1 \text{ TeV})^2}$$

$$\frac{|C_{4,cjk}|}{\Lambda^2} \lesssim \frac{1}{(\text{few TeV})^2}$$

$\mu \rightarrow e$ conversion: same diagrams as K_L decay,
Limits from SINDRUM II similar to or weaker than K_L 's.





Results

Operator	$t \rightarrow u(c)e^\pm\mu^\mp$	B , 2-body	B , 3-body	K_L
$\mathcal{O}_{1(2),ijk}$	2.1	-	-	-
$\mathcal{O}_{3,ijk}, i = u$	1.5	-	-	-
$\mathcal{O}_{3,ijk}, i = c$	1.5	-	-	-
$\mathcal{O}_{4,ijk}, i = u$	4.0	-	-	1
$\mathcal{O}_{4,ijk}, i = c$	4.0	-	-	-
$\mathcal{O}_{5(6),ijk}, i = u$	2.1	3.7	8	60
$\mathcal{O}_{5(6),ijk}, i = c$	2.1	1.8	16	40
$\mathcal{O}_{7,ijk}, i = u, j = e, k = \mu$	1.5	10	3	-
$\mathcal{O}_{7,ijk}, i = u, j = \mu, k = e$	1.5	16	3	-
$\mathcal{O}_{7,ijk}, i = c$	1.5	-	1	-
$\mathcal{O}_{8,ijk}, i = u$	4.0	-	7	-
$\mathcal{O}_{8,ijk}, i = c$	4.0	-	3	-



Class One operators largely unconstrained.



Conclusions

- Operators which can give $t \rightarrow u(c)e^\pm\mu^\mp$ probe \sim few TeV range.
- Several operators constrained by B , K decays.
- Some operators currently not constrained.
- $t \rightarrow u(c)e^\pm\mu^\mp$ could occur at LHC!
- arXiv: 0807.4199 (hep-ph), accepted to PRD.

